

## THE 19th **INTERNATIONAL CONFERENCE ON MAGNETISM**

July 8 - 13, 2012 Bexco, Busan, Korea www.icm2012.org

KPS The Korean Physical Society

The Korean Magnetics Society



Hosted by

International Union of Pure and Applied Physics

### CONTENTS

Sponsor • 26

## THE 19<sup>th</sup> **INTERNATIONAL CONFERENCE ON MAGNETISM**



- Welcome Address 3 Organizing Committee • 4 Int'l Advisory Committee • 5 Program Committee • 7 Editors & Speakers • 10 Program at-a-Glance • 11 Venue Layout • 18 Ceremony & Event Program • 20 General Information • 22 Exhibition • 25



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Oral Presentation • 27
- Plenary Lecture • 28
- Half-Plenary Lecture • 30
- Invited & Contributed Presentation • 33
Poster Presentation • 115
- July 9 (PA~PO) • 116
- July 10 (QA~QP) • 182
- July 12 (RA~RR) • 249
- July 13 (SA~SO) • 316
Author Index • 373
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### LIST OF ABSTRACTS

Half-Plenary Lecture		
nvited & Contributed Presentation		
uly 9 (Mon)		
AA: Multiferroics I - mainly manganites	AB: Non-fermi liquids and quantum phase transitions I	AC: Low-dimensional / Frustrated spin systems
D: Surface and interface effects I	AE: Electric field effect on magnetic systems	AF: Advanced methods of spin structure determination
AG: Arrays of magnetic nanostructures I	AH: Magnetic transducers in biomedicine	Al: Semiconductor spintronics I - group IV materials
U: Crystalline, nanocrystalline and amorphous materials IC: Organic and molecular magnetism / Spin ladder	BA: Superconductivity I - cuprate and other superconductors	BB: Valence fluctuations I
F: 3d transition metal oxides	BD: Exchange bias BG: Energy assisted magnetic recording	BE: Magnetic semiconductor BH: Interdisciplinary technology
I: STT MRAM and magnetic logic	BJ: Ferrites, garnets and other materials	bri, interdisciplinary technology
uly 10 (Tue)		
A: Superconductivity II - cuprate and other superconductors	CB: Magnetic nanoparticles I	CC: Spin liquid / Spin ice
D: Heavy fermions I	CE: Spin transfer oscillators	CF: Actinides and lanthanides
G: Semiconductor spintronics II - group III~V materials	CH: Heusler alloys etc	CI: Multiferroics II – scattering
J: Magneto-dielectric materials or meta-materials	DA: Superconductivity III - Fe-based superconductors	DB: Kondo systems I
C: Spin-orbit / Spin-lattice / Spin-oribital physics	DD: Diluted magnetic semiconductors and others	DE: Magnetic memories and logics
F: Chiral magnet and magnetic skyrmions	DG: Magnetic nanowires	DH: Oxide
I: Spin caloritronics I	DJ: Applications	EA: Non-Fermi liquids and quantum phase transitions II
B: SCES theory I	EC: Electronic structure / Spintronic materials	ED: Magnetic thin films and nanostructures I
E: Spin-orbit spin torque H: Novel materials and devices l	EF: Intermetallic compounds EI: Perpendicular magnetic anisotropy materials	EG: Metal spintronics l EJ: Rare-earth hard magnetic materials
		-
uly 11 (Wed) A: Spin caloritronics II	FB: Heavy fermions II	FC: Ultrafast switching I
2: Vortex dynamics I	FE: SCES theory II	-
uly 12 (Thu)		
A: Superconductivity IV - Fe-based superconductors	GB: Multiferroics III - nonreciprocal effect and elecgtronic ferroelectricity	GC: Heavy fermions III
iD: Ultrafast switching II	GE: Domain wall motion I	GF: Spin glasses and diluted magnets
iG: Arrays of magnetic nanostructures ll	GH: Novel materials and devices II	GI: Organic spintronics and carbon-based spintronics
J: Intermetallic and other hard magnets	HA: Superconductivity V - Fe-based superconductors	HB: [Symposium] High performance soft magnetic materials and their app
IC: Magnetism in s,p electron systems	HD: Spin waves I	HE: Metal spintronics II
F: Spin transfer torque switching	HG: Magnetometery in nano-scale	HH: Magnetometry in macro-scale
II: Topological insulators I	HJ: 4d and 5d compounds	IA: Non-fermi liquids and quantum phase transitions III
3: [Symposium] High performance soft magnetic materials and their applications II	IC: Magnetic phase transition	ID: Vortex dynamics II
E: Domain wall motion II H: Surface and interface effects II	IF: Magnetic tunnel junctions II: Topological insulators II	IG: Valence fluctuations II IJ: Ferrites and other materials
uly 13 (Fri)	JB: Multiferroics IV - noncollinear magnets	IC: Horse farmions N/
A: Superconductivity VI - Fe-based and other superconductors D: Magnetism theory / Simulation of quantum and classical systems	JE: Domain wall motion III	JC: Heavy fermions IV JF: Metal spintronics III
G: Spin waves II	JH: Nanostructured and composite hard magnetic materials	JI: Strong magnetic anisotropy materials
: Magnetocaloric effects / Magnetoelastic materials	KA: Kondo systems II	KB: Magnetic nanoparticles II
C: Magnetic thin films and nanostructures II	KD: Characterization of magnetic properties	KE: Domain walls and spin ice system
(F: Novel spintronic devices and materials	KG: SCES theory III	KH: Coercivity mechanism
I: Theoretical calculation	KJ: New developments	
had a December for the se		
oster Presentation		
A: Multiferroics I	PB: Superconductivity I	PC: Superconductivity V
D: Heavy fermions I	PE: Kondo Impurity and kondo lattice systems	PF: Theory of strongly correlated matter I
G: Magnetic materials and characterization methods	PH: 3d transition metal oxides I	PI: 3d transition metal oxides II
J: Spin-dependent transport l	PK: Perpendicular magnetic anisotropy and strong anisotrorpy	PL: Surface and interface effects including exchange bias
M: Soft magnetic materials I	PN: Dilutedmagnetic semiconductor/nano-composite l	PO: Interdisciplinary topics
uly 10 (Tue)		
A: Multiferroics II	QB: Superconductivity II	QC: Heavy fermions II
D: Valence fluctuations	QE: Frustrated systems, kagome, triangular systems	QF: 1D, low-dimensional systems
)G: Intermetallic compounds I	QH: Intermetallic compounds II	QI: Lanthanides I
): Lanthanides II	QK: Spin-dependent transport II	QL: Diluted magnetic semiconductors and others
M: Magnetic characterization	QN: Soft magnetic materials II	QO: Novel magnetic materials and devices II
P: Magnetic recording and memories /		
uly 12 (Thu)		
A: Multiferroics III	RB: Superconductivity III	RC: Topological insulators I
D: Heavy fermions III	RE: Non-fermi liquids and quantum phase transitions I	RF: Theory of strongly correlated matter II
G: Theory of strongly correlated matter III	RH: Theory, spin, magnetic materials	RI: Phase transition
U: Vortex dynamics	RK: Ultrafast dynamics	RL: Spin electronics I
M: Theoretical calculation	RN: Magnetic nanoparticles RQ: Measuring techniques and Instrumentation II	RO: Hard magnetic materials l RR: Industrial applications
P: Measuring techniques and instrumentation I		
P: Measuring techniques and instrumentation I uly 13 (Fri)		SC: Superconductivity VI
P: Measuring techniques and instrumentation I uly 13 (Fri) A: Multiferroics IV	SB: Superconductivity IV SE: Heavy fermions IV	SC: Superconductivity VI SF: Non-fermi liquids and quantum phase transitions II
Reasuring techniques and instrumentation I     Iuly 13 (Fri)     SA: Multiferroics IV     SD: Topologicai insulators II	SB: Superconductivity IV	SC: Superconductivity VI SF: Non-fermi liquids and quantum phase transitions II SI: Spin waves
RP: Measuring techniques and instrumentation I July 13 (Fri) SA: Multiferroics IV SD: Topologicai insulators II SG: New developments SJ: Modeling SJ: Magnetic thin films and others	SB: Superconductivity IV SE: Heavy fermions IV	SF: Non-fermi liquids and quantum phase transitions II

SN: Hard magnetic materials II

SO: Novel magnetic materials and devices I

### Dear Colleagues,

On behalf of the Organizing Committee, and all of those who were involved in preparing for the 19th International Conference on Magnetism (ICM2012), we wish to extend a warm-hearted welcome to all participants of ICM2012. It is our great honor and privilege to host the ICM2012 in Korea.

The major scientific societies in Korea, the Korean Physical Society (KPS) and the Korean Magnetics Society (KMS) are pleased to cohost the ICM2012, under the auspices of the International Union of Pure and Applied Physics (IUPAP). The ICM2012 incorporates the International Conference on Strongly Correlated Electron Systems (SCES) held annually.

As the most highly acclaimed conference, a considerable number of abstracts have been submitted from 52 countries around the world. In this conference, we expect over 1,700 participants around the world. For the scientific program, we have planned 7 plenary lectures including 3 Nobel laureates' lectures, 14 Half-plenary lectures, 509 oral presentations including 135 invited lectures, along with 1,502 poster presentations. In addition, 7 satellite symposia will be held before or after the ICM2012.

I wish to take this opportunity to thank all the sponsors for their generous support for the ICM2012. Also, I would like to convey my sincere gratitude to the international advisory members for their valuable advices and to the members of the ICM2012 organizing committee for their tremendous efforts in making this conference a success.

We wish you all a fruitful meeting and hope that you will benefit from the rich scientific programs, and your visit to wonderful Busan will last forever as a pleasant memory.



Prof. Sung-Chul Shin

2 | The 19<sup>th</sup> International Conference on Magnetism

SM: Magnetic thin films and others

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E

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6 | The 19<sup>th</sup> International Conference on Magnetism

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July 8-13, 2012 Busan, Korea | 7

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2. Magnatic Ctu	wetween and laterations			11. Measuring Tech	niques and Instrument	
3. Magnetic Str	ructures and Interactions	DOCTECII	Karaa		Derac Son	Hannam University
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8. Soft and Har	d Magnetic Materials and th	neir Applications				
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	Noidha wiesenddriger		Gernariy

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Time	July 8 (Sun)	July 9 (Mon)	July 10 (Tue)	July 11	(Wed)	July 12 (Thu)	July 13 (Fri)
08:30 09:00 —		Opening & Award		Plena	ny A		
05.00		Plenary 1	Half-Plenary	PPO		Lielf Dienen	Invited 9
10:00		FFUI	<b>1~6</b> HP11~HP32	Plena	rv 5	Half-Plenary 9~14	Invited & Contributed
10.00		Coffee Break		PPO		HP51~HP72	II~AI
11:00 —		Plenary 2	Coffee Break	Coffee	Break	Coffee Break	Coffee Break
11.00		PP02		Invited &			
12:00 —		Plenary 3 PP02	Invited & Contributed CA~CJ	FA~FE	Half-Plenary 7~8 HP41~HP42	Invited & Contributed GA~GJ	Invited & Contributed KA~KJ
13:00 —		Lunch	Lunch			Lunch	Lunch
14:00 —	Registration	Poster	Poster			Poster	Poster
15:00 —		Presentation I PA~PO	Presentation II QA~QP			Presentation III RA~RR	Presentation IV SA~SO
- 16:00 —	KPS 60th Anniversary Nobel laureate in Physics Public Lecture	Invited & Contributed	Invited & Contributed	Excur	sion	Invited & Contributed HA-HJ	Plenary 6 PP06
17:00 —		AA~NJ	DA~DJ			пА~пј	Plenary 7
_		Break	Break			Break	PP07
18:00 — _	Welcome Reception	Invited & Contributed BA~BJ	Invited & Contributed EA~EJ			Invited & Contributed IA-IJ	Closing
19:00 — - 20:00						Banquet	

### PROGRAM AT-A-GLANCE

# Program at-a-Glance

### PROGRAM AT A GLANCE

	م مناطبي <del>ا</del>	EXNIBITION		Registration		
	ЗF	Room 301				
		Room 206				
		Room 205				
	2F	Room 204				
	2	Room 203				
July 8 (Sun)		Room 202			Welcome Reception (2F, Lobby)	
} vlu(		Room 201			Welcome (2F, L	
		Room 109~10				
	1F	Room 104~5				
	-	Room 106~8				
		Room 101~3				
	AAttoria	Auditorium		KPS 60th Aniversary Nobel laureate in Physics Public Lecture 1) Albert Fert 2) Maus von Kitzing		
			08:00	15:00 14:00 15:00 15:00 15:00	18:00	20:00



### PROGRAM AT A GLANCE

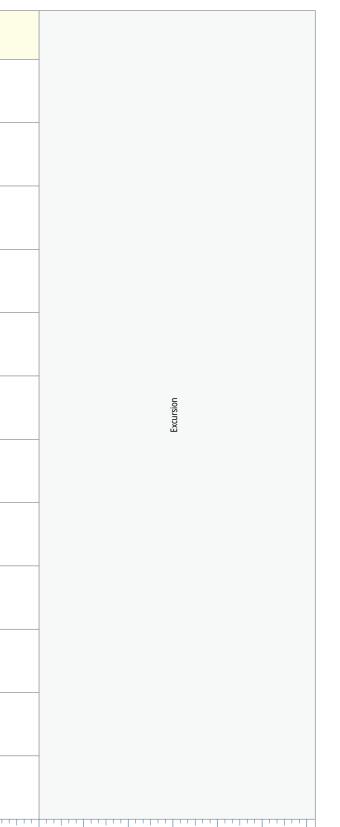
14:00											PA-PO Poster Presentation
16:00 2.50	AA Multiferroics I - mainly manganites	AB Non-fermi liqui and quantur phase transitior	AC Low-dimensional / Frustrated spin systems	AD Surface and interface effects I	AE Electric field effect on magnetic systems	Advanced methods of spin structure determination	AG Arrays of magnetic nanostructures I	AH Magnetic transducers in biomedicine	AI Semiconductor spintronics I - group IV materials	AJ Crystalline, nanocrystalline and amorphous materials	
- 00.7					Bre	Break					Exhibition
18:00	BA Superconductivity I - cuprate and other superconductors	BB Valence fluctuations I	BC Organic and molecular magnetism / Spin ladder	BD Exchange bias	BE Magnetic semiconductor	BF 3d transition metal oxides	8G Energy assisted magnetic recording	BH Interdisciplinary technology	BI STT MRAM and magnetic logic	BJ Ferrites, garnets and other materials	
19:00											

Program at-a-Glance

### PROGRAM AT A GLANCE

	Exhibition						Exhibition		QA~QP Poster Presentation		Exhibition		
ЭС	ΔF	Room 301								-			
		Room 206					C Magneto-dielectric materials or meta-materials			DJ Applications		EJ Rare-earth hard magnetic materials	
		Room 205		HP31 Half-Plenary 5 Yoshinori Tokura	HP32 Half-Plenary 6 Bemhard Keimer		CI Multiferroics II - scattering			DI Spin caloritronics I		EI Perpendicular magnetic anisotropy materials	
75		Room 204					CH Heusler alloys etc			DH Oxide		EH Novel materials and devices I	
ſ	V	Room 203					CG Semiconductor spintronics II - group III-V materials			DG Magnetic nanowires		EG Metal spintronics I	
July 10 (Tue)		Room 202				Coffee Break	Coffee Break CF Actinities and Lanthanides	Lunch		DF Chiral magnet and magnetic skyrmions	Break	EF Intermetallic compounds	
July 1		Room 201					CE Spin transfer oscillators			DE Magnetic memories and logics	Bre	EE Spin-orbit spin torque	
		Room 109~10					CD Heavy fermions I			Diluted magnetic semiconductors and others		ED Magnetic thin films and nanostructures	
Ļ		Room 104~5					CC Spin liquid / Spin ice			DC Spin-orbit / Spin-oribital physics		Electronic structure Magnetic thin films / Spintronic and nanostructures materials	
Ţ	_	Room 106~8		HP21 Half-Plenary 3 Hyun-Woo Lee	HP22 Half-Plenary4 Piers Coleman		CB Magnetic nanoparticles I			DB Kondo systems I		EB SCES Theory I	
		Room 101~3		HP11 Half-Plenary 1 Yoshichika Onuki	HP12 Half-Plenary 2 Tae-Won Noh		CA Superconductivity II - cuprate and other superconductors			DA Superconductivity III - Fe-based superconductors		EA Non-fermi liquids and quantum phase transitions II	
	Auditorium												
			08:00		10:00	0.1	12:00	13:00	14:00	16:00	- 00.71	18:00	19:00

	LO						E		
	EXNIDITION						Exhibition		
ЗF	Room 301								
	Room 206								
	Room 205						HP41 Half-Plenary 7 Stuart Parkin	HP42 Half-Plenary 8 Gabriel Aeppli	
2F	Room 204								
2	Room 203						FE SCES theory II		
July 11 (Wed)	Room 202						Pltrafast switching I Vortex dynamics I		
July 1	Room 201					Coffee Break	FC Ultrafast switching I		
	Room 109~10								
1	Room 104~5						FB Heavy fermions II		
-	Room 101~3 Room 106~8 Room 104~5 Room 109~10								_
	Room 101~3						FA Spin caloritronics II		
	Auditorium		PP04 Plenary 4 Klaus von Klitzing	PP05 Plenary 5	Roland Wisendanger				
		08:00	00:60	10:00 10:00		e e e	2. 	12:00	



17:00 -

16:00

14:00

15:00

### PROGRAM AT A GLANCE

Program at-a-Glance

19:00 -

20:00 -

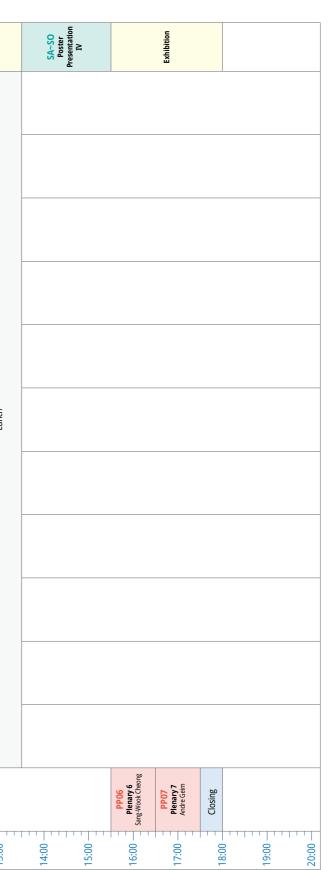
18:00 -

### PROGRAM AT A GLANCE

	Evhihition			Exhibition				RA~RR Poster Tresentation		Exhibition			
;	ЗF	Room 301											Banquet
		Room 206					G Internetallic and other hard magnets			HJ 4d and 5d compounds		J Ferrites and other materials	
		Room 205		HP71 Half-Plenary 13 Shoucheng Zhang	HP72 Half-Plenary 14 Claudia Felser		GI Organic spintronics and carbon-based spintronics			HI Topological insulators I		II Topological insulators II	
	2F	Room 204					GH Novel materials and devices II			HH Magnetometry in macro-scale		IH Surface and interface effects II	
	5	Room 203					GG Arrays of magnetic nanostructures II			HG Magnetometery in nano-scale		IG Valence fluctuations II	
(nul) 71 (Inu)		Room 202				Coffee Break	GF Spin glasses and diluted magnets	Lunch		HF Spin transfer torque switching	Break	IF Magnetic tunnel junctions	
1 July 1		Room 201				Coffee	GE Domain wall motion I	ΓΠ		HE Metal spintronics II	Bre	IE Domain wall motion II	
		Room 109~10					GD Ultrafast switching II			HD Spin waves I		ID Vortex dynamics II	
	1F	Room 104~5					GC Heavy fermions III			HC Magnetism in s,p electron systems		IC Magnetic phase transition	
	-	Room 106~8		HP61 Half-Plenary 11 Xiaozhong Zhang	HP62 Half-Plenary 12 LJ Heyderman		GB Multiferroics III - nonreciprocal effect and elecgtronic ferroelectricity			HB [Symposium] High performance soft magnetic materials and their applications I		IB [Symposium] High performance soft magnetic materials and their applications II	
		Room 101~3		HP51 Half-Plenary9 Caroline A Ross	HP52 Half-Plenary 10 Hideo Ohno		GA Superconductivity IV - Fe-based superconductors			HA Superconductivity V - Fe-based superconductors		IA Non-fermi liquids and quantum phase transitions III	
	Auditorium												
			08:00		10:00	11.00	12:00	13:00	14:00	16:00	- 00:/1	18:00	19:00

July 13 (Fri) 2F 3F	Room 201 Room 202 Room 203 Room 204 Room 205 Room 206 Room 301	JF JG JH JI Chronometer		Coffee Break	KF         KG         KH         K1         K1         K3           Novel spintronic         SCES theory III         Coercivity         Theoretical         New developments           and materials         and materials         calculation         calculation         calculation	-
July 13 (Fri)	Room 202 Room 203 Room 204 Room 205	JF JG JH JI Chronometer	Spin waves u Nandszuckured suong magneux and composite anisotropy hard magnetic materials materials	se Break	KG KH KI SCES theory III Coercivity Theoretical mechanism calculation	-
July 13 (Fri)	Room 202 Room 203 Room 204	JG JA Maral criterize Criterized	Spin waves II Nanostructured and composite hard magnetic materials	se Break	KG KH SCES theory III Coercivity mechanism	-
July 13 (Fri)	Room 202 Room 203	Matal cointronite Crin waves IT	TI SAARA LIIde	se Break	KG SCES theory III	-
July 13 (Fri)	Room 202	JF Matal contronice		se Break		-
	H			se Break	F bintronic ices iterials	
	101 Jone	Hew o		Coffe	KF Novel spintroni devices and materials	Lunch
15			motion III		KE Domain walls and spin ice system	
Ť	04~5 Room 109~10	JD Magnatics theory	magneusin ureory / Simulation of quantum and classical systems		KD Characterization of magnetic properties	
-	Room 104~5	JC Heavyfermions IV	<ul> <li>neavy reminions in magnetism mediy / Ruitiation of quantum and classical systems</li> </ul>		KC KD Magnetic thin films Characterization of Dc and nanostructures magnetic s properties	
	Room 101~3 Room 106~8 Room 1	JB Multiferroice IV	multurerroics IV noncollinear magnets		KB Magnetic nanoparticles II	
	Room 101~3	JA Superconductivity	superconductivity VI - Fe-based and other superconductors		Kondo systems II	
	_					
	Auditorium					

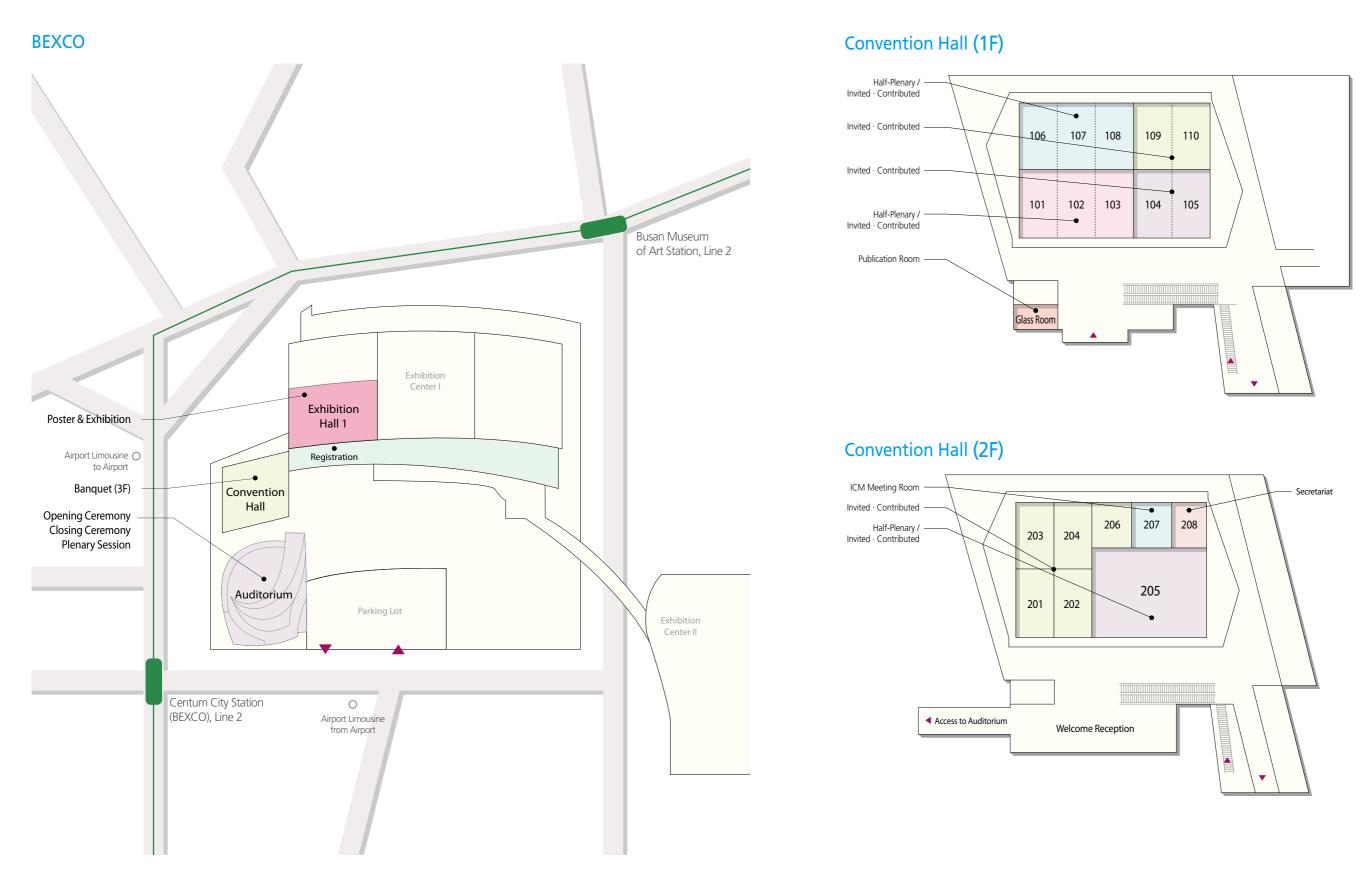
### PROGRAM AT A GLANCE



Program at-a-Glance

### VENUE LAYOUT

Venue Layout



### VENUE LAYOUT

Venue Layout

### KPS 60th Anniversary Nobel Laureate Public Lectures in Physics

In celebration of the KPS's 60th anniversary, the Public Lecture of Nobel laureate in Physics will be a program allowing the general public to learn about the interesting aspects of Physics. This public lecture program will be an exciting and enjoyable experience. Please do not miss this change to take the lectures on free.

This public lecture program will be an exciting and enjoyable experience for Korean audience interested in Physics.

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Date & Time: July 8 (Sun), 15:00~17:00
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    Venue: Auditorium
```

- Lecture: 1. Prof. Albert Fert Spintronics and its impact on inforamtion and communication technologies
  - 2. Prof. Klaus von Klitzing New applications of my Nobel prize

### Welcome Reception

You will experience a warm welcome from the host of ICM2012. All participants are highly welcome. Light refreshments and beer will be provided free of charge.

Date & Time: July 8 (Sun), 17:00~19:00
 Venue: Lobby of Convention Hall, 2F

### **Opening Ceremony**

ICM2012 will officially get started with a ceremony at Auditorium in Bexco. All registered participants are cordially invited to join us and celebrate the official opening.

Date & Time: July 9 (Mon), 08:30~09:00
 Venue: Auditorium

### ICM Award 2012

The IUPAP Magnetism Award and Néel Medal are presented every three years at the International Magnetism Conference to a scientist in recognition of an outstanding contribution to the field of magnetism. The award is sponsored by Elsevier Science B.V.. The IUPAP Young Scientist Medals in the field of magnetism are presented every three years at the International Magnetism Conference. The medals are sponsored by IUPAP.

Date & Time: July 9 (Mon), 08:30~09:00
 Venue: Auditorium

### Magnetism Award and Néel Medal 2012

- Sadamichi Maekawa (Japan Atomic Energy Agency, Japan)
   Heat and spin
- · Yoshinori Tokura (University of Tokyo, Japan) Electrodynamics of skyrmions

### Young Scientist Medals in the Field of Magnetism

• Suchitra E. Sebastian (University of Cambridge, UK) Nodal pocket revealed by quantum oscillations in an underdopedcuprate superconductor

### Banquet

Please join us to share an unforgettable evening. A delicious dinner with a traditional Korean music show, Samulnori (traditional percussion quartet) and 'B-boy', is combined to recreate the wonderful excitement performance.

Date & Time: July 12 (Thu), 19:00~20:30
 Venue: Room 301 (3F), Convention Hall

Have the opportunity to say farewell to friends and colleagues and to preview the next venue of ICM2015.

Date & Time: July 13 (Fri), 17:30~18:00
 Venue: Auditorium

### **Industrial Tour**

Participants those who applied for industrial tour of Samsung Heavy Industries Co., Ltd. should be gathered at 13:20 on July 11 (Wed) at the lobby. We will leave 13:30 on time. Application can be acceptable until 10:00 on July 9 at the information desk (Max. 100 people).

• Date & Time: July 11 (Wed), 13:30~18:00

· Destination: Samsung Heavy Industries Co., Ltd.

### CEREMONY & EVENT PROGRAM

### **Scientific Program**

ICM2012 program will consist of 7 plenary talks (1 hr), 14 half-plenary talks (45 min) and 135 invited talks (30 min). Over 2,000 contributed papers will be presented, 353 of them have been selected for oral presentation (15 min). In addition, poster presentations with ample time for discussion will be conducted. The official conference language is English. The program will focus on following topics;

1. Strongly Correlated Electron System (SCES)

- 2. Ouantum and Classical Spin Systems
- 3. Magnetic Structures and Interactions
- 4. Magnetization Dynamics and Micromagnetics
- 5. Spin-Dependent Transport
- 6. Spin Electronics
- 7. Magnetic Thin Films, Particles and Nanostructures
- 8. Soft and Hard Magnetic Materials and Their Applications
- 9. Novel Materials and Device Applications
- 10. Magnetic Recording and Memories
- 11. Measuring Techniques and Instrumentation
- 12. Industrial Applications
- 13. Interdisciplinary Topics

### Registration

All attendees will be required to wear the ICM2012 badge to access to all session.

- Venue: Lobby, Exhibition Hall 1
- Operation: July 8 (Sun) / 13:00~19:00
  - July 9 (Mon) ~ 12 (Thu) / 08:00~19:00
  - July 13 (Fri) / 08:00~13:00

### **Registration Fee**

	Category	On-site Registration
	Regular Participant	KRW 750,000
Registration Fee	Student/Retired Participant	KRW 400,000
	Accompanying Person	KRW 250,000
Banquet Fee (July 12)	All Participants	KRW 60,000

\* Participant's registration includes: Welcome reception, coffee breaks, admission to all scientific sessions, and a conference bag including abstract book \* Accompanying person's registration includes: Welcome reception, coffee breaks, banquet coupon and conference bag. Admission to scientific sessions is not included.

### Certificate of Attendance

The certificate of attendance is provided at the information desk on request or available for download via the website after the conference (www.icm2012.org)

### Internet Lounge

Internet access will be available during the conference at the Exhibition Hall where a PC computer pool will be provided.

### **Publication Room**

Authors can check the status of their manuscripts in the Publication Room, located in the Glass Room on the first floor of convention hall. Office hour of the Publication Room for authors will be as follows.

Operation Hours: July 9 (Mon), 15:00~17:00

July 10 (Tue), 12 (Thu), 10:00~12:00, 15:00~17:00 July 11 (Wed), 13 (Fri), 10:00~12:00

### **Oral Presentation Guideline**

Authors are expected to bring their presentations on their own laptop computer, and to have it powered up and ready to connect to the projector. Only standard PC-style VGA connections to the LCD projector will be supplied, therefore you must supply any required adaptor to connect up your computer.

### **Poster Presentation Guideline**

Posters are displayed in the Exhibition Hall 1 (1F). Poster should be posted by 08:30 and dismantled after 18:00 on the allotted date. The secretariat will not be held liable for any lost or damaged posters. All poster presenters are encouraged to be at their poster panels for discussion with participants during the time. All posters will be eligible for nomination for the best poster awards in each day.

- Venue:
- Exhibition Hall 1 • Operation: July 9, 10, 12, 13 (4days), 13:30~15:30
  - Affixation: 08:30~13:00 / Presentation: 13:30~15:30 / Removal: 18:00~19:00
- · Affixation & Removal: All presenters are requested to affix their posters and remove them after their presentation according to the above schedule. The secretariat will not be held liable for any posters lost or damaged.
- · Best Poster Award: There is a competition for the best poster. This award is given to recognize excellence in research and presentation. There will be two awards for each day. Each session chair is to nominate a single poster. The final review will be run by program executive members and the best awards are announced 30 minutes before the closing of the session.

### **Coffee Break**

Enjoy your break with a cup of coffee or tea that will be prepared as below;

	July 9 (Mon)	July 10 (Tue)	July 11 (Wed)	July 12 (Thu)	July 13 (Fri)
	10:00~10:20	10:30~11:00	10:30~11:00	10:30~11:00	10:30~11:00
Morning	Lobby Auditorium	1F, 2F Lobby Conventional Hall			
Afternoon	13:30~15:30	13:30~15:30	-	13:30~15:30	13:30~15:30
Afternoon	Exhibition Hall	Exhibition Hall	-	Exhibition Hall	Exhibition Hall

### **Complimentary Shuttle Service**

Shuttle bus will run between the conference venue and hotels. You may check the schedule and shuttle bus stop. The bus stop will be marked with banner stands at BEXCO.

Llatal	July 8	(Sun)	July 9	(Mon)	July 10	) (Tue)	July 11	(Wed)	July 12	2 (Thu)	July 1	3 (Fri)
Hotel	H→B	B→H	H→B	B→H	H→B	B→H	H→B	B→H	H→B	B→H	H→B	B→H
The Westin Chosun Busan	13:00	18:00	08:00	19:10	08:00	19:10	08:00	13:00	08:00	19:10	08:00	18:30
	14:00	19:00	08:30	19:40	08:30	19:40	08:30	14:00	08:30	21:00	08:30	19:00
Paradise Hotel Busan	13:00	18:00	08:00	19:10	08:00	19:10	08:00	13:00	08:00	19:10	08:00	18:30
Seacloud Hotel	14:00	19:00	08:30	19:40	08:30	19:40	08:30	14:00	08:30	21:00	08:30	19:00
Hotel Riviera Haeundae	13:00	18:00	08:00	19:10	08:00	19:10	08:00	13:00	08:00	19:10	08:00	18:30
	14:00	19:00	08:30	19:40	08:30	19:40	08:30	14:00	08:30	21:00	08:30	19:00
Novotel Hotel	13:00	18:00	08:00	19:10	08:00	19:10	08:00	13:00	08:00	19:10	08:00	18:30
Busan Ambassador	14:00	19:00	08:30	19:40	08:30	19:40	08:30	14:00	08:30	21:00	08:30	19:00
Lotte Hotel Busan	14:00	19:00	8:00	19:10	8:00	19:10	8:00	13:00	8:00	19:10	8:00	18:30
Hanwha Resort	13:00	18:00	08:00	19:10	08:00	19:10	08:00	13:00	08:00	19:10	08:00	18:30
	14:00	19:00	08:30	19:40	08:30	19:40	08:30	14:00	08:30	21:00	08:30	19:00

 $H \rightarrow B$ : From Hotel to Bexco,  $B \rightarrow H$ : From Bexco to Hotel

\* Participant who is staying in Haeundae Grand Hotel, please take a shuttle bus at the Westin Chosun Busan.

### GENERAL INFORMATION

\* Shuttle Service won't be provided for Haeundae Centum Hotel and ARPINA Buasn Youth Hostel located in walking distance.

### **GENERAL INFORMATION**

### **Cloak Room**

The cloakroom will be located in the exhibition hall so that you could keep your luggage during the conference.

### Water Station

One or two bottles of water will be provided each day. The Water Coupons will be given to all participants when you register. It will be contributed at the cloak room.

### Venue: BEXCO (Busan Exhibition Convention Center)

BEXCO, a landmark in the world-famous port city of Korea, has emerged as the most competitive exhibition and convention center in Northeast Asia.

Address 43 APEC-ro, Haeundae-gu, Busan 612-740, Ko	Corea
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- ۰Tel +82-51-740-7300
- Fax +82-51-740-7320
- Website www.bexco.co.kr

### ICM2012 Secretariat

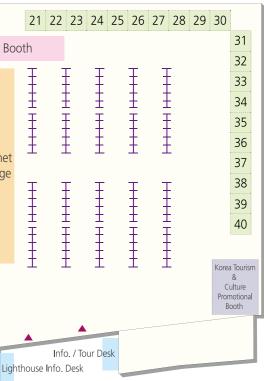
• Onsite Secretariat Office: Room 208, Convention Hall, Tel: 051-740-3730

· After the Conference: 1f Haeoreum Bldg., 16, Yeoksamro 17 gil, Gangnam-gu, Seoul 135-925, Korea Tel: +82-2-557-8422 Fax: +82-2-566-6087 Email: icm@icm2012.org Website: www.icm2012.org

Exhibition Hall (1F) 11 12 13 14 15 16 17 18 19 20 10 Exhibition Booth 9 8 Ŧ 7 6 5 4 3 2 1 Ŧ Ŧ Internet Lounge Т Т Poster Session Chair's Table Water Station Cloak Room

Exhibitor	Booth No.	Exhibitor	Booth No.	Exhibitor	Booth No.
Aaron. Co., Ltd.	16	Hyundai Motor Company	21, 22, 23	Quantum Design	24, 25
Ask Co.	9	ICM2015-Barcelona	40	ReC-SDSW(Seoul National University)	7
Carl Zeiss	27	Korea I.T.S.	17,18	Rigong International Inc.	13
ChangSung Co.	19, 20	Lake Shore Cryotronics	8	Semi-Ence Co., Ltd.	32
Coxem	31	Namotec	15	SmartTip BV	11
Cryogenic Ltd.	30	Nanomagnetics Instruments	10	Springer	36, 37
DGIST	1, 2	NT-MDT ANT Co.	26	Surface Systems Korea	12
Effucell	33	Oxford University Press	38	The Physical Society of Japan	39
ExaTech	29	PANalytical Korea	34	Top Techology Ltd.	14
HANARO @ KAERI	6	Park Systems Co.	35	UNIST	3
Hinds Instruments Inc.	28	Pohang Accelerator Laboratory	4,5		

### EXHIBITION



# Exhibition

### Registration

### Venue: Exhibition Hall 1

### Exhibition Schedule

 Installation: Shell Scheme Booth: July 8 (Sun) 08:00~ Exhibits & Display: July 8 (Sun) 18:00~24:00 • Exhibition Hours: July 9 (Mon)~12 (Thu), 09:00~19:00 July 13 (Fri), 09:00~16:00 • Removal: July 13 (Fri), 16:00~20:00

### SPONSOR





R I IGONG INTERNATIONALINC (주)이공교역

Solve Systems Korea

### ORAL PRESENTATION

Plenary Lecture • 28 Half-Plenary Lecture • 30 Invited & Contributed Presentation • 33

### PP01

#### Recent developments and emerging directions in spintronics Albert Fert

Unité Mixte de Physique CNRS/Thales, Palaiseau, and Université Paris-Sud, Orsay, France

Spintronics develops today in many promising directions. We focus on three topics. First: "Microwave generation by spin transfer", a field with fast recent advances anticipating short term applications. Second topic: "Spintronics with graphene or carbon nanotubes (CNT)", with fascinating long term prospects for "beyond CMOS". Finally: "Oxitronics", overview of results on devices associating magnetic, ferrogets for using spin transfer to induce magnetization precessions or gyrations of magnetic vortices. Large powers ( $\mu$ W range) and narrow emission widths ( $\approx 0.1$  MHz) can be today obtained at zero field by vortex gyration [1]. Synchronization is an additional way to increase the power and reduce the emission width [2], 2) Spintronics with graphene and CNT: Several concepts of spintronics (logic gates, "spin only logic circuits", etc) are based on spin transport in lateral channels between magnetic contacts. Recent experiments show the outstanding potential of graphene and CNT for spin transport to long distance (above 100  $\mu$ m) in such devices [3]. 3) Oxitronics: We illustrate the potential of oxitronics by results on tunnel junctions with freroelectric barriers: giant electro-resistance [4], interplay between ferroelectric and spin polarizations, prospects for ferroelectric memories and memristors.

[1] A. Dussaux, B. Georges, J. Grollier, V. Cros, A.V. Khvalkovski, A. Fukushima, M. Konoto, H. Kubota, K. Yakushiji, S. Yuasa, K. A. Zvezdi, K. Ando, A. Fert Nature Communications 1, 8, 2010. [2] A. Dussaux, A. V. Khvalkovskiy, J. Grollier, V. Cros, A. Fukushima, M. Konoto, H. Kubota, K. Yakushiji, S. Yuasa, K. Ando, and A. Fert Appl. Phys. Lett 98, 132506, 2011. [3] B. Dlubak, P. Seneer, A. Anane, M-B. Martin, C. Deranlot, B. Servet, S. Xavier, R. Mattana, M. Sprinkle, C. Berger, W. A. de Heer, F. Petroff, and A. Fert Nat. Phys. DOI: 10.1038 [4] V. Garcia, S. Fusi, K. Bouzehouane, S. Enouz-Vederme, N.D. Mathur, A. Barthélemy, M. Bibes, Nature 600, 81, 2009.

### **PP02**

Heat and spin Sadamichi Maekawa Japan Atomic Energy Agency, Japan

When metals and semiconductors are placed in a temperature gradient, the electric voltage is generated. This mechanism to convert heat into electric energy, the so-called Seebeck effect, has attracted much attention as the mechanism for utilizing wasted heat energy. The Seebeck effect is due to the entropy carried by the electric current so that it may be enhanced by the internal degrees of freedom of electrons, i.e., spin and orbital [1]. Ferromagnetic insulators are good conductors of spin current, i.e., the flow of electron spins [2]. When they are placed in a temperature gradient, generated is the spin voltage [3], i.e., spin accumulation. Once the spin voltage is converted into the electric voltage by spin Hall effect in attached metal films such as Pt, the electric voltage is obtained from heat. This is called the spin Seebeck effect. Here, we discuss the Seebeck effect and spin Seebeck effect based on the fluctuation-dissipation theorem and introduce a variety of the devices.

[1] S. Maekawa et al, Physics of Transition Metal Oxides (Springer, 2004). [2] S. Maekawa: Nature Materials 8, 777 (2009). [3] Concept in Spin Electronics, eds. S. Maekawa (Oxford University Press, 2006).

### PP03

#### Heavy electrons and superconductivity Zachary Fisk\* University of California, Irvine, USA

Heavy electron materials are the only class of materials in which we know where to look for superconductivity. We discuss the relation between single ion and dense Kondo heavy electron physics, what sets the energy scale in these materials, and the ways in which their low temperature exotic superconductivity is a prototype for all highly correlated electron superconductivity.

### PP04

### Correlated electrons in quantum hall systems

Klaus V. Klitzing\* Max-Planck-Institut fur Festkorperforschung, Heisenbergstr:1, D-70569 Stuttgart, Germany

The narrow energy bands of a two-dimensional electron system in strong magnetic fields lead to characteristic electron-electron interaction phenomena related to quantum Hall physics. The talk summarizes some of the most prominent correlation phenomena in quantum Hall devices including exciton condensation/superfluidity in bilayers and roton minima in the excitation spectrum of fractional quantum Hall states. \* In cooperation with W.Dietsche, X. Huang, I. Kukushkin, and J.H. Smet

### PP05

### From single-atom magnetometry to tailored nanomagnets and atomicscale spintronic devices Roland Wiesendanger\*

Institute of Applied Physics, University of Hamburg, Germany

The developments of novel magnetic materials as well as spin-based electronics are hot topics of current research in magnetism. Both research fields could profit tremendously from atomic-scale insight into magnetic properties and spin-dependent interactions at the atomic level. Based on the development of spin-polarized scanning tunneling microscopy (SP-STM) we have recently established the novel method of single-atom magnetometry which allows the measurement of magnetization curves and the determination of magnetic moments on an atom-by-atom basis. While the sensitivity level of single-atom magnetometry is below one Bohr magneton, it can easily be combined with the atomic-resolution imaging and manipulation capabilities of conventional STM, thereby offering a novel approach towards a rational material design based on the knowledge of the atomic-level properties and interactions within the solid state. Moreover, an atom-by-atom design and realization of all-spin logic devices has recently been demonstrated by our group based on the combined knowledge derived from surface physics, nanoscience, and magnetism. Alternatively, self-assembly of atomic magnetic chains on nanostructured substrates has been employed in order to create model-type systems for atomic-scale information transfer based on the concept of vector-spin chirality.

### PP06

### Multiferroic vortex network with $Z_{2x}Z_3$ symmetry Sang-wook Cheong\*

Rutgers University, USA

Hexagonal REMnO<sub>3</sub> (RE= rare earths) with RE=Ho-Lu, Y, and Sc, exhibit unique patterns of trimerization and ferrelectric domains, which meet in cloverleaf arrangements that cycle through all six-possible domain configurations. Occurring in pairs, the cloverleafs can be viewed as vortices and antivortices, in which the cycle of domain configurations is reversed. Vortices and antivortices are topological defects: even in a strong electric field they won't annihilate readily. These ferroelectric vortices/antivortices are also associated with a vortex configuration of magnetism. The seemingly irregular configurations of a zoo of multiferroic vortices and antivortices in h-REMnO<sub>2</sub> can be neatly analyzed in terms of graph theory and this graph theoretical analysis reflects the nature of self-organized criticality in complexity phenomena. These numerous multiferroic vortices/antivortices can be understood as an arrested Kosterlitz-Thouless phase. We have also discovered the emergence of Z2×Z3 symmetry in the seemingly-random network of numerous ferroelectric vortices and electric poling or self-poling due to a surface charge boundary condition induces global topological condensation through breaking of the Z2 part of the Z2×Z3 symmetry. The opposite process of restoring the Z2 symmetry can be considered as topological evaporation.

### **PP07**

#### Graphene's magnetism Andre Geim University of Manchester, United Kingdom

Graphene - a free-standing atomic plane of graphite - has turned out to be a wonder material. Like a magnet, it has attracted a large and still rapidly growing research community. After a short introduction describing reasons for this interest, I will focus on the main topics of this conference (namely, magnetism and correlations) and overview our recent work on magnetic properties of graphene, which are still poorly understood, and pronounced correlation effects observed in graphene double-layer heterostructures. As for magnetism, I will discuss the spin Hall effect [1] and (para) magnetism of pristine and defected graphene [2,3]. As for correlations, I plan to touch on renormalization of graphene's linear spectrum [4] and Coulomb drag between two Dirac-like systems in the limit of maximum interaction (unpublished but see [5]).

 D. Abanin et al, Science 332, 328 (2011) 2. M. Sepioni et al, PRL 105, 207205 (2010) 3. R. Nair et al, Nature Phys. 8, 199 (2012) 4. D. Elias et al, Nature Phys. 7, 701 (2011) 5. L. Ponomarenko et al, Nature Phys. 7, 958 (2011)

### PLENARY LECTURE



### HALF-PLENARY LECTURE

### HP11

### Heavy fermions and unconventional superconductivity in high-Quality single crystals of rare earth and actinide compounds

Yoshichika Onuki<sup>1</sup>, Rikio Settai<sup>1</sup>, Tetsuya Takeuchi<sup>2</sup>, Kiyohiro Sugiyama<sup>1</sup>, Fuminori Honda<sup>1</sup>, Yoshinori Haga<sup>3</sup>, Etsuji Yamamoto<sup>3</sup>, Tatsuma Daruma Matsuda<sup>3</sup>, Naoyuki Tateiwa<sup>3</sup>, Dai Aoki<sup>4</sup>, Ilya Sheikin<sup>5</sup>, Hisatomo Harima<sup>6</sup> and Hiroshi Yamagami<sup>7</sup> <sup>1</sup> Graduate School of Science, Osaka University, Japan <sup>2</sup> Low temperature Center, Osaka University, Japan <sup>3</sup> Advanced Science Research Center, JAEA, Japan <sup>4</sup> IMR, Tohoku University, Japan <sup>5</sup> LNCMI-G, CNRS, France <sup>6</sup> Department of Physics, Kyoto Sangyo University, Japan

The f electrons in rare earth and actinide compounds exhibit variety of characteristic properties including heavy fermions and unconventional superconductivity. Fermi surface properties in rare earth and actinide compounds such as CeSn3, USi3, NpGe3 and PuIn3 are clarified by the de Haas-van Alphen experiments on the basis of the results of energy band calculations. An abrupt nonlinear increase of magnetization, namely metamagnetic behavior is found in the heavy fermion compounds including YbT2Zn20 (T: Co, Rh, Ir). An effect of pressure on the electronic state of YbIr2Zn20 and antiferromagnets CeRhIn5 and CeIrSi3 are also studied in magnetic fields. The electronic instability including unconventional superconductivity occurs at about 2.4 GPa in CeRhIn5 and 2.6 GPa in CeIrS3.

### **HP12**

### Effects of spin-orbit-coupling in the electronic structures of 5d transition metal oxides

T. W. Noh

Physics and Astronomy, Seoul National University, Korea

In 5d transition metal oxides, the magnitude of spin-orbit (SO) coupling becomes comparable to those of other fundamental interactions, such as electron-phonon and on-site Coulomb interactions. Recent investigations of 5d transition metal oxides indeed reveal that the SO coupling can modify their electronic and magnetic structures significantly. Especially, it produces a novel quantum state that has never been observed in 3d and 4d transition metal oxides. Namely, the large SO coupling combined with the on-site Coulomb repulsion results in an unusual Jeff=1/2 Mott insulating state in Sr2IrO4. By changing the dimensionality in Ruddlesden-Popper series, the irridates change from insulator (n=1 and 2) to metal (n=∞). Further studies on double perovskite A2FeReO6 and perovskite (Ca,Sr)IrO3 thin films showed that their electronic structures are indeed determined by subtle interplay of the SO coupling, electron correlation, and lattice distortion. We will present the interesting role of SO coupling by combining the spectroscopic tools with first-principles band calculations. In addition, the layered honeycomb lattice Na2IrO3 and pyrochlore A2Ir2O7 were recently suggested as possible topological insulators, though topological insulators with transition metal d electrons have not been fully investigated. Spectroscopic result on these irridate systems will be also discussed.

### HP21

#### Magnetization dynamics of rashba ferromagnet

Kyoung-whan Kim<sup>1</sup>, Soo-man Seo<sup>2</sup>, Jisu Ryu<sup>1</sup>, Jung-hwan Moon<sup>2</sup>, Kyung-jin Lee<sup>2</sup> and Hyun-woo Lee<sup>1</sup>\*

<sup>1</sup> Department of Physics, POSTECH, Korea

<sup>2</sup> Department of Materials Science and Engineering, Korea University, Korea

Rashba spin-orbit coupling (RSOC) arises generically when structural inversion symmetry is broken. Surface electronic structure of heavy atomic elements such as Bi/Ag alloy [1] is known to exhibit large RSOC with characteristic RSOC parameter of the order of 1 eV•A. Topological insulator [2,3] is another class of material with large RSOC. Recently it was reported [4] that ultrathin magnetic layers may be subject to large RSOC when in contact with proper heavy metallic elements in strongly asymmetric environment. In this talk, we explore magnetization dynamics of such ultrathin magnetic layers with strong RSOC. It will be demonstrated that RSOC induces the deviation of conduction electron spin direction from local magnetization direction and thus modifies the spin torque [5,6] and current-driven magnetization dynamics properties considerably. We also discuss the phenomenon of the spin-dependent electric field induction [7,8,9] by magnetization dynamics. It will be illustrated that RSOC can strengthen the spin-dependent electric field more than one order of magnitude, so that the field can induce spin current sufficiently large enough to modify the magnetization dynamics itself. Thus ultrathin magnetic layers with strong RSOC are good systems to test various aspects of RSOC effects.

C. R. Ast et al., Phys. Rev. Lett. 98, 186807 (2007). [2] K. Ishizaka et al., Nature Mater. 10, 521 (2011). [3] M.
 S. Bahramy, B.-J. Yang, R. Arita, and N. Nagaosa, arXiv:1110.6846. [4] I. M. Miron et al., Nature Mater. 9, 230 (2010). [5] J. C. Slonczewski, J. Magn. Mag. Mater. 159, L1 (1996). [6] L. Berger, Phys. Rev. B 54, 9353 (1996).
 [7] L. Berger, Phys. Rev. B 51, 1572 (1986). [8] G. E. Volovik, J. Physc C 20, L83 (1987). [9] S. E. Barnes and S. Maekawa, Phys. Rev. Lett. 98, 246601 (2007).

### 30 | The 19<sup>th</sup> International Conference on Magnetism

### HP22

#### Giant Ising anisotropy and hastaticorder in URu<sub>2</sub>Si<sub>2</sub> (1,2) Piers Coleman\*

Dept of Physics and Astronomy, Rutgers University, United Kingdom

The hidden order that develops below 17.5K in the heavy fermion compound URu<sub>2</sub>Si<sub>2</sub> has eluded identification for twenty five years. We show that the recent observation of Ising quasiparticles in URu<sub>2</sub>Si<sub>2</sub> suggests a novel two-component order describing hybridization between electrons and the Ising 5I<sup>2</sup> states of the uranium atoms. This ``hastatic order" is distinct from conventional magnetism, breaking both single and double time-reversal symmetry operations, mixing states of different Kramers parity. It accounts for the magnetic anomalies in torque magnetometry and the pseudo-Goldstone mode in neutron scattering. Hastatic order is predicted to induce a basal-plane magnetic moment of order \$0.01\mu\_1B}S, a gap to longitudinal spin fluctuations that vanishes continuously at the first-order antiferromagnetic transition, a giant Ising anisotropy in the non-linear susceptibility and a narrow resonant nematic feature in the scanning tunneling spectra.

(1) R. Flint, P. Chandra and P. Coleman to be published. (2) Work supported by the Simons Foundation and NSF grant 0907179. (3) M. M. Altarawarneh et al, Phys. Rev. Lett. 108, 066407 (2012).

### HP31

### Electrodynamics of skyrmions

Department of Applied Physics, University of Tokyo, Japan

Skyrmions are nanometric spin-swirling vortex-like objects in solids, typically observed in the helimagnets of B20 type (FeSi type) transition-metal silicides and germanides. Due to their topological nature, the spin chirality or skyrmion number can give rise to the fictitious magnetic field acting on conduction electrons; in the strong coupling limit one skyrmion works as a single flux quantum. There, the topological Hall effect of conduction electron is observed in proportion to skyrmion density, while as its counteraction the Hall motion of skyrmions is anticipated in addition to the current-driven sliding motion. Here, we overview the late progress in experimental observations of skyrmions and skyrmion lattices in bulk and thin film forms as well as their spectacular topological transport phenomena. The materials family is now being extended to some multiferroic oxides, in which electric-field drive of the skyrmion would be possible.

### HP32

#### Spin and charge fluctuations in cuprate superconductors Bernhard Keimer

Max Planck Institute for Solid State Research, Germany

I will present recent inelastic neutron scattering and resonant inelastic x-ray scattering (RIXS) data that throw new light on the role of low-energy spin and charge fluctuations in the mechanism of high-temperature superconductivity and its interplay with competing spin- and charge-density-wave order. I will also discuss parallels and contrasts with iron pnictide superconductors.

M. Le Tacon, G. Ghiringhelli, J. Chaloupka, M. Moretti Sala, V. Hinkov, M.W. Haverkort, M. Minola, M. Bakr, K. J. Zhou, S. Blanco-Canosa, C. Monney, Y. T. Song, G. L. Sun, C. T. Lin, G. M. De Luca, M. Salluzzo, G. Khaliullin, T. Schmitt, L. Braicovich and B. Keimer, Nature Physics 7, 725 (2011).

### HP41

#### The spin on domain walls! Stuart Parkin IBM Almaden Research Center, USA

The formation and manipulation of magnetic domain walls (DWs) is of considerable interest both from a scientific as well as a technological perspective. A number of very interesting and potentially useful memory and logic devices based on the controlled manipulation of DWs in magnetic nano-elements have recently been proposed and are under intensive investigation. Using spin torque transfer from spin-polarized currents a series of domain walls can be moved in lock-step at high speed along magnetic nano-wires, enabling the Racetrack Memory [1]. We contrast the current induced motion of DWs in nano-wires with in-plane and out-of-plane magnetization and demonstrate an additional interface induced driving force in the latter case, which leads to very high domain wall velocities of up to almost 1 km/sec. Even higher efficiencies result from temperature gradient induced spin currents which we show are ~1,000 times more efficient in moving DWs than electrically generated spin polarized currents. Many of these devices depend critically on the detailed structure of the DW: we discuss methods by which the chirality and polarity of a vortex DW can be deterministically controlled and we use this to create a chirality based biplexer and ultra-stable domain wall topological bound states.

[1] S.S.P. Parkin, US Patent 6,834,005 (2004); S.S.P. Parkin et al., Science 320, 190 (2008); S.S.P. Parkin, Scientific American (June 2009); M. Hayashi, L. Thomas, R. Morya, C. Rettner and S.S.P. Parkin, Science 320, 209 (2008); L. Thomas et al., Science 330, 1810 (2010); X. Jiang et al. Nat. Comm. I:25 (2010) and Nano Lett. 11, 96 (2011); L. Thomas, M. Hayashi, R. Morya, C. Rettner and S.S.P. Parkin, Nat. Commun. (accepted).

### HP42

### Orbitronics in silicon

Gabriel Aeppli University College London, United Kingdom

We describe the control and observation of coherent superpositions of defect orbitals in silicon using both scanning tunneling microscopy and pulsed THz radiation generated by the Dutch-UK free electron laser FELIX. The results are contrasted with those for the microwave control of defect spins, and implications for future silicon-based information processing are discussed.

PNAS 105 10649-10653 (2008) Nature 465 1057-61 (2010) Nature Materials 9 725-9 (2010)

### HP51

#### 360 degree domain walls in magnetic nanowires Caroline A Ross

Department of Materials Science and Engineering, Massachusetts Institute of Technology, USA

Domain walls are important in a wide range of magnetic memory and logic devices. Most work has focussed on isolated 180 degree domain walls in magnetic wires, but impingement of two 180 degree walls can create 360 degree walls which exhibit facinating behavior. 360, 540 and higher order walls are created by successive injection of 180 degree domain walls is qualitatively different from that of 180 degree walls into a wire or during reversal of a thin film ring. Current-driven motion of 360 degree domain walls is qualitatively different from that of 180 degree walls. The velocity of a 360 degree walls cases with dc current but is independent of applied field along the stripe, and an annihilation process occurs at a field-dependent critical current density. 360 degree walls oscillate at a characteristic field-tunable frequency in the gigahertz range. The stray field from a 360 degree wall in one layer of a multilayer can affect the reversal of an adjacent soft magnetic layer, allowing the wall to be used as a gate for the reversal of the soft layer. 360 degree walls can be created, moved, annihilated to deliver spinwave energy, and detected, and could serve as a data token in a magnetoelectronic device.

### HP52

### Perpendicular CoFeB-MgO for spintronics devices

Hideo Ohno\* CSIS/RIEC, Tohoku University, Japan

bcc(100) CoFeB-(100) MgO is a preferred system for spintronics applications, because it offers a large tunnel magnetoresistance (TMR) of > 100% in magnetic tunnel junctions (MTJs) through symmetry filtering of wavefunctions, which is required for nonvolatile memory cells. Integration of high performance MTJs in the CMOS back-end not only enables non-volatile, high density, and fast stand-alone and embedded RAMs, but also makes it possible to realize nonvolatile logic-in-memory CMOS VLSIs [1]. To this end, we showed that a perpendicular interface anisotropy at the CoFeB-MgO interface [2, 3], strong enough (Ki = 1.3 mJ/m<sup>2</sup>) to overcome the demagnetization, can be used to realize 40 nm perpendicular MgO-CoFeB MTJs with high TMR (>100 %) and low switching current of 49  $\mu$ A [3]. We further showed that "activation volume" of the system plays an important role in determining the thermal stability of the MTJs [4]. It was also shown that the CoFeB-MgO system withstands a high annealing temperature of 400 °C required for logic integrated circuit applications [5]. If time allows, I will touch upon current-induced domain wall motion and related phenomena in perpendicular CoFeB-MgO films [6, 7], where perpendicular anisotropy is required to reduce the critical current density for domain wall motion.

This work was supported by the FIRST program from JSPS. [1] S. Ikeda, et al. IEEE Trans. Electron Devices, 54, 991, 2007. [2] M. Endo, et al. Appl. Phys. Lett., 96, 212503, 2010. [3] S. Ikeda, et al. Nature Mat., 9, 721, 2010. [4] H. Sato, et al. Appl. Phys. Lett. 99, 042501, 2011. [5] H. D. Gan et al. Appl. Phys. Lett. 99, 252507, 2011. [6] S. Fukami et al. Appl. Phys. Lett. 98, 082504, 2011. [7] T. Suzuki et al. Appl. Phys. Lett. 98, 142505, 2011.

### HP61

#### Geometric enhancement of low field magnetoresistance in silicon Xiaozhong Zhang

Department of Materials Science and Engineering, Tsinghua University, China

The magnetic sensing industry are greatly dependent on the use of both GMR and TMR devices. Both of them are made of magnetic/rare earth materials. However, the rare earth materials are getting more and more difficult to obtain. Magnetoresistance (MR) reported in some non-magnetic semiconductors particularly silicon has triggered considerable interest owing to the large magnitude of the effect. Here we showed that MR in lightly doped n-Si can be significantly enhanced by introducing a p-n junction and proper design of the carrier path [1]. We designed an MR device whose room-temperature MR ratio reaching 30% at 0.065 T, 5000% at 0.4 T, and 2000% at 1.2T, respectively, approaching the performance of commercial MR devices. The combination of high sensitivity to low magnetic fields and large high-field response should make this device concept attractive to the magnetic field sensing industry. Moreover, because our device is based on a conventional silicon platform, it should be possible to integrate it with existing silicon devices and so aid the development of silicon-based magnetoelectronics. Also combining MR devices and semiconducting devices in a single Si chip may lead to some novel devices.

[1] Caihua Wan, Xiaozhong Zhang, et al., Nature, 477, 304 (2011).

### HP62

### Artificial spin ice systems: exploring frustration and emergent magnetic monopoles with nanomagnets

Laura Heyderman\*

Paul Scherrer Institute, Switzerland

Artificial spin ice systems, consisting of two-dimensional arrangements of dipolar coupled singledomain nanomagnets, are in the focus of scientific interest since they allow the study of the effects of frustration. Our work has concentrated on artificial kagome spin ice, with elongated nanomagnets forming an array of hexagonal rings, using synchrotron x-ray photoemission electron microscopy to directly image the magnetic states. Focussing on finite building blocks comprising one, two and three hexagonal rings, allows the full characterization of the energy levels and indicated that an effective thermal anneal via demagnetisation does not lead to the ground state for the larger structures [1]. More recent observations demonstrate the existence of emergent magnetic monopoles in a quasi-infinite nanomagnet array [2]. In an applied magnetic field, monopole-antimonopole pairs nucleate and separate in an avalanche-type manner along one-dimensional Dirac strings consisting of overturned dipoles, and the behaviour is quantitatively explained by Monte Carlo simulations. This work opens the way to making use of the multiple states in coupled nanomagnet systems [3] and the controlled manipulation of magnetic charges that may lead to new spintronic devices.

[1] E. Mengotti, L.J. Heyderman, A. Fraile Rodriguez, A. Bisig, L. Le Guyader, F. Nolting, and H.B. Braun, Phys. Rev. B 78, 144402 (2008) [2] E. Mengotti, L.J. Heyderman, A. Fraile Rodriguez, F. Nolting, R.V. Hugli, H.B. Braun Nature Physics 7, 68 (2011) [3] L.J. Heyderman, T. Jung, E. Mengotti, A. Bisig, A. Fraile Rodriguez, F. Nolting, H.B. Braun, T. Schreff United States Patern US 8,085,578 B2, 27.12.2011

### HALF-PLENARY LECTURE

### HP71

**Topological Insulators** Shoucheng Zhang

Physics, Stanford University, USA

In this talk I plan to give an overview of the recent progress on topological insulators. I will focus on the interplay between topological and magnetic orders in novel materials. and discuss the topological magneto-electric effect.

X. L. Qi, S. C. Zhang, Phys. Today 63, 33 (2010). X. L. Qi, S. C. Zhang, Rev. Mod. Phys. 83, 1057

### **HP72**

#### Heusler compounds: from semiconductors to spintronics Claudia Felser\*

Max Planck Institute for Chemical Physics of Solids, Germany

Heusler compounds are a remarkable class of intermetallic materials with 1:1:1 (often called Half- Heusler) or 2:1:1 composition comprising more than 1500 members [1]. New properties and potential fields of applications emerge constantly; the prediction of topological insulators is the most recent example [2]. Surprisingly, the properties of many Heusler compounds can easily be predicted by the valence electron count or within a rigid band approach. The subgroup of more than 250 semiconductors is of high relevance for the development of novel materials for energy technologies. Their band gaps can readily be tuned from zero to 4 eV by changing the chemical composition. Thus, great interest has been attracted in the fields of thermoelectrics and topological insulator research. Ternary materials based on multifunctional properties, i.e. the combination of two or more functions such as superconductivity and topological edge states will revolutionize technological applications. The wide range of the multifunctional properties of Heusler compounds is reflected in extraordinary magnetooptical, magneto-electronic, and magneto-caloric properties. Tetragonal Heusler compounds Mn2YZ as potential materials for STT applications can be easily designed by positioning the Fermi energy at the van Hove singularity in one of the spin channels [3].

[1] Simple Rules for the Understanding of Heusler Compounds, Tanja Graf, Stuart S. P. Parkin, and Claudia Felser, Progress in Solid State Chemistry (2011), doi:10.1016/ j.progsolidstchem.2011.02.001, invited review [2] Tunable Multifunctional Topological Insulators in Ternary Heusler Compounds, S. Chadov, X.-L. Qi, J Kubler, G. H. Fecher, C. Felser, S.-C. Zhang, Nature Mat. 9 (2010) 541. arXiv: 1003.0193

### AA01

#### Orbital and spin states of bi-layered manganites La<sub>2-2x</sub>Sr<sub>1+2x</sub>Mn<sub>2</sub>O<sub>7</sub> Jae-Hoon Park

Pohang University of Science and Technology, Korea

The bi-layer manganites La<sub>2.2</sub>, Sr<sub>1+2</sub>, Mn<sub>2</sub>O<sub>2</sub> display variety of ordering behaviors involving charge, spin, and orbital degrees of freedom and exhibits various intriguing physical properties. Here I will present soft x-ray scattering study results on the spin states of bi-layer manganites with x = 0.5 and x = 0.3. The x = 0.5 half-doped system exhibits two types of antiferromagnetic spin orders, A-type and CE-type, which both involves the CE-type charge and orbital orders. The results show coupling behaviors between the spin and orbital ordering and also between different spin ordered states. On the other hand, the x = 0.3 system, known as a colossal magnetoresistance manganite shows a large magneto-elastic response. We found that the spin state is coupled with the orbital state and the magnetic field affects both states resulting in the large magneto-elastic response.

### AA02

#### First-principles calculation of multiferroic bilayer manganite Kunihiko Yamauchi1\* and Silvia Picozzi2 <sup>1</sup> ISIR-Sanken, Osaka University, Japan <sup>2</sup> CNR-SPIN, L'Aquila, Italy

'Improper multiferroics', materials where ferroelectricity is driven by spin ordering, charge ordering or orbital ordering, constitute a playground for the physics of crosscorrelation between the different long-range orderings [1]. Recently, a complex mechanism of ferroelectricity driven by spin- and charge-orderings has been proposed in hole-doped manganites, supported by theoretical studies [2]. In such system, it is considered that slightly charge-disproportionated two magnetic ions (i.e. Mn<sup>3</sup> and Mn4+) form dimers via double exchange interaction, resulting in sizable electric dipoles. In the particular case of bilayer-manganite Pr(Sr<sub>0.1</sub>Ca<sub>0.9</sub>)<sub>2</sub>Mn<sub>2</sub>O<sub>7</sub>, it has been experimentally suggested that the transition between two different CO phases is accompanied by the rotation of orbital stripes, in turn related to ferroelectricity [3]. In this paper, we provide insights into the cross-correlation phenomena and the ferroelectric instability in the bilayer manganite via first-principles approaches. Along this line, we also discuss the magnetic property and the stability of charge- and orbitalorderings

[1] S. Picozzi, and C. Ederer, J. Phys.: Cond. Mat. 21, 303201 (2009), [2] D. V. Efremov, J. van den Brink, and D. Khomskii, Nature Matt. 12, 856 (2004). [3] Y. Tokunaga et al., Nature Mater. 5, 937 (2006)

#### AA03

#### Coupling between lattice and spin degrees of freedom in multiferroic h-RMnO<sub>3</sub>

Xavier Fabreges1\*, Sylvain Petit2 and Isabelle Mirebeau2 LNCMI, Toulouse, CNRS, France <sup>2</sup> Laboratoire Leon Brillouin, CEA Saclay, France

Hexagonal RMnO3 forms a class of triangle-based multiferroic materials, which have been widely studied in the recent years [1]. Magnetic frustration combined with a striking magneto-elastic coupling seems to be at the origin of their properties, a cocktail that has a strong potential for novel physics. Here, we report on neutron scattering experiments carried out to study the structure and spin dynamics in this series of compounds. The spinwave spectra can be accounted for by numerical calculations, giving precious informations on exchange and anisotropy parameters [2-4]. Meanwhile, the 120° Neel order due to geometric frustration is accompanied by an isostructural transition [4.5]. This transition is a systematic feature in this series as shown by our high-resolution neutron diffraction and inelastic neutron scattering. In addition, we have established a correlation between the atomic positions, the type of magnetic structure, and the nature of the spin waves, depending on the compound and its magnetic structure. The key parameter is the Mn position within the triangular plane, which tunes the sign of the exchange interaction [4]. Our inelastic neutron scattering experiments also reveal a strong coupling of the Mn sublattice with the rare-earth magnetic moments. The cases of HoMnO3 and YbMnO3 are presented.

[1] Cheong S.W and Mostovoy M. Nature materials, 6, 13 (2007) and references therein, [2] S. Petit el al. PRL 99, 266604 (2007), [3] S. Pailhes et al, PRB 79, 134409 (2009), [4] X. Fabreges et al, PRL 103, 067204 (2009), [5] S. Lee et al. Nature 451,805 (2008)

### AA04

### Magnetoelectric coupling in hematite amplified by the collective transition

J. L. Musfeldt<sup>1</sup>\*, P. Chen<sup>1</sup>, N. Lee<sup>2</sup>, S. Mc Gill<sup>3</sup> and S. W. Cheong<sup>2</sup>

- <sup>1</sup> Department of Chemistry, University of Tennessee, USA
- <sup>2</sup> Department of Physics, Rutgers University, USA National High Magnetic Field Laboratory, USA

We investigated the magneto-optical properties of alpha-Fe2O3 in order to understand the interplay between charge and magnetism in a model transition metal oxide. We discovered that hematite appears more red in applied magnetic field than in zero field conditions, an effect that is amplified by the presence of the spin flop transition. Analysis of the exciton pattern on the edge of the d-d color band reveals C2/c monoclinic symmetry in the high field phase. These findings advance our understanding of magnetoelectric coupling away from the static limit and motivate spectroscopic work on other iron-based materials under extreme conditions.

### AA05

### Time dependence of multiferroic switching

Max Baum<sup>1</sup>, Thomas Finger<sup>1</sup>, Simon Holbein<sup>1</sup>, Jonas Stein<sup>1</sup>, Jeannis Leist<sup>2</sup>, Gotz Eckold<sup>2</sup>, Paul Steffens<sup>3</sup>, Arno Hiess<sup>3</sup>, Karin Schmalzl<sup>3</sup>, Petra Becker<sup>4</sup>, Ladislav Bohaty<sup>4</sup> and Markus Braden<sup>1</sup>

<sup>1</sup> II. Physikalisches Institut, Universitat zu Koln, Germany

<sup>2</sup> Institut fur Physikalische Chemie, Georg-August-Universitat Gottingen, Germany

<sup>3</sup> Institut Laue-Langevin (ILL), Grenoble, France

<sup>4</sup> Institut fur Kristallographie, Universitat zu Koln, Germany

Neutron scattering with spherical polarization analysis gives direct access to the chiral components of the magnetic structure which is directly linked to the electric polarization and thus may be controlled by an electric field. We applied stroboscopic techniques in order to investigate how fast the magnetic chirality adapts to an instantaneously switched electric field in MnWO4 finding quite long relaxation times in the range of 10 milliseconds. We will discuss the dependence of the multiferroic relaxation on electric field strength, sample thickness and temperature. Furthermore we present the first results on controlling ferrotoridal domains in LiFeSi<sub>2</sub>O<sub>6</sub> by vertical magnetic and electric fields

### AB01

### Quantum criticality, non-fermi liquid and unconventional superconductivity

Oimiao Si

Department of Physics and Astronomy, Rice University, USA

Competing orders frequently arise in strongly correlated electron systems, and the accompanying quantum criticality provides a route towards both non-Fermi liquid behavior and unconventional superconductivity. Among the many important questions is whether the theoretical description of quantum criticality conforms to, or goes beyond, the Landau framework of order-parameter fluctuations. Using heavy fermions as a concrete setting, I will summarize the theoretical aspects of local quantum criticality, which is based on the "beyond-Landau" physics of Kondo destruction, as well as the experimental evidence for this picture [1]; the latter includes anomalous dynamical scaling, multiple scales of collapsing energies, and Fermi surface jump. I will also discuss the theoretical reasoning for the emergence of novel magnetic phases and the corresponding global phase diagram [2,3], and attempt to place various quantum-critical heavy fermion compounds in this global phase diagram. Finally, I will outline some promising issues to explore in the future, including the interplay among electron localization (in the form of Kondo destruction in the present context), magnetism and unconventional superconductivity, both in connection with existing experiments [4,5] and as a general theoretical issue.

[1] Q. Si and F. Steglich, Science 329, 1161 (2010). [2] Q. Si, Physica B378, 23 (2006); Phys. Status Solidi B247, 476 (2010). [3] P. Goswami and Q. Si, Phys. Rev. Lett. 107, 126404 (2011). [4] T. Park et al., J. Phys.: Condens. Matter 23, 094218 (2011); H. Shishido et al., J. Phys. Soc. Jpn. 74, 1103 (2005). [5] O. Stockert et al., Nature Phys. 7, 119 (2011).

### AB02

#### Ternary compounds with ZrFe<sub>4</sub>Si<sub>2</sub> structure type: A new playground for ferromagnetic and antiferromagnetic quantum criticality

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The ZrFe<sub>4</sub>Si<sub>2</sub> structure type is promising for the search of quantum critical systems because both the Fe and the Zr sublattices present a quasi-one-dimensional character with geometrical frustration. In the last years we prepared several compounds crystalizing in this structure and found two very different kinds of systems close to a quantum critical point (QCP) separating a magnetic ordered ground state from a non-ordered one: in YbNi<sub>4</sub>P<sub>2</sub> a strong Kondo interaction acting on the magnetic trivalent Yb-lons (Ni atoms are non-magnetic) results in a strongly reduced ferromagnetic (FM) ordering temperature of only 170 mK, locating this system close to a FM QCP. In contrast in the AFe<sub>4</sub>X<sub>2</sub> family of compounds (A = Sc, Y, Lu, Zr, X = Si, Ge) it is the 3d metal Fe which is magnetic while the A-element is non-magnetic. Although the energy scale of magnetic interactions in these compounds state with a rather low T\_N of 32 K in LuFe<sub>4</sub>Ge<sub>2</sub> to a ground state dominated by critical fluctuations in ZrFe<sub>5</sub>Si<sub>2</sub>. We studied these compounds using different kinds of techniques and shall present and discuss our results.

### **AB03**

### Anomalous thermoelectric effects in the heavy fermion superconductor Ce,PdIn\_s

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Ce<sub>2</sub>PdIn<sub>8</sub> is a recently discovered [1,2] ambient-pressure heavy-fermion superconductor that in many aspects closely resembles the structurally related compound CeCOIn<sub>5</sub>. In particular, the superconducting gap that opens at Tc = 0.7 K [3-5] exhibits a nodal character and a magnetic-field-induced quantum critical point (QCP) resides near the upper critical field Hc2(0) = 2.4 T [3,4]. Moreover, the necessary conditions for the formation of modulated FFLO phase are met in the superconducting state of Ce<sub>2</sub>PdIn<sub>8</sub>, and this unique phase possibly appears in the low-temperature-high-field comer of the H-T phase diagram [3]. The reported Nemst effect (v) and thermoelectric power (S) studies were performed for Ce<sub>2</sub>PdIn<sub>8</sub> with the main aim at probing QCP in the temperature region, in which distinct non-Fermi liquid features were previously revealed in the electrical resistivity and heat capacity data [1-5]. In a manner markedly similar to CeCoIn<sub>5</sub>, the Nemst coefficient of Ce<sub>2</sub>PdIn<sub>6</sub> becomes greatly enhanced and field-dependent in the coherent Kondo state. Both thermoelectric coefficients behave highly anomalously at low temperatures, namely the ratios S/T and v/T diverge logarithmically towards T = 0. The observed behaviors strongly support a 2D antiferromagnetic spin-densitywave scenario of the quantum criticality in Ce<sub>2</sub>PdIn<sub>6</sub>, recently inferred from the thermodynamic data [4].

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### AB04

### Pressure driven quantum critical point in CeNiAsO

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We perform a systematic investigation on the pressure effect on the antiferromagnetic dense Kondo lattice CeNiAsO[1]. Both hydrostatic pressure and chemical pressure (i.e. P-for-As doping) are successful in tuning the electronic state to get a pressure driven quantum critical point. We found that ST\_{N1}^{1} decreases slowly with initial pressure, but 'suddenly' disappears around 4 kbar. In contrary to the first-order-like quantum phase transition of ST\_{N1}^{1} S\_{S}^{1}\_{1} (N2)^{2} is continuously suppressed with increasing pressure. A QCP like phase transition is observed near 6.7 kbar, where ST\_{N2}^{1} N2 is totally suppressed, and non-Fermi-liquid (non-FL) behavior is observed, accompanied with a 'divergent' density of state (DOS) and a sign change of Hall coefficient. In the highly pressurized condition, it exhibits as a typical Kondo lattice metal. In P doped CeNiAs, Px samples, the chemical pressure leads to similar phenomena, with the critical doping level  $x_{-}$  c=S0.4. Magnetic susceptibility measurement points to the delocalization of Ce-4f electrons. Band structure and Fermi surface topology of CeNiAsO and CeNiPO were calculated by the method of LDA+DMFT. A new hole pocket contributed by Ce 4f band emerges in CeNiPO. The calculated Kondo scales for CeNiAsO and CeNiPO are  $T_{k} = 15$  K and 527 K respectively, which qualitively agrees with the experiment.

Yongkang Luo. at al. CeNiAsO: an antiferromagnetic dense Kondo lattice. J. Phys.: Condens. Matter 23, 175701 (2011).
 Philipp Gegenwart, et al., Quantum criticality in heavy-fermion metals. Nature Physics 4, 186-197 (2008).

### AC01

#### HgTe as a topological insulator Laurens Wigbolt Molenkamp

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HgTe is a zincblende-type semiconductor with an inverted band structure. While the bulk material is a semimetal, lowering the crystalline symmetry opens up a gap, turning the compound into a topological insulator. The most straightforward way to do so is by growing a quantum well with (Hg,Cd)Te barriers. Such structures exhibit the quantum spin Hall effect, where a pair of spin polarized helical edge channels develops when the bulk of the material is insulating. Our transport data provide very direct evidence for the existence of this third quantum Hall effect, which now is seen as the prime manifestation of a 2-dimensional topological insulator. To turn the material into a 3-dimensional topological insulator, we utilize growth induced strain in relatively thick (ca. 100 nm) HgTe epitaxial layers. The high electronic quality of such layers allows a direct observation of the quantum Hall effect of the 2-dimensional topological surface states. Moreover, on contacting these structures with Nb electrodes, a supercurrent is induced in the surface states.

### AC02

Spin liquid and spin glass states in frustrated magnets Seung-hun Lee

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In frustrated magnets, the topology of the magnetic interactions between magnetic moments leads to macroscopic ground state degeneracy. The nature of the ground states that emerge in frustrated magnets is yet to be understood. In this talk, I will discuss a few examples that exhibit spin liquid and spin glass behaviors at low temperatures. I will present some simple and eloquent ways of thinking about the complex frustrated magnetic interactions that help us greatly to understand the nature of the ground states and their low energy excitations.

### AC03

A novel magnetic order of ZnCr<sub>2</sub>O<sub>4</sub> revealed by magneto-optical measurements in ultra-high magnetic fields of up to 600 T Atsuhiko Miyata<sup>1\*</sup>, Shojiro Takeyama<sup>1</sup> and Hiroaki Ueda<sup>2</sup> <sup>1</sup> Institute for Solid State Physics, University of Tokyo, Japan

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 $Cr^{3+}$  ions in  $ACr_2O_4$  (A = Zn, Cd, Hg) form a pyrochlore lattice and these compounds are regarded as geometrically frustrated magnets. CdCr2O4 and HgCr2O4 exhibit diverse magnetic phases including a 1/2 plateau phase under magnetic fields, which can be described by a theory taking account of a spin-lattice coupling by Penc et al.. On the other hand, ZnCr2O4 shows a quite-small spin-lattice coupling, and there are high possibilities that another novel magnetic phases exist in ZnCr2O4. We investigated magnetic phases of ZnCr2O4 by the Faraday rotation and magneto-optical absorption spectral measurements up to 600 T by the electro-magnetic flux compression method. A ferromagnetic phase transition was observed at 410 T and 4.6 K. Furthermore, an unambiguous anomaly was found in magneto-optical absorption spectra at lower magnetic field side of the ferromagnetic phase. This indicates existence of a novel magnetic phase which could only be explained by a theory beyond that by Penc et al... A magnetic structure of the novel phase was inferred from analogical consideration with quantum phases in 4He, based upon similarity of a symmetry breaking. As a result, we concluded that the novel phase corresponds to a magnetic "superfluid" state with an umbrella-like magnetic structure.

### AC04

### Origin and signatures of magnetic chirality in the frustrated multiferroic Ba,NbFe,Si,O<sub>14</sub>

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Geometrical frustration is the key promoter of exotic collective magnetic ground states in condensed matter, including spin liquids, spin ice and unconventional spin orders. On frustrated lattices, the dominant isotropic exchange interactions are often unable to raise the degeneracy of the ground state. Then, minute magnetic-anisotropy terms can become crucial for selecting the ground state. The acentric Fe-Langasite Ba<sub>3</sub>NbFe<sub>5</sub>Si<sub>2</sub>O<sub>14</sub> features a frustrated arrangement of spin-5/2 Fe<sup>3+</sup> ions on vertices of equilateral triangles arranged into 2D triangular lattice [1]. It has been drawing considerable attention due to its multiferroic properties and its unique magnetic ground state. The long-range magnetic order that develops simultaneously with the electric polarization below TN = 26 K is a unique double-chiral single-domain state, whose realization appears mysterious [1]. Recently, we have extended research of the chiral properties of Ba<sub>3</sub>NbFe<sub>5</sub>Si<sub>3</sub>O<sub>14</sub>. Our inelastic neutron scattering study has revealed that one of the two excitation branches is completely chiral over the whole energy spectrum, which is a unique dynamical fingerprint of the chiral ground state [2]. Moreover, our electron spin resonance investigation has unveiled the origin of the mysterious selection of the ground state [3]. We will highlight the crucial role of minute Dzyaloshinsky-Moriya magnetic anisotropy in this system.

[1] K. Marty, V. Simonet, E. Ressouche, R. Ballou, P. Lejay, and P. Bordet, Phys. Rev. Lett. 101, 247201 (2008). [2] M. Loire, V. Simonet, S. Petit, K. Marty, P. Bordet, P. Lejay, J. Ollivier, M. Enderle, P. Steffens, E. Ressouche, A. Zorko, and R. Ballou, Phys. Rev. Lett. 106, 207201 (2011). [3] A. Zorko, M. Pregelj, A. Potocnik, J. van Tol, A. Ozarowski, V. Simonet, P. Lejay, S. Petit, and R. Ballou, Phys. Rev. Lett. 107, 257203 (2011).

### AD01

#### Investigating magnetic dipolar interactions between Co nano-islands with spin-polarized scanning tunneling microscopy

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In application of spin-polarized scanning tunneling microscopy (SP-STM), we investigated magnetic interaction in nanometer scale between triangular Co nanoislands formed in the sub-monolayer Co deposited on Cu(111) substrate. The SP-STM images the magnetization or the relative spin configuration of the Co islands. The results indicate a competition between exchange coupling and dipolar interaction in nanometer scale among the conjoint Co islands. The parallel magnetization of a set of conjoint Co nano-islands is anticipated due to the ferromagnetic exchange coupling. However, significant numbers of sets with anti-parallel ordering can still be observed. In order to explain the anti-parallel data, we calculate the dipolar interaction between conjoint Co islands to realize if it is comparable with the exchange coupling strength, which could be correlated to the length of contact region. As a result, the SP-STM images and the exchange coupling in nanometer scale. The anti-parallel configuration, in which dipolar energy dominates, is shown to be limited by the contact length of two conjoint islands, which scales the strength of exchange coupling.

### AD02

### Temperature-driven oscillatory magnetic anisotropy in ultrathin ferromagnetic films

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A contribution of quantum well states (QWS) to magnetic anisotropy energy (MAE) can be expected if the spread of the Fermi function is much smaller than the energy difference between the two states of each QWS pair contributing to the MAE [1]. Such a temperature effect on MAE is shown and discussed e.g. for fcc-Co films on vicinal surfaces of Cu(001) [2] and for bcc-Fe films on vicinal surfaces of Ag(001) [3]. In particular, the role of QWS in the vicinity of a spin reorientation transition is analyzed in the context of a model describing the anisotropy of ferromagnetic (FM) films grown on vicinal surfaces. The model reproduces not only oscillatory behavior of the uniaxial step-induced anisotropy, but also of the oscillatory tilting angle of magnetization. Interestingly, a small amount of e.g. Au on top of an Fe film completely changes how the magnetic anisotropy evolves with varying thicknesses of an Fe film. More details on the depth profile of the FM magnetization due to QWS are obtained from soft x-ray resonant magnetic reflectivity (SXRMR). In particular, it is found that the magnetization is not homogeneous in all Fe atomic lyers.

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### AD03

### Magnetism of ultrathin Fe films on BaTiO<sub>3</sub>(001)

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Pristine (< 1 monolayer (ML)) Fe films can be grown on BaTiO<sub>3</sub>(001) template [1] around 10 K by thermal evaporation. X-ray magnetic circular dichroism revealed superparamagnetism of the Fe films due to their cluster formation. Their saturation magnetic moments increase with the increase of the coverage, reaching beyond that of bulk Fe for 0.7 ML, which indicates that the clusters coalesce to form extended film, while the thickness of the clusters are thicker than 1 ML as found in our scanning tunneling microscopy. Fe films show in-plane easy axis, in contrast to the prediction of perpendicular magnetic anisotropy of the strain, multilayer formation of the Fe films and also atoms at highly populated edge sites. Fe films are found to induce magnetic moments of Ti, but none of Ba ions as predicted by first principles calculations [1]

[1]X. Chen, S. Yang, J.-H. Ki, H.-D. Kim, J.-S. Kim, G. Rojas, R. Skomski, H. Lu, A. Bhattacharya, T. Santos, N. Guisinger, M. Bode, A. Gravermann, and A. Enders, submitted to New J. Phys. 13 083037 (2011). [2]C. -C. Duan, S. S. Jaswal, and E. Y. Tsymbal, Phys. Rev. Lett. 97 047201 (2006).

### AD04

### Thickness-dependent exchange splitting of EuO ultrathin films

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EuO is a ferromagnetic semiconductor with the Curie temperature (TC) of about 70 K [1]. In the electron doping case by the substitution of trivalent ions, such as La<sup>3+</sup>, TC increases to the temperature as high as 200 K [2]. To apply EuO to spintronics devices such as spin filter tunnel barriers, it is important to fabricate ultrathin films with a few atomic layers and to clarify the electronic structure including the exchange splitting energy as well as the TC. Recently, we succeeded to fabricate single-crystalline EuO ultrathin films by a molecular beam epitaxy method [3]. Using the three-dimensional angle-resolved photoemission spectroscopy, we observed the electronic structure of the EuO ultrathin films. Then we could determine the thickness-dependent exchange splitting energy and clarify the mechanism of the changing magnetic property with decreasing thickness. Both of TC and the exchange splitting energy gradually decrease with decreasing thickness. Both of TC and the exchange splitting energy gradually decrease and the decreasing thickness. The decrease of the hybridization intensities between the Eu 4f and O 2p states and between the Eu 4f and Eu 5d states due to the decrease of the neighbor atom number are the origin of the change of the magnetic properties of EuO ultrathin films.

[1] A. Mauger et al., J. Phys. (Paris) 39, 1125 (1978). [2] H. Miyazaki et al., Appl. Phys. Lett. 96, 232503 (2010). [3] H. Miyazaki et al., Submitted to J. Phys. Conf. Series.

### AE01

### Magnetoelectric control of magnetic anisotropy in ultrathin Fe films using a charge-trap heterostructure

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Magnetoelectric switching of the magnetization vector could enable new low-power logic devices and non-volatile memory cells. Magnetoelectric switching typically requires complex multiferroic oxides or strain coupled magnetostrictive/piezoelectric composites. However, recently it has been shown that surface magnetic anisotropy in ultrathin ferromagnetic metal films can be directly controlled by application of a gate voltage across an adjacent oxide layer [1]. Unfortunately, magnetoelectric effects in metals require a relatively large electric field and are inherently volatile. In this work [2], we examine magneto-electric switching in ultrathin epitaxial Fe films on Ag(001). We show that a charge-trapping layer integrated into the gate dielectric can provide the missing non-volatility to the magnetoelectric effect and enhances its efficiency by an order of magnitude. We report the largest voltage-induced change to surface magnetic anisotropy yet demonstrated for a metallic thin film, and directly correlate this change with the density of trapped charge in an adjacent charge storage layer. We also demonstrate a novel optical charge pumping mechanism, which offers the unique possibility of optical imprinting of the magnetic state in a continuous Fe film. Supported by the National Science Foundation through arant ECCS-1128439.

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### **AE02**

### The origin of electric-field effects on magnetic anisotropy in FePd ultrathin film

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The electric-field (EF) (assisted) magnetization reversal technique that has the advantage of lowpower consumption has been suggested. The EF dependence measurement on the coercivity or the magnetic anisotropy in FePt[1] and FePd[12] thin films has been reported. We have performed theoretical analysis for EF effects on magnetic anisotropy energy (MAE) in Pt/Fe/Pt(001)[3] and Pd/Fe/Pd(001)[4] thin films and found that these systems show the inverse EF effect between the Pt and Pd systems. It was also found that this effect corresponds well to the experimental observations [1,2] The microscopic origin for the inverse issue has not been discussed deeply, although the systems of Pt and Pd are similar to each other. The distinct property of Pd against Pt is magnetic polarization. In order to clarify origins for the inverse effect in Pd/Fe/Pd(001), by using the density functional calculation, we have further estimated MAE and its EF effect in the several sets of spinorbit coupling choice for atoms. As a result, we found that the contribution from the Pd substrate is very important in the inverse issue. In the presentation, we will show results of MAE density map in real space and discuss comparison with the Pt system.

[1] M. Weisheit et al., Science 315, 349 (2007). [2] F. Bonell et al., Appl. Phys. Lett. 98, 232510 (2011). [3] M. Tsujikawa and T. Oda, Phys. Rev. Lett. 102, 247203 (2009). [4] S. Haraguchi et al., J. Phys. D: Appl. Phys. 44, 064005 (2011).

#### AE03

### Ferroelectric control of spin polarization

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A current drawback of spintronics is the large power required for magnetic writing. Aiming at a nonvolatile low-power electrical control of spintronics properties we have combined two ferromagnetic electrodes (namely La<sub>2/3</sub>Sr<sub>1/3</sub>MnO<sub>3</sub> and Fe or Co) with a ferroelectric BaTiO<sub>3</sub> tunnel barrier. We have found that the TMR depends on the direction of the ferroelectric polarization in the barrier. This demonstrates a local, large, and nonvolatile control of spin polarization by electrically switching ferroelectric polarization. Interestingly, just as ferroelectricity in the BaTiO<sub>3</sub> influences the spin polarization of the adjacent Fe or Co magnetic electrodes, magnetism is induced by the electrodes in the ferroelectric. We have used soft X-ray resonant magnetic scattering to evidence that, in addition to being ferroelectric, the BaTiO<sub>3</sub> barrier possesses a hysteretic magnetization at room temperature. Ab initio calculations of realistic interface structures provide insight into the origin of the induced moments and perspectives towards new multifunctional interfaces for spintronics.

V. Garcia et al. Science 327, 1106 (2010) : S. Valencia et al. Nature Mater. 10, 753 (2011) : L. Bocher et al, Nano Lett. 12, 376 (2012) ; Financial support from ERC advanced grant no. 267579

### **AE04**

### Voltage controlled spin transport channel

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Two-terminal metal-oxide-based devices in which memory-resistive(memristive) switching occurs show great potential for use in next generation technologies such as non-volatile memories and neural networks[1,2]. It has been reported that TaOx-based devices combined with Cu and Pt electrodes show memristive behavior utilizing metal precipitation to create conductive channel (tens of nanometers in diameter) inside of the oxide thin film when an electric voltage (only around 1V between two electrodes sandwiching the oxide layer) is applied[3]. This unique mechanism can deliver an opportunity to explore a new type of technological application, an electron spin transport channel controlled by an external voltage. We have explored this possibility by investigating devices with a TaOx thin film sandwiched between two ferromagnetic metal electrodes. Three different sets of devices, Co/TaO,/Cu/Py, Py/TaOx/Cu/Co, and Co/TaO,/Py were prepared in cross-bar shape. Only devices with a Cu layer show memristive behavior having a typical OFF/ON resistance ratio of 105. Magnetoresistance measurements performed by sweeping an external magnetic field display evidence of spin transport in the low-resistance ON-state at 77 K. No magnetoresistance was detected in the OFF-state indicating spin transport vanishes. Our study illustrates that we experimentally realized a continuous electron spin transport channel controlled by a low power voltage[4]

[1] J. J. Yang et al., Appl. Phys. Lett. 97, 232102 (2010). [2] S. H. Jo et al., Nano Lett. 10, 1297 (2010). [3] T. Sakamoto et al., Appl. Phys. Lett. 91, 092110 (2007). [4] H.-J. Jang et al., Appl. Phys. Lett. 100, 043510 (2012).

### AE05

### Zeeman-type spin splitting controlled by an electric field

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Electrically manipulating electron spins based on Rashba-type spin-orbit interaction (SOI) is a key path-way for applications of spintronics and spin-based quantum computing. Two-dimensional electronic system (2DES) is a particularly important SOI platform where the in-plane spin polarization can be tuned with an electric field perpendicular to the 2DES (also spins). In contrast, electric manipulation of outof-plane polarized spins remains elusive and a great challenge though it is known beneficial for longer spin lifetime. Here, we present a new concept, in which an out-ofplane spin polarization is realized and further controlled with the external perpendicular electric field (parallel to polarized spins), by combining the first-principle calculations and magnetotransport experiments in a field effect transistor made of the transition metal dichalcogenide WSe<sub>3</sub>. Such an out-of-plane spin splitting is shown to take place at the corners of WSe<sub>2</sub> Brillouin Zone in a form of Zeeman-type spin polarization, due to the exposure of a large but hidden in-plane electric dipole of WSe, layers, thereby enabling new degrees of control for quantum confined spintronic devices.

#### **AF01**

#### Femtoscale magnetically induced lattice distortions in multiferroic TbMnO<sub>2</sub>

Helen Walker1\*, Francois De Bergevin2, Federica Fabrizi3, Luigi Paolasini2, Andrew Boothroyd<sup>3</sup>, Dharmalingam Prabhakaran<sup>3</sup> and Desmond Mcmorrow<sup>4</sup> Resonant Scattering and Diffraction Beamline P09, PETRA III, HASYLAB at DESY, Germany <sup>2</sup> European Synchrotron Radiation Facility, France <sup>3</sup> Department of Physics, University of Oxford, United Kingdom 4 London Centre for Nanotechnology, University College London, United Kingdom The discovery of the canonical spin-cycloid multiferroic TbMnO<sub>3</sub> [1], in which the onset of a spontaneous ferroelectric

polarization is concomitant with a magnetic phase transition, has generated considerable interest in the control of electric polarization by magnetic fields, and vice versa. Although comprehensive, microscopic descriptions of the magnetic structure have been obtained, our understanding of the ferroelectric state is still developing. Competing theoretical models proposed to explain the ferroelectricity are associated with either ionic displacements, arising due to antisymmetric exchange interactions [2], or charge transfer [3]. By exploiting the magneto-electric coupling in TbMnO3, we used an electric field to produce a single magnetic domain state, and a magnetic field to induce ionic displacements. Under these conditions, interference charge-magnetic x-ray scattering arose, encoding the amplitude and phase of the displacements. When combined with a theoretical analysis, our data allow us to resolve ionic displacements at the femtoscale which contribute a significant fraction of the zero-field ferroelectric moment, and reveal that both symmetric and antisymmetric Dzvaloshiskii-Moriva interactions need to be included into the microscopic models [4]. Thus our results represent an important step forward towards the goal of realising the technological potential of these materials.

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### AF02

### Direct measurement of the interatomic distance dependence of the magnetic exchange interaction

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Recently, we demonstrated the feasibility to utilize the short-ranged magnetic exchange interaction to map spin textures of surfaces with atomic resolution on insulators [1] as well as on metals [2] with magnetic exchange force microscopy. Even more promising is the spectroscopic mode, i.e., magnetic exchange force spectroscopy (MExFS), which enabled us to directly probe the distance dependence of magnetic exchange interaction [3]. The experimental data were acquired with a magnetically coated tip on the antiferromagnetic Fe monolayer on W(110). First, the distance dependence of the total tip-sample interaction was measured on Fe surface atoms of oppositely oriented magnetic moments. Thereafter, the magnetic contribution was obtained by subtracting both curves from each other. This approach guarantees that all non-magnetic interactions, e.g., the short-ranged chemical interaction and the long-range van der Waals interaction, cancel out. The result, i.e., the magnetic exchange interaction, exhibits excellent agreement with ab-initio calculations based on density functional theory performed for the same tip-sample system.

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### AF03

### Magnetic nanodomains in manganites revealed by Lorentz TEM and small-angle electron scattering

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La1-xSrxMnO3 (LSMO) shows a complex variety of magnetic and electronic ground states, which depend on the Sr doping concentration (x). Here, we have thoroughly examined magnetic microstructures in single crystals of LSMO by using lowtemperature Lorentz transmission electron microscopic (TEM) and small angle electron scattering experiments. The magnetic stripe domains with periodic separation of approximately 200 nm are formed in the orthorhombic structure of La1, Sr, MnO3 for  $x \sim 0.175$  in the zero magnetic field. An in-situ Lorentz TEM observation revealed that, when the vertical magnetic field is applied by utilizing the magnetic field of object lens, the magnetic stripe domains change into a form of the magnetic vortex with the 100 nm size. These magnetic nanodomains with chiral spin configuration (magnetic vortex) are new kinds of magnetic ground states (spin textures) in manganites. In the presentation, we will explain detailed responses of magnetic vortices to external magnetic fields and discuss the nucleation and growth mechanism of magnetic vortices in the magnetic stripe domains.

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### AF04

### Morin transition control of antiferromagnetic α-Fe<sub>2</sub>O<sub>3</sub> films with epitaxial strains

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It is well known that hematite(α-Fe<sub>2</sub>O<sub>3</sub>) undergoes a spin reorientation transition at around 263 K. It is Morin transition only found in hematite and some orthoferrites. J.O. Artman, J.C. Murphy, and S. Foner suggested that it is caused by the competition between the single ion anisotropy and the magnetic dipole anisotropy[1]. Their model has been confirmed with the fine particles and pressurized single crystal and it becomes clear that the lattice constants in c-plane are critical to determining the Morin temperature [2,3]. However, when Fe<sub>2</sub>O<sub>3</sub> are grown in the form of epitaxial films, no one has observed the transition despite the expectation that the epitaxial strain can be controlled better in the form of films [4.5]. In this work, we have grown a series of epitaxial α-Fe<sub>2</sub>O<sub>3</sub> films on Al<sub>2</sub>O<sub>3</sub> substrates. The lattice constants of them are successfully controlled with the Cr2O2 buffer layer and the strain relaxation. We found that the Morin temperature changes consistently with the lattice constants and it reaches 360 K for thinnest film. The results are well in accordance with the theoretical model. This research was supported by NRF funded by MEST (2009-0088969, 2010-0018733).

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### AF05

### Surface plasmons and magneto-optical activity in hexagonal Ni antidot arrays

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The interaction between the surface plasmon polaritons (SPPs) and a magnetic field in a pure magnetic metal is presented in this work. By utilizing nano patterned structures one can create plasmonic resonances on magnetic surfaces. We have investigated how hexagonal arrays of circular holes (250-275 nm) in a Ni matrix (anti-dot arrays) couple to plasmonic resonances and therefore how it changes the optical and magnetooptical spectra. We have performed angularly resolved scans of the reflectivity and transmission with and without application of magnetic field and compared them to theoretical calculations. We present signatures of SPPs as troughs in reflectivity and extraordinary transmission through the holes. In the plasmonic regimes a big enhancement of the magneto-optical response (polar moke) is observed. Application of a transverse magnetic field influences the excitation of SPPs as it can be seen by the big increase of magnetic contrast measured in the transverse Moke configuration. These effects pave the road for the development of new optical components and sensors with great application potential

E. T. Papaioannou, V. Kapaklis, E. Melander, B. Hjorvarsson, S. D. Pappas, P. Patoka, M. Giersig, P. Fum-agalli, A. Garcia-Martin and G. Ctistis, "Surface plasmons and magneto-optic activity in hexagonal Ni anti-dot arrays", Opt. Express 19, 23867 (2011)

### AG01

Jian Shen

#### Magnetic nanodots induced novel magnetic phenomena

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Study of Magnetic nanodots is at central in the field of nanomagnetism. Besides the interesting properties caused by dimensionality effect, magnetic nanodots can induce many novel phenomena when forming heterostructures with other materials. In this work, I will use several examples to demonstrate their effect. These examples include collective ferromagnetism of nanodot arrays on 2-dimensional electron gas, colossal Coulomb blockade magnetoresistance in tri-layers, and dramatic enhancement of metal-insulator transition temperature in manganites. All these fascinating phenomena originate directly from the presence of magnetic nanodots. Their underlying mechanisms, while somewhat understood, need further theoretical studies

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### AG02

### Magnetic properties of Fe-(Pt,Pd) thin films patterned by selfassembling of polystyrene nanospheres

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Arrays of nanoparticles with high magnetic anisotropy are studied for magnetorecording. Nanolithography techniques are employed to prepare arrays of magnetic dots/holes. Polystyrene nanosphere (PN) lithography has been recently exploited for magnetic thin films nanostructuring, due to its low-cost and large covering area [1,2]. Dot arrays of Fe53Pt47 (MgO(100) substrate heated at 400 °C) were obtained by assembling PN monolayers on a continuous film (10 nm), followed by sputter etching with Ar<sup>+</sup> ions (final dots diameter 80 ÷ 400 nm) [1]. Antidot arrays of Fe<sub>50</sub>Pd<sub>50</sub> (50 nm, Si substrate, post deposition annealing at 550 °C) have been created by exploiting PN as diffraction masks in combination with a mercury lamp [2]. In this alternative lithography technique, the PN behave as optical lenses to generate regular diffraction patterns (holes, diameter = 300 nm) on a photoresist. In both cases, the annealing induces the orderdisorder transformation towards the L10 tetragonal phase. Sample microstructure was studied with SEM and AFM microscopy. Parallel and perpendicular hysteresis loops were measured by AGFM, confirming the presence of the tetragonal phase. The complex multidomain magnetic pattern was studied in both compositions with MFM and related to dot/hole diameter and their mutual distance.

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### AG03

#### Position dependence of vortex core oscillation in polygonal nanomagnets

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A magnetic vortex structure formed in a micron-scale ferromagnetic dot has been gathering attention because of its superior performance for a unit cell in spintronic devices[1] and intriguing magnetic responses such as bloch point reversal[2]. Control and further understanding of vortex properties are required for the practical application. Previously, we showed that the vortex chirality can be controlled simply by the application of the in-plane magnetic field in a regular polygonal-shape nanomagnet[3]. This enables us to control the core position precisely by applying the external magnetic field. Since the potential distribution for the core in the polygonal nanomagnet should be different from a simple harmonic potential formed in a circular magnetic dot, we expect that the resonant frequency depends on the core position. In the present work, we investigate the vortex core dynamics in a triangular nanomagnet under the external magnetic field using an electrical homodyne detection technique and the transmission impedance measurements using a vector network analyzer. The resonant frequency was found to increase with decreasing the distance between the core and the vertex.

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### AG04

### Huge magnetic anisotropy and coercivity in Fe island and atomic wire

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We have investigated huge magnetic anisotropy energy and coercivity of Fe nano islands and nano stripes on flat and stepped W(110) using x-ray circular dichroism (XMCD) and scanning tunneling microscope, which has not been directly decided so far due to its large anisotropy [1]. The anisotropy energy is 1 meV per Fe atom as determined from magnetization curve by XMCD at Fe L edge, almost constant both for the nano island and the nano stripe below one monolayer. This anisotropy energy corresponds to an anisotropy field of 17 T. The coercivity, Hc, for the nano islands at 0.25 ML is very large, ~5 T, but it rapidly decreases to 1 T at 0.6 ML, which follows Hc ~ 1/r (radius of island). On the other hand, the coercivity for the nano stripes is almost constant, Hc ~ 3 T below 0.5 ML. The obtained coercivity is much smaller than the anisotropy field, suggesting that the magnetization reversal process is via a domain wall creation, not a coherent rotation. This contrasted behavior is explained by the shape of deposited Fe on the stepped and flat W surfaces and the creation of domain

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### AG05

### Oscillation of critical fields in highly dense arrays of magnetic nanodisks

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Magnetostatic interaction between magnetic nanostructures composing a highly dense array influences on its magnetization reversal and micromagnetic configuration. In this paper we demonstrate the effect of nanodisk array dimension N on magnetic properties. The nanodisk patterns with 600 nm in diameter and distance between disk's centers d from 700 nm to 1um were fabricated by electron-beam lithography. The number of disks on a boundary N was varied from 2 to 20. After sputtering of Py films and lift-off process the nanodisks with thickness t =9, 22 and 30 nm were prepared. Using magnetic hysteresis curves we found out that with an increase of N the fields of magnetic vortex nucleation Hn and annihilation Ha decreased. At the first time we demonstrate the oscillations of function Hn = f(N) in arrays. The maximum amplitude of oscillations was observed in case of nanodisks with t=22 nm and d=800 nm. Magnetic force microscopy shown that at t=22 nm the magnetization reversal was occurred mainly via vortex core motion. We were carried out the calculation of critical fields Hn and Ha with in "rigid" vortex model for different N. The theoretical result is in a good qualitative agreement with the experimental data.

### AH01

### Magnetic tools for molecular diagnosis

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This talk shows examples for the application of magnetic transducers such as magnetoresistive sensors ( $\rightarrow$ 'Lab-on-a-chip') or multifunctional magnetic nanoparticles (-> 'Lab-on-a-bead') as tools for molecular diagnosis Regarding Lab-on-a-chip systems we demonstrate in-vitro studies of magnetic particle uptake (phagocytosis) by cell cultures (e.g. NHDF or DU145 cells). The method is based on distance changes of the magnetic particles to embedded magnetoresistive sensors during phagocytosis and allows analysis of factors that are relevant for applications like magnetofection or hypothermia. As a second Lab-on-a-chip application. we show the concept and progress of a CMOS-integrated tunneling magnetoresistance sensor chip currently under development for a point-of-care device for sepsis diagnosis. Our Lab-on-a-bead approach focuses on homogeneous immunodiagnostics based on manipulation and observation of multi-functional nanoparticles. This new biosensor concept is based on highly sensitive plasmon-optical detection of the rotational dynamics of anisotropic magnetic nanonarticles immersed in the sample solution, which changes on the binding of analyte molecules to the nanoparticles. We present model calculations of the optical magnetic and hydrodynamic properties of both the required nanoparticles and measurement conditions. Experimentally, we show transmission measurements on plain ferromagnetic nanorod dispersions in rotating magnetic fields. Schrittwieser X, Ludwig F. Dieckhoff J. Saulantica K, Viau G. Lacroit I-M. Lenijo SM, Bouheler R, Maynadie J, Huetton A, Brueckl H, Schotter J. Modeling and Development of a Biosensor Based on Optical Relaxation Measurements of Hybrid Nanoparticles. ACS Nano, 2012;6:791

### AH02

### Control of the living cell machinery with nanomagnets

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We present proof-of-concept experiments on magnetically controlled cellular endocytosis of functionalized magnetic nanoparticles. The possibility of remote control of the cellular endocytosis and/or exocytosis rates by an externally applied magnetic field without using various types of inhibitors seems to be a very intriguing direction in nanomedicine [1]. By pharmacological and in vitro knockdown approaches, the principal cellular uptake mechanism for two types of magnetic nanoparticles was identified [2]. A model of the uptake process that allows determination of key parameters of endocytosis, including the rate of uptake, the number of nanoparticles per cell in saturation, the mean uptake time, and the correlation between the number of internalized nanoparticles and their extracellular concentration was developed. We also describe our recent progress in the applications of nanomagnets for magnetic hyperthermia treatments and targeted drug delivery [3]. In particular, we will propose a new mode of hyperthermia that we have analyzed. The obtained results have important implications for the fundamental knowledge on intracellular nanomechanics and open ways for magnetic controlling of the living cell machinery.

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### AH03

#### Bio-functionalized magnetic nanoparticles for in-vitro diagnosis of colorectal cancer

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Colorectal cancer (CRC) is the most popular cancer. Carcinoembryonic antigen (CEA) has been used as bio-markers for CRC. The signals for CEA solutions of various concentrations, or with interfering materials are detected via immunomagnetic reduction (IMR), the assay technology involves the utilities of bio-functionalized magnetic nanoparticles to label CEA. By using the bio-functionalized magnetic nanoparticles and the analyzer, the detection threshold and interference tests for assaying CEA are examined. Several serum samples from normal people and CRC patients are used for the detections of CEA concentration via IMR. The results of assaying CEA with IMR and enzyme-linked mmunosorbent assay (ELISA) were compared. As compared to ELISA, the detection threshold for assaying CEA with IMR is lowered by at least two orders of magnitude. For the interfering materials noted commonly in serum such as hemoglobin, bilirubin, triglyceride, and vascular endothelial growth factor, there is no detectable interfering effect for assaying CEA with IMR. These results reveal the feasibilities of assaying CEA in the blood using IMR, as well as achieving a high-sensitive and high-specific assay for CEA

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#### **AH04**

### Multiplexing capabilities of multi-frequency magnetic ratchets

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Through a combination of theory, numerical simulation, and experiment, we investigate the motion of magnetic beads on the surface of a magnetic ratchet driven by multifrequency fields. We demonstrate that in multifrequency driving fields, two different bead types experience differential motion by modulating the phase difference between the two driving frequencies. We use perturbation analysis to demonstrate that static forcing terms are the origin for why two different bead types can move in opposite directions on the same ratchet notential. Based on these theoretical results, we conduct experimental investigations that explore the effects of bead size and static forcing coefficient on the direction of head motion, which confirm most of the expected trends These results shed light both on past experimental work both by ourselves and others as well as elucidate the more general multiplexing capabilities of ratchets.

### AI01

#### Quantum control of single spins in diamond and silicon carbide David D. Awschalom<sup>3</sup>

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A motivation behind modern research into quantum physics has been to identify robust quantum states that can be easily controlled for future use in advanced information technologies. Recently, a defect in diamond known as the nitrogen-vacancy (N-V) center has attracted interest because it possesses an atomic scale electronic spin state that can be used as an individually addressable, solid state quantum bit (qubit) even at room temperature. The N-V center's optical transitions become coherent, spindependent, and may be precisely controlled with gate voltages in micron-scale devices [1]. Moreover, engineered coupling of this electron spin with the proximal single nitrogen nuclear spin enables a scalable quantum memory [2]. These exceptional coherent properties have motivated theoretical efforts to predict similar defects in other semiconductors with expanded functionality [3]. We show that spin states in various polytypes of SiC can be optically addressed and coherently controlled at temperatures ranging from 20 to 300 K with coherence properties comparable to diamond N-V centers. These spins are optically active near telecom wavelengths within an industrial scale material having advanced microfabrication techniques. This makes them promising candidates for photonic, spintronic, and quantum information applications that merge quantum degrees of freedom with classical electronic and optical technologies [4].

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### AI02

#### Dynamical spin injection into p-type Si using the spin pumping and spin transport at room temperature

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For observation of spin-transport in p-type Si, ferromagnetic Ni<sub>80</sub>Fe<sub>20</sub> and paramagnetic Pd were separately formed with various gap distances on a p-Si with doping concentration of 1×10<sup>19</sup> cm<sup>-3</sup>. The sample was placed near the center of a microwave cavity in ESR system. In a ferromagnetic resonance condition of the Nin-Ferro the steady magnetization precession was maintained by the absorption of the microwave and the angular momentum was transferred from the local spins to conduction electrons. Then conservation of spin angular momentum induced spin pumping into the p-Si, which generated pure spin current in the p-Si. The spin current propagated in the p-Si, and was absorbed into the Pd, where the spin current was converted to a charge current by the inverse spin-Hall effect(ISHE) in Pd. When the gap length was 490 nm, the output voltage of 1.4 µV due to the ISHE was obtained. Using the results, the spin diffusion length and spin coherence time in p-Si at room temperature were estimated to be at least 130 nm and 94 ps, respectively[1]. In the presentation, we will discuss the detail of the research.

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### AI03

### Studying the optical spin orientation in Ge by exploiting the spin filtering in Fe/MgO/Ge photodiodes

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Spin-polarized carriers can be photo-excited in a semiconductor by monochromatic circularly polarized light. To date, several techniques have been introduced to detect the excited spin polarization. In this work, the spin filtering effect at the Fe/MgO interface has been employed to measure the degree of spin polarization of both electrons and holes in Ge. We have fabricated photodiodes by means of optical lithography, starting from fully epitaxial Fe/MgO/Ge(001) heterostructures [1,2]. Spin-detection experiments are performed by illuminating photodiodes with circularly polarized light in a magnetic field applied parallel to the light helicity. Photo-generated carriers in Ge are spin-filtered by the MgO barrier depending on the direction of the spin. The degree of spin polarization (DSP) is related to the variation of the photocurrent due to full reversal of the Fe magnetization. As predicted by Rioux and Sipe [3], the DSP is maximum when the photon energy is resonant with the direct energy gap of Ge and decreases with the photon energy for both holes and electrons. Moreover, as the DSP is different from zero in the frequency range from 1300 nm to 1550 nm, the device is suitable for integrated detection of light helicity at room temperature in novel spin-optoelectronics systems

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### A104

Transition from spin injection into interface states to the channel in n-Ge

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Spin injection into semiconductors is crucial for exploring spin physics and new spintronic devices. Unlike GaAs or Si, very few studies have dealt with spin injection in Ge [1,2,3,4]. This material is of great interest for high carrier mobility, long spin diffusion length and large spin-orbit coupling to perform electric field spin manipulation through Rashba interaction. However the exact role of interface states in spin injection mechanism in n-Ge has not been clarified yet and except in ref. [1] no clear evidence of spin accumulation in the channel has been given. In this paper, we show a clear transition from spin injection into interface states to the channel. For this purpose, we have grown CoFeB/MgO spin injector on GOI [2]. We observe spin signal amplification at low temperature due to spin accumulation into interface states. At 200 K, we observe a clear transition to spin injection in the channel up to room temperature: the spin signal is reduced to a value compatible with spin diffusion model and more interesting we could demonstrate spin signal modulation applying a back gate voltage and spin-pumping by the ferromagnetic resonance of the CoFeB layer which are clear manifestations of spin accumulated in the channel

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### A105

### Tunneling anisotropy in crystalline Si/MgO/Fe devices

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Key advances in the electrical creation of spins in silicon at room temperature have recently been made, using ferromagnetic tunnel contacts to inject a spin accumulation, and to detect it [1,2]. This enables the systematic study of the various parameters that govern spin injection into silicon. Here we demonstrate that spin transport in ferromagnet/oxide/silicon tunnel devices is anisotropic, i.e., depends on the absolute orientation of the magnetization direction of the ferromagnetic electrode. In particular, we show that the magnitude of the spin accumulation induced in the Si is anisotropic and changes when the magnetization is rotated either in the in-plane or the out-ofplane direction. We present a systematic study of the anisotropic tunnel conductance of crystalline Si/MgO/Fe spin tunnel contacts at 300 K. We observe an in-plane anisotropy that reflects the crystal symmetry of the tunnel contact. The out-of-plane anisotropy shows a bias-dependent, non-trivial variation with magnetization angle that arises from a superposition of different contributions to the anisotropy, one of which is the anisotropy of the tunnel spin polarization. The results will be compared to anisotropy data on contacts with an amorphous Al2O3 tunnel barrier, n-type and p-type silicon, and Ni and Fe as the ferromagnet.

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### AJ01

#### Micromagnetic models in glass-coated microwires with circumferential anisotropy

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The magnetic behaviour of amorphous glass-coated microwires, with ultrasoft magnetic behaviour, has been modelled by quasi-analytical 1D micromagnetics. Co-based alloys with negative magnetostriction present a noticeable circumferential anisotropy where an axial core responds to the minimization of exchange energy. The theoretical approach is based on one-dimensional micromagnetic model where exchange, magnetoelastic and magnetostatic energies are considered. We assume that magnetization profile varies only along radial direction and magnetoelastic anisotropy does not depend to the radial coordinate. Solving firstly the variational problem that leads to Euler-Lagrange equation, we estimate the size of the core and the axial hysteresis loops as a function of radius, exchange coefficient A and anisotropy constant K [1]. The critical nucleus radius below which the magnetization is fully axial was analytically obtained. The model was extended to high frequency dynamics under axial ac field. The microwave permeability has been computed as a sum of eigenmodes. These modes are obtained as solution of a Schrodinger equation where the potential derives from static solution. Finally, comparison with experimental measurements shows the important role of the skin effect for thicker wires.

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### AJ02

### Manipulation of domain wall dynamics in microwires by transverse magnetic field

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Recently fast domain wall (DW) propagation of thin magnetic wires gained considerable interest owing to possibility of applications for data storage devices and extremely fast DW propagation in composite glass-coated amorphous microwires. A number of interesting effect, such as collisions between two DWs, trapping and injecting of DWs by local fields and DW dynamics manipulation through the magnetoelastic energy have been reported. In this paper we studied effect of transversal magnetic field on DW propagation in microwires. In order to activate DW propagation from the determined wire end in our experiment we placed one end of the sample outside the magnetization solenoid. Under certain transversal magnetic field, Ht, we were able to create additional DW on the opposite wire end. Changing the Ht we were able to tailor propagation field of this additional DW and observe DW collision in different parts of microwire. Previously we observed spontaneous DW nucleation on local defects limiting single DW dynamics regime in microwires. The collisions realized by us by application of transversal magnetic field can be used to release such additional DWs to achieve higher DW velocity extending single DW regime to higher external magnetic fields.

#### AJ03

### Effect of process parameters on the microstrucutre and magnetic properties of electrodeposited FeCo thin films

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FeCo alloys have been extensively studied as soft magnetic materials due to their superior properties relative to Fe-Ni alloys as write head core materials in harddisk-drives. Electrodeposition represents a simple, cost-effective way of fabricating thin film recording heads, although stable plating baths are needed for commercial processing. It has been a challenge to fabricate FeCo films that have high saturation magnetizations (>2 T) with good soft magnetic properties (low coercivities Hc) by electrodeposition methods. In this study, the magnetic properties of FeCo thin films plated, under a number of different processing conditions, are investigated. The aim is to optimize the microstructure and magnetic properties such as saturation magnetic flux density and coercivity. The effect of current density (15-50 mA/cm2), cobalt/ iron ions concentration and bath temperature (room temperature to 45oC) on the films microstructure and properties were studied. The microstructure of electrodeposited films was characterized by XRD, TEM and scanning probe microscope. The magnetic properties of electrodeposited films were investigated by VSM and show dependence on the deposition conditions. After optimizaiton of process parameters, high saturation magnetization (~2.1T) and low coercivity (~100e) were obtained for electrodeopsited FeCo thin film in present study.

#### AJ04

#### Structure and magnetic properties of FeCo alloy synthesized by a onestep polyol process

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Most of the chemical methods for the synthesis of FeCo employ complex experimental conditions, expensive chemicals and carbonyl precursors which are less preferable [1]. Moreover each chemical method results in FeCo of various composition, particle size, orderdisorder behaviour and oxide fraction which subsequently affect its magnetic properties [2]. In this report, we present the synthesis of FeCo nanoparticles of varying composition by a unique one-step polyol process and report its magnetic properties. The synthesis of FeCo alloys were undertaken in ethylene glycol using hydrated FeCl<sub>2</sub> and Co(II)acetate and NaOH. The samples were prepared at various temperatures upto 190 oC. The samples are characterized using XRD, TEM-EDX and VSM. The variation of Curie temperature is determined using Thermo-Magnetic Analysis and is found to be in accordance with the composition of the FeCo alloy. The fcc to be phase transition temperatures is also evident from the Differential Thermal Analysis. The morphology of the particles is studied using TEM and it showed spherical or cubic nature depending on the hydrous precursors used in the synthesis. The present synthesis method can be employed for large scale synthesis of FeCo alloys with controlled structure and magnetic properties which shall be discussed in detail.

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### AJ05

#### Investigations of the magnetic and structural properties of a metalloidfree Co<sub>80</sub>Zr<sub>10</sub>V<sub>10</sub> amorphous alloy

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Nanocrystalline ferromagnetic materials used in magnetic applications are obtained from Fe- and Cobased amorphous alloys containing metalloid elements. In these systems, the magnetic properties are optimized by adjusting the size of the nanocrystals forming from the metastable amorphous phase upon annealing. This study was undertaken to investigate the correlation between the variations in the magnetic properties upon annealing of a metalloid-free Co-based amorphous alloy. For that purpose the measurements of the magnetic properties have been complemented by a microstructural analysis by means of nuclear magnetic resonance. The hysteresis loops of as-cast and annealed samples were measured at 50 Hz (Hmax = 1 kA/m) to get information on the dependence of the coercive field (Hc) and saturation induction (B1000) with the annealing treatment. Zero field 59Co NMR spectra have been recorded with an automated broadband spectrometer. The lowest value of coercivity was found to be 5.8 A/m and highest saturation induction of 0.7 Tesla in the sample annealed at 425oC for 20 minutes. The NMR investigation suggested an important contribution of pure Co atoms. The dependency of the magnetic properties upon the annealing treatment will be discussed based on local atomic arrangement in the amorphous structure.

#### AJ06

### Magnetocaloric effect in Fe-Ni-Zr alloys prepared by rapidly quenched method

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 $Fe_{90:x}Ni_zZr_{10}(x=0, 5, 10, 15, 20 and 25)$  ribbons with various thickness were prepared by melt-spinning technique. Curie temperature,  $T_{\rm C^*}$  of the alloys is dramatically decreased from ~960 K to room temperature region by high quenching rates. When the alloys are in amorphous structure, their TC strongly depends on the Ni-concentration. Maximum entropy change,  $|\Delta S_m|_{max}$ , with  $\Box H$  = 12 kOe, of the alloys is around 1 J.kg  $^1{\rm K}^{-1}$  in room temperature region. On the other hand, the full width at haft the maximum peak of the entropy change is quite large, FWHM  $\sim$  85 K, making it possible for application in magnetic refrigerators at room temperature.

PACS number: 75.30.Sg, 07.20.Mc, 75.50.Kj

Keywords: Magnetocaloric effect, Refrigeration, Amorphous magnetic materials

#### **BA01**

### Nodal pocket revealed by quantum oscillations in an underdoped cuprate superconductor

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Quantum oscillations have proved a vital tool in uncovering the electronic structure underlying the normal state of the underdoped cuprates. Our experiments have established that quasiparticles in this regime are governed by fermi dirac statistics, yet a comprehensive description of the fermi surface geometry remains elusive. We use an array of quantum oscillation measurements in the underdoped cuprate  $YBa_2Cu_3O_{6+x}$  in tandem with results from complementary experimental techniques to distinguish between alternative Fermi surface possibilities. We present evidence including quantum oscillations that persist down to low magnetic fields, a low doping quantum critical point that mirrors the nodal Fermi velocity collapse in photeenission, and oscillations of the chemical potential, all of which strongly point to a Fermi surface comprising solely nodal pockets. Possibilities for the origin of the antinodal gapping of density of states at the Fermi level are discussed.

S. E. Sebastian et al. Nature Communications 2, Article number: 471

### BA02

#### Low-Dimensional Superconductivity in δ -Doped SrTiO<sub>3</sub> Harold Hwang\*

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SrTiO3 is the lowest density bulk superconductor, which when incorporated in heterostructures provides an opportunity for band-structure engineering this superconducting semiconductor. We take this opportunity by selectively doping a channel of finite thickness in an undoped matrix, which is analogous to  $\delta$ -doping in semiconductors. By varying the thickness of a Nb doped channel, we have studied the crossover from 3D to 2D in the superconducting and normal state properties, as the thickness of the doped layer is decreased. A notable feature of these structures is that the mobility strongly increases in the 2D limit. This aspect suggests that a new regime of 2D superconducting phase transitions can be experimentally accessed approaching the clean limit. Prospects and progress in understanding the spin-orbit coupling and coupling bi-layer structures will be discussed.

### BA03

#### Resonant x-ray scattering from YBCO family

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Motivated by the search for the mechanism of high-temperature superconductivity, an intense research effort has been focused on the evolution of the spin excitation spectrum upon doping from the antiferromagnetic insulating to the superconducting states of the cuprates. Here we take advantage of the recent developments of high-resolution resonant inelastic x-ray scattering [1]to show that a large family of superconductors, encompassing the model compounds YBa2Cu4O8 and YBa2Cu3O7, exhibits damped spin excitations - or paramagnons - with dispersions and spectral weights closely similar to those of magnons in undoped, antiferromagnetically ordered cuprates over much of the Brillouin zone[2]. A numerical solution of the Eliashberg equations based on the experimental spin excitation spectrum of YBa2Cu3O7 reproduces its superconducting transition temperature within a factor of two and strongly supports magnetic Cooper pairing models for the cuprates. Charge excitations are also investigated and reveal evidences for charge order in the pseudogap state of underdoped compounds [3].

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#### **BA04**

### Evolving electronic structures of high-tc cuprates studied by compton scattering

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The evolution of electronic structures with carrier doping is essential to understanding the mechanism of hightemperature superconductivity. Several spectroscopies, such as photoemission [1], x-ray absorption [2] and Compton scattering [3] have succeeded in probing the evolving nature of high-Tc cuprates and provided information about the electronic dispersions, Fermi surfaces, and orbital characters of doped carriers. In this presentation, we provide a new piece of information on the evolving electronic structure of La<sub>2</sub>,Sr<sub>4</sub>CuO<sub>4</sub> (LSCO) by means of Compton scattering. Compton scattering is complementary to other spectroscopies and uniquely measures the ground-state electron momentum distribution, i.e. Compton profile, of all electrons in a compound. This is bulk-sensitive because of a high-energy photon-in and photon-out technique. We have performed Compton profiles measurements of LSCO with x=0, 0.08, 0.15 and 0.3, and obtained two-dimensional electron momentum density (EMD) and k-space electron occupation density (EOD) projected along the c-axis. The EMD and EOD show a doping x dependency and their evolution is discussed in the light of some scenarios.

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### **BB01**

### Valence fluctuations and their possible role in stabilizing the correlated electron state in the system Ce<sub>1-x</sub>Yb<sub>x</sub>CoIn<sub>5</sub> M. Brian Maple\*

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The extraordinary correlated electron phenomena found in compounds containing lanthanide ions with an unstable valence (Ce, Pr, Sm, Eu, Tm, Yb) can be traced to the hybridization of localized 4f and conduction electron states. The physics underlying the correlated electron phenomena is associated with the Kondo effect for moderate hybridization in which the valence (4f shell occupation) is nearly integral and "valence fluctuations" for strong hybridization where the valence is nonintegral. After a brief review of "valence fluctuation" phenomena, we describe recent experiments on the superconducting heavy fermion system Ce<sub>1x</sub>Yb<sub>2</sub>Coln<sub>5</sub> which reveal that the correlated electron state is stabilized throughout the range  $0 \le x \le 0.8$ , apparently due to cooperative behavior of Ce and Yb ions involving their unstable valences [1]. Interestingly, the superconducting critical temperature decreases linearly with x from 2.3 K at x = 0 towards 0 K at x = 1. Non-Fermi liquid behavior is observed at low temperature, although there is no readily identifiable quantum critical point [2]. The intermediate valence of Yb in Ce<sub>1x</sub>Yb<sub>2</sub>Coln<sub>5</sub> has been verified by ARPES, EXAFS, and XANES measurements [3,4]. Research supported by U.S. DOE Grant No. DEFG02-04ER46105 (crystal growth) and NSF Grant No. 0802478 (low temperature measurements).

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### **BB02**

### Electronic structures of novel Ce-based systems via photoemission spectroscopy

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Ce-based intermetallic compounds exhibit novel physical properties due to the unique character of the localized Ce 4f electrons that have a large hybridization with the conduction-band electrons. In this talk we will describe the photoemission spectroscopy (PES) study of the electronic structures of the Ce compounds that show diverse physical properties, such as superconductivity, pseudo-gap formation, chargedensity-wave (CDW) formation, and colossal magneto-resistance. We will discuss on CeNiSn, CeTe<sub>2</sub>, CeRu<sub>2</sub>, and La<sub>0.7</sub>Ce<sub>0.3</sub>MnO<sub>3</sub>, for which the valence states of Ce ions vary from being trivalent (Ce3+) to being mixed-valent and tetravalent (Ce4+), depending on the strength of the hybridization. Our study shows the different roles of Ce 4f states in determining their ground states. In particular, we will report our recent angle-resolved PES (ARPES) study of the quasi-two dimensional CDW system CeTe2 in comparison with the band-structure calculations [1,2]. Our ARPES study provides the information on the role of the Ce 4f states in the CDW formation, the carriers near the Fermi level, the effect of band-folding due to the interaction with Ce-Te(2) layers on the Fermi surface, and the CDW modulation vector in CeTe<sub>2</sub>, Keywords: ARPES, CDW, electronic structure, quasi-two dimensional system

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### **BB03**

#### Valence state and spin state of Fe in SrFe<sub>1-x</sub>(Sc,Sn)<sub>x</sub>O<sub>3</sub> perovskites

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The coexistence of Fe<sup>3+</sup> and Fe<sup>4+</sup> in SrFeO<sub>3</sub>-d, in relation with oxygen vacancies, influences the electronic conductivity and is at the origin of the giant negative magnetoresistance effect [1,2]. With the aim to investigate the influence of the valence state of Fe on the physical properties (magnetism, transport) in SrFe<sub>1-x</sub>M<sub>x</sub>O<sub>3</sub>-d perovskites (M=Sc,Sn and O<sub>2</sub>x<1) were prepared by solid state synthesis method and investigated by x-ray diffraction, 57Fe Mossbauer spectrometry and susceptibility measurements. Replacing Fe ions with ions with a fixed valence state (Se<sup>2+</sup> or Sn<sup>4+</sup>) modifies the Fe<sup>3+</sup>/Fe<sup>4+</sup> ratio and the oxygen content. Introducing Sc<sup>3+</sup> ions leads to a decrease of both Fe<sup>4+</sup> and oxygen contents without affecting the crystal structure. However, the substitution is limited. In the SrFe<sub>03</sub>Sc<sub>03</sub>O<sub>25</sub> limit compound, only Fe<sup>3+</sup> ions are detected, in two different spin states. Introducing Sn<sup>4+</sup> ions leads to a decrease of the oxygen content. The substitution is not limited Fe<sup>4+</sup> and Fe<sup>4+</sup> ions and Fe<sup>3+</sup> and Fe<sup>4+</sup> ions is discussed. This work was supported by the French Agence Nationale de la Recherche (project ANR-08-BLAN-0005-01).

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### **BB04**

### Physics of cerusn studied on a single crystal

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The recently reported hysteretic anomalies of electrical resistivity and magnetization of CeRuSn between 160 K and 290 K are most probably associated with crystal structure transitions intimately coupled with charge ordering at a Ce ion of unstable valence [1,2]. To shed light on the intriguing phenomena in this monoclinic compound, we have grown a single crystal and measured single crystal X-ray diffraction, thermal expansion, magnetization, and electrical resistivity in a wide temperature and pressure range. We also investigated the role of the lattice volume. Contrary to reports our high-quality single crystal shows two distinct anomalies at ~ 290 and 220 K (when cooling), respectively. The anomalies are strongly anisotropic and show a large temperature hysteresis (~ 15-20 K). The crystal shrinks by 0.7 % along the c-axis between 300 and 200 K while it changes little along a- and b- Exerting hydrostatic pressure shifts both anomalies to higher temperatures by more than 11 K kbar-1 This corroborates the scenario attributing the anomalies to transitions coupled with Ce-ion valence changes. Antiferromagnetism below TN = 2.7 K has been suggested [1]. We will present the low-temperature magnetization, resistivity and specific heat behavior of the single crystal and discuss the antiferromagnetism in detail

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### **BC01**

### **Ouantum effects in molecular single-ion magnets**

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Crystals of single-molecule magnets provide unique opportunities for the investigation of quantum magnetism. Generally, these systems are formed by clusters made of a polynuclear magnetic core. However, it has been recently shown that molecules encapsulating an individual magnetic ion can also exhibit SMM behavior. In these systems, known as single-ion magnets, the quantum effects are much more pronounced than in the more complex cluster-type SMMs. In this paper we will show that the spin dynamics, and in particular the quantum tunneling effects of these systems, can be controlled by chemical design (symmetry control of the magnetic site, magnetic dilution of the crystals,...). On the other hand, we will show that these molecules can be potential candidates to act as spin qubits as they possess a well defined ground state doublet with a high quantum coherence.

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### **BC02**

### Direct observation of a ferri-to-ferromagnetic transition in a fluoridebridged 3d-4f molecular cluster

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Single-molecule magnets are exchange-coupled molecular clusters exhibiting slow relaxation of magnetization. Introduction of rare-earth ions into such clusters may boost the magnetic anisotropy and thus increase the barrier for magnetization reversal. In this contribution we present a X-ray magnetic circular dichroism (XMCD) study of a novel fluoride-bridged molecular nanomagnet DyCrDy. By measuring element-specific magnetization curves we are able to directly observe the field-induced transition from a ferrimagnetic to a ferromagnetic arrangement of the Dy and Cr magnetic moments. Using a spin-Hamiltonian model we obtain excellent fits to the XMCD data and, further, we are able to extract the strength of the Dy-Cr exchange coupling. The value found from XMCD is consistent with SQUID magnetometry and inelastic neutron scattering measurements. Furthermore, based on this approach we investigate other 3d-4f molecular nanomagnets.

J. Dreiser, K. S. Pedersen, C. Piamonteze, S. Rusponi, Z. Salman, Md. E. Ali, M. Schut-Magnussen, C. Aa. Thuesen, S. Piligkos, H. Weihe, H. Mutka, O. Waldmann, P. Oppeneer, J. Bendix, F. Nolting, H. Brune, Chemical Science, DOI:10.1039/C2SC00794K, in press.

### BC03

### Magnetic -filed and angular dependence of magnetism in the triangular mott insulator κ-(BEDT-TTF)2Cu<sub>2</sub>(CN)<sub>3</sub> investigated by 13C NMR

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The organic conductors, ĸ-(BEDT-TTF)2X, are prototype for investigating Mott physics and spin frustrations. The X=Cu[N(CN),]Cl is a Mott insulator which undergoes an antiferromagnetic phase transition[1]. On the other hand, X=Cu2(CN)3 which has a triangular lattice does not show a long range magnetic ordering[2]. This suppression is believed to deeply relate to strong spin frustrations Both salts change to non-BCS-type superconductors via band-width control Mott transition[3] While X=Cu<sub>2</sub>(CN)<sub>2</sub> does not show magnetic ordering, we previously reported anomalous behaviors in 13C NMR around 6 K[4], where heat capacity, thermal conductivity and lattice constant show anomalies as well[5]. NMR line broadening is enhanced below 6 K. The 1/T1 monotonically decreases down to 6 K, and then becomes temperature-insensitive, followed by a decrease below 1 K under 8 T. So, the 6 K anomaly is a key phenomenon for understanding the origin of absence of magnetic order. We have measured external-filed angular dependence of 13C NMR under magnetic fields up to 15 T for clarifying the origins of the line broadening and the 1/T1 anomaly around 6K At room temperature, angular dependence of spectra is well explained by a crystal structure. We will show the detailed experimental results and discuss the low temperature states.

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### **BC04**

### NMR study of quantum spin liquid in an organic triangular lattice antiferromagnet EtMe<sub>3</sub>Sb[Pd(dmit)<sub>2</sub>]<sub>2</sub>

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A quantum spin liquid state is realized in an organic spin 1/2 antiferromagnet on the triangular lattice, EtMe<sub>3</sub>Sb[Pd(dmit)<sub>2</sub>], [1, 2, 3]. No classical magnetic ordering exists down to 19 mK due to the quantum fluctuation and spin frustration, although the exchange interaction is 240 K [2, 4]. The enriched 13C nuclear spin-lattice relaxation rate 1/T1 under a field of 7.65 T shows an exotic transition at 1.0 K accompanied with symmetry breaking or topological ordering. The rate 1/T1 indicates that the excitations have no spin-gap above 1.0 K [1, 2], while below 1 K 1/T1 shows the appearance of an excitation gap, which is a nodal one rather than a full gap [3]. We have confirmed that our NMR data surely reflect the spin state of this system by using two kinds of enriched samples [5]. The transition temperature has the field dependence, which is opposite to that of an ordinary spin system accompanied with a spin-gap formation owing to dimerization. Our recent measurement for a single crystal has revealed that 1/T1 depends on the field direction below 5K. We will discuss the nature of quantum spin liquid, comparing with another thermal properties of this material.

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### BC05

### Wilson ratio of a Tomonaga-Luttinger liquid in a spin-1/2 Heisenberg ladder

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The Wilson ratio R w one of the crucial parameters that characterize Fermi liquids has served as a powerful tool to classify heavy fermion systems by quantifying the strength of spin fluctuations. This parameter remains crucial even in one dimension. In particular in one-dimensional antiferromagnets R K is related to the Tomonaga-Luttinger parameter K through the relation R W=4K. We have tested this hitherto overlooked relation in (C<sub>4</sub>H<sub>12</sub>N)<sub>2</sub>CuCl<sub>4</sub> (DIMPY), an S=1/2 quantum spin ladder with strong leg interactions [1]. Using micromechanical force magnetometry, we have measured the magnetization at temperatures down to 45 mK in magnetic fields up to 9 T. Both the temperature and field dependences are in good agreement with quantum Monte Carlo results, evidence for the excellent onedimensionality of the material. Particularly, a minimum of magnetization as a function of temperature is found, which is an onset of a Tomonaga-Luttinger liquid state predicted by theory[2]. The magnetic susceptibility in conjunction with specific heat yields R W [3] which, along with density-matrix renormalization-group calculations, supports the above relation between R W and K.

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### **BD01**

#### Exchange bias; where are the pinned uncompensated spins Ivan K Schuller1\*, M. Erekhinsky2, R. Morales3, I. V. Roshchin4, M. Kovylina5, A. Labarta5, X. Batlle5, M. Fitzsimmons6, S. Bar-ad7 and S. K. Sinha8

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Exchange bias in ferromagnetic-antiferromagnetic (FM-AFM) nanostructures has resisted a universal explanation in spite of its importance in the sensors and storage industries. Unusual phenomena such as; i) asymmetric FM reversal, ii) changes of the superparamagnetic transition temperature of magnetic nanostructures, iii) large exchange bias in nominally fully compensated surfaces, iv) positive exchange bias emergence for certain classes of bilayers, v) anomalous reversal at fast time scales (<300psec), vii) effect of the bulk AFM magnetic structure of the AFM and vii) noticeable change in the microscopic spin configuration appear. We have used nanostructured systems prepared by Molecular Beam Epitaxy, electron beam lithography, self-assembly and focused ion beams to investigate the common origin of all these phenomena. Comprehensive magnetization, magnetotransport, Kerr effect, neutron scattering and synchrotron radiation studies together with theoretical investigations provide a universal explanation regarding the origin of the exchange bias and location of the crucial pinned uncompensated spins. This opens up the possibility for the design of new systems and provides a road for the understanding of the mechanism controlling Exchange Bias. Work supported by the US Department of Energy, TAMU-CONACyT and the Air Force Office of Scientific Research.

### **BD02**

#### Atomic diffusion in (Pt/Co)3/IrMn multilayers

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The magnetic properties of exchanged-coupled multilayers strongly depend on the structure and chemistry of the interfaces. This is the case in particular for exchange-bias multilayers, in which a ferromagnetic layer is coupled with an antiferromagnetic layer, leading to a shift of the hysteresis loop [1]. Recent developments of the 3D Atom Probe Tomography (APT) in the Groupe de Physique des Materiaux allowed investigating the interfaces in exchange-bias (Pt/Co)3/IrMn magnetic multilayers with perpendicular magnetic anisotropy [2,3]. The APT analysis allows determining the chemical composition of the nanometric layers, thus providing an accurate structural characterization of the layers, leading to an estimation of their thickness, roughness, atomic concentration and width of their interfaces. The magnetic properties can thus be correlated with the structural properties. It has been showed that the presence of a Pt spacer at the Co/IrMn interface plays a determining role on the exchange-bias phenomenon, in relation with the Mn/Co interdiffusion process. We present here an investigation of the diffusion process at high temperature (300-700°C) by APT. The results are discussed in relation with the thermodynamic properties of the Pt-Co-Ir-Mn system.

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### BD03

### Magnetic properties of ferromagnetic-antiferromagnetic bi-layers with different spin configuration

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We investigated the effect of different spin direction of anti-ferromagnetic(AFM) layer on the magnetic properties of ferromagnetic(FM) layer in Fe-NiO and Fe-CoO bi-layer systems. For Fe-NiO system, we prepared successfully Fe/NiO/Ag(001) and Fe/NiO/MgO(001) systems on a single MgO(001) substrate. We examined magnetic properties of the bi-layer system using the surface magnetic optical Kerr effect(SMOKE) and X-ray magnetic linear dichroism(XMLD). From SMOKE measurement we observed the coercivity enhancement due to the set-up of AFM order of NiO films in both of the Fe/NiO/MgO(001) and Fe/NiO/Ag/MgO(001) system. The most remarkable results in our observation is that the coercivity enhancement of Fe/NiO/ Ag/MgO(001) is much larger than that of Fe/NiO/MgO(001). XMLD experiments confirmed the out-of-plane spin direction of NiO layers in Fe/NiO/MgO(001) and in-plane spin-direction of NiO layers in Fe/NiO/Ag/MgO(001), and we concluded that the origin of large enhancement of coercivity is due to the strong parallel coupling between Fe layers and NiO layers. We also confirmed that this strong parallel coupling maintained across the thin Ag layer inserted between Fe and NiO layers. With this Ag inserted Fe/NiO system, we could estimate the Neel temperature of the NiO lavers. For Fe-CoO system, we observed much larger coercivity enhancement in Fe/CoO/Ag(001) than in Fe/CoO/MnO(001).

### **BD04**

Tuning exchange bias in Ni/FeF<sub>2</sub> heterostructures using antidot arrays

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The transition from negative to positive exchange bias (NEB/PEB) can be systematically tuned with antidot arrays artificially introduced into Ni/FeF, ferromagnetic (FM)/antiferromagnetic (AF) heterostructures [1]. This is a consequence of the energy balance between the Zeeman coupling of AF spins and the AF coupling at the FM/AF interface. The nanostructure plays a key role [1,2]: the antidot carving produces regions of locally pinned uncompensated spins throughout the antidot faces of the FeF, and these non interfacial magnetic moments favor the onset of PEB at lower cooling fields, by increasing the Zeeman energy of AF domains and favoring the alignment with the external field. Those non interfacial AF spins, and the pinned uncompensated interfacial AF spins responsible for the EB (loop shift), align simultaneously with the cooling field since they belong to the same domain and become pinned below the Neel temperature. As a result, the cooling field required for PEB decreases one order of magnitude for the antidot arrays, as compared to the unpatterned heterostructures. This is relevant concerning the applicability of FM/AF bilayers as multi-state storage media [3], since the minimum cooling field necessary to obtain PEB establishes the magnetic field required for multi-state storage writing.

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### **BD05**

### Correlation between training effect and hysteretic behavior of angular dependence of exchange biasing in polycrystalline ferromagnet/ antiferromagnet bilayers

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EB training effect is important to the spintronic devices, where the exchange field HE and the coercivity HC decrease during consecutive hysteresis loops. Alternatively, hysteretic behavior of angular dependence of EB has also been studied extensively. AFM spins are believed to play a crucial role in above two EBrelated phenomena. Simultaneous studies of these two phenomena in polycrystalline FM/AFM bilayers will be helpful to further reveal the effect of AFM spins. In this work, we report on the EB training effect and hysteretic behavior between clock-wise (CW) and counter clock-wise (CCW) rotations in polycrystalline NiFe/FeMn bilayers. Using the vector VSM, we have measured the pinning direction at remanent state after each hysteresis loop. When the external magnetic field is off-aligned the initial pinning direction, HE and HC are reduced. At the same time, the pinning direction is rotated during consecutive hysteresis loops. During CW and CCW rotations, the angular dependence of either HE or HC, or pinning direction does not overlap with each other. The hysteretic behavior and training effect become weak with increasing FeMn layer thickness. At ultrathin FeMn layers, these two effects vanish. These experimental results can be explained in the framework of the thermal activation model.

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### **BE01**

### Interlayer exchange coupling in ferromagnetic semiconductor GaMnAs-based multilavers

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The study of ferromagnetic (FM) dilute magnetic semiconductors (DMSs) continues to be of great interest because of their potential for spin-electronic device application. The nature of interlayer exchange coupling (IEC) in GaMnAs-based multilayers is an important issue from the point of view of possible spin-electronic ("spintronic") applications. The ability to manipulate such IEC enables us to control the relative direction of magnetization in one magnetic layer relative to the other, which lies at the heart of giant magnetoresistance (GMR) devices in wide use in information processing. In this connection, it is important to establish under what conditions the IEC between adjacent GaMnAs layers is antiferromagnetic (AF) or ferromagnetic (FM), since manipulation of the IEC can then be directly applied to achieve GMR devices based on this material. In this talk we will describe magneto-transport, magnetization, and neutron reflectometry experiments applied to GaMnAs-based multilayer structures consisting of GaMnAs layers separated by non-magnetic GaAs spacers. These measurements serve to identify conditions under which AFM coupling will occur in such GaMnAs/GaAs multilaver systems, thus providing us the information which can be used for manipulating magnetization (and thus also GMR) in structures based on the ferromagnetic semiconductor GaMnAs.

### **BE02**

Experimental probing of the magnetic order in ultrathin (Ga,Mn)As M. Sawicki<sup>1</sup>, D. Chiba<sup>2</sup>, O. Proselkov<sup>3</sup>, A. Korbecka<sup>4</sup>, Y. Nishitani<sup>5</sup>, F. Matsukura<sup>6</sup>, J. A. Majewski<sup>4</sup>, J.

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Colossal negative magnetoresistance and the associated insulator-to-metal transition are arguably the most characteristic features of magnetic and dilute magnetic semiconductors. Particularly intriguing is the case of (Ga,Mn)As and related systems in which spin-spin coupling is mediated by holes. In order to find out how magnetism evolves when the carrier density is diminished, we have probed magnetization changes induced by: (i) electric field in metal/insulator/(Ga,Mn)As structures [1] and (ii) thinning of the channel by sequential chemical etching. Our findings show that both the channel depletion and thinning results in a monotonic decrease of Curie temperature and spontaneous magnetic moment, with no evidence for the maximum expected within the impurity-band models but explained theoretically in terms of the p-d Zener model, modified for the thin film case [1]. They also confirm the presence of the theoretically predicted spin reorientation transition. We have found that the transformation to a non-magnetic state proceeds via the emergence of a hitherto non-revealed superparamagnetic-like spin arrangement, driven by critical fluctuations in the local density of states, specific to the Anderson-Mott quantum transition. Within our model, these nanoscale fluctuations of magnetization give rise to critical scattering which vanishes when localized spins are ordered by an external magnetic field.

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#### BE03

#### Origin of ferromagnetism in Ga<sub>1-x</sub>Mn<sub>x</sub>N

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Despite Ga<sub>1x</sub>Mn<sub>x</sub>N is the model system for a broad class of dilute magnetic semiconductors and oxides, diverging opinions about its magnetic ground state have been reported: this material was found to be either spin glass, or an antiferromagnet or a ferromagnet with rather disperse values of the Tc, ranging from below 10 K up to over 300 K [1,2]. By employing a range of ion beam, diffraction, and spectroscopy methods we have demonstrated [3,4] that particular magnetic signatures can be linked to a specific distribution of the magnetic ions and to their charge state, both affected by growth conditions and by the presence of shallow impurities. By combining extensive epitaxy and nanocharacterization programs with high sensitivity mK-SQUID measurements, we demonstrate that Ga. Mn N with randomly distributed  $Mn^{3+}$  ions is a ferromagnet with Tc below 5 K for x < 5%. To clarify the origin of ferromagnetism in this material with no band carriers, we examine  $Tc(x) \sim x^m$ and show that the value m = 2 is consistent with ferromagnetic superexchange. We discuss our results in the light of the available ab initio studies and compare our system to Cr chalcogenide and oxide spinels, where ferromagnetic and antiferromagnetic superexchange coexist.

Spinchs, whice Enformagnetic and anner formagnetic supercontinge obexist. [1] A. Bonanni and T. Dietl, Chem. Soc. Rev. 39, 528-539 (2010) [2] T. Dietl, Nature Mat. 9, 965-974 (2010) [3] W. Stefanovicz, D. Sztenkiel, B. Faina, A. Grois, M. Rovezzi, T. Devillers, A. Navarro-Quezada, T. Li, R. Jakiela, M. Savicki, T. Dietl, and A. Bonanni, Phys. Rev. B 81, 235210.1-14 (2010) [4] A. Bonanni, M. Sawicki, T. Devillers, W. Stefanowicz, B. Faina, Tian Li, T. E. Winkler, D. Sztenkiel, A. Navarro-Quezada, M. Rovezzi, R. Jakiela, A. Grois, M. Wegscheider, W. Jantsch, J. Suffezynski, F. D'Acapito, A. Meingast, G. Kothleitner, and T. Dietl, Phys. Rev. B 84, 035206.1-11 (2011).

### **BE04**

#### The effects of non-magnetic dopant on semiconductor materials Caihong Zhang\* and Dickon H. I. Ng

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# WITHDRAWN

### **BE05**

I-Mn-V room temperature antiferromagnetic semiconductors

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Recent observation of a large magnetoresistance in an antiferromagnet (AFM) based spin-valve opens the prospect for utilizing AFMs in spintronics [1]. This motivates a search for new materials which may be suitable for spintronics and are room-temperature AFMs. The desired control of devices via electrical fields could be further exploited if the materials involved would be also semiconductors. We present our results [2] in the growth and characterization of a new family of room-temperature antiferromagnetic semiconductors and semimetals. We show X-ray diffraction, magnetization, transport, and differential thermal analysis measurements of chemically synthesized bulk CuMnAs and CuMnSb. Experiments are complemented by theory calculations of band structure and magnetic structure. The progress in bulk I-Mn-V materials encourages also the investigation of thin layers. We will show the epitaxial growth of LiMnAs [3] and CuMnAs on conventional GaAs. Scanning tunneling microscopy has shown that in the case of LiMnAs there is a bandgap [4].

[1] B. G. Park et al. Nature Materials 10, 347 (2011) [2] F. Maca et al. J. Maen. Mater. 324, 1606 (2012) [3] T. Junewirth, et al. Phys. Rev. B 83, 035321 (2011) [4] I. Wijnheijmer, et al, submitted

### **BF01**

#### The verway phase of magnetite - a long-running mystery in magnetism I Paul Attfield

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Magnetite (Fe<sub>3</sub>O<sub>4</sub>) undergoes a complex structural distortion and becomes electrically insulating below 125 K. Verwey proposed in 1939 that this transition is driven by a charge ordering of Fe<sup>2+</sup> and Fe<sup>3+</sup> ions, but the low temperature state has remained contentious as twinning of crystal domains hampers diffraction studies of the structure. A variety of ground state models has been proposed over seven decades of study. We have recently reported the full low temperature superstructure of magnetite, determined by high energy x-ray diffraction from a 40 µm grain, and we identify the emergent order. The acentric structure is described by a superposition of 168 frozen phonon modes all with amplitude <0.24 A. Distortions of the FeO<sub>4</sub> octahedra show that Verwey's hypothesis is correct to a first approximation. However, anomalous shortening of some Fe-Fe distances shows that the localised electrons are distributed over linear three-Fe-site units that we describe as 'trimerons' - this description is supported by band structure calculations. The charge order and trimeron distortions induce substantial off-centre atomic displacements and couple the resulting large electrical polarisation to the magnetisation so that multiferroic behaviour is expected. Trimerons may be significant quasiparticles above the Verwey transition and in other transition metal oxides.

M.S. Senn, J.P. Wright and J.P. Attfield, Nature 481, 173 (2012). M.S. Senn, I. Loa, J.P. Wright and J.P. Attfield, Phys. Rev. B, in press (2012).

#### **BF02**

### Possible link of a structurally driven spin flip transition and the insulator-metal transition in the perovskite La1-xBaxCoO3 Despina Louca and Peng Tong

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The intricate nature of the magnetic ground state near the insulator-metal transition (IMT) in La1-xBaxCoO3 was investigated via neutron scattering. For x less than the critical concentration, xc~0.22, a commensurate antiferromagnetic (AFM) phase initially appears. As x approaches xc, the AFM component continuously weakens while ferromagnetic (FM) order sets in the rhombohedral lattice. The two magnetic phases appear to be growing in different domains and have different ordering temperatures, with the FM order parameter setting in first at higher temperatures while the AFM order parameter occurs at lower temperatures. At xc, a spin flip to a new FM state occurs while the crystal transforms to an orthorhombic (Pnma) symmetry. The magnetic Pnma phase coincides with the minimum saturation reached in the resistivity. It is proposed that the orbital overlap in the Pnma phase is the most conducive to charge hopping. The structurally driven spin flip creates metallic FM droplets and the percolation of these droplets drives the system to become metallic. This, in turn, suggests that the FM domains in the rhombohedral phase are in fact insulating. This is fundamentally distinct from other magnetoresistive perovskites where only one FM transition is observed prior to the IMT.

### BF03

Slow magnetic crossover in the frustrated magnet Ca<sub>2</sub>Co<sub>2</sub>O<sub>4</sub> Stefano Agrestini1\*, Martin R Lees2, Catherine L Fleck2, Oleg A Petrenko2, Laurent C Chapon3, Claudio Mazzoli4 and Alessandro Bombardi5 1 Max-Planck Institut CPfS, Dresden, Germany <sup>2</sup> Department of Physics, University of Warwick, Coventry, United Kingdom <sup>3</sup> ILL Grenoble France 4 Dipartimento di Fisica, Politecnico di Milano, Milano, Italy <sup>5</sup> Diamond Light Source, Didcot, United Kingdom

Time-dependent phenomena play a crucial role in the properties of many magnetic systems including spin-glasses, single-molecule magnets and superparamagnets. However time-dependent magnetism is not expected to be observed in the presence of long-range magnetic order. Here we challenge this view by reporting the observation of a new time-dependent behavior where a transition from one long-range magnetically ordered state to another occurs in zero magnetic field over a time scale of several hours. We performed an extensive neutron diffraction study of the spin chain system Ca3Co2O6 whose step-like magnetization versus magnetic field aroused great interest in the scientific community. Our data show that for T< 14K< TN =25K the spin-density-wave order observed immediately below TN, becomes unstable and a commensurate antiferromagnetic phase appears via a very slow transformation process. As the temperature is reduced the characteristic time of the transition process increases rapidly and at low temperatures the magnetic states become frozen. The transition is also noteworthy because the two phases have different propagation vectors. Very rarely transitions between two magnetically ordered phases involve a change of translational symmetry. This discovery sheds new light on the peculiar dynamic properties of Ca3Co2O6. S. Agrestini et al., Phys. Rev. Lett. 106, 197204 (2011).

### **BF04**

### Incommensurate magnetic states in itinerant systems

#### Marat Timirgazin and Anatoly Arzhnikov

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We study the ground state of the Hubbard model within the Hartree-Fock approximation taking into account both the incommensurate magnetic phases (spiral and collinear spin-density waves) and their possible participation in the phase separation (PS). We construct magnetic phase diagrams for 2D and 3D lattices with different values of the parameter of the electron hopping between the next-nearest neighbors [1]. We show that the incommensurate magnetic state is the ground one in a wide range of the model parameters. There is a large region of the PS in the vicinity of half-filling. The data obtained agree well on a qualitative level with experimental data on the magnetic state of a superconducting compound La2-pSrpCuO4 [2]. We consider the metal-insulator transition (MIT) at half-filling taking into account incommensurate magnetic states. Calculations show that, for the square lattice, MIT coincides with the first-order magnetic transition from a spin-spiral to antiferromagnetic state. To our knowledge this result is the first example of MIT between two magnetically ordered states in the 2D lattice. Previously, the critical dimension for this transition was considered to be in a range from 2 to 3. Thus, our investigation proves this value to be helow 2

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### **BF05**

#### Canted spins of Mn<sub>3</sub>0<sub>4</sub> investigated by 55Mn<sup>2+</sup> and 55Mn<sup>3+</sup> nuclear magnetic resonance in magnetic field

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For a long time, huge magnetic anisotropy in Mn<sub>3</sub>O<sub>4</sub> has attracted a lot of interest from researchers. The magnetization along the c-axis is smaller than that along the ab-plane even in 30 Tesla. It has been believed that this phenomena occurs due to a uniaxial anisotropy of Mn3+ along the c-axis. However, there has been no experiment to verify this theory. We measured 55Mn<sup>2+</sup> and 55Mn<sup>3+</sup> Nuclear Magnetic Resonance of a Mn<sub>3</sub>O<sub>4</sub> single crystal in 7 Tesla external magnetic field. The angle between each magnetic moment and the field was calculated from the spectral shifts obtained for various magnetic field directions between the a, b and c-axes. From the fact that Mn<sup>3</sup> moments are tied to the ab-plane when the magnetic field is in the direction along the c-axis, we concluded that the Mn3+ moment has a large in-plane anisotropy contrary to the previous reports. In contrast, Mn2+ moments were tilted toward the field direction, implying that the Zeeman energy of a Mn2+ moment is competing with the energy of the exchange interaction between Mn<sup>2+</sup> and Mn<sup>3+</sup> moments and the magnetic anisotropy energy of Mn2+

### **BG01**

### L10 ordered FePt based granular films for thermally assisted magnetic recording

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L10 ordered FePt is considered to be one of the candidate for ultrahigh density perpendicular recording media beyond 1 Tb/in2 due to its high magnetocrystalline anisotropy. We recently reported well controlled microstructure of FePtAg-C granular film with an average particle size of about 6 nm and small size distribution of about 20 % [1]. Furthermore, we demonstrated areal density of 550 Gb/in2 by a static test using a head for thermally assisted magnetic recording [2] Although we achieved a particle size of 6 nm with narrow size distribution, the FePt-C system has a weakness in attaining small surface roughness with a columnar structure due to the strong phase separation tendency between FePt and C in not only the lateral direction but also the growth direction. To attain columnar structure with smooth surface, we chose TiO<sub>2</sub> as a matrix. In this presentation, we will show FePt granular film with columnar structure and smooth roughness.

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### **BG02**

### Nanogranular FePt films for thermally assisted recording

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In order to extend the data storage density in hard disk drives beyond 1 Tb/in2, using a high crystalline anisotropy (Ku) material as the magnetic media layer is required to reduce the grain size further while maintaining thermal stability. However, the magnetic field applied by a conventional write head is not strong enough to reverse the magnetization of the high-Ku media. A promising approach to recording using high-Ku materials is thermally-assisted recording (TAR), in which the media is locally heated near the Curie temperature by an optically-generated heat pulse in the presence of a magnetic field, both applied by a laser-integrated write head, in order to write a bit. Given its moderate Tc, high Ku, corrosion durability, and high dKu/dT near Tc, FePt alloy with L10 crystalline order is a potential candidate for TAR media. We deposit FePt nano-granular films by sputtering on glass substrates at elevated temperature. Underlayer materials are selected for heat-sinking and to obtain high, out-of-plane L10 order Structural and magnetic characterization and transmission electron microscopy of the media show that we can achieve properties that are promising for TAR, such as isolated grains with an average grain size <7.5nm, a low size distribution and  $Ku > 4.5 \times 10^7 \text{ erg/cm}^3$ 

### **BG03**

### Single crystalline isolated grains of L10-ordered FeCuPt prepared by combination of rapid thermal annealing with rapid cooling and additional annealing

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We reported1),2), that the Rapid Thermal Annealing (RTA) of Fe/Cu/Pt multilayered continuous films are effective to obtain perpendicularly magnetized small L10-FeCuPt grains on thermally oxidized Si substrate. In this report, we introduced a rapid cooling process (RCP) into RTA and compared with the natural cooling processes. FeCuPt grains were fabricated under the variety of heating up rate (TR) and the maximum temperature (Tm). Isolated grains are obtained at higher TR and Tm. With RCP, growing of grains are prevented, however new peaks in XRD (X-ray diffractometer) profile are appeared at slightly lower angle of (002) and it may be correspond to the disordered structure of FeCuPt. From electron diffraction patterns and dark field images by TEM, mostly L10-ordered polycrystalline structure was observed. Therefore, we preformed additional annealing for above isolated FeCuPt grains (600 degree C, 1 hour). Annealing condition was decided from an estimation of atomic diffusion length3). After additional annealing, grains kept almost similar size and the new peak was banished in XRD profile. Furthermore, we observed single crystalline L10-ordered structure from electron diffraction pattern of a single grain. As a result, single crystalline L10-FeCuPt isolated grains are obtained by RTA with rapid cooling and additional annealing

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### **BG04**

### Oscillation characteristics of spin-torque oscillator calculated using integrated simulator with spt writer

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Microwave-assisted magnetic recording (MAMR) requires a spin-torque oscillator (STO) to induce magnetic resonance in the recording medium. The oscillation characteristics of an STO have previously been investigated using micromagnetic simulators based on spatially isolated models [1]-[3], which neglect the magnetostatic interactions between an STO and a Single-Pole-Type (SPT) writer. We have developed a new simulator that takes into account the magnetostatic interaction between an SPT writer and an STO, and we refer to this as an integrated MAMR simulator. This interaction was found to have considerable influences on the oscillation behaviors of an STO. Simulations were carried out for an STO with a lateral length of 30 nm and a thickness of 10 nm inserted in the gap between a main-pole and a floating-type trailing shield. The stray field in the gap increased by approximately 30 % by inserting the STO and its oscillation frequency was 8 GHz. For comparison, the oscillation frequency calculated using a conventional isolated STO simulator was found to be 16 GHz. This decrease in frequency may be the result of an energy loss in the main-pole of an SPT writer that is excited by the high-frequency field from an STO.

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#### **BH01**

#### Magnetic nanotechnology for cancer treatment

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Often the conventional therapies needed to annihilate cancer cells have high systemic toxicity effects. We are developing systems which utilise magnetic nanoparticles loaded with therapeutic agents, directed and guided towards the tumour site and preferentially attached to the tumour to destroy it. The cancer cells are killed by the combined action of the drugs and a local temperature increase in the area around the tumour site using externally applied electromagnetic irradiation (hyperthermia). The aim is to avoid most side effects since the nanoparticles target specifically the tumour area and leave most of the surrounding tissues unaffected, therefore higher doses of the drug reach the target area even if lower overall doses are administered. Our preliminary clinical trials and pilot studies have demonstrated significant effects in the reduction of the tumour size and a 50% increase in the survival rates (on average) of patients with breast, lung, endometrial, colorectal, malignant melanoma, etc. For example, in patients with lung cancer IIIA stage we have demonstrated a tripling in the 5-year survival rates. The unique aspects of our platform technology are (i) synthesis of nanoparticles with high magnetization, (ii) less tendency to agglomerate, (iii) generation of spatially nonhomogeneous electromagnetic fields.

### **BH02**

#### The low magnetic field effect of sanals of primo vascular system

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The motion features of sanals inside of the primo vascular system (PVS), that is so-called the Kyungrak system, are investigated under a low static magnetic field by using the anatomy technology and optical microscope. The sanals with a size of about 1 selected and separated from the primo vessel and node of the real PVS inside of the surface of the internal organs are observed from rabbits' abdominal wall and dipped with PBS liquid inside of petri dish. The sanal's moving velocity along the direction of magnetic field (xdirection) and perpendicular to the direction of magnetic field (y-direction) under the low magnetic field of 0 Oe, 20 Oe, 40 Oe, 60 Oe, and 80 Oe, respectively, is observed below a internal temperature of 38 oC. Ten sanals' moving velocities versus magnetic field are shown two differently dominant tendencies with an average velocity of 0.9 pixel/s and a random velocity according to the x-direction and y-direction, respectively. This experimental results imply that the rotating motion of sanal with nuclei DNA composed of many inorganic magnetic materials of Mn and Co is monotonically weakened by the increase of applied magnetic field.

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### **BH03**

#### Magnetic targeting of mesenchymal stem cells in the spinal cord Vaclav Vanecek1, Jiri Ruzicka1, Serhiy Forostiak1, Michal Babic2, Vit Herynek3, Alexandr Dejneka4, Vitalii Zablotskii4\*, Sarka Kubinova1, Pavla Jendelova1 and Eva Sykova1 Institute of Experimental Medicine, Czech Republic

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Mesenchymal stem cells (MSC) are currently under intensive study as a possible therapeutic agent in spinal cord injury [1]. MSCs labeled with superparamagnetic iron oxide nanoparticles (SPIO) can be manipulated by an externally applied magnetic field and successfully targeted to different organs. We transplanted SPIO-labeled MSC into rats with a magnetic implant or a non-magnetic substitution inserted above the spinal cord intrathecally via lumbar puncture one week after spinal cord injury. The fate of the grafted cells was monitored by means of magnetic resonance imaging (MRI) [2] and histologically. In our experiments, the transplanted cells were present mostly on the surface of the spinal cord (in the subarachnoid space) in both the magnetic and non-magnetic (control) groups. In the lesion area of the spinal cord, there was a significant difference in cell distribution between the magnetic and non-magnetic groups. The SPIO-loaded cells circulating in the subarachnoid space were gathered under the magnetic implant near the site of injury. Additionally, we suggest how magnetic cell targeting can be improved by using a magnet specially designed to produce spatially modulated multigradient stray fields [3]. The proposed magnetic delivery system might be attractive for clinical application development. Supported by P304/12/1370 AVOZ503905703

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### **BH04**

#### Tilted bianchi type - I magnetised viscous fluid cosmological model Subrata Kumar Sahu

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Tilted Bianchi type - I cosmological model for viscous fluid distribution embedded in a magnetic field is investigated. The role of the bulk viscosity and electromagnetic field in getting an inflationary phase and in establishing the tilted universe is discussed.

### BH05

### Growth of highly uniform graphene for spintronic applications

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Graphene has attracted wide attention for nanoelectronics and spintronics. In this study, in-situ analysis were performed on the graphene growth in ultrahigh vacuum chemical vapor deposition (UHV-CVD) by exposing the epitaxial Ni(111) thin film to benzene vapor at 873 K. It is shown that the uniform single- and bi-laver graphenes can be synthesized by the control of benzene exposure in the range of 10-105 langmuirs, reflecting a change in the graphene growth-rate by three orders of magnitude in between the first and second layer. It is also suggested that the chemical interaction between bi-layer graphene and Ni(111) is weak in comparison with that between single-layer graphene and Ni(111). Ex-situ micro-Raman analysis on the large area graphene film transformed on a SiO<sub>2</sub> substrate makes it clear that the structural and electrical uniformities can be improved remarkably under the specific exposure conditions at which the growth of the respective graphene layers are completed. The present results demonstrate that the UHV-CVD method enables the growth of highly uniform graphene which would be necessary for controlling the spin transport process as well as realizing a long spin-relaxation length in the graphene-based spintronic devices by optimizing the dosage of precursors and the growth temperature.

### **BH06**

### Depth-resolved XMCD spectroscopy on single-layer graphene / Ni structure

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The graphene-metal interaction is not only of fundamental interest but also of technological importance in view of spintronic applications. For example, a recent theoretical calculation predicts a spin filtering effect of graphene attached on the ferromagnetic metal surface under the currentperpendicular-to-the-plane configuration [1]. For the purpose of clarifying the electronic and magnetic states of the single layer graphene (SLG)/FM structure especially at the interface region, depth-resolved X-ray absorption and X-ray magnetic circular dichroism (XMCD) spectroscopy are performed on the SLG/Ni(111) structure grown by using the ultrahigh vacuum chemical vapor deposition and molecular beam epitaxy techniques. An intense XMCD signal is observed in the C K-edge XMCD spectrum taken under the surface sensitive condition. This indicates that there exists the Ni-induced spin polarization in the orbitals of SLG even at room temperature. It is also found that the XMCD intensity shows a dramatic change with the probing depth. The depth-resolved analysis of the Ni L-edge XMCD spectrum makes clear that the magnetization of Ni is decreased by 20-30 % within a few atomic layers from the interface with SLG, possibly associated with the formation of C-Ni bonds.

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### **BI01**

#### Magnetoresistance and spin-transfer torque in magnetic tunnel iunctions

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A magnetic tunnel junction (MTJ) consisting of a tunnel barrier sandwiched between two ferromagnetic electrodes exhibits the tunnel magnetoresistance (TMR) effect due to spin-dependent electron tunneling. Since the discovery of room-temperature TMR in the mid-1990s, MTJs with an amorphous aluminun oxide (Al-O) tunnel barrier have been studied extensively. Such MTJs exhibit a magnetoresistance (MR) ratio of several tens of percent at room temperature (RT) and have been applied to magnetoresistive random access memory (MRAM) and the read heads of hard disk drives. MTJs with MR ratios substantially higher than 100%, however, are desired for next-generation spintronic devices. In 2001, firstprinciple theories predicted that the MR ratios of epitaxial Fe/MgO/Fe MTJs with a crystalline MgO(001) barrier would be over 1000% due to the coherent tunneling of specific Bloch states. In 2004, MR ratios of about 200% were obtained for MgO-based MTJs [1]. MTJs with a CoFeB/MgO/CoFeB structure were developed for practical application and found to have MR ratios of above 200% and other practical properties [1,2]. The talk focuses on the physics of magnetoresistance and spin-transfer torque in MTJs and the application of MTJs to various spintronic devices such as magnetic sensors and spin-transfertorque MRAM especially with perpendicularly magnetized electrodes.

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### **BI02**

#### MTJ based non volatile logic for ultimate power management Tetsuo Endoh, Takashi Ohsawa, Takahiro Hanvu and Hideo Ohno Tohoku University, Japan

Recently, the achievement of the silicon-LSI's target performance is becoming difficult only by scaling technologies. First, the speed-gap between the operation speed of logic and each memory has expanded. Moreover, its power consumption increases rapidly, too. In this paper, we describe both MTJ based Non Volatile (NV-) Embedded SRAM and MTJ based NV-Logic circuit. Firstly, we show the proposed MTJ based SRAM that consists of 4 MOSFETs and 2 MTJs. This STT-SRAM achieves an excellent static noise margin that is larger than the resistive load SRAM counterpart by taking advantage of the MTJ's switching characteristics. The cell size becomes comparable with the conventional 6-transistor SRAM and even smaller as the MTJ size shrinks along with the MOSFET scaling. Next, from the viewpoint of super-low power consumption systems, MTJ based NV-Logic, that fuses the STTRAM with semiconductor CMOS logic, is discussed. Especially, we describe the proposed MTJ based NV-Latch circuit with high speed operation of 600MHz. It is shown that with using our novel MTJ based NV-Latch circuit, the NV-F/F type NV-Logic can be realized. Finally, we discuss the Logic-in-Memory type NV-Logic. From all, we show novel ultimate power management technique with both MTJ based Embedded NV-Memory and MTJ based NV-Logic.

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### **BI03**

### Spin gating a transistor and spintronics with antiferromagnets

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Spintronics is among leading technologies for reading, writing, and storing information. Spintronic sensors and memory bits rely on spin-dependent electron transport through the device and the same applies to the currently explored concepts for spintronic transistors. Here we demonstrate a new route for using spin in transistors by moving the spin functionality from the transport channel to the capacitively coupled gate. In our transistor, the spin state of the gate controls the charge transport through a conventional non-magnetic electron channel. We explain the phenomenon using the relativistic, spin-orbit coupling theory framework employed in recent studies of anisotropic magnetotransport in spintronic nanodevices. The relativistic magnetic anisotropy phenomena are an even function of the microscopic magnetic moment vector. This means that they can be equally strong in antiferromagnets as in ferromagnets which opens the possibility to realize spin dependent microdevices based on antiferromagnets. As a demonstration we present our experimental observation of >100% tunneling anisotropic magnetoresistance in a device with an antiferromagnetic tunnel electrode

B. G. Park et al. Nature Mater 10, 347 (2011) X. Marti et al. Phys. Rev. Lett. 108, 017201 (2011)

### **BJ01**

### Co2y-nicuzn ferrite composites with high permeability

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Hexagonal structure magnetoplumbite ferrites have revealed a higher dispersion frequency than that of nickel ferrites because of the magnetoplumbite's magnetic anisotropy. The magnetoplumbite ferrite densification temperature always exceeds 1000oC and the initial low temperature firing permeability of magnetoplumbite ferrites with added glass is too low ( $\mu i = 2-4$ ). Therefore, it is desirable to develop a material that has a higher permeability at above 300 MHz and can be densified at temperatures below 900°C. The Bi<sub>2</sub>O<sub>2</sub>-B<sub>2</sub>O<sub>2</sub>-ZnO-SiO<sub>2</sub> (BBSZ) glass addition effects on the densification and magnetic properties of Co2Y-NiCuZn ferrite composites with various Co<sub>2</sub>Y/NiCuZn ferrite ratios were investigated. The densification of Co<sub>2</sub>Y-NiCuZn ferrite composites was enhanced by the addition of glass at low sintering temperatures (<900oC) due to the liquid phase sintering. Co2Y-NiCuZn ferrite composites with 4 wt% BBSZ glass sintered at 900°C show a relative density above 90% a high-initialpermeability of 5-6, a quality factor of above 30 in the 200-300 MHz frequency and a resonance frequency above 1 GHz, which can be used in high frequency multilayer chin inductors

### **BJ02**

### Investigation of Fe<sub>2</sub>YZ (Y=Ni, Cu; Z=Sn, Ga): The Heusler compounds with tetragonal structure

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Heusler compounds, a remarkable class of intermetallic materials with 1:1:1 (often called Half-Heusler) or 2:1:1 composition, have attracted considerable interest since their discovery, mainly because of their versatile magnetic properties. Recently has been published a computational study on X2YZ full Heusler alloys. According to this work, the predicted enthalpies of formation evidence 27 phases to be thermochemically stable against the elements and the regular X2YZ type. A chemical-bonding study yields an inherent tendency for structural distortion in a majority of these alloys, so they predict the existence of the new tetragonal phase Fe<sub>2</sub>CuGa (P42/ncm; a = 5.072 A, c = 7.634 A;  $c/a \approx$ 1.51) with a saturation moment of  $\mu = 4.69 \ \mu\text{B}$  per formula unit. In this work is reported the attempt to synthesize experimentally this new phase. Fe\_2.xMn\_xNiSn and Fe\_2CuGa (x = 0, 1, 2) alloys were prepared by arc-melting and annealed at different temperatures in the interval 800-1300K. The Curie temperature (TC) was determined from the high temperature magnetization measurements carried out in a field of 0.05 T using a vibrating sample magnetometer (VSM). The indexing of all XRD patterns of annealed samples and microstructure characterization are in progress and will be present in conference.

Supported by the project REFREEPERMAG of the EU

### **BJ03**

Crystallite growth kinetics and microwave properties of Fe-Ti substituted (La,Sr)MnO<sub>3</sub> prepared by mechanical alloying Nastiti Elwindari, Hinu Pramuji and Azwar Manaf\*

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### **BJ04**

### Electromagnetic characteristics of Cu substituted Co<sub>2</sub>Z-type ferrites $Ba_3Co_{2,x}Cu, Fe_{24}O_{41}$

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Cu substituted Z-type hexaferrites (Ba<sub>3</sub>Co<sub>2-x</sub>Cu<sub>x</sub>Fe<sub>24</sub>O<sub>41</sub> with x=0-1.0, Z-hex) were prepared by two step calcinations [1]. The stoichiometric mixture for M-type ferrite was precalcined at 900°C, and then single phase Z-type ferrite powders were synthesized at 1300°C by replacing CoO with various amounts of CuO. The frequency dependence of initial permeability of all samples was measured from 0.01 to 3GHz. The permeability varied with Cu contents, a maximum permeability of  $\mu$ <sup>l=</sup> 9.8 at 10MHz was observed for x=0.8. The permeability behavior was dependent on the content of Z-type ferrite phase and microstructure [2]. For x=0.8 the content of Z-type ferrite phase was estimated to be 89%, and the average grain size was around 2µm. The density of (Cu<sub>1.2</sub>Co<sub>0.8</sub>)Z-hex showed almost complete densification. However, the resonance frequency decreased to 700MHz as the permeability of Cu substituted Z-type hexaferrites (Ba<sub>3</sub>Co<sub>2.x</sub>Cu<sub>x</sub>Fe<sub>24</sub>O<sub>41</sub>) increased. Cu substituted Co<sub>2</sub>Z hexagonal ferrites could be a candidate material for wireless device applications.

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### **BJ05**

#### Electrical and magnetic properties of nickel and magnesium cosubstituted lithium ferrites

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Lithium ferrites are widely used in electronics and telecommunication applications owing to their low dielectric loss, high resistivity coupled with good magnetic properties. There is a need to further improve their performance by incorporating different cationic substitutions for microwave device applications [1,2]. Most of the reports in the literature focused only on single ion substitutions for enhancement of one property at the cost of other. However, only a scanty research work has been reported so far on cationic co-substitutions in these ferrites. It is therefore interesting to study and find out the influence of co-substitution of nickel and magnesium in place of lithum and iron in lithium ferrites. For this purpose,  $Li_{0.5,x}Ni_xMg_xFe_{2.5x}O_4$  where x = 0.00 to 0.25 in steps of 0.05, are prepared by conventional ceramic method. Structural, magnetic, resistivity and dielectric measurements have been carried out on all the samples. Lattice parameter has been observed to increase with the concentration where as Curie temperature, grain size and initial permeability gradually decreased. High resistivity and low dielectric loss apart from reasonable magnetization and Curie temperature are obtained at concentrations. The results are explained on the basis of cation influences on the electromagnetic properties at different crystallographic sites.

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### BJ06

### Resistivity and complex permeability dependence on isochronal recovery in polycrystalline yttrium iron garnet ( $Y_3Fe_5O_{12}$ )

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We propose to establish whether or not crystal atomic defects have any influence on the value of its complex permeability and electrical resistivity after undergoing heat treatments. A technique involving defects creation and their subsequent elimination within equal time-length durations was used: this is a technique for establishing isochronal recovery behavior of any atomic defects-dependent properties. YIG was prepared via mechanical alloving with a mixture of Y<sub>2</sub>O<sub>2</sub> and Fe<sub>2</sub>O<sub>2</sub> involving a 24 hours milling time. They were then sintered at different temperatures between 900°C to 1350°C and subsequently heated for 2 hours at 1000°C and quenched in cooking oil. Then, annealing of the samples at 1000°C for 2 hours was carried out. The magnetic and electrical properties before and after quenching and also after annealing were studied. The permeability and loss factor showed an isochronal recovery behaviour in which a parameter's value was decreased after quenching and increased back even higher after annealing suggesting that magnetic permeability and energy loss in YIG could significantly show an atomic defect-dependent behaviour. The resistivity value of the samples was observed to have an isochronal recovery behavior which the value dropped after quenching and increased almost gradually back even higher after isochronal annealing series.

### CA01

### Spin and charge excitations in cuprates and iron pnictides revealed by simulated resonant inelastic x-ray scattering

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We first examine Cu L3-edge resonant inelastic x-ray scattering (RIXS) spectra for cuprates based on the exact diagonalization results of the t-t'-t"-J model. Both spin and charge excitations are included in the calculations on the same footing. The calculated results are compared with recent experimental data [1]. We next examine the orbital excitations coupled to the spin degree of freedom in the iron pnictides, based on the calculation in a five-band itinerant model [2]. The calculated Fe L3-edge RIXS spectra disclose the presence of spin-flip excitations involving several specific orbitals. Magnon excitations predominantly composed of a single orbital component can be seen in experiments, although its spectral weight is smaller than spin-flipped interorbital high-energy excitations. The detailed polarization and momentum dependence is also discussed.

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### CA02

Discovery of fermi surface near anti-node in pseudogap phase of the under-doped Bi-2212 Chung Koo Kim<sup>1</sup>, Jhinhwan Lee<sup>2</sup>, Kazuhiro Fuiita<sup>3</sup>, Hiroshi Eisaki<sup>4</sup>, Shinichi Uchida<sup>4</sup>, J. C. Seamus

Davis<sup>5</sup> and Jinho Lee<sup>74</sup> <sup>1</sup> BNL/Cornell, USA <sup>2</sup> KAIST, Korea <sup>3</sup> Cornell, USA <sup>4</sup> AIST, Japan <sup>5</sup> Tokyo U., Japan <sup>6</sup> Cornell/BNL, USA <sup>7</sup> BNL/SNU, Korea

Quasi-Particle Interference (QPI) measured by spectroscopic-imaging scanning tunneling microscope (SI-STM) became an extremely useful tool in the study of correlated electron systems (CES). Particle-hole (p-h) symmetric QPI observed in the superconducting cuprates revealed many interesting phenomena including the disappearance of the QPI signal around the reduced zone boundary [1]. Recently, the most dominating band of Sr,Ruo, was identified[2] and p-h asymmetric QPI in the parent compound of the ferro-prictide superconductor revealed a nematic like features [3]. On the other hand, identification of the exact Fermi surface (FS) topology in "pseudo-gap" (PG) phase of the cuprate superconductors is one of the central issues of the modern condensed matter physics. ARPES data indicates "Fermi arc" develops in PG phase while quantum oscillation (QO) results show evidence of the electron pockets. To resolve this PG mystery, we investigated deeply under-doped Bi<sub>2</sub>Sr<sub>2</sub>CaCu<sub>2</sub>O8+6 above and below Tc using SI-STM and found that there is a QPI showing continuous, p-h asymmetric dispersion through EF above Tc which originates from anti-nodal FS which is disconnected from the nodal FS. This observation implies existence of Fermi surface with very low quasiparticle(QP) weight Zk in anti-node in PG phase. We also try to explain nodal anti-nodal anti-nodal in terms of Zk.

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### CA03

### Feedback effect on high-energy magnetic fluctuations in the model high-temperature superconductor HgBa,CuO<sub>4+d</sub>

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It is widely accepted that magnetic excitations make important contribution to the Cooper-pairing mechanism in cuprate high-temperature superconductors. Contribution from excitations below ~60meV is manifested by a generic 'resonance' feature observed by neutrons, which signifies a feedback effect of pairing on the magnetic excitations. However, the insufficient spectral weight of the resonance is unable to explain the high superconducting temperature (Tc). Recent research has demonstrated that intense excitations above 100meV are available as a possible resource for pairing, but it remains largely unknown whether this resource is actually utilized. Here we present a systematic electronic Raman scattering study of the model single-layer cuprate superconductor HgBa<sub>2</sub>CuO<sub>4+d</sub>. In an overdoped sample, we observe a pronounced enhancement of a high-energy peak related to two-magnon excitations in insulating cuprates upon cooling below Tc. This is accompanied by the appearance of the superconducting gap and a pairing peak above the gap in the Raman spectrum, and it can be understood as a high-energy feedback effect which evinces a direct involvement of the high-energy excitations in the Cooper-pairing mechanism. All of these effects occur already above Tc in two underdoped samples, demonstrating a related feedback mechanism associated with the pseudogap.

[1] Y. Li et al., arXiv:1112.2725

### CB01

### Ferromagnetism of Au nanoparticle assemblies: Role of chemical and structural parameters in magnetic properties

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Ferromagnetism induced by the surface atoms and protective organic molecules in Au nanoparticles has been investigated for both fundamental and application interests of magnetism.[1] Possible explanations of the ferromagnetism could be attributed to 5d localized holes generated by the surface atoms and the high spin-orbit coupling of Au. Further understanding of this mechanism requires to study the magnetic properties of the particles related to their relative distance (change of the ligand length) and the nature of the ligand-particle binding. In this work, we explored the effect of these parameters for Au nanoparticles assemblies which were cast in the form of compact pellet made by cold-pressing technique. The magnetic behaviors in thiol and octylamine-capped Au nanoparticle assemblies were investigated by different spectrometers and magnetometers. The evidence of ferromagnetic behavior at 300K was seen clearly in ~ 6nm size dodecanethiolated Au nanoparticles: surprisingly, for the diameter much lager than 2nm, a typical ferromagnetic magnetization curve with the saturation value of 0.25emu/g observed, while superparamagnetism with a blocking temperature around 50K shown in octylamine-capped Au nanoparticles. The surface analyses by XPS, ESR and FTIR indicated that the magnetic properties could be tuned upon not only the particle size but also their surface environments.

[1] P. Crespo et al. Phys. Rev. Lett. 98, 087204 (2004)

#### **CB02**

### Air-stable Fe@Au nanoparticles synthesized by the microemulsion's method

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Magnetic particles covered by gold are very important in many biological applications. However, there are not simple methods to produce small (< 5-10nm) nanoparticles. One of the main reasons for that is the general use of iron oxides as magnetic cores, which have a large crystalline mismatch with gold. The use of Fe would be more appropriate, but its high tendency to oxidation has largely precluded it from being used as a core. Here, we will show that using a simple "one-pot" successive reaction method in microemulsions, can avoid such problems and is able to produce very stable core-shell Fe@Au nanoparticles. By this successive "one-pot" procedure, nanoparticles of ~ 6 nm with an Fe core of 1.5 nm can easily be obtained. These Fe@Au nanoparticles, with a saturation magnetization of 1.7 emu/g, are very stable even in air after magnetic separation from the solution, which shows the good covering of the Fe core by the Au shell. In this contribution we will report the key parameters, which have to be taken into account, to prepare such stable Fe@Au dispersions and analyze their optical and magnetic properties, as well as their possible anolications as biosensors, tareeted magnetic separation etc.

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### CB03

### Numerical study of the exchange bias effects in assemblies of core/shell nanoparticles

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The study of the magnetic behaviour of random assemblies of nanoparticles with core / shell structure is a relatively new field of experimental and theoretical interest. We study numerically using the Monte Carlo simulations technique with the implementation of the Metropolis algorithm the role of the magnetostatic and exchange interactions on the magnetic behaviour of composite nanoparticle assemblies. The aim of our work is to determine the factors that influence the exchange bias behaviour in the assemblies. We have developed a simplified model that describes the complex internal structure of the nanoparticles and also includes the interparticle interactions. We have calculated the hysteresis loops and the Zero Field Cooled magnetization curves of the assemblies for different sizes of the shell thickness for each nanoparticle and as a function of the particle concentration and the size of the dipolar and interparticle exchange strength in the assembly. Our simulations show that for small shell thickness increase in the exchange bias field (Hex) and the coercive field (Hc). As the shell thickness increase the effect of the exchange interparticle interactions is Our results are compared with experimental findings [1.2].

I. J. Nogues, V. Skumryev, J. Sort, S. Stoyanov, D.Givord, Phys. Rev. Lett. 97 (2006) 157203 ; 2. N.Domingo, A. M. Testa, D. Fiorani, C. Binns, S. Baker, J.Tejada J. M. M. M. 316 (2007) 155-158

### **CB04**

### Exploring the effect of Co doping in the magnetic and magneto-optical properties of fine magnemite nanoparticles

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Nanosized spinel ferrites are the subject of increasing interest due to their remarkable magnetic, catalytic, electric and optical properties, which, in most cases, differ from bulk. Cobalt ferrite is the ferrite with larger magnetic anisotropy that can be directly controlled by varying the cobalt content. However other factors, like chemical state or inversion degree, determine the final magnetic properties. We present a study of the structural, magnetic and magneto-optical properties of a series of Co substituted ferrite nanoparticles prepared by thermal decomposition. The structural characterization, carried out by using several techniques (TEM, XRD, XAS) showed all the samples are high crystalline, 5-6 nm spherical NPs with the cubic spinel structure typical of ferrites. The evolution of latice parameters with cobalt content suggests that the material is Co-substituted magnetine, also confirmed by XAS and Magneto optical spectroscopy. The investigation of the magnetic and magneto-optical properties displays peculiar trends with the cobalt content, the main features being the large increase of the saturation magnetization and the anomalous dependence of magnetic anisotropy which reaches its maximum values for intermediate compositions. The large tuneability of this material makes it possible to implement the performances of devices used in biomedical and sensing applications.

### CB05

### Novel technique for self assembly of magnetic nanoparticles by cluster beam deposition

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Mono-layer dispersion and self-assembly of nanoparticles are of a key importance in many fields of science and industry. In chemical synthesis [1] and recent high-energy ball milling studies [2] self assembled nanoparticles have been obtained by the aid of surfactants and wet processing. Another promising technique for nanoparticle production is cluster beam deposition which provides excellent material purity. However, it lacks the ability to form nanoparticles in the self-assembled array form. In recent studies, particles have been deposited onto poly(vinyl alcohol) [3] and phospholipid [4] coated substrates to form arrays. Nevertheless, these coatings either lack the continuous array formation or need time consuming coating steps and post deposition treatments. In this study, we have chosen Oleic Acid (OA) as a cost effective and easier alternative to the above mentioned coatings. SmCo<sub>3</sub> nanoparticles have been deposited for 10 sec. at 25 W sputtering power onto both bare and OA-coated single crystalline Si substrates for comparison. Mono dispersed particles with self-assembly formation have been observed on OA coated substrates while a non-uniform distribution has been observed on bare substrates. Such dispersion can reduce the exchange interactions between the nanoparticles which is expected to effect the coercivity of the nanoparticles.

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### CC01

### Magnetricity and magnetic monopoles in spin ice

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The analogy between spin configurations in spin ice materials like Ho<sub>2</sub>Ti<sub>2</sub>O<sub>7</sub> and proton configurations in water ice, H<sub>2</sub>O, has been appreciated for many years (see Ref. [1] for a review). However it is only in the last few years that this equivalence has been extended into the realm of electrodynamics [2]. In this talk I shall describe our recent experimental work that identifies emergent magnetic charges ('monopoles'), transient magnetic currents ('magnetricity') and the universal properties expected of an ideal magnetic Coulomb gas (magnetic electrolyte - 'magnetolyte'). These universal properties include the Onsager-Wien effect, 'corresponding states' behaviour, Debye-Huckel screening and Bjerrum pairing [4-6]. I will describe experimental results for both traditional spin ice materials (Ho<sub>2</sub>Ti<sub>2</sub>O<sub>7</sub>, Dy<sub>2</sub>Ti<sub>2</sub>O<sub>7</sub>) and a recently discovered system (Dy<sub>4</sub>Ge,O<sub>7</sub>).

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### CC02

#### Recent developments in quantum spin liquid candidates Luis Balicas\*

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We discuss a few geometrically frustrated magnetic systems exhibiting ground states characterized by correlated spin excitations. The first system,  $Pr_2Ir_2O_7$  is a metallic pyrochlore, which does not exhibit a spin-ice ground state but a state which is claimed to be a chiral spin-liquid where the interaction between the Pr moments and the tinerant Ir carriers leads to an anomalous Hall-effect (AHE), which remains finite even in the absence of an external magnetic field and of long range magnetic order 1. We show that this spontaneous AHE is hysteretic and anisotropic depending strongly on the previous application of an external field along specific crystallographic directions 2. We will argue that our observations suggest the possibility that Dirac monopole and antimonopole pairs are deconfined acquiring a finite quantum-mechanical average in the zero-field state of  $Pr_2Ir_2O_7$ . Finally, we briefly discuss and compare two new triangular lattice systems 4,5, Ba<sub>3</sub>CuSb<sub>2</sub>O<sub>9</sub> and Ba<sub>3</sub>NiSb<sub>2</sub>O<sub>9</sub> having S=1/2 and S=1 moments respectively, which despite their insulating character, exhibit large electronic specific theat coefficients of 43.4 and 168 mJ/molK<sup>2</sup>, respectively. This can be understood if the spin excitations behave as coherent fermionic (or charge-like) excitations which in these systems form a Fermi surface at low temperatures.

In collaboration with S. Nakatsuji (ISSP-Tokyo) and H. D. Zhou (NHMFL - Tallahassee) I. Y. Machida et al., Nature 463, 210 (2010). 2. L. Balicas et al., Phys. Rev. Lett. 106, 217204 (2011). 3. H. D. Zhou, E. S. Choi, G. Li, L. Balicas, C. R. Wiebe, Y. Qiu, J. R. D. Copley, and J. S. Gardner Phys. Rev. Lett. 106, 147204 (2011). 4. J. G. Cheng, G. Li, L. Balicas, J. S. Zhou, J. B. Goodenough, Cemke Xu, and H. D. Zhou, Phys. Rev. Lett. 107, 197204 (2011).

### CC03

### Static and dynamic properties of a strong-leg spin ladder

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Static and dynamic properties of the strong-leg S=1/2 Heisenberg spin ladder system (C<sub>7</sub>H<sub>10</sub>N)<sub>2</sub>CuBr<sub>4</sub> are studied using inelastic neutron scattering and bulk magneto-thermodynamic measurements. The leg-odd excitation channel is dominated by long-lived single-magnon states in the entire Brillouin zone, which supports a symmetric ladder model with exchange constants J\_leg = 1.42(6) meV and J\_rung = 0.82(2) meV [1,2]. In the leg-even channel, a considerable fraction of the spectral weight is contained in a novel long-lived two-magnon bound state [2]. In applied magnetic fields we observe a Bose-Einstein condensation of magnons that manifests itself in 3D long-range antiferromagnetic ordering emerging beyond H\_c = 2.8 T. The field-temperature phase diagram showing the spin liquid, Tomonaga-Luttinger spin liquid and BEC phase is mapped out by specific heat measurements. The experimental results are in spectcular agreement with state-of-the art DMRG calculations [2]. The latter provide additional insight on certain spin ladder properties specific to the strong-leg regime. This work was partially supported by the Swiss National Fund and MANEP.

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### CC04

### Nonstationary processes in the spin-ice materials $Dy_2Ti_2O_7$ and $Ho_7Ti_2O_7$ investigated by ultrasound

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We report on a low-temperature comparative study of two spin-ice materials,  $Dy_2Ti_2O_7$ , and  $Ho_2Ti_2O_7$ , by use of longitudinal and transverse ultrasound waves. In both materials the sound velocity and the sound attenuation exhibit a number of anomalies versus applied magnetic field for temperatures below the freezing temperature. The anomalies show a broad hysteresis with a very clear signature of the transition to the saturated state for magnetic fields applied along the [111] direction. Additionally, in  $Dy_2Ti_2O_7$  the sound velocity and the sound attenuation exhibit a number of well-defined spikes which are highly asymmetric due to the fundamentally distinct nonequilibrium mechanisms involved: the release of the Zeeman energy from the spins and the transfer of energy out of the sample [1]. In case of  $Ho_2Ti_2O_7$ , both acoustic characteristics demonstrate much smoother anomalies. We discuss some possible reasons for the observed differences in the ultrasound features of these spin-ice materials.

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### CC05

Study of low-temperature magnetism in a pr-based pyrochlore magnet Kenta Kimura<sup>1</sup>\*, Satoru Nakatsuji<sup>1</sup>, Agung Nugroho<sup>2</sup>, Yoshitomo Karakt<sup>3</sup>, Kazuyuki Matsuhira<sup>4</sup>, Yasuyuki Shimura<sup>1</sup> and Toshiro Sakakibara<sup>1</sup> <sup>1</sup> ISSP, Univ. of Tokyo, Japan <sup>2</sup> Bandung Inst. Tech. Indonesia <sup>3</sup> Ryukyu Univ., Japan <sup>4</sup> Kyusyu Inst. Tech., Japan

Pr-based pyrochlore magnets have been recognized as a fascinating venue to search unconventional magnetism as a result of frustration and quantum effects. For example,  $Pr_3Ir_2O_7$  exhibits a novel spin liquid state, which has the "2-in, 2-out" short-range correlation, the same as a spin ice system, and simultaneously breaks time reversal symmetry [1, 2]. Here we report the detailed physical properties of single crystals of the related insulating compound  $Pr_2Zr_2O_7$ . Magnetic susceptibility fits to the Curie-Weiss law below 10 K, yielding the effective moment of 2.5µB and Weiss temperature of -0.4 K, consistent with the previous powder results [3]. Furthermore, magnetization measurement at very low temperature confirmed a magnetic anomaly signaling spin ice correlation. In the presentation, we will discuss the low temperature magnetism based on the results including specific heat and AC susceptibility measurements.

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### **CD01**

### A materials-based global phase diagram for heavy-fermion quantum criticality

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Heavy fermion compounds have in recent years emerged as prototypical quantum critical systems. Studies in the anisotropic heavy-fermion compound YbRh\_Si<sub>2</sub> have shown that different types of quantum critical points (QCPs) are induced by variations of the magnetic field, chemical or external pressure, raising the question of the extent to which heavy-fermion quantum criticality is universal. We have identified a cubic heavy-fermion material, Ce<sub>3</sub>Pd<sub>29</sub>Si<sub>6</sub>, as exhibiting a field-induced quantum phase transition as the lower of two consecutive phase transitions is suppressed to zero. Thus, here the QCP separates two different ordered phases. This transition is accompanied by an abrupt change of Fermi surface, reminiscent of what happens across the field-induced antiferromagnetic to paramagnetic transition in YbRh\_2Si<sub>2</sub>. From these results we have proposed a materials-based global phase diagram that points to the importance of dimensionality - and may serve as guide in the search for a unified theoretical description [1]. To weaken the lower phase transition temperature and possibly detach the Kondo destruction scale from the phase transition scale, we are now studying (Ce\_La)<sub>2</sub>Pd<sub>29</sub>Si<sub>6</sub> and Ce<sub>2</sub>Pd<sub>29</sub>Si<sub>6</sub>, under hydrostatic pressure. We acknowledge financial support from the European Research Council (ERC Advanced Grant No 227378).

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### CD02

#### Anomalous metals with strong valence / orbital fluctuations

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In strongly correlated electron systems, quantum criticality normally emerges on the border of magnetism accompanied by critical spin fluctuations. In particular for the 4f based heavy fermions, it has been discussed using Doniach type picture based on the Kondo lattice with integer valence. Here we discuss unconventional type of quantum criticality and anomalous metallic phases found with strong valence and orbital fluctuations. The first material is the Yb based heavy fermion superconductor  $\beta$ -YbAlB<sub>4</sub> [1]. Interestingly, this exhibits a zero field quantum criticality without tuning as the first example in a metal [2]. Doping and pressure effects indicate no magnetic phase in the immediate vicinity of the ambient pressure quantum criticality. The instability associated with electronic structure or valence may be the origin of the novel quantum criticality. Secondly, we discuss the anomalous metallic phase found in the Pr based cubic compounds PrTr\_Al<sub>20</sub> (Tr: transition metal) [3]. These allow us to study the competition between the strong hybridization effects and quadrupolar ordering. Orbital fluctuations under strong hybridization with conduction electrons suppresses orbital ordering and lead to prominent non-Fermi liquid behavior, suggesting the possibility of a nonmagnetic type of Kondo effect.

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### CD03

### Ce-based iron-pnictides: Intermediate valence and heavy-fermion behavior versus magnetism and superconductivity

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In the last few years layered iron pnictides have attracted considerable interest due to their superconducting properties characterized by critical temperatures, TC, up to 55 K. In Ce-based compounds of these series superconductivity competes with heavy Fermion behavior, spin density wave formation and magnetic order. High-resolution angle-resolved photoemission allows for direct insight into the correlated electronic structure of these systems. We show that the dispersion of Fe-3d derived bands close to the Fermi energy is strongly modified by hybridization with localized Ce 4f states what affects possible participation of these bands in other many-body interactions. Spectral 4f signatures as probed by resonant photoemission consist on a weakly dispersive peak just at the Fermi energy (in addition to an ionization peak at about 2 eV binding energy) and admixtures to the valence bands. The spectroscopic data are discussed in the light of band-structure calculations and a novel approach to the periodic Anderson model.

### CD04

### Observation of the quantum critical point in ${\rm CeRhSi}_3$ with the muon spin rotation technique

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The interplay of superconductivity and magnetism has been a central issue of research for many years. A variety of these systems have been found to display a so-called quantum critical point which marks the transition point between a nonmagnetic and a magnetic metal phase. The external tuning parameter can be doping, an external magnetic field or external pressure. We conducted a muon spin rotation study on the heavy fermion and noncentrosymmetric compound CeRhSi, which is an antiferromagnet at ambient pressure and displays superconductivity at high pressure. In this system, we were able to unravel the quantum critical point hidden inside of the pressure-induced superconducting phase of the pressure-temperature phase diagram. We directly observed the continuous suppression of the internal fields, and thus the magnetic moments at the critical pressure \* = 23.6 kbar. Simultaneously the Neel temperature goes to zero and we conjecture a spin density wave type quantum critical point.

### CE01

### Self-modulation in perpendicular anisotropy Co/Ni based spin-torque oscillators

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Nano-contact spin-torque oscillators (NC-STO) are nanoscale microwave signal generators based on the transfer of spin angular momentum from a spin polarized current to the local magnetization of a thin magnetic layer. While standard, NiFe based, NC-STOs require large external fields of about 1 T for operation, we here present low-field operation of Co/Ni based NC-STOs comprising a 50-300 nm diameter nano-contact on top of a Co8nm(IP)/Cu8nm/Co0.3nm-[Ni0.8nm/Co0.4nm]×4(PMA) orthogonal spin valve. Frequencies above 12 GHz are observed for close to zero-field (0.02 T) operation, which further increase up to 24 GHz in moderate fields of 0.4T [1]. When the applied is further increased, we find a discontinuity in STO frequency around 0.68 T. The STO frequency drops by almost 10 GHz and at the same time sidebands appear with the main STO signal. These sidebands indicate a self-modulation of the main high-frequency mode with a low-frequency (~1 to 2 GHz) oscillatory mode. We interpret this transition as the formation of a magnetic bubble underneath the NC. The low-frequency magnetodynamics of the bubble then modulates the NC-STO precession frequency.

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### **CE02**

### Magnetization tilt angles in [Pd/Co]/Cu/[Co/Pd]-NiFe pseudo spin valves

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Future spintronic devices based on spin-transfer torque require a delicate engineering of the constituent magnetic layers. A fixed layer magnetization that is tilted out of the plane allows for zero-field operation, as well as increased output efficiency and quality, e.g. linewidth. We recently proposed a [Co/Pd]-NiFe exchange spring for use as a spin polarizing layer with a tilted magnetization for applications in spin torque oscillators (STOS) [1]. We have shown that it is possible to easily tune the spin configuration by engineering the effective exchange coupling between a NiFe layer with in-plane anisotropy and [Co/Pd] multilayer with perpendicular anisotropy. Here, we report on magnetization tilt angles in complete [Pd/Co]/Cu/[Co/Pd]-NiFe(t<sub>NiFe</sub>) pseudo spin valves (PSVs) using angular dependent transport measurements. The tilt angle of the [Co/Pd]-NiFe is increased, the tilt angle of the [Co/Pd]-NiFe increases. These results provide a first observation of the tilt angle in a PSVs structure, which are similarly used in real STOs. This work provides meaningful insights for the future realization and optimization of tilted STO devices.

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### **CE03**

Temperature dependence of microwave voltage emission associated to spin-transfer induced vortex oscillations in magnetic tunnel junctions

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The growing market of telecommunication applications demands reduction of energy consumptions and devices miniaturization. Spin transfer-torque nanoscillators (STO) can fulfill such necessities and, given their rapidly improving characteristics, will might be able to replace conventional oscillators. In the past, their drawbacks were the low emitted power and large linewidth. We recently proved[1,2] that these issues can be overcome by considering a MgO-based magnetic tunnel junction[3,4] with a vortex free layer, i.e., spin transfer-torque vortex nano-oscillator (STVO). The obtained signal has large output power (>50 nW) and small linewidth (<1 MHz). In order to further optimize its characteristics, we have to understand the mechanisms underlying the linewidth generation. Thus, we consider the effect of temperature on such STVO by ranging T from 20 to 300 K. We first observe a linear decrease of the linewidth with decreasing temperature and a bottom limit of the linewidth (700 kHz) at lower temperature[5]. Thermal effects and spin torque non-linearities are not sufficient to explain such behaviour and one additional source of phase noise, frequency and time-domain measurements are compared. Support from ANR VOICE PNANO-09\_P231-36, EU MASTER NMP-FP7-212257, CANON-ANELVA for MTJ film are acknowledged.

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### **CE04**

### Zero external-field microwave oscillations in MgO magnetic tunnel junctions

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We report a potential solution to overcome major challenges for the implementation of spin transfer nano-oscillators (STNOs) in real microwave applications, namely a weak output power and broad spectral linewidth of STNO microwave signals. In addition to the challenges, an external magnetic field is often required for the microwave generation from STNOs. We have observed high-power microwave oscillations having a narrow linewidth, even without applying external magnetic fields, from STNOs based on MgO magnetic tunnel junctions (MTJs). The MTJs consist of SiO<sub>2</sub> substrate/ buffer layer/ synthetic antiferromagnet/  $Co_{60}Fe_{20}B_{20}$  (2)/ capping layer (in nm). A typical STNO has a resistance of 45 Ohm in the parallel state and 75 Ohm in the antiparallel state, showing a tunnel magnetoresistance of 67%. The STNO exhibits an oscillation peak of 70 nV/Hz<sup>34</sup> at 3 GHz at zero external magnetic field, and the full width at half maximum of the peak is less than 60 MHz. The microwave spectra as a function of bias current and in-plane external magnetic field (H<sub>ext</sub>) will be discussed in detail.

### **CE05**

### NCMR based spin-torque microwave generator and detector with high signal purity

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Nano-contacts magnetoresistive (NCMR) devices have many ferromagnetic ohmic nano-contacts (NCs) with the diameter of around 1nm, which connect the free and reference layer in the spacer layer called nano-oxide layer (NOL). In each NC, the domain wall is geometrically confined and produces the magnetoresistance by spin accumulation [1] and miss-tracking effect [2]. Furthermore, in injecting current, these NCs can increase the density of current and provide the high emission power with narrow linewidth. In this study, we systematically investigated the spin-transfer induced microwave oscillation [3], unique frequency modulation characteristics [4] and spin-torque ferromagnetic resonance (ST-FMR) [5] for the NCMR devices with thick free-layer with the thickness of ~15nm. In the microwave measurement, the high emission power (up to 100nW) with the narrow linewidth (down to 6.5MHz) was excited at the frequency of around 1GHz caused by auto-oscillation at multi-domain state, which is the intermediate state between parallel and antiparallel state. In the same magnetic state, ST-FMR signal also showed high sensitivity of ~250µV/100µW at almost same frequency as c current induced microwave oscillation. As a result, we confirmed that NCMR device has the possibility to become the nano-sized microwave generator and detector.

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### **CE06**

### Conditions for zero field spin transfer induced vortex oscillations with a perpendicular spin polarizer

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Despite the theoretical and experimental progress achieved, power, linewidth and zerofield emission remain critical issues for spin transfer nano-oscillators (STNOs) [1,2] in telecommunication applications. Here we investigate spin transfer induced vortex dynamics in specific STNOs [SAF / MgO / NiFe (vortex magnetization) / Cu / [Co/Ni] (perpendicular polarizer)]. Such a hybrid GMR/TMR system has been designed to obtain large amplitude spin torque induced magnetization oscillations at zero field owing to the perpendicular polarized spin current. In the best MTJ we measure large emitted power (600 nW) and narrow linewidth (590 kHz) at zero magnetic field, making these hybrid devices very promising for radiofrequency applications. In addition we investigate in detail the strong impact of vortex configuration (polarity, chirality) on the frequency and the linewidth of the signal [3]. These results are analysed in the frame of a nonlinear analytical model for the vortex dynamics [4]. We thank CANON ANELVA for preparing the MTJ films, the ANR VOICE grant and MASTER grant and E.G. acknowledge financial support from CNES and DGA.

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### CF01

### Spin-orbital short-range order on a honeycomb based lattice

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Strong antiferromagnetic interactions indicated by  $\Theta CW$ =-47 K, fail to produce static spin correlations above T=0.1 K in Ba<sub>2</sub>CuSb<sub>2</sub>O<sub>2</sub>. Magnetic neutron scattering at T=1.5 K is inelastic with a broad peak at 5.5 meV that shifts to higher energies upon warming and wave vector dependence indicating near neighbor singlet formation. While these features are difficult to reconcile with the previously inferred triangular lattice structure, short-range correlations of CuSbO<sub>9</sub> dumbbells resulting from structural frustration as revealed by synchrotron x-ray diffraction and EXAFS provides the essential lead. The resulting magnetic sub-lattice forms a nano-structured honeycomb lattice with nearest neighbor interactions. Based on the new structure we propose the unusual magnetic properties and the absence of a coherent, static Jahn-Teller transition indicate a twodimensional random singlet phase or a spin-orbital quantum liquid.

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### **CF02**

#### Quadrupolar waves in uranium dioxide Paolo Santini

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In solids with d- and f-electrons, under certain conditions the orbitals align to form an ordered pattern. Collective excitations breaking this arrangement can take the form of oscillations of electric quadrupoles, so-called quadrupolar waves. These represent a propagating pattern of charge densities implying a modulation of quadrupolar moments. We show that quadrupolar waves constitute a major component of the dynamics of uranium dioxide in its magneto-quadrupolar ordered phase [1,2]. The distinct roles of Jahn-Teller and superexchange mechanisms as sources of quadrupolar interactions are identified. The theory is fully consistent with earlier and recent [3] inelastic neutron scattering data.

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### **CF03**

### Magnon gap formation and charge density wave effect on thermoelectric properties in SmNiC<sub>2</sub> compound

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We studied the magnetic, electrical, and thermal properties of polycrystalline compound of SmNiC<sub>2</sub>. The electrical resistivity and magnetization measurement show the interplay between the charge density wave at T(CDW) = 157 K and the ferromagnetic ordering of Tc = 18 K. Below the ferromagnetic transition temperature, we observed the magnon gap formation of 4.3 ~ 4.4 meV by  $\rho$ (T) and Cp(T) measurements. The charge density wave is attributed to the increase of Seebeck coefficient resulting in the increase of power factor S<sup>26</sup>. The thermal conductivity anomalously increases with increasing temperature as the whole measured temperature range which implies the weak attribution of Umklapp phonon scattering. The thermoelectric figure-of-merit ZT significantly increases due to the increase of power factor at T(CDW) = 157 K. Here we argue that the competing interaction between electron-phonon and electron-magnon couplings exhibits the unconventional behavior of electrical and thermal properties.

### CF04

### Effect of R ion size variance on spin and orbital order in RVO<sub>3</sub> (R=rare earth and Y)

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Perovskite RVO<sub>3</sub> has orbital degrees of freedom between dyz and dzx orbitals in V<sup>3+</sup> ions, and shows two types of spin/orbital order (SO/OO): C-type SO (C-SO)/G-type OO (G-OO), and G-SO/C-OO. In this system, the transition temperature of each SO/OO depends on the R-site ionic radius1 and structural randomness caused by the size mismatch of R ions2. We have investigated the single crystals of  $Y_{1,a}(La_{0,195}Lu_{0,80})_xVO_3$ ,  $Eu_{1,a}(La_{0,222}Y_{0,678})_xVO_3$ , Sm<sub>1</sub>.  $_{3}(La_{0,522}Y_{0,678})_xVO_3$ , and Nd<sub>1,a</sub>(La\_{0,41}Y\_{0,490})\_xVO\_3 in order to clarify the effect of R ion size variance on SO/OO in RVO<sub>3</sub>. In the present systems, the average ionic radius of R-site is constant, while the size variance of R ion is proportional to x. In R=Y<sub>1,a</sub>(La\_{0,195}Lu\_{0,80})\_xVO\_3, Eu<sub>1</sub>.  $_{3}(La_{0,522}Y_{0,678})_x$ , and Sm<sub>1,a</sub>(La\_{0,522}Y\_{0,679})\_x systems, the R ion size variance suppresses C-SO/G-OO, while it stabilizes the other SO/OC. By the neutron and resonant X-ray scattering measurements, other order is confirmed as G-SO/C-OO. In Nd<sub>1,a</sub>(La\_{0,491}Y\_{0,509})\_xVO\_3 with larger R-site ions, however, C-SO/G-OO is so stable that the R ion size variance cannot induce the other SO/OO.

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### CG01

#### Spin Hall effects in n-GaAs near the metal-insulator transition Paul A Crowell

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We have measured the direct and inverse spin Hall effects in GaAs and In<sub>x</sub>Ga<sub>1-x</sub>As in samples doped near the metal-insulator transition. The direct spin Hall effect is measured through the spin accumulation induced at the edges of the semiconductor channel in the presence of a charge current. The inverse spin Hall effect is measured as a Hall voltage induced by injection of a spin current. We find that the magnitude and temperature dependence of the direct spin Hall effect is consistent with simple models for spin-orbit scattering from ionized impurities. The inverse spin Hall effect shows both an anomalously large magnitude and an unusual magnetic field dependence, suggesting that hyperfine interactions play an important role. I will review this work with an emphasis on our attempts to understand spin transport in GaAs near the metal-insulator transition.

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#### CG02

### The effect of an inhomogeneous interface on the transport properties across Fe/GaAs(001) films

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Efficient electrical spin-injection is the greatest challenges for semiconductor spintronics. Fe/GaAs remains one of the leading candidate systems for achieving this feat due to the small lattice mismatch. We report the first observation of different interface structures in Fe/GaAs(001) films. Both the ideal and mixed interfaces were observed leading to a distribution of Schottky barrier properties [1]. Regions with mixing have a lower barrier height creating pin-hole type areas where transport will be focused, leading to heating. This increases the contribution to the transport from thermionic emission, therefore reducing the spin injection efficiency. Using ab-initio calculations we show that resonance states form in the minority spin DOS for the mixed interface causing a polarization inversion. We show it is possible to create well behaved reproducible devices if an abrupt interface is created. Transport properties of 3-terminal devices which have a predominantly abrupt interface give spin lifetimes of over 30 ns at 5 K with a spin diffusion length of 16 µm at 10 K. Full details of the interface structures and transport properties of Fe/GaAs(001) films will be presented in the full paper.

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### CG03

#### Spin accumulation and decoherence mechanisms at ferromagnetic/ tunnel barrier/semiconductor interfaces

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The ability to inject spins into a semiconductor (SC) by electrical means and the possibility to convert a spin accumulation into an electrical signal is expected to bring new functionalities. However, it has been shown that the injection of spin-polarized carriers from a ferromagnetic transition metal (FM) into a SC requires the addition and the fine tuning of a spin-conserving resistance at the FM/SC interface. Following the work on Co/Al<sub>2</sub>O<sub>2</sub>/ GaAs interface [1], we now carry out a systematic study on the enhancement of spin accumulation involving modification of the interfacial properties by probing several semiconductor materials involving GaAs, Ge and InP as well as MgO and Al<sub>2</sub>O<sub>2</sub> barriers. Inverted Hanle effect measurements reveals the influence of a fluctuating magnetic field leading to decoherence of the spin accumulation at the FM/SC interface whose origin resides in stray magnetic field due interface roughness [2-3], or in hyperfine interaction in case of transport through localized states of GaAs. Both 3-terminal Hanle and Inverted Hanle measurements highlight effects of spin accumulation at the FM/SC interface, few orders of magnitude larger than the expected. The amplitude of the effects can be understood by a sequential injection mechanism through localized interfacial states at the barrier/SC interface.

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### CG04

### Electrical spin accumulation and detection in $Fe_3O_4/MgO/GaAs$ systems

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The Hanle effect is one of the effective and direct methods for studying the spin injection from a ferromagnet and it is useful for understanding the decay of these injected spins in the semiconductor [1,2]. We have chosen Fe<sub>3</sub>O<sub>4</sub>/MgO/GaAs (n-doped) system for our study, where the spin polarized electrons are injected from Fe<sub>3</sub>O<sub>4</sub> (10m) to GaAs through thin MgO layer (2nm). This is the first ever attempt in studying the spin injection and decay of spins in the above mentioned system. Fe<sub>3</sub>O<sub>4</sub> is a room temperature ferrimagnet which is reported to have 100% spin polarization. The 3-terminal non local Hanle measurement [3,4] is usually carried out in these systems. From our studies, the I-V measurements on the structured junctions of size varying from 200x150µm<sup>2</sup> to 200x25µm<sup>2</sup> exhibit tunnel behavior with no considerable rectification. The 3-terminal non local voltage  $\Delta V$  as a function of perpendicular magnetic field for the biasing current I≥1µA. Also, the Hanle plots show the dependency on the extraction and injection of the bias electrons to the above mentioned devices accordingly.

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### CG05

### Spin relaxation in defect-free InGaN/GaN quantum wells

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Due to ferromagnetic (FM) properties it presents at room temperature in its n-doped state (with rare-earth metals), and combination of wide-gap and low spin-orbit coupling (SOC), GaN has been identified as a promising compound for spintronics applications. In this work [1], we present numerical analysis of spin relaxation time in defect-free two-dimensional In doped GaN quantum wells grown on GaN nanowires [2]. Data, of wurtzite crystalline structure, has been collected in a range of temperatures, from 100 K to room temperature and at different doping concentrations and energy levels using photoluminescence techniques. In the analysis, spin relaxation mechanisms considered are Elliot-Yafet (EY) and D'yakanov-Perel'(DP). Relative importance and relevance of EY vs. DP mechanism [3] has been investigated. It was found that EY mechanism is very ineffective while DP mechanism dominates as the source of spin relaxation. We have also calculated Rashba and Dresselhaus coefficients for each doping concentrations and shown that increase in doping also changes Rashba and Dresselhaus coefficients in accordance with accepted values for IN and GaN.

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### CH01

#### Magnetic, magnetotransport and magnetocaloric properties of quaternary Ni-Mn-In-Z Heusler alloys

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We present recent results on structural, magnetic, magnetotransport, and magnetocaloric properties of  $N_{isgl}Mn_{34,8}In_{15,2}$ ,  $Z_x$  (Z=In, Ge, Al, B, Si) and  $N_{isgl}Mn_{34,8}In_{15,2}$ , B<sub>4</sub> Heusler alloys in the austenitic and martensitic state focusing on behavior of magnetocaloric effect, ordinary and anomalous Hall effect (OHE and AHE) in the close vicinity of phase transitions. The adiabatic changes of temperature  $\Delta$ Tad at the phase transitions have been found by direct measurements using an adiabatic magnetocalorimeter as well as by indirect method using heat capacitance measurements and thermomagnetic curves M(H,T). Negative and positive values of ?Tad of about -2K and +2K have been observed in the vicinity of the martensitic transition (MT) and Curie temperatures, respectively, for applied magnetic field  $\Delta$ H=1.8T. To find OHE and AHE coefficients we used asymptotic values of the Hall resistance in weak and strong magnetic fields as well as thermomagnetic curves M(H,T) and magnetorsistance. It was shown that OHE and AHE coefficient is weak. It clearly indicates that both Berry phase contribution and side-jump mechanism are not responsible for AHE in inhomogeneous alloys.

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### **CH02**

### Composition dependence of magnetic properties in tetragonal Heuslerlike Mn-Ga alloy films with large perpendicular magnetic anisotropy

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Magnetic films with perpendicular magnetic anisotropy (PMA) are advantageous to spin torque applications using magnetic tunnel junctions, such as magnetic random access memory, because PMA reduces switching current and increase thermal stability. It has been predicted that a tetragonal Heusler-like alloy Mn<sub>3</sub>Ga exhibits large spin polarization as well as large uniaxial magnetic anisotropy [1]. We have obtained Mn<sub>3</sub>Ga (x=0.5) epitaxial films and reported a large perpendicular magnetic anisotropy energy Ku ~ 12 Merg/cc and also low saturation magnetization Ms ~ 250 emu/cc, so far [2]. Furthermore, Gilbert damping for the alloys were relatively smaller than those in other PMA films [3]. We report here composition dependence of magnetic properties for Mn<sub>3</sub>Ga alloy films. The Mn<sub>3</sub>Ga films were fabricated using UHV magnetron sputtering on Cr buffered MgO substrates. X-ray diffraction showed all the films were of e-axis oriented epitaxial structures. Magnetization curves for the films show perpendicular magnetizations with squareness close to unity. The Ms for the films decreased linearly with increasing Mn content, whereas Ku show no remarkable decreases [4]. The results will be discussed with the ab-initio calculations [5]. This work was partially supported by the Strategic International Cooperative Program ASPIMATT (IST) and WPI program (MEXT) in Japan.

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### CH03

### Optical spectroscopy of half-metallic and thermoelectric Heusler compounds

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Heusler compounds are versatile material class allowing to design desired properties through various constituent atoms, i.e. properties intrinsically related to the modification of the atomic electronic structure. Among the studied classes are half-metallic and thermoelectric materials. Although Heusler compounds are throrughly studied by number of techniques, optical spectroscopies are often omitted. However, they can provide information of phonon energies, gap size and its purity, shape of critical points, presence of quasi-particles such as exciton, conductivity estimation from IR absorption, etc. Knowledge of complex refraction index is important for mass production control check of thin films, to check e.g. their thickness and roughness. Also, half-metallic Heusler compounds are promising materials for magneto-optical applications, due to low photon absorption for energies below half-metalic gap. Within this contribution we present complex refractivity index, determined in wide spectral range from far-infrared to near-ultraviolet, of half-metallic Heusler compounds Co<sub>2</sub>FeSi, Co<sub>2</sub>FeAl<sub>10</sub>, Si<sub>10</sub>, and Co<sub>2</sub>FeGa<sub>10</sub>, Ge<sub>10</sub>, and thermoelectric on CoTiSb and NiTiSn (later also doped with Sc and V) measured by ellipsometry and infrared reflectometry. The features of optical spectra are related with details of the electronic structure. The experimental optical properties are compared with ab-initio calculations.

### CH04

### Verification of band structure calculations for the Heusler compound Co<sub>2</sub>MnGa

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Among ferromagnetic Heusler materials many compounds are predicted to be half metals, which makes them highly attractive for spintronics applications. However, direct experimental evidence for the validity of band structure calculations for this class of materials is scarce. Here, the electronic density of states is probed by spin averaged and spin resolved ultraviolet photoemission spectroscopy for the Heusler compound Co<sub>4</sub>MnGa using a novel spin filter [Ko110]. The spectroscopy is performed in-situ on epitaxial thin films which show almost complete absence of oxidation after the measurement. This means a further reduction of sample degradation compared to our recent experiments [Hah10]. The experiments show several characteristic features such as a clear Fermi egde and two pronounced emission intensity maxima in the spin averaged data. In spin resolved photoemission a Heusler record polarization of 55% at the Fermi egde is followed by a sign change at E-E\_F=-0.7eV. Calculations of the band structure accounting for the many-body effects and the related photoemission spectrum were performed using the LSDA+DMFT approach in the framework of a fully relativistic spin-polarized Korringa-Kohn-Rostoker (SPRKKR) formalism. Good agreement observed between the experimental data and the calculations of

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#### **CH05**

### Field-driven domain-wall ratchet shift register

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Magnetic domain wall shift registers are considered for application in memory [1] and logic circuits [2]. However, to shift all domain walls in a nanowire in the same direction, one requires either impractically high current densities [1] or rotating magnetic fields [2]. Other issues include strict geometrical restrictions, and randomness in the domain wall displacement. We propose a new domain wall shift register suitable for magnetic materials with perpendicular anisotropy (PMA). Using focused ion beam irradiation [3], we engineer an energy landscape for the domain walls that is asymmetric along the propagation direction. This favors ratchet propagation of domain walls in one direction through a nanotrack. Whereas it is often believed that domain walls shift register. In a proof of principle experiment, indefinite propagation of two domain walls along a closed loop is shown.

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### CI01

### Solitonic lattice and Yukawa forces in the rare earth orthoferrite TbFeO<sub>3</sub>

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The control of domains in ferroic devices lies at the heart of their potential for technological applications. Multiferroic materials offer another level of complexity as domains can be either or both of a ferroelectric and magnetic nature. Here we report the discovery of a novel magnetic state in the orthoferrite TbFeO<sub>3</sub> using neutron diffraction under an applied magnetic field. This state has a very long incommensurate period of 340 A at 3 K and exhibits an anomalously large number of higher-order harmonics, allowing us to identify it with the periodic array of sharp domain walls separated by many lattice constants. These domain walls are formed by Ising-like Tb spins. They interact by exchanging spin waves propagating through the Fe magnetic sublattice. The resulting force between the domain walls has a large but finite range that determines the period of the incommensurate state. It is analogous to the pion-mediated Yukawa interaction between protons and neutrons in nuclei.

### CI02

#### Ferroelectricity from magnetic helicity in ferroaxial crystals Laurent C Chapon Institut Laue-Langevin, France

Spin-driven ferroelectricity in most non-collinear magnets, such as TbMnO<sub>3</sub>, is induced by the so-called inverse Dzyalonshinskii-Moriya mechanism and requires a cycloidal magnetic structure, an ordered magnetic state that is not truly chiral (or

lacks helicity). Conversely, in a truly chiral magnetic state (proper helix), the pseudoscalar helicity can not couple directly to the electric polarization, and therefore can't induce ferroelectric order. However, in systems of specific crystal symmetry, named here "ferroaxials," the presence of collective structural rotations mediates an indirect coupling between magnetic helicity and ferroelectricity. I will review our recent experimental results for new compounds of this class, obtained by magnetic X-ray and neutron diffraction techniques, including a clear demonstration that the magnetic helicity can be controlled by an electric field in RbFe(MoO<sub>4</sub>)<sub>2</sub>, and the existence of 'giant' improper ferroelectricity in CaMn<sub>7</sub>O<sub>12</sub>.

### CI03

### Chemical-doping control of magnetoelectric multiferroics

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In magnetoelectric multiferroics, electric polarizations can establish without ionic displacements via spin-orbit coupling that involves spin moments with noncollinear spatial arrangements. In this spin-current multiferroics the direction of electric polarization is determined by the spatial arrangement of interacting magnetic moments, for which inversion symmetry should be absent. In our most recent works, we have demonstrated the ability to control the magnetoelectricity in two different multiferroic materials via chemical doing, namely Mn1-xCoxWO4 and Ba05Sr15Zn2(Fe1-xAlx)12O22. It is found that the substitution of Co<sup>2+</sup> ions on Mn<sup>2+</sup> sites not only stabilizes the incommensurate spiral phase, but also changes the orientation of its easy plane. The direction of the ferroelectric polarization subsequently changes from the b to the a axis. In Ba<sub>0.5</sub>Sr<sub>1.5</sub>Zn<sub>20.5</sub>Sr<sub>1.5</sub>Zn<sub>2</sub>(Fe<sub>1.5</sub>Al<sub>5.12</sub>O<sub>22</sub>, on the other hand, the substitution of nonmagnetic Al3+ ions on Fe3+ sites induces the transition from planar proper screw to conical spin structures by additionally introducing commensurate axial components. This change in magnetic structure is ascribed to the reduction in planar magnetic anisotropy, which eventually leads to the reduction of low critical field for the fieldinduced magnetoelectric phase. These results indicate that it is possible to manipulate the magnetoelectricity of certain multiferroics materials by selective and precise control of chemical doping.

### CI04

### Magnetic x-ray scattering studies on multiferroic SmFe<sub>3</sub>(BO<sub>3</sub>)<sub>4</sub>

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In recent years rare-earth ferroborates, RFe<sub>3</sub>(BO<sub>3</sub>)<sub>4</sub>, have been identified as potential multiferroic materials from both, application as well as fundamental physics point of view. Due to the presence of two magnetic sublattices these materials offer a platform for understanding several interesting physics problems related to structural and magnetic anisotropy. We here present resonant and non-resonant x-ray scattering measurements on single crystalline SmFe<sub>3</sub>(BO<sub>3</sub>)<sub>4</sub>. This compound crystallizes in trigonal structure (s.g. R32) and preserves this structure at all temperatures. Below  $T_N$ , magnetic reflections are found at (h k l)  $\pm$  (0 0 3/2). Resonant x-ray measurements are performed using photon energies close to the Sm L-edges and Fe K edge. Non-resonant measurements are performed using photon energies above and below the edges as well as at 100 keV photon energy. The temperature dependence of the magnetic peak at the Sm L<sub>2</sub> edge reveals the polarization of the Sm moments by the Fe moments in form of an accelerated ordering of the Sm moments upon decreasing temperature Our measurements confirm the easy-plane magnetic anisotropy for the two magnetic sublattices of Sm and Fe. Non-resonant measurements performed at 100 keV photon energy, with and without magnetic field indicate spin rotations in the basal plane.

### CJ01

### Spin excitations and transformation of domain structure in nanocrystalline CoFeB-SiO<sub>2</sub> films with growth induced anisotropy Alexander Grishin\*

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We survey results on strong in-plane magnetic anisotropy of electrical transport and broad frequency band spin dynamics in (CoFeB)<sub>1x</sub>-(SiO<sub>2</sub>), films. Samples with less metallic content show higher anisotropy of electric conductivity and isotropic giant magnetoresistance (GMR), while films with smaller x have less anisotropic resistivity and show a mixture of GMR and anisotropic magnetoresistance (AMR). Oppositely, field of magnetic anisotropy Hp increases with "metallization" reaching 535 Oe for the film with x = 0.235. "Soft" magnonic modes were observed in a whole range of oblique orientations of magnetic field perpendicular to the "easy" magnetic axis. FMR frequency nullifies at the line of "spinodal decomposition". At descending branch of magnetization curve we found the sequence of domain structure transformations. Uniformly magnetized film breaks to stripe domains separated by low-angle Neel domain walls (DWs). Then, spontaneous nucleation of kinks of Bloch-type DWs occurs. DWs have different chirality and are separated by segments of Neel DW. When  $H \rightarrow 0$ , they shrink converting to the vertical Bloch lines (VBL). After field reversal, Bloch DW loaded with VBLs becomes energetically unfavorable and transforms instantaneously into Neel DW which gradually disappears with film's saturation.

### CJ02

### GMI effect of amorphous microwires with enhanced magnetic softness

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Amorphous microwires exhibit outstanding magnetic properties such as magnetic bistabilty, fast domain wall propagation, magnetically soft properties and GMI effect. We studied hysteresis loops and GMI effect (GMI ratio,  $\Delta Z/Z$ , impedance tensor components and pulsed GMI effect) at frequencies, f, till 4 GHz in amorphous microwires with different ratios, p, of metallic nucleus diameter, d, and total diameter, D. This allowed us to control the magnetoelastic anisotropy, since the strength of internal stresses is determined by the  $\rho$ -ratio. GMI effect and hysteresis loops exhibited strong sensitivity to the p-ratio. Hysteresis loops of Co-rich low-magnetostrictive microwires exhibit low coercivity (<10A/m). Magnetic anisotropy field increases when ρ decreases. Field dependence of the off-diagonal voltage response measured in pulsed scheme exhibits anti-symmetrical shape with monotonic growth within the certain field range. Annealing (under stress and/or magnetic field) significantly affects magnetic anisotropy and GMI effect of microwires. After stress annealing GMI effect of Fe-rich microwires (which did not exhibit detectable GMI effect in as-prepared state  $\Lambda Z/Z \le 1\%$ at f=10 MHz) can be considerable improved ( $\Delta Z/Z=60\%$  at f=10 MHz) as a result of stress-annealing induced magnetic anisotropy. Low field GMI hysteresis, related with helical magnetic anisotropy, can be suppressed by the circular bias magnetic field.

### CJ03

### Micro-fabricated silicon spiral spring based electromagnetic energy harvester

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In this study, a micro-fabricated electromagnetic energy harvester was newly designed, fabricated, and characterized to generate electrical energy from ultra-low ambient vibrations under 0.3g. The proposed energy harvester was comprised of a highly miniaturized Neodymium Iron Boron (NdFeB) magnet, silicon spiral spring, multi-turned copper coil, and polydimethylsiloxane (PDMS) housing. When an external vibration directly moves the magnet mounted as a seismic mass at the center of the spiral spring, the mechanical energy of the moving mass is transformed into electrical energy through the 183 turns of the solenoid copper coils. The silicon spiral spring was used to generate a high electrical output power by maximizing the deflection of the movable mass in response to low level vibrations. The fabricated energy harvester exhibited a self resonant frequency of 36Hz and an optimal load resistance of 99Ω, respectively. It generated an output power of 29.02 $\mu$ W and load voltage of 107.3m V at a vibration acceleration of 0.3g. It also exhibited a power density and normalized power density of 48.37 $\mu$ W/cm<sup>3</sup> and 537.41 $\mu$ W/cm<sup>3</sup> g<sup>2</sup>, respectively. The total volume of the fabricated energy harvester was low no 0.6cm (height).

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### **CJ04**

### Magneto-optical study of magnetization reversal in sub-micrometric glass covered wires

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With the purpose of the miniaturization of basic elements of magnetic sensors the magneto-optical Kerr effect (MOKE) [1] investigation of the magnetization reversal has been performed in Fe-rich sub-micrometric amorphous wires. The axial tensile stress and torsion stress have been applied during the experiments. The series of the microwires with different values of radius of metallic nucleus (400 nm, 700 nm, 1000 nm) has been studied. In the extremely thin sub-micrometric wires, the magnetic bistable behavior is observed. It confirms the existence of the Surface Large Barkhausen Jump in submicrometric glass covered wires like in glass covered wires of micro-scale. The highest value of the surface coercive field is observed for the smallest value of the radius. It has been attributed to the increasing of the strength of internal stresses as increasing the glass coating thickness. The performed analysis of the tensile and torsion stresses transformation of surface hysteresis loop demonstrates that about one order decrease of the wire scale does not abolish the basic effects observed earlier in thicker wires. It permits to reduce considerably the size of basic elements of magnetic sensors.

[1] J. Gonzalez et al., Journ. of Physics: A, 208 (2011) 502.

### CJ05

### The magnetic transition and large magnetoresistance effect in perovskite Nd<sub>1-x</sub>Sr<sub>x</sub>MnO<sub>3</sub> system

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The magnetic transition of  $Nd_{1,x}Sr_{x}MnO_{3}$  (with x = 1/3, 0.40 and 0.50) has been investigated with zero field cooling (ZFC) and field cooling (FC) measurements in the temperature range of 77 K - 350 K. It is indicated that the temperature dependence of magnetization of all samples complied with Block's law in the area of  $\hat{T} < TC$ . The magnetoresistance of Nd<sub>1-x</sub>Sr<sub>x</sub>MnO<sub>3</sub> system has been determined with magnetic field of 0.0-0.4T. It is found that the obtained magnetoresistance ratio (CMR%) decreased with increasing doped Sr concentrations. The maximum value of CMR(%) is out of the Curie temperature (TC) and metal-insulator transition temperature (Tp ) regions. The magnetoresistance ratio nearly linearly increased by decreasing temperature at x = 0.50compound. The dependence of CMR(%) on the doping Sr in this system has been fitted with hyperbol function.

### **DA01**

### Superconducting symmetry of Fe-based systems studied by impurity effects and neutron inelastic measurements

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We have focused on the superconducting symmetry of Fe-pnictides, because it is directly connected with the pairing mechanism. Effects of nonmagnetic-impurities on Tc, magnetic excitation spectra  $\chi$ " and NMR 1/T1-T curve, all of which are sensitive to the relative signs between the order parameters on their disconnected Fermi surfaces around  $\Gamma$ and M points, have been studied: If the signs are opposite (symmetry S±) the magnetic mechanism is relevant, while the same signs (symmetry S++) indicate a novel mechanism, because the ordinary phonon mechanism cannot realize Tc as high as ~55 K in Ln1111 (Ln=lanthanides). Results are as follows. (a)Verv small rates of Tc-suppression by M impurities in LnFe1-yM,AS0.89xF011+x (M=Ni, Co, Ru) can be explained only by S++[1, 2]. (b) x2 data for Ba(Fe, Co)2As2 (Tc-23 K) and Ca-Fe-Pt-As crystals (Tc~33 K) can be well explained by S++ rather than by S±[3]. (c)The NMR data can be understood by S++, too. They strongly support S++ and the existence of a novel pairing mechanism. The symmetry is consistent with the observed elastic softening of C66[4]expected from the orbital-fluctuation (OF) mechanism. Anomalous behavior of phonons by neutron studies are also presented in relation to the OF.

[1] M. Sato et al. J. Phys. Soc. Jpn. 79 (2010) 014710. [2]T. Kawamata et al. J. Phys. Soc. Jpn. 80 (2011) 084720. [3] J. Phys. Soc. Jpn. 80 (2011) 093709. [4] T. Goto et al. J. Phys. Soc. Jpn. 80 (2011) 073702.

### **DA02**

### Carrier doping versus impurity effects in transition metal-substituted iron-based superconductors revealed by ARPES

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In the electron-doned Fe pnictides, superconductivity is realized by substitution of transition-metal atoms for Fe. However, it has been controversial whether itinerant carriers are indeed doped into the system and how the impurity potential affects the electronic properties [1]. In order to address these issues, we have performed a systematic ARPES study of Ba(Fe,M),As,, where M = Mn to Zn. For M = Co, Ni, and Cu the substitution expands the electron Fermi surfaces and shrinks the hole Fermi surfaces, as expected. However, the number of doped electrons estimated from the Fermi surface volumes is generally smaller than the nominal value, indicating that part of the doped electrons do not become itinerant. Such a tendency becomes stronger in going from Co, Ni to Cu. Concomitantly, energy bands are distorted and M 3d-like states are split off below the Fe 3d bands. As for Mn and Zn, on the other hand, ARPES shows characteristic features of the antiferromagnetic state without any signature of carrier doping. This indicates that the half-filled and fully occupied 3d shells are stabilized in the divalent state of Mn and Zn respectively.

[1] H. Wadati et al., Phys. Rev. Lett. 105, 157004 (2010).

#### **DA03**

Specific heat measurements on fepn in fields up to Hc2 - a probe of nodal structure Greg Stewart\*

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We report the specific heat gamma as a function of field up to 30 T in Co- and Ni-doped 122, as well as in LiFeAs. The overdoped Co-doped sample displays qualitatively different behavior than the underdoped. Preliminary gamma as a function of angle data will also be discussed.

### **DA04**

#### NMR study on high temperature Fe-pnictide superconductor Ln-Fe-As-O with Tc=50 K

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Superconducting(SC) transition temperature (Tc) of Ln-FeAsO (1111)-based compounds (Ln: Rare earth) rises to 55 K when the Ln=Sm. In NMR investigation, however, the high-Tc Ln-1111 compounds are affected by the magnetism of rare earth ions, which prevents us from probing the intrinsic nuclear relaxation rate(1/T1) of the quasi-particles in the SC state. In order to reveal the mechanism of high Tc more than 50 K in Ln-1111 system, we have investigated the normal-state and SC characteristics of (La<sub>0.05</sub>Y<sub>0.95</sub>)FeAsOy (La<sub>0.05</sub>Y<sub>0.95</sub>1111) with Tc=50 K by means of 57Fe/75As-NMR measurements. In the SC state, the measurements of 1/T1 have revealed in terms of a multiple fully-gapped s(+-)-wave model that the SC gap becomes larger than that in optimally-doped LaFeAsOy (La111(OPT)) with Tc=28 K. In the normal state, the increase of 1/T1T upon cooling indicates that the antiferromagnetic(AFM) spin fluctuations develop toward Tc, which is more significant than in La1111(OPT) and (La<sub>0.8</sub>Y<sub>0.2</sub>1111) (Tc=34K)[3]. However, these antiferromagnetic spin fluctuations were weaker than in Ba<sub>0.6</sub>K<sub>0.4</sub>Fe<sub>2</sub>As<sub>2</sub> (Tc=38K) and (Ca<sub>4</sub>Al<sub>2</sub>O<sub>6</sub>)(Fe<sub>2</sub>As<sub>2</sub>) (Tc=27K), which have lower Tc values than in La<sub>0.05</sub>Y<sub>0.95</sub>1111, suggesting that the AFM spin fluctuations are not an unique factor to increase the Tc in Fe-pnictide superconductors.

### **DB01**

### Ultrahigh-resolution and time-resolved laser photoemission study on kondo materials

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Recently, photoemission spectroscopy has been developed very rapidly and found to be very powerful for the study of Kondo materials. For example, 1, The resolution of 150 ueV at 1.8K is achieved by the development of laser photoemission. 2. High resolution hard X-ray photoemission becomes possible using high brilliant synchrotron radiation. 3. The angle resolved photoemission in soft X-ray region also shows beautiful results using high brilliant synchrotron radiation 4 Time resolved photoemission becomes possible very recently using Femto-second laser. I would like to talk on these new experimental results on YbAlB<sub>4</sub>[1] and YbS and Yb metal [2] as well as several Ce, U, and Pr compounds using ultra-high resolution laser-photoemission and high resolution hard X-ray photoemission. I will also talk on the possibilities of time resolved photoemission spectroscopy on heavy electron systems.

[1] Okawa et. al., Phys Rev. Lett., 104(2010) 247201. [2] Matsunami et. al., PRB 78(2008)195118.

### **DB02**

### Magnetic moment screening in the correlated Kondo lattice model

Peter Thalmeier<sup>1</sup>, Mohammad Siahatgar<sup>1</sup>, Burkhard Schmidt<sup>1</sup> and Gertrud Zwicknagl<sup>2</sup> <sup>1</sup> Max Planck Institute for Chemical Physics of Solids. Germany <sup>2</sup> Technical University Braunschweig, Germany

The magnetic correlations, local moments and susceptibility in the correlated 2D Kondo lattice model at half filling are investigated. We calculate their systematic dependence on the Kondo coupling J\_K and Coulomb repulsion U. Exact diagonalisation (ED) approach for ground state properties as well as finite temperature Lanczos method (FTLM) for specific heat and the uniform susceptibility are employed for small tiles on the square lattice. The competition of on-site screening and induced inter-site correlations are monitored and a phase diagram is constructed. In particular it is shown that the screened local moment exhibits nonmonotonic behavior as function of U for weak Kondo coupling J K. In the large U limit the model is equivalent to the 2D Kondo necklace model with two types of localized spins. In this limit the numerical results are compared to those of the analytical bond operator method in mean field treatment and excellent agreement for the total paramagnetic moment is found. The temperature dependence of the susceptibility allows to extract the dependence of the Kondo temperature scale on the correlation strength U. A monotonic increase for small U is found in agreement with earlier analytical impurity calculations.

### **DB03**

### Influence of magnetic anisotropy on the underscreened Kondo effect in the presence of ferromagnetism

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The prominent role of magnetic anisotropy (MA) in formation of the Kondo effect has recently been demonstrated experimentally [12]. In particular, it turned out that in systems of spin S>1/2, such as magnetic adatoms (i.e. Fe. Co or Mn) or magnetic molecules, the Kondo effect can be tuned by modifying the system's MA. Furthermore, theoretical studies also indicate that MA can be a key factor determinin spin-polarized transport through a magnetic nanosystem [3]. Motivated by the recent experiment [2] in this communication we address how the MA affects the underscreened Kondo effect (i.e. partial compensation of the molecular spin by conduction electrons) in the case of an artificial molecule of spin S=1 coupled to a reservoir of spin-polarized conduction electrons [4]. The crucial ingredient of the model is the presence of uniaxial MA. The problem is analyzed by means of Wilson's numerical renormalization group (NRG) method, which allows for calculating the spectral function of the molecule. We show that the interplay of MA and ferromagnetism has a fundamental significance for occurring of the Kondo effect. Most importantly, despite the presence of the effective exchange field [5] the Kondo effect can be restored by adjusting the magnitude of MA.

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### **DR04**

### Quantum criticality out of equilibrium in the pseudogap Kondo model

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### **DB05**

A spin-selective kondo-insulator - cooperation between ferromagnetism and kondo-effect

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The Kondolattice model has been intensively studied for the last thirty years as a fundamental model for heavy fermions, Kondo insulators, and transition metals, e.g. the manganites. Yet, the understanding is far from complete. Taking ferromagnetic heavy fermion compounds as motivation we analyze the mechanism stabilizing the ferromagnetic state in the antiferromagnetically coupled Kondo lattice model. We find that even for this ferromagnetic state Kondo screening plays an essential role in stabilizing the ferromagnetic state at zero temperature leading to very interesting properties: while the majority-spin electrons are metallic, the minority-spin electrons form an insulating state. We clarify that this state is due to partial Kondo screening, so that parts of the local moments are bound to the electrons, resulting in a dynamicallyinduced commensurability which is essential for producing the gap in the minority-spin electrons. We believe that the mechanism proposed here, the dynamically generated commensurability, is generic for the ferromagnetic phase in the antiferromagnetically coupled Kondo lattice model, thus providing new insights into the zero temperature physics for the Kondo lattice model

### **DC01**

#### Spin-orbit entangled ground states and excitations in iridium oxides Ginivat Khaliullin

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A relativistic spin-orbit coupling may drive unusual interactions and orderings in Mott insulators. This coupling entangles the spin and orbital subspaces leading to a rich variety of effective Hamiltonians and exotic phases depending on the lattice geometry and orbital structure. Particular examples to be discussed are: (i)Iridium perovskite Sr<sub>2</sub>IrO<sub>4</sub> where the magnetic order and excitations closely resemble those of high-T<sub>c</sub> cuprates [1,2]; (ii) The honeycomb lattice Li<sub>2</sub>IrO<sub>2</sub> and Na<sub>2</sub>IrO<sub>3</sub> where the celebrated Kitaev model with spin liquid ground state might be at work [1 3]: (iii)Vanadate Sr<sub>2</sub>VO, which is predicted to exhibit a magnetically hidden octupolar order [4].

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### **DC02**

Elementary magnetic excitations of iridates and cuprates probed by resonant inelastic X-ray scattering

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Resonant Inelastic X-ray Scattering (RIXS) provides direct access to elementary charge, spin and orbital excitations in complex oxides. As a technique it has made tremendous progress with the advent high-brilliance synchrotron X-ray sources. From the theoretical perspective the fundamental question is to precisely which low-energy correlation functions RIXS is sensitive. Depending on the experimental RIXS setup the measured charge dynamics can include charge-transfer, phonon, d-d and orbital excitations [1]. The focus of this talk will be on RIXS as a probe of spin dynamics, in particular magnon and bi-magnon dispersions [2,3]. Based on the experimental observations, the novelties that RIXS reveals on the spin dynamics of high-Tc cuprates [4,5] and strongly spin-orbit coupled iridium-oxides will be discussed [6].

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### **DC03**

Resonant Inelastic X-ray Scattering study of Na<sub>2</sub>IrO<sub>3</sub>

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Recently there are intensive research efforts to elucidate physics of materials with 5d electrons, for which spin orbit coupling (SOC) is comparable to other energy scale of the system, such as Coulomb repulsion and crystal field splitting [1-3]. The layered iridate Na<sub>2</sub>IrO<sub>3</sub>, with Ir<sup>4+</sup> spins sitting on a distorted honeycomb lattice, is such an example [4-7]. Depending on the value of the Coulomb repulsion in such a lattice, the SOC can give rise to a topological insulator or a Kitaev-Heisenberg system for the possible spin-liquid phase [5]. The electronic structure of Na<sub>2</sub>IrO<sub>2</sub> has been investigated using resonant inelastic x-ray scattering (RIXS) and density functional theory calculations. Crystal field split d-d excitations are resolved in the high-resolution RIXS spectra. We observe a large splitting between the two eg orbitals, which arises from the dimerization of Ir atoms due to the metal-metal bonding. The low energy d-d excitations exhibit small momentum dependence, which can be described as due to spinorbital coupled excitation.

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### DC04

### Magnetization plateaus in generalized Shastry-Sutherland models

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The Shastry-Sutherland model is known to exhibit a novel sequence of magnetization plateaus in the presence of an applied field. It has been shown in the recent past that inclusion of anisotropic and/or additional interactions to the standard Sahstry-Sutherland model can change the magnetization profile completely - resulting in different sequences of magnetization plateaus. In the present work, we use large scale numerical simulations to systematically study the effects of Ising-like exchange anisotropy and a 3rd nearest neighbor ferromagnetic interaction on the field induced magnetization behavior. In particular we follow the evolution of the m/m s=1/3 and 1/2 plateaus as the strength of the additional terms are varied. The results are valuable in understanding the behavior of the rare-earth tetraboride family of quantum magnets whose magnetic properties are described by generalized Shastry-Sutherland models of the kind considered in this study.

### **DD01**

#### Modulated spinodal decomposition and magnetotransport in (Ge,Mn) films grown on GaAs(001)

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The field of ferromagnetic semiconductors evolves very fast nowadays for their potential use in spintronic devices. Up to now efforts have mainly focused on Diluted Magnetic Semiconductors but Curie temperatures in these materials still remain modest. One possible route to increase at least locally transition temperatures is to use spinodal decomposition leading to transition metal-rich high-TC nanostructures. We focus here on (Ge.Mn) considered as a model system for spinodal decomposition and compatible with Si-based microelectronics. While the growth of (Ge,Mn) films on Ge substrates leads systematically to Mn-rich self-assembled nanocolumns exhibiting high-TC, we demonstrate the fine control of spinodal decomposition in (Ge,Mn) films grown on GaAs. Using different surface preparations, we clearly identify the role of surface morphology and impurity diffusion from the substrate (Ga or As) on the nanocolumns growth and the electrical properties (MR and AHE)[1]. In particular holes exhibit an anomalous Hall effect, and electrons exhibit a tunneling magnetoresistance, both with a clear dependence on the magnetization of the Mn-rich inclusions; holes exhibit orbital MR, and electrons show only the normal Hall effect, and an additional component of magnetoresistance due to weak localization, all three being independent of the magnetic state of the Mn-rich inclusions.

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### **DD02**

#### Homogenous and heterogeneous magnetism in (Zn.Co)O

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For more than a decade ZnO doped with Mn and Co has remained as one of the most prospected DMS for the room temperature (intrinsic) ferromagnetism resulting in numerous reports on room temperature ferromagnetism conflicted with those which denv existence of any coupling at all. In order to clarify this issue we investigate ZnCoO layers grown by ALD at low temperatures [1]. We employ and relay on wide range of extensive material characterisation which in combination with SQUID magnetometry allow us decisively exemplify the growth temperature Tg as the key factor discriminating between paramagnetic (obtained at Tg=160oC) and various forms of ferromagnetic responses (seen when Tg>200oC). When the ferromagnetism is found, our data indicate presence of nearly temperature independent and highly anisotropic response which we unambiguously associate with few nm thin metallic Co-mesh located at the (Zn.Co)O/substrate interface. It explains why the magnitude of the ferromagnetic-like signal is virtually independent of the film thickness as well as elucidates the origin of magnetic anisotropy, as observed by us and others [2]. Furthermore, it makes it possible to understand significant deviations from the standard superparamagnetic behavior visible in our samples as well as in many high-temperature DMSs and DMOs [3].

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#### **DD03**

#### Magnetic and optical studies of hydrogenated Cu-doped ZnO film Tong Li, Wen Xiao, Tun Seng Herng, Nina Bao and Jun Ding\*

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ZnO doping with Cu has provoked broad interest recently 1,2 Some interesting findings, such as surface ferromagnetism (FM) in hydrogenated ZnO3 and hydrogen enhanced green emission (GE) in ZnO film4 were reported previously. In this work, we focused on hydrogenated ZnO:Cu film This study reported hydrogen enhanced FM and strong GE in ZnO:Cu (2 at%) film and possible mechanism. ZnO:Cu films (002) with different thickness were deposited on quartz (110) by PLD at 600 oC. Prior to H2 treatment, the areal saturation magnetization Msa (1×10-5 emu/cm<sup>2</sup>) was found to be insensitive to thickness, suggesting surface/interface magnetism5. After H2 treatment at up to 500 oC. Ms of ZnO:Cu (20 nm) reached  $3.1 \times 10-5$  emu/cm<sup>2</sup> (15.4 emu/cm<sup>3</sup>) which was three times stronger than hydrogenated ZnO<sub>2</sub>. Surprisingly, unlike hydrogenated ZnO, this hydrogen induced Msa drastically decreased with increasing thickness. It thus indicated that apart from surface FM resulting from OH attachment4, Cu dopants play a dominant role in FM enhancement in response to H2 treatment. Furthermore, although light emission of ZnO:Cu showed an obvious blue shift compared to ZnO a similar strong GE was observed after 500 oC-H2 treatment. This GE is attributed to defect complexes and formation of porous structure that creates optical cavities.

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### **DD04**

#### Formation and investigation of structural and magnetic properties of Ni-Mn-In Heusler alloy thin films

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In this work we report on the different methods of formation of the Ni-Mn-In Heusler alloy thin films by pulsed laser deposition and the dramatic dependence of martensitic transition temperatures (Tm) and Curie temperatures (Tc) on the alloy composition Using two-lasers co-deposition approach we have implemented various methods of the thin films formation with controlled stoichiometry: by ablation of one target (Ni<sub>50</sub>Mn<sub>35</sub>In<sub>15</sub>), by simultaneous ablation of two targets (Ni<sub>50</sub>Mn<sub>35</sub>In<sub>15</sub> and In) or by simultaneous ablation in the appropriate ratio of three pure metals targets (Ni, Mn and In). The phase and chemical composition, structural and magnetic properties of asgrown Ni-Mn-In thin films have been probed by X-ray Diffraction (XRD), Raman scattering, Auger electron spectroscopy, Rutherford backscattering spectrometry and vibrating sample magnetometery. Temperature dependences of magnetic and structural properties are investigated by Physical properties measurement system (PPMS) and high-resolution XRD. It was shown that the films are not susceptible to degradation. We have observed that increasing of In concentration from 14 to 20% leads to Tc decreasing by 30K (from 330 to 300K). At the same time temperature hysteresis, corresponding to martensitic transitions changes from temperature range Tm=250-310K to Tm= 260-370K. The low-temperature phase transition has been found for all samples.

### DD05

#### Magnetic excitations in rare earth based nanosystems

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### **DD06**

### Exchange coupled L10 FePt (hard)/ soft (A1 FePt or Co) nanocomposites

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L10 ordered FePt is one of the most promising materials for ultra high density perpendicular magnetic recording, however its high coercivity greatly exceeds the writing field of available heads. Thus new media structures proposed in order to decrease the required writing field named as exchange coupled composites (ECC) or so-called exchange spring media. The ECC media consists of two or more exchange coupled sub-regions, in which the hard region stores the information and the soft region(s) promotes the magnetization reversal of the hard region, while keeping high thermal stability. In this work structural and magnetic properties of core-shell type hard/soft (L10 FePt / A1 FePt or Co) exchange coupled nanocomposites are presented. Semi spherical core-shell nanocomposites with L10 FePt core and A1 FePt (fcc) or Co shell were obtained by depositing A1-FePt or Co cap layers on type L10 FePt nanoparticles in order to understand the influence of the soft magnetic layer thickness on the magnetic properties of the system. Epitaxial growth is confirmed by x-ray diffraction and TEM, while the coercivity decreases dramatically for the L10/A1-FePt or Co system when the thickness of the A1-FePt or Co cap layers is increased

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### **DE01**

### Micromagnetic simulation of magnetic nanostructures

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Magnets are key elements of modern society. Examples of the wide spread use of magnets are magnetic data storage, permanent magnets for energy applications, wireless sensors, and biomedical systems. The design and development of magnetic devices heavily relies on simulations. Recent advances in hard- and software make it possible to bridge the length scales. In magnetic recording simulations the input and output are the write current and the read back voltage while magnetization dynamics is treated at the nano-scale. Nano-scale non-coherent magnetization reversal decouples the switching field and the thermal stability of magnetic recording media. In the talk we will show the potential of bit patterned magnetic recording systems for storage densities reaching 5 Tbit/in2. In addition to the physics of magnetization reversal, we will review the methods and algorithms that are applied in numerical micromagnetics. Further examples of micromagnetic simulations will include low-rare earth containing permanent magnets and wireless magnetostrictive sensors.

### DE02

### Multi-bit magnetic memory based on the extraordinary Hall effect.

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We propose a principle of multi-bit magnetic random access memory in which each memory cell is split among several multilevel dots exhibiting perpendicular anisotropy, and the extraordinary Hall effect (EHE) is used to extract the stored information. Foureight- and sixteen- memory state cells have been realized in the proof-of-concept study. In addition to multiplicity of states, probably the most important advantage of the split cell architecture is freedom in positioning dots of the same cell at separate locations and different levels, thus building an effective three-dimensional memory.

### DE03

### Epitaxial Fe/MgO/Fe tunnelling junctions on BaTiO<sub>3</sub> (001)

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Novel schemes for magnetic recording and reading are based on spintronic devices where a determinant role is played by materials or interfaces displaying magnetoelectric coupling. In this context the study of the Fe/BaTiO<sub>2</sub> (BTO) interface is particularly interesting in view of a possible magnetoelectric coupling at the interface as suggested by Duan et al [1] and Sahoo et al [2]. After an initial study of the magnetoelectric coupling at the interface between thin Fe films and BTO single crystal substrates [3] epitaxial Fe/BTO//Nb/SrTiO<sub>2</sub>(001) and Fe/  $BTO/La_{a}Sr_{a}MnO_{a}//SrTiO_{a}(001)$  interfaces have been grown by combined use of molecular beam epitaxy and pulsed laser deposition. Finally, fully epitaxial Co/Fe/MgO/Fe/BTO heterostructures have been deposited and magnetic tunnelling junctions (MTJs) fabricated via optical lithography. I(V) curves clearly indicate that tunnelling is the dominant mechanism in our MgO junctions. Preliminary experiments testing the electric control of the TMR have been performed at different temperatures. A modulation of the TMR on the order of 10%. induced by application of an electric field across the BTO template, has been detected at 150 K. This result attests the great potential of this system for the electric control of magnetization of ferromagnetic electrodes in spintronic devices

[1] C. Duan et al. Phys. Rev. Lett. 97 (2006) 047201 [2] S. Sahoo et al. Phys. Rev. B 76 (2007) 092108 [3] S. Brivio et al. Appl. Phys. Lett. 98 (2011) 092505

### **DE04**

### Experimentally performed periodic NOT/AND/OR magnetic quantum dots cellular automata gate

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A magnetic quantum dots cellular automata (MQCA) [1, 2] can perform NAND/ NOR logic operation [3] via magnetostatic interaction between the dots. The MOCA is composed of elliptical dots, and digital information of "0" and "1" are stored with a direction of dot magnetization. In general, the logic operations are executed by applying a uniform external magnetic field (clock field). However, in arrayed MQCA structures it is difficult to control a data flow direction between the MOCA gates with the uniform clock field. Here we demonstrate a NOT/AND/OR logic operation in arrayed structure with a localized clock field from magnetic force microscopy (MFM) tip. As the NOT/AND/OR MQCA gate, four elliptical Ni-20at.%Fe dots with long/short axis of 110 nm/65 nm thickness of 20 nm were fabricated on a thermally oxidized Si(100) substrate with electron-beam lithography, ion beam sputtering, and lift-off technique. To write digital information to the dots, we use a magnetization manipulation method with MFM. With all possible initial states of the NOT/AND/OR MOCA gate we confirm that the gate can perform expected logic operation at room temperature. With arrayed MLG structures, high functional circuits, such as adder circuit will be realized in the near future.

[1] R. P. Cowburn and M. E. Welland 2000 Science 287 1466 [2] A. Imre, G. Csaba, L. Ji, A. Orlov, G. H. Bernstein and W. Porod 2006 Science 311 205 [3] H. Nomura and R. Nakatani 2011 Appl. Phys. Express 4 013004

### **DE05**

### Energy-efficient control of vortex-core polarizations by tailored orthogonal pulse currents in cross-point architecture

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<sup>2</sup> Center for X-ray Optics, Lawrence Berkeley National Laboratory, Berkeley CA 94720, USA

Magnetic vortices have been considered to be a potential candidate for information-storage applications due to their core binary states, as well as energetically stable configurations. Recent findings [1] of vortex-core switching with low-power consumption have stimulated further studies towards implementation of the magnetic vortices into actual norvotable random access memory [2]. In this presentation, we are going to report a low-power-consumption and reliable manipulation of vortex-core magnetizations in patterned magnetic disks. To achieve an alternative to switch the vortex-core magnetizations, we applied rotating fields that are generated by two orthogonal Gaussian-pulse currents. The threshold field strength required for vortex-core switching can be significantly reduced by optimizing the pulse widths and the time delay between the orthogonal pulses [3]. We also demonstrated a remarkable reduction in the threshold strength using a coherent train of rotating fields with their optimized time interval. For reliable memory-bit selection and information recording, we fabricated  $2 \times 2$  ovtrex-state disks with crossed electrodes and demonstrated selective vortex-core switching at the intersection of the two crossed electrodes [4]. These results imply that the optimized pulse-type rotating fields in the basic cross-pointarchitecture are energy-efficient and reliable mans of selective information recording based on magnetic ordine based on magnetic areas of selective information recording based on magnetic ordine based on magnetic vortex-core such as the size cross-pointarchitecture are energy-efficient and reliable means of selective information recording based on magnetic vortex core

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### **DF01**

### Skyrmion dynamics in metallic chiral ferromagnet

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<sup>3</sup> Applied Physics, The University of Tokyo, Japan

<sup>4</sup> Zernike Institute for Advanced Materials, University of Groningen, Netherlands

Chiral ferromagnetic materials have recently been shown to be important platforms for the generation of novel topological spin texture called Skyrmions. Their easy generation by moderate magnetic field at relatively high temperature is beginning to open up the potential for device and memory applications. In this talk, I will review some history of its recent rapid development and present basic theory of thermodynamics and dynamics of Skyrmions. To conclude, a possible means of generating Skyrmions by purely electrical means, without relying on magnetic field, will be discussed.

[1] Dynamics of Skyrmion Crystals in Metallic Thin Films, Jiadong Zang, Maxim Mostovoy, Jung Hoon Han, and Naoto Nagaosa, Phys. Rev. Lett 107, 136804 (2011) [2] Skyrmion Lattice in Two-Dimensional Chiral Magnet, Jung Hoon Han, Jiadong Zang, Zhihua Yang, Jin-Hong Park and Naoto Nagaosa, Phys. Rev. B 82, 094429 (2010) [3] Real-space observation of a two-dimensional skyrmion crystal, X. Z. Yu, Y. Onose, N. Kanazawa, J. H. Park, J. H. Han, Y. Matsui, N. Nagaosa and Y. Tokura, Nature 465, 901 (2010)

### DF02

### Long-range crystalline nature of the skyrmion lattice in MnSi

Tim Adams<sup>1</sup>, Sebastian Muhlbauer<sup>2</sup>, Christian Pfleiderer<sup>3</sup>, Florian Jonietz<sup>3</sup>, Andreas Bauer<sup>3</sup>, Andreas Neubauer<sup>3</sup>, Robert Georgii<sup>3</sup>, Peter Boni<sup>3</sup>, Uwe Keiderling<sup>4</sup>, Karin Everschor<sup>5</sup>, Markus Garst<sup>5</sup> and Achim Rosch<sup>5</sup> <sup>1</sup>TU Munchen, E21, Germany <sup>2</sup> Institut fur Festkorperphysik, ETH Zurich, Zurich, Switzerland, Germany <sup>3</sup> Physik-Department E21, Technische Universitat Munchen, D-85748 Garching, Germany <sup>4</sup> Helmholtz Zentrum Berlin, BENSC, D-14109 Berlin, Germany

<sup>5</sup> Institute of Theoretical Physics, Universitat zu Koln, D-50937 Koln, Germany

We report small angle neutron scattering of the Skyrmion lattice in MnSi using an experimental setup that minimizes the effects of demagnetizing fields and double scattering. Under these conditions, the Skyrmion lattice displays resolution-limited Gaussian rocking peaks that correspond to a magnetic correlation length in excess of several hundred micrometers. This is consistent with exceptionally well-defined long-range order. We further establish the existence of higher-order scattering, discriminating parasitic double scattering with Renninger scans. The field and temperature dependence of the higher-order scattering arises from an interference effect. It is characteristic for the long-range orystalline nature of the Skyrmion lattice as shown by simple mean-field calculations. A peculiar change of the pinning and a meandering around the direction parallel to the magnetic field suggests a symmetry breaking transition and a more complex interplay of the skyrmion lattice with the crystal structure.

T. Adams et al., Phys. Rev. Lett. 107, 217206 (2011) S. Muhlbauer et al., Science 323, 915 (2009) A. Neubauer et al., Phys. Rev. Lett. 102, 186602 (2009)

### **DF03**

#### Magnetic textures and electron transport in chiral helimagnets Jun-ichiro Kishine<sup>1</sup> Alexander Ovchinnikov<sup>2</sup> Joor Proskurin<sup>2</sup> Yoshihiko Togawa<sup>3</sup> Yusuke Kousaka<sup>4</sup> and

Jun Akimitsu<sup>4</sup> <sup>1</sup> Graduate School of Arts and Sciences, The Open University of Japan, Japan <sup>2</sup> Department of Physics, Ural Federal University, Russia <sup>3</sup> N2RC, Osaka Prefecture University, Japan <sup>4</sup> Department of Physics, Aoyama Gakuin University, Japan

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A coupling of free electrons with non-trivial spin textures has recently attracted a great attention because of an ability to manipulate magneto-transport properties through a control of the background spin subsystem. In this presentation, based on our recent two papers[1,2], I report magnetic texture formation and anomalous mangetotransport phenomena in magnetic crystal belonging to chiral space groups. In this class of magnets, the left- or right-handed helical spin texture along the crystallographic axis is stabilized by a competition between symmetric exchange and antisymmetric Dzyaloshinkii-Moriya (DM) exchange interactions. Then so-called chiral soliton lattice (CSL) is formed under weak magnetic field applied perpendicular to the helical axis. The CSL behaves as a magnetic superlattice potential and results in Bragg scattering of conduction electrons. Tuning of the weak magnetic field enables us to control a size of the superlattice Brillouin zone and gives rise to a series of resistivity anomalies originating from resonant Bragg scatterings. I will also mention CSL-related phenomena from theoretical viewpoints [3-5]. This talk is given from theoretical sides. A report from experimental sides will be given in the same session by Prof. Yoshihiko Togawa from Osaka Prefectural University.

[1] J. Kishine, I. Proskurin and A. S. Ovchinnikov, Phys. Rev. Lett. 107, 017205 (2011). [2] Y. Togawa, T. Koyama, K. Takayanagi, S. Mori, Y. Kousaka, J. Akimitas, S. Nishihara, K. Inoue, A. S. Ovchinnikov, and J. Kishine, and p. Proskarin, Phys. Rev. Lett. [3] J. Kishine, A. S. Ovchinnikov, and I. Proskarin, Phys. Rev. B 20, 064407 (2010). [4] J. Kishine and A. S. Ovchinnikov, Phys. Rev. B 79, 220405(R) (2000). [5] I. Bosrem, J. Kishine and A. S. Ovchinnikov, Phys. Rev. B 79, 201405(R) (2000). [5] I. Bosrem, J. Kishine and A. S. Ovchinnikov, Phys. Rev. B 79, 201405(R) (2000). [5] I. Bosrem, J. Kishine and A. S. Ovchinnikov, Phys. Rev. B 70, 204452 (2008).

### **DF04**

### The hexagonal spin structure of A-phase in MnSi

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Inspired by recent work [1] we have revisited the cubic helimagnet MnSi to address the question of the origin of A-phase in the (H-T) phase diagram. In the present experiments we used the experimental geometry described in [1] when the magnetic field was applied along the neutron beam. We have mapped the A-phase boundaries in the field-temperature (H-T) phase diagram for three principal crystal-to-field orientations (H || [111], H || [110], H || [100]). The A-phase revealed itself on the neutron diffraction map as a hexagonal pattern of Bragg spots in a narrow range of the fields [0.12-0.20] T close to  $ST_CS = 29$  K. The orientational and translational orders (the directions and value of structure wavevectors k) are well preserved within the A-phase over the whole crystal of the size of 100 mm<sup>3</sup>. The small angle neutron scattering ascribed to the orientationally disordered hexagon spin structure was observed beyond the A-phase boundaries in the field range from H\_D1 = 0.1 T to H\_D2 = 0.25 T at temperatures down to 25 K. Contrary to the A-phase boundaries, the values of H\_D1 and H\_D2 are temperature independent for all the field-to-crystal orientations.

[1] S. Muhlbauer, B. Binz, F. Jonietz, C. Pfleiderer, A. Rosch, A. Neubauer, R. Georgii, P. Boni, Science 323 (2009) 915-919.

### **DG01**

#### Preparation and analysis of ni nanowires on si gratings

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Ni was e-beam evaporated onto a lithographically prepared Si grating under a shallow angle of incidence. By geometrical self-shading effects, Ni nanowires of 10nm x 10nm cross-section, a spacing of 750nm and a length of several cm, could be deposited homogeneously on a surface of approx. 4cm. 10nm of Al were sputter deposited as capping layer. The structural and magnetic properties of this sample will be presented: Initial structural analyses of the prepared Ni-nanowires were performed by SEM imaging. To confirm these results and to probe the magnetic character of the Ni nanowire sample, AFM and MFM micrographs were taken. The lateral periodicity over a macroscopic distance and the buried sample structures were probed by off-specular x-ray scattering and a Distorted-Wave Born Approximation based analysis was performed: first the specular reflectivity was extracted and a corresponding theoretical model was found by computer aided fitting. The obtained transverse layer structure with corresponding averaged scattering length densities was then modified by lateral variations and used to simulate the off-specular scattering behavior with excellent agreement between the measured and simulated intensities. These outstanding off-specular x-ray scattering nesults motivated polarized off-specular neutron scattering measurements which will also be presented.

### DG02

#### Elaboration and characterization of Cu/Co multilayered nanowires Julien Bran<sup>1</sup>, Malick Jean<sup>1</sup>, Rodrigue Larde<sup>1</sup>, Jean-marie Le Breton<sup>1</sup>\* and Alain

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The Giant Magnetoresistance (GMR) effect discovered in 1988 [1] had inspired many researches on magnetoresistive materials. It has been predicted by the Valet-Fert model [2] that this effect is higher when the current is perpendicular to the layers (Current-Perpendicular-to-Plane CPP geometry). Over the past ten years, a lot of research works were focused on multilayered nanowires. Indeed, these structures, exhibiting a large aspect ratio (length/diameter), provide an ideal opportunity to investigate the CPP-GMR effect and represent good candidates for the development of new technologies in the high density storage domain. In our study, Cu/Co multilayered nanowires were elaborated by a low cost electrochemical process consisting of reducing alternatively the Cu<sup>5-</sup> and Co<sup>5-</sup> ions on a metallic substrate through a nanoporous template. Transmission Electron Microscopy (TEM) with Energy Dispersive Spectrometry was used to determine the roughness and the composition of each layer. Atom Probe Tomography (ATP) [3] was also used, additionally to conventional techniques, to characterize our samples. The (Cu20nm) canowires exhibit a ferromagnetic behaviour, with a strong anisotropy. The magnetoresistance reaches 6% after a heat treatment at 450°C for 1 hour. The correlation between nanostructure and both magnetor: and magnetoresistive properties of the nanowires will be discussed.

 M. N. Baibich, J. M. Broto, A. Fert, F. Nguyen Van Dau, F. Petroff, P. Etienne, G. Creuzet, A. Friederich, and J. Chazelas, Phys. Rev. Lett, 61(21) (1988) 2472-2475. [2] T. Valet and A. Fert, Phys. Rev. B, 48 (1993) 7099-7113. [3] D. Blavette, B. Deconihout, A. Bostel, J.M. Sarrau, A. Menand, Rev. Sci. Instrum, 64 (1993) 2911.

### DG03

### Microstructure and magnetic properties of as-deposited and annealed FeCo-based nanowires

Cristina Bran<sup>1</sup>\*, Javier Garcia<sup>2</sup>, Victor Prida<sup>2</sup>, Rafael Perez Del Real<sup>1</sup> and Manuel Vazquez <sup>1</sup>Institute of Materials Science of Madrid, CSIC. 28049 Madrid, Spain <sup>2</sup> Dept. Fisica, Universidad de Oviedo. 33007 Oviedo, Spain

Ordered arrays of ferromagnetic nanowires are recently becoming the subject of intense research due to their potential as magnetic memory and sensor devices, as well as in spintronic and medical applications [1-2]. FeCo nanowires exhibit the necessary capability to be employed in novel generation of rareearth free nanocomposite permanent magnets due to their high Curie temperature, large saturation magnetization, and relatively strong magnetorystalline and inherent large shape anisotropies. Magnetic properties of FeCo alloys can be further tailored by adding other elements [3]. Anodic Aluminum Oxide membranes have been employed as templates to prepare highly ordered FeCo, FeCoCu alloy single nanowire and FeCo/Cu multilayer nanowire arrays with different diameters and alloy concentrations. As-deposited nanowire arrays were annealed in the temperature dependence of magnetic behavior have been investigated revealing the role of the effective magnetic anisotropy in the coercivity mechanism. With increasing the annealing temperature hey during the deposition is investigated to result in the improvement of the alignment of easy axis and induction of magnetic anisotropy.

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### DG04

### FMR behavior of Co nanowire arrays

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We characterize and analyze the microwave behavior of columnar arrays of Co nanowires (NW). The 50-60 nm diameter wires were electrochemically grown with different lengths from 5 to 50 microns within anodic alumina templates. The NW arrays present a high Ms, close to the Co bulk value and variable remanence states. Depending on the combination of shape, magnetocrystalline anisotropy and dipolar interactions, FMR frequencies from 7 to 25 GHz at 0 applied field were observed. The wire composition, structure and the morphology were determined by XRD and FESEM. The magnetic properties were measured using by AGFM, with field applied perpendicularly or parallel to the wire axis. The macroscopic magnetic behavior appears to be dominated by long-range magnetostatic interactions, but the magnetization state of each nanowire, as determined by the AFM/MFM at the membrane surface[1], is saturated in the up or down direction. A complete characterization and analysis of the field-and-frequency modulated microwave behavior of the arrays was performed using an electromagnet to sweep the magnetic field from -1.3 T to + 1.3 T along the wire axis, while measuring the absorption at single frequency values, from 1 to 40 GHz[2]. The results are interpreted with reference to measured magnetization curves.

[1] J. De La Torre Medina, et al., PRB 81, 144411 (2010) [2] B.K. Kuanr JMMM 286 (2005) 276-281

### DG05

#### In situ magnetic field dependent Lorentz microscopy in Co nanowires grown by focused electron beam induced deposition

Luis Alfredo Rodriguez<sup>1</sup>, Cesar Magen<sup>\*</sup>, Luis Serrano-ramon<sup>3</sup>, Etienne Snoeck<sup>4</sup>, Rosa Cordoba<sup>1</sup>, Jose Maria De Teresa<sup>3</sup> and Manuel Ricardo Ibarra<sup>1</sup>

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The accurate characterization and control of the magnetic configuration in nanostructures, namely domain walls (DWs) or vortex structures, by external parameters such as the magnetic field is essential for applications such as information storage, sensing or magnetic logic [1]. Transmission electron microscopy (TEM) allows the observation of the magnetic nanostructures with nanometer-range spatial resolution by Lorentz microscopy (LM). In this work, we performed in 1situ characterization and magnetic-field manipulation of DWs by LM in Co nanovires (NWs) in a FEI Titan Cube 60-300 TEM. These NWs were fabricated by focused electron beam induced deposition (FEBID) with purity higher than 90% on Si<sub>2</sub>N<sub>4</sub> membranes [2]. We determined the nucleation and propagation field of DWs by direct observation of the magnetic structure by LM of curved L-shape Co NW with varying width (w=125-1000 nm) and thickness (t=5-30 nm), demonstrating that for specific dimensions they are good DW propagators [3]. Focal series in LM were acquired in order to map the in-plane magnetic induction by solving the Transport-of-Intensity Equation. The nucleation process gives rise to transversal DW in the thinner NWs (t < 13 nm). Above this value, the crossover to complex structures was investigated, including multiple vortex walls on the thicker NWs.

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### **DG06**

### Morphology and magnetic properties of GaAs/(Ga,Mn)As core-shell nanowires on Si (111) synthesized by self-catalyzed method

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Diluted magnetic semiconductors, of which (Ga,Mn)As as a representative, and semiconductor nanowires (NWs) both have been hot issues of research in the past decade, due to their potential to enlarge and strengthen semiconductor based technology. Much work has been done on both systems respectively while few efforts are invested in their combined system until recently. It is of advantage to integrate such system on silicon, the most commonly used semiconductor, which our work features. Here, we present our work on GaAs/(Ga,Mn)As core-shell NWs synthesized on Si (111) substrate. Another characteristic of our work is self-catalyzed growth which utilizes no foreigner particles as seeds. We have attempted to incorporate different manganese content but it is found only a narrow growth window is permitted for successful fabrication. Magnetic properties of GaAs/(Ga,Mn)As core-shell are measured by SQUID. Magnetic hysteresis loops (M-H) along in-plane and out-of plane direction at 5 K show robust ferromagnetism and temperature dependence of remanent magnetization (M-T) reveals the Curie temperature of 20 K.

### DH01

Novel functionality and devices via complex oxide heteroepitaxy Yuri Suzuki<sup>1\*</sup>, Franklin Wong<sup>2</sup>, Chunyong He<sup>2</sup>, Brittany Nelson- Cheeseman<sup>2</sup> and

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 <sup>3</sup> Advanced Light Source, Lawrence Berkeley National Laboratory, USA

Interfaces of complex oxides materials provide a rich playground for the exploration of novel properties not found in the bulk constituents but also for the development of functional interfaces to be incorporated into technological applications. In this talk, I will present two recent examples of our work. We have demonstrated metallicity in epitaxial LaTiO<sub>3</sub> and LaVO<sub>3</sub> thin films under compressive epitaxial strain, although both are Mott insulators in the bulk; films under little to no epitaxial strain exhibit insulating behavior. The metallicity in LaTiO<sub>3</sub> can be understood in terms of a modification of the electronic structure due to anisotropic lattice distortions. The metallicity in LaVO3 is an interface effect that may be attributed to electronic reconstruction. We have also developed a new class of spin filter devices composed entirely of magnetic layers. In these devices, spin polarized conduction is dominated by interface scattering or localized states depending on temperature. Through careful investigation of the magnetism at the interfaces, we have found that there is no magnetic coupling at one interface and strong interface coupling at the other. The nature of magnetism at these interfaces is essential in the realization of our allferromagnet junction device.

### **DH02**

### Quantum oscillations and subband properties of the LaAlO<sub>3</sub>/SrTiO<sub>3</sub> heterointerface

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Heterostructures of transition metal oxides offer myriad possibilities for multifunctional materials and device applications, as well as demonstrating new and complex fundamental properties [1-4]. Much of their potential stems from the possibility of interface properties that differ dramatically from those of the bulk components, and the discovery of a 2-dimensional electron gas at the interface between band insulators LaAlO, and SrTiO;[1] is a striking realisation of this. Understanding and controlling this 2D electron gas is currently the central problem in the field. Here I will present the results of recent magnetotransport measurements of the LaAlO/SrTiO; interface, performed at millikelvin temperatures and in magnetic fields of up to 30 T. Quantum oscillations observed in the resistivity allow us to probe the bandstructure of the interface electron gas, and extract important information about the band energies and the properties of the mobile charge carriers. The strong dependence of the oscillations on magnetic field, temperature, and orientation with respect to the plane of the interface, reveals a complex 2D system of closely-spaced multiple electronic subbands with different carrier effective masses and mobilities. Measurements in tilted magnetic fields show evidence of inter-subband scattering and diamagnetic shift of the 2D confinement energies.

[1] A. Ohtomo et al., Nature 427, 423 (2004). [2] A. Brinkman et al., Nature Materials 6, 493 (2007). [3] N. Reyren et al., Science 317, 1196 (2007). [4] H.Y. Hwang et al., Nature Materials 11, 103 (2012).

### DH03

### Electronic ordering in sodium cobaltate

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Sodium cobaltate (NaxCoO<sub>2</sub>) has emerged as a material of exceptional scientific interest due to the potential for thermoelectric applications, and because the strong interplay between the magnetic and superconducting properties has led to close comparisons with the physics of the superconducting copper oxides. The density x of the sodium in the intercalation layers can be altered electrochemically, directly changing the number of conduction electrons on the triangular Co layers. Previous measurements have determined the long range ordering of sodium superstructures for different values of x and numerical calculations have found that this ordering has a profound effect on the charge carrying cobalt layer. Multiple valence states have been detected by NMR, but here we report new spatial measurements using Resonant X-ray Scattering on the Materials and Magnetism beamline 116 at Diamond. The sodium vacancies are found to condense into tri-vacancy clusters, and these order long range into stripes. We find resonant x-ray scattering with the same periodicity as the sodium superstructure directly demonstrating that the electronic ordering in these cobalt layers is controlled by the sodium ordering. We are able to reproduce the energy, polarisation and azimuthal dependencies of the resonant x-ray scattering in calculations using the FDMNES code

### 62 | The 19<sup>th</sup> International Conference on Magnetism

### **DH04**

#### **Electrical switching of the magnetic phase in semiconductor oxides** Antonio Ruotolo<sup>1</sup>, Xiao Lei Wang<sup>1</sup>, Chi Wah Leung<sup>2</sup> and Rolf Lortz<sup>3</sup>

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The magnetic properties of thin-film ferromagnetic semiconductors can be altered by applying an electric field[1]. The effect is reversible but volatile. However, for technological applications, it would be desirable to have a film that retains the induced magnetic state until a reset command is given. We here demonstrate that the magnetic moment of a ferromagnetic oxide can be altered in a non-volatile and reversible manner. The studied system consists of  $Zn_{0.98}Mn_{0.02}O$  (ZMO) film sandwiched between a highly conductive Niobium Strontium Titanate (NSTO) substrates and a Platinum (Pt) top electrode. The current-voltage curves revealed that the electrical behavior was dominated by the Schottky interface between the NSTO and the ZMO. Reversable resistive switching[2,3] could be induced in the samples with high reliability and good retention time. By measuring the full magnetization loop, we found that a switching of the resistance corresponds to a switching of the magnetic phase. Remarkably, the threshold switching current was ~ 1 A/cm2 in film devices, as well as in patterned devices. This is orders of magnitude smaller than that reported in similar, non-magnetic systems[2,3].

[1] Ohno, Nature 408, 944-946 (2000). [2] Yang, Nat. Nanotech. 3, 429 (2008). [3] Lee, Nat. Mater. 10, 625 (2011).

#### DH05

### An approach to achieve layered spintronics material using Brownmillerite compound $Ca_{2.5}Sr_{0.5}GaMn_2O_8$

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Naturally occurring layered magnetic materials can play a significant role in shaping the next generation technology in the area of magnetic storage. We have, therefore, studied the magnetic and magnetotransport properties of the La-substituted compounds  $Ca_{2.5x}La_xSr_{0.5}GaMn_2O_8$  (x = 0, 0.05, 0.075, and 0.1). We have succeeded to introduce the double exchange (DE) interaction in the pure antiferromagnetic system  $Ca_{2.5}Sr0_{0.5}GaMn_2O_8$  by varying the ratio between  $Mn^{3+}$  and  $Mn^{4+}$  ions. The experimental results of dc magnetization, magnetotransport and neutron diffraction studies were interpreted on the basis of an electronic phase separation model, where a formation of ferromagnetic clusters (due to DE interaction) in the La-substituted compounds is proposed. Significantly, the present study shows that the magnetic and electronic properties of the layered system  $Ca_{2.5}Sr0_{0.5}GaMn_2O_8$  can be tuned/optimized by appropriate chemical substitution to achieve a new spintronic material based on naturally occurring layered system for practical applications.

### DI01

### Magnon-drag thermopile

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Thermoelectric effects in spintronics are gathering increasing attention as a means of controlling spin information by using heat flow. Thermal magnons (spin-wave quanta) are expected to play a major role, however, the coupling between electrons and magnons in ferromagnetic metals remains poorly understood. We demonstrate a conceptually new device that enables us to gather information on magnonelectron scattering and magnon-drag effects[1]. The device resembles a thermopile formed by a large number of pairs of ferromagnetic wires placed between a hot and a cold source and connected thermally in parallel and electrically in series. By controlling the relative orientation of the magnetization in pairs of wires, the magnon drag can be studied independently of the electron and phonon drag thermoelectric effects. Measurements as a function of temperature reveal the effect on magnon drag following a variation of magnon and phonon populations. These results demonstrate the feasibility of directly converting magnon dynamics of nanomagnets into an electrical signal and could pave the way to novel thereelectric devices for energy harvesting. This research was supported by the Spanish Ministry of Science (MICINN/ MAT2010-18065) and by the European Community (*PP*/NANOFUNCTION-257375).

[1] M.V. Costache, G.A. Bridoux, I. Neumann and S.O. Valenzuela, Nat. Mater. Advanced Online Publication (2011). http://dx.doi.org/10.1038/NMAT3201

### DI02

#### Seebeck spin tunneling in silicon Ron Jansen

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The combination of thermoelectrics and spintronics offers unique possibilities. On the one hand, it provides a new, spin-based approach to thermoelectric power generation and cooling. On the other hand, it provides a thermal route to create and control the flow of spin in spintronic devices that make functional use of heat and temperature gradients. Here we describe and report the demonstration of Seebeck spin tunneling a thermally driven spin flow, of purely interfacial nature - generated in a tunnel contact between electrodes of different temperatures. Seebeck spin tunneling arises from the spin dependence of the Seebeck coefficient of a tunnel junction and is thus the thermoelectric analog of spin-polarized tunneling. By exploiting this in ferromagnetoxide-silicon tunnel junctions, we observe a thermal flow of spin angular momentum from the ferromagnet to the silicon without a charge tunnel current. The spin accumulation induced in the silicon scales linearly with heating power and changes sign when the temperature differential is reversed. This thermal spin current can be used by itself, or in combination with electrical spin injection. The results highlight the engineering of heat transport in spintronic devices and enable the (re-)use of heat to increase device efficiency and reduce energy consumption.

J.C. Le Breton, S. Sharma, H. Saito, S. Yuasa and R. Jansen, Nature 475, 82 (2011). R. Jansen, A.M. Deac, H. Saito and S. Yuasa, Preprint at http://arxiv.org/abs/1112.3430 (2011).

### DI03

#### Tunneling magneto Seebeck effect

Andy Thomas<sup>1</sup>, Markus Munzenberg<sup>2</sup> and Christian Heiliger<sup>3</sup> <sup>1</sup> Physics, Bielefeld University, Germany <sup>2</sup> Physics, Gottingen University, Germany <sup>3</sup> Physics, Gieβen University, Germany

Creating temperature gradients in magnetic nanostructures has resulted in a new research direction, that is, the combination of magneto- and thermoelectric effects. Here, we demonstrate the observation of one effect of this class: the tunneling magneto-Seebeck effect. It is observed when a magnetic configuration changes the charge-based Seebeck coefficient. In particular, the Seebeck coefficient changes during the transition from a parallel to an antiparallel magnetic configuration in a tunnel junction. In this respect, it is the analog of the tunnelling magnetoresistance. The Seebeck coefficients in parallel and antiparallel configurations are of the order of the voltages known from the charge Seebeck effect. The size and sign of the effect can be controlled by the composition of the electrodes' atomic layers adjacent to the barrier and the temperature. We realized up to 40% tunneling magneto-Seebeck effect and up to 330% TMR ratio in MgO based magnetic tunnel junctions. We show the temperature dependent investigations of this effect, supported by TEM studies of the microstructure, numeric simulations of the temperature profile in the junction and possible devices utilizing the tunneling magneto Seebeck effect.

[1] M. Walter et al., Seebeck effect in magnetic tunnel junctions, Nature Materials 10 (2011) 742

#### DJ01

#### Novel clathrate-based composite materials for energy-efficient magnetic refrigeration

Anurag Chaturvedi, Stevce Stefanoski, George S. Nolas, Hariharan Srikanth and Manh-huong Phan\*

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Magnetic refrigeration based on the magneto-caloric effect (MCE) is an area of exciting research. The grand challenge in this technology is to seek new material that exhibits large MCE over a wide temperature range, namely, the large refrigerant capacity (RC). In this talk, we will present our recent discoveries of the large MCE and RC in  $Eu_8Ga_{16}Ge_{30}$  type-1 and type-VIII clathrates and  $Eu_8Ga_{16}Ge_{30}$  type I clathrate leads to the broadened MCE curve thus enhancing the RC in this material. This new finding may provide a route for improving the cooling efficiency in type-I clathrate materials for active magnetic refrigeration (AMR). The large table-like MCE and enhanced RC have been achieved in the  $Eu_8Ga_{16}Ge_{30}$  type I - EuO composite materials, making them one of the best candidates for AMR in the liquid nitrogen temperature range.

### DJ02

### Magnetocaloric properties of doped $La_{0.7}Sr_{0.3}MnO_3$ bulk ceramic and thick films

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The LaSrMnO<sub>3</sub> manganite has attracted research interests due to its colossal magneto resistance effect, half-metallic behavior and magnetocaloric effect [1-3]. The magnetocaloric effect is the reversible temperature change of a magnetic material associated with external magnetic field change in an adiabatic process. These properties are strongly related with the interactions between the rareearth and transition metal's charge states which can be varied by doping elements [4, 5]. In this study, the physical properties of magnetocaloric La<sub>0</sub>, Sr<sub>0</sub>, MnO<sub>3</sub> (LSMO) ceramic system, in both bulk and thick film, have been investigated according to various B-site elements doping. The substituted B-site elements were transition metal elements, i.e. Cr, Co, Ni, Cu with its contents of  $0 \sim 3$  %. The ceramic samples were prepared by conventional solid-state reaction process. The thick films were grown by Aerosol Deposition (AD) method [6, 7] on various substrates. The magnetocaloric properties measured by physical property for the system (PPMS) in magnetic field up to 30 kOe and temperature range up to 400 K. In this presentation, the magnetocaloric properties of the entropy change, TC variation, Banerjee criterion [8] etc. according to the B-site doping and substrate effects will be discussed in detail.

 Physical Review Letters, 71 (1993) 2331-2333. [2] Reports on Progress in Physics, 68 (2005) 1479-1539. [3] J Magn Magn Mater, 308 (2007) 325-340. [4] Physical Review Letters, 74 (1995) 5144-5147. [5] Journal of Physics Condensed Matter, 20 (2008).
 [6] Sensor Actuat a-Phys, 153 (2009) 89-95. [7] Materials Letters, 65 (2011) 2762-2764. [8] Physics Letters, 12 (1964) 16-17.

### DJ03

### Structural ,magnetic and magnetocaloric properties of $Ni_{50}Mn_{37.5}Sn_{12.5}$ ribbon Heusler alloys

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We present a detailed study of the magnetic and microstructural behavior of  $Ni_{50}Mn_{375}Sn_{125}$  Heusler ribbon alloy, which valence electron concentration per atom = 8.12. This ribbon, as flakes, around 1.5 - 2.0 mm width and 4 -5 mm length, was obtained after quenching the master alloy by the meltspinning technique in Ar atmosphere. Scanning Electron Microscopy (SEM, JEOL 6100) revealed a columnar structure of the Heusler-Based alloy. The averaged composition was determined by EDX (Inca Energy 200) from measurements taken at different points of the sample, obtaining  $Ni_{eq}AM_{36}Sn_{139}$ . Differential Scanning Calorimetry (DSC) and thermomagnetic curves confirm an austenite-martensite structural transformation below 400K. The start and finish temperatures of the martensitic phase transformation for the ribbon alloy were Ms = 279 K and Mf = 260 K, while the ones for austenite were As = 272 K and Af = 288 K, in agreement with previously reported values in similar compositions [1]. Moreover, it has been performed a study of magnetocaloric effect (MCE) by thermomagnetic measurements from which a direct and inverse magnetocaloric effects near Curie point and martensitic transformation temperature have been respectively observed.

[1] S.E. Muthu et al., J. Phys. D : Appl. Phys., 43 (2010) 425002

### **DJ04**

### Dependence of the magnetocaloric effect in ferromagnetic shape memory Heusler alloys on measurement protocol

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Results of direct measurements of the magnetocaloric effect in ferromagnetic shape memory alloys (FSMAs) Ni-Mn-X (X = Ga, In, Sn) showed that the adiabatic temperature change  $\Delta$ Tad is (i) small (typically smaller than in prototypical magnetocaloric material Gadolinium) and (ii) depends on thermal as well as on magnetic history of FSMAs samples. Here we discuss these features of  $\Delta$ Tad and show that the measurement protocol determines whether the adiabatic temperature change is reversible or irreversible. These characteristic of  $\Delta$ Tad are easy to understand considering thermodynamics of temperature- and magnetic field-induced thermoelastic martensitic transformations in the vicinity of the magnetostructural phase transition.

### DJ05

Magnetocaloric effects in manganites with perovskite structure Abdelwaheb Cheikhrouhou\*, Wissem Cheikhrouhou-koubaa and Mohamed Koubaa Materials Physics Laboratory, Faculty of Sciences of Sfax, Tunisia

Magnetocaloric refrigeration is a promising alternative technique to replace the conventional gas expansion-compression one based on the use of toxic gases. This novel technology presents several advantages, in fact besides the property of being environmentally friendly (it does not use coolants or emit greenhouse gases), it undergoes a higher efficiency and it is compact in size. Magnetocaloric materials, when subjected to a magnetic applied field at their phase transition temperature, undergo a magnetic transition which results in the entropy change of the magnetocaloric materials. This phenomenon is known as the magnetocaloric effect (MCE). The family of magnetic oxides with perovskite structure known as manganites have also shown interesting MCE, for a large temperature range around room temperature. In this work, we were interested in the study of the MCE in manganites with several compositions. Several compounds with general formula La<sub>1-x</sub>M<sub>x</sub>MnO<sub>3</sub> (M is a divalent alkali-earth or monovalent elements) have been elaborated and characterized. The entropy change versus temperature and versus magnetic applied field will be presented. Our results, compared to that observed in Gd or GdSiGe, show that the manganites exhibit a giant MCE around or above room temperature and are suitable for applications in the magnetic refrigeration area.

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### EA01

### Quantum phase transitions in heavy-fermion systems

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The canonical heavy-fermion system CeCu<sub>6-A</sub>Au<sub>e</sub> exhibits a quantum phase transition (QPT) that can be accessed by varying the Au concentration x, hydrostatic pressure p, and magnetic field B. Its unusual features have prompted the suggestion of a novel type of QPT not envisaged by the extension of the Landau-Ginzburg-Wilson theory of (classical) phase transitions to T = 0 by Hertz, Millis and Moriya. We will review a range of new experiments on this system: (1) probing the Kondo resonance upon approaching the QPT by photoelectron spectroscopy, (2) establishing a link between the pressure and concentration dependence of the magnetic ordering vector as determined by elastic neutron scattering, (3) comparing the volume and magnetic Gruneisen parameters obtained from thermal expansion and magnetocaloric effect, repectively, and (4) determining the thermoelectric power close to the critical concentration  $x_c$ . CePdAI is a heavy-fermion system with partial geometric frustration of one of the three Ce atoms in the unit cell. We will report on first experiments exploring the possible role of frustration on quantum criticality which can be accessed by Ni doping.

### EA02

### Lifshitz transitions and non-fermi liquid behavior in heavy-fermion metals

Matthias Vojta\* Institut fuer Theoretische Physik, Technische Universitaet Dresden, Germany

The field-driven non-Fermi liquid behavior in YbRh<sub>2</sub>Si<sub>2</sub> has been proposed to arise from a reconstruction of the Fermi surface due to a breakdown of Kondo screening. Here we discuss an alternative scenario, namely a Zeeman-driven Lifshitz transition of a narrow heavy-fermion band, with Kondo screening remaining intact. Apparent non-Fermi liquid behavior emerges in a crossover regime above a small effective Fermi energy. As carrier doping provides a unique knob to tune the Lifshitz transition, we discuss the evolution of heavy-fermion bands with doping in a quantitative fashion and connect our results to recent data obtained on Fe-doped YbRh<sub>2</sub>Si<sub>2</sub>.

### EA03

#### Sequential spin polarization of the fermi surface pockets in URu<sub>2</sub>Si<sub>2</sub> and its implications for the hidden order

Neil Harrison<sup>1</sup>, Moaz Altarawneh<sup>1</sup>, Luis Balicas<sup>2</sup>, P H Tobash<sup>1</sup>, J D Thompson<sup>3</sup>, F Ronning<sup>3</sup> and E D Bauer<sup>3</sup>

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Using Shubnikov-de Haas oscillations measured in URu<sub>2</sub>Si<sub>2</sub> over a broad range in a magnetic field of 11-45 T, we find a cascade of field-induced Fermi surface changes within the hidden order phase I and further signatures of oscillations within field-induced phases III and V. A comparison of kinetic and Zeeman energies indicates a pocket-by-pocket polarization of the Fermi surface leading up to the destruction of the hidden order phase I at ~35 T. The anisotropy of the Zeeman energy driving the transitions in URu<sub>2</sub>Si<sub>2</sub> points to an itinerant hidden order parameter involving quasiparticles whose spin degrees of freedom depart significantly from those of free electrons.

Physical Review Letters 106 (14):146403 [2011]

### EA04

Hydrostatic pressure study of the nematicity of Sr<sub>3</sub>Ru<sub>2</sub>O<sub>7</sub> Dan Sun<sup>1</sup>, Wenlong Wu<sup>1</sup>, Santiago Grigera<sup>2</sup>, Robin Perry<sup>3</sup>, Andrew Mackenzie<sup>2</sup> and Stephen Iulian<sup>1\*</sup>

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A 'electron-nematic phase' occurs in the bilayer ruthenate  $Sr_3Ru_2O_7$ , at very low temperature in the vicinity of a metamagnetic transition. It has been seen when the applied magnetic field, B, makes a large angle with the ab-plane [1,2,3]. The nematic phase is characterized by a spontaneously broken rotational symmetry for resistivity within the ab-plane. The underlying physics is not understood, but it is believed to be associated with a 'quantum critical endpoint' (QCEP) for the metamagnetic transition [2]. We show, using transport and magnetic susceptibility measurements on ultra-pure crystals, that this metamagnetic transition does not produce a nematic phase when the QCEP is reached by applying hydrostatic pressure [4] with the field applied in ab-plane. Moreover, a nearby region of nematicity is found to be robust against the application of pressure. Taken together our results suggest that proximity to a quantum critical point may not be the crucial ingredient for the nematicity.

[1] S. A. Grigera et al., Science 306, 1154 (2004). [2] R. A. Borzi, S. A. Grigera, J. Farrell, S. J. S. Lister, S. L. Lee, D. A. Tennant, Y. Maeno, and A. P. Mackenzie, Science 315, 214 (2007), [3] A.W. Rost, R. S. Perry; J. F. Mercure, A. P. Mackenzie, and S. A. Grigera, Science 325, 1360 (2009), [4] W. Wu et al., Rev. B 83, 045106 (2011).

### EB01

#### Pump-probe response for correlated electron systems out of equilibrium T. Devereaux\*

Stanford University, USA

Pump-probe techniques have been used to extend equilibrium spectroscopic methods to the time-domain, revealing temporal dynamics associated with lattice degrees of freedom, collective modes of novel electronic phases and decay and recovery processes with typical time scales of several picoseconds. Recent advances have extended the temporal resolution of these techniques to femtosecond and attosecond time scales. In this talk I will present several severalKeldsyh-based calculations and simulations for pump-probe spectroscopy across a metal-insulator transition, across a charge-density wave transition, and for electron-phonon coupled lattices. These results illuminate the intricate dynamics of how correlated materials return to equilibrium after photo-excitation.

### EB02

### Correlated electrons in strong electric fields

Philipp Werner<sup>1</sup>, Martin Eckstein<sup>2</sup>, Naoto Tsuji<sup>1</sup>, Takashi Oka<sup>3</sup> and Hideo Aoki<sup>3</sup> <sup>1</sup> Department of Physics, University of Fribourg, Switzerland <sup>2</sup> Max Planck Research Department for Structural Dynamics, University of Hamburg - CFEL, Germany

<sup>3</sup> Department of Physics, University of Tokyo, Japan

We use the nonequilibrium extension of dynamical mean field theory to study the dynamics of the Hubbard model in the presence of external electric fields. Specifically, we discuss the transition from a Bloch oscillating to a direct current regime in weakly interacting systems subject to static electric fields [1], and the dielectric breakdown of the Mott insulator [2]. In periodically driven metallic systems, the electric field may lead to a dynamical band inversion and an associated conversion of the effective on-site interaction from repulsive to attractive [3].

[1] M. Eckstein and P. Werner, Phys. Rev. Lett. 107, 186406 (2011) [2] M. Eckstein, T. Oka, and P. Werner, Phys. Rev. Lett. 105, 146404 (2010) [3] N. Tsuji, T. Oka, P. Werner, and H. Aoki, Phys. Rev. Lett. 106, 236401 (2011)

### EB03

### Photoinduced charge order enhancement in one-dimensional extended hubbard model

Hantao Lu<sup>1</sup>, Shigetoshi Sota<sup>1</sup>, Hiroaki Matsueda<sup>2</sup>, Janez Bonča<sup>3,4</sup> and Takami Tohyama<sup>1</sup>

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We present a compelling response of a low-dimensional strongly correlated system to an external optical perturbation[1]. By using time-dependent Lanczos method, the nonequilibrium evolution of one-dimensional Hubbard model at half-filling driven by transient laser pulse is investigated. When the system is close to the phase boundary, by tuning the laser frequency and strength, a sustainable charge order enhancement is found that is absent in the Mott insulating phase. We analyze the conditions and investigate possible mechanisms of emerging charge order enhancement. From the spectrum analysis on small size systems, it can be shown that with increasing the nearest-neighbor interactions, charge-order favorite eigenstates proliferate in the low energy regime, leaving themselves prone to be picked up by laser pulse. Feasible experimental realizations are proposed.

[1]Hantao Lu, S. Sota, H. Matsueda, Janez Bonča, and T. Tohyama, arXiv:1204.1107.

### EB04

### Electric-Field Effects on Complex Oxide Interfaces: Possible Two-Band Superconductivity

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We examine the superconductivity produced at the complex oxide interface of LaAlO<sub>3</sub> and SrTiO<sub>3</sub> (Tc =  $\sim$ 0.3 K). Through a phenomenological and semi-classical approach, we investigate the effects of local and global electric fields on the superconducting order parameter, carrier density, and transition temperature for this system. We show that the general shape of the superconducting dome can be reproduced through basic electrostatics and a standard BCS formalism using the density of states. Given the similarity between the transition temperature for the interface and of bulk SrTiO<sub>3</sub>, which is also superconducting at -0.3 K upon optimal electron doping, we provide evidence that the interface superconductivity is intrinsic to SrTiO<sub>3</sub> and we infer that superconducting mechanism must also be similar. By examining bulk SrTiO<sub>3</sub>, we show that the superconductivity may be described through weakly-coupled, two-band model. Finally, we predict that the superconductivity at the interface is most likely produced through a two-band coupling of the surface states of SrTiO<sub>3</sub> provided by an electron doping from a carrier source (in this case LaAlO<sub>3</sub>).

### EC01

### **Electronic structure and phonons in the high pressure phases of cerium** Borje Johansson<sup>1,2</sup> \* and D.Y. Kim<sup>3</sup>

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The isostructural  $\gamma$ - $\alpha$  transition in cerium under pressure has attracted a very extensive interest both experimentally and theoretically over the years. Since long it has also been known that at a pressure around 6.5 GPa cerium adopts the  $\alpha$ -tranium crystal structure (oC4). Recent developments have made it possible to use inelastic X-ray scattering to determine the phonon dispersions of elemental cerium in these high pressure structures 1.2. This opened the possibility to compare the experimental data with theoretical calculations. When the 4f electrons in cerium are treated as itinerant a very good agreement was obtained. Very interestingly, also the observed pronounced phonon anomalies in Ce-oC4 were well described by theory and points out a very close similarity between uranium and Ce-oC4. At the low-pressure end of its stability range, theory suggests that Ce-oC4 is on the verge of a latticedynamical instability and possibly a charge density wave. When this success of the itinerant picture for the 4f behavior in the high pressure phases of cerium is combined with the well-known knowledge that the 4f electron in  $\gamma$ -cerium is well described by a localized picture, one arrives at the conclusion that the 4f electron in oreium is a Mott transition3.

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### EC02

### From SOC induced phenomena to non-collinear magnetism and electric field effects in magnetic systems

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Quantum computational electronic structure theory, based on accurate density functional theory and structural modeling, has taken a leadership role in providing insight into the underlying physics of magnetism. Challenges currently lie in understanding non-collinear magnetism and magnetotransport properties in magnetic domain walls [1] and, more recently, externally driven electric-field magnetocrystalline anisotropy (MA) [2] in transition-metals that arise not only through the spin-orbit coupling (SOC): these phenomena provide new pathways for future magnetic devices at the nanoscale with the promise of ultra-low energy power consumption. This talk will present our recent firstprinciple predictions for magnetism and transport properties for i) magnetic domain walls (modeled by a spin-spiral) and ii) electric-field-driven MA in transition-metals, carried out using our noncollinear magnetism full-potential linearized augmented plane-wave (FLAPW) method in which the external electric field is directly incorporated and the conductivity tensor is obtained by applying the Kubo formula of linear response theory. \*Supported at Northwestern by the U.S. D.O.E. and the Japan Society for the Promotion of Science.

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### EC03

### First-principles calculation of the A-site ordered perovskite CaCu<sub>3</sub>Fe<sub>4</sub>O<sub>12</sub>

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Oxides with anomalous-valence ion Fe<sup>4+</sup> have been studied for many years and such systems have been recently found among the so-called A-site ordered perovskites AA'<sub>3</sub>B<sub>4</sub>O<sub>12</sub>. In the AA'<sub>3</sub>B<sub>4</sub>O<sub>12</sub> perovskites, A and A' cations occupy the A-site of the original simple perovskite ABO<sub>3</sub> and the A' cations are coordinated by four O ions in a planar fashion instead of the original twelve coordination. The cation sites may consist of magnetic ions, leading to several possible magnetic interactions between them. CaCu<sub>3</sub>Fe<sub>4</sub>O<sub>12</sub> reveals charge disproportionation  $2Fe^{4+} \rightarrow Fe^{3+}+Fe^{4+}$  at 210K with a transition from paramagnetic to ferrimagnetic state. Quite recently, the magnetic couplings between cations in CaCu<sub>3</sub>Fe<sub>4</sub>O<sub>12</sub> are experimentally determined by x-ray magnetic circular dichroism (XMCD) measurements [1]. In this paper, the electronic structure of the A-site ordered perovskite CaCu<sub>3</sub>Fe<sub>4</sub>O<sub>12</sub> is studied by means of first-principles calculations. X-ray absorption spectroscopy and XMCD spectra at the Fe and Cu L2,3 edges and O K edge are calculated and compared with experiments to discuss the magnetic order and coupling.

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### EC04

Pr partial electron donation and Co spin state changes at the metal-insulator transition in  $(Pr_{1,y}Y_y)_{1-x}Ca_xCoO_3$  as seen by x-ray absorption and emission Javier Henrero-muttin<sup>19</sup>, Jose Luis Garciar-munoz<sup>1</sup>, Carlos Fronten<sup>1</sup>, Aura Janeth Baron-gorzalez<sup>1</sup>, Jessica Padilla<sup>1</sup>,

Javier Inerconnaum, Jose Luis Garcianunoz, Canos Ironieta, Auta Jarein Patorregonzatez, Jessiel Fauna Sergio Valencia<sup>2</sup>, Ralf Feyerherm<sup>2</sup>, Esther Dudzik<sup>2</sup>, Florin Radu<sup>2</sup>, Radu Abrudan<sup>3</sup>, Gloria Subias<sup>4</sup> and Javier Blasco<sup>4</sup> <sup>1</sup>Institute of Materials Science of Barcelona - CSIC, Spain <sup>2</sup>Helmholtz-Zentrum Berlin, BESSY, Germany

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The electronic and magnetic properties of cobalt-based perovskites are largely determined by the electronic filling of the Co 3d valence band. The characteristic spin state degree of freedom in many cobalities is correlated to the electron mobility and thus it can play an important role in metal-insulator transitions (MIT). MIT are found in a wide range of compositions of the (Pr<sub>1-y</sub>X<sub>y</sub>)<sub>t-x</sub>Ca<sub>x</sub>CoO<sub>3</sub> series, coupled to magnetic and structural changes [1,2]. At low temperatures, x-ray diffraction confirms a crystal cell volume contraction occurs due to the appearance of a certain quantity of Pt<sup>+</sup> ions. The electrons lost by Pr ions move to Co sites. The expected increase of the average Co-O distance is not observed, may be due to a compensation effect by a likely Co ions spin state change [3]. We present here a comprehensive study of these compounds by means of x-ray absorption and emission spectroscopies, which leads us to confirm these hypothesis also quantifying the intersite charge donation across the MIT [4]. The observed photoinduced phase transition in Pr<sub>0.5</sub>Ca<sub>0.5</sub>Co<sub>0</sub>, under illumination by laser radiation, where metallic domains can be stimulated [5], could thus be related to induced spin state changes.

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### EC05

### Spin transport in the anisotropic Heisenberg chain at finite temperature and momentum

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Magnetic transport in one-dimensional quantum spin systems displays many new features beyond a conventional Boltzmann description. Among them is the dissipation of spin currents. At zero momentum, integrability, such as for the one-dimensional Heisenberg (XXZ) model, can lead to ballistic transport. Transport at finite momentum is an open issue. Here we investigate the role of momentum for the transport of magnetization in the spin-1/2 Heisenberg chain above the isotropic point at finite temperature and momentum. Using numerical and analytical approaches, we analyze the autocorrelations of density and current and observe a finite region of the Brillouin zone with diffusive dynamics below a cut-off momentum, and a diffusion constant independent of momentum and time, which scales inversely with anisotropy. Lowering the temperature over a wide range, starting from infinity, the diffusion constant is found to increase strongly while the momentum space cut-off for diffusion decreases. Above the cut-off momentum diffusion breaks down completely.

Phys. Rev. Lett. 107, 250602 (2011)

### ED01

### Growth of metastable fcc-Fe film on Cu(100) single-crystal underlayer and phase transformation from fcc to bcc

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Theoretical studies suggest that the magnetic property of metastable fcc iron (Fe) varies in a wide range from nonmagnetic to high-spin ferromagnetic state depending on the distance between Fe atoms [1]. The structure of bulk Fe crystal is bcc up to 912 °C, whereas metastable fcc structure can be stabilized in a form of film at RT. However, the metastable fcc structure is reported to transform into more stable bcc structure with increasing the film thickness [2]. In the present study, Fe films were prepared on Cu(100) single-crystal underlayers at RT by UHV-MBE in order to investigate the growth behavior of fcc-Fe crystal and the fcc-bcc phase transformation process. In an early stage of film growth, metastable fcc-Fe(100) structure was stabilized through hetero-epitaxial growth. With increasing the thickness, the fcc structure started to transform into bcc structure. High-resolution cross-sectional TEM and polefigure XRD analyses showed that the fcc structure transformed into bcc structure by atomic rearrangement from fcc{100} to bcc{100}. The transformation orientation relationship was thus bcc{100}<011> || fcc{100}<001>. The transformation occurred not only along the perpendicular direction but also along the in-plane directions. The relationship between film thickness and magnetic property will also be discussed at the conference

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### ED02

### Neel temperature and the thickness of surface NiO

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The surface oxidation on Ni/Cu(001) or CoNi alloy on Cu(001) have been studied for the spin reorientation transition and other magnetic interactions, such as magnetic demolishment, canted state, coercivity enhancement, and antiferromagnetic induced ferromagnetism. It is found that a transition temperature of 200 K for these phenomena. This study is to confirm that the transition temperature of 200 K is the Neel temperature of NiO formed by oxygen treatment on the surface of NiCu(001) films. We deposited films of x ML Ni/10 ML Co/Cu(001), where x = 0, 2, 3, 4, and 6. The films were then treated by oxygen exposure of 1800 Langmuir for the formation of y ML Ni/0/(x-y) ML Ni/10 ML Co/Cu(001). The coercivity is enhanced for all x below 200 K after field cooling. For x of 3, 4, and 6, the coercivity is the same and the coercivity enhancement phenomenon disappears when the temperature is above 200 K. It indicates that the thickness of the surface NiO is 3 ML and the Neel temperature is 200 K. Thus, all the phenomena mentioned above is induced by the formation of antiferromagnetic ordering of the surface NiO.

### ED03

### Reduced exchange bias field in antiferromagnet-patterned FeF<sub>2</sub>/Ni stripes

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The exchange interaction in antiferromagnetic/ferromagnetic (AF/FM) bilayers gives rise to the exchange bias phenomenon. In addition to the shift of the FM hysteresis loop, the exchange coupling can also modify the magnetization reversal mechanism. Incomplete domains walls (IDWs) have been observed in FeF<sub>2</sub> spin structure due to the Strong AF anisotropy [3]. We show using variable thickness FM layers that this scenario changes as the AF loced temperature (1,2]. However, reversal of the FM moments is expected to have little effect on the FeF<sub>2</sub> spin structure due to the strong AF anisotropy [3]. We show using variable thickness FM layers that this scenario changes as the AF becomes patterned. A wedged substrate/FeF<sub>2</sub> (70mn)Ni (tNi=2-74 nm)/AI (4mn) sample was patterned into stripes by photolithography. Only the FM layer was patterned for the thickest tNi whilst the FeF<sub>2</sub> layer becomes gradually patterned as tNi decreases. The exchange-bias field-dependence Heb (tNi) shows an increasing difference compared to the value expected with decreasing Ni thickness. This result originates from spring-like DWs parallel to the AF/FM interface formed in both the Ni and FeF<sub>2</sub> layers. Simulations of IDWs account for the experimental values of Heb(tNi). Work supported by MICINN grants FIS2008-06249, MAT2010-20798, MAT2009-08667 and the US Department of Energy DE FG03-87ER-45332.

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#### ED04

### Experimental verification of the magnetic interactions between Co particles in $C_{60}$ -Co granular films

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Spintronics devices using molecular materials have attracted growing attention over the past decade because of their long spin coherence. Various researches have accomplished about the electrical conduction characteristic and the magnetic transportation characteristic of the  $C_{60}$ -Co granular film[1-5], which has a structure that Co nanoparticles are dispersed in  $C_{60}$  matrix. We observed the following two phenomena below the blocking temperature in the sample ( $C_{60}$ -Co = 9.2:1): 1) Hysteresis of the magnetization curve. (Ferromagnetic characteristic.), and 2) Magnetic field dependence of the coulomb blockade-electric field characteristics were assumed due to the dipole interaction between the Co nanoparticles, and identified by the following method. We calculated the two particles model with dipole interaction and compared the numerical results and magnetization curve of the samples( $C_{60}$ -Co = 6.0:1, 9.0:1). The measured M-H curves are good agreement with the theoretical curves. The above 1) and 2) can be explained with this model. In this way, the dipole interactions between Co nanoparticles in  $C_{60}$  matrix were identified.

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### ED05

### Emergent magnetic switching in spin glass $La_{0.7}Sr_{0.3}(Mn,Fe)O_3/La_{0.7}Sr_{0.3}MnO_3$ thin films

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The emergent magnetic switching has been investigated in exchange coupled bilayers composed of spin glass ( $La_{0.7}Sr_{0.3}Mn_{1.7}Fe_yO_3$ , y=0.18~0.20) and ferromagnet ( $La_{0.7}Sr_{0.3}MnO_3$ ) epitaxial thin films. It was found that, i) the magnetization reversal develops from single magnetic switching into double magnetic switching as the sample is cooled below the spin freezing temperature; ii) both the critical field and the transition region for the second magnetic switching (SMS) increase upon cooling but decrease with increasing the cooling magnetic field; iii) at 1.8 K the SMS is only exhibited in the descending branch of the hysteresis loop; however its presence in the ascending branch can be partially restored via enlarging the cooling field; iv) there occurs an obvious training effect for the SMS observed at 1.8 K. The rich features and the micro-origin of the emergent magnetic switching are discussed in terms of a phenomenological domain wall model which involves the exchange anisotropy depending on the alterable spin frustration due to the frozen and activated spins at the spin glass-ferromagnet interface.

### ED06

#### FM-AFM crossover in vanadium oxide nanomaterials.

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The magnetic properties of nanomaterials based on vanadium oxide (multiwall nanotubes, nanorods, whiskers and nanolayers) have been investigated in the temperature range of 1.8-220 K by high frequency (60 GHz) EPR. A transition from a ferromagnetic temperature dependence to an antiferromagnetic temperature dependence has been observed in nanorods, whiskers and nanotubes with a decrease in the temperature. The FM-AFM crossover observed near 100 K is accompanied by a low temperature increase in the Curie constant by a factor of 2.7-7. The comparison of the experimental data for various VO<sub>x</sub> nanoparticles indicates that the most probable cause of the change in the type of magnetic interaction is a change in the concentration of V<sup>4+</sup> magnetic ions due to temperature dependent electron localization [1].

1. S.V.Demishev et al., JETP Letters, 91, 11 (2010)

### EE01

#### Spin Hall effect from first principles

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Spintronics without magnetic materials would be an interesting alternative to the existing spintronics applications. The spin Hall effect creates spin currents in nonmagnetic materials and avoids the problem of spin injection. Future applications of the spin Hall effect require two properties of the materials used in the spin Hall device, a large spin Hall angle and a long spin diffusion length. Ab intio calculations based on density functional theory are a powerful tool to design the desired materials and to get insight into the underlying microscopic processes. We investigated the spin Hall effect in dilute alloys, in particular the intrinsic effect based on the Berry curvature as well as the skew scattering contribution. By means of a relativistic Korringa-Kohn-Rostoker method we calculated the scattering potentials of several impurities in Cu, Ag, Au, and Pt hosts self-consistently. We solve the linearized Boltzmann equation to calculate the conductivity tensor taking as well as the scattering-out and the scattering-in term into account. We identify systems, e.g. Cu(Bi) and Au(C) with a Hall angle larger than 0.06 and a spin diffusion length of about 100 nm.

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### EE02

### Spin transfer torques in magnetic bilayers with strong spin orbit coupling

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Current driven magnetic dynamics in ferromagnetic thin films on top of non-magnetic films with strong spin orbit coupling combines several topics of significant recent interest. The current flowing through the ferromagnetic layer gives rise to adiabatic and non-adiabatic spin transfer torques when the magnetization is non-uniform. The current flowing through the non-magnetic layer gives rise to a spin Hall current, leading to a spin current incident on the interface between the two layers. This spin current causes spin transfer torques similar to those that are important in magnetic multilayers with current flowing perpendicular to the plane. Recent experiments, models, and simulations raise the possibility of an additional torque due to the spin-orbit coupling at the interface where the inversion symmetry found in the bulk materials is broken. All of the currents and torques are coupled to each other and are non-uniform in the layers. We compute this coupling and spatial variation using a Boltzmann equation approach. We determine the dependence of the torque on the thickness of the layers and study how the torque changes as the contribution from the interfacial spin orbit coupling is increased.

### EE03

### Diffusive spin dynamics in ferromagnetic thin films with a Rashba interaction

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Injection of a charge current into an asymmetrically sandwiched thin magnetic layer generates an effective Rashba field that reverses the magnetization [1, 2] and improves current-driven domain wall motion [3]. In a ferromagnetic metal layer, we study theoretically the coupled charge and spin diffusion equations in the presence of both Rashba spin-orbit interaction and magnetism. The mis-alignment between the magnetization and the non-equilibrium spin density induced by the Rashba field gives rise to Rashba spin torque acting on the ferromagnetic order parameter. Using Keldysh formalism and gradient expansion, we find that the Rashba torque consists of both in-plane and out-of-plane components [4]. This provides an explanation to the magnetization switching mechanism in a single ferromagnet as observed in the recent experiments [5]. While providing a coherent framework to describe the diffusive spin dynamics in magnetic media with or without Rashba spin-orbit interaction, we also show particularly that the in-plane Rashba torque can be tuned by varying the width of the magnetic nano-wire and its materials parameters.

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### EE04

### Emergence of magnetic monopoles in magnetic systems with spin-orbit coupling

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In spintronics, a spin current plays an essential role. Today, various spin-related phenomena have been discovered and understood through the spin current. However, the spin current is not conserved, and the relation between spintronics and electromagnetism is still an open question. For overcoming this difficulty, we theoretically analyze the spintornic phenomena from the viewpoint of the Maxwell's equations. In particular, we focus on the electron transport driven by a dynamic magnetization or an electric field in the presence of spin-orbit interactions. We calculate an electric current and density, and derive the Maxwell's equations constructed by this way involve a magnetic monopole contribution [1]. We will discuss this magnetic monopole generated by the magnetization dynamics or the applied electric field in a ferromagnetic-nonmagnetic junction, and we will discuss also the magnetic monopole signal is expected in a similar geometry as the inverse spin Hall signal [2,3].

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### EE05

Generalization of Gilbert damping in Rashba systems

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Mutual interaction between spin current and magnetization dynamics is described by spin-transfer torque (STT) [1] and spin motive force (SMF) [2]. Spin-transfer torque results in magnetization dynamics due to spin current. In contrast, SMF is spin dependent motive force induced by magnetization dynamics. As a result of combination of these two, magnetization dynamics affects magnetization dynamics itself mediated by spin current. This self-feedback of magnetization dynamics give almost ignored since it is negligibly small. However, in an ultrathin magnetic system with strong Rashba spin-orbit coupling (RSOC) [4], which is of recent interest in spintronics, we demonstrate that SMF can be orders of magnitude enhanced so the selffeedback becomes now relevant. We derive the generalized Gilbert damping tensor including RSOC-induced SMF and discuss its effects on magnetization dynamics qualitatively. Our work implies that ultrathin magnetic films are not mere thin limits of thicker ferromagnets but systems governed by qualitatively different physics.

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### EF01

### Superconducting, antiferroquadrupolar, and structural transitions in caged compounds $PrT_2Zn_{20}$ (T=Ru, Rh, and Ir)

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There has been considerable interest in praseodymium-based intermetallic compounds, because an abundance of unusual phenomena has been intensively studied such as ferromagnetic order of hyperfine-enhanced nuclearspins in PNb[11], incommensurate antiferroquadrupolar order in PPbb[22] and heavy-fermion superconductivity in filled-skutterudite PrOs<sub>3</sub>Sb<sub>1</sub>[3]. We have recently focused on Pr-based compounds with caged structures, where large coordination number of the Pr ion leads to weak crystalline electric field effect, whereas hybridization of the f electrons with electrons of cage atoms is strengthened. Such conditions may enable the compounds to show strongly correlated electronic phenomena. Keeping this in mind, we have synthesized and studied a new family of caged compounds PrT\_Zn<sub>20</sub> (T=Ru, Rh, and Ir). PrRh<sub>2</sub>Zn<sub>20</sub> and PrIr\_Zn<sub>20</sub> with a nonmagnetic gamma3 doublet ground state undergo superconducting transitions at TSC=0.0 6K and 0.05 K, respectively, in the presence of antiferroquadrupolar order below TQ=0.06 K and 0.11 K[4,5]. The entropy releases at TQ are 6% and 20% of Rh<sub>2</sub> for PrRh<sub>2</sub>Zn<sub>20</sub> and PrIr\_Zn<sub>20</sub>, respectively, suggesting interplay between the quadrupolar fluctuations and formation of the superconducting pair. On the other hand, in PrRu<sub>2</sub>Zn<sub>20</sub>, the symmetry of the Pr ion is lowered by the structural phase transition at TS=138 K. Therefore, no phase transition manifests itself down to 0.04K

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### EF02

### Structural and magnetic phase separation in $PrMn_2Ge_{2\cdot x}Si_x$ and related compounds

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We report a comprehensive study of the thermal and magnetic properties of PrMn<sub>2</sub>Ge<sub>2.8</sub>Si<sub>x</sub> compounds along with neutron and synchrotron X-ray diffraction studies of their atomic and magnetic structure. Generally, the Mn atoms display the same antiferromagnetic configurations as the broad family of RMn<sub>2</sub>X<sub>2</sub> compounds (R = rare-earth, X = Si<sub>x</sub> Ge). However, between x= 1.0 and 1.4 we see phase separation at low temperatures, into Fmc and AFmc forms. This is caused by spontaneous magnetostriction which differs greatly between the two magnetic phases, and is linked to a subtle structural phase separation over a broader concentration range (from x= 0.4 to 1.6). Clearly two separate & unique structures exist depending on the local concentration of Silicon and Germanium on the mixed (X) site. Our interpretation brings into question whether a random substitution on the X site could produce such remarkable magneto-elastic phenomena or whether atomic short range order plays a significant role in the physical properties of the mixed 122 compounds. Indeed we find ample evidence in the broader family of RMn<sub>2</sub>X<sub>2</sub> compounds to support a claim that systematic variation in local order of the R & X sites drives the complicated magnetic and physical properties observed in this system

### EF03

Observations of magnetic and ferroelastic nanoclusters in RCo2

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New findings of short-range order (SRO) correlations in magnetic materials have recently been reported, even on compounds previously studied for decades, as a result of the use of new and improved experimental techniques. In particular, X-ray Magnetic Circular Dichroism (XMCD) and transverse susceptibility have revealed the existence of several new magnetic configurations in the phase diagram of the family of compounds RCo<sub>2</sub> (R = Er, Ho, Tb, Tm). The presence of magnetic SRO (magnetic nanoclusters) has been shown to be a prerequisite for the existence of these configurations. Relaxation of the magnetic nanoclusters leads to anomalies in the ac susceptibility, observed within the paramagnetic phase of RCo<sub>2</sub> compounds. Additionally, Resonant Ultrasound Spectroscopy (RUS) studies have shown the presence of fluctuations in local strain, also described as ferroelastic clusters, above the magnetic ordering temperature. Comparing RUS and magnetic masurements sheds light on the interrelationship between ferroelastic and magnetic nanoclusters.

### EF04

### First-principles molecular dynamics study on the magnetic structure of Mn<sub>3</sub>Pt

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The ordered Mn<sub>3</sub>Pt is a three-dimensional frustrated magnet with an octahedral configuration of Mn local magnetic moments, and exhibits two ordered phases below the Neel temperature: collinear antiferromagnetic (AF) phase for 400K < T < 475K, and noncollinear AF phase for T < 400 K. To investigate the origin and details of the magnetic structures of Mn<sub>3</sub>Pt, we have applied the first-principles molecular dynamics (MD) approach to this alloy. The theory is based on the first-principles TB-LMTO Hamiltonian, the functional integral method and the isothermal MD technique, and can determine automatically the magnetic structures of itinerant magnets at finite temperatures. Numerical results at 25 K using 4\*4\*4 FCC unit cell reveal the triangular shaped noncollinear magnetic structure, being consistent with experiment. It is found that with increasing temperature up to 300 K, the amplitude and the thermal average of Mn local moment decrease gradually, and the d-electron density of states (DOS) near the Fermi energy increases. Based on the numerical results for further increasing temperature above 400 K, a discussion is given as to the relationship between frustration, the d-electron DOS, the magnetic structures, and their temperature dependence.

#### EF05

### One-dimentional magnetism in metallic MnB<sub>4</sub>

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We have investigated from first-principles an electronic structure and magnetism in MnB<sub>4</sub> compound with experimentally observed orthorhombic C<sub>12</sub>/m1 structure. It is found that Mn tetra-borides (MnB<sub>4</sub>) is found to have metallic ground state with well defined local Mn magnetic moments. This conclusion was drawn from calculation within full potential Linear Augmented Plane Wave method and Korringa-Kohn-Rostocker method using Disordered Local Moment Approximation. We have shown using Lichtenstein Green function method for first-principle calculation of the exchange constants that magnetic exchange interactions between Mn moments are strongly ferromagnetic along 1D-chains of Mn atoms and they are practically vanishing between the chains. The metallic state appears to exhibit a strongly one-dimensional lattice dimension. Thus it is appears that MnB<sub>4</sub> is a novel one-dimensional metallic magnet.

### EG01

### Negative spin current polarization in amorphous CoFeB measured via the spin-wave Doppler effect

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We use the spin-wave Doppler technique[1-4] to measure and compare the spin polarization of the current (P) of  $Co_{72}Fe_{18}B_{10}$  before and after thermal annealing. The technique involves inductively launching and detecting spin waves in ferromagnetic wirrs carrying electrical current. The adiabatic spin transfer torque produces a transmission frequency Doppler shift of  $\Delta f = v/\lambda$  where  $\lambda$  is the spin wave wavelength and  $v = J_{\mu}BP(e, Ms)$  is the spin drift velocity, which is a function of the current density J, and the magnetization Ms. We find that the as-deposited amorphous film has a small negative value of P, in contrast to the positive values obtained via point-contact Andreev reflection (PCAR)[5]. Measurements on several other CoFeB samples with different composition also yield small or slightly negative values of P. On the other hand, the annealed film has a positive value of P, in agreement with other techniques[5,6], and as measured for  $Co_{80}Fe_{50}$  without boron[3]. Relatively little has been reported about the current polarization properties of amorphous CoFeB. The spin wave Doppler technique, with its immunity to interfacial effects, yields bulk polarization results.

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### EG02

### Switching the conductance of a magnetostrictive nanocontact by magnetic field

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The electrical conductance G of magnetostrictive nanocontacts fabricated from Dy break junctions has been investigated at 4.2 K where Dy is in the ferromagnetic state. In addition to the variation of the conductance while breaking the wire mechanically, the conductance can be changed by variation of an applied magnetic field H due to the large magnetostriction of Dy. The behavior of G(H) and its angular dependence can be explained by taking into account the magnetostrictive properties of Dy. This demonstrates the possibility of tuning the conductance of magnetic field on the contact shape created during mechanical training and find that the shape can be controlled by the applied field due to the strong magnetoselastic coupling of Dy.

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### EG03

### Spin wave and spin pumping in permalloy strips

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Microfabrication techniques have made it possible to characterize both the resonant as well as the travelling characteristics of spin waves in permalloy [1, 2]. Among variety of methods for measuring spin waves, pulse inductive microwave magnetometry [3-5] is the electrical technique for travelling spin waves. We investigate the effect of the pulse width of excitation pulses on the generated spin wave packets using both experimental results and micromagnetic simulations. We show that spin wave packets generated from electrical pulses are a superposition of two separate spin wave packets, one generated from the rising edge and the other from the falling edge, which interfere either constructively or destructively with one another, depending upon the magnitude and direction of the field bias conditions. A method of spin wave sresulting from two closely spaced voltage impulses for the modulation of the magnitude of the resultant spin wave packets. We also investigate the contribution of the surface spin waves to spin pumping. We show that spin waves can be used as a very efficient method for increasing the spin pumping signal compared to the ferromagnetic resonance case.

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### EG04

### Clarification of oxygen impurity effect on NCMR with the film resistivity and bulk scattering spin asymmetry for [FeCo/Natural Oxidation] multi-layers

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Nano-Contacts Magnetoresistance (NCMR) with Fe<sub>0.5</sub>Co<sub>0.5</sub> AlOx-NOL (Nano-Oxide-Layer) has been reported to show MR ratio of more 30 % at RA of 0.3  $\Omega\mu$ m2[1]. However, NCMR-SV is predicted to have potentially huge MR ratio of more than 100 % using high bulk spin scattering asymmetry  $\beta$  of ferromagnetic material Fe<sub>0.5</sub>Co<sub>0.5</sub> with 0.81[2]. It was surmised this difference of MR ratio between experiments and theory is caused by the impurity of oxygen in FeCo NCs. we investigated the quantitative evaluation of oxygen in asymmetry for [FeCo/Natural Oxidation] multi-layers. We estimated the resistivity and  $\beta$  of FeCo(O) with [FeCo 2 nm / N.O.]n were 37.4  $\mu$ Ocm and 0.75, and those of FeCo(O) with [FeCo 1 nm / N.O.]n were obtained 76.7  $\mu$ Ocm and 0.35. These results suggest the increase of the resistivity by oxygen impurity decreases the  $\beta$ . On the other hand, the resistivity of the single NC in our NCMR was estimated 60  $\mu$ Ocm by conductive-AFM. The cause of low NCMR ratio is considered to be the impurity of oxygen in NCs. So, reducing oxygen impurity in NC leads to higher NCMR ratio.

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### EG05

#### Room-temperature magnetoresistance properties of planar-type Ni nanostructures controlled from nanoconstrictions to nanogaps Jun Kitagawa, Ryutaro Suda and Jun-ichi Shirakashi\*

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Innovative devices called ferromagnetic single-electron transistors (FMSETs), which are based on the interplay of single-electron charging phenomena and spin-dependent tunneling effects, have attracted much interest. For the realization of FMSETs, we have reported the magnetoresistance (MR) properties of planar-type tunnel junctions with ferromagnetic nanogap system at 16 K, formed by electromigration methods. [1, 2]. In this paper, we study the room-temperature MR properties of planar-type Ni nanostructures controlled from nanoconstrictions to nanogaps using novel electromigration method, with the aim of fabricating FMSET with higher operation temperature. The stepwise feedback-controlled electromigration (SFCE) methods [3] were performed at room temperature in a vacuum chamber for Ni initial nanochannels with a constriction of a few hundreds of nanometers. We obtained anisotropic MR (AMR) of 0.14 % in the Ni initial nanochannel before performing SFCE procedure. Then, the resistance of the Ni nanochannel was gradually increased from 60  $\Omega$  to 76 k $\Omega$  by SFCE approach with atomic scale controlled biltructures at room temperature. These results strongly suggest the possibilities of room temperature operation of planar-type FMSETs with Ni/vacuum/Ni ferromagnetic tunnel junctions.

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### EH01

### New materials for enhancing device perpormance in spintronics

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Spin current creates significant effects such as tunneling magnetoresitance (TMR). current perpendicular to plane magnetoresistance (CPP-GMR) and spin transfer torque (STT), which are key factors in spintronics. They are anticipated to be applied for spintronic devices such as ultrahigh density HDD exceeding Tb/in2, STT-MRAM, high frequency oscillators, spin logics etc. Perfectly spin-polarized half-metals are essential materials for enhancing the polarization of spin currents and the resultant device performance in spintronics. Full-Heusler allys with L21 structure have been practically expected as half-metals, because they have a large Curie temperature. and their Fermi level can be tuned by the composition. We have demonstrated that Fermi-level-tuned Co<sub>2</sub>FeAl<sub>0.5</sub>Si<sub>0.5</sub> Heusler alloy provides half-metallicity even at room temperature [1]. MgO barrier has been used in MTJs because it exhibits a large TMR due to the coherent tunneling in enitaxials MTIs. Very recently we have discovered a new tunnel barrier, spinel (MgAl2O4), bringing a large TMR due to coherent tunneling as MgO barrier [2]. Spinel barrier has a benefit of very good lattice matching with bcc metals and Heusler alloys. In this talk we will show half-metallic full-Heusler alloys and their applications to MTJs with in-plane and out-of plane magnetizations [3], and MTIs with an MgAl<sub>2</sub>O<sub>4</sub> harrier

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#### EH02

Organic high temperature ferromagnetic compositions

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Developments of the high temperature organic ferromagnetic materials have been attempted by many researchers. To achieve the ferromagnetic ordering by using organic materials, the high spin multiplicity polymers are also proposed by Itoh et al. [1] in 1968, because they may realize the magnetic storage devices with a versatile processability. Recently, we found that high temperature ferromagnetism could be achieved by intercalating an organic discotic liquid crystalline (DLC) compound with an iron-phthalocyanine [Fe(III)Pc] complex[2] and also with stable organic fire radicals. This discovery is originated from our earlier studies[3-5] on the magnetic properties of DNA. We find an analogy between the ferromagnetic properties of inorganic dilute magnetic semiconductors and DNA and DLC compositions. Herein, we will discuss the high temperature ferromagnetic properties of DLC compositions doped with Fe(III)Pc complexes and also with organic fire radicals. We also will discuss a preliminary result on the room temperature ferromagnetic properties of polymeric composites consisting of some of polyconjugated polymers and paramagnetic dopants. We studied in details their magnetic properties by electron paramagnetic resonance and SQUID studies. These findings are expected to open a new research field as well as high technological applications of organic materials.

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#### EH03

## Structure and magnetic properties of the new ferrimagnetic AFe<sub>3</sub>O(PO<sub>4</sub>)<sub>3</sub> (A=Ca,Sr,Cd,Pb) compounds

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The crystal structure of the novel monoclinic oxyphosphate PbFe<sub>3</sub>O(PO<sub>4</sub>)<sub>3</sub> was characterized by single crystal (SC) XRD between 293 and 973 K.  $\chi$  measurements performed on SCs unveiled an unusual sequence of ferromagnetic-like divergences at Tc1=31.8 K, Tc2=23.4 K and Tc3-10 K, whereas Cp measurements performed on SCs confirmed only the two phase transitions at Tc1 and Tc3, and a slight cusp around Tc2. In powder samples however, the second  $\chi$  divergence at Tc2 is absent, but they do exhibit a  $\lambda$ -type peak on the Cp(T) curve. A set of  $\chi$ ac measurements, (logarithmic) long time remanent magnetization decay measurements, modified ZFC magnetization measurements, were performed on both powders and SCs. Mean-field critical exponents were obtained in powders and SCs above Tc1, consistently with the low development of magnetic correlations observed in  $\chi$  curves, but not with the asymmetric  $\lambda$ -type Cp peak. We have clear evidence for two dissipative mechanisms, giving peaks at 27 and 13 K on  $\chi$ ac measurements. While the former spectrum shows a Debye profile with a well defined Eact-240 K, the latter one shows a distribution of Eact typical of spin glass systems. Finally, FCZ/EC magnetization cycles and  $\chi$  measurements in the AFe<sub>2</sub>O(PO<sub>4</sub>)<sub>2</sub> (A=Ca,St;Cd,Pb) compounds will be presented.

#### **EH04**

## Magnonic metamaterials formed by arrays of Co antidots on continuous NiFe films

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Magnonic crystals periodic magnetic media represent one of the central objects in the research field of magnonics [1]. Usually, the periodic modulation of the effective magnetic field in magnonic crystals is achieved as a compositional modulation created either during growth or using lithographic processing. In this report, we demonstrate a method of producing magnonic crystals by periodically modulating the internal magnetic field inside a continuous thin film of NiFe (Permalloy). The modulation of the internal magnetic field is achieved via fabrication of a Co nanostructure (an array of antidots) on top of the Permalloy film, with the latter experiencing the stray magneto-dipole field from uncompensated magnetic charges in Co. The dynamical properties of the fabricated magnonic crystals are characterized using the time resolved scanning Kerr microscopy (TRSKM) [2]. The measurements reveal a strong modification of the magnonic spectrum the strength of which increases as the value of the bias magnetic field decreases. The finding opens new ways for creation of magnonic metamaterials. The research leading to these results has received funding from the EC's 7th Framework Programme (FP7/2007-2013) under GA 228673 (MAGNONICS) and from EPSRC of the UK under project EP/E055087/1.

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#### EH05

Formation of FeSi thin films and magnetic properties

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The deposition of transition metal layers on silicon and their reaction with substrate are important issues in semiconductor device technology. The interface between metal and semiconductor determines the device performance. whereby both an ohmic contact and a Schottky barrier, i.e., an active electronic element can be realized. The 3d transition metal monosilicides such as FeSi, CoSi, MnSi and CrSi have attracted much attentions because they are easily formed in the interface between transition metal and Si. FeSi is a narrow band gap semiconductor. The activation energy for exciting electrons from the filled band to the empty band is estimated to be 0.05 eV from the experiments on the temperature dependence of the Mossbauer effect, magnetic susceptibility and NMR shift. In this work, Fe thin films were grown on Si(100) substrates using MBE. The structure and surface morphologies of Fe films were determined by X-ray diffraction and atomic force microscopy measurements, respectively. We have observed that all samples grown at 100, 400 and 600 oc have a phase of CsCl crystal structure (c-FeSi). In order to investigate the correlation between magnetization and charge carrier transport, we performed magnetoresistance and Hall resistance measurements by using a physical property measurement system.

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#### **EI01**

#### Perpendicular magnetic anisotropy in Fe/Fe<sub>1-x</sub>Co<sub>x</sub> multilayers

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A uniaxial anisotropy of a volume character that can keep perpendicular magnetization in relatively thick films is rare. The concept relates to the lattice distortion that modifies the electronic band structure. A model system is provided by Fe<sub>1x</sub>Co<sub>x</sub> alloy films tetragonally distorted due to their pseudomorphic growth on fcc mismatching substrates like Rh(001) and Ir(001) [1]. In particular, the Fe<sub>0x</sub>Co<sub>0x</sub> films show perpendicular magnetic anisotropy, whereas Fe layers (x=0) are magnetized in-plane. The Fe<sub>1x</sub>Co<sub>x</sub> and Fe layers can be separated by a spacer that mediates a ferro- or antiferromagnetic exchange coupling [2]. As MOKE measurements provide only a signal from the average magnetization, a layer resolved knowledge of the spin structure is required. Soft x-ray magnetic circular dichroism allows us to probe the magnetization separately for Fe<sub>1x</sub>Co<sub>x</sub> (measurements at the Co edges enabled us to investigate the magnetization configurations can be produced also if there is no spacer, but there is a direct exchange coupling between the Fe and Fe<sub>1x</sub>Co<sub>x</sub> layers. A final magnetization distribution depends on the composition of the constituent layers, and how thick Fe is with respect to Fe<sub>1x</sub>Co<sub>x</sub>.

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#### EI02

## Effect of annealing temperature on L1o ordering and perpendicular magnetic anisotropy of FePd/CoFeB films

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Spin-transfer-torque random access memory (STT-RAM) using perpendicular magnetized L1o-ordered alloy for magnetic tunnel junctions (p-MTJs) has considerable interest because of advantages for scalability and stability of memory. However, tunnel magnetoresistance (TMR) ratio was a few percent at room temperature using L1o-ordered alloys. [1,2] A high TMR ratio as well as perpendicular magnetic anisotropy (PMA) was obtained using the thin CoFeB combined with MgO barrier due to interfacial effect. [3] In this study, we systematically investigated annealing temperature effects on PMA with the insertion of thin CoFeB layer between L1o-FePd and MgO films for use as free layer in the p-MTJs. Stacking structure is  $Cr_{40}/Pd_{10}/FePd_/CoFeB_0/MgO_2/Ta_3$  (unit in nm, and t=1.5-5.0) and these were epitaxially grown on MgO (100) substrates. FePd and CoFeB were deposited at 300°C and ambient temperature, respectively. After post-annealing, the perpendicular anisotropy was observed at Ta≥ 300°C and 250°C for t= 2.0 nm and 2.5 nm, respectively, whereas it was observed for all samples above t≥3.0 nm which might be attributed to the crystallization of CoFeB and increased of PMA. The results show that annealing is an effective method to improve the PMA of a FePd/CoFeB system. Acknowledgement: This study was supported by FIRST program from JSPS.

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#### EI03

**Magnetic properties of tetragonally strained Fe/(W,Re) multilayers** Cristina Bran<sup>1</sup>, Matthias Hudl<sup>2</sup>, Matts Bjorck<sup>1</sup>, Vassilios Kapaklis<sup>1</sup> and Gabriella Andersson<sup>1</sup>\*

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Following recent theoretical predictions on perpendicular uniaxial magnetocrystalline anisotropy in multilayers of 5 monolayers of Fe and 2 monolayers of (W,Re) with tetragonally strained components [1], we have investigated the magnetization and detailed atomic structure in a series of samples with 20-80 at% W. Experimental techniques used were SQUID magnetometry, Magneto-Optic Kerr Effect, Energy Dispersive X-ray Spectroscopy, and X-ray Diffraction including reciprocal space mapping around (002) and (103) body-centered-tetragonal Bragg peaks to measure inplane atomic distances. Although all the samples exhibit considerably smaller strain in the (W,Re) alloy layers than the previously modelled multilayers, and consequently in-plane easy magnetization axes, we do observe an enhancement of the saturation magnetization that corresponds well to the predicted hybridization-induced moments on individual Fe and (W,Re) atomic layers.

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#### EI04

## Alloying as a possible mechanism in annealing induced perpendicular magnetic anisotropy in alumina/Co/M (where M=Pd,Pt or Au) trilayers

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Since the work of Carcia[1], perpendicular magnetic anisotropy (PMA) in ferromagnetic/nonmagnetic trilayers or multilayers has attracted much attention. Spin-orbit interaction is the main cause for PMA. Recently it has been shown that in MgO/CoFeB/Ta<sub>2</sub> and Oxide/Co/Pt<sub>5</sub>, PMA can be induced by annealing. In particular, it has been argued in [4] that the origin of this PMA is the hybridization between Co and O orbitals enhanced by annealing. We demonstrate in this work that another explanation is alloying between cobalt and the capping layer. We have grown and annealed at different temperatures trilayers of alumina(3nm)/Co(tnm)/M(3nm) where t varies from 1 to 2 nm and M= Pd, Pt and Au. For Ta=300°C, PMA is observed in Co/Pd and Co/Pt for t<1.6 nm and 1.4 nm respectively. However no PMA is observed for any thickness in Co/Au. At Ta= 400 °C Co/Pd and Co/Pt show PMA for all thicknesses whereas for Co/Au no PMA is observed. It seems that the influence of Co-O bonds is not enough to drive the magnetization out-of-plane in Co/Au. We propose that the observed effect is due to alloying between Co and Pt or Pd, as Co and Au are immiscible no effect is observed in Co/Au.

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#### EI05

## High perpendicular magnetic anisotropy at $Co_xNi_{1,x}(x=0.0{\sim}1.0)/$ $\alpha\text{-}Cr_2O_3$ interface

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α-Cr<sub>2</sub>O<sub>3</sub>(0001) surface offers a high perpendicular exchange bias when it is attached to a ferromagnetic layer [1] because the topmost Cr spins on the α-Cr<sub>2</sub>O<sub>3</sub>(0001) surface are fully uncompensated and align normal to the surface. In particular to the perpendicular exchange bias, the perpendicular magnetic anisotropy (PMA) at a ferromagnetic(FM)/antiferromagnetic(AFM) interface is essential, which has not been reported. In this work, we have investigated the interface magnetic anisotropy of the  $Co_x Ni_{1-x}(111)/\alpha \cdot Cr_2O_3(0001)$  interface with the Co composition x = 0-100 at %. From the structural characterizations, we confirmed that the fabricated samples were in the epitaxial manner, i.e. Pt(111) || Co<sub>x</sub>Ni<sub>1-x</sub>(111) || α-Cr<sub>2</sub>O<sub>3</sub>(0001) || Pt(111) || α-Al<sub>2</sub>O<sub>3</sub>(0001)-sub. We precisely determined the PMA at the Co(111)/ $\alpha$ -Cr<sub>2</sub>O<sub>3</sub>(0001) interface through the change in the PMA energy with the number of stacking-period of Pt/[Co/Pt]n/α-Cr<sub>2</sub>O<sub>3</sub>/Pt epitaxial superlattices. The PMA energy increases with increasing the number of stacking-period, which indicates that the strong PMA is generated at the Co(111)/ $\alpha$ -Cr<sub>2</sub>O<sub>3</sub>(0001) interface and that the PMA energy of the Co(111)/a-Cr2O3(0001) interface is comparable to that of the Co/Pt interface. For the Pt/ Co<sub>x</sub>Ni<sub>1-x</sub>(111)/a-Cr<sub>2</sub>O<sub>3</sub>(0001)/Pt films, the PMA energy monotonically decreases with increasing Ni composition. The change in the PMA with the Ni composition is similar to that of Pt/Co,Ni1,/Pt multilaver [2].

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#### EI06

## Controlling domain wall motion by electric fields in perpendicularly magnetized materials

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Magnetic domain walls (DWs) are excellent candidates to store and transport information in future memory devices [1]. Up until now, control of DWs has only been possible through the application of a magnetic field or current. Here, we demonstrate a novel approach to control DW motion, namely by altering the magnetic properties with an electric field [2].  $Pt/Co/AIO_x$  wires exhibiting perpendicular magnetic anisotropy (PMA) are fabricated to investigate the influence of the electric field on DW propagation. By applying a voltage over the insulating  $AIO_x$  barrier charges are induced at the Co/AIO<sub>x</sub> interface, locally changing the magnetic properties. Depending on the sign of the voltage, it is observed that DW velocities can be increased or decreased up to an order of magnitude. By analyzing the experimental results with the DW creep law, we were able to attribute the voltage-induced changes in DW motion to the sensitivity of PMA to an electric field. Comparing the obtained changes in PMA to values in the literature [2], we estimate that controlling DW motion over five orders of magnitude should be well within experimental reach.

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#### EJ01

## Evaluation of interlayer exchange coupling in $\alpha$ -Fe(100)/Nd<sub>2</sub>Fe<sub>14</sub>B(001) Films

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Exchange-coupled hard/soft nanocomposite magnets have a potential to exhibit a large maximum energy product[1]. Recently, a new issue has been proposed by the first-principle calculations[2] for the  $\alpha$ -Fe/Nd<sub>2</sub>Fe<sub>4</sub>B system, which suggest that a magnitude of the hard/soft exchange-coupling depends strongly on the Miller indices of the crystallographic planes at the interface between the two phases, and that even antiparallel coupling would be stabilized in particular combination of planes. Recently, we fabricated the  $\alpha$ -Fe(100)/Nd<sub>2</sub>Fe<sub>4</sub>B(001) bilayer on the MgO(100) substrate by using UHV magnetron sputtering system. Q-band FMR measurements revealed the positive interlayer coupling in this sample [3]. In this work, similar films with different values of the  $\alpha$ -Fe layer thickness tFe were fabricated in order to evaluate the exchange-coupling constant more precisely. With increasing tFe, the shift in the  $\alpha$ -Fe resonance field decreased, as compared with the resonance field for the  $\alpha$ -Fe single layer film. These FMR spectra and the magnetization curves are found to be in consistent with the simulation assuming the positive exchange coupling in the  $\alpha$ -Fe(100)/Nd<sub>2</sub>Fe<sub>4</sub>B(001) interfaces.

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#### EJ02

**Morphology and magnetic properties of** S<sub>m</sub>Co<sub>3</sub>/Fe and Sm<sub>2</sub>Co<sub>7</sub>/FeCo nanocomposite magnets prepared via severe plastic deformation Narayan Poudyal, Nguyen Van Vuong, Ying Zhang and J. Ping Liu\* Department of Physics, University of Texas at Arlington, USA

SmCo5 and Sm2Co17 compounds have been widely applied as the first and the second generations of rare-earth permanent magnets respectively. There are other Sm-Co internmetallic phases with lower Co concentration that have been ignored as permanent magnetic materials for their relatively low magnetization and anisotropy. Those phases, including SmCo3 and Sm2Co7, have been systematically investigated in our recent research as the hard-phase component in hard/soft nanocomposite systems. It has been found that the nanocomposites magnets having SmCo3 and Sm2Co7 as the hard phases have excellent properties. The composites can incorporate more soft phase compared to SmCo<sub>5</sub> and Sm<sub>2</sub>Co<sub>17</sub>. Here, we report SmCo<sub>3</sub>/Fe and Sm<sub>2</sub>Co<sub>7</sub>/FeCo nanocomposites prepared by high energy ball-milling and subsequent heat treatments. The evolution of morphology and magnetic properties with ball milling time and the soft phase content has been studied. The a-Fe phase distributed in the matrix of hard magnetic Sm-Co phases has grain size less than 20 nm after the severe plastic deformation. Enhanced remanence and maximum energy product up to 13.5 and 17.6 MGOe in the nanocomposites with 25 and 40 % of the soft phase are obtained for SmCo3 and Sm2Co7 respectively, which is up to 300 % higher than that of single-phase counterparts.

#### EJ03

## Atomic scale investigation of Sm-Co/Fe nanocomposites: Influence of Fe/Co interdiffusion on the magnetic properties

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The nanostructure of SmCo<sub>3</sub>/Fe nanocomposites elaborated by high energy ball milling was investigated using atom probe tomography. This high resolution technique allowed obtaining in the real space a three dimensional mapping of the chemical species (Sm, Co and Fe) at the atomic scale [1]. 3D reconstructions of the nanostructure of the analysed samples were realised from the data and show that the majority of the SmCo<sub>3</sub>/Fe nanocomposite is composed of Fe(Co) soft magnetic nano-sized crystallites embedded in a Sm-(Co<sub>3</sub>Fe) hard magnetic matrix. These results confirm the formation of the Fe(Co) phase and the contamination of the initial SmCo<sub>3</sub> by Fe atoms. The interfaces between the Fe(Co) soft magnetic phase and the Sm-(Co<sub>3</sub>Fe) hard magnetic phase were also chemically analysed at the atomic scale. Large graded interfaces were evidenced. These observations were correlated with Monte Carlo simulations and confirm the benefit of the Fe/Co interdiffusion on the magnetic exchance coupling.

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#### EJ04

#### Effect of particle size on the coercivity of Nd-Fe-B and Sm-Co nanoparticles prepared by surfactant-assisted ball milling

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Anisotropic Nd<sub>2</sub>Fe<sub>14</sub>B and SmCo<sub>5</sub> nanoparticles have been produced by surfactant-assisted high-energy ball milling (HEBM) of nanocrystalline precursor alloys. A two-stage HEBM was performed to obtain the nanoparticles; first the coarse powders were made nanocrystalline by milling, and then were subjected to surfactant-assisted milling for different time. Nanoparticles with different shape (irregular, square and spherical shapes) and size have been obtained by varying the grain size of nanocrystalline precursors. The size of the nanoparticles is found to depend strongly on the grain size of nanocrystalline precursors. For a surfactant-free wet milling of 0.5 to 4 hr, anisotropic Nd<sub>2</sub>Fe<sub>14</sub>B and SmCo<sub>5</sub> nanoparticles have been found with a size from 25 to 2.8 nm. The nanoparticles had coercivity values at 50 K of 9 kOe for the 15 nm Nd<sub>2</sub>Fe<sub>14</sub>B and 15 kOe for the 25 nm SmCo<sub>5</sub>. The room temperature coercivities were lower with the values of 2.5 and 5 kOe, respectively. The coercivity was found to decrease with the particle size suggesting a decrease in the overall anisotropy because of surface effects. In the Nd<sub>2</sub>Fe<sub>14</sub>B nanoparticles, a lower spin reorientation temperature was observed (117 K) as compared to bulk (135 K).

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#### FA01

#### Dynamical generation of spin currents

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Spin current, a flow of electrons' spins in a solid, is the key concept in spintronics that will allow the achievement of efficient magnetic memories, computing devices, and energy converters. I here review phenomena which allow us to use spin currents in insulators [1]: inverse spin-Hall effect [2,4], spin pumping, and spin Seebeck effect [4-6]. We found that spin pumping and spin torque effects appear at an interface between an insulator YIG and Pt. Using this effect, we can connect a spin current carried by conduction electrons and a spin-wave spin current flowing in insulators. We demonstrate electric signal transmission by using these effects and interconversion of the spin currents [1]. Seebeck effect (SSE) is the thermal spin pumping [5]. The SSE allows us to generate spin voltage, potential for driving nonequilibrium spin currents, by placing a ferromagnet in a temperature gradient. Using the inverse spin-Hall effect in Pt films, we measured the spin voltage generated from a temperature gradient in various ferromagnetic insulators. This research is collaboration with K. Ando, K. Uchida, Y. Kajiwara, S. Maekawa, G. E. W. Bauer, S. Takahashi, and J. Ieda.

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#### FA02

#### Domain wall motion by the magnonic spin Seebeck effect

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In the last years it was demonstrated that in ferromagnetic materials spatial temperature grandients can lead to spin accumulation [1], even in magnetic insulators [2]. In this context, it is important to note that in addition to a spin polarized charge current also a chargeless angular momentum current driven by spin waves can exist. Here, we propose the existence of domain wall (DW) dynamics driven by magnonic spin currents due to temperature gradients. To get some insight into this effect we introduce two different approaches: the stochastic Landau-Lifshitz-Gilbert equation, applied to spin models, and the Landau-Lifshitz-Bloch equation describing the dynamics of the thermally averaged spin polarization on micromagnetic length scales. We show that these approaches describe this new type of DW motion, where chargeless spin currents following from a temperature gradient drag a DW into the hotter region [3]. Furthermore, for a better understanding of the relevant length scales, we investigate the propagation and frequency range of thermally induced magnons. We acknowledge financial support by the Deutsche Forschungsgemeinschaft through Schwerpunktprogramm SpinCaT.

K. Uchida et al., Nature 455, 778 (2008).
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#### FA03

#### Phonon-drag spin Seebeck effect

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Spin Seebeck effect refers to a thermal spin injection from a ferromagnet into an attached nonmagnetic metal that occurs over macroscopic scale of several millimeters [Uchida et al., Nature 455, 778 (2008)]. The spin Seebeck effect is now established as a universal aspect of ferromagnet since it is observed in a variety of materials ranging from a metallic ferromagnet (NiFe) and semiconducting ferromagnet (GaMnAs) to an insulating magnet (YIG). Recent theoretical and experimental efforts have clarified that the phonon degrees of freedom are of crucial importance in this intriguing phenomenon. Here we theoretically discuss the phonon-drag contribution to the spin Seebeck effect [Adachi et al., Appl. Phys. Lett. 97, 252506 (2010)]. The spin Seebeck effect is driven by nonequilibrium phonons that drag the low-lying spin excitations. This scenario explains simultaneously the local nature of the spin Seebeck effect [Jaworski et al., Nature Materials 9, 898 (2010); Uchida et al., Nature Materials 10, 737 (2011)] and the signal enhancement at low temperatures [Jaworski et al., Phys. Rev. Lett. 106, 186601 (2011)].

#### FA04

#### **Entanglement of spin Seebeck effect and anomalous Nernst effect** Chia-ling Chien

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Spin Seebeck effect (SSE) is a spin caloritronic effect that generates a pure spin current. The studies of SSE have often been made on a patterned Pt strip on a ferromagnetic (FM) thin metallic film on substrate, where the in-plane temperature gradient ( $\bigtriangledown_x$ T) generates in the FM film a spin current, which is detected by the inverse spin Hall effect as a voltage in the Pt strip. However, there is also the propensity of an out-of-plane temperature gradient ( $\bigtriangledown_x$ T), which gives rise to the anomalous Nernst effect (ANE) due to the larger thermal conduction through the substrate. I While SSE with  $\bigtriangledown_x$ T and ANE with  $\bigtriangledown_z$ T are two different effects, they can be simultaneously present. More importantly, the voltages of the two effects are additive with the same field dependence and symmetry, and therefore cannot be easily separated. In such cases, the establishment of the SSE is contingent upon the ANE measurements before the Pt strip is in place.

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#### FB01

#### Textured superconductivity in the heavy fermion CeRhIn5

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In the coexistence region of antiferromagnetism (AFM) and superconductivity (SC) of CeRhIn<sub>5</sub>, the superconducting transition temperatures determined resistively and thermodynamically display a significant difference while the difference disappears at higher pressures where AFM is suppressed and only SC exists. In the coexisting state of AFM and SC, anisotropic transport near the SC transition reveals the emergence of textured SC planes that appear without a change in translational symmetry of the lattice. Similar behaviors have also been observed in other strongly correlated electron systems where there are coexisting and competing orders with SC.

Tuson Park et al. Phys. Rev. Lett. in press.

#### FB02

#### Exotic superconductivity of heavy electrons in artificial twodimensional Kondo lattices Takasada Shibauchi

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When interacting electrons are confined to low-dimensions, the electron-electron correlation effect is enhanced dramatically, which often drives the system into exhibiting behaviours that are otherwise highly improbable. Superconductivity with the strongest electron correlations is achieved in heavyfermion compounds, which contain a dense lattice of localized magnetic moments interacting with a sea of conduction electrons to form a three-dimensional (3D) Kondo lattice. It had remained an unanswered question whether superconductivity would persist upon effectively reducing the dimensionality of these materials from three to two. We report on the observation of superconductivity in such an ultimately strongly-correlated system of heavy electrons confined within a 2D square-lattice of Ce-atoms (2D Kondo lattice), which was realized by fabricating epitaxial superlattices[1] built of alternating layers of heavy-fermion CeCoIn, and conventional metal YbCoIn, [2]. The field-temperature phase diagram of the superlattices exhibits highly unusual behaviours, including a striking enhancement of the upper critical field relative to the transition temperature. Possible origins of this enhancement will be discussed, including extremely strong coupled nature as a result of two-dimensionalization[2], formation of FFLO state, and local spatial inversion symmetry breaking [3].

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#### FB03

## Evolution of quasiparticle entropy in high-field superconducting phase in CeCoIn.

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Unambiguous detection of Fulde and Ferrell and Larkin and Ovchinnikov (FFLO) superconducting (SC) phase in solid state systems is still a challenge in the comunity of strongly correlated electron systems (SCES). A heavy fermion superconductor, CcCoIn<sub>5</sub>, is one of the best candidates for the formation of FFLO state. Existence of FFLO phase in this material has been extensively discussed in the past decade and is still one of the most debated current topics in the field of SCES[1]. Entropy as a function of field can capture the characteristic nature of FFLO state, namely the additional increase of quasi-particles, leading to a steep increase of entropy at the transition field from uniform to FFLO SC state. We have obtained the entropy, derived from combined specific heat and magnetocaloric effect measurements at temperatures ST/geq 1008-mK and fields SH/leq 125-T aligned parallel, perpendicular and \$18%/circ\$ off the tetragonal [100] direction. In any direction, we do not observe an additional quasiparticle entropy expected in a FFLO SC state. By contrast, for SH/parallel [100]S a negative entropy contribution, compatible with spin-density-wave (SDW) ordering, is found. Our data exclude the formation of a FFLO state in CeCoIn<sub>5</sub> for out-of-plane field directions, where no SDW order exists.

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#### FB04

#### Electronic structure of ferromagnetic heavy fermion YbNi4P<sub>2</sub>

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Quantum critical phenomena can be studied in great detail in lanthanide based heavy-fermion systems. Up to now, however, only materials with antiferromagnetic ground state were found among the 4f based heavy fermions. In the new heavy fermion material YbNi<sub>4</sub>P<sub>2</sub> a clear ferromagnetic transition is observed at TC = 0.17 K. YbNi<sub>4</sub>P<sub>2</sub> has the prospect to study ferromagnetic quantum criticality accessible by either composition tuning or transverse-field tuning. Above the transition temperature in YbNi<sub>4</sub>P<sub>2</sub> strong evidence is found for its vicinity to a quantum critical point, the specific heat diverges in a power-law form and the resistivity follows a linear temperature dependence. We present first Shubnikov-de Haas measurements in combination with electronic band structure calculations. The electronic structure is dominated by parallel disconnected sheets in accordance with the crystal structure featuring quasi-one-dimensional chains of Yb. The observed Shubnikov-de Haas oscillations are assigned to quasi-two-dimensional features of the Fermi surface. The experimentally observed mass enhancement of these orbits reflects the heavy-fermion character of YbNi<sub>4</sub>P<sub>2</sub>.

#### FC01

## Models of spin dynamics; ultrafast heat pulses as a sufficient stimulus for reversal in a ferrimagnet

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The response of magnetic materials to the heating effect of ultrafast laser pulses is a challenging problem which is giving important insight into spin dynamics at elevated temperatures and sub-picosecond timescales. We have developed a model of ultrafast spin dynamics using an atomistic model based on the Heisenberg exchange Hamiltonian and using the Landau-Lifshitz-Gilbert equation with Langevin dynamics to calculate the evolution of the magnetisation following a laser pulse. Here we describe studies of the properties of amorphous ferromagnetic Rare Earth (RE) Transition Metal (TM)alloys, which are shown to exhibit interesting and unexpected properties. Firstly, the model demonstrates differential sublattices dynamics, specifically a different rate of demagnetisation for the RE and TM sub-lattices, even though these are coupled by a strong exchange interaction. Even more surprising is the existence of a transient ferromagnetic state, stable for around 300fs. It is shown theoretically that this leads to magnetisation reversal, in the absence of a magnetic field, driven by a heat pulse alone. The underlying physics of this effect and its confirmation in recent experiments [1] will be discussed.

[1] T Ostler et al., Nat. Commun., 3:666 doi: 10.1038/ncomms1666 (2012)

#### FC02

## Coherent spin-photon interaction and ultrafast magnetism: From principles to applications

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How fast can one manipulate the magnetization of materials using ultrashort optical pulses? This question is at the heart of several fundamental and applied researches with the prospective of making ultrafast magnetic devices. An important related topic is the coherent time dependent interaction between spins and photons. In the present talk we review the state of the art understandings in that field, describing the spin dynamics from first principles and showing some applications. First, we show that the spin-orbit interaction is a crucial concept when manipulating the spins with a laser field. Second, we distinguish between different mechanisms resulting from the time ordering between a sequence of pump and probe pulses when performing a time resolved magnetooptical experiment. In particular we distinguish between the coherent and population components associated to the spins dynamics. Third, we demonstrate the validity of these theoretical concepts by considering the time dependent magneto-optical response of garnets or transition metals (Ni, CoPt) excited by femtosecond laser pulses. In the case of ferrimagnetic Bi-doped garnets the coherent contribution is most important, leading to the generation of efficient magneto-optical four-wave mixing signals. In the case of ferromagnetic transition metals, the population dynamics dominates with a weaker coherent response

#### FC03

Ultrafast switching of ferrimagnets Sonke Wienholdt<sup>1</sup>, Denise Hinzke<sup>1</sup>, Peter Openeer<sup>2</sup> and Uli Nowak<sup>1</sup> Department of Physics, University of Konstanz, Germany <sup>2</sup> Department of Physics, Uppsala University, Sweden

The ultrafast manipulation of magnetisation by fs laser pulses promises to become a real alternative to conventional techniques based on magnetic fields. It was demonstrated that a 80 fs, circularly polarised laser pulse is able to reverse magnetisation on a ps time scale, as if it acts as a short magnetic field pulse caused by the inverse Faraday effect [1]. In single-shot time-resolved imaging of magnetic structures [2] it has been shown that the magnetisation reverses via a linear pathway [3] without any precession. Even with linearly polarized light switching was demonstrated recently [4], probably on a purely thermal basis. These new types of switching have been demonstrated only in ferrimagnetic materials like GdFeCo, probably because of the antiferromagnetic coupling of the two different sub-lattices in these materials, leading to completely different dynamics as compared to a ferromagnet. To understand this new type of dynamics we perform atomistic spin model simulations of ferrimagnets and investigate their switching mechanisms in detail. Funded by the EU project FentoSpin and the CAP in Konstanz.

[1] C. D. Stanciu et al., Phys. Rev. Lett. 99, 047601 (2007). [2] K. Vahaplar et al., Phys. Rev. Lett. 103, 117201 (2009) [3] N. Kazantseva et al., Europhys. Lett. 86, 27006 (2009) [4] I. Radu et al., Nature 472, 205 (2011)

#### FD01

## Spin wave mediated magnetic vortex core reversal. Towards a 100 $\mbox{ps}$ V(ortex)MRAM

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The discovery of magnetic vortex core reversal by excitation of the sub-GHz gyromode [1] not only opened a new chapter in vortex dynamics but also enabled the use of the magnetic vortex core polarization p as memory bit for spintronic applications. Suggestions for V(ortex)MRAMs published so far [2-6] are based on the excitation of the vortex gyromode at sub-GHz frequencies and allow switching times of a few ns. Last year we have demonstrated [7] that much faster vortex core reversal can be achieved by excitting azimuthal spin wave modes with external alternating magnetic fields at multi-GHz frequencies. Depending on the sample geometry, lower limits exist [8] for the time needed for spin wave mediated vortex core reversal, which cannot be overcome by shortening the excitation length or by increasing the excitation amplitude. These limits are explained by a finite time needed for the energy transfer of the global excitation towards the center of the sample [8]. In the present talk it will be demonstrated by experiments [9] and micromagnetic simulations, how both, the duration of excitation and the time needed for undirectional vortex core reversal can be reduced to well below 100 ps, allowing a fast writing of V(ortex)MRAMs.

[1] B. Van Waeyenberge et al., Nature 444, 461 (2006), [2] S. Kim et al., Appl. Phys. Lett. 92, 022509 (2008), [3] S. Bohlens et al. Appl. Phys. Lett. 93, 142508 (2008), [4] B. Figean et al., Appl. Phys. Lett. 96, 132506 (2010), [5] Y-S. Yu et al., Appl. Phys. Lett. 98, 052507 (2011), [6] K. Nakano et al., Appl. Phys. Lett. 99, 262505 (2011), [7] M. Kammerer et al., Nature Communications 2, 279 (2011), [8] M. Kammerer et al., http://antixorg/abs/112.1903, [9] Performed by scanning transmission X-ray microscopy (STXM) at the MAXIVLS endstation at BESSY III. Berlin, Germany:

#### FD02

## Magnetic vortices and antivortices - From time-resolved imaging to the influence of temperature

Thomas Kamionka, Michael Martens, Andre Drews, Benjamin Krueger, Ole Albrecht and Guido Meier\*

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Subnanosecond dynamics and potential memory applications give rise to a broad scientific interest in properties of ferromagnetic microstructures with vortex and antivortex magnetization configuration. Such singularities can be excited to gyrate resonantly at a characteristic eigenfrequency. Dynamics are studied by time-resolved scanning transmission X-ray microscopy and broadband-ferromagnetic resonance measurements. The influence of temperature on the gyrotropic eigenfrequency of vortices in micrometer-sized permalloy squares is studied [1]. Ferromagnetic absorption measurements around room temperature show that the eigenfrequency is decreased by 5.4% when the temperature is increased by 100 °C. The lowering of the saturation magnetization and the associated modification of the potential of gyration is discussed as a reason for the frequency shift within the model of the harmonic oscillator. Micromagnetic simulations confirm a linear relation between the eigenfrequency and the saturation magnetization. Absorption measurements under application of static external magnetic fields yield the same percentage frequency shift as without external field when gyrations around equally displaced equilibrium positions are compared. Financial support by the DAAD via the Project 50725506 and by DFG via the SFB 668 and the GrK 1286 as well as the Forschungs- und Wissenschaftsstiftung Hamburg via the Exzellenzcluster "Nano-Spintronik" is gratefully acknowledged.

[1] T. Kamionka, M. Martens, A. Drews, B. Kruger, O. Albrecht, and G. Meier, Phys. Rev. B 83, 224424 (2011).

#### FD03

#### Oersted field contribution on the magnetic vortex core dynamics proved by homodyne detection

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The recent discovery that a spin-polarized current can induce magnetic vortex dynamics opened a new way to manipulate magnetization. Here, we report on the resonant excitation of a magnetic vortex core in a permalloy disc due to the injection of the high frequency current. The Oersted field contribution due to the inhomogeneous current distribution in the magnetic vortex core dynamics on a magnetic circular disc is experimentally investigated using a homodyne detection scheme. The homodyne technique allows us to obtain the resonance frequency, the amplitude of the vortex core gyration, and the phase shift between the microwave current and the magnetoresistance oscillation. By the vortex core position dependent measurements, the amplitude of the vortex core gyration obviously increases due to the enhancement of the Oersted field contribution. From systematic phase measurements as a function of microwave frequency, two phenomena are observed: (i) the trajectory of the vortex core gyration is distorted by the bounding effect due to the strong shape anisotropy close to the edge and (ii) the Oersted field contribution is dominant driving the vortex core close to the edge.

#### FD04

#### Vortex core switching driven by the novel inverse Faraday effect

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Many researchers have investigated the mechanisms of the ultrafast vortex core switching by electric current experimentally and theoretically [1]. In this presentation, we theoretically predict a novel mechanism of the core switching by an optical means. We show that for vortex structures, an unconventional inverse Faraday effect [2] without the spin-orbit interaction induces strong magnetic field when a circularly polarized light is applied and is useful for fast core reversal. The magnetic dynamics of the vortex core is calculated with the help of the simulation in the framework of the Landau-Lifshitz-equation, and the switching time is obtained to be 0.2 ns. This switching time is two orders shorter than one for the electrical switching, 20ns [1]. This ultrafast optical core switching is realized in the following process: first the orbital angular momentum, and then the spin chirality, which is a solid angle of the vortex spin structures, it to the spin angular momentum.

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#### FE01

#### Ab initio studies of strongly correlated electron systems Masatoshi Imada

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Recent trends of ab initio studies and progress in methodologies for electronic structure calculations of strongly correlated electron systems are discussed. The interest for developing efficient methods is motivated by recent discoveries and characterizations of strongly correlated electron materials and by requirements for understanding mechanisms of intriguing phenomena beyond a single-particle picture. A multi-scale ab initio scheme for correlated electrons (MACE) is developed by utilizing the hierarchical electronic structure in the energy space[1]. It provides us with a first-principles downfolding of the global band structure into low-energy effective models followed by accurate low-energy solvers for the models. The applications of MACE is illustrated with examples of several materials. In particular, we focus on electron correlations in iron-based superconductors[2], interplay of spin-orbit interaction in 5d systems such as Sr<sub>2</sub>IrO<sub>4</sub>[3], and Mott physics in organic conductors [4].

[1] M. Imada and T. Miyake, J. Phys. Soc. Jpn. 79 (2010) 112001. [2] T. Misawa, K. Nakamura and M. Imada, arXiv:1112.4682. [3] R. Arita, J. Kunes, A.V. Kozhevnikov, A.G. Eguiluz, and M. Imada, arXiv:1107.0835. [4] H. Shinaoka, T. Misawa, K. Nakamura, and M. Imada, arXiv:1110.6299.

#### FE02

## Coarse graining tensor renormalization by the higher-order singular value decomposition

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We propose a novel coarse graining tensor renormalization group method based on the higher-order singular value decomposition (HOSVD). This method provides an accurate but low computational cost technique for studying two- or three-dimensional (3D) lattice models. The method is demonstrated using the Ising model on the square and cubic lattices. By keeping up to 16 bond basis states, we obtain by far the most accurate numerical renormalization group results for the 3D Ising model.

#### FE03

#### Monte-carlo approach to stationary non-equilibrium of mesoscopic systems

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Calculating properties of correlated systems out of equilibrium is a challenging task, even if on targets only stationary situations. In particular, transport through nanoobjects like molecules or quantum dots is of strong interest, and a theory to calculate transport properties or merely local quantities in a reliable way for reasonably strong correlations very desirable. Based on a suggestion by Han and Heary [1] we use advanced quantum Monte-Carlo techniques to calculate quantities for stationary nonequilibrium of a single-impurity Anderson model with high accuracy [2]. Employing a two-dimensional analytical continuation based on maximum entropy, we are able to calculate different local quantities like current, magentization or double occupancy. We compare our results to those obtained by different methods like time-dependent numerical renormalization, real-time quantum Monte-Carlo or real-time density-matrix renormalization group.

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#### FE04

## SU(4) symmetry for strongly correlated electrons: Kondo and mixed-valence effects in terms of gell-mann matrices

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The concept of dynamical symmetries [1] is used for formulation of the renormalization group approach to the Kondo effect in the Anderson model with repulsive and attractive interaction U in the Kondo and mixed valence regimes. It is shown that the generic local symmetry of the Anderson Hamiltonian is determined by the SU(4) Lie group. The Anderson Hamiltonian is rewritten in terms of the Gell-Mann matrices of 4-th rank, which form the set of group generators and the basis for construction of the irreducible vector operators describing the excitation spectra in the charge and spin sectors. The multistage Kondo screening is also interpreted as a manifestation of the local SU(4) dynamical symmetry. It is shown that the similarity between the conventional Kondo cotunneling effect for spin 1/2 in the positive U model and the Kondo resonance for pair tunneling in the negative U model is a direct manifestation of implicit SU(4) symmetry of the Anderson/Kondo model. The transformation from local SU(4) symmetry to the global SO(4) symmetry in the Hubbard chain model is also discussed.

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#### GA01

Detection of orbital fluctuations above the structural transition temperature in iron pnictides and chalcogenides

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The electronic structure of AEFe<sub>2</sub>As<sub>2</sub> (AE=Ca,Sr,Ba) and Fe<sub>1+y</sub>Te is studied using point contact or quasiparticle scattering spectroscopy (QPS). [1] For AE=Sr,Ba and Fe<sub>1+y</sub>Te, a conductance enhancement reproducibly appears at To, well above the structural transition temperature, Ts. For Ba(Fe<sub>1+x</sub>Co<sub>2</sub>)<sub>2</sub> As<sub>2</sub>, the conductance enhancement exists only in the underdoped regime. Thus, we add a new region of strong correlations to the phase diagram: For x = 0 to 5.5, To ~ 175 K to ~ 150 K, respectively. These strong correlations arise from orbital fluctuations, and the prediction [2] that the excess conductance above Ts is only observed in detwinned crystals that exhibit resistance anitorpy above Ts [3] has been verified [1]. The role of orbital fluctuations and nematicity in quantum criticality of these materials will be discussed. [4] Supported by the Center for Emergent Superconductivity, an EFRC, US DOE, DE-AC0298CH1088; Cambridge EPSRC, and Ames Lab US DOE, DE-AC02-07CH11358.

I. H. Z. Arham et al, arXiv:1201.2479. 2. W.-C. Lee and P. Phillips, arXiv:1110.5917. 3. I. R. Fisher et al., Rep. Prog. Phys. 74, 124506 (2011) and references therein. 4. S. Kasahara et al, Phys. Rev. B 81, 184519 (2010); S. Kasahara et al., Bull. Am. Phys. Soc. 56, Z26.00010 (2011); Y. Matsuda, private communication.

#### GA02

## Nature of magnetic excitations in superconducting iron superconductors

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Since the discovery of the metallic antiferromagnetic (AF) ground state near superconductivity in iron-pnictide superconductors, a central question has been whether magnetism in these materials arises from weakly correlated electrons, as in the case of spin-density-wave in pure chronium, requires strong electron correlations, or can even be described in terms of localized electrons such as the AF insulating state of copper oxides. We use inelastic neutron scattering to determine the absolute intensity of the magnetic excitations throughout the Brillouin zone in a number of iron-based superconductors, which allows us to obtain the size of the fluctuating magnetic moment in absolute units. We find that superconducting BaFe<sub>1</sub>oNi<sub>0.1</sub>As<sub>2</sub> and AF BaFe<sub>2</sub>As<sub>2</sub> both have fluctuating magnetic moment similar to those found in the AF insulating copper oxides. The common theme in both classes of high temperature superconductors is that magnetic excitations have partly localized character, thus showing the importance of strong correlations for high temperature superconductivity.

#### GA03

## Universal microscopic description of the infrared conductivity of 122 iron arsenides

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We report the full complex dielectric function of high-purity  $Ba_{0.68}K_{0.32}Fe_2As_2$  single crystals with  $T_c=38.5K$  determined by wide-band spectroscopic ellipsometry at temperatures 10 < T < = 300K. We discuss the microscopic origin of superconductivity-induced infrared optical anomalies in the framework of a multiband Eliashberg theory with two distinct superconducting gap energies  $2\Delta_A = 6k_BT_c$  and  $2\Delta_B = 2.2k_BT_c$ . The observed unusual suppression of the optical conductivity in the superconducting state at energies up to  $14k_BT_c$  can be ascribed to spin-fluctuation-assisted in the clean limit of the strong-coupling regime. We further demonstrate that the same model provides a good description of the infrared conductivity of electron-doped compounds in this class of superconductors.

#### GA04

## Various fabricating conditions of potassium doped $BaFe_2As_2$ films by pulsed laser deposition system

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Among the various Fe-based superconductors, potassium doped BaFe<sub>2</sub>As<sub>2</sub> is favorable for application because of its higher transition temperature and a low anisotropy compared to other iron based superconductors. To study the superconductivity and applicable aspects, high quality thin films should be fabricated. However, it is difficult to fabricate thin films because of the high volatility of potassium. In this paper, we show the details of fabricating technique of Ba<sub>1,x</sub>K<sub>x</sub>Fe<sub>2</sub>As<sub>2</sub> films by ex-situ PLD method in various conditions and physical properties of these films. Barium ratio in target is controlled to make films with various potassium doping rate. An annealing temperature and amount of potassium are also controlled to find out optimal condition of fabricating films.

#### **GB01**

## Nonreciprocal directional dichroism and toroidalmagnons in multiferroic materials

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We investigate dynamical magnetoelectric effect at electromagnon resonances in multiferroic materials. One of the typical phenomena caused by interferences between the magnetic and electric responses is the nonreciprocal directional dichroism (NDD), where differences in the absorption spectra of electromagnetic wave are observed when propagation directions of the incident wave are inverted. We show that NDD is observed at a magnon excitation that is accompanied by a dynamical toroidal moment, i.e., at a toroidalmagnon excitation. As examples, we discuss toroidalmagnons in multiferroic materials RMnO<sub>3</sub> and Ba<sub>2</sub>CoGe<sub>2</sub>O<sub>7</sub>.

S. Miyahara and N. Furukawa, J. Phys. Soc. Jpn. 81 (2012) 023712/1-4. S. Miyahara and N. Furukawa, J. Phys. Soc. Jpn. 80 (2011) 073708/1-4.

#### **GB02**

Electronic ferroelectricity in correlated electron systems Sumio Ishihara\*, Makoto Naka and Akihiko Sekine Department of Physics., Tohoku University. Japan

Exotic ferroelectrics and mutiferroics are one of the attractive topics in recent correlated electron systems. It is widely known that ionic lattice distortions responsible for the electric dipole moments in conventional ferroelectric compounds. Recently new class of ferroelectricity, termed "electronic ferroelectricity" where electronic charge order brings about the electric polarization, has attracted much attention from fundamental and application view points. We study theoretically ferroelectric and multiferroelectric properties of the electronic ferroelectricity in correlated electron systems. We mainly focus on the layered iron oxides, LuFe<sub>2</sub>O<sub>4</sub>, and the low-dimensional organic salts, kappa-(BEDTTF)<sub>2</sub>Cu<sub>2</sub>(CN)<sub>2</sub>. In the iron oxide, we examine the electronic and dielectric properties by analyzing the model Hamiltonian. It is found that a low-lying charge fluctuation due to a geometrical frustration plays a key role on an origin of ferroelectricity. This fluctuation is expected to be detected by the inelastic x-ray scattering experiments. As for the molecular organic salts, we found that the ferroelectric charge order strongly competes with a Mott insulator based on molecular dimerization, termed a dimmer Mott insulating plase, as well as a metallic phase. A collective excitation mode for an electric polarization and superconductivity due to polarization

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#### GB03

#### Interplay between electronic ferroelectricity and magnetism in molecular TMTTF salts

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We theoretically investigate the interplay between charge ordering and magnetic states in quasi-one-dimensional molecular conductors TMTTF2X, motivated by the observation of a complex variation of competing/coexisting phases. We show that the ferroelectric-type charge order (electronic ferroelectricity) increases two-dimensional antiferromagnetic spin correlation, whereas in the one-dimensional regime two different spin-Peierls states are stabilized. By using first-principles band calculations for the estimation for the transfer integrals and comparing our results with the experiments, we identify the controlling parameters in the experimental phase diagram to be not only the inter-chain transfer integrals but also the amplitude of the charge order. We also discuss the effect of a diagonal inter-chain transfer integral, which causes spin frustration in the dimer-Mott insulating state, but conversely enhances the magnetic ordering in the charge ordered phase.

K. Yoshimi, H. Seo, S. Ishibashi, and S. E. Brown, arXiv:1110.3573 (to be published in Phys. Rev. Lett.); arXiv:1110.3575 (to be published in Physica B).

#### GB04

## Dielectric anomaly in dimer-Mott insulator $\beta^{\prime}\mbox{-}(BEDT\mbox{-}TTF)_2ICl_2$ with square lattice

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Organic dimer-Mott insulator  $\kappa$ -(BEDT-TTF)<sub>2</sub>Cu<sub>5</sub>(CN)<sub>5</sub> with triangular lattice has been extensively investigated as a candidate for spin liquid materials. Moreover, the recent discovery of a relaxorlike dielectric response [1] in this system paved the way to a new phenomenon of purely electronic ferroelectricity as well as the strong charge fluctuation or charge disproportionation on dimers [2, 3]. Here, we have investigated the dielectric property of a typical antiferromagnetic dimer-Mott insulator  $\beta$ -(BEDT-TTF)<sub>2</sub>[Cl<sub>2</sub> with square lattice, which is compared to that of a spin liquid candidate  $\kappa$ -(BEDT-TTF)<sub>2</sub>Cu<sub>2</sub>(CN)<sub>5</sub>. Anomalous broad peak structure in dielectric constant was observed in  $\beta$ -(BEDT-TTF)<sub>2</sub>Cl<sub>2</sub>, which shows a Curie type temperature dependence with higher Curie temperature than the antiferromagnetic transition and a strong frequency dependence similar to ferroelectric relaxors. The dielectric constant parallel to the BEDT-TTF firmer layer is larger than perpendicular to it, which suggests that the charge degrees of freedom within a dimer of BEDT-TTF molecule has significant role in the dielectric anomaly in dimer-Mott type organic insulator. The coupling between the charge and spin degrees of freedom will be discussed.

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#### **GB05**

## Multiferroic transition in a quasi-layered bismuth ferrite Chan-ho Yang

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Usually magnetic transition temperature is insensitive to structural modification by external strain, while ferroelectric transition temperature can be largely tuned by strain. As a counter-example of the generic tendency of the magnetic transition, we will focus on a highly elongated multiferroic BiFeO<sub>3</sub> thin film, which has been newly discovered as a quasi-layered perovskite phase stabilized through a misfit strain. Special attention has been paid to the phase competition between the new highly elongated phase and the normal phase leading to colossal electron-strain. However the magnetic properties of the new highly elongated phase has not been fully unveiled yet. Here we will provide experimental evidence for the large suppression of the magnetic order indicating strong spin-lattice coupling and discuss underlying physics inherent in the highly-elongated quasi-layered perovskite.

#### GC01

## Field-dependent Fermi surface and high-field superconductivity in URhGe

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We present the first observation of the Fermi surface in the ferromagnetic superconductor URhGe [1]. A small class of heavy fermion ferromagnetis including URhGe show bulk superconductivity deep within the ferromagnetic state. Existing results for URhGe [2] strongly suggest that the superconductivity arises from p-wave Cooper pairing of electronic quasiparticles with equal spin. Uniquely within this class, URhGe shows a second, field-induced superconducting phase with critical fields in excess of 32 tesla [3]. Our measurements of the Shubnkov-de Haas effect, made on the verge of field-induced superconductivity, reveal a small Fermi surface pocket that shrinks with increasing applied magnetic field, eventually disappearing at a topological transition of the Fermi surface (a Lifshitz transition). The quasiparticle mass decreases and remains finite, implying that the Fermi velocity vanishes due to the collapse of the Fermi wavevector. A simple model calculation of the orbitally limited critical field for the shrinking Fermi pocket reproduces the experimental phase diagram for high-field superconductivity, suggesting that it is the slowing down of quasiparticles at a Lifshitz transition that allows superconductivity or survive above 30 T, without a divergence of the quasiparticle mass. This represents an innotrant denature from all present theoretical treatments of the phase formation.

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#### GC02

#### Conventional quantum criticality in CeCu<sub>2</sub>Si<sub>2</sub> Oliver Stockert\*

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The heavy-fermion compound CeCu<sub>2</sub>Si<sub>2</sub> displays unconventional superconductivity around the magnetic quantum critical point (QCP) where antiferromagnetic order vanishes. The superconducting state is characterized by a spin excitation gap, and almost critical spin fluctuations are responsible for the Cooper pair formation. The energy and wave-vector dependence of the normal state spin fluctuations as measured by inelastic neutron scattering indicate a three-dimensional spin-density-wave QCP in agreement with macroscopic properties. Furthermore, muon spin rotation under pressure allowed to tune through the QCP and to study the interplay between magnetism and superconductivity in more detail. Both phenomena are connected by a first-order transition with the occurrence of phase separation. Small magnetically ordered regions are observed well below the superconducting Tc. The results are in line with the findings obtained by neutron scattering and will be compared to other quantum-critical heavy-fermion compounds.

#### GC03

#### Shubnikov-de Haas oscillation in PuIn3

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Highly correlated f-electron systems are extensively studied because of their unusual physical properties. While the electronic state of 4f rare earth compounds is understood based on the localized 4f state hybridized with conduction bands, that of 5f electrons in actinide compounds is not well defined. In this context, direct observation of Fermi surface provides strong implication whether the 5f electrons are localized or delocalized to participate in Fermi surface. In the present study, we report Shubnikov-de Haas (SdH) study in a plutonium compound Puln3 with the cubic AuCu3-type structure. We succeeded in observing a clear SdH oscillation as a function of inverse magnetic field, where the SdH frequency is proportional to the cross sectional area of the Fermi surface along the magnetic field direction. The obtained SdH frequency can be explained by the band calculations assuming the 5f electrons are delocalized. The SdH oscillation was only observed when the field is applied near the <111> direction and suddenly vanishes, which cannot be explained from the calculated Fermi surface properties and needs further investigations.

#### GC04

## Spin fluctuations and Lifshitz transition in UGe<sub>2</sub> probed by Larmor neutron diffraction under pressure

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We present high resolution measurements of the lattice constants of UGe<sub>2</sub> under pressure probed by a novel technique that utilizes Larmor precession of polarized neutrons to surpass the resolution of conventional scattering methods by an order of magnitude. At low temperature UGe<sub>2</sub> is ferromagnetic up to critical pressure p\_c but superconductivity is peaked at a lower pressure p\_ x coinciding with a less understood transition within the ferromagnetic state [1]. At ambient pressure we observed sharp anomalies in the lattice parameters at the Curie temperature, T\_C. At higher pressure sharp anomalies in the lattice parameters at the Curie temperature, T\_C. At higher pressure sharp anomalies in the lattice parameters at both T\_C and T\_X (the characteristic temperature for the transition that occurs at p\_x) shift to lower temperatures in agreement with the known phase diagram. We show that the magnetic thermal expansion is dominated by the critical fluctuations near T\_C and is well explained by the spin fluctuation theory at p<p\_x, however at p=p\_x we identify an additional contribution to the magnetic thermal expansion which is due to pressure driven Lifshitz transition [2] in spin bands of UGe<sub>2</sub>. Together with the magnetic correlation length diverging under pressure, this may play an important role in stabilising superconductivity.

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#### GD01

#### Ultrafast manipulation of magnetic order

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The interaction of sub-picosecond laser pulses with magnetically ordered materials has developed into an extremely exciting research topic in modern magnetism and spintronics. From the discovery of sub-picosecond demagnetization over a decade ago to the recent demonstration of magnetization reversal by a single 40 femtosecond laser pulse, the manipulation of spins by ultra short laser pulses has become a fundamentally challenging topic with a potentially high impact for future spintronics, data storage and manipulation and quantum computation. In addition, when the time-scale of the perturbation approaches the characteristic time of the exchange interaction (~10-100 fs), the magnetic dynamics enters a novel, highly non-equilibrium, regime where the exchange interaction might even become time dependent. Using ultrashort excitations, we may be able to manipulate the exchange interaction itself or use it strength to manipulate or even switch magnetization. Such studies require the excitation and probing of the spin and angular momentum contributions to the magnetic order at timescales of 10fs and below, a challenge that could be met by the future fs X-ray FEL's.

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#### GD02

#### Ultrafast emergence of nanoscale ferromagnetism far from equilibrium Hermann Andreas Durr

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Long-range magnetic order in solids is usually ascribed to the exchange interaction between electron spins. Close to equilibrium this leads to spontaneous magnetization when the system cools below the magnetic ordering temperature. We show that the far from equilibrium flow of angular momentum via spin currents can also achieve long-range ferromagnetic order even above the ordering temperature. To reveal this process, we use ultrafast x-ray diffraction at SLAC's Linac Coherent Light Source to probe the nucleation, growth and transient existence of ferromagnetic order on the nm length and fs timescale after fs optical laser excitation has brought a metallic 3d - 4f alloy system into a highly non-equilibrium chaotic spin state. The technological exploitation of this effect could pave the way for novel ways to manipulate and transport information within the smallest space-time dimensions.

#### GD03

Modeling of ultra-fast magnetisation dynamics

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The ultra-fast laser-induced magnetisation dynamics is currently an attractive topic of modern magnetism due to a new physical phenomena appearing during these strongly non-equilibrium processes and the possibility to extend magnetic switching speed down to femtoseconds. While in ferromagnetic metals such as Ni, Fe, Gd only ultra-fast demagnetisation has been observed, in ferrimagnetic materials such as FeCoGd the controllable magnetization switching in several ps has been reported [1]. Recently we have reported that this switching occurs under ultra-fast heat alone without the necessity of any other external stimulus [2]. Our modeling of the underlying processes is based on two approaches: (i) atomistic simulations using the Landau-Lifshitz-Gilbert equation and (ii) micromagnetic approach using classical and quantum Landau-Lifshitz-Bloch (LLB) equation. Both models assume thermal magnetization dynamics via the coupling to electron (and phonon for Gd) temperature bath. The comparison of our modeling based on the LLB approach with the experiments in Ni [3] and Gd shows an excellent agreement providing a validation of thermal magnetization dynamics mechanism. In agreement with the experiment the modeling reveals different timescales for demagnetisation: fast (100 fs) in Ni and slow (50 ps) in Gd. We will present our atomistic modeling on the switching of a ferrimagnetic CoFeGd compound with the only external stimulus provided by the ultra-fast heating. The atomistic results will be explained in terms of the dynamical system based on the recently-derived two-component LLB micromagnetic equation.

1. K.Vahaplar et al, Phys. Rev. Lett. 103, 117201 (2009) 2. T.Ostler et al., Nature Commun., 3, 666 (2012) 3. U.Atxitia et al., Phys. Rev. B 81, 174401 (2010)

#### **GD04**

#### Ultrafast inverse Faraday effect in paramagnetic dielectrics

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We present a critical and thorough analysis of the concept of a light induced effective magnetic field. The analysis reveals incontinency in the theoretical interpretation of ultrafast femtomagnetic phenomena in paramagnetic dielectrics. Furthermore, we discuss the validity of the assumption of equivalence of origins of opto-magnetic and magneto-optical effects, which has not been rigorously tested so far. We support our arguments by the all-optical femtosecond time-resolved measurements of the ultrafast inverse Faraday Effect in terbium gallium garnet crystal. The measurement of the transient rotation of polarization of the probe pulse induced by a pump pulse allows us perform a systematic study of the full nonlinear optical response of the sample. The magneto-optical contribution attributed to inverse Faraday effect is then extracted by thoroughly studying the dependence of the signal upon the pump polarization. The amplitude of the Faraday rotation was found to significantly exceed that predicted by phenomenological theory relying on the static Verdet constant value. The discrepancy demonstrates that the nature of the inverse Faraday Effect on femtosecond timescale is different from that of the quasi-static inverse Faraday Effect observed by Pershan and coworkers [1,2] and that it cannot be described in terms of the effective magnetic field.

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#### GE01

#### Detection of domain wall position and magnetization reversal in nanostructures using the magnon contribution to the resistivity

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We present a new magnetoresistance (MR) based on the contribution of the magnons to the resistivity, i.e., Magnon Magnetoresistance (MMR). We show that the MMR can be used to study the magnetization reversal in systems with either perpendicular (FePt) or in-plane anisotropy (NiFe), including domain wall (DW) position detection. MMR measurement exhibits a linear dependence of the resistivity on the applied field, originating from electron-magnon scattering. The drop of resistivity, due to the abrupt change of magnon density, corresponds to the magnetization switching. MMR measurement allows detecting precisely the position of a DW along a FePt nanowire, and following the dynamic of DW motion. In system with in-plane magnetization as NiFe nanowires, we show that the enhancement of the shape anisotropy in the narrowest widths leads to the disappearance of the AMR, the remaining contribution to the MR being that of the MMR. We also show that the 0MR signal allows detecting the position of a DW along a change of paradigm in the study of in-plane magnetized nanowires: for certain field directions the magnot contribution dominates the MR, clearly overcoming the AMR.

V. D. Nguyen et al., Phys. Rev. Lett. 107, 136605 (2011) V. D. Nguyen et al., Appl. Phys. Lett. 99, 262504 (2011)

#### **GE02**

#### Tunable resistivity of individual magnetic DWs

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Despite the broad interest in current-induced magnetic domain wall (DW) motion, the current-DW interaction is still not fully understood. A property strongly related to this interaction is the intrinsic DW resistivity. Measuring this property is therefore an important step in unraveling the physics behind DW motion and in particular the non-adiabatic spin torque. However, measurements of the DW resistivity as a function of the main property of the DW, its width, are still lacking. Here, we investigate in a unique way how the resistivity of individual DWs depends on their width [1], by using focused ion beam irradiation to tune the DW width D in perpendicularly magnetized Pt/Co/Pt strips. We use a very direct way to measure the resistivity of individually appearing domain walls, by monitoring in real-time the domain structure in a Kerr microscope. We find that convincingly, the DW resistivity scales with  $1/D^2$ , in quantitative agreement with the theory of Levy and Zhang [2].

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#### GE03

Observation of domain-wall capacitance in permalloy nanowires Kulothungasagaran Narayanapillai, Mahdi Jamali, Ajeesh Sahadevan and Hyunsoo

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Studies on domain walls have attracted interest due to its potential applications on storage and logic devices. Domain walls can be detected by magneto-resistance effect such as the giant magneto-resistance effect and by magnonic effect [1]. We here report the detection of a domain wall by employing the capacitive effect of a single domain wall in a Permalloy nanowire. Spin-capacitance has been reported as a phenomenon in magnetic tunnel junctions and interfacial oxides where the spin configuration affects not only the resistance across the tunnel junction, but also the capacitance known as magneto-capacitance [2]. Semi-circular nanowires are patterned with the film structure of Ta (5 nm)/ NiFe (30 nm)/ Ta (5 nm) and the capacitance and magneto-resistance is measured. We observe that the magneto-capacitance is a reciprocal of the magneto-resistance. In addition to the resistive component, the capacitance due to the domain wall spin structure enables us to detect the presence of a domain wall. In angular measurements with externally applied magnetic fields, large changes in the field dependent capacitance are discussed.

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#### GE04

## Proposal new type of low current driven spin logic in PMA TbFeCo wire

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We have performed a current driven magnetic domain AND logic circuit by designing a TbFeCo wire which has a large perpendicular magnetic anisotropy and high domain wall motion seed in 'Y' shape with two inputs and an output. The TbFeCo wire thickness is 4.8 nm. Well-defined starting states of as single domain wall from two inputs were recorded by pulse currents. The domain walls propagate into the output as a combined domain wall and detected by anomalous Hall voltage through a pair of electrodes (J=3.8×10<sup>10</sup> [A/m<sup>2</sup>]). In the case of two input domain walls are recorded as a couple bits of '1' /1', a sharp-changed anomalous Hall voltage is detected as a bit of '1'. In another case of only one input domain wall is recorded, there is no change in anomalous Hall voltage which corresponds to a bit of '0' due to a different width of input and output wires which causes annihilated the domain wall. These results indicated that the circuit of TbFeCo wire work as a logic AND function. Moreover, OR and NOT function can be also designed. This technique expectedly promises lowerpowered, higher-speed, nonvolatile devices with which to build the next generation of computing technology.

#### **GE05**

## Current-induced domain wall motion in perpendicularly magnetized nanowire

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Electrical displacement of a domain wall (DW) is a prospective method for information processing in new type of magnetic non-volatile memories and logic devices. Such novel spintronic devices require a low DW drive current and a high DW de-pinning field for stable information retention. We show that Co/Ni multilayer with perpendicular magnetic anisotropy is a promising material that satisfies these requirements. An electric current can drive a DW in one direction regardless of the polarity of a magnetic field in a Co/Ni nano-wire with perpendicular magnetization, i.e., the current can drive a DW against a magnetic field. Furthermore, both the DW velocity and the threshold current density for the DW motion show almost no dependence on the external magnetic field. These counter-intuitive behaviors can be interpreted as the consequence that the intrinsic pinning mechanism determines the threshold current, and the adiabatic spin torque dominates the DW motion in this system [1-6]. The established field-insensitivity of the electrical DW motion is promising for future spintronics applications based on the DW motion. This work was partly supported by the NEDO Spintronics nonvolatile devices project and a Grant-in-Aid for Scientific Research (S) from the Japan Society for the Promotion of Science.

T. Koyama et al., Nature Materials, 10 (2011) 194. [2] T. Koyama et al., Appl. Phys. Lett. 98 (2011) 192509. [3]
 K. Ueda et al., Appl. Phys. Express 4 (2011) 063003. [4] D. Chiba et al., Appl. Phys. Express 3 (2010) 073004. [5] H. Tanigawa et al., Appl. Phys. Express 1 (2008) 101303.

#### **GF01**

#### Melting spin ice

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Geometric magnetic frustration is a ubiquitous phenomenon within condensed matter physics. The magnetic ground states are highly degenerate, leading to competing low temperature phases and enhanced quantum fluctuations. The nature of the magnetic fluctuations in these systems is currently the subject of lively investigation. "Spin ice" model magnetic systems combine the rich phenomenology associated with geometric frustration with the theoretically tractable nature of Ising models, drawing on a close analogy with the statistical mechanics of proton disorder in water ice. Theory predicts the low-temperature magnetic excitations in spin ices consist of deconfined magnetic charges, or monopoles. A recent transverse-field (TF) muon spin rotation (µSR) experiment [S T Bramwell et al, Nature 461, 956 (2009)] reports results claiming to be consistent with the temperature and magnetic field dependence anticipated for monopole nucleation - the so-called second Wien effect. We demonstrate via a new series of µSR experiments in Dy.[Ti<sub>2</sub>O, that such an effect is not observe spin fluctuations which become temperature independent at low temperatures, behaviour which dominates over any possible signature of thermally nucleated monopole excitations.

#### **GF02**

#### Artificial spin ice: Dimensional reduction, avalanches and disorder Remo Viktor Hugli, Gerard Duff and Hans - Benjamin Braun\* Physics. University College Dublin. Ireland

Regular arrays of nanomagnets have proven to serve as ideal models for the study of spin-ice phenomena [1]. In contrast to their 3D analogues, the pyrochlore systems, the magnetic moments in 2D artificial spin ice systems exhibit high anisotropies which renders them stable at room temperature. This allows imaging with various experimental techniques in real space and at room temperature, which makes such metamaterials ideal candidates for comparison with theoretical models. Here we explore the nonequilibrium properties of a model that combines long range dipolar interaction, frustration and disorder which explains experiments on 2D artificial kagome spin ice [2]. We show how frustration induced by dipolar interaction leads to unusual 1D avalanches, an example of dimension reduction due to frustration. In particular, magnetization reversal proceeds via nucleation and subsequent avalanche-type dissociation of emergent monopole-antimonopole pairs along 1D Dirac strings. We show how the flow of these emergent monopoles or charges can be controlled via judicious modification of the islands' anisotropy and magnetic fields, thus paving the way for magnetic information processing devices.

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#### GF03

## Low temperature magnetic studies of geometrically frustrated $SrHo_2O_4$

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SrHo<sub>2</sub>O<sub>4</sub> is a geometrically frustrated compound as its structure allows the magnetic ions to be arranged in a network of triangles and hexagons [1]. Magnetic susceptibility measurements on single crystals of SrHo<sub>2</sub>O<sub>4</sub> show a disparity between the Curie-Weiss constant (-16.9 K) and the establishment of Neel order at 0.68 K [2], which is typical for frustrated systems. Low-temperature measurements of magnetisation [3] show the highly anisotropic nature of this crystal, and more intriguingly, the dM/dH curve for H// b shows double peak behaviour indicating a rich H-T phase diagram. We have performed powder neutron diffraction experiments using the GEM diffractometer (ISIS) [2]. At low temperatures, both k=0 magnetic Bragg reflections and a strong diffuse scattering feature - a broad peak near the position (0 0 1/2) - were observed. Recent experiments on single crystals of SrHo<sub>2</sub>O<sub>4</sub> using the D7 instrument (ILL), reveal that this feature can be attributed to a one-dimensional collinear ordering with the spins pointing along the b-axis. This suggests a highly unusual magnetic ground state for SrHo<sub>2</sub>O<sub>4</sub>, as it is made up of two distinct components: one of the crystallographically inequivalent Ho<sup>3+</sup> sites orders in a long-range magnetic structure, while the other orders in a short-range one-dimensional structure.

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#### **GF04**

#### Spin densities in manganese molecular cluster : [Mn<sub>3</sub>L<sub>4</sub>](ClO<sub>4</sub>)<sub>2</sub>(H<sub>2</sub>O)<sub>2</sub>

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The molecular cluster  $[Mn_3L_4](ClO_4)_2(H_2O)_2$ , (HL=2-methoxy-6-(pyridine-2-ylhydrazono-methyl)-phenol) has been described as an almost lineal trimer of Mn(II) in a high spin configuration with two isotropic magnetic interactions: nearest neighbours and next nearest neighbours interactions1. One striking fact is the experimental EPR g factor of 2.14 for such an isotropic system. Due to its isotropic and not degenerated nature, the system is an ideal benchmark to investigate the interplay of the magnetic interactions in systems with several magnetic centers and more than one electron per magnetic center. In this contribution we will show how an in deep study of the system through the combination of several experimental and theoretical methods (polarized neutron diffraction at ILL and LLB, magnetic measurements, Density Functional Theory and spin Hamiltonian models) has allowed us to demonstrate the isotropic nature of the system and the anti-ferromagnetic character of the intra-molecular magnetic interactions.

[1] J. Tang et al., 2007 Eur. J. Inorg. Chem 4119

#### **GF05**

#### Electronic structure and magnetic properties of Cr-doped rutile TiO<sub>2</sub>: Charge and magnetic state of Crimpurity

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We report theoretical and experimental studies of Cr-doped rutile TiO<sub>2</sub>. In experiment we observe that the electronic and magnetic properties of Cr-doped rutile TiO<sub>2</sub> are highly dependent on the growth conditions. The ferromagnetic component in magnetic susceptibility and the Cr defect concentration are enhanced for the sample grown under the oxygen flow in comparison with those grown under the Ar flow. To understand the charge state of Cr dopants and their role in response to the external magnetic field, we carried out density functional theory (DFT) calculations for the Cr-doped rutile TiO<sub>2</sub>. From the results of the formation energy calculations for the Cr atom substituting the Ti site and the O vacancy, assuming the thermodynamic equilibrium, we demonstrate that the Cr<sup>3+</sup> state is a source of Curie-Weiss-type magnetic response in the O-poor limit, whereas the Cr<sup>4+</sup> defect states contribute to the ferromagnetic component in the moderately reduced condition. We also provide the electronic structures of various defect configurations and attempt to explain the optical and electronic properties of the Cr-doped system.

#### **GG01**

#### Magnetic nanoparticle arrays by nanomasking pattern transfer

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In nanomasking, self-assembled nanoparticle monolayers are used as etch masks to pattern underlying thin films [1]. Here we apply this technique to prepare two types of magnetic nanodot arrays, and characterize their structural and magnetic properties. Large area nanoparticle arrays were fabricated by a Langmuir film method and transferred to a solid substrate. After electron irradiation curing and brief exposure to an oxygen plasma, most of the surfactant has been removed yet the ordered pattern of the particle cores has been preserved. We examined direct pattern transfer by methanolbased reactive ion etching into a FePt film. A second strategy prepares nanodot arrays by filling an antidot lattice. Here 2 nm of titanium is deposited on the array of bare nanoparticle cores, which are then removed with phosphoric acid to leave an alumina antidot mask. Here the mask was prepared on a silica film. RIE was then used to generate an array of pits, and permalloy was sputtered to overfill the pits. Magnetic measurements as a function of methanol etching show the transition from a continuous film to an array of nanodots.

1. C. R. Hogg, S. A. Majetich, and J. A. Bain, IEEE Trans. Magn. 46, 2307-2310 (2010)

#### **GG02**

#### Spin wave bands and bandgaps in a two-dimensional ferromagnetic antidot arrav

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The spin wave band structure of a two-dimensional square array of NiFe circular antidots having diameter of 120 nm and periodicity of 800 nm has been investigated by using Brillouin light scattering and micromagnetic calculations based on the dynamical matrix method. The external magnetic field was applied in the plane and perpendicularly to the transferred wave vector. Extended and localized spin modes having a propagative nature were found. Opening of bandgaps is interpreted in terms of Bragg diffraction of spin waves from the antidot lattice and this effect is explained by studying the behavior of the internal field. According to an analytical model, the relevant scattering potential for Bragg reflection is not provided by the holes themselves, but by the concomitant internal field inhomogeneity between holes [1]. This is in contrast to antidots in photonics and electronics where the back-reflection is directly caused by the presence of holes. The research leading to these results has received funding from the European Community's Seventh Framework Programme (FP7/2007-2013) under Grant Agreement n228673 (MAGNONICS).

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#### **GG03**

#### Ratchet effect in magnetic domain wall motion induced by 2D arrays of triangular submicrometric holes

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Controlling magnetic domain wall motion in ferromagnetic materials is essential for the development of spintronics. The introduction of asymmetric pinning potentials has proved itself as an efficient tool to induce ratchet effects in domain wall motion, as observed in nanowires with triangular section [1] or asymmetric notches [2]. Similar effects can be induced in extended ferromagnetic thin films when introducing 2D arrays of asymmetric holes[3], although the study is rather more complex due to the higher dimensionality of domain walls [4]. In the present study, antidot arrays have been fabricated on amorphous thin films of Co73Si27 with in-plane uniaxial anisotropy. Arrays of triangular antidots of size 500nm and lum and edge-to-edge separation of three, four and five times their size have been studied. Equivalent arrays with rhombohedral antidots have been fabricated for comparison. A Kerr effect microscope has been used for the magnetic imaging. Magnetic domain walls are observed to move through the triangular antidot arrays only in the forward direction, defined from the base of the triangles to the vertex. The cause of the ratchet effect is confirmed to be the holes asymmetry, as domain wall motion is symmetric when crossing rhombohedral (symmetric) antidot arrays

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#### **GG04**

#### Tailored magnetic anisotropy of Py /Co bilayer ordered nanohole arrays

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Highly ordered arrays of magnetic nanoholes show a great potential for technological applications, as high-density magnetic storage media [1] or magnonic devices [2]. Magnetism of ordered nanohole arrays is determined by composition and particularly by geometry characteristics [3]. In this work, magnetic nanohole arrays have been sputtered onto hexagonally self-assembled anodic alumina membranes (AAM) [4] with 35 nm hole diameter and 105 nm interhole distance. Thickness of Py/Co bilayer nanohole arrays were: 10nm, 20nm and 43 nm for the Py layer, with a constant thickness of 28 nm for Co layer. Their magnetic properties were determined by vibrating sample magnetometer and magneto optical Kerr effect. The in-plane angular dependence of coercivity shows that the magnetic behaviour of the Py layer strongly depends on its interaction with the Co, which reminds that of a spring-like system. The sample with thinnest Py layer exhibits highest coercivity (Hc=140 Oe) which decreases with Py thickness (Hc=16 Oe). A bimagnetic-like behaviour is observed in some samples indicating the relevance of local magnetostatic interactions. The overall results have been compared to those of continuous Py/Co bilayer films, showing an increase of five times its coercivity by the only presence of the nanoholes.

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#### **GG05**

#### Tailoring magnetic properties of Co thin films through antidot arrays: crossover from antidot to dot regime

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Magnetic properties of thin films are hardly controllable since they depend strongly on the microstructure and the presence of defects. Nanopatterning of dot [1] and antidot [2] arrays are two different approaches to tailor magnetic properties only as a function of geometry. In this study, a series of antidot arrays have been fabricated with square symmetry, varying the periodicity from 500nm, that is in the diluted regime to 105nm, almost in the percolation limit. The holes were etched using a focused gallium ion beam, that allows the control of the geometrical parameters with nanometric resolution. Magnetic characterization by MOKE and magnetic imaging by XPEEM show a change in the dependence of coercivity with iodicity below 140nm that can be correlated with a change in the magnetic domain configuration HRTEM chemical and morphological analysis show gallium ion implantation and amorphization in the cobalt around the holes, that could have turned it into non-magnetic. Then, the effective diameter of the non-magnetic area becomes greater than the antidot one itself. As antidot-antidot distance decreases, the magnetic percolation is lost and a crossover from antidot to dot regime is encountered.

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#### GH01

#### Bio-functional magnetic nanoparticles in biomedical applications Herng-er Horng\*

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With distinctive characteristics of magnetic properties and a variety of bio-functions by conjugated some specific bio-probes, bio-functional magnetic nanoparticles (MNPs) have been widely investigated in the biomedical applications currently. In this report several applications using MNPs will be mentioned as examples of such applications. For the immunoassay, an innovative method, ImmunoMagnetoReduction (IMR)(1), will be introduced at first. Secondly, the cancer images by using low field MRI(2-3) will be demonstrated to reveal the possibility of this instrument in the cancer study Next, the hyperthermia with MNPs, including metabolism of MNPs study(4), the instrument and the temperature manipulation etc., will be declared. Finally, some other examples, such as magnetofection(5) etc., will be alao reviewed in this talk.

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#### **GH02**

Spin resolved measurements of single molecular magnets on surfaces Jens Brede1\*, Jorg Schwobel1, Regis Decker1, Andrew Dilullo2, Germar Hoffmann1, Svetlana Klyatskaya3, Mario Ruben3 and Roland Wiesendanger <sup>1</sup> Institute of Applied Physics, University of Hamburg, Germany <sup>2</sup> Department of Physics and Astronomy, Ohio University, Germany <sup>3</sup> Institute of Nanotechnology, Karlsruhe Institute of Technology, Germany

The use of magnetic molecules opens a gateway to a flexible design of novel spintronic devices to store, manipulate, and read spin information at the nanoscale. Crucial is the precise knowledge of molecular properties at the interface towards an electrode. Progress in this field relies on resolving and understanding the physics at the relevant interfaces. In particular the role of individual molecular constituents and the impact of the atomic environment on molecular properties, determine device relevant parameters, such as conductance and spin polarization. Here, we applied spin-polarized scanning tunneling microscopy to resolve the physics of the molecule-ferromagnet interface. The analysis focuses on different phthalocyanine molecules. The phthalocyanine constitutes of an organic macrocyclic ligand and can be functionalized with various metal ions in order to modify, e.g. the molecular spin state. We will discuss the spindependent transport from magnetic surfaces through such molecules. In particular, spin-split molecular orbitals were resolved with sub-molecular spatial resolution. The magnitude of the exchange splitting is directly determined by spin-resolved tunneling spectroscopy

#### **GH03**

MgO tunnel junction magnetic field sensors at high frequencies Mustafa Arikan1\*, Matthew Carter2, Gang Xiao3 and Snorri Ingvarsson4 <sup>1</sup> Science Institute, University of Iceland, Iceland <sup>2</sup> Micro Magnetics, Inc., USA <sup>3</sup> Department of Physics, Brown University, USA 4 Science Institutte, University of Iceland, Iceland

Micron sized MgO magnetic tunnel junctions with 1.7 nm tunnel barrier were fabricated and characterized at DC and high frequencies to be used as magnetic field sensors. Both DC and AC properties of the sensors were measured by tunneling magnetoresistance (TMR) and impedance spectroscopy between 100 Hz and 40 MHz as a function of applied external magnetic field. Several different types of sensors were investigated: Single and multiple MTJ-sensors. In the first case, single junctions were measured. A simple RLC circuit model was applied to impedance spectroscopy results to investigate magnetocapacitance properties in order to evaluate and compare different sensing schemes such as resistive and capacitive sensing. Contrary to the previous reports in the literature we didn't observe field dependent spin canacitance despite excellent agreements in other parameters such as resistance of the junction and interface capacitance. We attribute this discrepancy to the size of our junctions and reach a conclusion that limits the applicability of the spin-capacitance concept to large area devices. Then multiple tunnel junctions in serpentine geometry were measured. Using the same circuit model, we measured non-zero magnetocapacitance in multiple junction sensors unlike single junction devices. We explain this as a result of the sensor geometry.

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#### GH04

#### LaSrVMoO<sub>6</sub>: a compensated half metal or not?

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LaSrVMoO<sub>6</sub> was first synthesized in 2004. The crystal structure is uncertain and was suggested to be cubic (Fm-3m). compensated half metal (i.e., half metallic compound with zero magnetization) is suggested. Theoretical study in 2005 indicated that it is a normal metal and ferrimagnet. In 2009, LaSrVMoO6 has been revisited by the experimental study. The crystal structure is determined to be cubic (Fm-3m). It was concluded that LaSrVMoO6 is a compensated half metal. Theoretical study in 2010 based on the cubic structure indicated that LaSrVMoO<sub>4</sub> is a ferrimagnetic half metal, in contrast to compensated half metal suggested experimentally. Recently in 2010, LaSrVMoO<sub>6</sub> was found to be orthorhombic (Pnma) and metallic. Further theoretical study in 2011 indicated that LaSrVMoO6 in disordered states for V and Mo is compensated half metal. So the studied system poses great challenge for both experiment and theory.

#### GH05

Pressure effects on the magnetic properties of Emim[FeCl<sub>4</sub>], a magnetic ionic liquid with antiferromagnetic ordering

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The magnetic ionic liquids (MILs) are a new class of materials which can favourably combine the properties of ionic liquids with the magnetism that is originated from the metal incorporated in the complex anion [1]. Most of the MILs are paramagnetic. being the Emim[FeCl4] the first one in which a long-range magnetic order was found at low temperatures [2]. The physical properties of 1-ethyl-3-methylimidazolium tetrachloroferrate Emim[FeCl4] clearly show antiferromagnetic ordering when it is solidified (in polycrystalline state at 285 K), with a Neel temperature TN ~ 3.8 K [3]. Recent Mossbauer and muon spin relaxation spectroscopies as well as neutron powder diffraction experiments [4], confirm the three-dimensional magnetic ordering. In this work, we present the effects of pressure on the X(T) and M(H) curves up to 10 kbar. It is observed that the pressure modifies the magnetic interactions, increasing the order temperature and inducing a ferrimagnetic behaviour. Additionally, from the X(T), it is observed that the solid-liquid transition is very sensitive to pressure, being above 300 K for pressures larger than 0.72 kbar.

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#### GI01

#### Magnetic proximity and spin behavior at organic semiconductor/ ferromagnet interfaces towards molecular spintronics

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Information storage bit size and computing is inching towards molecular level. Spin transport in organic semiconductors (OS) has the potential in realizing this goal in a straightforward way. The charge and spin transport at OS and ferromagnetic (FM) metal interface, although complex, has the needed ingredients for achieving to reach molecular level spintronics. Spin tunneling through a few monolayers to 30nm thick films of several types of OS, all the way to room temperature combined with spectroscopic and polarized neutron reflectometry studies show unique morphology driven signatures of weakened magnetic behavior of OS/ferromagnet interfaces. However, our recent studies show surprising magnetic interaction of the phenalenyl molecules over the surface of a ferromagnet (Co, permalloy etc), driving the molecular layer to be magnetic and coupled. In addition, the second molecular layer shows spin-filtering property, serving as a spin-sensing layer. As such the FM surface layer and the adjacent two molecular layers of phenalenyl combination demonstrates a well-defined and large interface magnetoresistance (IMR), persisting to near room temperature. With spin valve structures such as Co/Zn-methyl phenaleny/permalloy an IMR of 50% is measured at 4K which only reduced to ~22% by 250K. With this study the potential for molecular level spintronics is demonstrated.

This work was done with Karthik Raman\*1.217, Tiffany Santos, Jenny Shim, Shannon Watson, Julie Borchers, Alexander Kamerbeek, Nicolae Atodiresei, Arup Mukherjee, Tamal Sen, Predrag Lazic, Vasile Caciuc, Reent Michel, Dietmar Stalke, Swadhin Mandal, Stefan Blugel and Markus Munzenberg. Research is supported by grants from ONR, NSF and KIST-MIT program.

#### **GI02**

#### Spin specific transport properties of chiral molecules

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Spin effects normally are seen either in magnetic materials or in systems containing heavy atoms that facilitate spin-orbit coupling. We report spin-selective transmission of electrons through self-assembled monolavers of double-stranded DNA. The spin polarization exceeds 60% at room temperature. The spin filtration efficiency depended on the length of the DNA. In addition, we show that conduction through double stranded DNA oligomers is spin selective, demonstrating a true organic spin filter. The selectivity exceeds that of any known system at room temperature. The spin dependent resistivity indicates that the effect cannot result solely from the atomic spin-orbit coupling and must relate to a special property resulting from the chirality symmetry. A theoretical model indicates that the spin selectivity is related to relative narrow resonances through which the electrons are transmitted Based on the theory new type of spintronics devices are constructed. The results may reflect on the importance of spin in determining electron transfer rates through biological systems.

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#### GI03

Reversible and deterministic spin state switching of individual spincrossover molecules on a surface

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A nano-scale molecular switch composed of a single molecule mainly utilizes its conductance change by external stimuli. Adding spin switching functionality to molecular switches is the key concept for realizing molecular spintronics devices with ultimate density. For this purpose, spincrossover (SCO) molecules are promising since their spin state is also switchable between a high spin (HS) state and a low spin (LS) state, in addition to the conductance. We show scanning tunneling microscopy (STM) study of individual SCO molecules, Fe(phen)2(NCS)2 (phen = 1,10-phenanthroline). While the molecules lose their switching functionality on a metallic Cu(100) surface due to a strong coupling of the NCS-groups to the surface, reducing the interaction with the surface by introducing thin CuN layer enables them to switch. Injecting a current with a tip of STM changes the conformation of the molecules between two states, resulting in on (HS-state) and off (LS-state) of a Kondo resonance accompanied with a conductance change. Observed reversible and deterministic switching behaviors of the molecules, which are vital for memory operation, give perspective on future molecular spiontronics devices with the smallest unit.

#### **GI04**

#### Graphene-based spintronic components

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Using non-equilibrium Green's function method combined with density functional theory, we propose a complete set of basic spintronics devices including bipolar spin diode, transistor and logic gates based on zigzag graphene nanoribbons (ZGNRs). Nearly  $\pm 100\%$  spin-polarized current can be generated and tuned by a source-drain voltage in the bipolar spin diode. This transport property is attributed to the intrinsic transmission selection rule of the spin subbands near the Fermi level in ZGNRs. The bias voltage and magnetic configurations of the two-terminal ZGNR-based spin diodes provide a rich variety of ways to control the spin current which can be used to design three-terminal spin transistors and logic gates. These spintronic components make possible the manipulation of spin-polarized current, such as rectification, amplification and logic operation.

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#### **GI05**

#### Detection and manipulation of spin currents in graphene with nonmagnetic electrodes

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We present an interplay between charge currents and spin currents which allows the use of nonmagnetic contacts for both detection [1] and manipulation of spin accumulation in a paramagnet. These effects are the result of a nonlinear interaction between spin and charge transport solely given by the energy dependence of the conductivity [2], analogous to the case of thermoelectrics. A conductivity spin polarization is induced by the presence of a spin accumulation, even in a paramagnet, when the conductivity is energy dependent. We use graphene as a model system to study these effects in nonlocal spin valve devices, including its dependence on charge carrier density and spin amplification under nonmagnetic contacts and in bipolar graphene p-n junctions. The general concept is applicable to the field of semiconductor spintronics.

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#### GJ01

#### Science and technology of modern permanent magnet materials George C. Hadiipanavis\*

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Permanent magnets (PMs) are indispensable for the electric, electronic and automobile industries, information technologies, automatic control engineering and many other commercial and military applications. In most of these applications, an increase in the magnetic energy density of the PM, usually presented via the maximum energy product (BH)max, immediately increases the efficiency of the whole device and makes it smaller and lighter. Worldwide demand for high performance PMs has increased substantially in the past few years driven by hybrid and electric cars, wind turbines and other power generation systems. This talk will cover the major principles guiding the development of PMs, including the important role of microstructure on coercivity and overview state-of-the-art theoretical and experimental research. Recent progress in the development of nanocomposite PMs, consisting of a fine (at the scale of magnetic exchange length) mixture of phases with high magnetization and large magnetic hardness will be discussed. Fabrication of such PMs is currently the most promising way to boost the (BH)max, while simultaneously decreasing, at least partially, the reliance on the rare earth elements. Current efforts in the development of high performance non-rare earth magnets and their future prospects will also be discussed.

Work supported by DOE and ARPA-E

#### GJ02

**Differential thermal analysis on MnBi in high magnetic fields up to 45 T** Keiichi Koyama<sup>1</sup>\*, Yoshifuru Mitsui<sup>2</sup>, Eun Sang Choi<sup>3</sup>, Yuki Ikehara<sup>2</sup>, Eric Palm<sup>3</sup> and Kazuo Watanabe<sup>2</sup> <sup>1</sup> Kagoshima University, Japan

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Differential thermal analysis was carried out for hard magnetic material MnBi in the temperature range 300-773 K in magnetic fields up to 45 T to investigate the effect of high magnetic fields on its decomposition process and corresponding phase diagram. The decomposition temperature Tt (MnBi  $\rightarrow Mn_{1.08}Bi$  + liquid Bi) increases from 632 K (at a zero field) to 714 K by applying a magnetic field of 45 T. Furthermore, the magnetocaloric effect of MnBi is observed in 11.5-45 T in the vicinity of 689 K, showing that a field-induced composition process occurs. The obtained results show that the equilibrium state of MnBi can be controlled by a high magnetic field [1].

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#### GJ03

#### Magnetization of Dy<sub>2</sub>Fe<sub>17</sub> in fields up to 85 Tesla

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 $Dy_2Fe_{17}$  is an easy-plane rare-earth ferrimagnet. Applying sufficiently high magnetic field one can break the antiparallel alignment of the sublattice moments and drive the system towards the ferromagnetic state via intermediate non-collinear phases. In this particular case, if the magnetic field is applied within the easy plane, the moment direction of the Fe sublattice deviates only slightly from the field direction. The Dy-moment instead starts from the direction opposite to the field and the Fe moment, and subsequently passes several energetic minima. This manifests itself in corresponding jumps on the magnetization curve. Measurements in fields up to 67 Tesla have shown a field-induced jump on the magnetization curve for field aligned along the a axis. According to a theory developed for hexagonal ferrimagnets and having the field and the jump height along the a axis, we expected the jump along the b direction to appear at 78 T. To check this prediction, we have measured the high-field magnetization curve of a  $Dy_2Fe_{17}$  single crystal for fields up to 85 Tesla along the b axis . We have observed a sharp magnetization jump at 77.8 Tesla, which agrees perfectly with the theory.

#### GJ04

#### Electrodeposited FePt films on Ag underlayer with high coercivity Sirikanjana Thongmee

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FePt films with a composition close to  $Fe_{30}Pt_{30}$  were fabricated following with post annealing. The films were deposited onto Si substrate with a under layer of Ag. The phase formation temperature, microstructure evolution by annealing and the magnetic properties were investigated in detail. The L10 FePt phase started to form after annealing at 500 °C. From TEM results, as-deposited FePt films consisted of randomly oriented nanograin. The average grain sizes or FePt films on Ag underlayer were about 50 nm after the film was annealed at 800 °C and SAED result was consistent with XRD. High coercivity of 15 kOe was obtained after annealing at 600 °C. After annealing at 800 °C, the coercivity increased to 18 kOe with a significant perpendicular anisotropy of FePt films. Out of plane anisotropy was observed in relatively thick film (800 nm). The work has shown that FePt film with a high coercivity can be achieved by electro-deposition and the thickness, anisotropy can be tuned, which may provide a method to fabricate FePt film for the potential applications both recording media and MEMS.

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#### GJ05

#### Magnetic properties of BaMg<sub>0.4</sub>Al<sub>0.4</sub>Fe<sub>11.2</sub>O<sub>19</sub>+SiO<sub>2</sub> nanocomposites for high frequency applications K Sadhana and K Praveena

Materials Research Centre, Indian Institute of Science, Banglore-560012, India

The nanomcomposites of Ba<sub>0.5</sub>Mg<sub>0.4</sub>Al<sub>0.4</sub>Fe<sub>112</sub>O<sub>19</sub>/SiO<sub>2</sub> were prepared using Microwave-Hydrothermal method at 160°C/45min. The as synthesized composites were characterized using X-ray diffraction and Transmission Electron microscopy (TEM). The average particle size is found to be 62nm. The present composites were densified using microwave sintering method at 900°C/60min. The frequency dependent complex permeability of the composites were measured in the range of 1MHz to 1.8GHz. The real part of permeability of the composites found to decreases with an addition of SiO<sub>2</sub> but the resonant frequency is shifted towards higher frequency side. The magnetic properties such as saturation magnetization and corecive field of sintered composites were calculated based on M-H curves. A possible relation between the magnetic hysteresis curves and the microstructure of the sintered composites was investigated.

#### HA01

Coexistence of competing orders in unconventional superconductors Setsuko Tajima\*, E. Uykur, K. Tanaka, T. Masui and S. Miyasaka Dent. of Physics. Osaka University. Japan

In my talk, I will focus on the problem of the coexistence of competing orders which is probably inevitable in strongly interacted systems. A typical example is the pseudogap state in high Tc superconducting cuprates, which has been a long-standing puzzle In order to address this problem, we have carefully measured the c-axis polarized optical spectra of YBa2(Cu,Zn)2Oy. Removing an additional spectral feature due to the transverse Josephson plasma by Zn-doping, we could unambiguously discuss the spectral weight transfer. The result clearly showed that the pseudogap originates from some competing order but not a precursor of superconductivity. Moreover, we found that the pseudogap persists even below Tc, which becomes pronounced by Zn-substitution. This indicates that the pseudogap and the superconducting gap are coexisting, presumably in a phase separated form. The most interesting problem is whether or not such a coexistence of competing order plays a positive role in superconductivity mechanism. One of the interesting facts is that the onset temperature for superconducting fluctuation increases with decreasing doping. Its energy scale is unusually large (~ 0.4eV). This suggests that the competing order with a large energy scale contributes to superconducting fluctuation and possibly to pairing mechanism.

#### HA02

## Ultrafast transient generation of spin density wave order in the normal state of BaFe,As, driven by coherent lattice vibrations

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The proximity of magnetic ground states to the superconductivity in cuprates and pnictides as well as in other unconventional superconductors has put the magnetism at the heart of the discussion concerning the mechanism of high temperature superconductivity. The recent development of the pump-probe time domain spectroscopy, which was typically limited in the terahertz region below 10 meV, has extended its spectral window to the infrared (IR) region allowing a resonant probe of important elementary excitations such as phonons, excitations across a superconducting gap as well as a density wave gap. We study the ultrafast dynamics of the spin density waves (SDWs) in BaFe<sub>2</sub>As<sub>5</sub>, a parent compound of pnictide high temperature superconductors. We monitor the SDW gap resonantly by the near-infrared (NIR) pump-IR probe technique. In the SDW state, we observe that the NIR excitation of the carriers breaks the SDW order. The excitation with yet higher fluence launches strong coherent lattice vibrations. Stithus(p), these coherent oscillations give rise to a transient SDW order even in the normal state. We discuss that the strong electron-phonon/spin-phonon coupling effects in pnictides result in the monulations of the SDW order parameter in real time in response to the coherent phonon vibrations.

K. W. Kim et al., Nature Materials (2012, to be published)

#### HA03

## High- and low-energy ARPES study of spin-density wave order in FeTe single crystals and FeTeOx films

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We have performed a high-resolution (synchrotron- and laser-based) ARPES investigation of FeTe single crystals films, as well as thin films of the novel superconductor FeTeOx [1]. Our results from the single crystals reflect the previously reported Fermi surface pocket around the X-point [(r, 0)], possibly connected to a spin-density wave (SDW) order [2]. Unlike such previous report, our results also reveal the presence of an energy gap ( $\Delta$ ), which would be expected from the SDW order. The temperature dependence shows that  $\Delta$  closes in the vicinity of the magnetic transition temperature, supporting its interpretation as reflecting the SDW state. We also present bulk sensitive results acquired using a soft x-ray ARPES technique. Photon-energy (i.e. kz) dependent measurements reveal a strongly 3D electronic structure that can be uniquely connected to the SDW-gap symmetry found in the low-energy data. Our results clearly illustrate the importance of acquiring not only the true intrinsic bulk-electronic structure but also that it is imperative to consider the full 3D band dispersion in iron-based superconductors and their parent compound in order to obtain an adequate theoretical understanding. Finally, we also acquired novel ARPES data from cleaved FeTe and FeTeOx thin films.

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#### HA04

## Magnetic fluctuations - a driving force for superconductivity Neutron scattering investigations in Fe-based superconductors.

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One of the most exciting findings in investigations of High-Tc superconductivity has been the discovery of a magnetic resonance in the superconducting state, associated with a spin excitation with a characteristic position in wavevector and energy space [1]. This resonance has been investigated across the families of Fe-based superconductors with several common features observed such as the scaling of the resonanc energy with Tc (Er ~ 4.3 Tc in (Ba,K)Fe<sub>2</sub>As<sub>2</sub> [1]) and the position of the resonance either at, or incommensurate about, the antiferromagnetic wavevector [1,2,3]. This work presents results of inelastic neutron scattering experiments on a broad range of Fe-based superconducting materials in order to address the question - are these features really common to all these materials? Results from the characteristic 122-arsenide family [4] are contrasted with LiFeAs [2], CsxFe<sub>2-a</sub>Se<sub>2</sub> and phosphorus containing materials.

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#### **HA05**

#### Low-energy quasiparticles probed by heat transport in the iron based superconductor lafeno

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Measurements of ultra-low temperature heat transport are an extremely effective probe of the presence of nodes in the superconducting order parameter. In this study we report the thermal conductivity of the iron pnictide superconductor LaFePO down to temperatures as low as T=60mK and in magnetic fields up to 5 T. The data shows a large residual contribution that is linear in temperature, consistent with the presence of low-energy electronic quasiparticles [1]. We interpret the magnitude of the linear term, as well as the field and temperature dependence of thermal transport in several pairing scenarios, with the aim of distinguishing between accidental or symmetry imposed nodes. The presence of an unusual supralinear temperature dependence of the electronic thermal conductivity in zero magnetic field, and a high scattering rate with minimal Tc suppression argues for a sign-changing nodal s+/- state. We compare these results to recent thermal transport studies in related iron-based superconductors.

M. Sutherland et. al. Phys. Rev. B 85, 014517 (2012)

#### **HB01**

#### Domains and magnetization processes in electrical steel

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Understanding and lowering of energy loss in electrical steel requires knowledge about the magnetic domains and magnetization processes. In this presentation some aspects of domains and processes will be addressed, based on domain analysis by Kerr microscopy and neutron imaging. The former provides information on the surface domains and requires the support of domain theory to come up with conclusive models. Examples will be shown for oriented and non-oriented material, including the effects of mechanical stress [1]. The volume domains, which are relevant for key magnetic properties like susceptibility, energy loss, noise or coercivity, can be analysed by the novel method of neutron dark-field microscopy [2,3]. By this grating interferometry technique, the domains and magnetization propagation in bulk FeSi material can be imaged non-destructively and even tomographic analysis becomes possible [4] with a lateral resolution down to 10 um. The method has the unique potential to visualize the magnetic field penetration in bulk material [3].

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#### HB02

#### Iron loss behaviors in 6.5 wt% grain-oriented silicon steel

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6.5 wt% silicon steel has great potential for electrical appliances and devices, such as transformer and motor cores, due to its high electrical resistance, extremely low magnetostriction, low magnetocrystalline anisotropy, and low iron loss [1-2]. However, the presence of ordered phases, i.e. B2, D03, deteriorates the workability of the 6.5 wt% silicon steel, which impedes the mass production using the conventional rolling technique [2]. In addition, there is a lack of research on the effect of the ordered phases on the magnetic behaviors of high silicon steels. In this paper we therefore investigated the relationship between the ordered phases and the iron loss behaviors in high silicon steels using transmission electron microscopy and Kerr microscopy. In order to obtain the high silicon steels, 3 wt% grain-oriented silicon steels inserted between SiO<sub>2</sub> textiles were annealed for the decomposition of the textiles and the diffusion of Si. The size, shape and amount of the ordered phases were controlled by thermo-mechanical treatment. The microscopical analyses showed that the antiphase boundaries of the ordered phases cut magnetic domains, which strongly affected the iron loss behaviors of the high silicon steels.

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#### HB03

#### FeCoB films with large saturation magnetization and high magnetic anisotropy field to attain high ferromagnetic resonance frequency Shigeki Nakagawa

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FeCoB films prepared on a Ru underlayer using the Oblique incidence of sputtered and backscattered particles have a high in-plane magnetic anisotropy field of 500 Oe Facing targets sputtering system is suitable to attain such deposition condition. In-plane X-ray diffraction analysis clarified that there is anisotropic residual stress, which is the origin of the high in-plane magnetic anisotropy. The directional crystalline alignment and inclination of crystallite growth were also observed. Such anisotropic crystalline structures may affect the anisotropic residual stress in the FeCoB laver. The B content of 6 at.% was appropriate to induce such anisotropic residual stress. The anisotropic residual stress develops according with the film growth. The film with large saturation magnetization of 22 kG and high magnetic anisotropy field of 500 Oe exhibited high ferromagnetic resonance frequency of 9.2 GHz. The large saturation magnetization and the low coercivity along the hard axis direction assure that the film can be operated in the high frequency range up to 8 GHz with maintaining high relative permeability.

#### HC01

#### Magnetism where you least expect it Priva Mahadevan

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The picture of magnetism that has prevailed over the years has centered around the existence of localized electrons and their ordering leading to different types of magnetic order. Recently several unconventional examples of magnetism have been emerged such as p-band magnetism in alkali metal oxides - a system with no traditional magnetic ions in them. Unfortunately high magnetic ordering temperatures were not possible and this was traced to an orbital ordering transition [1]. While orbital ordering is usually found to reinforce one type of magnetic order over another, in the present case one found it led to the loss of any possibility of high magnetic ordering temperature. In another example, we consider the example of a 4d oxide SrTcO<sub>3</sub> where recent experiments [2] have found a Neel temperature of 1023 K far higher than its 3d counterpart SrMnO<sub>2</sub> A detailed analysis [3] within a microscopic multiband Hubbard model reveals a rather surprising result of higher bandwidth and smaller Hund's exchange coupling strength resulting in higher magnetic ordering temperatures in 'SrTcO3-type' oxides. This is work done in collaboration with A.K. Nandy, S. Middey, P. Sen and D.D. Sarma.

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#### HC02

#### Exotic magnetism of s-electron cluster array: Ferromagnetism, ferrimagnetism and antiferromagnetism

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Alkali metal nanoclusters can be stabilized in the regular cages of zeolite crystals by the loading of guest alkali metals. Cages are connected by the sharing of windows of aluminosilicate framework, and arrayed in simple cubic, diamond and body centered cubic structures in zeolites A, X and sodalite, respectively The s-electrons have the localized nature of nanoclusters with magnetic moments and the mutual nteraction through windows. They show exotic magnetism depending on the structure type of zeolite alkali metal and the average loading density of alkali atoms per cage, n. In zeolite A, potassium clusters are formed in alpha-cages. They exhibit spontaneous magnetization at  $n \ge 2$ . Ferromagnetic properties are explained by the canted antiferromagnetism of the Mott insulator, where the canting provided by Dzyaloshinsky-Moriya interaction is strongly enhanced by the 1p-like degenerate molecular orbitals of clusters[1] and the superlattice structure. Rb clusters generated at alpha- and beta-cages of zeolite A exhibit ferrimagnetic properties[2]. Na-K alloy clusters generated at supercages and beta-cages of lowsilica X zeolite exhibit ferrimagnetism at specific loading densities of alkali metals, where clusters are in the metallic phase[3]. Alkali metal clusters in sodalite show the Heisenberg antiferromagnetism of the Mott insulator[4]

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#### HC03

#### Spin-dependent molecular arrangement of O<sub>2</sub>-O<sub>2</sub> dimer in nanoporous metalorganic solids

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O2 is a molecular magnet with S = 1. The characteristic magnetic properties related to the molecular arrangement can be expected since the magnetic exchange interaction between O<sub>2</sub> is comparable to the electric interaction between the quadrupole moments. [1] Recently, for a model system of isolated O2-O2 dimer confined in the nanopores of Cu-cvclohexanedicarboxvlic acid, the molecular arrangement structure relating to its magnetic properties was investigated by precise x-ray diffraction analysis. The obtained charge density level structures revealed that the significant orientational change of the molecular axes occurs with changing temperature, coupling with the spin states of O2-O2 dimer. The temperature dependence of magnetic susceptibility and the magnetization process characterized by the absence of half-plateau (S\_dimer = 1) can be consistently explained by a simple model taking into account the orientational change of molecular axes. Consequently, the present results experimentally proved that the molecular orientation depends on the spin state, the rectangular parallel H-geometry for the singlet ground state (S dimer = 0) and the shifted parallel S-geometry for the thermally excited quintet state (S\_dimer = 2), due to the spin-dependent intermolecula potential which has not been discussed for solid O2.

R. Kitaura et al. Science 298, 2358 (2002).

#### HC04

#### Indications for a field-induced 2d collectively-coupled dimer state in nitronyl-nitroxid biradicals

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Arrays of coupled spin S = 1/2 dimers, exposed to a magnetic field, are of general interest as they can be considered as the source for generating a gas of interacting bosons [1]. We present a novel coupled spin-dimer system with a double-layer structure based on nitronyl-nitroxid stable radicals bridged via tolan molecules. The intra-dimer coupling constant is around 10 K, whereas the inter-dimer couplings, mediated via H-bonds, are of the order of 1 K. The moderate size of these couplings enables us to study the field-induced ordered phase forming for T < 0.3 K and 8 T  $\leq$  B  $\leq$  9.5 T. The anomaly observed in the ac-susceptibility  $\chi(B)$  differs significantly from ordinary B-induced 3d ordering, tantamount to Bose-Einstein condensation of magnetic excitations, where  $\chi(B)$  shows sharp peaks at the on- and offset fields. Rather we find a rounded double-peak structure where the peaks remain rounded down to 75 mK, the base temperature of our experiment. These features indicate the formation of a 2d collectively-coupled dimer state, implying topological order for  $T \rightarrow$ 0. In order to analyze the nature of the B-induced ordered state, the experimental data are compared with the results of Ouantum-Monte-Carlo calculations.

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#### HD01

#### Magnon caloritronics

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Eigen excitations of a magnetic system, e. g. spin waves and their quanta, magnons, are expected to play an important role in spin-caloric effects. The recently observed spin Seebeck effect in a magnetic insulator demonstrates the strong connection between magnon and heat transport processes. I will address the interplay between magnon and heat transport phenomena and present recent Brillouin light scattering experiments from magnons in an yttrium iron garnet (YIG) waveguide subject to a temperature gradient. Two issues are addressed in a thermal gradient field: first the propagation of spin waves and second the distribution of thermal magnons. We find that first magnons adapt their wavelength in a thermal gradient field and stopping effects may occur. Second, we report direct measurements of the magnon distribution and show, that it does not follow the phonon temperature.

#### HD02

Temperature dependence of spin wave resonance frequency in a magnetostatic surface wave mode

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Propagating spin waves in a 20 nm NiFe film, in the magnetostatic surface wave (MSSW) mode, have been measured at various temperatures in a time domain. Pulse inductive microwave magnetometry (PIMM) is used for the detection of spin wave packets. Impulse voltages were applied to the excitation conlanar waveguide (CPW) by a pulse generator, while the resultant dynamic responses were observed through a sampling oscilloscope at the detection CPW. The spin wave resonance frequencies (fR) increases with decreasing temperature. The saturation magnetization (MS) increases at low temperature and consequently fR increases in accordance with the dispersion relation of MSSWs. The change in fR from 300 K to 12 K is approximately 400 MHz at a bias field of 280 Oe. A second mode is also observable below 100 K, and becomes progressively clearer at lower temperatures. At higher temperature, there is significant magnon-phonon interaction. However, the lattice vibrations become subsided at low temperatures, and magnon-phonon interaction reduces, while electron-electron scattering decreases. This causes a change in the nature of the damping phenomenon. The nature of the intrinsic damping has been investigated as a function of temperature and has been correlated to the difference in the damping mechanisms.

#### HD03

#### Nanoscale spin wave switches and phase shifters

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The short wavelength of spin waves propagating in nanoscale magnonic waveguides at microwave frequencies could potentially provide a device miniaturization opportunity via accomplishing certain tasks with a reduced number of physical devices as compared to semiconductor circuitry. To utilize this promise of magnonics, it is important to learn to manipulate the spin wave phase and perhaps magnitude, which is highly related to data encoding and logical operation, etc. Here, we use OOMMF simulations to demonstrate a new method of achieving controlled phase shift of spin waves that relies on the reverse of the recently discovered spin wave emission mechanism [1]. The method could be highly useful for magnonic applications since it depends heavily on the static magnetization of the control element in the structure. Hence, by changing static magnetization direction in the element, it is possible to control the magnitude and phase of propagating spin waves. These switchable, non-volatile devices could serve as either spin wave valves or phase shifters, all at the nanoscale and low energy cost. The research leading to these results has received funding from the EC's Seventh Framework Programme (FP7/2007-2013) under GAs 233552 (DYNAMAG) and 228673 (MAGNONICS) and from the EPSRC of the UK (EP/E055087/1).

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#### HD04

#### Optically induced tunable magnetization dynamics in nanoscale Co antidot lattices

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We present the fabrication of high quality Co antidot lattice structures with 100 nm antidot diameter and with varying lattice constants S between 200 and 500 nm by using focused ion beam lithography. We have excited and detected the magnetization dynamics in those antidot lattice structures by an all-optical timeresolved Kerr microscope. The dynamics shows two prominent magnonic bands for sparsely packed lattices with a clear bandgap. The bandgap increases with the decrease in the lattice constant and at the lowest value of the lattice constant S = 200 nm, four distinct magnonic bands appear. The observations are qualitatively reproduced by micromagnetic simulations and the resonant mode profiles are calculated and interpreted. The lower frequency mode corresponds to the magnetostatic surface wave mode while the higher frequency modes are quantized modes with varying quantization number. We further demonstrate the tunability of magnonic spectra in composite structures constructing of antidot lattices with varying lattice constants arranged in varying geometry. The observations are important for potential applications of the antidot lattices in nanoscale magnonic crystals in the form of composite antidot structures with tunable bandgaps.

I. Kruglyak, V. V.; Demokritov, S. O.; Grundler, D., Magnonics. J. Phys. D: Appl. Phys. 2010, 43, 264001. 2. Neusser, S.; Duerr, G.; Bauer, H. G.; Tacchi, S.; Madami, M.; Woltersdorf, G.; Gubbiotti, G.; Back, C. H.; Grundler, D., Anisotropic Propagation and Damping of SpinWaves in a Nanopatterned Antidot Lattice. Phys. Rev. Lett. 2010, 105, 067208. 3. Tacchi, S: Madami, M.; Gubbiotti, G.; Carlotti, G.; Adeveve, A. Q.; Neusser, S.; Botters, B.; Grundler, D., Angular Dependence of Magnetic Normal Modes in NiFe Antidot Lattices with Different Lattice Symmetry. IEEE Trans. Magn. 2010, 46, 1440-1443

#### HD05

#### **Plasmonic and quantum plasmonic enhancement of magneto-optics** Alexey P. Vinogradov<sup>1</sup>\*, Denis G. Baranov<sup>1</sup> and Alexander A. Lisyansky<sup>2</sup>

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We consider two scenarios of magneto-optical transmission lines build on the basis of an array of plasmonic nanoparticles. The first scenario corresponds to the wellknown transmission line of the plasmonic nanoparticles embedded into a magnetooptical medium. We show that the propagation of the dipole-plasmonic mode travelling along the array is accompanied by the rotation of the polarization (the Faraday effect). The angle of rotation is dozens times greater than that in the uniform magneto-optical material. In the second case, we consider a chain of metallic nano-particles exhibiting magneto-optical properties embedded into a dielectric matrix. When the matrix is passive, in spite of the plasmonic resonance, the Faraday rotation enhancement is small due to the small propagation length of the dipole-plasmonic mode. In an active matrix, the metal nanoparticles are surrounded by active centers (quantum dots or active molecules) forming magneto-optical spasers. The dipole-plasmonic mode in such an array of spasers exhibits high values of the Faraday rotation and the propagation length.

#### HE03

#### Extrinsic SHE induced by small impurities in copper Yoshichika Otani\*, Yasuhiro Niimi, Yohei Kawanishi and Dahai Wei ISSP, University of Tokyo, Japan

Spintronic devices manipulating pure spin currents, flow of spins with no net charge flow, must play an important role in low energy consumption electronics of next generation. This explains the current interest for the spin Hall effect (SHE) which provides a purely electrical way to generate spin currents without ferromagnets and magnetic fields. The inverse SHE (ISHE) is the only means to generate electrical signals from spin currents. It is indispensable for practical applications to have a large spin Hall (SH) angle which characterizes the efficiency in conversion between charge and spin currents. Today platinum is adopted as the best material for this conversion. The ISHE of Pt, for instance, has been used for the detection of spin currents produced from heat flows by the spin Seebeck effect. Thus it is a prime task to find materials with large SH angles and preferably lower cost than Pt. Here we discuss how a giant SH angle (-12%) can be obtained by simply doping copper with a small amount of bismuth impurities.

#### HE01

## Pure spin current generation using highly spin polarized Co<sub>2</sub>FeSi electrodes

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Pure spin current provides attractive means for the efficient operation of spin-based electronic devices. Efficient generation of the pure spin current may be a key for developing the future spin-based devices. However, the generation efficiency of the pure spin current using the conventional ferromagnetic metal electrodes such as NiFe and Co is extremely low. This is because most of the injected spin currents return back to the spin injector with the low spin polarization. Such back flows of the spin currents can be effectively reduced by using highly spin polarized ferromagnetic electrodes. In the present study, we fabricated lateral spin valves with Co<sub>2</sub>FeSi electrodes, which is one of the Heusler compounds expecting half metalicity. Large nonlocal spin valves signals were observed at room temperature, revealing the high spin polarization of the Co<sub>2</sub>FeSi electrodes. From the analysis based on one-dimensional spin diffusion model, the generation efficiency of the pure spin current is found to be significantly enhanced compared to those with the conventional ferromagnetic electrodes. [1] Moreover, from the systematic study of numerous number of the devices, we found the correlation between the electrical resistivity and the spin polarization. [2]

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#### HE02

## Highly reproducible lateral spin valves for the study of spin injection in metals

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Creation and transport of pure spin currents is a key ingredient in spintronics, which can be achieved by electrical spin injection using lateral spin valves (LSV) [1]. We fabricated highly reproducible LSV devices by using a two-step electron-beam lithography process. In the first step, ferromagnetic (FM) electrodes (Co or Py) were deposited and, in the second one, the non-magnetic (NM) metal (Cu, Au or Ag) was put on top. In these devices, the spin signal was measured at different distances and temperature (T), from which the spin polarization of the FM ( $\alpha$ ) and the spin diffusion length of the NM ( $\lambda$ ) are obtained as a function of T. Whereas  $\lambda$ (T) shows the role of surface spin-flip scattering [2] in the various NM, the results in  $\alpha$ (T) for the various FM metals help understanding the role of the FM/NM interface in the electrical spin injection. Furthermore, due to the high reproducibility of our LSV devices,  $\lambda$ (T) for Co can be independently obtained. The latter parameter has been controversial for its anomalously large value [3]. These results are relevant to control the role of FM metals in electrical spin injection and of NM metals in spin transport.

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#### HE04

#### Spin Injection at the LaAlO<sub>3</sub>/SrTiO<sub>3</sub> Interface

Nicolas Reyren, Manuel Bibes, Edouard Lesne\*, Jean - Marie George, Cyril Deranlot, Sophie Collin, Agnes Barthelemy and Henri Jaffres Unite Mixte de Physique CNRS/Thales, France

Future spintronics devices will be built from elemental blocks allowing the electrical injection, propagation, manipulation and detection of spin-based information. Owing to their remarkable multifunctional and strongly correlated character, oxide materials already provide such building blocks for charge-based planar devices such as ferroelectric field-effect transistors, as well as for spin-based vertical devices like magnetic tunnel junctions, with giant responses in both cases. In an attempt to bridge these two areas, we report results of electrical spin injection at the high-mobility quasi-two-dimensional electron system (2-DES) that forms at the LaAlO<sub>3</sub>/SrTiO<sub>3</sub> interface [1]. In a non-local, three-terminal measurement geometry, we analyze the voltage variation associated with the precession of the injected spin accumulation driven by perpendicular or transverse magnetic fields (Hanle and inverted Hanle effect) [2]. The influence of bias and back-gate voltages reveals that the spin accumulation signal is amplified by resonant tunneling through localized states in the LaAlO<sub>3</sub> strongly coupled to the 2-DES by tunneling transfer.

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#### HE05

#### Coherence in collective spin precession in lateral spin valves

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Pure spin current is a diffusive flow of spins with no charge flow, which has drawn much attention worldwide due to spintronic application as well as fundamental interests. Non-local spin injection in a lateral spin valve generates a pure spin current, of which polarization direction can be modulated via Larmor precession [1]. During the diffusive transport between injector and detector, however, spins lose phase coherency with frequent collisions, because these events lead to a broad distribution of dwell time. Here we show that the lateral spin-valves with dual injector Permalloy/MgO/Ag enable us to detect highly coherent spin precession over a distance of 10 µm with keeping spin accumulation vector in plane against out of plane magnetization process, which results in genuine signal of the in-plane precession. We also experimentally found that the better the phase coherency becomes the longer the spins travel. This tendency appears to fall on a material independent universal curve when the coherence is plotted against the reduced parameter of the injector-detector distance over spin diffusion length. This is useful for the material design of spintronics devices based on pure spin current.

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#### HF01

#### Bias-dependence of the spin-transfer torques in MgO-based magnetic tunnel junctions

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Understanding the bias dependence of the two spin-transfer torques (IP in-plane and OOP out-of-

plane) in Magnetic Tunnel Junctions (MTJs) is important from the fundamental point of view and for applications. For example, back-hopping phenomena at large bias, unwelcome for MRAM applications, has been ascribed to a competition between IP and OOP torques [1]. There is still a debate on the expected bias dependence of both torques. In symmetric tunnel junctions (same thickness and composition for both ferromagnetic electrodes), some results report a OOP torque linear with bias [2], and others quadratic [3]. It has also been shown that by voluntarily introducing an asymmetry in the MTJ, it was possible to control the sign of the linear part of the OOP torque [4,5]. The presented work is a systematic study of the bias-dependence of the two torques in MTJs with different compositions for the electrodes, ranging from symmetric CoFeB/MgO/CoFeB to asymmetric CoFeB/MgO/NiFe MTJs. We will present the switching plase diagrams of the different MTJs, from which we infer the evolution of the two torques by comparing our experimental results to numerical simulations. Financial support by JSPS Postdoctoral Fellowships For Research Abroad and the ERC 2010 Stg 259068 is acknowledged.

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#### HF02

## Spin torque assisted magnetization switching in thermally activated region

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Thermal stability is an important quantity of ferromagnets because it determines performances of spintronics devices. The value of the thermal stability has been determined by analyzing the spin torque switching in the thermally activated region with the switching probability formula,  $P=1-\exp[-ft\exp[-\Delta(1-I/Ic)^b]]$ , where  $\Delta$ , I, and Ic are the thermal stability, current, and spin torque critical current at zero temperature, respectively. The value of the exponent, b, has been assumed to be unity. However, recently, several authors, including us, argued that b=2. The determination of b is important because it significantly affects the estimation of the thermal stability. Recently, Heindl et al. (JAP 109) experimentally found that the current pulse duration time dependence of the switching current shows nonlinearity, while the theory with b=1 predicts linear dependence. They concluded that the discrepancy arises from the breakdown of the thermal model. In this study, we theoretically studied the spin torque switching in the thermally activated region. We employed the mean first passage time approach which reduces the exact solution of the switching probability in the whole range of the thermally activated region and gives b=2 in several limits. Our theory reproduced the nonlinearity of the switching current obtained by Heindl et al.

#### HF03

Joule heating and spin-transfer torque investigated on the atomic scale Stefan Krause\*, Gabriela Herzog, Anika Schlenhoff, Andreas Sonntag and Roland Wiesendanger

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Reorienting the magnetization of a nanostructure by injecting a spin-polarized current is in the focus of ongoing research because of its relevance for future spintronic and magnetic memory devices. Recent experiments using spin-polarized scanning tunneling microscopy (SP-STM) demonstrated current-induced magnetization switching (CIMS) across a vacuum barrier, driven by Joule heating, spin-transfer torque and Oersted field [1,2]. In our study, we specially utilize a superparamagnetic Fe/W(110) nanoisland to explore CIMS with SPSTM. By simultaneously observing and manipulating its switching behavior with a spin-polarized tunnel current, we separate and quantitatively determine the individual contributions of Joule heating and spin-transfer torque, with the nanoisland serving as a combined local thermometer and spin-transfer torque analyzer. Comparing our results to experiments performed on nanopillar MTJs reveals a very high spin-transfer torque efficiency for SP-STM MTJs. Our studies allow for a detailed investigation of Joule heat generation and spin-transfer torque switching on the atomic scale, thereby providing new insight into the details of CIMS.

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#### HF04

#### Spin-transfer torque and joule heating of field-emitted electrons

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Applying a high electrostatic field results in the emission of electrons from a solid surface into vacuum. These field-emitted electrons are used for microscopy purposes [1]. For magnetic emitters the spin polarization of the field-emission current has been detected since the 1960's [2]. However, the microscopic details of the interaction of spin-polarized field-emitted electrons with magnets remain to be discovered. Here, we utilize spin-polarized scanning tunneling microscopy (SP-STM) in the field emission mode [3] for the direct observation and manipulation of atomic-scale superparamagnets. Injecting high spin-polarized currents into the first field emission resonance (FER) of individual magnets reveals the switching behavior being strongly affected. Telegraphic noise experiments allow for a current-dependent lifetime analysis [4], thereby quantifying the spin-transfer torque and Joule heating of the field-emitted electrons. Switching the magnetization of quasistable nanomagnets is realized by ramping field emission currents. Our experiments demonstrate the capability of field-emitted electrons for magnetic observation and manipulation on the atomic scale at nm distances.

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#### HF05

## Perpendicular spin torque at high bias in MgO-based magnetic tunnel junctions

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This talk will consist of two parts. One is an experimental study on bias-dependence of perpendicular spin torque in MgO-based magnetic tunnel junctions (MTJs) and the other is a numerical study on current-induced self-resonant switching of magnetization. In the first part, we show an experimental method to estimate the bias dependence of the perpendicular spin torque in MTJs based on the switching phase diagrams [1]. A special attention is paid for high-bias range. Our method utilizes the pulse-width-dependent change in the switching phase diagram and addresses the perpendicular-torque-driven precessional switching at high-bias ranges where the in-plane spin torque at sa an additional damping. The bias dependence of the perpendicular spin torque at high bias was found to be linear for a voltage polarity. This method will be useful to address spin torque effects in MTJs at high bias. In the second part, we propose a new type of current-induced magnetization switching in multilayers. It utilizes spin torques caused by DC current injection and consequent RF fields caused by magnetization dynamics. We show that this way of switching magnetization is more efficient than the conventional ways. It allows a large reduction in the switching current with keeping the thermal stability.

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#### HG01

#### Recent progress in spin sem

Kazuyuki Koike

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Spin-polarized scanning electron microscopy (spin SEM) is a magnetic domain observation method. It has excellent capabilities, including high spatial resolution, high sensitivity to a thin or even a monolayer film, a large dynamic observation range, and detection of magnetization direction. Since the initial development of spin SEM by Koike et al. in 1984[1], its capabilities have been greatly improved. The spatial resolution is now 3 nm[2], the sample can be cooled down to 10 K, and images of three magnetization components and surface topography can be obtained simultaneously[3]. Moreover, images of elements and crystallographic orientation can be obtained for the same area as the magnetic domain image. Taking advantage of these capabilities, we have studied recorded bits of high density magnetic recording media, magnetic coupling between a ferromagnetic Fe thin film and a single-crystal NiO antiferromagnet[4], a patterned magnetic island array[5], and the temperature-dependent spin structure of a strongly correlated magnetic material: La<sub>22x</sub>Sr<sub>1+2x</sub>Mn<sub>4</sub>O<sub>7</sub> (x = 0.30, 0.32)[6][7]. In this talk, I will discuss the above topics in detail, and demonstrate the applicability of spin SEM to various fields related to magnetism.

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#### **HG02**

#### Asymmetries in the formation process of magnetic vortex states in a permalloy nanodisk

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Magnetic vortices in nanopatterned elements are attracting increased interest because they offer fascinating topological spin structures and great potential as novel concepts for data storage technologies [1]. A-priori, one would assume that a magnetic vortex is degenerate with four energetically equal ground states and therefore the formation of a magnetic vortex state (VS) should exhibit perfect symmetry. In our work, the formation process of VSs in Ni<sub>80</sub>Fe20 nanodisks has been investigated by direct imaging of in-plane and out-of-plane magnetic components in vortex structures utilizing high resolution magnetic transmission soft X-ray microscopy (MTXM). We observe a symmetry breaking in the formation process of VS in individual nanodisks. Micromagnetic simulations confirm that the Dzyaloshinskii-Moriya interaction on the full surface of nanodisks and surface-related inhomogeneities cause the observed asymmetric phenomenon [2]. In addition, we have expanded the study of the formation process of VS to include nanodots with asymmetric geometry. Based on our experimental observation, we will discuss the possibility to control the generation process of VS by directed modifications of nanodot geometry.

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#### **HG03**

#### Magnetization switching utilizing the magnetic exchange interaction Rene Schmidt\*, Alexander Schwarz and Roland Wiesendanger

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Magnetization reversal is usually induced by applying sufficiently large external magnetic fields. More recently, spin transfer torque using a spin polarized current has been successfully introduced as an alternative approach [1]. Here, we demonstrate the feasibility to utilize the magnetic exchange interaction with a Magnetic Exchange Force Microscopy (MExFM) set-up to reverse the magnetization direction of a magnetic specimen. MExFM opened up the possibility to map spin structures of insulating [2] as well as conducting surfaces [3] in real space. We use the distance dependence of the magnetic exchange interaction, measured with Magnetic Exchange Force Spectroscopy (MExFS) [4], to induce magnetization reversals. In our case the magnetic tip apex switches its magnetization direction during scanning an antiferromagnetic surface with atomic resolution. The apex actually behaves like an independent superparamagnetic cluster with uniaxial anisotropy [5]. We are able to study switching lifetimes, exchange interaction dependent energy barriers and the influence of magnetic fields. Our findings suggest that the magnetic exchange interaction can be employed to monitor and control magnetization reversal processes with atomic resolution on conducting as well as on insulating surfaces.

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#### HG04

#### X-ray spectroscopy in pulsed high magnetic fields

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The combination of synchrotron x-ray spectroscopy and pulsed high magnetic fields provides new opportunities for the study of magnetic materials. We have developed X-ray magnetic circular dichroism (XMCD) [1] for element [2] and orbital [3] selective studies of magnetism, and Nuclear Forward Scattering (NFS) of synchrotron radiation [4] to perform isotope selective hyperfine spectroscopy in high magnetic fields. A miniature coil developed by the ESRF Sample Environment Support Laboratory reaches fields of 30 T with a total duration of 1 ms at repetition rates of 5 pulses per minute, and the integrated sample cryostat provides optical access to the sample along the field direction and covers temperatures from 5 K to 250 K [5]. In this contribution we present the coil [5] and detection schemes [12,4] along with application examples such as element selective magnetometry in Erbium Iron Garnet [2], the mapping of a magnetostructural phase transition in double Perovskites [3], the high field magnetization of alpha iron [4] and a study of the spin structure in Copper Iron Oxide [6].

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### 88 | The 19<sup>th</sup> International Conference on Magnetism

#### **HH01**

#### Recent developments in magnetic measurements: From technical method to physical knowledge

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Novel magnetic materials and applications call for enhanced measurement methods, where the goal of increased physical insight is pursued against the requirements of reproducibility and traceability. We discuss a few examples of measuring techniques where advances have been recently achieved. We shall deal in particular with the following topics: 1) High-speed magneto-optical microscopy where nanosecond time-resolved Kerr microscopy reveals the high-frequency domain wall dynamics in a variety of samples: thin films, amorphous ribbons, patterned magnets. It concurrently allows for pulsed inductive magnetometric techniques, by which the spin precessional modes and their damping behavior can be observed 2) Calorimetry under magnetic field a technique permitting one to to directly measure the entropy s(Ha T) as a function of temperature T and applied magnetic field Ha. It is especially attractive for the characterization of magnetocaloric materials and the determination of first order phase transitions with hysteresis. The discussion will focus on the magnetic field induced transition in La-Fe-Co-Si alloys Heusler alloys and materials with spin reorientation transition like Er<sub>2</sub>Fe<sub>14</sub>B and W-type Co-Zn substituted Ba ferrites. 2) Pulsed field magnetometry in permanent magnets, an industrially attractive method of characterization of extra-hard rare-earth based magnets, overcoming the intrinsic fieldstrength limitations of the standard electromagnet-based characterization

#### **HH02**

So, you need reliable magnetic measurements you can use with confidence? How the magnetic measurement capabilities at NPL can help

#### Michael Hall

National Physical Laboratory, United Kingdom

Magnetic quantities have an impact on many aspects of our everyday life; in health and safety, the magnetic field levels produced by the domestic appliances we use so regularly is carefully monitored by the manufacturers to meet EMC regulations (we provide measurements that are compliant up to 120 kHz with EN 50366). In the aerospace industry, soft magnetic materials are used in extreme conditions (temperatures up to 450 °C and stress levels up to 450 MPa) and the material can be characterised at NPL using unique measurement techniques. In the car industry, the material properties of permanent magnets have become very important; some cars now contain in excess of 70 magnets. The demand for traceable electrical and magnetic measurements is broadening as new materials and techniques have brought the need for these measurements to new areas of industry. In this talk the magnetic measurement facilities, capabilities and field standards available at NPL will be presented. How these are being developed to meet the demanding conditions of use of instruments and materials will be presented and illustrated with case studies.

#### HH03

#### Application of pulsed eddy current technique to inspect the pipeline of nuclear plants

D. G Park

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The local wall thinning is a point of concern in almost all steel structures such as the pipe lines. The present research describes the application of Pulsed Eddy Current (PEC) for wall thickness monitoring. We have built a PEC system which has the ability to detect the wall-thinning or corrosion, and sub-surface defects. In our PEC probe an exciting coil in conjunction with a Hall-sensor has been used. The PEC response to varying metal thickness was measured at various thicknesses of insulations on the tested sample. Excitation coil in the probe is driven by a rectangular current pulse; the time domain features of the detected pulse, such as 'peak value' and 'time to zero' were used to describe the wall thinning in the tested sample. A real time LabVIEW program was developed for the data acquisition and scanning the probe on the insulated sample. The scanning results were continuously displayed on the computer monitor. Signal processing techniques such as energy and power spectrum density (PSD) were devised to infer the PEC response signal. The results show that the proposed PEC differential probes

#### HI01

#### Probing the exotic surface states in topological insulators and superconductors Yoichi Ando\*

ISIR, Osaka University, Japan

A topological state of matter is characterized by a nontrivial topological structure of the quantum-mechanical wavefunctions in the Hilbert space. Due to the bulk-edge correspondence, a gapless surface state always accompanies a topologically nontrivial bulk state. In topological insulators (TIs), a nontrivial Z2 topology of the bulk state leads to the emergence of Dirac fermions on the surface Similarly topological superconductors (TSCs) are accompanied by surface Andreev bound states that often consist of Majorana fermions. In this talk, I will present our experiments to address those exotic surface states. For TIs, we discovered that the chalcogen-ordered tetradymite Bi<sub>2</sub>Te<sub>2</sub>Se presents a high bulk resistivity allowing one to observe clear surface quantum oscillations [1]: more recently we demonstrated that in a related solid-solution system Bi<sub>2</sub> .Sb, Te<sub>2</sub> .Se, it is possible to achieve a surface-dominated transport even in a bulk crystal [2]. For TSCs, in a superconducting doped-TI material Cu Bi<sub>2</sub>Se<sub>2</sub>, we have succeeded in observing an unconventional surface Andreev bound state, which gives evidence for a new type of topological superconductivity associated with helical Majorana fermions [3]. These works were done in collaboration with A. A. Taskin, Z. Ren, S. Sasaki, and K. Segawa,

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#### HI02

#### Electronic structure study of bulk HgTe via angle resolved photoemission spectroscopy

Chang Liu<sup>1</sup>, Suyang Xu<sup>1</sup>, Madhab Neupane<sup>1</sup>, Helin Cao<sup>2</sup>, Ireneusz Miotkowski<sup>2</sup>, L. Andrew Wray<sup>3</sup>, Hsin Lin<sup>4</sup>, R. S. Markiewicz<sup>4</sup>, A. Bansil<sup>4</sup>, Yong P. Chen<sup>2</sup> and M. Zahid Hasan<sup>1</sup> <sup>1</sup> Joseph Henry Laboratory and Department of Physics, Princeton University, USA <sup>2</sup> Department of Physics, Purdue University, USA Advanced Light Source, Lawrence Berkeley National Laboratory

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#### HI03

#### Hidden topological order in URu<sub>2</sub>Si<sub>2</sub>

Tanmov Das Theoretical Division Los Alamos National Laboratory USA

Recently breakthrough studies have discovered series of new state of matters, including topological insulators, topological Mott insulators, topological Kondo insulators, topological antiferromagnetic insulators and others. Here we present an analogous, vet fundamentally novel topological order which arises due to spin-orbit coupling density wave in the heavy-fermion material URu2Si2.[1] We show via ab-initio calculations that an incommensurate Fermi surface 'nesting' in the partially-filled f-states causes a staggered spin-orbit coupling in the hidden-order state. In this cause neither the spin (S) nor the orbital (L) alone causes ordered state, rather a modulated spin-momentum locked density wave propagates along the unidirectional nesting direction with a polarized total angular momentum  $mJ = \pm 2$ , in excellent agreement with neutron scatterings. This is a radically new order parameter and a novel phase of matter. It breaks spontaneous rotational symmetry, but not the time-reversal symmetry and is thus immune to any time-reversal invariant perturbation such as pressure, whereas magnetic field will destroy it. We compute the topological quantum number to show that the hidden-order gap opening causes a trivial to non-trivial topological phase transition, and hence defines a novel 'topological quantum critical point' as a function of magnetic field

#### HI04

#### Robustness of 1D topological superconductors with Majorana edge states against lattice modulation

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Recently there has been a substantial amount of effort to observe Majorana edge states at the edges or vortex cores of topological superconductors. The robustness of such states against spatial inhomogeneity and electron correlation will be particularly important in observing and manipulating them. We have studied a one-dimensional quantum wire to analyze the condition for a topological phase with edge Majorana fermions to appear when an incommensurate (quasiperiodic) lattice modulation is applied. We have adopted the tight-binding model Hamiltonian by Stoudenmire and coworkers [1], which incorporates the Zeeman effect, the spin-momentum coupling, the proximity effect induced by a bulk superconductor, and the onsite electron-electron interaction U, and analyzed its low-energy behavior when the lattice modulation is applied by the density-matrix renormalization group method. For U=0 we have found multiple reentrant transitions between topological phases with edge Majorana fermions and non-topological phases [2]. The number of transitions increase as the modulation becomes stronger. For U>0, as in the case without the lattice modulation [1], the region with Majorana fermions is broadened in the phase diagram [2]. The effects of other types of spatial inhomogeneities, such as a harmonic trapping potential, will also be presented

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#### H105

#### Floquet theory of photo-induced topological phase transitions: Application to graphene

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The effect of strong lasers on the topology of electron systems is becoming a hot topic [1-3]. Recently, a theoretical proposal was made in two dimensional Dirac systems where an application of circularly polarized light was shown to turn the system into a quantum Hall state with a non-trivial photo-induced Chern number and an emergence of photo-induced edge channels [1,2]. One can see this as a dynamical realization of the Haldane model of a quantum Hall state without Landau levels [4], i.e., the circularly polarized light plays the role of the local magnetic fields. This proposal applies to a broad class of multi-band systems including graphene, graphite and surface states of topological insulators as well as cold atoms in optical lattices with synthetic gauge fields. Theoretically, this prediction is based on the Floquet extension of linear response theory and the TKNN formula[1] as well as the Landauer theory of transport[2]. We also explain possible experimental setups to detect this state using transport, photoemission, and orbital Kerr effect. \* This work was done in collaboration with H. Aoki, T. Kitagawa, T. A. Brataas, L. Fu and E. Demler.

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#### HJ01

#### Strong coupling of spin, orbital and lattice degrees of freedom in Ru oxides

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Being a 4d transition element, Ru has more extended wavefunction compared to those of 3d transition metal elements, and so is likely to have relatively smaller correlation effects. However, it turns out that this relatively broader wavefunction produces strikingly different physical properties from those of 3d transition metal oxides. In this talk, we will highlight the unique physical properties of two Ru oxides emphasizing its difference from 3d transition metal oxides. First example is Tl<sub>2</sub>Ru<sub>2</sub>O<sub>7</sub>, which exhibits a metal-insulator transition at 120 K. We will examine the nature of the metal-insulator transition both experimentally and theoretically, and are going to discuss how a spindriven spontaneous orbital ordering produces the experimentally observed structural transition and leads to the metal-insulator transition at the exactly same temperature. The second example is the ferromagnetic transition of SrRuO3 with Tc=165 K. Although it has been known over a half century, the true nature of the ferromagnetic transition is poorly understood. Here combining both neutron diffraction and inelastic neutron scattering we will show how the ferromagnetic ground state is intimately coupled to the unusual volume anomaly, i.e. Invar behavior at an atomic scale.

#### HJ02

Competing magnetic interactions in eta-carbide-type transition-metal compounds: New class of itinerant-electron frustrated magnets Hiroyuki Nakamura\*, Takeshi Waki and Yoshikazu Tabata

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Although a number of ternary transition metal compounds with the cubic eta-carbidetype structure,  $M_3T_3X$  and  $M_6T_6X$  with  $M = M_0$ , W, etc., T = Fe, Co, Ni, etc., and X = C or N have been known to exist, their magnetic properties have been less studied. The T sublattice forms the 'stella quadrangula' lattice, which is a newly found geometrically frustrated corner-shared lattice of two nested regular tetrahedrons. Although the stella quadrangula lattice is one of pyrochlore-derived lattices, we anticipate the unique nature of the frustration due to the competition between first- and second-neighbor interactions, J1 and J2, possibly of direct exchange type. Because of the delicate balance between J1 and J2, the eta-carbide-type compounds shows various magnetic ground states such as nonmagnetic non-Fermi liquid, ferromagnetism, antiferromagnetism, etc. [1]. In addition, various types of field-induced transitions are often observed [2]. Here we summarized comprehensively the magnetism of the eta-carbide-type transition metal carbides and nitrides together with resent experimental data.

[1] T. Waki et al., J. Phys. Soc. Jpn. 79 (2010) 043701, EPL 94 (2011) 37004, J. Alloys Compd. 509 (2011) 9451. [2] T. Waki et al., J. Phys. Soc. Jpn. 79 (2010) 093703, J. Phys.: Conf. Ser. 320 (2011) 012069.

#### HJ03

## Long-time variation of magnetic structure in (Ce-La)Ir<sub>3</sub>Si<sub>2</sub> : Effect of randomness

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Long-time variations of magnetic properties observed in spin-glasses have been regarded as due to the multi-valley structure of the free energy surface arising from random magnetic interactions. Therefore, in a system without randomness or imperfections we have not expected to observe a time variation of magnetic property within an attainable time scale. Contrary to all expectations, we found a long-time variation of magnetic structure in a non-diluted uniform magnet Celr<sub>3</sub>Si<sub>2</sub> by means of time-resolved neutron scattering measurements.[1,2] We examined other materials and have found similar long-time variations of magnetic structure in PrCo<sub>3</sub>Si<sub>2</sub>. TbNi<sub>2</sub>Si<sub>5</sub>. Ca<sub>2</sub>Co<sub>2</sub>O<sub>6</sub> and some other compounds [3,4] All these compounds exhibit successive magnetic transitions and multi-step metamagnetic transitions, which suggest that the frustrations coming from competing magnetic interactions causes the long-time variation of magnetic structure. In order to show that randomness caused by inevitable impurities or imperfections is not the main cause of the long-time variation, we have made macroscopic and neutron scattering measurements of a random system (Ce-La) Ir<sub>3</sub>Si<sub>2</sub>. The time variation behavior of this random system is basically identical to that of Celr<sub>3</sub>Si<sub>2</sub>. The see results have shown that the long-time variation of magnetic structure in Celr<sub>5</sub>Si<sub>4</sub>, conto cause show that the long-time variation fragmetic structure in the cause of the long-time variation.

I. K. Motoya et al. : J. Phys. Conference Series 200 (2010) 032048. 2. T. Mayoshi et al.: J. Phys. Soc. Jpn. 81 (2012) 014704. 3. K. Motoya et al. : J. Phys. Conference Series 273 (2011) 012124. 4. T. Moyoshi and K. Motoya: J. Phys. Soc. Jpn. 80 (20112) 034701.

#### HJ04

## Influence of symmetry on Sm magnetism studied in $SmIr_2Si_2$ polymorphs

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SmIr<sub>2</sub>Si<sub>2</sub> crystallizes in two structural polymorphs which make possible unique study of influence of symmetry on Sm magnetism. One polymorph is the primitive tetragonal structure of the CaBe<sub>2</sub>Ge<sub>2</sub> structure type stable at high temperatures (high temperature polymorph-HTP) and the other is a body centered tetragonal structure of the TCr<sub>2</sub>Si<sub>2</sub> structure type stable at low temperature (low temperature phase-LTP). The HTP exists at room temperature as a metastable phase in the quenched samples while the LTP can be obtained by proper annealing. We have found three magnetic phase transitions in both structure polymorphs in specific heat data. The transitions at temperature 1.7, 6.2 and 18.8 K have been found for the HTP and the LTP exhibits transitions at 1.9, 6.1 and 38.9 K. Magnetic moment of 0.034  $\mu$ B/f.u. has been found for both polymorphs (in the field of 7 T at 2 K). Magnetic features of both polymorphs are discussed on the basis of the strong influence of the trystal field and the admixture of the low lying excited multiplet J = 7/2 with the ground multiplet J = 5/2. All experimental results are supported by first principles calculations.

#### HJ05

## Unconventional thermal expansion of BaIrO<sub>3</sub> investigated by temperature dependent x-ray and neutron diffractions

Jinwon Jeong<sup>1</sup>, Bin Chang<sup>1</sup>, Dahee Jung<sup>1</sup>, Han-jin Noh<sup>1</sup>\* and Seongsu Lee<sup>2</sup> <sup>1</sup> Department of Physics, Chonnam National University, Korea <sup>2</sup> Korea Atomic Energy Research Institute, Korea

We have studied the temperature dependence of the unit cell volume of 9R-type BaIrO<sub>3</sub> by X-ray and neutron diffractions. The volume and the unit cell parameters show a strong correlation with the weak ferromagnetic charge-density-wave phase of this 5d transition metal compound. The volume vs. temperature curve is not well described by a Debye function, which indicates that the interplay between charge, spin, and lattice plays an important role in the thermal expansion behavior. A possible scenario will be presented in order to understand this unconventional behavior.

#### IA01

#### A quantum phase transition hidden beneath the superconducting dome of iron-pnictides Yuji Matsuda

Physics, Kyoto University, Japan

An enduring question in strongly correlated electron systems is whether high-Tc superconductivity is driven by an underlying quantum critical point (QCP) separating different ground states. In particular, whether a QCP lies beneath the superconducting dome or the criticality is avoided by the transition to the superconducting state has been a central issue. Among iron-based superconductors, the isovalent pnictogen substituted system BaFe<sub>2</sub>(As<sub>1</sub>, P<sub>2</sub>)<sub>2</sub> appears to be the most suitable system to discuss many physical properties, because BaFe<sub>2</sub>(As<sub>1</sub>, P<sub>2</sub>)<sub>2</sub> can be grown with very clean and homogeneous, as evidenced by the quantum oscillations observed over a wide doping range even in the superconducting dome giving detailed knowledge on the electronic structure. We report a sharp peak in the x-dependence of zero-temperature London penetration depth L(0) in BaFe<sub>2</sub>(As<sub>1</sub>, P<sub>2</sub>)<sub>2</sub> at the optimum composition x=0.30 (Tc=30 K). This peak structure most likely results from pronounced quantum fluctuations associated with the QCP which separates two distint superconducting phases, giving the first convincing signature of a second-order quantum phase transition deep inside the superfluid-4He value, implying a possible crossover towards the Bose-Einstein condensate limit driven by quantum criticality.

This work has been done in collaboration with T. Shibauchi, K. Hashimoto, S. Kasahara, Y.Mizukami, T. Terashima, H. Ikeda (Kyoto), A. Carrington (Bristol), K. Cho, R. Prozorov, M. Tanatar (Ames), N. Salovich, R.,W. Giannetta(Illinois at Urbana-Champaign), and H. Kitano (Aoyama Gakuin)

#### IA02

## Re-entrant quantum criticality in pressurized $Yb_2Pd_2Sn$ and $Yb_2Pd_2In_{1,x}Sn_x$

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<sup>5</sup> Max Planck Institute for Chemical Physics of Solids, Dresden, Germany <sup>6</sup> Dipartimento di Chimica e Chimica Industriale, University of Genova, Italy

Recently we reported on the discovery of two consecutive, pressure driven magnetic instabilities in Yb\_Pd\_Sn [1]. They emerge in a non-Fermi liquid environment at the initial and the final point of a dome-like, single magnetic phase at pressures pc1 ~ 1 GPa and pc2 ~ 4 GPa. This unique behavior of Yb\_Pd\_Sn is supposed to result from mutually competing, pressure modified energy scales, which in case of Yb\_Pd\_Sn cause a sign change in the pressure dependence of the Kondo temperature TK and magnetic ordering temperature TN. An earlier study [2] demonstrated the occurrence of long range magnetic order in a dome-like phase space for a narrow concentration range of Sn-rich samples in Yb\_Pd\_In<sub>1,x</sub>Sn,. Ordered moments of the order of 1 µB are attained and the transition temperature reveals values above 3 K, although both Yb\_Pd\_In and Yb\_Pd\_Sn exhibit paramagnetic ground states. The aim of the present work is a detailed study of the pressure and field response of Yb\_Pd\_Sn and of alloys with concentrations in the "magnetic dome" of Yb\_Pd\_In<sub>1,x</sub>Sn,. We show how the transition temperatures and the ground states modify when applying both pressure and magnetic fields.

[1] T. Muramatsu et al., Phys. Rev. B 83, (2011) 180404(R). [2] E. Bauer et al., J. Phys.: Cond. Mat. 17 (2005) S999.

#### IA03

Coupled fermi-bose renormalization group flow for a two-flavor spinfermion model close to its antiferromagnetic quantum critical point Junhyun Lee, Philipp Strack and Subir Sachdev\* Department of Physics, Harvard University, USA

The strong correlations emerging in proximity of antiferromagnetic quantum critical points (QCP) in low-dimensional electronic materials are believed to be one of the main causes for their striking non-Fermi liquid behavior. In this quantum-critical regime, metallic electrons are strongly coupled to gapless, collective spin fluctuations rendering theoretical approaches a challenging affair. In this talk, we will present first results of a renormalization group (RG) analysis of a new, prototypical two-flavor spin-fermion model in the vicinity of its quantum critical point. This model is free of the sign-problem and therefore ideally suited to compare results of analytical methods with quantum Monte Carlo simulations [1]. Our RG method is based on an exact flow equation for the effective Fermi-Bose action. It allows us to treat the electrons and spin fluctuations on equal footing, and to extract correlation functions consistently ordered along a joint energy scale. Solving our flow equations, we present results for renormalized observables such as the fermion self energy and the magnetic susceptibility when approaching the QCP from the paramagnetic phase.

[1] arXiv.org:1206.0742

#### IA04

#### Magnetism and filling-controlled mott transition in strongly spin-orbitcoupled iridium oxide

Kenya Ohgushi<sup>1</sup>\*, Jun-ichi Yamaura<sup>1</sup>, Hiroyuki Ohsumi<sup>2</sup>, Soshi Takeshita<sup>2</sup>, Hidenori Takagi<sup>3</sup>, Taka-hisa Arima<sup>3</sup> and Yutaka Ueda<sup>1</sup> <sup>1</sup> Institute for Solid State Physics, University of Tokyo, Japan <sup>2</sup> RIKEN SPring-8 Center, Japan <sup>3</sup> University of Tokyo, Japan

By performing resonant x-ray diffraction experiments for the post-perovskite-type compound CalrO<sub>3</sub>, we revealed that the magnetic structure is a striped-typed order, which is stabilized by the unique superexchange interaction across the edge-sharing bonds in the strongly spin-orbit-coupled Jeff = 1/2 state. Moreover, by introducing hole carriers into CalrO<sub>3</sub> through the chemical substitution, we have successfully developed a novel metallic state, where bandwidth, Coulomb repulsion, spin-orbit coupling, and Hund coupling are competing.

[1] K. Ohgushi, et al., arXiv:1108.4523v1. [2] K. Ohgushi, et al., Phys. Rev. B 74, 241104(R) (2006).

#### IA05

#### Magnetic field tuned QCP in YbPtBi

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Magnetic field dependence of thermodynamic and transport properties of single crystals of stoichiometric, heavy fermion ( $\gamma \sim 8$  J/mol K2), antiferromagnetic (TN  $\sim 0.4$  K) compound YbPtBi was studied down to 20 mK. The magnetic order is suppressed on application of  $\sim 4$  KOe magnetic field. The Fermi-liquid state appears to be separated from the quantum critical point. The H\* crossover line ("Hall line") marks the emergence of the Fermi-liquid state. The H-T phase diagram will be discussed in context of results on other Yb-based heavy fermions.

#### IB01

Field-annealed Fe-Ni-Nb-B amorphous and nanocrystalline alloys for magnetic sensor applications

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The amorphous and nanocrystalline alloys have attracted a great deal of technological interest due to their applicability in various types of magnetic sensors. These alloys often require specific annealing conditions to achieve the desired response to applied magnetic field. In this work, we report on the effects of both longitudinal and transverse magnetic field applied during the heat treatment on the soft magnetic behaviour and the giant magnetoimpedance (GMI) effect in the series of (Fe, Ni.), Nb<sub>2</sub>B<sub>12</sub> (x=0.25, 0.33 and 0.5) amorphous and nanocrystalline alloys. The melt-spun ribbons were isothermally annealed under high vacuum in the presence of transverse (TF) or longitudinal (LF) magnetic field with the magnitude up to 640 kA/m. The heat treatment under LF-conditions results in an appreciable reduction of the coercivity. Sheared loops with good field linearity were achieved after TF-annealing. The GMI measurements were performed in dc magnetic field up to  $\pm$  9600 A/m over a frequency range 0.1 - 10 MHz. The nanocrystalline samples show markedly larger values of GMI ratio  $(\Delta Z/Z)$  as compared to their amorphous counterparts. The field annealing results in the modified the GMI response of both amorphous and nanocrystalline ribbons, which can be used for tuning their characteristics for magnetic sensor applications.

#### IB02

## Amorphous and nanocrystalline magnetic materials: Research and production in china

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The planar-flow casting (PFC) is a continuous casting method for the production of a large quantity of rapidly solidified metallic ribbons. In particular, the PFC attracted great interests in the production of amorphous and nanocrystalline thin and wide ribbons because of the cost savings process in a single stage and the novel properties arising from the microstructure of rapidly solidified ribbons. Amorphous and nanocrystalline ribbons are widely used in electrical and electronic system such as distribution transformers and inductive devices because of low core loss. This paper seeks to summarize the recent progress in synthesis methods, properties, and applications in the field of amorphous and nanocrystalline soft magnetic materials in forms of ribbon, the emphasis is focused on some key technologies for producing wide amorphous ribbon and thin nanocrystalline ribbon. In the summary, the total market demands of amorphous and nanocrystalline sin China are predicted.

#### IB03

## Recent status of soft magnetic material applications for renewable energy and eco-friendly vehicle

In-bum Jeong

Changsung Corporation, Korea

Soft magnetic materials have been widely used in electrical and electronic devices such as inductor, transformer, sensor, motor, EMI absorber and others. Ideal soft magnetic material which is relevant to modern high power electronics should have the best combination of high permeability under DC excitation, low loss, good thermal stability, good frequency behavior, and reasonable cost. But, unfortunately, all the commercialized soft magnetic materials such as ferrite, laminated silicon steel sheet, wound amorphous and nanocrystalline ribbon, and compressed metallic alloy powder core do not meet all the requirements perfectly so far. In Renewable energy and electric vehicle applications, especially the soft magnetic materials are getting important for high efficiency and reliability. Therefore, there are strong market demands for strengthening each material's forte and overcoming each material's limitations and shortcomings. This presentation is focused on Recent Status of Soft Magnetic Material Applications for Renewable Energy and Eco-friendly Vehicle.

#### IC01

#### Random fields in molecular magnets

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Measurements of the magnetic susceptibility of single crystals of the usual form of  $Mn_{12}$ -acetate and of a new high-symmetry variant,  $Mn_{12}$ ac-MeOH, indicate a transition to dipolar ferromagnetism below ~ 0.9 K in both materials [1]. In the case of  $Mn_{12}$ ac-MeOH, the suppression of ferromagnetism by a magnetic field applied transverse to the easy axis is consistent with mean-field theory for an ordered single crystal:  $Mn_{12}$ ac-MeOH is a realization of a "Transverse-Field Ising Ferromagnet" (TFIFM). By contrast, the suppression of ferromagnetism by transverse field is considerably more rapid in  $Mn_{12}$ -acetate, resembling instead the behavior observed in randomly site diluted LiHo<sub>1-x</sub>Y<sub>x</sub>F<sub>4</sub>, the only heretofore known realization of a "Random-Field Ising Ferromagnet" (RFIFM) [2]. The randomness in  $Mn_{12}$ -acetate derives from the intrinsic distribution of locally tilted magnetic easy axes known to exist in these crystals, suggesting that this prototypical molecular magnet is a realization of the random field Ising ferromagnet (RFIFM) in which the random field can be tuned by a field applied transverse to the easy axis.

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#### IC02

#### Stability of incommensurate field-induced magnetic order via site-disorder

Francesco Casola<sup>18</sup>, Toni Shiroka<sup>1</sup>, Shuang Wang<sup>2</sup>, Christian Rüegg<sup>4</sup>, Henrik Moodysson Rønnow<sup>2</sup>, Michael Grbié<sup>4</sup>, Mladen Horvatić<sup>4</sup>, Stelfen Krämer<sup>4</sup>, Suirtha Mukhopadhyay<sup>4</sup>, Claude Berthier<sup>4</sup>, Hans Rudolf Ott<sup>4</sup> and Joël Mesot<sup>4</sup> <sup>1</sup>Laboratorium für Festkörperphysik, ETIH Hönggerberg, CH-8093 Zürich, Switzerland <sup>2</sup>Laboratory<sup>6</sup> Oyuantum Magnetism, Ecole Polytechnique Fédérale de Lassame, CH-1015 Lausanne, Switzerland

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Low-dimensional spin-gapped antiferromagnets (AF) can be driven out of their quantum-disordered (QD) regime not only through the application of pressure or magnetic fields, but also via the introduction of disorder. Despite previous observations of these phenomena in a variety of materials such as Haldane chains, spin ladders or spin-Peierls systems, a limited number of studies has considered the impact of defects on the field-induced magnetic order. It was found that disorder can affect the quantum-critical behaviour of these materials, induce new phases dominated by triplet confirmement or even destroy completely any phase transition in case of incommensurate correlations in the QD regime. Extending low-field 31P NMR experiments on the frustrated two-leg spin-1/2 ladder  $Bi(Cu_{1-x}Tn_x)_2PO_6$  [1] to magnetic fields up to 31T we observed for x=0 the onset of incommensurate field-induced magnetic order at Hc = 20.96 T. The critical exponent for the transition temperature  $Tc(H) \propto (Hc(T)-Hc(0))v$  is v = 0.42. Remarkably, the same critical behaviour is found for the compound with x=0.01, but with an enhanced Hc(0) of about 24 T. Site disorder introducing non magnetic Zn ions on regular Cu sites opens the possibility to study finite-size effects in quasi-1D spin systems by directly tuning the chain length via defect concentration.

[1] - F. Casola, T. Shiroka, S. Wang, K. Conder, E. Pomjakushina, J. Mesot and H.-R. Ott, Phys. Rev. Lett. 105, 067203 (2010).

#### IC03

## Low-temperature heat transport and field-induced quantum phase transitions of spin gapped quantum magnets

X. F. Sun<sup>1</sup>\*, C. Fan<sup>1</sup>, L. M. Chen<sup>2</sup>, W. P. Ke Ke<sup>1</sup>, Z. Y. Zhao<sup>1</sup> and X. M. Wang<sup>1</sup> <sup>1</sup> Hefei National Laboratory for Physical Sciences at the Microscale, University of Science and Technology of China, China

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Low-dimensional or frustrated quantum magnets were revealed to exhibit exotic ground states, magnetic excitations, and quantum phase transitions (QPTs). In the spin-gapped antiferromagnets, the external magnetic field can close the gap in the spectrum, which results in a QPT between a low-field disordered paramagnetic phase and a high-field long-range ordered one. An intriguing finding is that this ordered phase can be approximately described as a Bose-Einstein condensation (BEC) of magnons. In this work, we study the low-temperature and high-field thermal conductivity (kappa) of several spin gapped quantum magnets, including the quasi-one-dimensional S=1 chain compound NiCl<sub>2</sub>-4SC(NH<sub>2</sub>)<sub>2</sub> (DTN), the quasi-one-dimensional ladder compound (CH<sub>3</sub>)<sub>2</sub>CHNH<sub>3</sub>CuCl<sub>3</sub>(IPACuCl<sub>3</sub>), and the layered spin-dimer compound Ba<sub>3</sub>Mn<sub>2</sub>O<sub>8</sub> [1-3]. It is found that the magnetic excitations can contribute to the heat transport rather strongly in these materials, particularly at the field-induced QPTs, by either transporting heat or scattering phonons.

[1] X. F. Sun, W. Tao, X. M. Wang, and C. Fan, Phys. Rev. Lett. 102, 167202 (2009). [2] W. P. Ke, X. M. Wang, C. Fan, Z. Y. Zhao, X. G. Liu, L. M. Chen, Q. J. Li, X. Zhao, and X. F. Sun, Phys. Rev. B 84, 094440 (2011). [3] L. M. Chen, X. F. Sun et al., unpublished.

#### IC04

#### Review talk about spin superfluidity Yury Bunkov Institut NEEL, France

There is a complete analogy between the Bose-Einstein condensation of atomic gases and the Bose-Einstein condensation of magnons in antiferromagnets. Five different states of magnon condensation with Spin Supercurrent have been found in different states of antiferromagnetic superfluid 3He [1]. All these cases are examples of the Bose-Einstein condensation of magnons with the interaction potential provided by specific spin-orbit coupling. The BEC phenomenon in the gas of magnons is readily accessible owing to the possibility of modifying the spin-orbit coupling. In some cases the BEC of magnons corresponds to almost 100% condensation. Recently the magnon BEC state and Spin Supercurrent was found on a 55Mn nuclear in antiferromagnetic monocristals  $CsMnF_3$  and  $MnCO_3$  [2].

[1] Yuriy M. Bunkov and Grigoriy Volovik, J. Phys.: Condens. Matter 22 164210 (2010) [2] Yuriy M.Bunkov et all., Phys. Rev. Lett. 108, 177002 (2012).

#### IC05

## The spin-1/2 Heisenberg antiferromagnetic chain experimental confirmation of 2 and 4 spinon terms

Bella Lake<sup>1\*</sup>, D. Alan Tennant<sup>1</sup>, Jean-sebastien Caux<sup>2</sup> and Christopher D Frost<sup>3</sup> <sup>1</sup> Helmholtz Zentrum Berlin fur Materialien und Energie, Germany <sup>2</sup> Instituut voor Theoretische Fysica, Universiteit van Amsterdam, Netherlands <sup>3</sup> ISIS Facility, Rutherford Appleton Laboratory, United Kingdom

The Spin-1/2 Heisenberg antiferromagnetic chain is the fundamental example of a correlated quantum system and has been a long-standing problem in condensed matter physics. The problem was first tackled by Hans Bethe in 1931 who found that the ground state is quantum disordered down to lowest temperatures [1]. In 1981 Faddeev and Takhtajan realized that the basic excitations are spinons [2] which have fractional spin quantum number (S=1/2) and are necessarily created in pairs. In neutron scattering experiments the excitations are observed as a multi-spinon continuum consisting of 2,4,6,8,... spinon states. The dynamical structure factor which is measured by the neutrons has however only recently been calculated [3,4,5]. Caux, Hagemans and Maillet have determined it to high precision and were able to show that while the 2-spinon contribution accounted for 73% of the spectral weight, the 4-spinon contribution was also substantial at 25%. Here we use inelastic neutron scattering to obtain the dynamical structure factor of the spin-1/2 Heisenberg antiferromagnetic chain compound KCuF<sub>3</sub> to high accuracy and in absolute units. The data show much better agreement with this recent theoretical work than earlier theories (Muller Ansatz and field theory) and is able to confirm the presence of the 4-spinon term.

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#### ID01

#### X-ray microscopy of nanoscale spin dynamics Peter Fischer

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One of the scientific and technological challenges in nanomagnetism research is to image magnetism down to fundamental length and time scales with elemental sensitivity in advanced multicomponent materials. Magnetic soft X-ray microscopy is a unique analytical technique combining X-ray magnetic circular dichroism (X-MCD) as element specific magnetic contrast mechanism with high spatial resolution down to almost 10nm and a temporal resolution below 100ps [1]. This allows to investigate in detail precessional and relaxation phenomena, which occur e.g. in domain wall motion [2-3] and vortex dynamics [4-8]. In this talk I will review the achievements and future perspectives of using soft X-rays for magnetic imaging by selected examples from current research. The collaboration with many colleagues in particular M.-Y.Im, S.-K. Kim, D.-H. Kim, G. Meier, S. Kasai, T. Ono, D. Allwood, K. Buchanan, B. Mesler and S.-C. Shin is highly appreciated. This work is supported by the Director, Office of Sasic Energy Sciences, Materials Sciences and Engineering Division, of the U.S. Department of Energy under Contract No. DE-AC02-05-CH11231.

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#### ID02

#### Non linear spin transfer induced vortex dynamics

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A magnetic vortex can be excited into precessional motion around its equilibrium position through spin transfer torque[1,2]. The associated microwave emission is weak but very coherent. Here we present experimental evidences of high power, low linewidth spin-transfer induced vortex oscillations in MgO based magnetic tunnel junctions with a large TMR ratio[3]. Indeed, a clear picture of how the spin transfer torque influences the frequency and the amplitude of the gyrotropic vortex motion in a nano-pillar is still lacking. Here we derive an analytical model that takes into account the nonlinearity of the vortex core confining force and the non-linearity of the effective damping that will be compared to experimental results and micromagnetic simulations. Therefore it allows us to determine the evolution of the vortex core dynamics with perpendicular field and current. Finally, we show that a very efficient coupling[4] to an external microwave current can be obtained in such spin transfer vortex oscillators (STVOs) and compare these results to theoretical locking range, taking into account all possible components of spin transfer torques. Support from ANR through VOICE PNANO-09-P231-36 and EU MASTER No. NMP-FP7-212257 and Canon-Anelva for MTJ films is acknowledged. E. G. acknowledges support from CNES and DGA.

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#### ID03

## Study of spin transfer induced coupled vortices dynamics in a single spin-valve

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In order to address the problem of low coherence of actual spin-transfer nano-oscillators, recent studies have focused on phase locking phenomena, between an oscillator and an external source, or between assemblies of two or more oscillators. In the latter case, new collective modes have to be differentiated from the modes of isolated single oscillators, namely the chirality and polarity in case of a vortex[2,3]. In this work, we suggest to use two vortices, nucleated in both layers of a Py(15nm)Cu(10nm)/Py(4nm) spin-valve, as two interacting oscillators. Fundamental interests of this system lie in the identification of the excited coupled modes and their selection rules. Besides the observation of highly coherent vortex oscillations, we demonstrate that the dynamic behaviour is highly dependent on the vortices parameters[3]. We will compare the experimental results to analytical predictions and simulations, studying the importance of fufferent sources of coupling between the vortex. A careful study of spin-polarized transport also shows the strong influence of the fast magnetization gradients associated with vortices. This work will be extended further to conditions to synchronize arrays of interacting vortex oscillators.

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#### ID04

## Collective excitation of magnetostatically coupled two-adjacent magnetic vortices and their relative phase difference

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A magnetic vortex exhibits a core gyration which produces dynamic dipolar interaction between adjacent magnetic vortices. In our previous experiment, we detected the excitation spectra of magnetostatically coupled two gyrating vortices in a pair of magnetic disks by applying the radio frequency ac current only to one of the disks [1]. The excitation mode observed in the single disk splits into two eigenfrequencies. In disk pairs four modes appear depending on the combination of polarities, which are characterized by rotational directions and the phase difference in core gyration. We here investigate the effect of the phase difference in gyration on the excitation of magnetostatically coupled vortices by varying the phase difference between ac excitation currents separately applied to the both disks. Firstly the collective excitation is induced by applying the ac current to one of the disks, and then the ac current with variable phase is applied to another disk to detect its indirect excitation signal at lower frequency clearly changed with respect to that at higher frequency and also the signal changed depending on chiralities.

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#### IE01

## Spin-orbit coupling induced spin torques in diluted magnetic semiconductors

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It has been recently demonstrated that appropriately designed spin-orbit coupling (SOC) can be used to generate spin torque (SOC-torque) in a single ferromagnet, without the need of an external polarizer [1,2]. This effect has been observed experimentally in both metallic systems [2] and dilute ferromagnetic semiconductors (DMS) [3]. The case of DMSs is of particular interest due to their tunable bulk inversion asymmetry allowing for the generation of both cubic and linear Dresselhaus SOC [3,4]. A theoretical understanding of the microscopic nature of the SOC-torque in DMSs is still needed to accurately interpret the recently obtained experimental results [4]. The present study addresses the nature of SOC-torque in DMS in the framework of Luttinger Hamiltonian (GaMnAs, InMnAs etc...). Based on kinetic-exchange model, the non-equilibrium spin transport in DMS is studied theoretically within the first-order Born approximation [4]. Both cubic and linear Dresselhaus SOC are examined and the angular dependence and magnitude of the SOC-torque are studied for a wide range of parameters. Interestingly, the cubic Dresseulhaus SOC does generate a spin torque as long as the Fermi surface is not spherical. The role of the carrier concentration and shape of the Fermi surface are emphasized and experimental implications are discussed.

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#### IE02

#### Spin orbit torque assisted domain wall depinning in Pt/Co/Pt

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It has been shown that, in perpendicular magnetic anisotropy (PMA) materials, the injection of a current results in a torque on the magnetization (SO-torque). This torque could participate to the current-induced (CI) domain wall (DW) motion, which up to now has been described in terms of spin transfer torque (STT) effect. This experiment aims at a comparison of SO-torque and STT in a Pt/Co/Pt stack. Using a focused ion beam technique we create a soft magnetic domain in the center of the stripe, and we study its evolution under the application of current and magnetic field. The measurements are done using a pump-probe detection scheme, suitable for stroboscopic experiments with sub-1 ns time-resolution. Our results show that the effect of the current is a broadening/narrowing of the magnetic domain, indicating that the SO-torque is dominant over STT and Joule heating. Moreover, we observe some features that are signature of SO-torque, notably the dependence of the CI-effect on the sign of the in-plane magnetic field and on the relative thickness of the two Pt Utifer layers. Finally, we emphasize the potential of our set-up for time-resolved study of CI

#### IE03

#### Piezo-electric control of the motion of a single domain wall

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We report piezo-electric control of a single domain wall (DW) driven by electrical current in a microbar fabricated in a perpendicularly magnetized GaMnAsP epilayer and by a magnetic field in the unpatterned film. By applying the piezo-voltage we tune in situ the magneto-crystalline anisotropy and thus modulating the stability of the internal wall structure which determines the value of the Walker breakdown. The low defect density in our high-quality GaMnAsP epilayer allows to study DW propagation in both the viscous and precessional regimes, before and after the Walker breakdown critical current, respectively. This combined with the piezo-electric controlled DW stability allows us to experimentally identify the adiabatic torque as the dominant spin-transfer mechanism and to vary the DW velocity by up to 500% at a fixed driving current and by more than two orders of magnitude at a fixed applied magnetic field. Our experimental conclusions are supported by a one-dimensional DW motion model and detailed micro magnetic simulations taking into account the individually measured complex magnetic anisotropies of the specific microbars. We conclude that our results are generic and applicable to DW motion in magnetic films with tunable magnetic anisotropy and Curie temperatures above room temperature.

#### IE04

Domain wall manipulation by spin currents in magnetic tunnel junctions Julie Grollier<sup>1</sup>, Peter Metaxas<sup>1</sup>, Jaao Sampaio<sup>1</sup>, Rie Matsumoto<sup>1</sup>, Andre Chanthbouala<sup>1</sup>, Alexey Khvalkovskiy<sup>1</sup>, Vincent Cros<sup>1</sup>, Abdelmadjid Anane<sup>1</sup>, Albert Fert<sup>1</sup>, Konstantin Zvezdin<sup>2</sup>, Akio Fukushima<sup>1</sup>, Hitoshi Kubota<sup>1</sup>, Kay Yakushij<sup>1</sup> and Shinji Yuasa<sup>1</sup>

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The spin transfer effect allows the manipulation of magnetic domain walls (DWs) in ferromagnetic wires by injection of dc currents, which has led in recent years to the proposal of a number of potentially revolutionary devices [1,2]. In most of these systems, the current is injected along the direction of DW propagation, and the current density required to move the DW has generally remained prohibitively high for applications. Recently, it was predicted [3] that higher efficiency could be obtained by using vertical current injection through multilayer devices, specifically in magnetic tunnel junctions (MTI). Our group has experimentally verified the higher efficiency DW motion with perpendicularly injected dc currents in MTJs, which was attributed to the out-of-plane spin transfer torque and the Oersted field [4]. We now present time-resolved measurements of DW motion under vertically injected pulsed currents, and demonstrate that DW speeds in excess of 500 m/s can be attained using current densities that are up to 2 orders of magnitude less than those required in devices that use lateral current injection. As such, this work represents an important advance in the race towards energy-efficient, high speed domain wall devices. We acknowledge the European Research Council (Stg 2010 No.259068).

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#### IF01

Large magnetoresistance in antiferromagnet tunnel junctions

Byong-guk Park<sup>1</sup>\*, Joerg Wunderlich<sup>2</sup>, X. Marti<sup>3</sup>, J. Hayakawa<sup>4</sup>, H. Takahashi<sup>2</sup> and T. Jungwirth<sup>5</sup> <sup>1</sup> Material Science and Technology, KAIST, Korea

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Antiferromagnets (AFM's) have been used in spintronics devices so far only to pin the magnetization direction of a ferromagnetic electrode through the exchange-bias effect. Spintronics devices whose transport is governed by AFM's have been theoretically proposed, but remain a great challenge for experimental realization. Here, we demonstrate a large magnetoresistance effect in a tunnel junction with an AFM electrode of IrMn and a nonmagnetic counter electrode [1]. In the device, the tunneling resistance depends on the magnetization direction of the IrMn electrode, so-called tunneling anisotropic magnetoresistance (TAMR) effect, which is based on the spin-orbit coupling. This is quite different from the conventional spin-valves where the resistance depends on the relative magnetization directions of two ferromagnetic electrodes. The magnetization direction of the AFM IrMn layer was manipulated with a relatively small magnetic field of 50mT by the exchange spring effect of coupled soft NiFe. In addition, the AFM TAMR provides a means to study the exchange-bias effect by an electronic transport measurement [2].

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#### IF02

#### A first-principles study on spin-dependent tunneling conductance in magnetic tunnel junctions with spinel-type MgAl<sub>2</sub>O<sub>4</sub> barrier Yoshio Miura<sup>1\*</sup>, Shingo Muramoto<sup>2</sup>, Kazutaka Abe<sup>1</sup> and Masafumi Shirai<sup>1</sup>

Yoshio Miura'\*, Shingo Muramoto', Kazutaka Abe' and Masafumi Shirai' <sup>1</sup> RIEC & CSIS, Tohoku University, Japan <sup>2</sup> RIEC, Tohoku University, Japan

Recently, a normal-spinel type MgAl<sub>2</sub>O<sub>4</sub> was grown on and Co<sub>2</sub>FeAl<sub>05</sub>Si<sub>05</sub> and bcc-Fe electrodes to explore new materials as a barrier layer in magnetic tunnel junctions(MTJs) [1,2]. The MgAl<sub>2</sub>O<sub>4</sub> showed a great advantage in the epitaxial growth compared to the MgO because of a small lattice mismatch less than 1% with bcc-type ferromagnetic metals at (001) face. In this study, we investigated the transport properties of Fe/ MgAl<sub>2</sub>O<sub>4</sub>/Fe(001) MTJs using the first-principles calculations. We found that Fe/ MgAl<sub>2</sub>O<sub>4</sub>/Fe(001) MTJ showed the coherent tunneling conductance through the delta 1 evanescent state as well as Fe/MgO/Fe(001) MTJs. However, the calculated tunneling magnetoresistance (TMR) ratio of Fe/MgAl<sub>2</sub>O<sub>4</sub>(1nm)/Fe(001) MTJ was about 160%, which was much smaller than that of Fe/MgO(1nm)/Fe(001) MTJ (1600%). This can be attributed to appearance of new conductive channels in the antiparallel magnetization configuration by the band-folding in the two dimensional Brillouin zone of in-plane wave vector of bcc-Fe, because the in-plane cell-size of MgAl<sub>2</sub>O<sub>4</sub> is twice of that of bcc-Fe. These results indicate that suppression of band-folding is essential to obtain a large TMR ratio in Fe/MgAl2O4/Fe(001) MTJs. This work was partially supported by JSPS through the FIRST Program and by JST through ASPIMATT program

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#### IF03

#### Enhanced tunnel magnetoresistance in magnetic tunnel junctions with an epitaxial Mg-Al-O barrier

Hiroaki Sukegawa, Seiji Mitani, Tomohiko Niizeki, Tadakatsu Ohkubo, Koichiro Inomata and Kazuhiro Hono

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Achievement of defect-free heterostructure is essential to establishment of high performance spintronics devices such as magnetic tunnel junctions (MTJs), spintransfer devices and spin-transistors. A spinel MgAl<sub>2</sub>O<sub>4</sub> has a potential for a future tunnel barrier because of the small in-plane lattice mismatch (< 1%) between MgAl<sub>2</sub>O<sub>4</sub> and bcc ferromagnetic alloys such as CoFe alloys and Heusler alloys. In addition, MgAl<sub>2</sub>O<sub>4</sub> is a nondeliquescent material, which is favorable for practical applications. Recently a large tunnel magnetoresistance (TMR) ratio of 117% was observed at room temperature (RT) in a fully epitaxial Fe/MgAl<sub>2</sub>O<sub>4</sub>/Fe(001) MTJ, indicating the enhancement of TMR by an epitaxial MgAl<sub>2</sub>O<sub>4</sub> barrier as well as an MgO barrier [1]. In this study we achieved a fully epitaxial CoFe/Mg-Al-O/CoFe(001) MTJ using a sputter deposition and a plasma oxidation of Mg-Al alloy layers. We demonstrated the very large TMR ratio of 308% at RT in the MTJ even though the barrier layer had an off-stoichiometric MgAl<sub>2</sub>O<sub>4</sub> composition (Mg<sub>40</sub>Al<sub>60</sub>-O<sub>x</sub>). This result suggests that Mg-Al-O is promising for a tunnel barrier of future spintronics devices using coherent tunneling.

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#### IF04

#### The memristive magnetic tunnel junction as a nanoscopic synapseneuron system

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Memristors cover a gap in the capabilities of basic electronic components by remembering the history of the applied electric potentials, and are considered to bring neuromorphic computers closer by imitating the performance of synapses. We used memristive magnetic tunnel junctions based on MgO to demonstrate that the synaptic functionality is complemented by neuron-like behavior in these nanoscopic devices. The synaptic functionality ordiginates in a resistance change caused by a voltage-driven oxygen vacancy motion within the MgO layer. The additional functionality provided by magnetic electrodes enabled a current-driven resistance modulation due to spin-transfer torque. We report on memristive magnetic tunnel junctions characterized by the simultaneous occurrence of resistive switching and tunnel magnetoresistance. Since resistivity provides a natural measure of the synaptic strength, and because of the bipolar nature of the resistance change, long term potentiation and long term depression were emulated. Furthermore, we show that the flux is a good variable for describing voltage-induced resistance variation, which provides the scope for the emulation of spike timing dependend plasticity as well.

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#### IF05

#### Tunnel magnetoresistance in perpendicularly magnetized Co<sub>2</sub>FeAl/ MgO/CoFeB magnetic tunnel junctions

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Perpendicular magnetization of Co<sub>2</sub>FeAl (CFA) full-Heusler alloy thin film was achieved with the high perpendicular magnetic anisotropy (PMA), Ku of 2 ~ 3×106 erg/cm3 in the structures of CFA/MgO, which is comparable to that of the conventional PMA materials[1]. The CFA thickness dependence of Ku was investigated at different annealing temperatures, which indicated that the Ku of CFA is contributed by the interfacial anisotropy between CFA and MgO. Furthermore, we successfully fabricated perpendicularly magnetized tunnel junctions (p-MTJs) using CFA full-Heusler alloy. The out-of-plane tunneling magnetoresistance (TMR) ratio of 53% at room temperature was observed in CFA/MgO/Co20Fe60B20 p-MTJs on Cr-buffered MgO(001) substrates. By inserting 0.1 nm Fe (Co<sub>50</sub>Fe<sub>50</sub>) between MgO and Co<sub>20</sub>Fe<sub>60</sub>B<sub>20</sub>, TMR ratio was significantly enhanced to 91% (82%) due to the improved interface. The annealing temperature dependence of TMR ratio was also investigated. Comparison of the bias voltage dependence of the differential conductance was carried out between MTJs with the core structures of CFA(1)/MgO(1.8)/ Co<sub>50</sub>Fe<sub>50</sub>(5) and CFA(30)/MgO(1.8)/Co<sub>50</sub>Fe<sub>50</sub>(5) (unit: nm), which suggested that the improvements of B2 ordering of ultrathin CFA and interface structures are required for enhancing coherent tunneling in order to achieve a higher TMR ratio. This work opens up a way for perpendicularly magnetized MTJs using full-Heusler allovs

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#### IG01

#### Transport anomalies due to critical valence fluctuations Kazumasa Mivake

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It has been gradually turned out in this decade that some Ce- and Yb-based heavy fermion metals exhibit quantum critical phenomena associated with critical end point of valence transition [1]. In particular, temperature dependences of the specific heat, the uniform susceptibility, the longitudinal relaxation rates of NMR, and the resistivity have critical exponents quite different from those of conventional magnetic critical phenomena of a la Moriya-Hertz-Millis, giving possible explanations of behaviors observed in beta-YbAlB4, YbAuCu4, YbRh2Si2, CeCu2(Si,Ge)2, CeRhIn5 and so on [2,3]. In this talk, we briefly summarize previous studies on the above aspects, and discuss transport anomalies due to the critical valence fluctuations (CVF). Namely, the effect of collective modes beyond the Boltzmann transport due to quasiparticles is analyzed generalizing the case of superconducting fluctuations near the transition point [3]. It is shown that the Hall coefficient, the Nernst coefficient, the Seebeck coefficient, and the thermal conductivity are greatly affected by the CVF, explaining some anomalies of the Hall coefficient, the Nernst coefficient, and the Seebeck coefficient observed in CeCu<sub>2</sub>Si<sub>2</sub> and CeRhIn<sub>5</sub> under pressure. This talk is based on collaborations with Shinji Watanabe and Osamu Narikiyo who are gratefully acknowledged.

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#### IG02

#### Fluctuations and quantum criticality in Eu ternary pnictides

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We have studied the magnetic properties of the Zintl compounds EuZn<sub>2</sub>(P, As, Sb)<sub>2</sub> by ESR. The ESR was measured on frequency 9.3 GHz in TE102 rectangular cavity at 4.2 - 300 K. Above 150 K in all cases, we observed the symmetric resonance lines Eu<sup>2+</sup> with ideal Lorentzian lineshape. At the temperature decreasing well before the antiferromagnetic ordering temperature we have observed an increasing of lineswidth and a decreasing of the resonance fileds, which in our case very good descriebed by Landau's theory of magnetic fluctuations. The fluctuational nature of such behaviour is confirmed by changing the shape of the resonance line from Lorentzian to Gaussian. The paramagnetic temperature  $\theta$ p, obtained from the temperature dependence of the ESR signal intensity, has the positive sign. We belive such behaviour of  $\theta$ p is connected with the proximity to quantum critical point and with the relevant instability of magnetic and crystal structures (instability B<sub>2</sub>X<sub>2</sub>-Jayers, competition nonmagnetic 4f0 and magnetic 447 states of Eu) The obtained ESR data are interpreted in terms of Bloembergen-Rowland's indirect exchange interaction through the valence electrons and the Falicov-Kimball-Kugel model the hybridization of localized electrons with tinerant electrons. Work was supported by Grants of RFBR and RAS.

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#### IG03

## Synchrotron x-ray spectroscopy study on the valence state in a- and $\beta$ -YbAlB<sub>4</sub> at low temperatures and high magnetic fields

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Valence fluctuation phenomena in  $\beta$ -YbAlB<sub>4</sub> and its polymorph  $\alpha$ -YbAlB<sub>4</sub> have attracted much attention. [1, 2] The possible quantum criticality due to the valence transition is theoretically predicted. [3] In the present study, valence states of in  $\alpha$ - and  $\beta$ -YbAlB<sub>4</sub> have been investigated by the x-ray absorption spectra (XAS) in SPring-8 at temperatures from 2 to 280 K. The partial fluorescence yield mode and direct transmission method were used. High magnetic field XAS measurement was done on  $\beta$ -YbAlB<sub>4</sub> at 0 K in pulsed magnetic fields. Observed Yb valence is 2.8 ± 0.05 in  $\alpha$ - and  $\beta$ -YbAlB<sub>4</sub> at 5 and 2 K, respectively. This value is slightly higher than the previously reported results 2.73± 0.02 for  $\alpha$ -YbAlB<sub>4</sub> and 2.75± 0.02 for  $\beta$ -YbAlB<sub>4</sub>. [1] The valence is found to gradually increase with increasing temperature toward the trivalent state; the valence increment from the low temperature to 200 K is about 0.02 in both materials. We also found a small increase of Yb valence by a magnetic field of 31 T.

[1] M. Okawa et al, PRL 104, 247201 (2010). [2] Y. Matsumoto et al., Science 331, 316 (2011). [3] S. Watanabe et al., PRL 105, 186403 (2010).

#### IG04

Metal-Insurator crossover accompanied by the dual nature of 5f electrons with localized and itinerant character in US2

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Uranium dichalcogenide  $\beta$ -US2 is a low carrier system with metallic behavior at room temperature, while insulating with opening a charge gap Eg-90K at low temperatures with colossal magneto-resistance effect  $10^8$  for H~10 T. The magnetic properties can be well described by a localized model. In fact clear CEF excitations have been observed at ~7 meV, indicative of the localized character of 5f electrons at low temperatures. With elevating temperature ( > 100 K) above the conduction and magnetic excitation gap, however, we observed magnetic quasi-elastic response accompanied by the metal-insulator crossover with a large resistivity drop in the order of  $10^6$ . This suggests a new mechanism of metal-insulator crossover as a consequence of the change in 5f character between localized and itinerant. We succeeded to separate the magnetic susceptibility into two parts, the contribution from the localized and itinerant part, by carrying out Kramers-Kronig transformation of the imaginary part of the dynamical susceptibility measured by magnetic excitation with neutron inelastic scattering. We revealed a dual nature of 5f electron system, growing itinerant character with elevating temperature beyond the charge gap.

#### IG05

## Valence transition induced by pressure and magnetic field in antiferromagnet EuRh,Si,

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EuRh<sub>2</sub>Si<sub>2</sub>, which crystallizes in the tetragonal ThCr<sub>2</sub>Si<sub>2</sub> type structure, is an antiferromagnet with TN=23K and the Eu ion is in the divalent state (4f7, J=7/2). Recently, we have found that the application of pressure of 1.0GPa collapses the antiferromagnetism and simultaneously induces a first-order valence transition at Tv~30K [1]. Under higher pressure than 1.0GPa, the sample shows a temperature-independent Pauli-paramagnetic behavior (a non-magnetic phase) below Tv while a Curie-Weiss paramagnetism corresponding to Eu<sup>2+</sup> (a magnetic phase) above Tv. In this study, we have measured M-B curves under pressure to examine influences of magnetic field on the non-magnetic phase. A metamagnetic transition accompanied by a large hysteresis is observed at 20K under 1.0GPa. The magnetization jumps from 1.0 $\mu$ B/Eu at B=3.6 T to 1.6 $\mu$ B/Eu at B=5.0 T with increasing magnetic field. This phenomenon corresponds to a field-induced valence transition from the non-magnetic phase.

[1] A. Mitsuda et al. J. Phys. Soc. Jpn. 81 (2012) 023709

#### IH01

#### Orbital ordering and multiphase separation at manganite interfaces

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Interfacial effects between La<sub>23</sub>Sr<sub>13</sub>MnO<sub>3</sub> (LSMO) thin films and different capping layers are analyzed by using x-ray absorption techniques, x-ray diffraction and transport measurements. We make use of x-ray linear dichroism (XLD) to show that, independently of the capping layer, LSMO films exhibit a preferential occupation of Mn 3d 322-r2 eg orbitals at the interface. The strength of such an orbitally ordered phase does not depend on the capping layer material neither is related to the previously observed degradation of the Mn oxidation state nor to the magnetic properties. Moreover, its intensity is similar to that observed in uncapped films thus indicating its origin lies on the symmetry breaking at the LSMO surface. Transport measurements across LSMO/capping layer interfaces allow estimating the scale length of this effect showing that the disruption of the double exchange ferromagnetic (DE-FM) phase occurs only by the interface, in fact around 2 to 4 unit cells. Thus, XAS local techniques point to a complex scenario at the manganite's interfaces in which orbital ordering and multiphase segregation coexist.

#### IH02

## Conical spin-spiral state in an ultra-thin film driven by higher-order spin interactions

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The magnetic properties of transition-metal nanostructures are commonly explained based on the interplay of Heisenberg exchange, Dzyaloshinskii-Moriya (DM) interaction and magnetocrystalline anisotropy while higher order terms such as the biquadratic exchange and the four-spin interaction are typically neglected due to their small strength. Here, we demonstrate that higher-order terms can play a crucial role for the magnetic ground state and report as an example a transverse conical spin-spiral state in an ultra-thin film composed of two atomic layers of Mn on W(110) [1]. This spin structure is characterized by magnetic moments rotating on a cone that is perpendicular to the [001] propagation direction of the spin-spiral with a periodicity of 2.4 nm. The cones of nearest-neighbor Mn atoms point into opposite directions which results in nearly antiferromagnetic alignment. This intriguing spin structure has been resolved on the atomicscale using spin-polarized scanning tunneling microscopy and confirmed to be the ground state by firstprinciples calculations based on DFT. Our calculations also reveal that the canting of the spins is induced by higher-order exchange interactions, while the spiraling along the [001]-direction is due to frustrated Heisenberg exchange and DM interaction.

[1]Y.Yoshida et al., Phys. Rev. Lett. (in press)

#### IH03

#### Non-collinear magnetic ground state in finite metallic chains

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The investigation of one-dimensional magnetic structures on the atomic scale is a quite challenging task, since they are very prone to thermal and quantum fluctuations [1]. Here, we present spin-polarized scanning tunneling microscopy (SP-STM) measurements of bi-atomic Fe chains on  $(5\times1)$ -Ir(001). The chains exhibit a modulation along the chain axis with a three atom periodicity in an external magnetic field perpendicular to the surface which is not present without the magnetic field [2]. This is due to a spin spiral ground state, whose thermally induced magnetization switching leads to a time-averaged SP-STM signal, with the external field preferring one of the magnetization directions. However, due to the rigid coupling of the Fe atoms' magnetic moments along the chain axis, the spin spiral order is not affected by thermal fluctuations at the measurement temperature. Therefore, changes in the spin direction of one chain end can be probed tens of nanometers away, suggesting a new way of transmitting information about the state of magnetic objects on the nanoscale.

[1] N.D. Mermin and H. Wagner, Phys. Rev. Lett. 17, 1133 (1966). [2] M. Menzel et al., submitted.

#### **IH04**

## Magnetism and the thermodynamics of Fe-Pt surface alloy formed at Pt(110) surface

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Surface alloy formation at noble metal surface shows very interesting phenomena. If one of the constituent atoms is magnetic, then its magnetic behavior evolves depending on the ratio of composition and the temperature. We will show atomistic view of Fe-Pt surface alloy formation at Pt(110)-missing low surface using scanning tunneling microscopy. When the external Fe atoms are deposited on top of Pt(110)-(2x1) surface, they are located at the missing row position. When the temperature increases, they exchange the position with the Pt atoms nearby due to the increase of the entropy term to minimize the free energy. Pt atoms located at the missing row position is energetically unstable. The change of magnetic characteristics through this order-disorder phase transition has been probed with the surface magneto-optic Kerr effect. With the aid of density functional theory, all the nanoscopic description of this magnetic surface surface surface alloy formation will be discussed.

#### IH05

## A study of antiferromagnetic/ferromagnetic systems using x-ray magnetic dichroism

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Epitaxial CoO/Fe/Ag(001) system was investigated using X-ray Magnetic Circular Dichroism (XMCD) and X-ray Magnetic Linear Dichroism (XMLD). XMCD measurement was used to measure the Fe hysteresis loops and to determine the unidirectional (exchange bias), uniaxial, and the 4-fold anisotropy. XMLD measurement was used to determine the response of the CoO spins. We find that the CoO spins consist of rotatable and frozen spins with respect to the Fe spin rotation, and only the Fe uniaxial magnetic anisotropy follows the CoO frozen spins. We find that only uniaxial magnetic anisotropy is correlated to the frozen spins.

#### IH06

## Magnetic properties and microscopic structures of ultrathin Co/ $\sqrt{3}\times\sqrt{3}$ -R30°-Ag/Si(111) films

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Combined scanning tunneling microscopy and surface magneto-optic Kerr effect studies were employed to study the relation between magnetic properties and microscopic interfacial structures of ultrathin  $Co/\sqrt{3}\times\sqrt{3}$ -R30° -Ag/Si(111) films. On the top of  $\sqrt{3}\times\sqrt{3}$ -R30° -Ag/Si(111), pure Co clusters form without disrupting the  $\sqrt{3}\times\sqrt{3}$ -R30° structure of the Ag buffer layer. The great strain due to the large mismatch between Co and the substrate influences the nucleation of Co atoms to form large clusters. No magnetic hysteresis in the polar configuration was observed for films thinner than 10 monolayers. The easy axis of magnetization is in the surface plane. Capping  $\sqrt{3}\times\sqrt{3}$ -R30° -Ag on the top of Si(111) surface before the deposition of Co overlayers can efficiently reduce the nonferromagnetic Co-Si compound to zero thickness. For Co coverage between 2.9 and 4.2 monolayers, the lowered Curie temperature in ultrathin films is observed. Due to the existence of a smooth interface between Co and the  $\sqrt{3}\times\sqrt{3}$ -R30° -Ag buffer, the coercivity for Co/ $\sqrt{3}\times\sqrt{3}$ -R30° -Ag/Si(111).

#### II01

## Giant anomalous Hall effect in magnetic topological insulator Qi-kun Xue

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Topological insulator (TI) is a time-reversal-invariant and therefore magnetic field free version of the quantum Hall system. Breaking the time-reversal symmetry by ferromagnetic order in TIs can induce quantum Hall effect in the absence of external magnetic field, known as quantized anomalous Hall (QAH) effect. Here we report the realization of the carrier-independent ferromagnetism in thin films of Cr-doped TI ( $Bi_s Sb_{1,2} Te_5$ . Despite the opposite types of carriers and one order of magnitude change in carrier density, robust ferromagnetism with nearly constant Curie temperature is observed. Remarkably, anomalous Hall effect is significantly enhanced with decreasing carrier density, with the anomalous Hall angle reaching an unusually large value of 0.2 and the zero field Hall resistance reaching one quarter of the quantum resistance (h/e2), indicating the approaching QAH regime. The work paves the way to realize the QAH effect.

#### II02

A rich rashba system created on the surface of a topological insulator Phil King<sup>1\*</sup>, A De La Torre<sup>1</sup>, Felix Baumberger<sup>1</sup>, M. Bianchi<sup>2</sup>, R. Hatch<sup>2</sup>, Philip Hofmann<sup>2</sup>, M.S. Bahramy<sup>3</sup>, R. Arita<sup>3</sup>, N. Nagaosa<sup>3</sup>, J. I. Mi<sup>4</sup>, B. Iversen<sup>4</sup> and G. Balakrishnan<sup>5</sup> <sup>1</sup> School of Physics and Astronomy, University of St Andrews, United Kingdom <sup>2</sup> Department of Physics and Astronomy, Aarhus University, Denmark <sup>3</sup> Correlated Electron Research Group, RIKEN-ASI, Japan

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Topological insulators are famed for their unconventional surface states which span from the bulk valence to conduction bands [1]. Adsorbing n-type dopants at the surface of these compounds in ultra-high vacuum, analogous to the application of an external gate voltage in a device, we show that new electronic states emerge that co-exist with the topological surface state [2,3]. A strong downward band bending causes quantization of the bulk conduction band states into a two-dimensional electron gas (2DEG) [2], while the bulk valence band states also become quantized due to the narrow valence bandwidth at the zone centre. This provides a novel opportunity to probe the interplay of dimensionality and topological order. Moreover, the 2DEG states can be driven to develop a large Rashba spin splitting [3]. The resulting multi-subband 2DEG provides a very nich Rashba system. The inner-most Fermi surface sheets display the chiral in-plane spin texture characteristic of classic Rashba states such as the L-gap surface state on Au(111). However, the outermost Fermi pockets display increasing degrees of hexagonal warping, resulting in an increasing outof-plane spin canting. We investigate this system by angle-resolved photoemission and model calculations.

[1] M.Z. Hasan & C.L. Kane, Rev. Mod. Phys. 82 (2010) 3045 [2] M. Bianchi, et al., Nature Commun. 1 (2010) 128 [3] P.D.C. King et al., Phys. Rev. Lett. 107 (2011) 096802

#### 1103

From topological semimetals towards insulators. First-principles study Stanislav Chadov<sup>1</sup>\*, Claudia Felser<sup>1</sup>, Kristina Chadova<sup>2</sup>, Diemo Kodderitzsch<sup>2</sup> and Hubert Ebert<sup>2</sup>

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The presence of disorder is often seen as a destructive mechanism which must be reduced by any means. In present study we attempt to make it constructive using the robustness of the spin current in topologically non-trivial systems with respect to the time-reversal symmetric perturbations. Based on the fully-relativistic first-principles calculations involving the Coherent Potential Approximation (CPA) and Kubo-Greenwood formalism, we inspect the disorder-affected transport properties of random alloys between topologically non-trivial and trivial materials. The subsequent analysis encounters few interesting aspects. First, it straightforwardly leads to an alternative scheme improving the spin-Hall transport in a large class of topological semimetals and indicates the possibility of the topological Anderson insulator state. Second, we obtain the quantized measure of the spin-Hall conductivity in 3D without explicit calculation of the topological state of the material based on a purely bulk information.

#### **II04**

#### Engineering and manipulating topological qubits in 1D quantum wires Panagiotis Kotetes<sup>1</sup>, Alexander Shnirman<sup>2</sup> and Gerd Schon<sup>1</sup>

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Recently, the interest in topological quantum computing has grown due to the appearance of promising platforms for realizing the long sought Majorana bound states. Among the proposals that seem suitable for engineering Majorana bound states, the most prominent involves a 1D semiconducting quantum wire in proximity to a bulk s-wave superconductor, where in addition a Zeeman field is applied. In this work we explore novel routes for realizing topological qubits and we further investigate alternative possibilities for performing qubit operations without for instance the use of electrostatic gates. Specifically we focus on the interplay of the phases of the magnetic field and the superconducting order parameter, illuminating the significance of keeping intact the spin degree of freedom for engineering and braiding the topological qubits. Finally, by considering an appropriate junction setup, we demonstrate novel features in the Josephson effect of these systems, such as the magnetic control of the supercurrent.

#### IJ01

## Magnetic structure of iron borate ${\rm SmFe}_3({\rm BO}_3)_4$ : A neutron diffraction study

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Rare earth ferroborates  $RFe_3(BO_3)_4$  have recently attracted widespread interest due to their interesting magnetic, magnetoelectric and multiferroic properties [1]. The magnetic structure is determined by a low dimensional element in the crystallographic structure which sees well separated chains of helicoidally arranged Fe atoms running along the crystallographic c-axis. No direct R-R or R-O-R interactions exist and the development and arrangement of long range magnetism is determined by the strongly polarizing Fe-sublattices and the competition between the iron and the rare earth anisotropies [2]. We determined the magnetic structure of SmFe<sub>3</sub>(BO<sub>3</sub>)<sub>4</sub> by neutron powder diffraction using a double wall container and diluting the sample with Al-powder in order to reduce the strong natural neutron absorption of Sm. The antiferromagnetic structure developed below TN=33K is of the easy plane type as already found for the R=Er, Y and Ho (above 6K) compounds [3,4]. The fit of the magnetic peaks was best using a magnetic formfactor for Sm<sup>3+</sup> where C2=5.42 giving a total magnetic moment of  $\mu$ Sm=0.24 $\mu$ B composed of an orbital moment of  $\mu$ SmL=1.3uB and an spin moment of  $\mu$ Sm=21.06uB.

[1] e.g.: R.P. Chaudhury et al., PRB 81, 2010, 220402 [2] C. Ritter et al., J. Phys.: Condens. Matter 19, 2007, 196227 [3] C. Ritter et al., J. Phys.: Condens. Matter 20, 2008, 365209 [4] C. Ritter et al., J. Phys.: Condens. Matter 22, 2010, 206002

#### IJ02

## Preparation and characterization of $Sr_3Fe_2O_{7\mbox{-}x}$ for different oxygen contents

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Iron(IV)- and ruthenium(IV)-based perovskite phases have proven rich in novel physics. The SrFeO<sub>3-y</sub> system exhibits a wide range of unusual magnetic phases, from five distinct forms of helical magnetism when fully oxygenated [1] to a phase with frustrated, disordered Fe<sup>4+</sup> moments for an oxygen content of 2<sup>3</sup>/<sub>4</sub> [2]. The Ru-based intergrowth phase Sr<sub>3</sub>Ru<sub>2</sub>O<sub>7</sub> exhibits metamagnetic quantum critical points in applied field [3]. However, the magnetic phase diagram of its Fe<sup>4+</sup> analogue Sr<sub>3</sub>Fe<sub>2</sub>O<sub>7-x</sub> remains largely unexplored. We describe single crystal growth by the floating zone technique, oxygen annealing, and transport and magnetic measurements of large Sr<sub>3</sub>Fe<sub>2</sub>O<sub>7-x</sub> crystals suitable for neutron diffraction experiments. This work lays the foundation for comprehensive doping-dependent studies of Sr<sub>3</sub>Fe<sub>2</sub>O<sub>7-x</sub>'s magnetic phase diagram and magnetic excitations.

 S. Ishiwata et al., Phys. Rev. B 84, 054427 (2011) [2] M. Schmidt et al., J. Phys.: Condens. Matter 15, 8691 (2003) [3] R.A. Borzi et al., Science 315, 214 (2007)

#### IJ03

#### The magnetic structures of CoPS<sub>3</sub> and NiPS<sub>3</sub>

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The compounds CoPS<sub>3</sub> and NiPS<sub>3</sub> are candidates for two dimensional (2D) antiferromagnets, with the transition metal atoms lying on a honeycomb lattice. Magnetic structures for them were published in the 1980s, based on neutron powder diffraction data. Both materials were believed to be collinear antiferromagnets, although the antiferromagnetic axis differed between the compounds, with a magnetic propagation vector of [0 1 0]. We have made a careful study of the magnetization and neutron diffraction from single crystals, taking advantage of recent developments in instrumentation for neutron Laue diffraction has very few observable peaks and can, indeed, be fitted with the quoted magnetic structure. However, our single crystal data prove that the reported magnetic structures are erroneous, with the observation of many new Bragg peaks that could be indexed with [ $1/2 \ 1/2 \ 1/6$ ]. Our experiments have an impact on the interpretations of the magnetic Hamiltonians, and we will contrast our findings with two other members of the family: MnPS<sub>18</sub> which is a good example of a 2D leisenberg antiferromagnet, and structures and the reported of a 2D leisenberg antiferromagnet, and FePS<sub>18</sub>, which is a good example of a 2D leisenberg antiferromagnet, and

#### IJ04

**Magnetic properties in Fe-doped LnCo<sub>1-x</sub>Fe<sub>x</sub>AsO (Ln=La, Sm) systems** Yuke Li<sup>1</sup>, Chenyi Shen<sup>2</sup>, Yongkang Luo<sup>2</sup>, Chen Lv<sup>2</sup>, Qian Tao<sup>2</sup>, Jianhui Dai<sup>1</sup>, Guanghan Cao<sup>2</sup> and Zhuan Xu<sup>2</sup>\*

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We study the effect of Fe doping in Co-rich side of LnCo<sub>1-x</sub>Fe<sub>x</sub>AsO (Ln=La,Sm) systems. As Co is replaced partially by Fe ions, the FM transition temperature Tc was seriously suppressed to lower temperatures and it disappears around x = 0.3. In the SmCo<sub>1-x</sub>Fe<sub>x</sub>AsO system, the undoped SmCoAsO undergoes three magnetic phase transitions with decreasing temperatures, a FM transition around Tc of 80 K, and a FM to AFM transition below TN1 of about 45 K, and finally an AFM order of Sm<sup>3+</sup> ion at TN2 of 5.6 K. With increasing Fe content, both Tc and TN1 gradually decreases. Around x = 0.3, the FM order fof Co ions disappears above 2 K, and meanwhile, the FM-to-AFM transition cannot be observed either. However, the AFM order of Sm<sup>3+</sup> ion is robust, and TN2 slightly decreases in for x = 0.3. We concluded that in the 1111 type Co-based LnCo<sub>1-x</sub>Fe<sub>x</sub>AsO systems, 4f electrons. On the contrast, the magnetism of rare earth elements is robust to the variations of 3d electrons. A rich magnetic phase diagram of the LnCo<sub>1-x</sub>Fe<sub>x</sub>AsO systems is established.

#### IJ05

#### Symmetry argument of cyano-bridged copper-molybdenum complexes Jun Ohara\* and Shoji Yamamoto

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Photoreactive magnets of general formula Cu2[Mo(CN)8]·nH2O [1] are composed of ocatacyanomolybdate ions of square antiprism configuration assembled within a tetragonal lattice. CuII<sub>2</sub>[MoIV(CN)<sub>8</sub>]·nH<sub>2</sub>O irradiated with blue light turns into CuICuII[MoV(CN)8] nH2O exhibiting a spontaneous magnetization. The induced magnetization is stable for long time below 100 K and is photoreversible. We group-theoretically study competing magnetic phases in this intriguing material. While individual octacyanomolybdate ions of square antiprism configuration are D4d-invariant, the point symmetry of their assembly reads C4h. A single octacyanomolybdenum complex is invariant to polyaxial rotations but has no inversion symmetry. The assembly is merely invariant to uniaxial rotations but is further symmetric with respect to inversion. Besides the paramagnetism (PM) and ferromagnetism (FM), which correspond to the major states before and after photoirradiation, respectively, we find three antiferromagnetic (AFM) states as well as nonmagnetic charge-density-wave and bond-order-wave states of no major interest. The AFM states consist of a spin density wave within the molybdenum sublattice, that within the copper sublattice, and that spreading out in both sublattices, respectively, the latter two of which closely compete with both PM and FM states and therefore possibly have significant effects on the photoreactive magnetization mechanism.

[1] S. Ohkoshi et al., JACS 128, 270 (2005).

#### IJ06

## Phase diagram of the dzyaloshinskii-moriya helimagnet $Ba_2CuGe_2O_7$ in canted magnetic fields

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The evolution of different magnetic structures of the non-centrosymmetric tetragonal antiferromagnet Ba<sub>2</sub>CuGe<sub>2</sub>O<sub>7</sub> is systematically studied as function of the orientation of the magnetic field H. Neutron diffraction in combination with measurements of magnetic susceptibility and specific heat shows a virtually identical phase diagram of Ba<sub>2</sub>CuGe<sub>2</sub>O<sub>7</sub> for H confined in both the (1,0,0) and (1,1,0) plane. The existence of a recently proposed incommensurate double-k AF-cone phase [1] is confirmed in a narrow range for H close to the tetragonal c-axis. In contrast, for large angles enclosed by H and the c-axis, a complexly distorted non-sinusoidal magnetic structure has recently been observed [1]. We show that its critical field systematically increases for larger canting. Measurements of magnetic susceptibility and specific heat finally indicate the existence of a incommensurate/commensurate (I/C) transition at H =9 T, applied in the basal (a,b)-plane [2]. The observation of this transition in combination with both odd and even harmonics of the propagation vector agree with a non-planar, distorted magnetic structure, possibly caused by the stagegred component of the Dzyaloshinskii-Moriya vector.

[1] S. Muehlbauer et al. Phys. Rev. B 84, 180406 (2011) [2] S. Muehlbauer at al. in preparation (2012)

#### **JA01**

## Close relationship between superconductivity and the bosonic mode in $Ba_{0,6}K_{0,4}Fe_2As_2$ : A pairing glue for superconductivity Hai-hu Wen\*

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Center for Superconducting Physics and Materials, National Laboratory of Solid State Microstructures and Department of Physics, Nanjing University, Nanjing 210093, China Since the discovery of high temperature superconductivity in the iron pnictides and chalcogenides in early 2008, it remains unclear whether there is a glue for the pairing and what it is. There is a debate at this moment about whether the pairing is due to a retarded electron-boson interaction, as in the conventional phonon-mediated superconductors, although here the bosonic excitations may be the antiferromagnetic spin fluctuations. In this paper, we show the clear evidence of the bosonic mode with the energy identical to that of the neutron resonance, and its close relationship with superconductivity in a strongly coupled superconductor Ba<sub>0.6</sub>K<sub>0.4</sub>Fe<sub>2</sub>As<sub>2</sub> based on the measurements of scanning tunneling microscopy. Our data also indicate an interesting asymmetric feature of the tunneling spectrum, i.e., both the superconducting coherence peak and the peak related to the bosonic mode are always stronger on the positive-bias side, and vanish simultaneously within the vortex cores. Collaborated with Zhenyu Wang, Huan Yang, Delong Fang, Bing Shen

#### JA02

Broken time-reversal symmetry superconducting state in LiFeAs

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Here, and despite the current experimental support for s-wave superconductivity, we show clear evidence for broken time reversal symmetry within the superconducting state of LiFeAs through magnetic torque and magnetization measurements at high fields and in high quality single crystals close to optimal doping. Both the torque and the magnetization reveal a change in the sign of the magnetic inversibility/hysteresis as one approaches the upper critical Hc2, from a clear diamagnetic like response dominated by the pinning of vortices according to the Bean model, to a state with a 1000 times smaller paramagnetic like irreversibility which disappears at Hc2. If diamagnetism results from screening super-currents, the paramagnetic like response results from currents circulating in the opposite sense mimicking a field-dependent magnetic moment below Hc2. The quality of our samples is indicated by the observation of the de Haas van Alphen effect and the 75As NMR spectrum which reveals a single Lorentzian line, indicating the absence of magnetism or inhomogeneities. We conclude that the superconducting state in LiFeAs must undergo a field-induced phase-transition towards a chiral superconducting state as the smaller superconducting gap(s) is suppressed by the field. Small anomalies in the reversible component of the magnetization supports this scenario

#### JA03

#### Observation of anomalous magneto-resistance behavior near the inplane upper critical field in Sr(Fe,Ni)<sub>2</sub>As<sub>2</sub> single crystals

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Department of Physics and Astronomy, Seoul National University, Korea Department of Chemistry, Pohang University of Science and Technology, Korea

We have investigated Hc2 behaviors in various iron based SC including Sr(Fe,Co)<sub>2</sub>As<sub>2</sub>, Fe(Se,Te) and LiFeAs [1-3]. In this particular work, we measured the resistivity of Sr(Fe,Ni)<sub>2</sub>As<sub>2</sub> up to 35 T and report our unexpected findings in the magneto-resistance just above the upper critical fields around the H // ab direction at low temperatures. Apart from the transition in H // c direction, in which the resistivity is saturated after abruptly appearing from the zero resistive superconducting state with increasing magnetic field, the resistivity gradually increases after it rapidly appears until it reaches the value saturated for the H // c direction. This can be interpreted as 1) the rate of superconducting pair breaking changes or 2) the normal state resistivity is suppressed with emerging an abnormal state with the strong magneto-resistance with the in-plane field component. We will discuss the origin of this anomalous magneto-transport behavior based on the band calculation result, in connection with the possible Fulde-Ferrel-Larkin-Ovchinnikov (FFLO) ground state or the competing ground state.

[1] S. Khim et al., Physica C 470, S317 (2010).
 [2] S. Khim et al., Phys. Rev. B 81, 184511 (2010).
 [3] S. Khim et al., Phys. Rev. B 84, 104502 (2011).

#### JA04

Superconductivity in an Einstein solid:  $A_xV_2AI_{20}$  (A = Ga, Al) Atsuhi Onosaka, Junichi Yamaura, Yoshihiko Okamoto and Zenji Hiroi\* *ISSP, University of Tokyo, Japan* 

We study a series of intermetallic compounds  $A_rV_2Al_{20}$  with A = Ga, Al, Sc, Y, La, and Lu. They crystallize in the CeCr<sub>2</sub>Al<sub>20</sub> structure which contains a large atomic cage partially occupied by Ga/Al/Sc atoms and completely occupied by Y/La/Lu atoms. Low-energy local modes, which may be similar with what is now called rattling, have been observed for A = Ga and Al. Superconductivity was also reported at 1.6 K for A = Al and Ga, though experimental data was not given. In order to investigate the rattling and its effect on the superconductivity, we synthesized a series of polycrystalline samples and measured heat capacity, resistivity and magnetic susceptibility. Superconductivity is observed at  $T_c = 1.66$ , 1.50, 1.84, and 0.69 K for  $A = Ga_{0,2}$ ,  $Al_{0,3}$ ,  $Sc_{0,2}$ , and Y, respectively, and not above 0.5 K for A = La and Lu. It is suggested that the enhancement in Tc for the Ga, Al and Sc compounds is due to the rattling. In addition, we found a markedly large diamagnetism for A = Y and La, but not for A = Ga, Al, Sc, and Lu, despite the apparent similarity of the electronic structures among these compounds.

#### JA05

## Non-unitary triplet pairing in the centrosymmetric superconductor LaNiGa,

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Symmetry breaking is a central concept of physics for which superconductivity provides a paradigm. In a conventional superconductor gauge symmetry is broken, while unconventional superfluids and superconductors break other symmetries as well. Recently, we showed experimentally time-reversal symmetry (TRS) breaking in the noncentrosymmetrie superconductor LaNiC<sub>2</sub> [1]. Using group theory analysis we concluded that LaNiC<sub>2</sub> has only four possible symmetry states consistent with this observation, all of which are non-unitary [2]. This hinges on the low symmetry. C2v, of this material in which all of the irreducible representations are one dimensional. Here, we report the results from muon spin relaxation/rotation experiments on the intermetallic superconductor LaNiGa<sub>2</sub> (Te-2K) which, although centrosymmetric, has a similarly low-symmetry structure D2h. We find again that the onset of superconductivity coincides with the appearance of spontaneous magnetic fields, implying that in the superconducting state TRS is broken. Group-theoretic analysis of the possible pairing symmetries again suggests only four triplet states compatible with this observation, all of which are also non-unitary [3]. A comparison will be made between LaNiGa<sub>2</sub> and LaNiC<sub>1</sub> and we will propose that these superconductors represent a new family of paramagnetic non-unitary triplet superconductors.

 A. D. Hillier et al Phys. Rev. Lett. 102 117007 (2009) [2] J. Quintanilla et al Phys. Rev. B 82 174511 (2010) [3] A. D. Hillier et al submitted to Phys. Rev Lett.

#### JB01

Magnetoelectric effects and related phenomena in non-collinear spiralspin systems Tsuvoshi Kimura

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Among various multiferroics, extensive studies of ferroelectrics originating from magnetic orders, i.e., magnetically-induced ferroelectrics in which the inversion simmetry breaking and resultant ferroelectricity are induced by complex magnetic orders (e.g., noncollinear spiral order), have been triggered almost a decade ago by the discovery of multiferroic nature in a perovskite-type rare-earh manganites TbMnO<sub>3</sub>. Because the complex magnetic orders often arise from the competition between nearest-neighbor and further-neighbor magnetic interactions, systems containing competing magnetic interactions are promising candidates for magnetoelectric multiferroics [1]. Thus, it is no longer so difficult to find new magnetoelectric multiferroics related to noncollinear spiral-spin orders have been discovered in the past few years. In this presentation, I show recent progress on our studies of magnetoelectric effects and related phenomena in noncollinear spin-spiral systems [2].

[1] T. Kimura, "Spiral magnets as magnetoelectrics" Annu. Rev. Mater. Res. 37, 387 (2007). [2] T. Kimura, "Magnetoelectric hexaferrites", Annu. Rev. Condens. Matter Phys. (to be published).

#### **JB02**

## Electric field control of nonvolatile four-state magnetization at room temperature

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The control of magnetization by an electric field at room temperature remains as one of the great challenges in materials science. We demonstrate the realization of this long-sought capability in a multiferroic hexaferrite  $Ba_{0.52}Sr_{2.48}Co_3Fe_2AO_{41}$  single crystal. The electric polarization in this crystal rapidly increases in low magnetic fields (~5 millitesla), and its magnetoelectric susceptibility reaches the highest value (3200 picosecond per meter) among single-phase materials. The magnetization is then modulated up to 0.62 µB per formula unit in an electric field of 1.14 MV/m. Furthermore, this compound allows nonvolatile, magnetoelectric read-and-write operations entirely at room temperature. Four different magnetic/electric field writing conditions generate repeatable, distinct magnetization versus electric field curves without dissipation, offering an unprecedented opportunity for multi-bit memory- or spintronic device-applications.

#### **JB03**

## Low magnetic field reversal of electric polarization in Y-type hexaferrites

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Hexaferrites with spiral magnetic structures provide great opportunities for searching for new multiferroic materials with a high operating temperature and great magnetoelectric sensitivity. Here we report on the magnetically tunable ferroelectricity and the giant magnetoelectric sensitivity up to 250 K in a series of Y-type hexaferrite with the formula of BaSrZnCoFe<sub>12x</sub>Al<sub>Q02</sub>. Not only the magnitude but also the sign of electric polarization can be effectively controlled by applying low magnetic fields (a few hundreds of Oe) that modifies the spiral magnetic structures. The magnetically induced ferroelectricity is stabilized even in zero magnetic fields. Decayless reproducible flipping of electric polarization by oscillating low magnetic fields is shown. The maximum magnetoelectric coefficient reaches a high value of  $\sim 3.0 \times 10^3$  ps/m at 200 K.

#### **JB04**

## Nearest - next-nearest neighbor exchange frustrated quantum chain antiferromagnets: Recent results

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Lately, one-dimensional quantum chain antiferromagnets which feature magnetic frustration due to the presence of nearest and next-nearest-neighbor superexchange and super-superexchange interaction have attracted considerable attention. Such systems can e.g. be realized in compounds which contain  $MX_2$  ribbon chains where the magnetic species M are located in a square planar anion environment (X = an oxygen or a halogen atom). Some of these ribbon-chain antiferromagnets exhibit unusual magnetic ground states which e.g. support effective magnetoelectric coupling and multiferroic behavior. I review some of our recent investigations on the well-known nearest - next-nearest neighbor quantum chain antiferromagnetic LiCuVO<sub>4</sub> and report first results gained on some new systems, CuCrO<sub>4</sub>, TiPO<sub>4</sub> and CuCl<sub>2</sub>.

#### **JB05**

#### Multiferroic properties of layered triangular compounds

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One of the most puzzling families of multiferroics is characterized by stackings of triangular planes, on which geometrically induced frustration of antiferromagnetic exchanges leads to complex non-collinear magnetic structures, like in the delafosite CuCrO<sub>2</sub> [1, 2]. The proposed microscopic origin of this spin-driven ferroelectricity is based on the variation of the ligand-metal orbital hybridization, as a result of spin-orbit coupling [3]. In this context, tuning hybridization by changing the oxygen ligand ion could give clues in the understanding of the physics of multiferroics. The chalcogenide AgCrS<sub>2</sub> has a thombohedral structure at 300 K [4], in which the Cr<sup>2+</sup> ions form a perfect triangular lattice. Our recent neutron scattering study [5] has shown that a large magneto-elastic effect leads to a crystal symmetry lowering at  $T_{xy}$  which releases the geometric frustration of the magnetic lattice : below  $T_{x} = 41.5$  K, the antiferromagnetically. Drastic changes in the correlations between spino on cooling through TN are observed. Unlike CuCrO<sub>2</sub>, ferroicity in AgCrS<sub>2</sub> is related to the spin-lattice coupling of electric dipoles and magnetic moments, which results in a polar (P ~ 20 uClm<sup>2</sup> at 5 K [6]) structure below  $T_{x}$ .

M. Poienar et al., Phys. Rev. B 79, 014412 (2009), [2] S. Seki et al., Phys. Rev. Letters 101, 067204 (2008),
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 K. Singh et al., Chem Mater 21, 5007 (2009)

#### **JC01**

## STM and magnetotransport investigations on the heavy fermion metals YbRh<sub>2</sub>Si<sub>2</sub> and CeMIn<sub>5</sub> (M = Co, Ir)

Steffen Wirth<sup>1\*</sup>, Andrea Bianchi<sup>2</sup>, Stefan Ernst<sup>1</sup>, Christoph Geibel<sup>1</sup>, Zachary Fisk<sup>2</sup>, Stefan Kirchner<sup>4</sup>, Cornelius Krellner<sup>1</sup>, Joe D Thompson<sup>5</sup> and Frank Steglich<sup>1</sup>
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<sup>2</sup> Departement de Physique, Universite de Montreal, Quebec H3C 3J7, Canada
<sup>3</sup> University of California, Irvine, California 92697, USA
<sup>4</sup> Max Planck Institute for Physics of Complex Systems, Germany

<sup>5</sup> Los Alamos National Laboratory, Los Alamos, New Mexico 87545, USA Heavy fermion metals are often characterized by a variety of relevant energy scales and competine interactions

which may result in such fascinating phenomena as quantum criticality and unconventional superconductivity. Here, we focus on results obtained by Scanning Tunneling Microscopy and Spectroscopy (STM/S) and Hall effect measurements. YDRh<sub>2</sub>Si<sub>2</sub> and its doped counterparts YD(Rh<sub>1</sub>, J), (M = Co, Ir) exhibit a quantum critical point which results from a competition between RKKY and Kondo interaction [1]. The Kondo interaction is visualized by Scanning Tunneling Spectroscopy on YDRh<sub>2</sub>Si<sub>2</sub> [2]. The hybridization of conduction and 4f electrons results in a gap-like feature of the tunneling conductance. Importantly, the crystal field excitations are unambiguously reflected by STS. A strongly temperature dependent peak in tunneling conductance is attributed to a resonance resulting from the Kondo lattice. At even lower temperature, Hall effect measurements reveal a Kondo break-down of the heavy quasiparticles [3]. In the CeMIn, class of compounds the relation between superconductivity and antiferromagnetism will be discussed [4]. Magnetotransport measurements on Celfrln, indicated a precursor state to superconductivity. A gap detected by low-temperature STS in CeColn<sub>3</sub> is compatible with Sd<sub>1</sub> (x<sup>2</sup> - y<sup>2</sup>) S symmetry of the superconducting order parameter and is, again, consistent with a precursor state to superconductivity.

[1] S. Friedemann füt et al.), Nature Phys. fbf 5/ 2009) 465. [2] S. Ernst füt et al.), Nature fbf 474/ 2011) 362. [3] S. Friedemann füt et al.), Proc. Natl. Acad. Sci. USA (bf 107) (2010) 14547. [4] S. Nair füt et al.), Proc. Natl. Acad. Sci. USA (bf 107) (2010) 9537.

#### **JC02**

## Fermi surface of $URu_2Si_2$ in the hidden order state and in the antiferromagnetic state

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The heavy fermion compound URu<sub>2</sub>Si<sub>2</sub> is famous for its hidden order state appearing below T0 = 17.5 K at ambient pressure. Inside this phase superconductivity sets in at 1.5 K. Up to now, the order causing huge signatures in macroscopic quantities at the transition temperature has not been detected experimentally with microscopic probes. The proximity of the hidden order state to a usual antiferromagnetic state appearing under pressure allows a comparison of the two. In this talk a Fermi surface study of URu<sub>2</sub>Si<sub>2</sub> is presented. New Fermi surface sheets have been detected which puts constrains for theory. The pressure dependence shows that the Fermi surface is the same in the hidden order state and in the antiferromagnetic state. As a conclusion, the hidden order must reconstruct the Fermi surface in the same way the antiferromagneties.

#### JC03

## Switching of magnetic ordering near the quantum critical point of the heavy fermion superconductor CeRhIn<sub>5</sub>

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The nuclear quadrupole resonance(NQR) and the magnetic neutron diffraction experiments in CeRhIn<sub>s</sub> clearly indicate that the antiferromagnetic long-range order takes place uniformly over the entire sample where the bulk superconductivity appears. The most intriguing feature in this compound is the switching of magnetic ordering near the quantum critical point, where the spin-density waves (SDW) with different wave vectors are simultaneously observed in the NQR spectrum as well as the neutron diffraction profile. In order to understand the switching mechanism, we have calculated the spin excitation spectrum in the presence of antiferromagnetic and superconducting long-range order using the RPA and found that the concurrence of commensurate and incommensurate SDW is associated with the van-Hove singularity that occurs with the superconducting gap opening.

#### JC04

## Resonant magnetic exciton mode in the heavy-fermion antiferromagnet $CeB_{\boldsymbol{\varepsilon}}$

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Resonant magnetic excitations are widely recognized as hallmarks of unconventional superconductivity. Numerous model calculations have related these modes to the microscopic properties of the pair wave function, but the mechanisms underlying their formation are still debated. Here we report the discovery of a similar resonant mode in the non-superconducting, antiferromagnetically ordered heavy-fermion metal CeB<sub>6</sub>. Unlike conventional magnons, the mode is non-dispersive, and its intensity is sharply concentrated around a wave vector separate from those characterizing the antiferromagnetic order. The mode energy increases continuously below the onset temperature for antiferromagnetism, in parallel to the opening of a nearly isotropic spin gap throughout the Brillouin zone. These attributes bear strong similarity to those of the resonant modes observed in unconventional superconductors below their critical temperatures. This unexpected commonality between the two disparate ground states throws new light on the interplay between antiferromagnetism, superconductivity, and 'hidden' order parameters in correlated-electron materials.

#### JD01

#### Magnetostriction to 97.4T in frustrated Shastry-Sutherland compound SrCu<sub>2</sub>(BO<sub>3</sub>)<sub>2</sub> Marcelo Jaime<sup>1\*</sup>, Ramzi Daou<sup>2</sup>, Scott A Crooker<sup>1</sup>, Franziska Weicker<sup>1</sup>, Atsuko Uchida<sup>1</sup>, Adrian Feiguin<sup>4</sup>, Cristian D

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Strong geometrical frustration in magnets leads to exotic states, such as spin liquids, spin supersolids and complex magnetic textures. SrCu<sub>2</sub>(BO<sub>3</sub>)<sub>2</sub>, a spin-1/2 Heisenberg antiferromagnet in the archetypical Shastry-Sutherland lattice, exhibits a rich spectrum of magnetization plateaus and stripe-like magnetic textures in applied fields [1]. The structure of these plateaus is still highly controversial due to the intrinsic complexity associated with frustration and competing length scales [2,3]. We reveal new magnetic textures in SrCu<sub>2</sub>(BO<sub>3</sub>)<sub>2</sub> via magnetostriction and magnetocaloric measurements in fields up to 97.4 Tesla [4]. In addition to observing the low-field fine structure of the plateaus with unprecedented submicrostrain resolution, the data also reveal lattice responses at 82 T and at 73.6 T which we attribute, using a controlled density matrix renormalization group approach, to the long-predicted 1/2-saturation plateau, and to a new 2/5 plateau. -AF acknowledges NSF funding under grant DMRG-0955707. Experiments at the High Magnetic Field Laboratory Dresden (HLD) were sponsored by Euro-MagNET II under the EU contract 228043. Work at the NHMEL was supported by the National Science Foundation, the US Department of Energy trough the BES "Science at 100T program, and the State of Florida.

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#### JD02

## Unconventional spin-glass behaviors in pyrochlore Heisenberg antiferromagnets coupled with lattice distortions

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Spin glass, in which spins are frozen randomly, is one of low-temperature phases widely observed in geometrically frustrated magnetic materials. However, it is unclear so far how its nature is different from the canonical one driven solely by randomness. Experimentally, several puzzling spin-glass behaviors have been pointed out for pyrochlore-based magnets, such as R<sub>2</sub>Mo<sub>2</sub>O<sub>7</sub> and ACr<sub>2</sub>O<sub>6</sub>, i.e., the transition temperature appears to be almost independent of the strength of randomness, and the value is much higher than the theoretical prediction. Motivated by these puzzles, we investigate the effect of magnetoelastic coupling in a bond-disordered Heisenberg antiferromagnet on a pyrochlore lattice [1,2]. Through classical Monte Carlo simulations using the extended loop algorithm [3,4], we show that the magnetoelastic coupling largely enhances the spin-glass transition temperature is set by the magnetoelastic coupling, and becomes almost independent of bond disorder, the transition temperature is set by the magnetoelastic coupling. Base behaviors are ascribed to spin collinearity induced by the magnetoelastic coupling. We also compute the specific heat and (non)linear susceptibility and compare the results with experiments as well as the canonical spin-glass behaviors. The present theory qualitatively explains the puzzling spin-glass behaviors berved in experiments.

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#### JD03

#### Theory of spin liquids in integer spin pyrochlores

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Rare earth pyrochlores, with a chemical formula  $A_2B_2O_7$ , exhibit many interesting features. Depending on A site rare earth elements, spin ice and magnetically ordered phases are shown in several experiments. Moreover, they have been also focused as possible candidates of U(1) spin liquid. In order to explore such versatile phases, we study the pseudospin-1/2 model, which is quite generic to describe rare earth pyrochlores with integer spins, in the presence of spin-orbit coupling and crystalline electric field. Using a new 'gauge mean field theory', we show the possible ground states, corresponding to several phases listed above.

#### JD04

## Field-induced spin nematic and spin density wave orders in spatially anisotropic frustrated magnets

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Magnetic multipolar order including spin nematic (magnetic quadrupolar) order is one of the current topics in frustrated magnetism. Recently, frustrated spin-1/2 chains with ferromagnetic nearest neighbor coupling J1(-0) and antiferromagnetic next nearest neighbor one J2(-0) have been theoretically shown to exhibit multipolar quasi long range orders in the wide region of J1/J2 as an external magnetic field is applied [1,2]. In addition, it is known that several kinds of quasi one-dimensional (1D) cuprates [3] can be described by this J1-12 spin chain [4,5]. Particularly, recent experiments show that LiCuVQ<sub>6</sub>, one of the quasi-ID cuprates, possesses a new phase near saturation, and it is expected to be a spin nematic long-range ordered phase [6]. Motivated by these results, we have completed the field-temperature phase diagram for spatially anisotropic magnets consisting of weakly coupled J1-J2 spin chains, making use of the accurate results of the single J1-J2 spin chain [1,4]. The phase diagram contains spin mematic and spin-density-wave ordered phases; and these two orders compete with each other. We will discuss some universal features of the phase diagram and the relevance of our result to LiCuVQ<sub>6</sub>.

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#### **JD05**

## Spin liquids for spin 1/2 systems with strong charge fluctuation on the triangular lattice

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We describe a new class of spin liquids for organic spin 1/2 triangular lattice systems. Adopting the Fermionic-spinon representation, d-wave pairings between spinons are considered, which makes U(1) gauge symmetry down to Z2 one. Finite amount of energy gap of the Z2 gauge theory prohibits coupling between the emergent gauge field and external electro-magnetic field at low energy. Thus external magnetic field only couples to the Zeeman term and it is argued that the ground state does not show thermal Hall conductance. Also, the fermionic band structure allows both Dirac-fermi points and a quadratic band touching point. It is shown how the fermionic band structure could resolve various puzzling experiments, for example, metallic thermal conductivity without thermal Hall conductivity.[1,2,3]

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#### JD06

## Emergent criticalities and phase transitions in monomer-dimer mixture system on a honeycomb lattice

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The frustrated magnets such as  $Dy_2Ti_2O_7$  are called as the spin ice, and have received much attention [1]. In particular, a defect representing a breaking the ice rule can be regarded as the magnetic monopole in solids, so its property is a current interest [2]. Here, we study the spin ice in a magnetic field using a toy model, i.e., a monomerdimer mixture defined on a honeycomb lattice. In a low-doping region of monomers, its effective description is given by the dual sine-Gordon model [3]. In intermediate- and strong-doping regions, the Potts lattice gas theory can be employed [4]. Synthesizing these, we construct a renormalization-group flow diagram, which includes the stable and unstable fixed points corresponding to M5 and M6 in the minimal models of the conformal field theory. We perform numerical calculations to determine a phase diagram and also to proffer evidence to check our prediction [5].

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#### JE01

**Domain-wall physics and devices using focused electron and ion beams** Henk Swagten\*, Jeroen Franken, Christian Geurts, Mark Van Der Heijden, Mark Hoeijmakers, Tim Ellis, Elena Mure, Beatriz Barcones, Juergen Kohlhepp and Bert Koopmans

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In domain-wall (DW) devices with perpendicular magnetic anisotropy (PMA) using e.g. ultrathin Co, a number of exciting phenomena have been observed very recently. Ion irradiation is shown to be an effective tool to locally modify PMA properties, with which a kaleidoscope of new DW physics or device directions may be explored. For example, we have used ion irradiation in PMA strips for (1) controlled pinning of DWs (2) predictions of a novel domain-wall oscillator, (3) intrinsic DW resistivity as a function of the width of the DW, and (4) a novel ratchet DW memory device by sawtooth shaping of the DW potential. After shortly introducing these developments, the presentation will focus on the use of electron-beam-induced-deposition (EBID) for growing Co and Fe-rich nanopillars on top of our magnetic strips, with a diameter typically 50-150nm. After nucleation of a DW at an ion-irradiated part of the strip, the DW is moved towards the position of the pillar, at which it may be pinned/depinned using a variable strength of the applied perpendicular field. Apart from controlled (de) pinning due to the altered DW energy landscape induced by the pillar stray field, we could also uniquely extract the switching fields of these nanopillars.

#### **JE02**

#### Real time analysis of spinmotive forces due to domain wall motion Jun'ichi Ieda\*, Yuta Yamane and Sadamichi Maekawa

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Spinmotive force, which is induced by motion of a nonuniform magnetic nanostructure, reflects real time dynamics of the magnetization texture [1,2]. It is wellknown that domain wall motion in a permalloy nanowire exhibits nontrivial magnetic field dependence. In particular, in high field regime far above the Walker breakdown field, the domain wall shows structural deformation including annihilation and creation of magnetic vortex cores [3]. Using numerical simulations we study the spinmotive force in such a circumstance and find that its dc component scales with an applied magnetic field even in a field range where the wall motion is no longer associated with periodic angular rotation of the wall magnetization. This feature has been confirmed in a recent experiment [4]. As the field is increased, spikes in the voltage signals start to appear, which are mainly associated with vortex core nucleation and annihilation, and this tendency is enhanced with further increase in the field. At high fields, the slope of the generated dc voltage with the applied field is expected to be only dependent on the spin polarization of conduction electrons and thus can be used to accurately determine the degree of spin polarization in various materials.

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#### JE03

#### Spin-current induced magnetization dynamics

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Since conventional approaches of manipulating magnetization using magnetic fields exhibit unfavourable scaling and results in limited switching speed, alternative approaches based on spin-currents have emerged. The transfer of angular momentum ("spin transfer torque effect") leads for instance to current-induced domain wall motion (CIDM), which has become the focus of intense research in the last few years. We have comprehensively investigated CIDM and determined the acting adiabatic and non-adiabatic torque terms [1]. We find that the previously neglected diffusive torque term [2] can play an important role for vortex core displacement. For out-of-plane magnetized Co/Pt a large non-adiabatic torque is found. while in Co/Ni the adiabatic torque dominates [3]. By separating charge and spin transport we can generate pure diffusive spin currents with no associated net charge current. Using such diffusive spin currents we find large efficiencies for domain wall displacement due to strong spin accumulation absorption [4]. To increase the spin diffusion length, we use robust turbostratic graphene and find spin injection across transparent contacts. Beyond generating spin currents by spin injection, we have also used the Spin Seebeck Effect, where temperature gradients generate spin currents that then affect for instance domain wall propagation [5].

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#### **JE04**

#### External magnetic field dependence of the magnetic wall drive current density in a TbFeCo magnetic nanowire

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Low current driven domain wall by using amorphous TbFeCo wire has been reported (1), the current density was 5x10<sup>6</sup>A/cm<sup>2</sup>. Most reported magnetic wire samples of domain wall drive by current are prepared by using EB or Photo lithography and etching process. However, this process damages the edge of a magnetic wire and increase resistance for domain wall displacement. To reduce the resistance, we propose a new fabrication method of magnetic wire without any etching process. It is named as nano-imprint magnetic wire method. Using the method, the improved TbFeCo wire was fabricated by magnetron sputtering on to the plastic substrate with a groove shape like a wire pattern. Therefore, relatively smooth roughness nanowire can be made. This new fabrication method of nanowire can be reduced number of pinning site. Then, the current density for domain wall driving can be also decreased to 6x105 A/cm<sup>2</sup>. Moreover, applying an external field to the sample, the current density can be reduced more. Thus, this new fabrication method of nanowire is attractive for nanowire memory device.

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#### **JF01**

#### Disentangling and manipulating intrinsic and extrinsic contributions in the anomalous hall effect

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The anomalous Hall effect is one of the most prominent phenomena existing in magnetic materials. It has remained unsolved for more than a century because its rich phenomenology defies the standard classification methodology, prompting conflicting reports claiming the dominance of various processes. Working with epitaxial films of Fe, Ni, and Ni, Cu1,, we succeeded in independent controls of different scattering processes through temperature and layer thickness, which allows an unambiguous identification and control of the intrinsic mechanism as well as the extrinsic mechanisms of the anomalous Hall effect.

#### **JF02**

#### Spinmotive forces in spin-orbit coupling systems Yuta Yamane\*, Jun'ichi Ieda and Sadamichi Maekawa

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Spinmotive force in ferromagnetic conductors with spin-orbit interaction is theoretically studied. We have derived the expression for the spinmotive force including the contribution from a general form of spin-orbit interaction. We have found that, in addition to the conventional spinmotive force which reflects the time- and spatialdependent magnetization, time variation of the vector product of the magnetization and an electric field gives rise to a spinmotive force. Therefore, a time-varying electric field can induce a spinmotive force even in a static and uniform magnetization configuration. In order to detect this phenomenon, we propose an experimental setup with a metallic film which is thin enough for an electric field to penetrate. In addition considering a two dimensional electron system with the Rashba and Dresselhaus spinorbit couplings, the spinmotive force signals are different for each origin. Therefore this spinmotive force offers a method to determine the ratio of spin-orbit couplings.

#### **JF03**

#### Theory of the spin Hall effect in ferromagnetic metals: Nonlinear behaviors around the curie temperature

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We give a theory of the spin Hall effect in ferromagnetic metals near the Curie temperature. We do this by extending Kondo's theory of the Anomalous Hall Effect in ferromagnetic metals to include the short range spin-spin correlations. We find a novel relation between the Spin Hall Effect and a second order nonlinear spin fluctuation in ferromagnetic metals near Curie temperature Tc. In 1962 Kondo gave a relation between the Anomalous Hall Effect and the first order nonlinear spin fluctuation in pure ferromagnetic metals near Tc. Thus our results show an essential difference between the Anomalous and Spin Hall Effects in terms of the nonlinear spin correlations that are probed. Our theory can be compared to experiments in ferromagnetic alloys near Tc by Y. Otani et al.

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#### **JF04**

#### Minority band gap and magnetic properties of Co<sub>2</sub>(Fe,Mn)Z (Z=Al, Ga ; Si, Ge) in the context of CPP-GMR transport Faleev Sergev and Oleg N. Mrvasov\* Physics and Astronomy and MINT Center, Physics and Astronomy and MINT Center, USA



#### **JF05**

#### Anisotropy in the intrinsic anomalous Hall effect Lin Wu, Yufan Li, Jianli Xu and Xiaofeng Jin\*

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A new scaling for the anomalous Hall effect was recently proposed [Phys.Rev. Lett.103,087206 (2009)], which can disentangle the intrinsic and extrinsic contributions in the anomalous Hall effect. Based on the new scaling, we have investigated in experiment the anisotropic effect of the intrinsic anomalous Hall contribution. Comparing with Fe(001)/GaAs(001), we find the value of intrinsic AHE contribution is different in Fe(111)/GaAs(111), which is directly related to the electronic band structure of Fe

#### **JG01**

#### Propagation and scattering of spin waves in curved magnonic waveguides

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Usually, the modulation of the effective magnetic field in magnonic crystals and devices is achieved as a compositional modulation created either during growth or using lithographic processing. In this report, we develop a continuous medium theory of dispersion and scattering of spin waves propagating in thin cylindrical nanowires (therefore acting as magnonic waveguides) with curved regions. The theory predicts that, assuming that the static magnetization is aligned along the waveguide, the curvature makes a significant contribution to the effective magnetic field in the waveguide, scaling inversely with its radius of curvature squared. We calculate the spectrum and coefficients of reflection and transmission of spin waves that are created due to this topological nonuniformity and discuss problems arising from and opportunities offered by the discovered effect. The research leading to these results has received funding from the EC's 7th Framework Programme (FP7/2007-2013) under GA 247556 (NoWaPhen) and from EPSRC of the UK under project EP/E055087/1.

#### **JG02**

#### Theory of static and dynamic properties of magnetic dot arrays coupled by dipole-dipole interaction

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Arrays of magnetic nano-dots are promising candidates for applications in microwave signal processing and in magnetic recording technology. Since ground state of an isolated dot is double-degenerate, there are many possible ground states of an array. Different ground states, which can be controlled dynamically [1], correspond to different spin-wave (SW) spectra of an array. Thus, using dot arrays it is possible to develop new artificial magnetic materials: dynamic magnonic crystals. In this paper we present a general theory of static and linear dynamic properties of magnetic dots arrays coupled by a magnetodipolar interaction. In contrast with the majority of previous analytical methods the theory presented here is not limited to macrospin approach, as spatially nonuniform static and dynamic magnetization profiles are taken into account by introducing static-static, staticdynamic and dynamic-dynamic effective mutual demagnetization tensors. We show that all the possible ground states of a finite or infinite array with given periodicity can be found as solutions of a system of linear equations. The developed theory provides a simple way to calculate mode structure, damping rates and excitation efficiency of SWs, and microwave response of a magnetic dot array.

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#### **JG03**

#### Theoretical study on ferromagnetic resonance of FePt/Py bilayers

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NRI-AIST Japan

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Twisted magnetic structures such as a domain wall and a magnetic vortex attract much attention for applications in nano spiontornics. It has been widely known that we can make a twisted magnetic structure in a ferromagnetic bilayer consisting of materials with different magnetic anisotropies, which is called an exchange spring bilayer (ESB). We can control its twist angle by applying an external magnetic field. Recently the spin dynamics in ESB has attracted much attention since the microwave assisted switching was proposed by Fal et al. [1]. However the spin dynamics in ESB under oscillating magnetic field is not fully understood. We theoretically studied the ferromagnetic resonance of FePt/Py bilayers of which material properties were well studied experimentally in [2]. We solved the Landau-Lifshitz-Gilbert equation and found that the resonance frequency of the perpendicular standing spin wave [3] is lowered by twisting the spins in ESB. The decrease of the resonance frequency is enhanced when ESB is about to switch We shall show the details of our calculation and comparison with experimental results in our talk.

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#### **JG04**

#### NMR observations of level crossings in a Cr<sub>8</sub>F<sub>8</sub> pivalate single crystal: The solution to the structured enhancement of 1/T1

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Micotti et al. [1] measured 1/T1 of a single crystal of the antiferromagnetic ring [Cr<sub>8</sub>F<sub>8</sub>(piv)16] up to 15.6 T at 0.3 K. Interestingly enough 1/T1 is single- and double-peaked at the first and second level crossings, respectively. A single but twinned crystal would produce such a split peak, but this possibility is ruled out by x-ray measurements. There is no other observation of such a structured peak. 1/T1 observations of decanuclear ferric wheels [2] are obtainable beyond the third crossing field, but each structure, if any, could not be resolved due to the use of a powder sample. Those of hexanuclear [3] and octanuclear [4] ferric wheels are restricted to the first level crossing. Diagonalizing a well-equipped microscopic spin Hamiltonian of each cluster and then inquiring further into intercluster coupling effects, we reproduce the spin dynamics in the Cr<sub>8</sub> single crystal at sufficiently low temperatures. Here is a solution to the wide-split peak of 1/T1 at the second, rather than the first, level crossing,

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#### **JG05**

Spin state of ferric chloride investigated by Fe NMR Byeongki Kang, Changsoo Kim, Euna Jo, Sangil Kwon and Soonchil Lee\* Physics, KAIST, Korea

It was reported that ferric chloride(FeCl<sub>3</sub>) has helical spin structure with period of 15 in the ground state and undergoes two quantum phase transitions at 1.5 T and 4.0 T as magnetic field increases. Between 1.5 - 4 T, the spins of Fe ions are arranged in two directions canted along the axis parallel with magnetic field and above 4.0 T, spins flop. We have measured magnetization curve and magnetic field and temperature dependences of Fe NMR spectrum of polycrystalline FeCl<sub>3</sub>. The temperature dependence of the NMR resonance frequency in zero magnetic field reveals sublattice magnetization. The result fits well to Bloch's T2 law which is expected for a simple antiferromagnet having no anisotropy. The magnetic field dependence of the NMR signal line width shows that the directions of the spins hardly change below 2.5 T. From the field dependence of NMR frequency and line width, we conclude that the spins are almost perpendicular to the field above 5 T. The results are compatible with the previously reported transition at 1.5 T but not with that at 4 T. We estimated the exchange coupling constant from the measured angles between the magnetic field and the spins above 5 T

#### **JG06**

#### Spin dynamics of ferrite nanoparticles studied by 57Fe Mossbauer spectroscopy

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Due to its specific timescale, Mossbauer spectroscopy is highly suitable to investigate the dynamic properties of magnetic nanoparticles. The hyperfine magnetic spectra between the blocking temperature and very low temperatures may exhibit a broad variety of different shapes. Accordingly, to describe this rich behaviour a whole range of different, controversially discussed models can be found in literature [1-4]. We performed measurements on ZnFe2O4 nanoparticles, prepared by a non-aqueous sol gel method and characterized by x-ray scattering, dynamic light scattering and TEM. The spectra were taken on strongly and weakly interacting particles and the fits to the spectra with the different models are compared in order to gain information about the suitability of the applied models.

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#### **JH01**

#### Multi-layered nanocomposite thick film-magnet for power MEMS applications

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An isotropic nanocomposite thick film-magnet is one of candidates for a magnet used in power MEMS because of its high remanence and the easiness of multi-polar magnetization [1]. Previously, we developed a method of synthesizing multi-layered nanocomposite thick film-magnets composed of several hundred layers by using PLD (Pulsed Laser Deposition) method, and reported magnetic properties of Nd-Fe-B/Fe<sub>3</sub>B and Nd-Fe-B/α-Fe multi-layered thick film-magnets [2,3]. In this investigation, we improved magnetic properties of Nd-Fe-B/a-Fe multi-layered thick film-magnets and their reproducibility by controlling preparation conditions of films such as energy density of laser beam and stacking periods of Nd-Fe-B/α-Fe layers, tp. Superior magnetic properties were obtained for the films with tp of 10-20 nm. This tp value is consistent with the results of the previous computer simulation. It is also clarified that decrease in the number and/or size of droplets improve(s) magnetic properties, because the presence of droplets degrades a layered structure and resultantly reduces effective exchange interaction between soft and hard phases. The obtained (BH)max value, 112 kJ/ m3, is the larger than the values obtained in previous studies for isotropic film-magnets thicker than 10 µm, which suggests that multi-layered nanocomposite thick film-magnets are promising for power MEMS applications.

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#### **JH02**

#### Development of high performance micron-scaled hard magnetic structures for micro-system applications

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Micro-systems incorporating magnetic materials have many potential applications in the fields of bio-medicine, information technology, energy transformation/management. As the size of the magnetic element is downscaled, the strength of the magnetic field gradient it produces is up-scaled and values as high as 10 exp 6 T/m can be achieved with micron scaled structures. Hard magnetic materials offer the distinct advantage that they need neither an external magnetic field source nor power supply, and thus they are particularly well suited for use in autonomous, mobile devices or where space is limited. The development of hard magnet based systems has been hindered to-date by the challenges faced in integrating these materials at the appropriate scale using techniques compatible with today's micro-system technologies. In this presentation we will report on the preparation and patterning of high performance hard magnetic materials in film and polymer-powder composite form. We will present their magnetic structural and mechanical properties and will give examples of a range of bio-related applications now being developed using these micron-scaled hard magnetic structures.

#### **JH03**

#### Exchange spring magnet for rare earth free permanent magnet

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Using the full potential linearized augmented plane wave (FLAPW) method, we have investigated the thickness dependent magnetic properties of rare earth free exchange spring magnet FeCo/FePt(001). The FeCo adlayer thickness is increased from one monolayer (ML) to four ML coverage. It is observed that the FeCo adlayers and Fe atoms in FePt substrate show almost half metallic behavior, while an ordinary metallic feature is found in Pt atoms. The average magnetization increases with FeCo thickness and the estimated energy product reaches 65.61 MGOe in FeCo(4ML)/FePt(001). A giant perpendicular magnetocrystalline anisotropy (MCA) energy is found in pure FePt(001) and the MCA energy in FeCo/FePt(001) systems also show quite large values. For instance, the calculated MCA energy of pure FePt(001) is about 18.2 meV/ cell, and it is about 17.35 meV/cell for FeCo(4ML)/FePt(001). Both MCA energy and estimated energy product values of FeCo/FePt(001) exchange spring magnet structure are comparable with those found in rare earth substituted permanent magnet. This may imply that the FeC/FePt can be utilized for potential rare earth free exchange spring magnet material This work was supported by the KCAP located in Sogang University funded by MEST through (NRF-2011-C1AAA001-2011-0030278), and by KOSEF grant founded by MEST (No.R01-2008-000-20014-0).

#### **JH04**

#### Prediction of maximum energy product for exchange coupled coreshell nanomagnets

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Ferromagnetic τ-phase MnAl possesses a high magnetic moment of 2.4 µB per unit cell [1] and the Curie temperature (Tc) of 655 K [2], and is also rare-earth free. However, its maximum energy product, (BH)max, is about 12.3 MGOe at 300 K. This maximum energy product is not enough for high-energy magnetic device applications. Therefore, we have proposed core (7-phase MnAl hard magnet)/shell (soft magnet) nanostructure to enhance the (BH)max, which is higher than that of Sm<sub>2</sub>Co<sub>17</sub>. We have modified the Skomski's approach [3], that is exchange hardening in nanoscale combination of a soft phase and an oriented hard phase, to predict the theoretical limit of (BH)max for the τ-phase MnAl core-shell nanomagnet. It was found that (BH)max of core (70 nm τ-phase MnAl hard magnet)/shell (18 nm thick 1.6 T soft magnet) nanomagnet exceeds 40 MGOe at 300 K, where the soft shell thickness is thinner than two times of MnAl domain wall thickness. It is noted that the magnetization of soft shell plays a key role in enhancement of (BH)max. Detailed calculations for (BH)max will be discussed for various hard-core sizes and soft-shell thickness.

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#### JI01

#### Ferromagnetic properties of Co-Pd-SrTiO<sub>3</sub> alloy films with high magnetic anisotropy

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Developments of electromagnetic devices require higher frequency soft-magnetic materials with high magnetic anisotropy (Hk) and electrical resistivity (p). Early work has found Co-based nanocomposite films show good high-frequency magnetic properties[1]. However, it's difficult to obtain both high Hk and o simultaneously. In this work, Pd is induced to form Co-Pd alloy nanoparticles to enlarge Hk and SrTiO<sub>3</sub> (STO) is employed as nonmagnetic structural phase due to its high o. The effect of Pd on structure and magnetic properties has been investigated. The Co-Pd-STO alloy films were deposited onto Si and quartz substrates by reactive magnetron sputtering, using a composite target composed of a STO disk and Co, Pd chips. Film structures were investigated by XRD. The chemical composition of the films was analyzed by EDS and EPMA. The magnetization was measured with VSM. The Co-Pd-STO film structures and magnetic properties are strongly affected by Pd content. At the optimal composition, the films exhibit Bs of about 7kG,  $\rho$  in order of 103  $\mu\Omega$ cm, and Hk of around 150 Oe. The film shows good frequency response of permeability at GHz range. Co-Pd-STO alloy film is a good candidate for the GHz range electromagnetic devices.

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#### .1102

#### Effect of change in thickness on the structural and magnetic properties of L10-ordered FePd films with (001) texture

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L10-ordered films with strong perpendicular magnetic anisotropy (PMA), such as CoPt, FePt and FePd, have received considerable attention since they are a superior candidate that can overcome superparamagnetism as they become smaller [1, 2]. To form L10-ordered FePd on an MgO (001) substrate, we prepared samples having a repetition of a bi-layer composed both Fe and Pd monolayers by a dc ultra-high vacuum magnetron sputtering system at room temperature, and annealed them under a perpendicular magnetic-field of 4 kOe. The total thickness of [Fe/Pd]n was changed from 3 to 20 nm. As-deposited FePd had a (100) textured FCC disordered phase and showed in-plane magnetic anisotropy. Interestingly, after annealing at 500 °C for 2 h, FePd was transformed to a (001) textured FCT L10-ordered phase and showed PMA. Especially, 3- and 5-nm-thick FePd showed strong PMA with a squareness of unity. With increasing of thickness of FePd, the FCC (100) texture was gradually increased, resulted in a decrease in PMA. In fact, the 20-nm-thick L10-FePd showed mixed anisotropy with both in-plane and out-of-plane magnetic anisotropy components. This research was supported in part by the IT R&D program of MKE/KEIT. (KI002189, Technology Development of 30 nm level High Density Perpendicular STT-MRAM)

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#### JI03

#### Mechanism of large magnetic anisotropy of thin film m-DO19 Fe<sub>3</sub>Pt and analogous 3d-5d compounds

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#### JI04

#### Fabrication of highly ordered L10 type FePt thin films by rapid thermal annealing

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Large uniaxial magnetic anisotropy materials are extremely promising for the application to high-density magnetic storage devices or permanent magnets. In particular, an L10 type FePt alloy reveals extremely large magnetic anisotropy. therefore numerous studies on the fabrication of FePt films have been reported However, it is a fatal problem that a high growth-temperature or a high postannealingtemperature is needed to obtain an ordered FePt phase. In this paper, we employed a rapid thermal annealing to promote an L10 ordering of sputtered FePt films even at a relatively low temperature, and investigated their crystallographic and magnetic properties. Fe and Pt were deposited simultaneously on thermally oxidized Si substrates. The thickness of FeNi films was 6.5 nm. After the deposition films were rapidly annealed by a temperature elevation with a rate of 10°C/sec in high vacuum. It was found that a magnetic anisotropy increased by annealing, and a uniaxial magnetic anisotropy energy (Ku) was estimated to be 3.5×107 erg/cc from magnetization curves. Ku and coercivity drastically increased with higher rates up to 50°C/sec. On the other hand, it was revealed that a chemical order parameter estimated from X-ray diffraction pattern slightly decreased with the rate. Detailed structural analysis will be also discussed

#### JI05

#### Competing intrinsic and side-jump anomalous Hall effects in Isoelectric L10 FePtPd ternary alloy films

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We have for the first time studied the anomalous Hall effect for isoelectric L10 Fe<sub>0.5</sub>(Pt<sub>x</sub>Pd<sub>1-x</sub>)<sub>0.5</sub> ternary alloy films. When the spin-orbital coupling strength is feasibly controlled changing the concentration ratio of Pd to Pt atoms, anomalous Hall conductivity (AHC) can be continuously tuned from 300 S/cm for x = 0 (L10 FePd) to 930 S/cm for x = 1 (L10 FePt). Calculations have shown the intrinsic AHC increases significantly whereas the side-jump AHC decreases weakly. In particular, the present results of the side-jump AHC for L10 FePt and FePd films reproduces the calculated results based on scattering-independent model. Our results will generate new theoretical interest in the L10-ordered ternary alloys.

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#### JI06

#### Effect of deposition temperature on the crystallographic structure and first-order magnetic phase transition of FeRh thin films on glass substrate

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Ordered B2(CsCl)-type FeRh alloy and its nearby compositions have been subject of extensive theoretical and experimental studies due to a first-order magnetostructural phase transition from an antiferromagnetic (AFM) to ferromagnetic (FM) state when heating beyond a critical transition temperature. Recently, there are many literatures investigating the first-order AFM/ FM transition in FeRh-based thin films. Generally, deposition temperature is expected to influence the structure and consequently the magnetic properties of magnetic thin films. In this presentation, the effect of deposition temperature on the structure and AFM/FM phase transition is investigated in details. XRD results show that the films fabricated at Ts = 350 and 400 oC are grown with FeRh(001) planes parallel to the MgO(100) substrate. The temperature dependent magnetization curves indicate that the as-deposited thin film at Ts = 3500 did not show any AFM/FM phase transition although it has ordered B<sub>2</sub> structure as shown in XRD results. While for the FeRh film deposited at 400oC, it shows a clear AFM/FM phase transition when heating from low temperature to high temperature. But the transition is very broad (range from -50oC to about 100oC) due to the composition variation, internal stress, defects and so on in the thin films

#### JJ01

## Some new physics and magnetism of rare earth-rich $R_5T_4$ and $R_5T_3$ compounds

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Intermetallic compounds of the rare earth metals (R) with group 14 elements (T) at the R<sub>3</sub>T<sub>4</sub> and R<sub>3</sub>T<sub>3</sub> stoichiometries provide numerous opportunities to clarify structure-property relationships that are generally elusive in compounds formed by metals, and in the future, to exploit this knowledge. The uniqueness of R<sub>3</sub>T<sub>4</sub> compounds lies in well-defined, self-assembled layers composed of R and T atoms coupled with the flexibility to modify their arrangements in closely related structures using a variety of triggers. Interestingly, the R<sub>3</sub>T<sub>3</sub> compounds that also adopt distinctly layered but different crystal structures exhibit an interesting interplay between their crystal and magnetic sublattices. In this presentation we will be concerned with some recently discovered, extraordinarily interesting structural, magnetic, and electronic transport phenomena that are related to structural and microsctructural modifications that facilitate an unprecedented level of control over the physical behaviors of these compounds. This work is supported by the U.S. Department of Energy, Materials Sciences and Engineering Division of the Office of Basic Energy Sciences under contract No. DE-AC02-07CH11358 with lowa State University.

#### **JJ02**

## Morphotropic phase boundary in ferromagnets - A way leading to large magnetostriction

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Morphotropic phase boundary (MPB), a phase boundary separating two ferroelectric phases of different crystallographic symmetries in the composition-temperature phase diagram, is crucially important in ferroelectric materials, because MPB can lead to a great enhancement of piezoelectricity, the most useful property of this large class of functional materials. From the parallel physics between ferroelectricity and ferromagnetism, it is tempting to ask if similar MPB can exist in another large class of functional materials, and if it can lead to significant enhancement of the corresponding properties (magnetostrictive effect, or magnetic-field-induced distortion). Here we report the existence of a MPB in a ferromagnetic system TbCo<sub>2</sub>-DyCo<sub>2</sub> between a ferromagnetic thombohedral phase and a ferromagnetic tetragonal phase. Such a magnetic MPB involves a first-order magnetostrictive response compared with that of the off-MPB compositions. In-situ observation by high-resolution synchrotron x-ray diffractometry (XRD) reveals that the MPB region is composed of co-existing rhombohedral and tetragonal phases, indicating the bi-stability of both magnetization vector and lattice distortion. This accounts for the maximum magnetostrictive roperty at the MPB.

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#### JJ03

#### **Magnetostriction in geometrically frustrated Co<sub>3</sub>V<sub>2</sub>O<sub>8</sub> single crystals Ryszard Zuberek<sup>1</sup>**, Ritta Szymczak<sup>1</sup>, Jan Fink- Finowicki<sup>1</sup>, Victor Nizhankovskii<sup>2</sup> and Henryk Szymczak<sup>1</sup>

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In our previous papers (Phys.Rev.B,2006, 2009) we have shown that frustrated Co<sub>3</sub>V<sub>2</sub>O<sub>8</sub> single crystals are characterized by a strong magnetocrystalline anisotropy and that this anisotropy arises due to presence of Co<sup>2+</sup> ions strongly connected to the lattice. In order to check this single-ion mechanism we decided to perform magnetostriction measurements in Co3V2O8 single crystals. The high-resolution capacitive dilatometry was used to determine linear (longitudinal and transverse) and volume magnetostriction as a function of temperature in low temperature region and in magnetic field up to 12 T directed along the main crystallographic directions. The single crystals of Co<sub>3</sub>V<sub>2</sub>O<sub>8</sub> were grown by floating zone technique using an optical image furnace. The results of measurements are discussed in terms of the singleion model taking into account the competing first and second neighbor Heisenberg exchange interactions and magnetocrystalline anisotropy constants determined in our previous studies. The satisfactory agreement with experimental results shows that the magnetostriction in geometrically frustrated Co<sub>3</sub>V<sub>2</sub>O<sub>8</sub> is dominated by the single-ion mechanism. The paper was partly supported by the Ministry of Science and Higher Education (Poland) through grant No. N202 125135 .

#### JJ04

#### Magneto-volume anomalies and low-temperature inverse magnetocaloric effect in Er<sub>2</sub>Fe<sub>1</sub>.

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Most of the magnetic properties of  $Er_{z}Fe_{\tau\tau}$  intermetallic can be understood using a two-sublattice model, corresponding to the erbium and iron elements, respectively. This compound crystallizes in the hexagonal Th\_Ni<sub>x</sub>-type structure (PG5/ mmc space group) and exhibits ferrimagnetic behaviour with the Er sublattice moments antiparallel to those of the Fesublattice, and a Curie temperature TC-300C. Combining different experimental techniques, such as neutron powder thermo-diffraction, x-ray powder diffraction under high hydrostatic pressure (20GPa) and magnetization measurements as a function of temperature (2-350K), magnetic field (up to 80Kc) and hydrostatic pressure up to 1GPa), we show the presence of storn gmagneto-volume anomalies, including Invar effect, from low temperature up to 450 K, indicating the existence of short-range magnetic correlations well above TC. In addition, this compound exhibits a direct magneto-caloric effect (MCE) around room temperature as well as an inverse MCE below 100K. A mean field Hamiltonian, that incorporates both crystalline electric field and exchange interactions, has been used to describe the main experimental trends for the temperature dependences of the MCE and the atomic magnetic moments. Therefore, it seems that a correlation between the magneto-volume anomalies and the MCE is the key to understand the physical properties of this intermetallic compound.

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#### JJ05

## Neutron diffraction study of rare-earth compound $\mathrm{Ho}_{5}\mathrm{Pd}_{2}$ with large magnetocaloric effect

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In recent years, more efficient magnetic refrigerant materials are desired due to potential application

In energy-efficient environment-friendly refrigeration technology. Samanta et al. discovered that the rare-earth compound Ho<sub>3</sub>Pd<sub>2</sub> shows large magnetocaloric effect (MCE) [1]. The interesting point is that the large MCE might be originated from field-induced metamagnetic transition below antiferromagnetic transition temperature TN of 28 K. However, it is not easy to judge whether Ho<sub>3</sub>Pd<sub>2</sub> is antiferromagnetic or not only from magnetization data. In order to determine its magnetic structure, neutron diffraction is most powerful technique. We have performed neutron powder diffraction experiments at JRR-3, Japan using polycrystalline Ho<sub>3</sub>Pd<sub>2</sub>. Ho<sub>3</sub>Pd<sub>4</sub> was prepared by arcmelting and annealed at 850 degrees for 72 hours in a vacuum. In this paper, we will present powder neutron diffraction patterns as a function of temperature. The magnetic peaks with a progagation vector  $k = (\delta, \delta, \delta)$  ( $\delta = 0.18$ ) are gradually developing below about 100 K. The peak width is significantly broader than an experimental resolution even temperature is 5 K. These results indicate that Ho<sub>2</sub>Pd<sub>2</sub> is a short-range ordered antiferromagnet with a long wave length.

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#### KA01

## Probing the Kondo effect on the atomic scale by mapping the itinerant electrons

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Scanning tunneling spectroscopy (STS) has provided an approach to study the Kondo effect - one of the oldest many particle phenomena known in condensed matter physics - in real space [1]. In the present work we follow a novel route to investigate Kondo bulk impurities. It has been shown recently [2] that the anisotropy of the copper Fermi surface leads to a strongly directional propagation of quasi particles called electron focusing which gives access to individual bulk impurities in a metal. Using low temperature scanning tunnelling spectroscopy (STS) we have studied the energy-dependent scattering characteristics for single isolated atoms of Co and Fe buried under a Cu(100) surface. For both magnetic impurity atoms we observe a long range Kondo signature which is periodic with the distance to the impurity [3]. The comparison of Co and Fe atoms demonstrates that both impurity species show similar behavior on completely different energy scales, which is determined by the Kondo temperature. We investigate the scattering amplitude as well as the phase. A theoretical interpretation based on a combined approach of band structure and many-body numerical renormalization group calculations is able to describe the rich spatially and spectroscopically resolved experimental data.

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#### KA02

#### Coexistence of antiferromagnetic order and hybridization gap in Cehased kondo semiconductors

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The mechanism of the gap formation in so called Kondo semiconductors (KSs) such as SmB<sub>w</sub> YbB<sub>12</sub>, and CeRhAs remains elusive. The anisotropic gap structure is thought to originate from the k-dependence of hybridization of a crystal field ground state of the rare-earth ion with a conduction band. The ground states of all the KSs are paramagnetic except for CoS<sub>x</sub>Sb<sub>12</sub> which undergoes an SDW transition at 0.9 K. Recently, the compounds CeT<sub>2</sub>Al<sub>10</sub> (T = Ru, Os) have been found to have transport gaps of approximately 50 K and order antiferromagnetically at exceptionally high temperature, 28 K [1-5]. It is puzzling that the ordering temperature is much higher than those of the Gd counterparts in spite of the small ordered moment of 0.3-0.4 B/Ce. To answer this puzzle, we have performed bulk and microscopic measurements on single crystalline samples of CeT<sub>2</sub>Al<sub>10</sub> (T = Fe, Ru, Os) and alloys. We will present the results of magnetic, transport, and optical conductivity measurements [6] as well as neutron scattering. SR, and XAS experiments. The relation between the anisotropic charge/spin gaps, the wave functions of the crystal field ground state, and anisotropic electronic states will be discussed.

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#### KA03

#### Kondo scattering investigated by Nernst-effect measurements Peijie Sun<sup>1</sup>. Christoph Geibel<sup>2</sup> and Frank Steglich<sup>2</sup>

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The transverse thermoelectric response, known as Nernst effect, was investigated for (CeLa)Cu<sub>2</sub>Si<sub>2</sub> with varying Ce concentrations from 2 K to room temperature. Except for the non-magnetic LaCu2Si2, all compounds exhibit largely enhanced Nernst signal in the whole temperature range investigated. Within the standard Boltzmann transport theory, the Nernst effect measures the relaxation-time spectra of charge carriers as a function of energy at the Fermi level. This, however, is hardly accessible by the longitudinal thermoelectric effect, which is mostly dominated by the asymmetric feature of the electronic density of states. With these considerations in mind, we hope to gain new insights into the charge-carrier relaxation related to Kondo scattering by studying the Nernst effect, as well as its correlation with the enhanced thermoelectricity generally observed in heavy-fermion systems. Our observations of large Nernst signals in (CeLa)Cu2Si2 point to a remarkable asymmetric character of the chargecarrier relaxation time attributable to the Kondo effect. Furthermore, our analysis shows that the asymmetric relaxation can account for a major part of the enhanced thermopower observed in the current systems over a large temperature range, leading to the conclusion that the asymmetric density of states is of only minor importance in this respect for Kondo systems

#### KA04

#### Electron spin resonance in antiferro-quadrupolar ordered CeB<sub>6</sub> Pedro Schlottmann\*

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 $CeB_6$  is the first heavy fermion compound without ferromagnetic short-range correlations displaying an ESR signal [1]. The ferromagnetic correlations among Ce or Yb spins in all other compounds showing an ESR signal narrow the resonance width rendering it observable [2].  $CeB_6$  is a cubic compound with a Gamma\_8 ground-quartet. The orbital content of the quartet gives rise to an antiferro-quadrupolar ordered phase below 4 K. Single ions with a Gamma\_8 ground-multiplet are expected to display four transitions, however, only one was observed. Two fundamental questions arise: (1) why is only one transition seen, and (2) why is this transition observed if the Kondo temperature is larger than the linewidth and there are no ferromagnetic correlations between the ions? While for other Ce and Yb compounds with ESR-signal it is not possible to distinguish if the resonance is due to localized spins or conducting heavy electron spins, an itinerant picture within the antiferro-quadrupolar state is necessary for CeB<sub>6</sub>. Work supported by the Department of Energy under grant DE-FG02-98ER45707.

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#### KB01

#### Aligning and measuring the magnetic easy axis direction of superparamagnetic nanoparticles at temperatures much greater than the blocking temperature

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Cobalt-doped magnetite nanoparticles with diameter ~8nm and blocking temperature ~ 60K were fabricated using the protein ferritin as a template. Magnetic measurements of a solution of the nanoparticle-containing ferritin (magnetoferritin) in 50 mM Tris-HCl reveal unusual dynamic behaviour. The magnetoferritin is free to rotate at temperatures down to ~ 210K, although differential scanning calorimetry shows that the bulk of the solution starts to freeze at ~269K. Cooling the sample through 210K in a 5T applied field aligns the easy axes of the magnetoferritin, even though 210K is considerably higher than the blocking temperature. At low temperatures (5K) the magnetoferritin with aligned easy axes has significantly higher remanence and coercivity than when the easy axes are randomly oriented. For temperatures 60<<T<210K, the aligned magnetoferritin also has a much higher susceptibility than when the easy axes are randomly oriented. Warming the sample through 210K, the loss of easy axis alignment is observed as a decrease in the susceptibility from the aligned towards the randomly oriented value. The time constant for this decrease depends on the effective viscosity, and is extremely sensitive to temperature. This work has received funding from the European Community's Seventh Framework Programme (FP7/2007-2013) under Grant Agreement n°228673 (MAGNONICS).

#### KB02

#### Exchange-bias in iron-based nanoparticles

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Iron-based magnetic alloys are much investigated for their good magnetic properties, especially their high saturation magnetisation. However, large magnetocrystalline anisotropy is necessary to obtain a support for magnetic recording. This can be achieved in particular for FePt nanoparticles with the L10 structure [1,2]. In this work, we investigated the magnetic properties of FePt, FeAu and FeW nanoparticles embedded in a metallic matrix (Cr or W). The samples were elaborated by inert-gas condensation. The structural properties have been investigated by high resolution transmission electron microscopy (HRTEM) and laser assisted atom probe tomography (APT) [3]. Their magnetic properties were characterised by SQUID magnetometry and Mossbauer spectrometry. SQUID measurements revealed in some cases an exchange-bias effect which could be due to the occurrence of a core-shell structure in the nanoparticles or to the presence of RKKY interactions through the metallic matrix. This point will be discussed in relation with the nanostructure characterised by HRTEM and APT.

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#### KB03

## Photocontrolled magnetism through interface strain in core-shell prussian blue analogues

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Cubic heterostructured (BA) particles of Prussian blue analogues, composed of a shell of ferromagnetic KjNik[Cr(CN),]]  $\cdot$  nH<sub>2</sub>O (A), Tc ~ 70 K, surrounding a bulk core of photoactive ferrimagnetic RbaCob[Fe(CN),]  $\cdot$  nH<sub>2</sub>O (B), Tc ~ 20 K, have been studied. Below Tc ~ 70 K, these samples exhibit a persistent photoinduced decrease in low-field magnetization, and these results resemble data from other core-shell particles [1] and analogous ABA heterostructured films [2]. This net decrease suggests that the photoinduced lattice expansion in the B layer generates a strain-induced decrease in the magnetization of the A layer, similar to a pressure-induced decrease observed by others in a pure A-like material [3] and by us in our BA cubes. Photocontrol at higher temperatures is being pursued by applying this strain from the B material to other hexacyanochromate analogues, which possess higher ordering temperatures for some stoichiometries. Work supported by NSF DMR-1005581 (DRT) and the NHMFL.

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#### **KB04**

Confinement effect on the A1 to L10 phase transformation of FePt Andrew Gallagher\*, Levent Colak, Ozan Akdogan and George Hadjipanayis Physics and Astronomy, University of Delaware, USA

The major challenge for the application of chemically synthesized FePt nanoparticles (NPs) in magnetic storage media is the sintering problem encountered during the required high temperature annealing to obtain the high anisotropy L10 phase. In this work, we have used two different methods to avoid sintering; coating the NPs with a protective layer of silica (SiO<sub>2</sub>) and using porous aluminum oxide (Al<sub>2</sub>O<sub>3</sub>) as a template to embed the NPs. The NPs were synthesized via the synthesis method of Sun et al.[1] The NPs were embedded into the Al<sub>2</sub>O<sub>3</sub> by in-situ suctioning of the reaction solution into the porous Al2O3 template followed by annealing at high temperatures. Monodispersed FePt NPs with a size of 5.8 and 15 nm were coated with SiO<sub>2</sub> shells using a water-in-oil micro emulsion method. High room temperature coercivities were only obtained after annealing the samples at 900°C for long times (24-48 h) under forming gas flow as compared to the usual heat treatment at 600-700°C. Values of 4.7 and 7.8 kOe were observed in SiO2 and Al2O3 samples, respectively after annealing for 24 h at 900°C. This behavior suggests that the restricted geometry of the samples suppresses the atomic ordering and the phase transformation drastically.

Work supported by DOE DE-FG02-04ER4612 [1] S. Sun, C. B. Murray, D. Weller, L. Folks, A. Moser Science 2000, 287, 1989

#### **KB05**

Are small CoPt and FePt nanoparticles mono-L10 domain?

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Despite the many experimental and theoretical studies, open questions regarding the structure of small CoPt and FePt nanoparticles remain (threshold size for L10 ordering, preferential segregation, most favourable geometry). Moreover, while well known planar defects (c-domain boundaries, antiphase boundaries) are met in macroscopic materials or thin films, their relevance for particles of a few nanometers has not been examined. Therefore, an implicit belief is that small CoPt or FePt nanoparticles, if they can be chemically ordered, would be mono-L10 domain. CoPt and FePt nanoparticles, with a diameter between 2 and 5 nm, have been prepared by cluster deposition and annealed. For both alloys, we observe using transmission electron microscopy a coexistence of fcc nanocrystals and multiply-twinned particles with decahedral or icosahedral shapes. In addition to mono-L10 domain particles, we put into evidence that even small particles can display several L10 domains. In particular, the chemical order can be preserved among twinning which can give rise to spectacular chemically ordered decahedral particles made of five L10 domains. The stability of such structures, which had been recently predicted from theoretical simulations, is thus experimentally confirmed. These observations are then confronted to recet experimental results of magnetometry and synchrotron measurements (XMCD, EXAFS).

#### **KB06**

#### Magneto-structural correlations in antiferromagnetic and ferrimagnetic nanoparticles Nuno Silva<sup>1</sup>\*, Vitor S. Amaral<sup>1</sup>, Luis D. Carlos<sup>1</sup>, Ainhoa Urtizberea<sup>2</sup>, Rodney Bustamante<sup>2</sup>, Angel Millan<sup>2</sup>, Fernando Palacio<sup>2</sup>, Erik Kampert<sup>2</sup>,

Uli Zeitler<sup>3</sup>, Sophie De Brion<sup>4</sup>, Yuki Komorida<sup>5</sup>, Masaki Mito<sup>5</sup>, Oscar Iglesias<sup>6</sup>, Amilcar Labarta<sup>6</sup>, Ines Puente Orench<sup>7</sup> and Javier Campo<sup>7</sup> <sup>1</sup> Dep. Fisica and CICECO, Universidade de Aveiro, Portugal

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We present structural and magnetic studies performed in ferrimagnetic[1,2] and antiferromagnetic[3-7] nanoparticles. We focus our attention on the antiferromagnetic ferritin and cobalt oxide and on the ferrimagnetic maghemite. High-field studies in ferritin reveal that the temperature dependence of the antiferromagnetic susceptibility decreases with temperature when correctly determined from high-field data [5] High-field hysteresis loops also show that coercivity and loop-shifts are a function of the maximum field, which can be understood in the frame of coherent rotation models considering a field-dependent anisotropy energy.[7] Neutron diffraction is particularly effective in the study of structural and magnetic coherence sizes of antiferromagnetic nanoparticles. In the case of cobalt oxide, we found that the difference between both sizes is consistent with a core-shell model where the shell is structurally ordered but magnetically disordered with a thickness of ~2 nm.[6] This core shell model was also successfully applied to a series of maghemite nanoparticles, based on small-angle X-ray scattering and magnetization results, yielding, in this case, a magnetically disordered shell with about 1 nm.[1] Surface and core contribution to anisotropy energy in maghemite can be elegantly separated by using pressure, since the core/shell volume ratio changes with pressure inducing a change in the average anisotropy energy.[2]

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#### **KC01**

#### 360 degree domain walls in various magnetic ring thin films Chunghee Nam1 and Caroline A. Ross2

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Magnetic domain wall (DW), which is an interface between differently oriented magnetization regions, becomes an important physical property as the size of magnetic materials is reduced. DW can be regarded as a rigid particle, which has been verified in experiments and theoretical study, so that diverse applications have been introduced. In this presentation, the authors report that 360 degree magnetic domain walls(360DW) are formed in a circular or ellipse Co rings. In theoretically, 360DWs have been recently studied in terms of the formation and dynamic as a function of applied magnetic fields and electrical currents.[1,2] 360DWs are consisted of two transverse DWs with opposite polarities but the same chirality. Thus 360DWs are locally in a flux-closure state at remanence, which reduce the magnetostatic energy. Magnetic force microscopy(MFM) is used to see the formation of 360DWs. Based on the MFM measurements, the formation of 360DWs gives rise to monitoring the vortex chirality in various ring-structured thin films. Furthermore, 360DWs generate a low stray field, which is useful for high density magnetic devices without proximity effects. The interaction between 360DWs and a tip of MFM will be shown to verify properties of 360DWs in detail.

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#### KC02

#### Magneto-optical effect of rare earth doped zinc ferrite thin films prepared using PLD

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Rare earth elements have large orbital interactions. In addition, electron transition of trivalence rare earth ions exists in the ultra violet regions, and suppresses the absorption in the visible light region. Therefore, large magneto-optical effect is expected for ZnFe<sub>2</sub>O<sub>4</sub> thin films doped with rare earth elements even in the visible region. In this work, ZnFe2O4 thin films doped with rare earth elements (La, Nd, Sm, Eu, Gd, Dy, Ho, Yb) in the Fe site were prepared on fused SiO<sub>2</sub> substrate using pulsed laser deposition. All films were polycrystalline. Ho doped ZnFe2O4 thin film showed relatively high optical transmittance in the visible light region (400-800 nm). Faraday rotation angle changes with the kind of rare earth elements, and maximum angle (12000 deg/cm) was obtained for Ho doped ZnFe2O4 thin film. This results reflects the large spin orbital interactions

#### KC03

Electric field control of coercivity of Pt / Co / Al-O trilayer structures Tatsuro Ohashi<sup>1</sup>, Junichi Shiogai<sup>1</sup>, Tim Yang<sup>1</sup>, Makoto Kohda<sup>1</sup>\*, Takeshi Seki<sup>2</sup>, Kesami Saito<sup>2</sup> Koki Takanashi<sup>2</sup> and Junsaku Nitta <sup>1</sup> Department of Materials Science, Tohoku university, Japan <sup>2</sup> Department of Materials Science, IMR, Tohoku university, Japan

In recent years, much attention has been focused on studies of effective magnetization switching by electric current1-4 because of lower power consumption and zero-magnetic field operation. Due to its compatibility with conventional technology, the electric field control of magnetic properties is more attractive 5-8. Here we demonstrated the electrical gate control of coercivity in an ultrathin Co layer sandwiched between Pt and Al-O. A single wafer which consists of Pt (3 nm) / Co (0.6 nm) / Al-O (1.0 nm) trilayer on SiO, substrate, was prepared. In order to investigate magnetic properties by anomalous Hall effect (AHE), the wafer was processed into a Hall bar geometry. Subsequently, a 20nm-thick Al-O insulating layer was formed by atomic layer deposition at 200 °C. Finally, an AuGeNi / Au gate electrode was evaporated on top. The Curie temperature of our sample was determined to be 340 K at VG = 0 V from the temperature dependence of AHE. The coercivity change of  $\pm 3.0$  Oe was observed at 280K by applying the gate voltages of  $\pm$  8 V, which correspond to electric fields of  $\pm$  3.83 MVcm-1. We successfully modulated the coercivity of the Pt / Co layer by the electric field application.

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#### KC04

#### Microstructure and magnetic property of epitaxial Fe/MgO layer on GaAs and InAs (001) substrates

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Epitaxial growths of Fe/MgO layers on GaAs and InAs heterostructures have been extensively studied for efficient spin injection from ferromagnetic metals into semiconductors. Structural property related to magnetic anisotropy of the Fe layers is mainly determined by underlying strain of MgO tunnel barrier due to the large lattice mismatch between MgO and substrates. In this presentation, we investigated the microstructural evolution and the effect on in-plane magnetic anisotropy of epitaxial Fe/MgO lavers grown on GaAs and InAs(001) substrates with respect to the growth temperature and thickness of the MgO tunnel barrier. The Fe layer grew in three-dimensional(3D) islands with two different in-plane textures of [010] and [110] on the MgO layers grown at low temperatures while it grew in two-dimensional(2D) layer on the MgO layers grown at high temperatures. As the MgO growth temperature and the thickness increase, the strain of the MgO and the subsequent Fe lavers are simultaneously relaxed. The partial strain relaxation of the Fe leads to the transition from 3D islands of no magnetic anisotropy to 2D layer of in-plane cubic anisotropy. Furthermore, the islands coalesce as the nominal thickness of Fe increases, resulting in the continuous 2D layer of Fe with the cubic anisotropy.

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#### KC05

#### Magnetic and transport properties of epitaxial discontinuous Fe/MgO multilavers

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Discontinuous metal-insulator multilayers (DMIMs) are cermet structures prepared by repeated alternate deposition of a continuous insulating and a discontinuous metal layer [1]. However, little attention has been paid to the influence of the substrate nature, deposition temperature, and degree of crystallinity of the insulator layer on transport properties of DMIMs. In recent work [2] we reported tunnelling magnetoresistance (TMR) values for epitaxial Fe/MgO DMIMs deposited on single-crystalline MgO(001) substrates. Here we report a comparison of the magnetic and transport properties of DMIMs [Fe (0.6 nm)/MoO (3 nm)]10 denosited on two different substrates: Coming Glass (CG) and single-crystalline MoO(001) with deposition temperatures (Ts) ranging from 20 °C to 250 °C. Zero field cooled and field cooled magnetic susceptibility measurements were carried for H=50 Oe and 5K<T<300K. The samples exhibit superparamagnetic behaviour. For temperatures higher than the bifurcation temperature (TB) the susceptibility measurements were fitted using a Curie-Weiss law and the average size (Davg) of the magnetic grains was obtained. More precise determination of Davg, and their distribution was done by measuring magnetization isotherms for T>TB. We observed that the Fe grains volume is bigger for the samples prepared on MgO substrates. The relationship between transport and magnetic properties will be discussed.

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#### **KC06**

#### Magnetic and transport properties of submicron Gd strip

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Magnetic refrigeration utilizing magnetocaloric effects is expected as future environment-friendly refrigeration system. So far, the magnetocaloric effects have been investigated mainly in bulk systems. Recent nanofabrication techniques enable us to make the magnetocaloric materials with reduced dimensions. However, physical properties in such materials have not been investigated intensively. Here, we investigate the magnetic and transport properties of the Gd with reduced dimensions. Submiron-wide Gd strips have been prepared by the conventional lift-off technique combined with the electron-beam lithography. The magnetic properties have been evaluated by using a micro Hall cross based on GaAs/AlGaAs two dimensional electron gas (2DEG) system, and the transport properties have been evaluated by the magnetoresistance measurements. Clear magnetic hysteresis loops with the shape anisotropy were observed by the 2DEG Hall measurements, meaning that the Gd maintains still ferromagnetism even in the reduced dimension. On the other hand, the resistance change related to the magnetization was not observed. Since the f-electrons play an important role in the magnetocaloric effect of the Gd this result implies that the correlation between the conduction electrons and f-electrons is very small

#### **KD01**

Spin coupling, orbital angular momentum quenching, and electron localization in size-selected free transition metal clusters

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X-ray magnetic circular dichroism (XMCD) spectroscopy is a local and element specific technique to study spin and orbital magnetization. By applying XMCD spectroscopy to size-selected free clusters, we are able to resolve spin and orbit contributions to the magnetization in particles without any interaction with a support or matrix. This experimental technique allows us to study magnetic coupling in transition metal clusters, follow the quenching of the orbital angular momentum, and gain insight into the interaction of a magnetic impurity with a finite, few-atom free electron gas. Recent highlights will be discussed, including antiferromagnetic coupling of the central atom in Fe 13<sup>+</sup>, ferromagnetic spin coupling in Cr 2<sup>+</sup>, and the complex interplay of hybridization and on-site coulomb repulsion which determines the magnetic moment of impurity atoms in transition metal-doped gold clusters.

M. Niemeyer, K. Hirsch, V. Zamudio-Bayer, A. Langenberg, M. Vogel, M. Kossick, C. Ebrecht, K. Egashira, A. Terasaki, T. Moller, B. v. Issendorff, J. T. Lau, Phys. Rev. Lett. 108 (2012) 05720.

#### **KD02**

#### Nm-sized magnetic domains observed by small angle neutron scattering in exchange coupled superlattices

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#### **KD03**

#### Optical and magneto-optical characterization of Y<sub>0.5</sub>Bi<sub>2.5</sub>Fe<sub>5</sub>O<sub>12</sub> and Bi<sub>3</sub>Fe<sub>5</sub>O<sub>12</sub>thin films prepared by metal-organic decomposition

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In this work, optical and magneto-optical (MO) properties of Y05Bi25Fe5O12 (Bi:YIG) and Bi3Fe5O12 (BIG) thin films prepared by a metal-organic decomposition (MOD) method [1] have been studied. The Bi:YIG and BIG thin films were prepared on Gd<sub>3</sub>Ga<sub>5</sub>O<sub>12</sub> (100) substrates. After a spin-coating process, the samples were placed in a furnace with pre-annealing temperatures of 450°C and 400°C for 30 minutes and annealed at 750°C and 490°C for 3 hours to crystallize. Spectroscopic ellipsometry and magneto-optical spectroscopy were employed to derive a spectral dependence of permittivity tensor at photon energy range from 0.6 to 5 eV. Magneto-optical spectra displayed spectral features typical for Bi substituted garnets [2] reflecting fully developed garnet structure in all investigated films. At the energy range of 3.2~5 eV, Bi:YIG film showed higher amplitude of Kerr rotation than BIG, which is opposite to Faraday rotation measurements at the energy range 1.7~3.1 eV. This can be explained by higher surface roughness of BIG film according to atomic force microscopy measurements. Presented results showed high quality of investigated Bi:YIG films which demonstrates MOD as an effective technique to prepare highly substituted garnet films suitable for MO applications

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#### **KD04**

#### Fabrication of the epitaxial growth of (100) and (110) oriented Heusler alloy films for magnetic damping measurement.

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It has been reported [1] that the critical current density of MRAM is proportional to the damping parameter. Various investigations have reported a small damping parameter on the half metallic heusler alloy but the mechanism of magnetic damping is not yet clarified [2]. The present work aims at fabricating the epitaxial growth films for the damping measurement. Thus, Cr/Co<sub>2</sub>MnSi/Ru and Ta/Cr/ Co<sub>2</sub>MnSi/ Ta films were fabricated by sputtering respectively on MgO (100) substrate and on a-sapphire substrate at annealing temperatures ranging from 350 to 5500C. The pole figure analysis of the fabricated films show an epitaxial growth of Co<sub>2</sub>MnSi laver on MgO (100) and a-sapphire substrate. We thus, have successfully fabricated (100) and (110) oriented Heusler alloy films. The films contain also L21 order structure. In addition, the FMR and the VSM measurements were conducted in order to investigate the magnetic properties of those films. The line widths of ferromagnetic resonance spectrum depend on the annealing temperature of Co<sub>2</sub>MnSi layer. Those measurements displayed optimal conditions (a high Magnetization and a small damping constant) for the sample with Co<sub>2</sub>MnSi layer annealed at 4500C. The samples will have Co<sub>2</sub>MnSi laver annealed at 450 0C for further investigations.

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#### KE01

Thermalised and frozen magnetization dynamics in artificial spin ice Jason Morgan<sup>1</sup> Zoe Budrikis<sup>2</sup> Johanna Akerman<sup>3</sup> Aaron Stein<sup>4</sup> Paolo Politi<sup>5</sup> Sean Langridge<sup>6</sup> Robert Stamps7 and Christopher Marrows1\*

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Spin ices are rare earth pyrochlores where the crystal geometry leads to frustration of the rare earth moments [1], which meet at tetrahedra in the lattice. Nanotechnology allows many of the essential features of this physical system can be reproduced in arrays of nanomagnets [2]. This approach offers the opportunity to continuously tune the various parameters controlling the magnetic microstate, and also to inspect that microstate using advanced magnetic microscopy [3,4]. A significant difference with the naturally occurring spin-ices is that the change in symmetry gives rise to a true long-range ordered ground state, although the frustrated interactions in these athermal systems mean that its observation is extremely difficult [5]. Most attempts to achieve it rely on a rotating field demagnetisation protocol that produces states that show reasonably good ice rule fidelity but only short range correlations [2,5]. Here I will describe our recent work on an array of 250 nm × 80 nm Permalloy islands in the square ice geometry, including the achievement of a thermalised ground state during fabrication and the observation of the effects of fractionalised monopoles on excitations out of it [6], and athermal achievement of the ground state using a suitable field protocol [7].

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#### KE02

#### **Towards fully 3-dimensional MRAM**

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Spintronics could have a revolutionary impact on microelectronics and data storage if it provides the enabling step for transforming today's planar 2-dimensional devices into volume-filling 3-dimensional devices, where the data storage and processing capacity are related to the minimum feature size of the fabrication process. F. by F instead of F<sup>-2</sup>. Recently, we proposed a new approach to 3-dimensional spintronics in which topological kink solitons in multi-layered magnetic nanostructures are used to code and move data. Such solitons can be extremely stable at room temperature. highly compact, and easily injected, detected and synchronously propagated. As such, they are interesting candidates for use in ultrahigh density 3-dimensional MRAM devices. During this talk I will show the first experimental demonstration of a vertical shift register based on this principle and operating at room temperature. Eleven perpendicularly magnetised CoFeB layers coupled by Ru interlayers allow solitons coding binary '0' and '1' to be injected into the bottom layer and then propagated during eight successive clock edges up the stack, eventually exiting at the top layer.

#### KE03

#### Spin-transfer-torques-induced domain-wall motion in ferromagnetic Pt/Co/Pt nanowires with perpendicular magnetic anisotropy

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The spin-current-based electric controllability of ferromagnetic domain walls (DWs) has been of great interest because of its feasible applications. In particular nanowires with perpendicular magnetic anisotropy have been predicted as more suitable materials for practical applications. However, the experimental verification with multiple DWs control in such materials has not been explored vet. And also, the physical driving mechanism resulting from the competition between the spin transfer torques (STTs) adiabatic and nonadiabatic STTs and disorders is still controversial. By using Pt/Co/Pt nanowires, here we experimentally demonstrate that the purely current driven multiple DWs motion can be obtained at current densities less than 10<sup>7</sup> A/cm<sup>2</sup> allowing random magnetic bits recording and transferring. Furthermore, based on the DWs motion driven by magnetic field and/ or current, we can unambiguously distinguish the role of STTs on thermally activated magnetic domain wall motion. The adiabatic STT resulting in quadratic contribution to effective field is found to be dominant for large current densities, whereas the nonadiabatic STT playing the same role as magnetic field subsisted at low current densities. This finding will provide a step to the complete understanding of STT induced DW dynamics as well as the technological progress in DW based emerging nanodevices.

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#### **KF01**

#### Non-linear dynamics and high RF detection sensitivity in MgO-based spin-torque diode

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In the emerging field of 'spintronics', two electron attributes, namely, charge and spin, are simultaneously manipulated to obtain new, desirable characteristics, such as large magnetoresistance, spin-transfer induced magnetization switching, and precession in magnetic nanomaterials. In 2005, a phenomenon called the spin-torque diode effect [1] was demonstrated, which enabled the rectification of RF signals in magnetic tunnel junctions (MTJs); the rectification output, however, was several thousand times smaller than that of semiconductor diode detectors. Many efforts, such as the control of magnetic field direction [2,3] and the use of stochastic resonance [4], have significantly enhanced the sensitivity; nevertheless, the sensitivity of MTJs remain far lower than that of semiconductor diode detectors. In this study, we show that applying a DC bias current to MTJs while precisely controlling their magnetic potential profiles provides a high RF detection sensitivity of 12,000 V/W at room temperature, which exceeds that of semiconductor diode detectors (3,800 V/W). Analysis based on a macro-spin model reveals that this increase is caused by rotation of the precession axis that depends on RF input power. This rotation, a type of non-linear ferromagnetic resonance, causes a change in resistance and affords the high sensitivity.

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#### **KF02**

#### Finite tunnel magnetoresistance in junctions with a zero magnetization ferromagnetic electrode

Karine Dumesnil\*, Mathias Bersweiler, Mathieu Da Silva, Catherine Dufour, Michel Hehn, Daniel Lacour and Francois Montaigne Institut Jean Lamour - Lorraine University, France

# WITHDRAWN

#### KF03

Effects of mechanical rotation and vibration on spin currents Mamoru Matsuo\*, Jun'ichi Ieda, Eiji Saitoh and Sadamichi Maekawa Advanced Science Reserch Center, Japan Atomic Energy Agency, Japan

In the frontier of spintronics, much attention is paid on the control of spin currents. Due to the recent progress of nanoelectromechanics, mechanical manipulation of spins increases its importance. We discuss theoretically the generation of spin currents in both rotationally and linearly accelerating systems. The spin-orbit interaction argumented by inertial effects is derived from the low energy limit of the generally covariant Dirac equation. It is shown that the spin-orbit interaction is responsible for the generation of spin currents by mechanical rotation and vibration. We also investigate SU(2) x U(1) gauge theory in accelerating systems, which allows us to extend the spintronic theory in inertial frame to non-inertial frames.

M. Matsuo, J. Ieda, E. Saitoh, and S. Maekawa, Physical Review Letters 106, 076601 (2011); Applied Physics Letters 98, 242501 (2011); Physical Review B 84, 104410 (2011).

#### **KF04**

#### Design of self-organized nanostructures to achieve high blocking temperatures in MgO-based d<sup>0</sup> ferromagnets

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d<sup>0</sup> ferromagnetism, a new class of magnetism wherein the ferromagnets contain no magnetic elements, has created significant interest in the field of dilute magnetic semiconductors [1]. In this study, we present a materials design of MgO-based  $d^0$  ferromagnets with high blocking temperature ( $T_B$ ) based on multiscale simulations [2-4]. Chemical pair interactions between N atoms in Mg(O,N) and V<sub>MP</sub> in (Mg,V<sub>MP</sub>) O were calculated using a generalized gradient approximation and the VASP code. Using the Ising model, Monte Carlo simulations of the crystal growth were performed to predict the configurations of dopant distribution. In Mg(O,N), a strong attractive interaction was found between the 1st and 4th nearest neighbors (NN) in which the 1st-NN interaction is approximately twice as strong as the 4th-NN interaction. In (Mg,V<sub>MP</sub>)O, a strong attractive interaction only occurs between the 2nd NN. Our simulation showed that self-organized nanostructures can be formed both in Mg(O,N) and (Mg,V<sub>Me</sub>) O under layer-by-layer crystal growth, suggesting that these d<sup>0</sup> ferromagnets can obtain high  $T_{B}$  due to superparamagnetic blocking. Furthermore, depending on the simulation condition corresponding to the interval time between layers in deposition, various self-organized nanostructures can form, such as quasione-dimensional nanowires, two-dimensional nanoplanes and three-dimensional nanoclusters

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#### **KF05**

#### Negative spin-polarization of Fe<sub>4</sub>N observed by spin-resolved photoemission spectroscopy

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Ferromagnetic (Co,Fe), N has attracted a growing interest for application to spintronics devices because the large negative spin polarization at the Fermi level (EF) is expected. However, there have been no experimental reports confirming negative spin-polarization in (Co,Fe)<sub>4</sub>N by photoemission spectroscopy. In this study, we investigate the valence band (VB) electronic band structure and spin-polarization of Fe<sub>4</sub>N using hard x-ray photoemission spectroscopy (HX-PES) and spin resolved photoemission spectroscopy (SRPES) at room temperature. CaF2(1 nm)/Fe4N(5 nm) layered structure was epitaxially grown on SrTiO3(001) by molecular beam epitaxy. We measured VB structure of the Fe<sub>4</sub>N layer using HX-PES (hv = 5953 eV) at BL15XU of SPring-8, and compared the obtained spectrum with the calculated photoemission spectrum. After removing the CaF2 capping layer by Ar ion sputtering, we measured spin resolved density of states (DOS) below EF of the Fe<sub>4</sub>N using SRPES (hv = 21.2 eV) at HiSOR, Hiroshima University. The peak position of measured Fe 3d electron photoemission spectrum by HX-PES was well explained by the theoretical DOS. We deduced spin-polarization from the obtained SRPES spectrum by eliminating the background signals of secondary electrons caused by He II resonance line (hy = 40.8 eV), and confirmed negative spin-polarization at EF in Fe<sub>4</sub>N.

#### KG01

#### Hidden order pseudogap and hybridization modulation in URu<sub>2</sub>Si<sub>2</sub>

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A. Balatsky (Los Alamos) URu<sub>2</sub>Si<sub>2</sub> proved to be a compound that has exhibited similarities with other correlated materials like HighTc oxides in that it has strong correlations, competing phases and intense nanoscale inhomogeneity in Kondo lattice. We point out that Hidden order transition is not a mean field transition, in contrast to prevailing discussion to date. It has significant precursor effects, so called pseudogap. Through an analysis and modeling of data from various experimental techniques, we present evidence for the presence of a hidden order pseudogap in URu<sub>2</sub>Si<sub>2</sub> in the temperature range between 25 K and 17.5 K. We evaluate the effects that gap fluctuations would produce on observables like tunneling conductance, neutron scattering and nuclear resonance, and relate them to the experimental findings. We show that the transition into hidden order phase is preceded by the onset of non-coherent hidden order fluctuations. We also discuss nanoscale inhomogeneity seen in URu<sub>2</sub>Si<sub>2</sub> with STM as an evidence for hybridization modulation due to local defects and discuss the role of hybridization modulations in Hidden

Hidden order pseudogap in URu Six J. Haraldsen, 10.1103/PhysRevB.84.214410 (2011) Electronic inhomogeneity in a Kondo Thaten onder pseumogap in Charge 2: Induated, 10:11 parts 10:21 (2017) accurate (2017) accurate anomageneii) in a Komol Datarie, E. Bauer et al. PMAS 106, 6857 (2011) How Nondo Holes Create Intense Noncode Heary-Fermion Hybridization Disorder/PMAS 2011 108 (45) 18233-18237: published ahead of print October 17, 2011, doi:10.1073/pnas.1115027108 Anomalous femanescond quasiparticle dynamics of hidden order state in URu2S2 Georgi L. Dakovski et al. Phys. Rev. B 44, 161103 (2011) Hybridization Wave as the "Hidden Order" in URu2S12, Y. Dubi et al, Phys. Rev. Lett. 106, 086401 (2011) Incommensurate spin resonance in URu2Si2, A. V. Balatsky, et al, Phys. Rev. B 79, 214413

#### KG02

#### Strain-effect on topological quantum phase transition in Ir-oxides Jaeiun Yu

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We predict that a three-dimensional strong topological insulator (TI) phase can be realized by controlling the electron hopping terms in Na2IrO3. From a realistic tightbinding Hamiltonian, constructed from the results of first-principle calculations, we identified the crucial parameters of the effective low-energy Hamiltonian, which determine the topological nature of the spin-orbit coupled ground state. We present a phase diagram exhibiting a quantum phase transition from normal-to-topological insulators in terms of the spin-orbit coupling strength and the hopping parameter. The existence of the TI phase has also been verified by the parity analysis from the firstprinciples calculations with a modified crystal structure corresponding to the predicted tight-binding parameters. We clarified that the topological nature of the Na<sub>2</sub>IrO<sub>3</sub> band structure is sensitive to the structural variation. The long-range hopping and trigonal crystal field play a crucial role in addition to the strong spin-orbit coupling in Na2IrO3. The TI phase realized in Na<sub>2</sub>IrO<sub>3</sub> can be a "good" playground for the study of the interplay between spin-orbit coupling and on-site Coulomb interaction.

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#### KG03

#### Hollandites - theoretical aspects of their unique electronic properties

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Density-functional-theory-based electronic structure calculations are made to clarify the mechanisms of the observed anomalous electronic properties in transition-metal oxides with hollandite-type crystal structure. We first discuss some general aspects of the electronic structures of the 3d and 4d series of hollandites, focusing on the validity and insufficiency of the rigid-band filling of electrons in the series. Then, we in particular discuss (i) the Peierls mechanism of the metal-insulator transition in ferromagnetic hollandite K22Cr8O16 [1], (ii) superatomic-crystal-type electronic structure emerging in molvbdenum hollandites K<sub>2</sub>Mo<sub>6</sub>O<sub>16</sub> and Rb<sub>16</sub>Mo<sub>6</sub>O<sub>16</sub> [2], and (iii) Tomonaga-Luttinger-liquid properties emerging in ruthenium hollandite K2Ru8O16 [3]. We also consider the essential role of electron correlations played in the metal-insulator transition of vanadium hollandite K<sub>2</sub>V<sub>2</sub>O<sub>2</sub>/

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#### KG04

#### Inter-band pairing and inhomogeneous superconductivity in multiorbital systems

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In multi-band systems, electrons from different orbitals coexist at a common Fermi surface. An attractive interaction among these quasi-particles gives rise to inter-band or hybrid pairs, which eventually condense in a superconducting state. These quasiparticles have a natural mismatch of their Fermi wave-vectors,  $\delta kF$  which depends on the strength of the hybridization between their orbitals. The existence of this natural scale leads us to consider the possibility of inhomogeneous superconducting ground sates in these systems, even in the absence of an applied magnetic field. Furthermore, since hybridization V depends on pressure this provides an external parameter to control the wave-vectors mismatch at the Fermi surface. In this work, we study the zero temperature phase diagram of a two-dimensional, two-band superconductor with interband pairing. We show that as the mismatch between the Fermi wave-vectors of the two-bands is reduced, the system presents a normal-to-inhomogeneous superconductor quantum phase transition of unusual character at a critical value of the hybridization,  $Vc = \Delta 0$ , below which the system condenses in pair-wave superconducting state characterized by a wave-vector  $qc = \Delta 0/vf$  ( $\Delta 0$  is the homogeneous superconducting gap and vf the Fermi velocity.

#### **KH01**

#### High performance permanent magnets for energy applications

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A new energy paradigm, consisting of greater reliance on renewable energy sources and increased concern for energy efficiency in the total energy lifecycle, has accelerated research in energy-related technologies. Due to their ubiquity, magnetic materials play an important role in improving the efficiency and performance of devices in electric power generation, conversion and transportation. Magnetic materials are essential components of energy applications (i.e. motors, generators, transformers, actuators, etc.) and improvements in magnetic materials will have significant impact in this area, on par with many "hot" energy materials efforts. The lecture focuses on the synthesis, characterization, and property evaluation of high performance permanent magnets. Using multi-scale characterization and modeling the structure-property relationships will be examined. Options for the development of novel resource efficient magnets will be discussed in the context of drastically reducing the rare-earth contents or developing rare earth free magnets. Finally, considering future bottle-necks in raw materials and in the supply chain, options for recycling of rare-earth metals will be analyzed.

O. Gutfleisch, J.P. Liu, M. Willard, E. Bruck, C. Chen, S.G. Shankar, Magnetic Materials and Devices for the 21st Century: Stronger, Lighter, and More Energy Efficient (review), Adv. Mat. 23 (2011) 821-842.

#### KH02

#### Influence of surface anisotropy on orientation of crystal grain in rareearth permanent magnet

Chiharu Mitsumata\*, Hiroki Tsuchiura and Akimasa Sakuma Tohoku University, Japan

The coercive force of Rare-Earth permanent magnets has been investigated. In general, the nucleation and the domain wall pinning are discussed to understand the magnetization reversal process [1]. However, the large discrepancy between the ideal anisotropy field (Hk) and the coercivity (Hc) still remains in an unsolved problem of the permanent magnet. Recently, we proposed the possible mechanism of the reduction for Hc in NdFeB system, that was realized owning to the negative anisotropy of the rare-earth ion at the surface of a crystalline [2, 3]. In this paper, we discuss the influence of the surface anisotropy on the mis-orientation of the crystal grain in the permanent magnets. In the case of simultaneous rotation of the surface anisotropy effect and the magnetic easy axis increases. For comparing the ideal simultaneous rotation, we assume the polycrystalline magnet in which the surface anisotropy effects on the crystal grains. As the result, the surface anisotropy causes the large reduction of Hc from Hc, however, a tolerance of Hc reduction against the mis-orientation of crystal grains becomes promising.

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#### KH03

#### International comparison of the properties of permanent magnets measured using an electromagnet and a pulse field magnetometer Michael Hall

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An IEC TC 68 comparison on the measurement of the magnetic properties of permanent magnets was recently completed. With permanent magnets being used in safety critical situations it is necessary to perform 100% quality control at the operating temperature. An established method using an electromagnetic to measure the demagnetising curve is not practical for these measurements due to the time taken, the difficulty in raising the temperature and the operator skill required. A Pulsed Field Magnetometer (PFM) has been developed that can determine the full BH curve in as little as 100 ms. By comparing measurement made using electromagnet and pulse methods, the influence of the dynamic effects of the latter could be established. It was found that for the technology significant parameter of intrinsic coercivity a significant difference exist between the methods. The talk will present the results from measurements on 6 NdFeB magnets with intrinsic coercivities ranging from 1000 to 2600 kA/m by 8 institutes based in China, Japan, Italy, Belgium, Germany and the UK. For the electromagnet method, the Influence of local pole piece saturation on the measurement of the intrinsic coercivity will be discussed.

#### KH04

Ab initio calculations of magnetic moment and magneto-crystalline anisotropy: New ternary alloy Mn-Bi-Co permanent magnet

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There are two important issues to address for full utilization of permanent magnets in electric vehicle, windmill generator, and other motors. One is a technical issue such as low operation temperature of permanent magnet, and the other is rare-earth metal source and price. Rare-earth permanent magnets, including Nd<sub>2</sub>Fe<sub>14</sub>B and Dy-doped Nd<sub>2</sub>Fe<sub>14</sub>B, Sm<sub>2</sub>Co<sub>17</sub>, and Sm<sub>2</sub>Fe<sub>17</sub>N<sub>3</sub>, have been used for motors for hybrid vehicle and electric vehicle. The Nd<sub>2</sub>Fe<sub>14</sub>B possesses 60 MGOe of the highest maximum energy product (BH)max. All these permanent magnets contain rare-earth metals. Therefore, in order to address the second issue, we have proposed rare-earth-free Mn-Bi-Co ternary alloy and performed first principles calculations to investigate the effects of the alloying element Co on magnetic properties of conventional Mn-Bi binary alloy magnet. Our calculations predict that an addition of C to the Mn-Bi magnet will likely to increase the magneto-crystalline anisotropy by more than 300%, while the saturation magnetization will be increased by 20%. Accordingly, it is predicted that the (BH)max increases to 21 MGOe from 17 MGOe of the Mn-Bi permanent magnet. The proposed Mn-Bi-Co magnet exhibits a much higher maximum energy product than the conventional Mn-Bi magnet. Experimental verification of these results is in progress.

#### KI01

## First-principles prediction of large perpendicular magnetocrystalline anisotropy of 4d-monolayers on bcc-Fe(001) surface

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Using the highly precise full-potential linearized augmented plane-wave method based on general gradient approximation, we investigated the magnetism and magnetocrystalline anisotropy energy ( $E_{\rm MCA}$ ) of 4d Ru, Rh, and Pd monolayers (MLs) on a bcc-Fe(001) surface. Coupling with Fe ferromagnetically, the Ru, Rh, and Pd monolayers are induced to have the magnetic moments of 0.43, 0.92, and 0.36  $\mu_{\rm B}$ , respectively, in consistent with experimental [1,2] and previous theoretical results [3]. The calculated magnetic moments of Fe at the interface are significantly enhanced 2.8  $\mu_{\rm B}$  for Rh and 2.92  $\mu_{\rm B}$  for Pd, compared to the bulk bcc-Fe value, but Ru makes the Fe atoms to retain the bulk value (2.2  $\mu_{\rm B}$ ). We predict that the Ru and Rh MLs have the perpendicular huge  $E_{\rm MCA}$  of -0.2 meV/atom. The huge positive  $E_{\rm MCA}$  with induced magnetic moments of the Ru and Rh MLs on a Fe(001) surface are expected to be promising in application to devices such as high-density magnetic recording media and random access memory. The electronic origins of the induced magnetism and  $E_{\rm MCA}$  will be discussed with the density of states.

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#### KI02

## Ferromagnetic phase at the LaAlO<sub>3</sub>/SrTiO<sub>3</sub> (001) interface induced by SrTiO<sub>3</sub> lattice deformation

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We address the origin of the ferromagnetic phase formed at the LaO/TiO<sub>2</sub> interface in LaAIO<sub>3</sub>/SrTiO<sub>3</sub> (001) heterostructures [1]. Using first-principles density functional calculations we show that the charge transfer to the Ti t2g conduction band, induced by the polar mismatch at the interface, combined with the incipient ferroelectric character of the SrTiO<sub>3</sub>, lead to a substantial lattice deformation. The ferroelectriclike distortion of the SrTiO<sub>3</sub> slab generates the spontaneous spin-polarization of the lowest dxy orbital parallel to the interface and the emergence of a ferromagnetic ground state. Moreover, it yields differential occupation of the non-degenerated dxy and degenerated dx-yz orbital states. We predict the existence of three different types of electrons, magnetic two-dimensional (2D) dxy electrons and non-magnetic higherlying dxy and dxz-yz quasi-2D electrons. While the former, with a 2D-light effective mass (m\*), are confined to the TiO<sub>2</sub> interface plane, the non-spin polarized higher-lying dxy also with a light m\*, and the Bloch dxz-yz electrons spread over several SrTiO<sub>3</sub> layers. Their differences could explain the coexistence in the same sample of superconductivity and magnetism [2-5]. Furthermore, since fercelectric field.

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#### KI03

#### Interfacial magnetic anisotropy of junctions between Fe and transitionmetal nitrides or carbides: A first-principles study

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Recently, CoFeB/MgO-based perpendicular magnetic tunnel junctions (MTJ) with high thermal stability have been fabricated[1]. For further reduction of MTJ size, larger perpendicular magnetic anisotropy (PMA) is required. In the MgO/Fe(001) interface, the hybridization between O-2p and Fe-3d orbital decreases the in-plane magnetic anisotropy (IMA) component. This effect is also expected for Fe/nitride(carbide) interfaces. In order to design large PMA film, we investigated magnetic anisotropy of Fe(001) and NaCI-type transition-metal nitride(carbide) interfaces (Fe/MN(C): M=Sc, Ti, V, Cr, Zr, Nb, Hf, and Ta) by using a first-principles density-functional calculations. Large PMA energy is obtained for the Fe/ScN (1.21 mJ/m2) and Fe/ScC (0.89 mJ/m2) film. The origin of PMA component of these films is attributed to the surface PMA of Fe(001) (0.60 mJ/m2). In other films, weak PMA or IMA is obtained due to the reduction of PMA component caused by the strong hybridization between interfacial Fe-3d and transition-metal-d orbitals. We also discuss the influence of boron atoms at the interfaces. This work was supported by JSPS though the FIRST Program initiated by the CSTP and a Grant-in-Aid for Scientific Research (Grant No. 22360014) from JSPS.

[1] S. Ikeda et. al., Nature Mater. 9, 721 (2010).

#### KI04

#### Enhancement of exchange coupling by incoherent quantum resonance Ching Hao Chang

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By simulating the nanoscale Fe/Ag/Fe trilayer in the quantum well potential, the coupling strength is found to oscillate with Fe thickness. This is generated by the competition between quantum interferences for ma jority carriers in ferromagnetic configuration and for carriers in anti- ferromagnetic configuration. Furthermore, we fix the Fe side-layer thickness and study the coupling as the function of spacer width. The strength oscillates with the period determined by quantum well states of minority carriers in ferromagnetic configuration while the amplitude decided by this competition of interferences. This amplitude can be enhanced almost one time by choosing the proper side-layer thickness that satisfies (1) the conditions for incoherent quantum interferences and (2) much thinner than the spacer width.

#### KI05

#### Scrolling effects in vanadium oxide nanotubes and nanolayers

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The comparative analysis of the ESR together with the magnetic susceptibility data [1] has allowed elucidating effects of scrolling the VO<sub>x</sub> nanolayers, NLs, (stacks of 4 to 8 distorted layers of vanadium oxide separated by an organic template) into VO<sub>x</sub> nanotubes, NTs. First, scrolling leads to formation of the antiferromagnetic dimers consisting of almost half of the vanadium ions with a spin gap near 720 K. The second effect, specific to NTs, is the correlated change in the Curie constant for the nondimerized V<sup>4+</sup> quasifiee spins and in variations of the background formed by Van Vleck and Pauli terms: the decreasing temperature in the interval 70°T<120K induces an initial 1.8-fold growth of the Curie constant followed by a decrease, along with a change of the sign of the background. Third, it is found that scrolling results in reduction of the quasifiee electron concentration in the VO<sub>x</sub> plane by a factor of 2-4. It is found that the number of electrons in the V<sup>4+</sup> state is always less than the total electron concentration, and the concentration of the "excessive" electrons may vary in the range 0.11-0.46 per vanadium site.

1.S.V.Demishev et al., Phys. Rev. B 84, 094426 (2011)

#### KJ01

## THz Emission from High-T\_ Superconductor $Bi_2Sr_2CaCu_2O8_{*\,\delta}$ Intrinsic Josephson Junctions

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After the discovery of THz emission of radiation from high- $T_c$  superconductor Bi<sub>2</sub>Sr<sub>2</sub>CaCu<sub>2</sub>O<sub>8+0</sub> in 2007, intensive studies on the mechanism of radiation from high- $T_c$  superconductors have been carried out both theoretically and experimentally. So far, two important ingredients are found for the necessary condition of radiation: One is the ac-Josephson effect working in the individual intrinsic Josephson junction existing in a unit cell of the crystal structure of high- $T_c$  superconductors. This generates the electromagnetic waves at a fixed frequency according to the ac-Josephson effect and the power of the radiation at the frequency can be enhanced by the cavity resonance, which is determined by the geometry of the mesa structure. This mechanism works as a second important mechanism for the radiation. We show experimental results confirming these important ingredients for the emission of radiation to nearly 1 mW level by the array of three mesas synchronized together quantum mechanically. We will also show a prototype imaging machine for the applications.

#### KJ02

#### Electric field control of magnetization in spiral magnets

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The control ofmagnetization by an electric field remains as one of greatchallengesin materials science. Multiferroics, in which magnetism and ferroelectricitycoexist and couple to each other, could be the most plausible candidate to realize this longsoughtcapability. While recent intensive research on the multiferroics has made significant progress in the sensitive magnetoelectric (ME) effects, the electrical control of magnetization, the converse effect, has been observed only in a limited range far below room temperature. Here Iwill summarizerecent developments on how the spiral magnets, particularly based on the hexagonal ferrites, can realize the viable, converse ME effectin a broad temperature range including room temperature. Starting with the chemical aspect of the hexaferrite materials, I discuss the importance of the strong magnetoelectric(ME) coupling to achieve the significant converse ME effect. The three key material parameters, magnetic anisotropy, spin frustration, and spin exchange interaction, will be pointed out. By controlling those key parameters, we demonstrate that the four state magnetization curves can be obtained at room temperature by an electric field. We also show based on the similar mechanism that the magnetization reversal by an electric field can be realized with fast switching time.

Sae Hwan Chun et al., Phys. Rev. Lett. 108, 177201 (2012); ibid,104, 037204 (2010); Y. S. Chai et al., (unpublished).

#### KJ03

#### Novel Josephson effect in triplet Josephson junctions: The story begins Dirk Manske

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In the theoretical study of Josephson junctions, it is usually assumed that the properties of the tunneling barrier are fixed. This assumption breaks down when considering tunneling between two triplet superconductors with misaligned d-vectors in a TFT-junction (triplet ferromagnet triplet) [1,2]. Such a situation breaks time-reversal symmetry, which radically alters the behaviour of the junction, stabilizing it in a fractional state, i.e. the free energy minimum lies at a phase difference intermediate between 0 and  $\pi$ . Fractional flux quanta are then permitted at the junction [3]. A further consequence of the d-vector misalignment is the appearance of a Josephson spin current, which flows even in the absence of an equilibrium charge current. Not only do our calculations enhance the physical understanding of transport through triplet superconductor junctions, but they also open the possibility of novel spintronic Josephson devices [4].

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 P.M.R. Brydon and D. Manske, Phys. Rev. Lett. 103, 147001 (2009).

#### KJ04

#### Field-induced polarization of dirac valleys in bismuth

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The principal challenge in the field of "valleytronics" is to lift the valley degeneracy of electrons in a controlled way. In bulk semi-metallic bismuth, the Fermi surface includes three cigar-shaped electron valleys lying almost perpendicular to the high-symmetry axis known as the trigonal axis. The in-plane mass anisotropy of each valley exceeds 200 as a consequence of Dirac dispersion, which drastically reduces the effective mass along two out of the three orientations. According to our very recent study of angle-dependent magnetoresistance in bismuth [1], a flow of Dirac electrons along the trigonal axis is extremely sensitive to the orientation of in-plane magnetic field. Thus, a rotatable magnetic field can be used as a valley valve to tune the contribution of each valley to the total conductivity. As a consequence of a unique combination of high mobility and extreme mass anisotropy in bismuth, the effect is visible even at room temperature in a magnetic field of 1 T. At high temperature and low magnetic field, the three valleys are interchangeable and the three-fold symmetry of the underlying lattice is respected. As the temperature is decreased or the magnetic field increased, this symmetry is lost.

1. Z. Zhu, A. Callaudin, B. Fauque, W. Kang and K. Behnia, Nature Phys. 8, 89 (2012)



## POSTER PRESENTATION

- July 9 (PA~PO) 116
- July 10 (QA~QP) 182
- July 12 (RA~RR) 249
- July 13 (SA~SO) 316

#### PA01

## Magnetoelectric polarization in the field-induced commensurate phase of Y-hexaferrite $Ba_0$ , $Sr_1$ , $Zn_2$ (Fe<sub>1</sub>, $Al_y$ ), $O_{22}$

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Hexagonal ferrites containing  $Fe^{3+}$  (S = 5/2) ions are attracting much attention because it is expected that solutions for high-temperature low-field multiferroics can be found in these materials. Y-type Ba<sub>0.7</sub>Sr<sub>1.3</sub>Zn<sub>2</sub>Fe<sub>12</sub>O<sub>22</sub> possesses proper-screw antiferromagnetic order in zero field, which does not allow spin-current ferroelectric polarization. Electric polarization can, however, be induced when a strong external field is applied perpendicular to its c axis. The field-response can greatly be improved upon Al doping, which is ascribed to the formation of conical spin structures. In this work, we investigated the magnetic phase transitions of single crystal Ba0,7Sr13Zn2(Fe1  $_{x}Al_{x})_{12}O_{22}$  (x = 0.04 & x = 0.08) in external magnetic field. Both samples exhibited incommensurate helical antiferromagnetic order in zero-field cooling down to 10 K. When the external magnetic field was applied so that electric polarization was induced, however, the incommensurate order was replaced by a commensurate order in all samples. Our analysis shows that the commensurate order plays the role of breaking inversion symmetry, thereby allowing finite electric polarization. We argue that two different magnetic anisotropies, easy-axis and easy-plane, play an important role in establishing stable field-induced electric polarization in the observed commensurate magnetic phase

#### PA02

## Temperature- and field-tuning of magnetic phases in multiferroic $NdFe_4(BO_4)_A$

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We report the low-temperature coexistence of magnetic phases in NdFe<sub>3</sub>(BO<sub>3</sub>)<sub>4</sub>, which is a field-induced multiferroic material. The ground state incommensurate magnetic phase coexists with a strained commensurate magnetic phase that is primarily at the surface of the crystal. Increasing the temperature or magnetic field decreases the incommensurability and stabilizes the commensurate magnetic phase above  $T_{\rm g} \approx 14$  K or H<sub>sc</sub> = 0.9 T. A comparison to published studies indicates the onset of longitudinal magnetostriction and electric polarization at the field-induced transition, which may arise due to a basal plane spin-flop and canting of moments along the field direction. Use of the National Synchrotron Light Source, Brookhaven National Laboratory, was supported by the U.S. Department of Energy, Office of Science, Office of Basic Energy Sciences, under Contract No. DE-AC02-98CH10886.

#### PA03

## Control of magnetic anisotropies for stable electric polarization in multiferroics hexaferrites

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Hexaferrites are widely studied for the hopes of realizing multiferroics working at room temperature and under low field. It was shown recently that Al-doping can effectively enhance magnetoelectric susceptibility and low-field behaviors of Y-type hexaferrite Ba<sub>0.5</sub>Sr<sub>1.5</sub>Zn<sub>2</sub>Fe<sub>12</sub>O<sub>22</sub>. The key for such enhancement has been suspected to be the formation of heliconical spin ordering via reduction of planar anisotropy. Using single crystal neutron diffraction, we investigated in detail the Al-doping dependence of the magnetic phase transitions in Ba<sub>0.5</sub>Sr<sub>1.5</sub>Zn<sub>2</sub>(Fe<sub>1.4</sub>Al<sub>4</sub>)<sub>1.7</sub>O<sub>22</sub>. We confirm that the heliconical phases indeed appear upon Al doping, and finally collapse into commensurate phases at excessive doping concentrations. The observed evolution of magnetic order coincides with the doping dependence of the magnetoelectricity previously reported. Analysis of magnetic energy shows, however, that planar anisotropy alone will not allow the heliconical order to stabilize. We argue that it is important to have a competition between easy-axis and easy-plane anisotropies in order to stabilize the heliconical order, and thereby establish the optimum magnetoelectric

#### PA04

#### Chemical control of ferroelectric polarization in Mn<sub>1-x</sub>Co<sub>x</sub>WO<sub>4</sub>

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MnWO<sub>4</sub> is one of the best known magnetoelectric multiferroics, in which ferroelectric polarization appears due to long-range spin order. Whereas it exhibits three different antiferromagnetic phases at low temperatures, ferroelectric polarization along the b axis is observed only in the intermediate phase with an elliptical spiral incommensurate spin structure. In this work, we present the effect of Co doping on the magnetic phase transitions and multiferroics properties of Mn<sub>1-x</sub>Co<sub>x</sub>WO<sub>4</sub> using neutron diffraction, magnetization and ferroelectric measurements. The main results include the stabilization of the incommensurate phase at x  $\sim$  0.03 and the flop of its spiral plane for x > 0.05. Subsequently, the ferroelectric polarization appears mainly along the a axis in this range. These results indicate that the multiferroics properties of MnWO<sub>4</sub> can be controlled via Co doping.

#### PA05

## Multiferroic phase competition in orthorhombic RMnO<sub>3</sub>: Monte Carlo approaches

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For type-II multiferroic oxide compounds, the ferroelectricity is believed to be induced by specific spin orders which can break the spatial inversion symmetry, allowing an electric polarization to respond significantly to external magnetic field. These phenomena have been receiving attentions due to the underlying physics and promising application potentials. Nevertheless, conventional understanding of these multiferroics remains in an early stage since these materials exhibit typical features of strongly correlated electron systems in which the spin, charge, orbit, and phonon degrees of freedom are intrinsically coupled together. So far, rare models other than the Mochizuki-Furukawa (M-F) model specifically addressing orthorhombic RMnO3 are available. In this talk, we present our Monte Carlo simulations on the multiferroic behaviors of orthorhombic RMnO<sub>2</sub>, based on this model. Our simulations address the phase competition between the ab-plane spiral spin order and bc-plane spiral spin order as well as their coexistence, including the following cases: (1) nonmagnetic Mn-site substitution; (2) R-site substitution; (3) interaction between multiferroic RMnO3 thin films and magnetic substrates. We will further extend our simulation based on the semiquantum two-orbit double exchange model to the above issues

#### PA06

#### Spin-driven electric polarization in akermanite Sr<sub>2</sub>MSi<sub>2</sub>O<sub>7</sub> crystals Mitsuru Akaki<sup>1</sup>\*, Tomova Tadokoro<sup>1</sup>, Takumi Kihara<sup>2</sup>, Masashi Tokunaga<sup>2</sup> and Hideki

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Since the magnetoelectric effect was observed in akermanite  $Ca_2CoSi_2O_7$  [1] and  $Ba_2CoGe_2O_7$  [2,3], the multiferroicity in these materials has been attracting much attention. The akermanite materials show the magnetic-field-induced electric polarization which can be explained by the transition-metalligand (p-d) hybridization depending on the spin direction [4]. The multiferroicity of akermanite materials is, therefore, expected to be influenced by substitution on the transition-metal site. In this study, we have investigated magnetic and dielectric properties of akermanite  $Sr_2MSi_2O_7$  (M: transitionmetal ion) single crystals. In  $Sr_2CoSi_2O_7$ , the electric polarization was not observed in a zero magnetic field. By applying the magnetic field along the [110] direction, the electric polarization mechanism, the electric polarization is proportional to the spin quantum number S. However, the induced electric polarization in  $Sr_2MnSi_2O_7$  ( $Mn^{2+2}$ : S=5/2) is almost two orders of magnitude smaller than that of  $Sr_2CoSi_2O_7$  ( $Co^{2+2}: S=3/2$ ). These results suggest that the electric configuration of the transition-metal ion is of crucial importance for the appearance of electric polarization. We will also present and discuss the results of the comparative study of magnetoelectric properties in other  $Sr_2MSi_2O_7$  crystals.

[1] M. Akaki et al., Appl. Phys. Lett. 94, 212904 (2009). [2] H. T. Yi et al., Appl. Phys. Lett. 92, 212904 (2008). [3] H. Murakawa et al., Phys. Rev. Lett. 105, 137202 (2010). [4] T. Arima, J. Phys. Soc. Jpn. 76, 073702 (2007).

#### PA07

#### Magnetoelectric effects in antiferromagnet Ba<sub>2</sub>CoGe<sub>2</sub>O<sub>7</sub> Shin Miyahara<sup>1</sup> and Nobuo Furukawa<sup>2</sup> <sup>1</sup>JST ERATO-MF, Japan <sup>2</sup>JST ERATO-MF, Aovama Gakuin University, Japan

We have investigated magnetoelectric effects in Ba<sub>2</sub>CoGe<sub>2</sub>O<sub>7</sub>, which is a quasi twodimensional antiferromagnet[1]. By applying the external magnetic field along certain directions, electric polarization is induced[2,3]. We consider a spin-dependent d-p hybridization mechanism[4], where the electric polarization couples to a single spin structure via the spin-orbit interaction, and we can explain the novel magnetoelectric behaviors in Ba<sub>2</sub>CoGe<sub>2</sub>O<sub>7</sub>[3]. In this compound, it is recently reported that there exist magnetic excitations induced by electric components of electromagnetic wave at THz frequencies[5]. In addition, the magneto-electric excitation in Ba<sub>2</sub>CoGe<sub>2</sub>O<sub>7</sub> accompanies a non-reciprocal linear directional dichroism, where absorption coefficient depends on the direction of the incident electromagnetic wave (+k or -k). We theoretically clarify the origin of the peculiar features in these dynamical magnetoelectric phenomena[6]. The non-reciprocal dichroism observed in Ba<sub>2</sub>CoGe<sub>2</sub>O<sub>7</sub> can be explained by cross-correlated effects, i.e., the interference between the magnetic and the electric responses.

[1] A. Zheludev et al., Phys. Rev. B 68 024428 (2003), [2] H.T. Yi et al., Appl. Phys. Lett 92 212904 (2008), [3] H. Murakawa et al., Phys. Rev. Lett. 105 137202 (2010) [4] T. Arima, J. Phys. Soc. Jpn. 76 073702 (2007), [5] I. Kezsmrki et al., Phys. Rev. Lett. 106 057403 (2011) [6] S. Miyahara and N. Furukawa, J. Phys. Soc. Jpn 80 073708 (2011)

#### PA08

#### X-ray non-reciprocal effects in multiferroic single crystal of GaFeO<sub>3</sub> Andrei Rogalev\*, Fabrice Wilhelm and Alexei Bosak

European Synchrotron Radiation Facility, France

Ferroelectric and ferrimagnetic ordering coexist in gallium ferrate - GaFeO<sub>3</sub>. The multiferroic properties of this compound have been extensively studied since early sixties by many different experimental techniques. In order to study these properties on a microscopic level we have measured various X-ray reciprocal and non-reciprocal dichroisms at the Fe K-edge. The results of these experiments carried out at the ESRF beamline ID12 on a high quality untwined single crystal of GaFeO<sub>3</sub> are presented here. X-ray magnetochiral dichroism and X-ray non-reciprocal magnetic linear dichroism, which are a direct measure of the Fe orbital anapole and higher order magnetoelectric multipole, are disentangled experimentally for the first time. These two dichroic signals are compared to usual X-ray magnetic circular dichroism which is a measure of orbital magnetic moment carried by Fe atoms. Moreover, X-ray natural circular dichroism has been also measured on this biaxial non-enantiomorphous crystal. All these dichroisms are analyzed with the help of a set of the sum rules. This analysis allowed us to deduce the expectation values of different effective operators related to multiferroic properties of GaFeO<sub>3</sub>.

#### PA09

#### NMR study on Ba0.5Sr1.5Zn2(Fe0.92Al0.08)12O22

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We did zero-field NMR study on the annealed single crystal of multiferroic helimagnets  $Ba_{0.5}Sr_{1.5}Zn_2(Fe_{0.92}Al_{0.08})_{12}O_{22}$  (Al-BSZFO). This material shows extremely high magnetoelectric susceptibility so that the critical field for switching electric polarization is less than 1 mT below 90 K [1]. NMR frequency change by the temperature follows Bloch's T3/2 law which presents the low temperature excitation is ferromagnetic spin wave. The nuclear spin-lattice relaxation rate and the nuclear spin-spin relaxation rate were also measured. Both increased rapidly as the temperature increases above 60 K at which the spin structure changes from normal longitudinal cones (NLCs) to alternating longitudinal cones (ALCs) for the case of as-grown (not annealed) sample [2]. Due to rapid shortening of spin-spin relaxation time, the NMR signal intensity abruptly reduced above 60 K. We also studied rf pulse width and power dependence and revealed that most of the signal came from the domain walls.

S. H. Chun et al., Phys. Rev. Lett. 104, 037204 (2010).
 H. B. Lee et al., Phys. Rev. B 83, 144425 (2011).

#### PA10

#### Magnetic ordering in multiferroic TbFexMn $_{2-x}O_5$ with x=0.18

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Manganite compounds present a large range of functional properties from colossal magnetoresitor to multiferroicproperties. TbMn<sub>2</sub>O<sub>5</sub> is well known multiferroic materials that exhibit complex magnetic phase transitors [1] and multifunctional properties which are tunable by a magnetic field [2,3]. Below 38K, the long range incommensurate anti-ferromagnetic ordering induces the feroelectric behavior through either an exchange striction mechanism between neighboring spins wave (Si.Sj#0) [4] or inverse Dzyaloshinskii-Moriya interaction due to cycloidal spin arangement [5]. All the magnetic properties in these systems are controlled by frustrations. The frustration can be tuned in increasing or decreasing the interactions by substituting Mn<sup>3+</sup>/Mn<sup>4+</sup> with Fe<sup>3+</sup>/Fe<sup>4+</sup> (HS, S=52) without structural distortion. Fe doped materials in ceramics sample at similar composition have shown enhancement of Ps. Large single crystals of TbMn<sub>2-x</sub>Fe<sub>x</sub>O<sub>5</sub> (x=0.18) have been grown by flux method and characterized by (ac and dc) susceptibility measurements. Single crystal neutron diffraction on E4 ard TriCs diffractometer have been carried out respectively at the Helmoltz Zentrum Berlin and Sing. Paul Scherr Institut. Below 38K, Mn/Fe one and orders antiferromagnetically  $(\pi(0.5, 0, q2)$  qz–(0.21, 0.245)]. Below 15K, Tb sublatice is polarized by Mn/Fe one and orders antiferromagnetically be.

[1] E. F. Berthaut et al., Solid. Stat. Comm. 5,25 (1925) [2]I. Akinobu and K. Kay, Ferroelectrics 169, 75 (1995) [3] N. Hur et al., Nature 429, 392 (2004). [4] L. C. Chapon, Phys. Rev. Lett. 96, 976001 (2006). [5] H. Kimura, J Phys Soc. Jpn 76, 074706 (2006).

#### PA11

## Strain-induced ferroelectric instabilities in the epitaxial $RMn_2O_5\,(R{=}Dy$ and Tb) thin films

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Epitaxial thin films of DyMn<sub>2</sub>O<sub>5</sub> (DMO) and TbMn<sub>2</sub>O<sub>5</sub> (TMO) were successfully grown on Nb-doped TiO<sub>2</sub> (110) substrates by the pulsed laser deposition to investigate the effects of substrate-induced strains on multiferroic properties. Synchrotron X-ray azimuthal angle scans and reciprocal space mappings confirmed the epitaxial qualities of the films. Magnetic phase transitions of the films were observed at  $T_n = 43$  K, showing consistent values with those of the bulk samples. Dielectric constant showed a steplike anomaly around 16 K, where a magneto-capacitance effect about 10 % at 8 T was detected. Magnetization-induced ferroelectric phases below Tn appear to become unstable in the epitaxial thin films due to the substrate-induced tensile strains which make the magnetic exchange interactions between nearby manganese ions weaker.

#### PA12

#### Magnetically driven ferroelectric atomic displacements in perovskitetype YMnO<sub>3</sub> determined by single-crystal structure analysis

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Multiferroic materials have been intensively studied with a particular focus on the driving mechanism of cross-correlated interactions between the magnetism and electricity [11]. Many theoretical and experimental works have clarified the origin of the interplay between magnetism and electricity. However, the lack of information about the crystal structure in the multiferroic phase has prevented quantitative discussion of the ferroelectricity driven by magnetic order. In order to quantitatively discuss the microscopic origin of the ferroelectricity of multiferroic phase has prevented quantitative discussion of the ferroelectricity driven by magnetic order. In order to quantitatively discuss the microscopic origin of the ferroelectric polarization below 30 K [2], we performed the crystal structural analysis for a twin-free single crystal by using synchrotron x-ray diffractometer at BL/2021 in SPring-8, and successfully determined the ferroelectric displacements of the order of 0.001 Angstrom. The refined polar structure shows the characteristic bond alternation, which is consistent with that expected for the exchange striction mechanism giving rise to a ferroelectric polarization along the a-axis. From the experimental results of x-ray diffraction and a theoretical firstprinciples calculation, we quantitatively discuss the origin of the ferroelectricity in multiferroic YMnO<sub>4</sub>[3].

[1] Y. Tokura and S. Seki, Adv. Mater. 22, 1554 (2010). [2] S. Ishiwata, Y. Tokunaga, Y. Taguchi, and Y. Tokura, J. Am. Chem. Soc. 133, 15818 (2011). [3] D. Okuyama, S. Ishiwata, Y. Takahashi, K. Yamauchi, S. Picozzi, K. Sugimoto, H. Sakai, M. Takaha, R. Shimaon, Y. Taguchi, T. Arima, and Y. Tokura, Phys. Rev. B 84, 054440 (2011).

## POSTER PRESENTATION

#### PA13

Magnetic symmetry and electric polarization on Mn<sub>1-x</sub>Co<sub>x</sub>WO<sub>4</sub> multiferroics Irene Urcelay - Olabarria<sup>1\*</sup>, Eric Ressouche<sup>1</sup>, Jose Luis Garcia - Munoz<sup>2</sup>, Vassil Skumryev<sup>3</sup>, Alexander Mukhin<sup>4</sup> and Juan Manuel Perez - Mato<sup>5</sup> <sup>1</sup>Institut Laue Langevin, 38042 Grenoble, Cedex 9, France

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MnWO<sub>4</sub> exhibits magnetoelectric effects and belongs to the new class of frustrated multiferroics in which magnetic and ferroelectric order coexist [1-3]. It undergoes three successive magnetic phase transitions [1,2]. Below  $T_a = 13.5$  K the spins are collinear and sinusoidally modulated, (AF3, paraelectric). In the AF2 phase (7.5 K<T<12.5 K) a spontaneous polarization along b axis coexists with a cycloidal spin structure, k = (40.214, 1/2, 0.457). Below 7.5 K the system is in the AF1 collinear commensurate phase,  $k = (\pm 1/4, 1/2, 1/2)$ . Addition of Co+2 ions (S=3/2) modifies the overall magnetic interactions, produces changes in the magnetic anisotropy at 3d- site and stabilizes the multiferroic behaviour [3-5]. We have studied the temperature induced phase transitions of Mn<sub>1</sub>,  $C_0_xWO_4$  (x = 0, 0.1) multiferroics by single crystal neutron diffraction experiments, magnetic and polarization measurements, and we have analyzed the results in the light of the symmetry. All the magnetic models have been refined with JANA2006 [6] using the superspace formalism which allows predicting the crystal tensor properties, such as the electric polarization. We have observed that compared to pure MnWO4, the magnetic symmetry of the multiferroic phase changes in agreement with the experimentally observed flip of the electric polarization.

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#### PA14

uly 9 (Mon)

#### Crystal and magnetic structure of multiferroic Ba<sub>2</sub>CoGe<sub>2</sub>O<sub>7</sub>

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The crystal structure of Ba<sub>2</sub>CoGe<sub>2</sub>O<sub>7</sub> has been precisely determined by single crystal diffraction at four different temperatures. Room temperature as well as 90 K structure determinations were performed using x-ray synchrotron radiation. The observation of the superstructure reflections commensurate with the tetragonal cell and violating the 2-1 axis lead to the conclusion that P-421m is a rather good approximation of the average structure, but the real structure has lower symmetry and most probably corresponds to the Cmm2 SG [1]. Crystal structure above (10 K) and crystal and magnetic structures below (2 K) the antiferromagnetic transition temperature of 6.7 K have been determined using neutron diffraction on single crystal. The structural results show no significant changes with the temperature indicating no structural phase transition down to the multiferroic phase. Antiferromagnetic non-collinear magnetic order has been found in the ground state contrary to the previous literature data [2]. The canting angle and the magnitude of the ordered moment have been refined. The possibility of the existing of polar structure and week ferromagnetic component in the (a,b) plane shed a new light in the understanding of the multiferroicity in the Ba<sub>2</sub>CoGe<sub>2</sub>O<sub>7</sub>.

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#### PA15

## The role of Co-doping in $Mn_{0.85}Co_{0.15}WO_4$ studied by magnetic X-ray scattering

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MnWO<sub>4</sub> is known to show a multiferroic (MF) behaviour at low temperatures (~8-13 K) where the development of an incommensurate (ICM) Mn cycloidal spin structure (AF2 phase) is responsible for the appearance of a net electric polarization along the crystal b-axis [1-3]. The partial substitution of Mn by Co in small quantities stabilizes AF2 down to the ground state. A further Co doping produces the appearance of a new ICM phase, defined as AF2<sup>'</sup>[4,5]. We have performed resonant magnetic x-ray scattering (RMXS) with soft x-rays (Mn, Co  $2p \rightarrow 30$ ) on  $Mn_{685}Co_{113}WO_4$  to independently analyze the spins orientation in Mn and Co ions in antiferromagnetically (AFM) ordered phases. A full polarization analysis of the incident and scattered x-rays permits us to provide an accurate complete picture of the magnetic structure three the latter points to a strong influence of Co doping on the Mn spins alignment direction. We discuss these results in the frame of a large Co single ion anisotropy as reported in CoWO<sub>4</sub> [6].

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#### PA16

Magnetic and electronic properties of hexagonal  $RMnO_3$  (R = Y, Tb) quantum-wells

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#### PA17

**Magnetic properties of the pyroxenes** Sergey Streltsov<sup>1</sup> and Daniel Khomskii<sup>2</sup> <sup>1</sup> Institute of metal Physics, Russia <sup>2</sup> University of Cologne, Germany

Pyroxenes, silicates of composition AMSi<sub>2</sub>O<sub>6</sub>, comprise one of the most important classes of minerals both in the Earth's crust and upper mantle, and as typical (extra) terrestrial rock-forming materials (A=Li,Na,Ca, and M is a transition metal ion). Interestingly enough these materials exhibit very different magnetic properties. Some of them have a spin-gapped ground state, like NaTiSi<sub>2</sub>O<sub>6</sub>, others show AFM or FM (which is unusual for transition metal insulator) long range magnetic order, Fe-based pyroxenes were found to be multiferroic and possibly have the spin-state transition under a high pressure. In the present talk we will review magnetic properties of the pyroxenes and provide a theoretical description for some of their properties. This work was supported by grants RFBR 10-02-00046 and 10-02-96011.

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#### PA18

#### High-field study of multiferroic $Ni_3V_2O_8$

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#### PA19

## Control of coexisting ferroelectric phases in RMnO<sub>3</sub> crystals with fine tuning of 4f moment

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Recently, a new class of multiferroics with a coexistence of magnetic and ferroelectric ordering has been attracting revived interests. RMnO<sub>3</sub> (R: rare earth) with distorted orthorhombic structure is one of the most canonical examples for such a multiferroic material. In TbMnO<sub>3</sub> [1], the electric polarization lies along the c axis in the ferroelectric ground state, and the polarization does not flop without a magnetic field. Meanwhile, in Eu<sub>0.55</sub>Y<sub>0.465</sub>MnO<sub>3</sub> without 4f moment [2], the magnetoelectric properties are quite different from those of TbMnO<sub>3</sub>, although they have the same ionic radius of R-site. This result suggests that the magnetoelectric properties of the RMnO<sub>3</sub> strongly depend on the characteristics of the R-site 4f-moments [3]. In this study, we have systematically investigated magnetic and dielectric properties of Eu<sub>0.595</sub>(Y<sub>1.4HO</sub>)<sub>0.462</sub>MnO<sub>3</sub> (0 $\leq$ xc1) single crystals with control of 4f moments while keeping the average size of the R-site. In the x = 0.3 crystal, the ferroelectric polarization along the c axis (Pc) emerges at 25K, and then Pc is suppressed and Pa appears simultaneously with decreasing temperature. On further decreasing temperature, Pa disappears and Pc appears again. The observed reentrant behavior of the Pc phases indicates that Pa and Pc phases are strongly competing with each other.

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 M. Hitomi et al., J. Phys. Soc. Jpn. 79, 054704 (2010).

#### PA20

#### SmCr<sub>3</sub>(BO<sub>3</sub>)<sub>4</sub> - A new multiferroic?

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Rare-earth borates with general formula  $RM_3(BO_3)_4$  (R = Y, La-Lu; M = Al, Fe, Cr) are interesting for science and technology. Iron borates have rich magnetic properties stipulated by the presence of two interacting magnetic subsystems. It has recently been shown that  $RFe_3(BO_3)_4$ belong to a new class of multiferroics [1,2]. For this reason it is interesting to study rare-earth borates with other magnetic ion chromium. At present, little is known about these compounds. Antiferromagnetic phase transition has been found in NdCr<sub>3</sub>(BO<sub>3</sub>)<sub>4</sub> at  $T_{3/7}$ =8.0 K [3]. This work describes an optical study in a wide spectral and temperature range of the SmCr<sub>3</sub>(BO<sub>3</sub>)<sub>4</sub> undergoes a cascade of phase transitions at temperatures of 8.0 K, 6.7 K and 4.3 K. The first two ones are second-order phase transitions and, supposedly, correspond to the antiferromagnetic and (anti) ferroelectric ordering, respectively. First-order phase transition at 4.3 K is, most probably, the spin reorientation of the chromium magnetic moments. The presence of the magnetic and ferroelectric order allows us to talk about samarium chromium borate as a new multiferroic crystal.

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#### PA21

#### Magnetodielectric effect in the antiferromagnet $SrNdFeO_4$

Jungmin Hwang<sup>1</sup>, Eun Sang Choi<sup>2</sup>, Haidong Zhou<sup>2</sup>, Yan Xin<sup>2</sup>, Jun Lu<sup>2</sup> and Pedro Schlottmann<sup>1</sup> <sup>1</sup> Florida State University/NHMFL, USA <sup>2</sup> NHMFL (National High Magnetic Field Laboratory), USA

We investigated the magnetic phase diagram of single crystals of the compound SrNdFeO<sub>4</sub> by measuring the magnetic susceptibility, the magnetization, the specific heat and the dielectric constant. The system has two magnetically active ions, namely Fe<sup>3+</sup> and Nd<sup>3+</sup>. The Fe<sup>3+</sup> spins are antiferromagnetically ordered below 360 K with the moments lying in the ab-plane, and undergo a reorientation transition at about 35-37 K to an antiferromagnetic order with the moments along the c-axis. A short-range, antiferromagnetic ordering of Nd<sup>3+</sup> along the c-axis was attributed to the reorientation of Fe<sup>3+</sup> followed by a long-range ordering at lower temperature.[1] At low temperatures and magnetic fields above 8 T the Nd<sup>3+</sup> moments are completely spin-polarized. The dielectric constant also shows anomalies associated with spin configuration changes. Remarkably, the sign of the dielectric changes are dependent on the direction of the probing electric field. A possible magnetic structure explaining all the data is presented.

 S. Oyama, M. Wakeshima, Y. Hinatsu, and K. Ohoyama, J. Phys.: Condens. Matter. vol. 16, 1823 (2004).

#### PA22

#### Mössbauer studies of Y-type hexaferrite by Aluminum doping

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In recent years, the magnetoelectric effect observed in some of helicalmagnets has been of great interest. With Al substitution, the improvement of magnetoelectric effect has been reported due to reduction of the in-plane orbital moment originating from lowering magnetic anisotropy[1-2]. Therefore, it is important to understand the origin of magnetic properties in conjunction with the site distribution in hexagonal structure. The crystal structure of polycrystalline  $Ba_2Co_{1.5}Mg_{0.5}(Fe_{1.x}Al_x)_{12}O_{22}(x = 0, 0.01, 0.02)$ samples were determined to be rhombohedral with the space group of R-3m. From the temperature-dependent of magnetization curves under 100 and 300 Oe, all samples showed ferrimagnet-to-paramagnet and helimagnet-to-ferrimagnet transitions. Base on the applied-field dependent magnetization measurements up to 10 kOe at 295 K, the saturation magnetization (Ms) and coercivity (Hc) of  $Ba_2Co_{1.5}Mg_{0.5}(Fe_{1.x}Al_x)_{12}O_{22}(x =$ 0, 0.01, 0.02) samples were found to be Ms = 28.0, 24.7, 23.4 emu/g and Hc = 255.4, 222.0, 196.2 Oe, respectively. Also, from Mössbauer spectra of 295 K, Isomer shift values of all samples indicate that the charge states are Fe<sup>3+</sup>. We expect that the decrease in the Ms with Al ion doping is due to the fact that Al ions preferentially occupy the 3bVI, 18hVI and 3aVI octahedral sublattices of up-spin sites[3].

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#### **PA23**

## Optical spectroscopy of the triangular lattice antiferromagnets CuCrO<sub>2</sub> and alpha-CaCr<sub>2</sub>O<sub>4</sub>

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We will compare and discuss our results obtained by optical spectroscopy on CuCrO<sub>2</sub> and alpha-CaCr<sub>2</sub>O<sub>4</sub>. While CuCrO<sub>2</sub> is famous for its multiferroicity [1], in alpha-CaCr<sub>2</sub>O<sub>4</sub> a polarization can only be observed under the application of electric or magnetic field, despite having a closely related structure [2]. In alpha-CaCr<sub>2</sub>O<sub>4</sub> we can see a strong, broad absorption at 57 cm<sup>-1</sup> in the magnetically ordered state. The other compound shows two excitations at 11 and 48 cm<sup>-1</sup> below the magnetic ordering temperature. At near infrared and visible light frequencies we observe Cr<sup>3+</sup> crystal field absorptions and below  $T_{iv}$  excitons and exciton-magnon-transitions appear. The width of these exciton-magnon transitions is analyzed with respect to the existence of Z2 vortices as proposed by Kojima et al. [3].

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#### PA24

## Nuclear forward scattering in high magnetic fields; spin structures in the magnetic staircase of frustrated multiferroic CuFeO<sub>2</sub>

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Geometrically fustrated magnets often show a multitude of nearly degenerate ground states, which may be stabilized individually by the presence of small internal or external perturbations. In Copper Iron Oxide, the combination of geometrically fustrated interactions, with uniaxial anisotropy and spin-lattice coupling lead to a staircase of metamagnetic phases as a function of applied field [1-3]. We used Nuclear Forward Scattering of synchrotron radiation in magnetic fields up to 30 T [4, 5] to probe the spin structure of these phases up to the 1/3-plateau for fields parallel and perpendicular to the c axis [6]. The low field data and all transition fields agree with previous work. For the parallel configuration, we confirm the low field calinear spin structure. The data for the ferroelectric phase does not support a cycloid structure, while a 3D distribution approximating multi-domain transverse proper helix structure gives qualitative agreement. For the 1/3-plateau, the anticipated collinear 3-sublattice structure base trepresents the data. In the perpendicular configuration we directly evidence the stability of the canted 4-sublattice phase up to the 1/3-plateau, where a collinear three sublattice configuration provides the best fit when a small amount of the spins oriented in the H–c plane is allowed for.

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#### PA25

#### Magnetic dispersion of the quasi-1D, spin-1/2, multiferroic CuO Stephen Michael Gaw

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Cupric Oxide (CuO) exhibits a magnetoelectric multiferroic phase over the narrow but relatively high temperature range of 213-230K [1]. The mechanism for this magnetoelectric coupling is not fully understood and there have been various theoretical models proposed [2]. The excitation spectra of CuO in the collinear antiferromagnetic phase (T<213K) was measured using inelastic neutron scattering. We have observed three-dimensional magnetic behaviour transform into a spin-1/2 chain regime with increasing energy. Through comparison of the data to a semiclassical spinwave model and a quantum spinon model, we have extracted values of the magnetic exchange constants. With these values we can refine the existing theoretical models of the magnetoelectric mechanism as well as unravel the nature of one-dimension quantum magnetism in CuO.

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## Iuly 9 (Mon)

#### **PB01**

#### Hole-doped cuprate panorama and the second neighbor hopping Partha Goswami\*

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We consider the coexistent id-density wave (DDW) order, at the anti-ferromagnetic wave vector  $Q = (\pi, \pi)$ , representing the pseudo-gap (PG) state, and d-wave superconductivity (DSC), driven by an assumed attractive interaction, within the BCS framework for the two-dimensional (2D) fermion system on a square lattice starting with a mean-field Hamiltonian involving the singlet DDW and the DSC pairings. The second-neighbor hopping, which is known to be important for cuprates [K. Tanaka et al., Phys. Rev. B 70, 092503(2004)]and frustrates the kinetic energy of electrons, leads to Fermi surface sheets being not connected by Q. The particle-hole asymmetry in the single-particle excitation spectrum of the pure DDW state is reflected in the coexisting DDW and DSC states though the latter is characterized by the Bogoluibov quasiparticle bands- a characteristic feature of SC state. Quite significantly, we find that the coexistence is possible due to the non-nesting property.

#### **PB02**

Quantum oscillations from nodal bilayer magnetic breakdown in the underdoped high temperature superconductor YBa2Cu3O6+x Neil Harrison

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We report quantum oscillations in underdoped YBa2Cu3O6.56 in magnetic fields over a significantly large range in magnetic field extending from ≈-25 to 94~T, enabling three well-spaced low frequencies at ≈-462~T, 532~T, and 602~T to be clearly resolved. We show that a small nodal bilayer coupling that splits a nodal pocket into bonding and antibonding orbits yields a sequence of frequencies, F\_0- $\Delta$  F, \$F\_0\$ and F\_0+ $\Delta$ F and accompanying beat pattern similar to that observed experimentally on invoking magnetic breakdown tunneling at the nodes. The relative amplitudes of the multiple frequencies observed experimentally in quantum oscillation measurements are shown to be reproduced using a value of nodal bilayer gap quantitatively consistent with that measured in photoemission experiments in the underdoped regime. A clear consistency between the experimental data and a Fermi surface consisting of nodal pockets is revealed

#### **PB03**

#### Influence of BaSnO<sub>3</sub> nanoparticle dispersions on flux pinning property of GdBa2Cu3O7.8 multilayered thin films

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Flux pinning property of GdBa<sub>2</sub>Cu<sub>2</sub>O<sub>7.8</sub> (GdBCO) thin films with embedding BaSnO<sub>2</sub> (BSO) nanoparticle pseudo-layers structured as (BSO/GdBCO) x N layers was investigated The multilavered films whose entire thickness of approximately 200 nm were fabricated on SrTiO<sub>3</sub> (STO) substrate by using the pulsed laser deposition (PLD) technique. Number of BSO pseudo-layers was varied from 1 to 5 in order to increase the density of pinning centers along the c-axis of the films BSO nanonarticle density per layer was estimated to be approximate 114/um2 by using an atomic force measurement (AFM). A minor decrease in the transition temperature Tc was observed for (BSO/GdBCO) multilayered films. Magnetization data showed that BSO pseudo-lavers provides enhanced in-field critical current density Jc up to certain magnetic fields. Jc enhancements were attributed to the dispersions of BSO nanoparticle pseudo-layers confirmed by cross sectional TEM image.

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#### **PB04**

#### Highly anisotropic dielectric behavior of insulating Bi2Sr2RECu2O8+6 (RE = Dy, Y)

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The low-frequency dielectric measurements on lightly-doped Bi<sub>2</sub>Sr<sub>2</sub>RECu<sub>2</sub>O<sub>848</sub> (RE = Dy, Y) single crystals were carried out in order to investigate the electronic structure and carrier dynamics near the insulator-to-metal transition. We found the unusually large dielectric polarization in the ab plane, whereas the out-of-plane dielectric constant was reasonably small at low temperatures. With increasing temperature above about 230 K, the out-of-plane dielectric constant shows a rapid increase followed by a sharp peak. The large anisotropy and unusual temperature dependence of the dielectric response probably reflect the hopping polarization induced by localized charge carriers and the spatial charge distribution peculiar to the insulating Bi2212. Our results seem to imply that the carrier movement is limited within the CuO<sub>2</sub> planes at low temperatures below about 230 K, and the confined carriers start to hop toward the perpendicular direction at higher temperatures [1, 2]. We also present data for Zndoped Bi2Sr2RECu2O8+8 samples to discuss the relationship between the enhanced dielectric polarization and a long-range antiferromagnetic ordering of Cu magnetic moments in the CuO<sub>2</sub> planes.

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#### **PB05**

#### Pressure effect on the superconductivity and crystal structure for Hg cuprate

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18 years ago, Hg cuprate recorded the highest Tc=164 K under high pressure (31 GPa) in 1994 [1]. But the record has not broken yet. We try to make new world record of  $T_c$  with the high purity  $Hg_{1223}$  (HgBa<sub>2</sub>Ca<sub>2</sub>Cu<sub>3</sub>O<sub>8+ $\delta$ </sub>) samples and the high purity  $Hg_{1234}$ (HgBa2Ca3Cu4O10+6) sample. We measured electrical resistance and X-ray diffraction under the pressure with a diamond-anvil cell (DAC). We measured two different Hg<sub>1223</sub> samples, which were under doped sample (Hg1223 UD) and optimal doped sample (Hg1223 OPT), and one Hg1234 sample which was optimal doped sample. We measured two different pressure conditions, using pressure medium (NaCl) or not (None). The resistance drop of superconducting transition becomes broader by increasing pressure. At 80 GPa, we can not find the drop of superconducting transition above 10 K. The T<sub>c</sub> have a maximum at 20-30 GPa accompanied with the anomaly in crystal structure.

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#### **PB06**

#### Hole doping effect for the T'-Ln<sub>2</sub>CuO<sub>4</sub> (Ln = La,Nd) cuprate Wataru Ito, Kenji Kawashima, Suguru Igarashi, Michinori Fukuma and Jun Akimitsu Physics and Mathematics, Aoyama Gakuin University, Japan

The Ln<sub>2</sub>CuO<sub>4</sub> (Ln=La,Nd) cuprate has two different structure types and shows the superconductivity for hole doping of La2CuO4 and electron doping of Nd2CuO4. From the Madelung energy calculation in hole doping cuprate,  $T_c$  is enhanced with increasing the distance between Cu2+ ions in CuO2-layer and apical oxygens. Nd2CuO4 with T'type structure has exhibited superconductivity by electron carrier doping in CuO2layer. We focused on our attention to synthesize the hole doping to T'-Ln<sub>2</sub>CuO<sub>4</sub>, in which there is no oxygen atom at apical site of CuO2 layer. We successfully synthesized the polycrystalline samples of T'-(Ln3+,M2+)2CuO4 (Ln=La,Nd, M=Ca,Sr). In particular, T'-(La<sup>3+</sup>,Sr<sup>2+</sup>)<sub>2</sub>CuO<sub>4</sub> was synthesized using CaH<sub>2</sub>, which is a strong reducing agent at annealing temperature of 573 K [1]. The powder X-ray diffraction patterns showed that the samples were a single phase, and diffraction peaks indexed by a tetragonal symmetry of T'-type cuprate of Nd2CuO4. The lattice constants a and c show the systematic changes by increasing the M2+ concentration. We determined the copper ion's valence state using iodometry method for substituted sample of T'-(Ln3+,M2+)2CuO4 and confirmed that the valence of Cu ion is closed to 2+. These facts indicate that there is oxygen deficiency in Ln2O2 or CuO2-layer of T'-(Ln3+,M2+)2CuO4 and canceled hole carriers.

T. Takamavsu, M. Kato, T. Noii, and Y. Koike, Physica C 471 (2011) 679-681

#### **PB07**

#### Thermodynamic properties of copper oxide Cu<sub>6</sub>O<sub>8</sub>YCl<sub>1-x</sub>Br<sub>x</sub>

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Copper-oxide family Cu<sub>6</sub>O<sub>8</sub>MCl (M=cation) has a cubic structure (Fm-3m) with Cu<sub>6</sub>O<sub>8</sub>-cage including Cl ion in its crystal structure [1,2]. Cations of M ions are located in the cuboid space between Cu<sub>6</sub>O<sub>8</sub>-cages. The Cu<sub>6</sub>O<sub>8</sub>-cages share their face to form a three-dimensional network. Cu<sub>6</sub>O<sub>8</sub>YCl shows the metallic and paramagnetic behavior at low temperature. Moreover, Y and Cl ions have large atomic displacement parameters, suggesting the anharmonic vibrations like rattling phenomena of other cage structured compounds. We synthesized the polycrystalline samples of Cu<sub>6</sub>O<sub>8</sub>YCl<sub>1</sub> Br, as a single phase to clarify relationship between the thermodynamic properties including its ground state and Cu6O8 cage size. In thermo-electronic measurements of Cu<sub>6</sub>O<sub>8</sub>YCl, seebeck coefficient is relatively high value (-14 µV/K at 350K), and thermal conductivity exhibits of low value (~0.3W/Km at 350K). The lattice constant a increases with increasing Br concentration, indicating expansion of Cu6O8-cage. The temperature dependence of electrical resistivity of Cu<sub>6</sub>O<sub>8</sub>YBr exhibits semiconducting behavior, suggesting that the ground state of Cu<sub>6</sub>O<sub>8</sub>YCl<sub>1-x</sub>Br<sub>x</sub> changes from metallic to semiconducting by increasing Br concentration. Thermal conductivity undergoes very little change. However, S is enhanced by substituting Br. These facts indicate that the ground state and physical properties are sensitive on the Cu<sub>6</sub>O<sub>8</sub>-cage size in Cu<sub>6</sub>O<sub>8</sub>YCl<sub>1-x</sub>Br<sub>x</sub> system.

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#### PB08

#### Nodal superconducting gap in Bi<sub>2201</sub> investigated by low temperature specific heat measurements

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In order to examine the gap magnitude near nodes in Bi2201 cuprate superconductors, we have measured the magnetic field dependence of low temperature electronic specific heat Cel(T, H), which is much more bulk-sensitive than standard ARPES and STS. In optimally doped Bi2201, the coefficient of T-linear term in Cel at T<<Tc, which is almost vanished in zero field, shows a square-root dependence of magnetic field H as is expected in d-wave superconductors. We determined the nodal gap slope v, reflecting the gap magnitude near nodes, from the square-root dependence of H; v~4.4x10^5 cm/ s. In heavily underdoped Bi2201, the coefficient of T-linear term is not zero even in zero field, and is nearly independent of H up to 6 T. The comparison of these results with

recent ARPES and STS results will be discussed in detail.

#### **PB09**

#### Chemically introduced disorder effects on the critical current density and pinning force of YBaCu<sub>2</sub>O<sub>7</sub>-δ single crystals

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We report isotherm dc magnetization hysteresis loops of a series of  $YBa_{2x}SrxCu_3O_7$ - $\delta$  (x = 0, 0.1, 0.25, 0.37 and 0.5) single crystals with the aim to study the influence of the chemically introduced site disorder on the YBa<sub>2</sub>Cu<sub>3</sub>O<sub>7</sub>-δ critical current density, JC and flux pinning properties, FP. [1,2] The Sr ion size chemical disorder is introduced in YBa2Cu3O7 & structure by lattice distortion. Transmission electronic microscopy (TEM) observations of our samples structure revels the existence of a high density of twins, probably decorated by many small local precipitates. The JC were determined by application of the modified Bean critical state model.[3-5] Preliminary magnetization data analysis shows a significant and a systematically decrease of the critical current density and pinning force intensities as Sr concentration level increases up to 25%. The normalized FP curves showed no reasonable scaling as a function of the Sr doping level or temperature range.[4,5] The application of the Dew Hughes law to the normalized FP curves revels that pinning mechanism in the doped samples are strongly temperature and level doping dependent. We consider the particular defect structure introduced by Sr substitution as a responsible to the flux pinning properties.

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#### **PB10**

#### The correlation between the magnetic irreversibility and the zero resistance temperatures in granular YBa<sub>2</sub>Cu<sub>3</sub>O<sub>7.8</sub> single crystals

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We report on dc magnetization and ac magnetoresistance measurements of the Y009Ca001 Ba2Cu3O7.5. YBa198Sr002Cu3O75 and YBa2Cu297Zn003O75 single crystals with the propose of to study the role of the superconducting granularity on the correlation between the magnetic irreversibility, Tirr(H) and the zero resistance, Tr(H) temperatures.[1-3] The results show that in the low field region of the H-T diagram that the Tr(H) data falls systematically underneath of the Tirr(H) data and at H-T high field region the Tr(H) data matches the Tirr(H) data. In this scenario, Tirr(H) and Tr(H) do not depend on the same parts of the sample. While the Tirr(H) depends on wellcoupled grain clusters, the Tr(H) depends on grain arrays traversing the whole sample. Along such long range paths, the Tr(H) can be attained only at some temperature below the Tirr(H). On the other hand, in fields above several kOe, for which the magnetic field penetrates the grains, the Tirr(H) is dominated by the intragrain Abrikosov flux dynamics while the Tr(H) is still ruled by the grain junctions. It can vanish only after long-range coherence is achieved. The granular aspects of our results are discussed in terms of the superconducting glass theories.[1-4]

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#### **PB11**

#### Paramagnetic Meissner effect and strong time dependence at high fields in melt-textured high-tc superconductors

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The conventional Meissner effect is characterized by a diamagnetic response of the superconducting material when a magnetic field is applied, but in several cases this magnetic response in field-cooled experiments can be paramagnetic. This effect is known as paramagnetic Meissner effect (PME) [1]. Sometimes the PME presents a time dependence of the FC magnetization and a paramagnetic signal can be observed when the field-cooled moment relaxes at constant magnetic field and temperature [2,3]. In this work we report on systematic field-cooled magnetization experiments in melt-textured YBaCuO samples containing Y211 precipitates. Magnetic fields up to 14 T were applied either parallel or perpendicular to the ab planes and a strong paramagnetic response related to the superconducting state was observed. The magnitude of the PME increases when the field is augmented. This effect shows a strong paramagnetic relaxation, such that the paramagnetic moment increases as a function of the time. The pinning by the Y211 particles plays a crucial role in the explanation of this effect and our results suggest that the pinning capacity [4] can produce a strong flux compression [5] into the sample, originating the PME and the strong time dependence.

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#### PB12

#### Enhancement of Ti-Se bonding length in CuxTiSe<sub>2</sub>

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Titanium diselenide (TiSe<sub>2</sub>) undergoes a charge-density-wave (CDW) instability characterized by a 2x2x2 real-space superstructure [1]. TiSe<sub>2</sub> is not superconducting at low temperature, but CDW is suppressed and superconductivity stabilized either by Cu intercalation (maximum Tc =4.5 K) [1] or pressure (maximum Tc = 1.8 K) [2]. The resulting phase diagram looks similar to that of cuprates with the difference that the spin-density wave order has been replaced by a CDW one. The mechanism at the origin of the CDW and superconducting phases in TiSe<sub>2</sub> is currently unknown. The charge-order instability in pure TiSe<sub>2</sub> is known to involve both exciton formation and a strong electron-phonon coupling, it may be expected that the superconducting phase will likewise involve a combination of both excitonic [3] and phonoic [4] contributions. In order to elucidate the role of Cu atoms, we have investigated the Cu-doped TiSe<sub>2</sub> by employing both x-ray absorption near-edge structure (XANES) and extended x-ray absorption fine-structure (EXAFS). The Cu atoms enhance the bonding length of Ti and Se atoms with lowening temperature. Concurrently, the Debye-Waller factors change at the transition temperatures of CDW and superconductivity, respectively.

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#### PB13

## Para-conductivity of $(Bi_{0.25}Cu_{0.25}Li_{0.25}Tl_{0.25})Ba_2Ca_2Cu_3O_{10-\delta}$ superconductors

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Excess conductivity analyses of as-prepared and oxygen post-annealed ( $B_{10,2}Cu_{0,2}Li_{0,2}Tl_{0,2}$ )  $Ba_2Ca_2 Cu_3 O_{10,5}$  samples were carried out by following Lorentz-Danish (LD) model beyond the transition region and Ginzburg-Landau (GL) equations in the critical regime. In these studies the effect of oxygen annealing on the thermodynamic fluctuations of the Cooper pairs just above the Tc is performed by using Aslamazov-Larkin (AL) theory. We have observed a major increase in the width of 3D AL and 2D LD regimes after post-annealing in oxygen atmosphere. Moreover, with the oxygen post-annealing the coherence length along the c-axis has also been found to increase along with inter-plane coupling. The thermodynamic critical magnetic field, the ouper critical magnetic field, the critical current density and penetration depth are also calculated from the analysis of four regions namely critical (cr), three dimensional (3D), two dimensional (2D) and zero dimensional (0D) region that appeared as the temperature increase. The values of Fermi velocity, phase relaxation time, penetration depth, x values and energy required to break apart the cooper pairs are increased with the annealing of oxygen. This shows that in  $(B_{10,2}Cu_{0,2}T_{10,2})Ba_2(a_2Cu_0,D_{10,4})$  sample, the optimum density of carriers in the conducting CuO2 is essential for enhanced superconductivity in these compounds.

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#### **PB14**

## Charge transfer instability governs unconventional behavior of doped cuprates

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A large body of experimental data points towards an unique charge transfer (CT) instability of parent insulating cuprates. True CT gap in these compounds is believed to be as small as 0.4-0.5 eV as derived from the midinfrared absorption measurements [1] rather than 1.5-2.0 eV as usually derived from the fundamental absorption measurements. Actually we deal with a competition of the conventional (3d9) ground state and a CT state with formation of electron-hole dimers which evolves under doping to an unconventional bosonic system [1-3]. Making use of a quantum Monte-Carlo (QMC) technique we study the evolution of the phase state of CuO<sub>2</sub> planes in a model CT unstable cuprate La<sub>2,x</sub>Sr<sub>x</sub>CuO<sub>4</sub>. Nonisovalent doping gives rise to a suppression of parent antiferromagnetic phase with the nucleation of the inhomogeneous supersolid (OB) order parameters which competition results in a generic T-x phase diagram with a distinct pseudogap regime due to a charge ordering and emergence of a local superconductivity. Our QMC simulation does reproduce all the main features of the T-x phase diagrams for doped cuprates.

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#### PB15

## The effect of CdO nanoparticles doping and sintering time on the structure and critical temperature of Bi<sub>2223</sub> superconductor

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The effect of CdO nanoparticles doping and sintering time on the structure and critical temperature of Bi<sub>2223</sub> superconductor M. Zargar shoushtari1, S. E. Musavi Ghahfarokhi, N. Hossinzadeh 1 Department of physics, Shahid Chamran University of Ahvaz, Ahvaz, I. R. Iran Presenting author: zargar@scu.ac.ir In this paper, Bi<sub>1.64</sub>Pb<sub>0.36</sub>Sr<sub>2</sub>Ca<sub>2</sub>,  $\chi$ Cd,  $\chi$ Cu<sub>3</sub>O<sub>y</sub> (x= 0.0, 0.01, 0.02, 0.03, 0.04, 0.1) superconductor is prepared by using the solid state method. The effect of doping CdO nanoparticles for different sintering times (t= 90, 180 and 270 h) on the structure and critical temperature has been studied. The structural analysis was carried out using XRD and SEM measurements. The critical temperature was measured by the standard four-probe method. The results showed that the more sintering time increases the critical temperature such that the samples prepare at 270h have maximum critical temperature. Also, our results reveal that, the Bi<sub>1.64</sub>Pb <sub>0.36</sub>Sr<sub>2</sub>Ca<sub>1.97</sub>Cd<sub>0.05</sub>Cu<sub>3</sub>O<sub>y</sub> sample has the highest Tc (Tc= 137.5k) in between the doped

#### PB16

#### Charging/Discharging and overcurrent characteristics of GdBCO coils using various partial insulation winding methods

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Eliminating the insulation in high temperature superconducting (HTS) windings can allow a system to be compact and well protected by bypassing the excessive current flow that occurs through the turn-to-turn contact between uninsulated turns. However, the time constant of coils without insulation is much higher than that of completely insulated coils due to the absence of insulation resistance, resulting in a coil chargedischarge rate that becomes quite slow. Therefore, this study examined the charging, operating, and discharging behaviors of Gd-Ba-Cu-O (GdBCO) HTS coils with Kapton insulation every 3, 6, or 9 turns to investigate the effect of changes in insulated turns on the coil's time constant. Furthermore, to verify the demagnetization phenomena of the coils in the excessive current condition, the overcurrent characteristics of partially insulated GdBCO coils were evaluated and discussed in detail.

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#### PB17

## Effect of filter shape on the capture efficiency of a high gradient magnetic separation (HGMS) system

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High gradient magnetic separation (HGMS) using a ferromagnetic filter enables the capture of fine particles of even weakly magnetic materials by the high gradient of the magnetic field. In general, such a field gradient, which can determine the capture capability of the HGMS system, is mainly affected by the filter's mesh size, wire diameters, and materials. Therefore, this study examined the capture efficiency of the HGMS system with respect to mesh size and wire diameter to investigate the effect of filter shape on gradient magnitude. Furthermore, the capture efficiency of various filter shapes was also examined and compared with the simulation.

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#### PB18

## Effects of various epoxy impregnations on the electrical properties of GdBCO-coated conductor racetrack pancake coils

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A high temperature superconducting (HTS) field coil in a wind turbine generator is generally impregnated with epoxy resin to provide high mechanical integrity against the rotational vibrations. However, if cooled down, the epoxy-impregnated coil wound with second generation (2G) HTS tape usually encounters superconducting property degradation and delamination due to the thermal contraction mismatch between the epoxy and the HTS tape. Therefore, to investigate the effect of various epoxies on the coil's electrical properties, the GdBCO racetrack pancake coils impregnated with Stycast 2850-FT, Epikote 828R, or Epon 815 were characterized using Ic testing. Furthermore, degradation of the coils' superconducting properties was also examined using repetitive cooling tests.

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#### PB19

#### Crystal structure of (Ru<sub>0.5</sub>Cu<sub>0.5</sub>)(Sr<sub>1.47</sub>Ba<sub>0.2</sub>Nd<sub>0.33</sub>)(NdCe)Cu<sub>2</sub>O<sub>10-δ</sub> compound H.K. Lee<sup>1</sup> and Y.I. Kim<sup>2</sup> <sup>1</sup> Physics, Kangwon National University, Deptartment of Physics, Korea

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#### PB20

## Absence of broken time reversal symmetry below the surface of (110)-oriented YBCO superconductors

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Low-energy muon spin rotation (LE- $\mu$ SR) is a very sensitive local magnetic probe to study magnetism on a nanometer length scale near the surface of materials and thin films. Here, we report the results of a search for spontaneous magnetization due to a time reversal symmetry breaking phase in the superconducting state of (110)-oriented YBCO films, expected to develop near the surface in this orientation. Zero field and weak transverse field measurements performed using LE- $\mu$ SR few nm inside optimallydoped YBCO-(110) films showed no appearance of spontaneous magnetization below the superconducting temperature up to 2.9 K, contrary to tunneling measurements. Our results give an upper limit of 0.01 mT for spontaneous internal fields in these films.

#### PB21

## Transport properties of the twin boundary of $YBa_2Cu_3O_7$ thin films on LaAlO\_3 substrates

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We have fabricated superconducting YBa<sub>2</sub>Cu<sub>3</sub>O<sub>7</sub> (YBCO) twin-boundary junctions and studied their critical transport properties. The twin boundary junctions were fabricated by using a focusd-ion-beam (FIB) etching technique and of the size of nominally 150 nm in width, 100 - 120 nm in length, and 100 nm in thickness. The junctions showed Tc = 91 K with 1 K transition width and a high critical current density of Jc = 15 MA/ cm<sup>2</sup> at 77 K. Measured current-voltage (I-V) curves showed the resistively-shunted-junction (RSJ) characteristics near the transition with excess current associated with flux-flow effect. High-resolution X-ray diffraction (HR-XRD) spectra strongly indicated the possibility that the epitaxially-grown YBCO film was also twinned in commensurate with the twinning of the LAO substrate. Misorientation of the c-axis of YBCO at the twin boundary is believed to be the major cause of the weak link behavior of the bridge across the twin boundary.

#### **PB22**

## Displacement waves of oxygen atoms in the Bi, $Pb_{2223}$ lattice of superconducting composites annealed in an oxygen reduced atmosphere

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The Bi, Pb<sup>2223</sup>/Ag composites annealing in oxygen-nitrogen atmosphere is believed to reduce the number of the accompanying phases, to make contacts between crystallites closer and to increase the critical current. The goals of this study were determining the 2223 lattice changes at annealing in the reduced oxygen atmosphere, revealing the reasons of these changes and their effect on the ceramics superconductivity. After such annealing the transversely-polarized waves of oxygen atoms displacement in [010] direction are observed in the 2223 phase by electron diffraction analysis. These waves may appear due to the lack of oxygen in the 2223 lattice or to the nitrogen penetration in it. As demonstrated by the X-ray photo-electron spectroscopy and nuclear microanalysis, nitrogen does not interact with the 2223 lattice, and the oxygen index decreases to 9.67, which is lower than the stoichiometric. Thus, the atomic displacement waves result from the lack of oxygen in Bi-O bilavers.

#### **PB23**

## Irreversibility line in the CNT and carbon doped YBCO superconductors

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Irreversibility Line in the CNT and Carbon Doped YBCO Superconductors S. Dadras<sup>1\*</sup>, N. Manivannan<sup>2</sup>, V. Daadmehr<sup>1</sup>, and K. H. Kim<sup>2</sup> <sup>1</sup>Magnetic and Superconducting Res. Lab., Department of Physics, Alzahra University, P.C. 1993891176, Tehran, Iran. 2eXtreme Multifunctional Physics Laboratory, PRD, School of Physics and Astronomy, Seoul National University, Seoul 151-747, Republic of Korea. \*dadras@alzahra.ac.ir s\_dadras2001@yahoo.com Abstract. We investigated the irreversibility line (IL) of the carbon nano tube (CNT) and carbon doped YBa<sub>2</sub>Cu<sub>3</sub>O<sub>7.5</sub> superconductor samples using magneto-resistive measurement results. Irreversibility line moves to higher magnetic fields for CNT doped and to lower magnetic field for carbon doped YBCO samples respect to the undoped sample.

#### **PB24**

Nonlocal excitations and 1/8 singularity in cuprates Yoshiro Kakehashi\*, M. Atiqur R. Patoary and Sumal Chandra Department of Physics, University of the Ryukyus, Japan

Low energy excitations and unusual behaviors of electrons in the 2D cuprates have been the question under debates. Especially the non Fermi liquid behaviors and the pseudogap remain unresolved theoretically because of the limited range of the intersite correlations and the limitation of the momentum and energy resolutions in the theory. We present here our results for the momentum-dependent single-particle excitation spectra of the 2D Hubbard model, which are based on the full self-consistent projection operator method. The theory self-consistently takes into account the longrange intersite correlations with use of the off-diagonal effective medium, and allows us to calculate the excitation spectra with high energy and momentum resolutions. We obtained the excitation spectra being consistent with the OMC results and found at T=0 that the nonlocal excitations enhance the Mott-Hubbard gap and create the shadowband excitations due to strong antiferromagnetic correlations. The present theory yields the characteristic hole concentration  $1-n^* = 0.123$  (=1/8) at which the van Hove type singular peak appears just on the Fermi level at T=0, while it is  $1-n^* = 0.15$  at finite temperatures in the 4x4 DCA. The present result suggests the emergence of the 1/8 stripe phase due to the Fermi surface instability.

#### **PB25**

#### Magnetic memory in a ceramic YBCO superconductor composed of sub-micron size grains

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The nonequilibrium nature of a ceramic YBCO composed of sub-micron size grains is investigated by measurements of the dc magnetizations. The ceramic YBCO is considered as random Josephsoncoupled networks of 0 and r junctions and shows successive phase transition. The first transition occurs inside each grain at Tc1=81 K and the second transition occurs among the grains at Tc2=47 K. The magnetic glass behavior similar to those of spin-glasses is observed below Tc2.[1] Theoretically, the frustration effect due to the random distribution of  $\pi$  junctions should lead to the chiral-glass state, as predicted by Kawamura and Li.[2] The memory phenomena are investigated by recording the zero-field-cooled and thermoremanent magnetization temperature dependence measured on heating after the cooling process with and without field switches.[3] There are memory effects of the halt at Ts=44 K imprinted in the system on the re-heating the sample. In the case without field switches at Ts, the influence of the halt is confined to a narrow temperature region near the halt temperature, whereas the memory effects of the halt employing a field switch is extended over a wide temperature region below Ts. The results suggest that chiral-glass ordering occurs at Tc2 in the ceramic YBCO.

[1] H.Deguchi et al., J. Phys. Conference Series 320 (2011) 012076. [2] H.Kawamura and Li, J. Phys. Soc. Jpn. 66 (1997) 2110. [3] E.L.Papadopoulou and P.Nordblad, Eur. Phys. J. B 22 (2001) 187.

#### **PB26**

## Phase diagram of high-tc superconductivity and antiferromagnetism revealed by Cu-NMR in multilayered cuprates

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We present site-selective nuclear magnetic resonance (NMR) studies on multilayered cuprates, which have uncovered the intrinsic phase diagram of antiferromagnetism (AFM) and high-temperature superconductivity (HTSC) for a disorder-free CuO<sub>2</sub> plane with hole carriers. We revealed the existence of the AFM metallic state in doped Mott insulators, the uniformly mixed phase of AFM and HTSC, and the emergence of d-wave SC with a maximum Tc just outside a critical carrier density, at which the AFM moment on a CuO<sub>2</sub> plane disappears. These results can be accounted for by the Mott physics based on the t-J model, where the superexchange interaction J among spins within the plane plays a vital role as a glue for Cooper pairs or mobile spin-singlet pairs.

#### PB27

#### Magnetism and superconductivity in CeCu<sub>2</sub>Ge<sub>2</sub> under high pressures and magnetic fields

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CeCu<sub>2</sub>Ge<sub>2</sub> with the ThCr<sub>5</sub>Si<sub>2</sub> type tetragonal structure is an antiferromagnet with TN = 4.2 K (AF1) and shows two kinds of superconducting phases named a low-Tsc phase (SC1) and a high-Tsc phase (SC2) under high pressures [1]. We have grown CeCu<sub>2</sub>Ge<sub>2</sub> single crystals by Sn-flux method and studied an electronic state under high pressures and magnetic fields by means of the electrical resistivity measurement. We have found two anomalies at TN = 4.1 K and TN<sup>=</sup> 1.3 K at 2.6 GPa for example, implying a Neel temperature and a change of magnetic structures with the first order. In the temperature-pressure phase diagram, two kinds of magnetic phases named AF1 and AF2 appear and merge at 5 GPa. At pressures higher than 5 GPa, the AF2 phase becomes dominant. The AF2 phase can be another antiferromagnetic phase, since the metamagnetic transition to the AF1 phase is observed around 3 T for the H || [001] direction. At 8.2 GPa, pressure-induced superconductivity (SC1) is induced below Tsc(onset) = 0.5 K. The upper critical field is 30 kOe for H || [001]. The AF2 phase.

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#### **PB28**

## Doping and temperature dependence of Fermi arc in cuprate superconductors

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The pseudogap observed in the excitation spectrum as a suppression of spectral weight in the normal-state of cuprate superconductors is thought to be key to understanding the mechanism of superconductivity. Within the framework of the kinetic energy driven superconducting mechanism [1], we [2] study the evolution of the Fermi arc with doping and temperature by considering the interplay between the superconducting gap and normal-state pseudogap [3]. It is shown that the system in the underdoped regime is a nodal liquid, and the length of the Fermi arc increases with increasing temperatures, in qualitative agreement with the experimental observation on cuprate superconductors [4]. Our results also show that the unusual behavior of the quasiparticle dispersion and Fermi arc in cuprate superconductors is intriguingly related to the effect of the normalstate pseudogap.

[1] Shiping Feng, Phys. Rev. B 68, 184501 (2003); Shiping Feng et al., Physica C 436, 14 (2006). [2] Huaisong Zhao, Lulin Kuang, and Shiping Feng, unpublished. [3] Shiping Feng, Huaisong Zhao, and Zheyu Huang, arXiv:1109.3973. [4] A. Kanigel et al., Nature Phys. 2, 447 (2006); U. Chatterjee et al., Nature Phys. 6, 99 (2010).

#### PB29

#### Static spin correlation in $Pr_{2x}Ca_xCuO_4$ Studied by neutron scattering

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It is well known that a small amount of hole-doping into La2CuO4 having K<sub>2</sub>NiF<sub>4</sub>type (attributed as T) crystal structure destroys the antiferromagnetic (AF) order and superconducting phase appears with further hole doping. On the other hand, recent experimental study on Nd<sub>2</sub>CuO<sub>4</sub>-type (attributed as T') structured La<sub>2</sub>CuO<sub>4</sub> with no apical oxygen clarified contrasting properties to that in T-structured La<sub>2</sub>CuO<sub>4</sub>, that is Neel temperature is different in two structured systems. [1] Therefore, the crystal structure is an important factor to determine the physical properties. In order to gain further insight into the physics of doped Mott insulator, we performed neutronscattering measurement on T'-Pr<sub>2</sub>,Ca<sub>4</sub>CuO<sub>4</sub>. The as-grown and annealed x=0.10 samples shows the long-range AF order below TN-290 K. This result suggests a weak Ca doping effect on the AF order, which is contrast to the results in T-La<sub>2x</sub>Sr<sub>x</sub>CuO<sub>4</sub>. In my poster, the details of doping dependence of magnetic order will be presented.

[1] R. Hord et al., Phys. Rev. B 82, 180508 (2010).

#### PB30

#### High-energy neutron scattering study of spin excitation in slightlyoverdoed La<sub>1 87</sub>Sr<sub>0.18</sub>CuO<sub>4</sub>

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Antiferromagnetism in the doped cuprate Mott-insulator has been investigated due to its rich physics and close connection with the high-Tc superconductivity. A finding of a generic form of spin excitations showing the "hourglass" shape in the different class of superconductors La<sub>2.4</sub>(SrBa),CuO<sub>4</sub> (La<sup>214</sup>) [1] and YBa<sub>2</sub>Cu<sub>0</sub>G<sub>ue1</sub> [2] suggests a common role of spin fluctuations in the mechanism of high-Tc superconductivity. Recently, Vigsolle and co-workers revealed the existence of two energy scales in the spin excitations of optimally-doped La<sup>214</sup> [3] and pointed out the possibility of different origins for the excitation separated by energy. To study the dual nature of spin excitations in more detail, we have performed a high-energy neutron scattering measurement on slightly-overdoped La<sub>2.15</sub>(srBa),CuO<sub>4</sub>. (La<sup>214</sup>) (Du<sub>0</sub>, The constant-energy spectrum at low temperature is a broad single-peak in a wide energy range below ~150meV, although the peak-slitting was observed below ~010meV. This feature is somehow similar to the dispersion of spin excitation in the overdoped La<sup>214</sup>. On the other hand, the integrated intensity shows sharp and broad maximums at ~18meV and ~45meV, respectively, as is the case of optimally-doped system. In the poster, we will present the temperature dependence of the overall excitation spectrum and discuss the dual nature of the hourglass-shaped excitation.

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 B. Vignolle et al., Nature Physics, 3, 163 (2007).

#### PB31

#### Quantized massive gauge fields around the doped holes in high-tc cuprates and the relation to iron pnictides Ikuzo Kanazawa\*

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Recently Yazdani[1] has suggested strongly that the high-energy(up to about 400meV) hole-like excitations of the normal state are a direct predictor of the strength of pairing, although he cannot present a model for the excitations. The present author has proposed the mechanisms of evolution of the Fermi arc with increasing temperature[2] and with increasing of hole-doping[3] in high-Tc cuprates. In this study, we present one model for the high-energy excitations, which play an important role on strength of Cooper pairing, and show that quantized massive gauge fields induce short-range spin-fluctuations and short-range orbital fluctuations[4,5].

[1] A. Yazdani, J.Phys.Condense Matters 21,164214(2009) [2]I.Kanazawa, J.Phys.Chem.Solids 66,1388(2005) [3]I.Kanazawa, Physica C 470,5183(2010) [4]H.Kontani and S.Onari,Phys.Rev. Lett.104,157001(2010) [5]Y.Yanagi et al. J.Phys.Soc.Jpn. 79,123707(2010)

#### PB32

**Ho-doping effect on the static stripe order in La<sub>214</sub> superconductor** Masaki Fujita<sup>1\*</sup>, Masanori Enoki<sup>2</sup>, Satoshi Iikubo<sup>2</sup>, Kenji Tsutsumi<sup>3</sup>, Kentaro Sato<sup>3</sup>, Masato Matsuura<sup>1</sup> and Kazuyoshi Yamada<sup>4</sup> <sup>1</sup> Institute for Materials Research, Tohoku University, Japan <sup>2</sup> Graduate School of Life Science and Systems Engineering, Kyushu Institute of Technology, Japan

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We have performed elastic neutron scattering experiments on the superconducting  $La_{1.84}Ho_{0.04}Sr_{0.12}CuO_4$  to study the substation effect of cation with the large magnetic moment at La site on the static spin correlation. In the Ho-free sample with the hole concentration of 0.12, spatially modulated magnetic order is known to be stabilized below T\_spin ~ 30K [1]. In the present Ho-doped sample, we found the appearance of magnetic peaks at (0.5, 0.5±0.118, 0) positions below T\_spin ~ 35K, which is similar to the observation in the pristine sample. Furthermore, the volume corrected intensity in the Ho-free and Ho-doped samples is comparable, suggesting a negligible effect of Ho-substitution at La site. These results are quite contrast to the huge enhancement of magnetic intensity and the increase of T\_spin by substituting magnetic Fe<sup>3+</sup> ions onto CuO<sub>2</sub> planes [2]. Therefore, the stability of stripe order induced by the cation substitution is sensitive to the substituted site.

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#### PB33

## High field paramagnetic meissner effect in Ca doped $YBa_2Cu_3O_{7\cdot\delta}$ single crystals

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We report on dc magnetization of a series of  $Y_{1x}Ca_xBa_zCu_yO_{7,6}$  (x = 0, 0.0025, 0.05, 0.1) single crystals with the goal to study the role of the hole doping on the paramagnetic Meissner effect (PME). [1] The magnetization as a function of the temperature, under constant magnetic field, was recorded adopting the field-cooled procedure, Mfc(T) as well as magnetic relaxation, Mfc(t) for a measurement time up to 50.000s. The Mfc(T) data of our single crystals shows the usual superconductor diamagnetic response when magnetic fields up to 0.5 kOe are applied. In contrast, when H > 0.5 kOe are applied, the Mfc(T) data displays a systematic reduction of the superconductor diamagnetic response until it became predominantly paramagnetic at higher fields not showing tendency to a temperature saturation behavior in the doped samples. This behavior is a signature of the high field paramagnetic Meissner effect (HFPME) in HTSC. [2,3] The Mfc(t) data reveals that the paramagnetic moment of the samples increase continuously as a function of the time. In specific, the Ca doping promotes an enhancement of the HFPME intensity as compared with no doped sample. We suggest the flux compression scenario to explain the HPME in your samples. [4]

V. N. Vieira et al. Phys. Rev. B 76 (2007) 024518-1. [2]F. T. Dias, et al. Physica C 470, S111 (2010).
 F.T. Dias et al. Phys. Rev. B 70 (2004) 224519. [4] A.K. Geim et al. Nature 396 (1998) 144.

#### **PB34**

#### Growth of a-axis oriented thin films of infinite-layer $Sr_{1\mbox{-}x}La_xCuO_2$

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Electron-doped infinite-layer compound Sr<sub>1-x</sub>La<sub>x</sub>CuO<sub>2</sub> has the simplest crystal structure, and is suitable for fundamental studies of high-Tc superconductivity as well as device applications. The epitaxially grown c-axis oriented Sr1, La, CuO2 thin films with Tc of around 40 K were fabricated on the lattice matched substrates like DyScO<sub>3</sub> and KTaO<sub>3</sub>[1]. Since the in-plane coherence length ξab(4.5nm) of Sr<sub>1-x</sub>La<sub>x</sub>CuO<sub>2</sub> is larger than the out-of-plane coherence length  $\xi c(0.3nm)[2]$ , the a-axis oriented film is favorable for junction fabrication. In this paper, we will make the first report on the preparation and properties of a-axis oriented thin films. Sr<sub>09</sub>La<sub>01</sub>CuO<sub>2</sub>(a0=0.3949nm, c0=0.341nm) thin films were prepared by magnetron sputtering. KTaO3(a0=0.3989nm), SrTiO<sub>3</sub>(a0=0.3905nm), (La0.18Sr0.82) (Al0.59Ta0.41)O<sub>3</sub>(LSAT)(a0=0.3868nm) and LaAlO<sub>3</sub>(a0=0.3791nm) substrates were used to investigate the effect of the lattice mismatch. On KTaO<sub>3</sub>, LSAT and SrTiO<sub>3</sub> substrates, c-axis oriented films were obtained. In contrast, on LaAlO<sub>3</sub> substrates with a large lattice mismatch (-4.2%), a-axis oriented films were obtained. One possible reason for the a-axis growth is due to a higher-order lattice matching, that is, lattice mismatch between ten unit cell of Sr<sub>0.9</sub>La<sub>0.1</sub>CuO<sub>2</sub>(001)(3.10nm) and nine unit cell of LaAlO<sub>3</sub>(100)(3.412nm) is 0.03%. As-grown thin films showed a semiconducting behavior. Post deposition annealing resulted in the a-axis oriented films with Tconset of around 30 K

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#### PB35

## Thermal stability of an epoxy-impregnated HTS racetrack coil without turn-to-turn insulation for rotating machines

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Since use of the no-insulation (NI) technique creates a compact system with better thermal and electrical stability, use of a high temperature superconducting (HTS) coil without turn-to-turn insulation is now being proposed for the field coil in HTS rotating machines. Epoxy impregnation of field coils is generally needed to protect the coil against mechanical disturbances by time-varying magnetic fields and to provide high mechanical integrity against the rotational vibrations. Therefore, to use the NI technique on the practical HTS field coil of a rotating machine, it is essential to obtain sufficient information on the thermal and electrical behaviors of the epoxy-impregnated NI coil. In this study, the quench characteristics of epoxy-impregnated HTS racetrack coils without turn-to-turn insulation were investigated using overcurrent tests. Furthermore, the thermal stability of the epoxy-impregnated NI coil at various current charging rates was also investigated and discussed.

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## POSTER PRESENTATION

#### PB36

**PB37** 

#### Design, fabrication, and testing of a cooling system using solid nitrogen for a 3 T/60-mm RT bore superconducting HGMS

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To develop a compact and stable superconducting high gradient magnetic separator (HGMS) for the removal of copper and silica from chemical mechanical polishing wastewater, a 3 T/60-mm room temperature bore HGMS using solid a nitrogen cooling system was fabricated. The HGMS's magnet was assembled using 18 GdBCO-coated conductor double pancake coils, each of which was wound without turn-to-turn insulation. The cooling capabilities of the solid nitrogen cooling system for the superconducting HGMS were examined using cool-down, warm-up, and charging-discharging tests in the temperature range of 10-20 K. The capture efficiency of the superconducting HGMS using the solid nitrogen cooling system was also evaluated using a filtration test.

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#### PB39

## Removal of silica and copper ions from CMP wastewater via magnetic seeding aggregation using superconducting HGMS

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With the increased demand for semiconductor devices, the purification of chemical mechanical polishing (CMP) wastewater in the semiconductor industry has become of great importance for worldwide environmental protection. Superconducting high gradient magnetic separation (HGMS) is one promising technique due to its compact design and efficient and accurate removal of particles. For the HGMS, various ferromagnetic materials were required to coagulate non-magnetic heavy metal ions and various nanoparticles in CMP wastewater. In this study, we suggested the use of a coagulation process optimized by magnetic seeding aggregation at various pH values for the separate yet simultaneous removal of copper ions and silica nanoparticles. The filtration efficiency of the flocculated substances in an HGMS was evaluated using turbidity measurements.

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## Purification of chemical mechanical polishing wastewater using a 2G HTS high gradient magnetic separation system

Dong Gyu Yang, Jung-bin Song, Young-gyun Kim, Jongseok Lee, Yeonjoo Park and Haigun Lee\*

Department of Materials Science and Engineering, Korea University, Korea

Magnetic separation efficiency is related to the strength of the applied magnetic field and the field gradient used to trap the magnetic particles. The use of a magnet wound with high temperature superconducting (HTS) tape enables the magnetic separation system to more efficiently remove wastewater contaminants than their conventional counterparts due to the facile generation of a higher magnetic field due to the absence of Joule heating losses. In this study, we fabricated and tested the second generation (2G) HTS high gradient magnetic separation (HGMS) system for the purification of chemical mechanical polishing wastewater. The capturing efficiencies of the 2G HTS HGMS system were investigated in terms of the wastewater flow velocity and turbidity using repetitive separation tests. Furthermore, the appropriate purification process for enhancing removal efficiency was also discussed.

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#### **PB38**

## Effect of liquid cryogen on a 2G HTS magnet using a mixed cryogen cooling system

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A cooling system employing a mixture of solid cryogen with small amounts of liquid cryogen was reported for use in high temperature superconducting (HTS) applications. Although the liquid cryogen in the mixed cryogen cooling system can improve the thermal contact between the solid cryogen and the HTS magnet, it is essential that further studies be performed because there are insufficient data about the thermal/ electrical behaviors of the HTS magnet that has been subjected to mixed cryogen. Therefore, in this study, the quench/recovery characteristics of mixed cryogen-cooled 2G HTS magnets with respect to liquid cryogen amount were examined to investigate the influence of the solid/liquid ratio on the thermal/electrical stabilities of the HTS magnet.

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## PB40

## Joint characteristics of ReBCO-coated conductors using various fusion splicing techniques

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For superconductor splicing, it is essential that joint resistance be minimized between superconducting tapes to operate with the persistent current mode and lengthen them. We have already proposed the concept of RE-Ba<sub>2</sub>Cu<sub>3</sub>O<sub>7-6</sub> (ReBCO) tape fusion splicing technology that is based on the unique ReBCO thermal and structural properties. The melting point of ReBCO decreases as the surrounding oxygen partial pressure decreases, allowing for low temperature splicing and recovery of the superconductivity lost during the joining process by a subsequent oxygenation annealing process. In this study, the fusion splices with changing joining parameters, such as PO<sub>2</sub>, heating rate, peak temperature and holding time, cooling rate, oxygen flow rate, temperature and holding time for oxygenation annealing, tape clamping pressure, and so on were performed using the following two splicing methods: 1) the ReBCO-ReBCO interface was directly connected by fusion, and 2) the Ag layer was diffused in a face-to-face manner without solder. The joint characteristics were evaluated in terms of joint resistances, critical currents, and index numbers. Furthermore, the sample morphologies and structures were analyzed by scanning electron microscopy (SEM) and X-ray diffraction (XRD), respectively.

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#### PC01

#### Bulk electronic structure of LaRu<sub>2</sub>P<sub>2</sub> probed by soft X-ray angleresolved photoemission spectroscopy (SX-ARPES)

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#### PC02

 $\begin{array}{l} \textbf{Superlattice quantum critical point in the cubic metal (Sr/Ca)_{3}Ir_{4}Sn_{13} \\ \textbf{Lina Esther Klintberg}^{l}, Swee Kuan Gohl*, Patricia Alireza^{l}, Paul Saines^{l}, David Tompsett^{l}, Peter Logg^{l}, Jinhu Yang^{2}, Bin Chen^{2}, Kazuyoshi Yoshimura^{2} and Malte Grosche^{l} \\ \end{array}$ 

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Structural self-organisation is a central theme in condensed matter physics. Often, the symmetry of a given parent structure is lowered by subtle structural variations which decrease the electronic degeneracy and thereby the total energy. Examples include Jahn-Teller and Peierls distortions and, more generally, modulated lattice distortions, or superlattices. The quasi-skutterudite compounds R3T4X13, where R is an earth alkaline or rare-earth element, T is a transition metal and X is a group-IV element, form either in a simple cubic structure, called the I-phase, or in the variant I'-phase, which can be viewed as a superlattice distortion of the I-phase. Here, we investigate the borderline system  $(Sr/Ca)_2Ir_4Sn_{13}$ , which we report forms in the I-phase, but transforms into the I'-phase on cooling. Hydrostatic pressure is used to suppress the second order structural transition enabling the first comprehensive investigation of a structural quantum critical point in a three-dimensional charge density wave driven superconductor.

#### PC03

**Unconventional superconductivity in PuCoIn<sub>5</sub>: An NQR investigation.** Hiroshi Yasuoka<sup>1</sup>, Georgios Koutroulakis<sup>1</sup>\*, Hiroyuki Chudo<sup>2</sup>, Eric D. Bauer<sup>1</sup> and Joe D. Thompson<sup>1</sup> <sup>1</sup> Los Alamos National Laboratory, USA <sup>2</sup> Advanced Science Research Center, Japan Atomic Energy Agency, Japan



#### PC04

## Pressure effect on the structural and superconducting transitions in a caged compound $PrRh_2Zn_{20}$

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PrRh<sub>2</sub>Zn<sub>20</sub> crystallizes in a cubic CeCr<sub>2</sub>Al<sub>20</sub>-type structure, where the Pr<sup>3+</sup> ion is encapsulated in a polyhedron formed by 16 zinc atoms. The crystalline electric field ground state of 4f2 electrons in the Pr ion was found to be a nonmagnetic T3 doublet [1]. The resistivity of this compound shows a superconducting transition at TSC = 0.060 K and a hysteretic behavior in the temperature range from 140 K to 470 K, which originates from a structural phase transition. In order to reveal the interplay between the structural and superconducting transitions, we have measured the electrical resistivity  $\rho$  for a single crystalline sample under pressures up to 2 GPa and at temperatures down to 0.052 K. With increasing pressure up to 2 GPa, the anomaly in  $\rho(T)$  at TS = 160 K shifts to higher temperatures by the ratio of 3 K/GPa. This indicates that the lowtemperature phase (T<TS) is stabilized under pressure. At P = 0.85 GPa, TSC was not observed down to 0.052 K, indicating that the superconducting state is suppressed under pressure. The opposite responses to the pressure imply that the superconducting state is in competition with the low-temperature phase below TS.

[1] N. Nagasawa, to be published in the proceedings of SCES 2011.

#### PC05

#### Phonon raman scattering of YB6

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YB6 becomes superconducting state below 7.5K. Both tunneling and photoemission measurements suggested that YB6 is the strong coupling superconductor with  $2\Delta\!/$ kBTc =3.8 and phonons with 8meV play an important role for superconductivity. We once reported the results of Raman scattering using 514.5nm laser in SCES2010. Very recently we have measured the Raman scattering spectra with the high resolution using 674 1nm laser. Three Raman active phonons are similar with the previous report, that is, the doublet spectra of the T2g and Eg modes suggest that the crystal symmetry of YB6 is F4/mmm as already reported. The remarkable change has been observed in the low energy spectra below 200cm-1, where the vibration due to Y is dominant. Below 100K, the peak at about 100cm-1 disappears and two new peaks appear. At this temperature the resistivity deviated from the linear dependence. This implies the structural change at low temperature. Furthermore, the energy of two lowest-energy peaks at 50 and 70cm-1 decreases with decreasing temperature. Below Tc, their energies become constant. These results are the experimental evidences that the Y vibration is highly anharmonic and that the anharmonic vibration or rattling mode due to Y is important for superconductivity.

#### PC06

## Vortex lattice structures in spin-triplet superconductors with weak spin-orbit coupling

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The oxide  $Sr_2RuO_4$  was discovered to be superconducting by Maeno et al. in 1994[1].  $Sr_2RuO_4$  is considered to be a spin-triplet superconductor from both theoretical and experimental point of view. Recently, the direction of d-vector in the superconducting state of  $Sr_2RuO_4$  subjected to the magnetic field is discussed. In order to the determine the direction of d-vector, we investigate the vortex lattice structure of  $Sr_2RuO_4$  with a weak spin-orbit coupling for the magnetic field along c-axis. We will report the difference in vortex lattice structures between chiral state and helical state.

[1]Superconductivity in a layered perovskite without copper, Y. Maeno, H. Hashimoto, K. Yosida, S. Nishizaki, T. Fujita, J. G. Bednorz and F. Lichtenberg, Nature (London) 372, 532 (1994)

#### PC07

## Effect of the magnetic trapped flux on the heat capacity of the low-temperature superconductors: Pb, La, Sn

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Effect of the magnetic prehistory on the temperature dependence of the heat capacity of the superconducting Pb, La, Sn has been studied. The intermediate state is characterized by a subdivision of the specimen into superconducting and normal regions that the relative total volumes of these two types of region are determined by the mean induction. The trapped flux was produced after the magnetic field was turned off. We observed a difference of the heat capacity between zero field cooling (ZFC) and field cooling (FC) states in zero magnetic field for the ring specimen. It was found that the FC heat capacity has a smaller value than the heat capacity both in the normal and in superconducting phases and gives rise to the change of the heat capacity of the specimen. This result is the first experimental evidence of the surface energy contribution to the heat capacity. The work was supported by Russian Foundation For Basic Research Grant No. 10-02-96019-r-ural-a.

#### PC08

A novel superconductivity in Ir oxides with large spin-orbit coupling Hiroshi Watanabe\*, Tomonori Shirakawa and Seiji Yunoki *RIKEN, Japan* 

Recently, the 5d transition metal oxides such as Sr<sub>2</sub>IrO<sub>4</sub> and Ba<sub>2</sub>IrO<sub>4</sub> have attracted much attention as a novel Mott insulator. In these materials, three t2g orbitals of Ir atoms are hybridized with each other by the spin-orbit coupling of 5d electrons. As a result of the quantum entanglement of spin and orbital degrees of freedom, an anomalous Jeff=|L-S|=1/2 state is realized, which causes interesting properties [1]. To clarify the properties of this system, we have studied the ground state of the three-orbital Hubbard model with a spin-orbit coupling term using variational Monte Carlo method and variational cluster approximation [2]. Here, we study the electronic states when carriers are doped in this three-orbital system and discuss the possibility of superconductivity. The obtained ground state phase diagram reveals the antiferromagnetic state, stable around the electron density n=5, is destabilized by carrier doping and the ground state turns to be superconducting under a certain condition. Similar to the high-Tc cuprates, a large asymmetry between electron doping (n>5) and hole doping (n<5) is also observed. Due to the large spin-orbit coupling, the spin is no longer a good quantum number. Instead, the pseudospins form a singlet pairing and a "d-wave like" superconductivity is realized.

[1] B. J. Kim et al., Science 323, 1329 (2009). [2] H. Watanabe et al., Phys. Rev. Lett. 105, 216410 (2010)

#### PC09

**Fermi surface studies of Sr<sub>3</sub>Ir<sub>4</sub>Sn<sub>13</sub> via the Shubnikov-de Haas effect** Swee K. Goh<sup>1</sup>\*, Lina Klintberg<sup>1</sup>, David A. Tompsett<sup>1</sup>, Sven Friedemann<sup>1</sup>, Stan Tozer<sup>2</sup>, Jinhu Yang<sup>3</sup>, Bin Chen<sup>3</sup>, Kazuyoshi Yoshimura<sup>3</sup> and Malte Grosche<sup>1</sup> <sup>1</sup> University of Cambridge, United Kingdom <sup>2</sup> National High Magnetic Field Laboratory, Florida State University, USA <sup>3</sup> Kyoto University, Japan

The ternary stannide system  $(Ca,Sr)_3Ir_4Sn_{13}$  is recently reported to feature a superlattice quantum critical point, where the superlattice distortion transition temperature T\* is suppressed to 0 K [1]. Sr\_3Ir\_4Sn\_{13} lies on the part of the phase diagram where T\*-147 K, and hence the low temperature phase is the superlattice phase. Using the Shubnikov-de Haas effect, we directly study the Fermiology of this superlattice phase, and compare the results with bandstructure calculations performed both with and without the superlattice distortion. The effective mass and the mean free path associated with various Shubnikov-de Haas frequencies will be presented. The relevance of the present study to superconductivity and quantum criticality will be discussed.

[1] L. E. Klintberg, S. K. Goh, P. L. Alireza, P. J. Saines, D. A. Tompsett, P. W. Logg, J. Yang, B. Chen, K. Yoshimura and F. M. Grosche, submitted (2012)

#### **PC10**

#### Magnetism in CeIr(Si<sub>x</sub>Ge<sub>1-x</sub>)<sub>3</sub> compounds

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The CeMX<sub>3</sub> intermetallics adopting the noncentrosymmetric tetragonal BaNiSn<sub>3</sub>-type crystal structure recently attracted much attention mainly owing to unconventional superconductivity observed in some of them [1]. CeIrSi<sub>2</sub> and CeIrGe<sub>2</sub> were reported to exhibit antiferromagnetic ordering with TN = 5.0 K and 8.5 K, respectively, CeIrGe, also undergoes an order-order transition at T1 = 4.7 K. When applying hydrostatic pressure TN gradually decreases and a superconducting dome appears in the vicinity of OCP for the pressure Pc = 2.2 and 23 GPa, respectively [1, 2]. Our work has been motivated by the desire to know how the magnetism and superconductivity evolve with the composition of the SiGe sublattice. For this purpose we have prepared polycrystalline samples of selected compositions of the SiGe sublattice by arc melting, annealed and characterized them and measured magnetization, resistivity and specific heat in a wide temperature range and various magnetic fields. The solutions keep the crystal symmetry of the parent compounds and the lattice parameters follow Vegard's law. The TN vs. x dependence is not entirely monotonous; a minimum TN is observed for x between 0.75 and 1. The results of powder neutron diffraction will be presented as well

[1] Ch. Pfleiderer, Rev. Mod. Phys. 81, 1551 (2009) and ref. therein [2] F. Honda et al., PRB 81, 140507 (2010)

#### PC11

#### **Superconductivity at 5.2K in ZrTe<sub>3</sub> polycrystals** P. L. Paulose\* and C. S. Yadav

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Interplay of the Charge Density Wave (CDW) and Superconducting (SC) states is a subject of significant interest. Low dimensional chalcogenides of transition metals are one of the well studied systems that show the coexistence of these competing phenomena. Recently Cu and Ni intercalated single crystals of ZrTe<sub>3</sub> are reported to show bulk SC at 3.8K and 3.1K respectively. However single crystalline ZrTe<sub>3</sub> show only filamentary SC below 2K, along with CDW at 63K. We have studied the polycrystalline ZrTe<sub>3</sub> and the effect of intercalation of Cu and Ag on the superconducting and CDW states of the compound. Our main finding is the occurrence of filamentary superconductivity at enhanced temperature of 5.2K in the polycrystalline ZrTe<sub>3</sub> at ambient pressure, and the SC coexists with the CDW phase. Enhancement in Tc mimics the effect of external pressure on SC in ZrTe<sub>3</sub> single crystals. The strains between the agglomerated small single crystallites in polycrystalline samples could be one possible reason. The intercalation of Cu or Ag, does not affect the transition temperatures but suppresses the CDW state. We have analyzed the resistivity data to estimate the relative loss of carriers and reduction in the nested Fermi surface area upon CDW formation in these

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#### PC12

Coexistence of superconductivity and antiferromagnetism in CeNi<sub>0.8</sub>Bi<sub>2</sub> Soo-whan Kim<sup>1</sup>, Soohyun Kim<sup>1</sup>, Kyujoon Lee<sup>1</sup>, Adrian Hiller<sup>2</sup>, Devashibhai Adroja<sup>2</sup> and Myung-hwa Jung<sup>1</sup>\* <sup>1</sup>Department of Physics, Sogang University, Korea

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There have been some reports on RNiBi<sub>2</sub> (R=rare-earth elements) with ZrCuSiAstype structure[1,2]. Both superconductivity and antiferromagnetism were observed in CeNi<sub>0.8</sub>Bi<sub>2</sub>. The previous studies suggested two types of carriers with different masses; a light electron responsible for superconductivity and a heavy electron for antiferromagnetism[3]. However, the antiferromagnetic transition was not found in specific heat data. Thus, we have also studied RNi<sub>0.8</sub>Bi<sub>2</sub> (R = La, Ce) in order to take a close look at the possible coexistence of superconducting and antiferromagnetic order parameters. As observed previously, both samples are type-II superconductors with the superconducting transition temperature TC = 4 K. For CeNi<sub>0.8</sub>Bi<sub>2</sub> the lower and upper critical fields are found to be 62 Oe and 660 Oe. Two antiferromanetic transitions are observed at TN1 = 5 K and TN2 = 7 K. These results well agree with the magnetization versus field data showing two metamagnetic transitions around 5 kOe and 40 kOe. We suggest that they correspond to commensurate and incommensurate antiferromagnetic transitions, respectively. These results are also in good agreement with the neutron experiments.

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 M. H. Jung, A. H. Lacerda, and T. Takabataka, Phys. Rev. B 65, 132405 (2002) 3. H. Mizoguchi, S. Matsuishi, M. Hirano, M. Tachibana, E. Takavama-Muromachi, and H. Karwaji, Phys. Rev. Lett. 106, 057002 (2011).

#### PC13

#### Optical studies of superconducting InN thin film

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The indium nitride (InN) thin film (thickness about 220 nm) was grown on (0001) sapphire substrate by metal-organic vapour phase epitaxy technique. The superconducting transition temperature determined by dc resistivity and magnetization gave a Tc about 2 K. Near-infrared photoluminescence from InN thin film can be clearly observed at room temperature. The optical reflectance was measured over a wide frequency range (50 ~ 52000 cm-1) and at temperature between 10 and 340 K. To extract the optical constants of the films, the Drude-Lorentz model was used to fit all of the layers of this thin-film structure. From the parameters obtained, the optical constants were computed. The vibrational spectrum of InN thin film consists of a A1(TO) and a E1(TO) at about 448 cm-1 and 478 cm-1. The temperature variation of the phonon parameters can be explained by the anharmonic effect. Notably, the Drude plasma frequency (~ 3200 cm-1) and the position of the band absorption (~ 6500 cm-1) exhibit little temperature dependence, whereas the scattering rate monotonically decreases with decreasing temperature. All of these observables suggest the unusual normal-state behavior in this superconducting thin film.

#### PC14

#### Lateral Josephson junction induced by inverse proximity effect Lu-kuei Lin<sup>1</sup>, Ssu-yen Huang<sup>1</sup>, Jin-hua Huang<sup>2</sup> and Shang-fan Lee<sup>1\*</sup> <sup>1</sup>Institute of Physics, Academia Sinica, Taiwan <sup>2</sup>Materials Science and Engineering, National Tsing Hua University, Taiwan

A Josephson junction is characterized by a phase coherent transfer of the Cooper pairs across a weak link between two superconducting electrodes. The proximity effect at the interface between a superconductor and a ferromagnetic metal occurs, implying that the Cooper pairs penetrate into the ferromagnet. It results in spatial oscillations of the superconducting order parameter in the ferromagnet. Recently an opposite effect so-called "inverse proximity effect", i.e. the induction of a magnetic moment in a superconductor in contact with a ferromagnet attracted much attention. We utilize the inverse proximity effect in superconductor-ferromagnet bilayer to generate lateral Josephson junctions. The weak link is created by a strong ferromagnet Py strip across a superconductor Nb bridge due to inverse proximity effect and forms a region S'. Samples with different size of weak link region S' were fabricated by varying the Py strip width from 0.5 to 2 µm. The junctions exhibit a modulation of the critical current in perpendicular magnetic field similar to a Fraunhofer interference pattern even when the S' has such long distance as 2 µm, which proves the dc Josephson effect and spin triple component taking place.

#### PC15

#### 13C NMR study of charge fluctuation induced superconductivity in beta''-(BEDT-TTF)4[(H<sub>3</sub>O)Ga(C<sub>2</sub>O<sub>4</sub>)<sub>3</sub>]-C<sub>6</sub>H<sub>5</sub>NO<sub>2</sub> Yoshihiko Ihara\*, Harumi Seki and Atsushi Kawamoto

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Superconductivity emerging near antiferromagnetism has been thoroughly investigated. The antiferromagnetic fluctuations in the vicinity of quantum critical point are believed to induce such superconductivity. In this paper, we report experimental evidence that points to the novel possibility of superconducting state, that is, charge fluctuation induced superconductivity. The molecular superconductor beta"-(BEDT-TTF)4[(H<sub>3</sub>O) Ga(C<sub>2</sub>O<sub>4</sub>)<sub>3</sub>]-C<sub>6</sub>H<sub>5</sub>NO<sub>2</sub> demonstrates a charge instability at 100 K, before showing the superconducting transition at 7.5 K [1]. Besides, this salt is intriguing because of its high upper critical field of 33 T, which is almost three times the Pauli limiting field [2]. We performed the 13C NMR experiment and revealed on the basis of the Knight shift measurement that the spin-singlet superconductivity is realized in the charge disproportionate state. The nuclear spin-lattice relaxation rate measurement detects the strong electron-electron correlations slightly above superconducting transition temperature. Since 13C NMR technique can probe only the magnetic properties, we measured in-plane resistivity as a complementary probe and found an abrupt increase below 10 K. These results lead us to address the importance of charge fluctuations to induce superconductivity

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#### PC16

## Ac susceptibility components of a thin type-II superconducting annulus carrying a radial current

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By assuming a spatial dependence on the sheet-current density, we have investigated the real  $(\chi')$  and imaginary parts  $(\chi'')$  of the susceptibility of a thin type-II superconducting annulus in the absence and also in the presence of a transport radial current, using the Bean critical state model in which the critical current density is assumed to be independent of the local magnetic field. The results of our calculations on the components of the magnetic susceptibilities in two cases and for the different aspect ratios are compared with each other. The comparison shows that by applying a radial current to the sample, the imaginary part of the susceptibility is increased. We have also studied the variations of the  $\chi''$  with respect to  $\chi'$  for a several aspect ratios.

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#### PC17

#### Pressure evolution of superconductivity in $\beta$ -pyrochlore oxides

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The  $\beta$ -pyrochlore oxides AOs<sub>2</sub>O<sub>6</sub> (A = K, Rb and Cs) have attracted much interest because of their unusual electronic properties including the superconductivity[1]. A unique feature of this system is the presence of the remarkable anharmonic vibration of an A ion, called 'rattling' motion. Although it has been pointed out by various studies that the rattling motion plays an important role for the occurrence of superconductivity in this system, its concrete role for superconductivity have not been understood yet. In order to clarify the interplay between the rattling motion and superconductivity, we study the pressure (P) evolution of superconductivity in magnetic field (H) at low temperatures (T). The rattling motion could be affected by pressure through the modification of the crystal lattice. In this paper, we report the T-H-P phase diagram of AOs<sub>2</sub>O<sub>6</sub> and K<sub>0.8</sub>Rb<sub>0.2</sub>Os<sub>5</sub>O<sub>6</sub> obtained by specific heat measurements and discuss the pressure variation of the electron-phonon coupling and phonon frequency. We demonstrate that the non-monotonous pressure dependence of the transition temperature can be explained in terms of the competition between the enhancement of the electronphonon coupling and the suppression of the phonon frequency.

[1]Y.Nagao et al., J.Phys.Soc.Jpn. 78 064702 (2009)

#### **PC18**

## Inverse magnetic proximity effects in superconducting Sn-Ni nanoparticle assemblies

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Although the superconductivity of the element Sn, in their bulk form, is believed to be associated with the spin-singlet s-wave pairing, it is now known that the superconducting parameters are strongly dependent on the physical size of the system. Here, we report on the results of magnetization, magnetic susceptibility and resistivity measurements made on Sn, Sn-Ni and Sn-Au nanoparticle assemblies. The thermal profiles of the magnetic susceptibility of Sn nanoparticle assemblies can be described by Scalapino's expression. Inverse magnetic proximity effect is observed, as the TC of the superconducting nanoparticles is found to be noticeably increased, when magnetic Ni nanoparticles are introduced into the vicinity of the superconducting nanoparticles. Closing up the spatial separations between Sn and Ni nanoparticles leads to a further increase in TC. There is a critical Ni composition and a critical spatial separation between the magnetic and superconducting nanoparticles that must be reached before their magnetic proximity will suppress the superconductivity. On the other hand, the inverse proximity effect of the In-Ni nanoparticle assemblies is not observed in the (Sn-C)14(Au-A)86 nanoparticle assembly, where TC of Sn remains essentially unaltered by cold pressing. A qualitative mechanism is proposed to understand the present observations.

#### PC19

## Pressure study on anisotropic electrical resistivity of Hg-doped $CeRhIn_{s}$

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The stoichiometric compound CeRhIn<sub>5</sub> is a prototypical antiferromagnet where Ce 4f moments align below 3.8 K. With increasing pressure, the antiferromagnetic ordering of CeRhIn<sub>5</sub> disappears and the superconducting state emerges. When doped with Hg, the antiferromagnetic transition TN initially decreases and develops a different magnetic structure with further increasing Hg concentration. In this research, we focused on a 0.45 % Hg-doped CeRhIn<sub>5</sub>, where TN is suppressed from 3.8 K to 3.4 K and the magnetic structure is same as that of the undoped compound with Q=(1/2, 1/2, 0.298). By applying hydrostatic pressure, we suppressed TN to zero and measured the anisotropic electrical transport measurements for the electrical current applied within and perpendicular to the Ce-In plane. Evolution of the anisotropic transport property across the quantum critical point will be discussed.

#### **PC20**

Control of superconductivity in parity mixing superconductors Li<sub>2</sub>T<sub>3</sub>B(T:Pt,Pd) by non-magnetic impurity and defect doping

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Recently, superconductors with noncentrosymmetric crystal structures like Li,T,B (T-Pd, Pt) have been the focus

of in-denth research with its parity mixing nature. Previous studies of NMR and penetration denth measurements suggested that Li-Pt-B is a spin triplet dominant and Li-Pd-B is a s-wave spin singlet dominant superconductor. approximately 2:3 and 4:1 ratio of spin singlet and spin triplet order parameters, respectively [1-3]. It is known that s-wave superconductor without the sign inversion of the order parameter on the Fermi surface is not affected by non-magnetic impurity and defects doning as contrasted to that of a non s-wave superconductor. In this research, we prepared different quality samples by substituting B with Al. Sample qualities were estimated by residual resistivity. Li<sub>2</sub>Pd<sub>3</sub>B exhibits the small Tc suppression attributed by the non-magnetic impurities and defects, while Hc<sub>2</sub>(0) value increased. This behavior is similar in ordinary s-wave superconductor. While results for Li<sub>2</sub>Pt<sub>2</sub>B show that Tc and Hc<sub>2</sub>(0) value were found to be suppressed by disorder. However, the rate of Tcsuppression by disorder has been found to be not so large to be explained by the pair-breaking effect expected for the non s-wave superconductor. Further details will be outlined during the actual presentation.

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#### PC21

#### Pressure-induced metal-insulator transition of Mott insulator Ba<sub>2</sub>IrO<sub>4</sub>

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Ba<sub>2</sub>IrO<sub>4</sub> is an antiferromagnetic Mott insulator, because the large spin-orbit interactions of the 5d electrons in Ir atoms lift the energy levels. The magnetic moment is suppressed to ~ 0.34  $\mu$ B/Ir atom because of the spin fluctuation, and the Neel temperature TN is ~ 240 K. These magnetic properties are similar to ones of La<sub>2</sub>CuO<sub>4</sub>, which is the parent material of a high-Tc superconductor [1-3]. The crystal structures are also similar. Here, in order to inhibit the insulating state and to search the superconducting phase, the electrical resistivity of Ba2IrO4 has been performed under high pressure. According to the present results, metal-insulator transition was observed at 28 GPa, but not the superconducting transition above 50 mK

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#### PC22

#### Charge and spin order of charge stripe ordered La2-xSrxCoO4

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Since the discovery of charge stripe order in the cuprates, the role of stripes in high temperature superconductivity, and whether charge stripes are ubiquitous to the cuprates has been under intensive debate[1]. Recent neutron scattering studies of cobalt and manganese based charge-stripe ordered materials have demonstrated how the striking hourglass excitation spectrum of the cuprates is consistent with charge-stripe order[2,3]. One important experimental observation in charge-stripe ordered La2, Sr, CoO4 is however presently missing, direct observation of stripe charge order itself. Although alternative forms of charge order has been observed in cobaltates in the form of a checkerboard charge ordered state and a proposed charge river state[4,5]. Here we use the technique of resonant soft x-ray scattering to study by diffraction the spin and charge ordering of the charge-stripe phase of La2, Sr, CoO4. We report our analysis of the diffuse order recorded in our study this study completes the evidence for the charge-stripe model of this system that is used to explain the hourglass spin excitation spectrum.

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#### **PC23**

#### Strong enhancement of superconductivity in inorganic electride 12CaO•7Al<sub>2</sub>O<sub>3</sub>:e- under high pressure

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The nanoporous main group oxide  $12CaO \cdot 7Al_2O_3$  (C<sub>12</sub>A<sub>7</sub>) is a compound in the system CaO-Al<sub>2</sub>O<sub>3</sub> and is used widely as a constituent of aluminous cements. It can be transformed from a wide-gap insulator to a metal with substituting electrons for anions in cages constituting a positive frame. A superconducting transition was found at Tc = 0.2-0.4 K for  $C_{12}A_7$  electride ( $C_{12}A_7$ :e-) doped with anionic electrons to 2\*10^21 per cubic centimeter[1]. We carried out the ac-susceptibility measurement of single crystalline C12A7:e- under high pressure using a piston-cylinder cell and a diamond anvil cell. With increasing pressure. Tc monotonically increases up to 1.0 K at 2.4 GPa. In addition, the upper critical field Hc<sub>2</sub> (T) and -dHc<sub>2</sub>/dT becomes larger. The density of states at the Fermi level is estimated to increase under high pressure, in agreement with dTc/dP > 0

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#### **PC24**

#### Superconductivity in conical magnets

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The influence of spatially modulated exchange fields on superconducting properties is discussed. Of particular interest are conical magnets where the non-collinear magnetic structure permits to control the symmetry of the pair wave function. The limiting cases - homogeneous ferromagnetic and anti-ferromagnetic exchange field - have been extensively studied. The central focus of the present work is the variation with modulation length of the superconducting properties. I discuss the evolution of the induced triplet components and analyze their orbital structures. Spinresolved tunnelling spectra highlight the cross-over from typical pair-breaking to pairweakening which occurs when the magnetic modulation length is comparable to the superconducting coherence length in the absence of exchange fields.

#### PC25

#### Observation of Bose-metallic phase in Ta Films Sungvu Park\* and Eunseong Kim

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Superconductor-insulator transition has been induced by tuning film thickness or magnetic field [1-3]. Recent electrical transport measurements of Ta thin films revealed an interesting intermediate metallic phase which intervened superconducting and insulating phases in the zero temperature limit [4]. The resistance of Ta films in this regime exhibited a sharp drop at the transition temperature but a finite saturated value at low temperatures. In addition, IV characteristic curves showed non-linear response, indicating possible appearance of a new metallic phase. The intriguing quantum metallic phase can be interpreted as a consequence of vortex dynamics or dislocation response in the vortex lattice [5]. Here we present systematic studies on the superconductor-metal-insulator transition and nonlinear transport in two dimensional Ta films

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#### PC26

#### Non-trivial vortex dynamics in a superconducting Corbino disk Masaru Kato and David E. Fujibayashi

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Using the molecular dynamics method for vortex motion and the finite element method for heat transport, we investigate the vortex motion in a superconducting corbino disk, where an external current is injected at the center of the disk and flows toward the perimeter of the disk. For this geometry, vortices move circularly, but velocity of the vortices is large around the center and small at the edge of the disk, according to the magnitude of the external current density. Moving vortices cause resistivity and heat, and therefore there appear non-uniform temperature distribution. This non-uniform temperature distribution makes the vortices out of their circular motion through entropy force. Therefore the motion in the superconducting corbino disk becomes non-trivial one. We focused on the how heat transport to the outside of superconductors, such as a substrate under the superconductor, affect the vortex motion. And we find that firstly the vortices move toward the edge of the disk, but finally they form laminar flow for large heat resistance to the substrate but they show a peculiar ordered state for small heat resistance. We show details of these vortex motions

#### **PC27**

#### Coexistence of ferromagnetism and superconductivity in single-phase Bi<sub>3</sub>Ni nanostructures

Thomas Herrmannsdoerfer1\*, Richard Skrotzki1, Rico Schoenemann1, Yurii Skourski1, Joachim Wosnitza<sup>1</sup>, Daniel Koehler<sup>2</sup>, Regine Boldt<sup>2</sup> and Michael Ruck<sup>2</sup> Dresden High Magnetic Field Laboratory, Helmholtz-Zentrum Dresden-Rossendorf, Germany <sup>2</sup>Department of Chemistry and Food Chemistry, TUDresden, Germany

We have demonstrated the coexistence of superconductivity and ferromagnetism in Bi<sub>3</sub>Ni nanostructures which have been prepared by making use of novel chemical-reaction paths [1]. We have characterized their magnetic and superconducting properties by means of SQUID magnetometry, ac susceptometry, and electrical-transport measurements in a wide field and temperature range. Here, we also present recent experiments on novel nanostructures, such as monodisperse spherical clusters with a diameter of 8 nm as well as nanofibers. Pulsed-field susceptibility data up to 60 T allow for a determination of the saturation magnetization of Bi<sub>3</sub>Ni nanostructures. Resistivity measurements performed on moderately compacted Bi<sub>3</sub>Ni nano fibers down to 40 mK have shown clear evidence for the existence of an isosbestic point. Superconductivity in confined Bi<sub>3</sub>Ni emerges in the ferromagnetically ordered phase and is stable up to remarkably high magnetic fields. This coexistence would most likely be possible in the case of triplet pairing. The absence of an inversion center of the lattice of Bi<sub>3</sub>Ni nanostructures would allow for the formation of an antisymmetric spatial component of the electron-wave function and could lead to a significant admixture of a spin-triplet component of the order parameter. Part of this work was supported by EuroMagNET, EU-contract No. 228043.

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#### PC28

#### Observation of twofold symmetry breaking in the gap function of heavy-fermion superconductor UPt<sub>3</sub>

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#### PC29

#### Investigation of the three-dimensional electronic structure of MgB<sub>2</sub> by soft X-ray angle-resolved photoelectron spectroscopy

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#### PC30

#### Vortex channeled effect in Nb thin film with artificial pinning array

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The vortex pinning effect is a fundamental characteristic in type II superconductors Vortex dynamics in the mixed state in type-II superconductors is strongly influenced by the pinning centers. We have studied a channeled-like pinning potential in Nb thin films which is formed by the pinned vortices. Nb films with periodic pinning arrays show matching effects in the magnetic field dependence of the resistivity and Hall Effect. It is suggested that the motion of interstitial vortices has great influence by the channel. When vortices propagate through these arrays, both the longitudinal and transversal voltages show cusp-like anomalies at matching fields. Vortices can be guided by the channel formed between pinned vortices.

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#### PC31

#### Role of the third dimension on the spectral property and transport behaviour in layered cuprates

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In the present work, we have analyzed the influence of the third dimension coupling on the normal state electronic spectra of bilayer cuprates employing Green's function equations of motion approach within tight binding extended Hubbard model and model, which necessarily includes the three site exchange interaction (J3) and the inter unit cell resonant tunneling (T12). The influence of inter unit cell resonant tunneling on the out of plane (c-axis) conductivity has also been analyzed for cuprates employing Kubo formula and Green's function technique. On the basis of numerical computation, it is found that the inter cell resonant tunneling (T12) introduces a broadening in the spectral function and increases the c-axis conductivity. It is also found that the J3 term suppresses the bilayer splitting in the electronic spectra. These results are viewed in terms of recent ARPES measurements

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#### PD01

#### Ground dielectric state of the Mott-doped material

Vladimir Gavrichkov\* Krasnoyarsk, Kirensky Institute of Physics, Russia

Using analytic calculations, within the framework of the Russell-Saunders scheme we investigated the dependence of the number of valence states on a doping level x in the full package of the valence bands for the Mott-doped material. According to our calculations the total number of valence states Nv(x) coincides with the number of electrons per unit cell (N-x) if there is a ban (both spin S-, and orbit L-) on the first removal quasiparticle states (frs). Thus is formed an insulating ground state of the doped material at the T=0K. We believe that such materials with the forbidden states with the S-forbidden states.

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#### **PD02**

## Temperature-dependent phonon anomalies in uranium and plutonium compounds

Peter S Riseborough\* Physics, Temple University, USA

There is evidence that a number of heavy-fermion/mixed-valent materials show strongly renormalized hybridization gaps either at the Fermi-energy or close to the Fermi-energy. In the former case, a heavy-fermion semiconducting state ensues and in the later case, the system remains metallic at low temperatures. Due to the temperature-dependence of the electronic correlations, the magnitudes of the hybridization gaps decrease with increasing temperatures. The existence of a temperature-dependent low-energy electronic energy scale opens up the possibility that the Born-Oppenheimer approximation may fail and that there may be a resonant coupling between the phonons and the electronic excitations. It is argued that such a mechanism may be the cause of the anomalous temperature-dependence of the phonon spectrum in Pu and the anomalous phonon mode observed in  $\alpha$ -uranium.

#### PD03

Non-linear conductivity of resistive oxides: truths and myths B. - Fisher\*, J. - Genossar, L. - Patlagan and G. M. Reisner *Physics, Technion, Israel* 



#### **PD04**

#### The modulated spin liquid and hidden order in URu<sub>2</sub>Si<sub>2</sub> Sebastien Burdin

Bordeaux University, France

We have shown that near a Kondo breakdown critical point, a spin liquid with spatial modulations is likely to be formed. Unlike its uniform counterpart, we find that this occurs via a second order phase transition. The amount of entropy quenched when this ordering is manifest is of the same magnitude as for an antiferromagnet. Moreover, these two states are in direct competition with each other, and at low temperatures they are separated by a first order phase transition. This suggests that the modulated spin liquid is indeed a viable candidate for the unique phase of matter which is observed in the hidden order phase of URu<sub>2</sub>Si<sub>2</sub>. We investigate the modulated spin liquid proposal taking full account of the tetragonal-body-centered lattice structure which characterizes URu<sub>2</sub>Si<sub>2</sub>. We find that the physical quantities predicted from the model are in very good qualitative and quantitative agreement with recent experiments, including thermodynamic properties, inelastic neutron scattering, and Fermi surface measurements. We also present an overview of other f-electron compounds where similar phases have been observed.

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#### PD05

#### Fermi-surface evolution in Yb-substituted CeCoIn<sub>5</sub>

Joachim Wosnitza<sup>1\*</sup>, Andrey Polyakov<sup>1</sup>, Andrea D. Bianchi<sup>2</sup>, S. Blackburn<sup>2</sup>, B. Prevost<sup>2</sup>, G. Seyfarth<sup>2</sup> and Michel Cote<sup>2</sup>

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In CeCoIn5, substitution on the rare-earth site has been found to influence the Kondolattice coherence and Cooper pairing in a rapid and uniform way. Yb substitution, however, shows a different behavior in that Tc and the Kondo-coherence temperature do not scale, in contrast to other heavy-fermion compounds. We performed a comprehensive dHvA study on Ce1-xYbxCoIn<sub>5</sub>. For small Yb concentration, x =0.1, the band-structure topology and the ef-fective masses remain nearly unchanged compared to CeCoIn<sub>5</sub>. This contrasts clearly modified Fermi surfaces and light, almost unrenormalized effective masses for x = 0.2 and above. These observations cannot explain the heavy-fermion physics observed in specific-heat and resistivity data even for high Yb concentrations. Thus, we suggest the existence of heavy quasiparticles with short mean free paths, not detectable by dHvA experiments. However, the mechanism by which superconductivity can emerge from these charge carriers remains elusive. Work done in cooperation with K. Gotze, D. Hurt, Z. Fisk, R. G. Goodrich, I. Sheikin, M. Richter. Part of this work was supported by EuroMagNET II (EU contract No. 228043).

#### PD06

Spin exciton formation inside the hidden order phase of CeB<sub>6</sub> Alireza Akbari and Peter Thalmeier MPI-CPIS Germany

The heavy fermion metal CeB6 exhibits hidden order of antiferroquadrupolar (AFQ) type below  $T_0=3.2K$  and subsequent antiferromagnetic (AFM) order at  $T_x=2.3K$ . It was interpreted as ordering of the quadrupole and dipole moments of a  $\Gamma_8$  quartet of localised Ce 4f<sup>4</sup> electrons. This established picture has been profoundly shaken by recent inelastic neutron scattering[1] that found the evolution of a feedback spin exciton resonance within the hidden order phase at the AFQ wave vector appears and is stabilized by the AFM order. We develop an alternative theory based on a fourfold degenerate Anderson lattice model, including both order parameters as particle-hole condensates of itinerant heavy quasiparticles. This explains in a natural way the appearance of the spin exciton resonance and the momentum dependence of its spectral weight, in particular around the AFQ vector and its rapid disappearance in the disordered phase. Analogies to the feedback effect in unconventional superconductors are nointed out

G. Friemel, et al, Nature Communications 3, 830 (2012).
 A. Akbari and P. Thalmeier, Phys. Rev. Lett. 108, 146403 (2012).

#### PD07

#### Antiferromagnetic fluctuation in hidden order phase of U(Ru,Rh)2Si2

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The origin of so-called hidden-order (HO) transition observed at To=17.5K in URu<sub>2</sub>Si<sub>2</sub> has been one of the longstanding issues in physics of heavy-fermion systems. In the mixed compounds U(Ru1-xRhx)<sub>2</sub>Si<sub>2</sub>, it is found that a large-moment antiferromagnetic (AF) order replaces HO for  $0.02 \le < 0.03$  via a first-order phase transition, and both orders disappear above x≥0.04. To clarify microscopic features of magnetic fluctuation in both HO and AF phases, we have performed inelastic neutron scattering experiments for U(Ru1-xRhx)<sub>2</sub>Si<sub>2</sub> (x≤0.03). At the AF wave vector Q=(1,0,0), the AF excitation peak is clearly observed in HO phase for  $0 \le < 0.02$ , but it disappears in the AF phase for  $0.02 \le < 0.03$ . The excitation gap at 1.4K is found to be weakly reduced from 2.4meV (x=0) to 1.8meV (x=0.015) as x is increased. Temperature variations of the staggered susceptibility  $\chi$ (Q,T) at x=0 estimated from the inelastic-peak intensities show a weak enhancement in HO phase becomes pronounced, but a cusp-like anomaly at To is reduced. These features indicate that the AF fluctuation develops on the verge of the first-order phase boundary in HO phase, while it is suppressed in the AF phase.

#### PD08

## Pressure effect on the field-induced ordered phase in heavy fermion compound $YbC_{0,2}Zn_{2,0}$

Tetsuya Takeuchi<sup>1</sup>, Yuki Taga<sup>2</sup>, Shingo Yoshiuchi<sup>2</sup>, Masahiro Ohya<sup>2</sup>, Yusuke Hirose<sup>2</sup>, Fuminori Honda<sup>2</sup>, Rikio Settai<sup>2</sup> and Yoshichika Onuki<sup>2</sup> <sup>1</sup>Low Temperature Center, Osaka University, Japan <sup>2</sup>Graduate School of Science, Osaka University, Japan

We succeeded in growing high-quality single crystals of heavy fermion compounds  $YbT_2Zn_{20}$  with T=Co, Rh, and Ir, and found a metamagnetic behavior at  $H_{\rm m}$  below the characteristic temperature Tymax, where the magnetic susceptibility takes a broad maximum. For  $YbCo_2Zn_{20}$ ,  $H_{\rm m}=6$  kOe and Tymax=0.32 K, for example. Hm and Tymax in these compounds were found to satisfy a universal relation  $H_{\rm m}$  (kOe) = 15Tymax (K). Recently, we found a field-induced ordered phase (FIOP) above  $H_{\rm Q}=60$  kOe for H // <111> in  $YbCo_2Zn_{20}$  below  $T_{\rm Q}=0.6$  K and suggested an antiferroquadrupole order for the FIOP. In this work, we performed pressure experiments in  $YbCo_2Zn_{20}$  with H // <111> in order to study the pressure effect on the FIOP as well as the pressure-induced antiferromagnetic phase (PIAF) above  $P_c-2$  GPa. When the pressure up to around  $P_c$ . Above  $P_c$ , there is a pressure region where the PIAF and FIOP disappears or merges with the PIAF, further pressure experiments are in progress under higher pressures.

#### PD09

#### 29Si-NMR Study of antiferromagnet CeRh<sub>2</sub>Si<sub>2</sub> using single crystals Hironori Sakai, Yo Tokunaga, Shinsaku Kambe, Yuji Matsumoto, Tatsuma D Matsuda and Yoshinori Haea

Advanced Science Research Center, Japan Atomic Energy Agency, Japan

29Si-NMR Study of antiferromagnet CeRh<sub>2</sub>Si<sub>2</sub> has been performed using single crystals. Heavy fermion material CeRh<sub>2</sub>Si<sub>2</sub> is known to be an antiferromagnet with strong uniaxial magnetic anisotropy along the crystallographic c-axis under ambient pressure. From our measurements of Knight shifts and relaxation rates, the spin dynamics of this material have been investigated. At ambient pressure, as seen in the static susceptibility, the Ising-type antiferromagnetic spin fluctuations along the c-axis has been found to be enhanced. In this presentation, we will show our latest data and discuss about the anisotropy of hyperfine couplings and spin dynamics.

#### **PD10**

#### Bulk compressibility of orthorhombic YbFe<sub>2</sub>Al<sub>10</sub>-type CeRu<sub>2</sub>Al<sub>10</sub>

Yukihiro Kawamura<sup>1</sup>\*, Kazuki Matsui<sup>1</sup>, Keiichi Yamamoto<sup>1</sup>, Yusuke Hori<sup>1</sup>, Junichi Hayashi<sup>1</sup>, Keiki Takeda<sup>1</sup>, Chihiro Sekine<sup>1</sup> and Takashi Nishioka<sup>2</sup> <sup>1</sup> Muroran Institute of Technology, Japan

<sup>2</sup>Kochi University, Japan

Orthorhombic YbFe<sub>2</sub>Al<sub>10</sub>-type CeRu<sub>2</sub>Al<sub>10</sub> shows antiferromagnetic transition at T0–27 K[1]. This antiferromagnetic transition has various unusual features. One of them is reduced magnetic moment; The neutron experiment reveals that the magnetic moment of Ce in the ordered phase on CeRu<sub>2</sub>Al<sub>10</sub> is 0.34 µB/Ce[2]. Considering this reduced magnetic moment, T0 should be much lower. Another unusual feature is abrupt disappearance of antiferromagnetic transition by applying pressure[3]. Structural modulation might be occurred at around the pressure of disappearing T0. Thus, we performed synchrotron X-ray study at room temperature under pressure by using diamond anvil pressure cell and investigated bulk compressibility of CeRu<sub>2</sub>Al<sub>10</sub>. We revealed that lattice parameters a, b and c montonically decreases with increasing pressure, which indicates that there is no structural modification up to 9.2 GPa at room temperature. We also revealed that the decrease ratio of b is smaller than that of a and c.

[1] A. M. Strydom, Physica B 404 (2009) 2981-2984 [2] D. D. Khalyavin et al., Phys. Rev. B 82 (2010) 100405 [3] T. Nishioka et al., J. Phys. Soc. Jpn 78 (2009) 123705

#### PD11

#### Ultrasound measurements on the skutterudite compound SmOs<sub>4</sub>P<sub>12</sub>

Yoshiki Nakanishi<sup>1</sup>\*, Gen Koseki<sup>1</sup>, Dai Tamura<sup>1</sup>, Kohei Kurita<sup>1</sup>, Takeshi Saito<sup>1</sup>, Minoru Koseki<sup>1</sup>, Mitsuteru Nakamura<sup>1</sup>, Masahito Yoshizawa<sup>1</sup>, Masahito Yoshizawa<sup>1</sup>, Yuya Koyota<sup>2</sup>, Chihiro Sekine<sup>2</sup> and Takehiko Yagi<sup>3</sup>

<sup>1</sup>Iwate University, Japan

<sup>2</sup> Muroran Institute of Technology, Japan

<sup>3</sup> ISSP The University of Tokyo, Japan

We present an ultrasound study on the Sm-based filled skutterudite compound SmOs<sub>4</sub>P<sub>12</sub> for the first time. The measurements were performed on a polycrystalline sample using a phase-sensitive detection technique. The ternary compound SmOs<sub>4</sub>P<sub>12</sub> were prepared at high temperature and high pressures using a wedge-type cubic-anvil high-pressure apparatus. SmOs<sub>4</sub>P<sub>12</sub> exhibits antiferromagnetic ordering below 4.5 K. A clear decrease was observed at 4.5 K in the temperature dependence of both longitudinal and transverse elastic constants. However, less or no elastic softening toward the transition temperature was observed in the both elastic constants. Absence of the precursor behavior reminds us a multipolar ordering, possibly octupolar ordering observed in the isostructural system SmRu<sub>4</sub>P<sub>12</sub>, reported previously by our group. The transition is robust against the external magnetic field. To be more interesting, the elastic constants show a minimum around 180 K, possibly related to a rattling motion due to weakly bounded Sm ion in an oversized P cage. We argue a low-lying degenerated levels derived from 4f-multiplets ground state of Sm ion split by crystalline electric field effect in SmOs<sub>4</sub>P<sub>12</sub>. Furthermore, we discuss the phononic properties in comparison with other isostructral systems which include an ionic degrees of freedom of the rattling motion.

#### PD12

#### YbRh<sub>2</sub>Si<sub>2</sub>: Fermi surface and crystal-field splittings of a heavy-Fermion compound

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YbRh,Si, is a well known heavy Fermion system with a Kondo temperature of 25 K and a mean valence of 2,9. Due to the mixing of 4f14 and 4f13 configurations in the ground state, the 4f13 state is reproduced as a final state in a photoemission experiment allowing for direct observation of near ground state properties. Angleresolved photoemission spectra reveal a general dispersion of the 4f13 state with a Fermi-level crossing around the Γ-point. A clear wave-vector dependence of the crystal-field split states is observed that leads to a variation of the energy splittings across the Brillouin zone and even to an interchange of the sequence of states. The bulk Fermi surface reveals strong similarities with the one expected for a stable trivalent compound, but is slightly larger, reveals strong 4f character, and deviates from LDA results by a larger region without states around the F-point. These properties are well described in the framework of a simple fd-hybridization model.

## POSTER PRESENTATION

#### PD13

## Single crystal growth and various electronic states in Yb-based compounds

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To clarify the various electronic states of Yb compounds, we grew single crystals of several Yb compounds such as YbTin<sub>5</sub> (T : Co, Rh, Ir), YbCoGa,, YbT<sub>2</sub>Zh<sub>30</sub> (T : Co, Rh, Ir) and Yb<sub>2</sub>Pt<sub>2</sub>Pb by the flux method, and YbPdGe and YbPd<sub>2</sub>Al, by the Bridgman method. YbTin<sub>3</sub> and YbCoGa, are Pauli paramagnets, but the cyclotron effective mass of YbCoGa, is by four times larger than that of YbTin<sub>5</sub>, and YbCo<sub>30</sub> (T : Co, Rh, Ir) compounds are heavy fermion compounds. In fact, the electronic specific heat coefficient  $\gamma$  is 8000 mJ/k2+mol in YbCo<sub>2</sub>Zn<sub>30</sub>, very close to the quantum critical point. YbPdGe is an Ising-type ferromagnet with the Curie temperature TC = 11.4 K and a saturated magnetic moment  $\mu s = 1.7 \, \mu$ B/Yb, but possesses a relatively large value of  $\gamma = 150 \, \text{mJ}/$ K2-mol. Yb\_2Al, of which the analogue is a heavy fermion superconductor NpPd\_Al, is a well-known heavy fermion superconductor[1] is an antiferromagnet with a very small Neel temperature TN = 0.19 K.

[1] D. Aoki et al. : J. Phys. Soc. Jpn. 76 (2007) 063701.

#### **PD14**

## Physical properties under pressure in a heavy fermion superconductor $CeIrIn_{\scriptscriptstyle \rm E}$

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Celrh, is a heavy fermion superconductor with Tc = 0.4 K and the electronic specific heat  $\gamma$  = 700 mJ/(K2-mol) [1]. To investigate the nature and the pressure effect of the Kondo effect on Celrfn5, electrical resistivity  $\rho$  and magnetization M under pressure in Celrfn, have been measured using the single-crystalline Celrfn, which were grown by the In-flux method. The temperature dependence between 2.0 K and 300 K of the a-axis and e-axis electrical resistivity  $\rho$  of Celrfn, under pressure up to 8.0 GPa, which were obtained using a cubic anvil device, exhibits the typical dense Kondo behavior: on cooling from room temperature at P = 1.0 GPa, the  $\rho$ increases and shows a maximum of the coherence peak, and then goes down into the coherent Kondo state. We also find that the coherence peak moves to higher temperatures with pressure, indicating that the Kondo temperature of Celrfn, increases by applying the pressure. In the conference, we will also report the M data under pressure up to 1.5 GPa and compare the  $\rho$  data with our recent thermoelectric power S measurements, which exhibit a large positive value up to 90  $\mu$ /K with a sharp maximum in its temperature dependence. [2]

[1] C. Petrovic et al., Europhys Lett., 53, 354, (2001). [2] Y. Takaesu et al., J. Phys.: Conf. Series, 273, 012058, (2011).

#### PD15

#### Periodic Anderson model with correlated conduction electrons

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We investigate the so-called periodic Anderson-Hubbard model with the aim to understand the role of interaction between conduction electrons in the formation of the heavy-fermion and mixed-valence states [1]. We perform variational calculation using the Gutzwiller wave function and exact diagonalization of the Hamiltonian for short chains. The f-level occupancy and the renormalization factor of the quasiparticles are calculated as a function of the energy of the f-orbital for a wide range of the interaction parameters. The results obtained by the two methods are in reasonably good agreement for the periodic Anderson model. The agreement remains even for the Anderson-Hubbard model, except for the half-filled case. We find that due to this interaction the energy range of the bare f-level, where heavy-fermion behavior can be observed, shifts and widens. The Gutzwiller method indicates a robust transition from Kondo insulator to Mott insulator in the half-filled model, while this interaction enhances the quasiparticle mass when the filling is close to half filling.

[1] I. Hagymasi, K. Itai and J. Solyom Periodic Anderson model with correlated conduction electrons: variational and exact diagonalization study, arXiv::1106.4999

#### PD16

#### Thermoelectric study of the metamagnetic behavior in $YbCo_2Z_{20}$

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Despite the intensive studies on the metamagnetic behavior which has been found in a wide variety of heavy fermion materials such as CeRu<sub>2</sub>Si<sub>2</sub>, CeCu<sub>4</sub> and UPt<sub>2</sub> the origin remains under debate. The recent discovered heavy fermion compound YbCo<sub>2</sub>Zn<sub>20</sub> exhibits the metamagnetic behavior under unexpectedly small magnetic field  $\sim 0.6$  T, providing a new platform for elucidation of the metamagnetism. We have addressed this issue by means of the thermoelectric transport, which has been a promising probe to unveil the low-energy itinerant excitations and the electronic transformation across the metamagnetic transition. We measured the Seebeck and Nernst coefficients of YbCo<sub>2</sub>Zn<sub>20</sub> down to very low temperature. Remarkably, we found a huge Seebeck coefficient ~ -250 uV/K^2 under zero field, as an indication of the "super" heavy fermion state. Moreover, on passing through the metamagnetic field the Seebeck coefficient shows dramatic suppression, being distinct from that observed in CeRu<sub>2</sub>Si<sub>2</sub> in which the Seebeck coefficient is largely enhanced at the metamganetic field. In our presentation, we will discuss the thermoelectric response of YbCo<sub>2</sub>Zn<sub>20</sub> and its implication to the metamagnetism. The comparison to the other metamagnetic materials will be also made

#### **PD17**

#### Raman scattering study of the hidden order state of URu<sub>2</sub>Si<sub>2</sub> Jonathan Buhot<sup>1</sup>, Marie-aude Measson<sup>1\*</sup>, Yann Gallais<sup>1</sup>, Maximilien Cazavous<sup>1</sup>, Alain

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For over 25 years, researchers attempt to understand the microscopic origin of the Hidden Order that appears below T0=17.5 K in the heavy fermions compound URu<sub>2</sub>Si<sub>2</sub>. Mainly, optical measurements on URU<sub>2</sub>Si<sub>2</sub> are far-infrared reflectance experiments[1,2]. Only few Raman scattering studies have been performed [3,4] and only for c axis. We have studied URu<sub>2</sub>Si<sub>2</sub> for both axes, down to low temperature (2K). We have probed the phononic and electronic properties and their interaction. Softening of phonon modes together with Fano-shape were observed, implying that a phonon-electron coupling is at play. At low energy, electronic Raman signal shows either fluctuations or more conventional Drude shape for different symmetry. We have studied both signals through the Hidden Order temperature.

D.A. Bonn et al., PRL 61 1305 (1988), J. Levallois et al., PRB 84, 184420 (2011) S. Cooper et al., PRB 36 5743 (1987), D. Lampakis et al., PHYSICA B 378-380, 578 (2006)

#### PD18

## The evolution of superconductivity and magnetism in Pd-doped CeRhIn<sub>5</sub> and Ce<sub>2</sub>RhIn<sub>8</sub>

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CeRhIn<sub>1</sub> and Ce<sub>2</sub>RhIn<sub>8</sub> are heavy fermion antiferromagnets. Unconventional superconductivity can be induced by hydrostatic pressure or doping [1,2]. By contrast, Ce<sub>2</sub>PdIn<sub>8</sub> is a heavy fermion superconductor (Te=0.7K) at ambient pressure. We have examined the influence of Pd doping in single crystals of CeRh<sub>1</sub>, aPdxIn<sub>8</sub> and Ce<sub>2</sub>Rh<sub>1\*4</sub>PdxIn<sub>8</sub>. Since the synthesis of Ce<sub>2</sub>PdIn<sub>8</sub> is connected either to intergrowth of CeIn<sub>1</sub> or variation of critical temperature, growth conditions were studied by differential thermal analysis. By this, we have succeeded in obtaining high-quality single crystals free of CeIn<sub>3</sub> and with a critical temperature Te-0.7K. The crystals have been subjected to specific heat, resistivity and AC susceptibility measurements. In addition, resistivity experiments under hydrostatic pressures up to 3GPa were performed. Surprisingly, at ambient pressure the Neel temperature of CeRh<sub>1\*</sub>PdxIn<sub>8</sub> has only slightly decreased with increasing Pd content. And moreover, the pressure influence on magnetism in CeRh<sub>0\*8</sub>Pd<sub>10\*3</sub>In<sub>5</sub> does not significantly differ from CeRhIn<sub>5</sub>. Contrary, the superconducting transition increases significantly compared to its 115-counterpart. The less two-dimensional Ce<sub>2</sub>Rh<sub>1\*</sub>PdxIn<sub>8</sub> behaves different. Here, the antiferromagnetic ordering is immediately lowered upon Pd substitution and vanishes completely at x-0.35 already. The obtained magnetic field-temperature phase diagrams will be discussed in context of the physical properties of other CenTIn3n+2 compounds.

[1] H. Hegger et al., Phys. Rev. Lett. 84, 4986 (2000). [2] M. Nicklas et al., Phys. Rev. B 67, 020506 (2003).

#### PD19

## Magnetic phase diagram of the new heavy fermion compound Ce<sub>2</sub>PtIn<sub>81</sub>

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CenTmln<sub>3ter2m</sub> (n=1,2; m=1; T=transition metal) type compounds are subject of intense interest in the condensed-matter community [1]. They are predestinated to investigate the interplay between magnetism and superconductivity. Recently, two new compounds from this heavy fermion family have been found. Ce<sub>2</sub>PdIn<sub>4</sub> is an ambient pressure superconductor while CePt<sub>2</sub>In<sub>7</sub> orders antiferromagnetically [2]. We report on the existence of a new compound Ce<sub>2</sub>PdIn<sub>4</sub>, Similar to Ce<sub>2</sub>PdIn<sub>4</sub>, the synthesis of Ce<sub>2</sub>PdIn<sub>4</sub> is rather complicated. However, from our studies of solution growth of Ce<sub>2</sub>PdIn<sub>8</sub> and the support by differential thermal analysis, optimal conditions for growing Ce<sub>2</sub>PtIn<sub>4</sub> were deduced. Single-crystal X-ray diffraction confirmed that Ce<sub>2</sub>PtIn<sub>8</sub> crystallizes in Ho<sub>2</sub>CoGa,+type structure with lattice parameters a=4.699A and c=12.185A. We will present specific heat, resistivity (ambient and under hydrostatic pressure) and magnetization measurements. Ce<sub>2</sub>PtIn<sub>9</sub> orders magnetically below 2.1K. A second transition, likely into a magnetic commensurate ordering, is observed just below at 2K. Contrary to Ce<sub>2</sub>Rhln<sub>8</sub>, the magnetic transitions shift to lower temperatures in applied magnetic field along both principal crystallographic axes suggesting different character of magnetic ordering. Specific heat measurements reveal an enhanced Sommerfeld coefficient ( $\gamma$ -740 m//mol.K2). The magnetic field-temperature phase diagram will be discussed in the context of superconductivity and magnetism evolution in related compounds.

[1] Ch. Pfleiderer, Rev. Mod. Phys. 81, 2009 [2] E. D. Bauer, et al., Phys. Rev. B 81, 180507 (2010)

#### PD20

## Coexistence and competition of superconductivity, magnetism and charge density waves in rare-earth tri-telluride TbTe<sub>3</sub>

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Rare-earth tri-telluride RTe3 (R = Y, La-Sm, Gd-Tm) forms a weakly orthorhombic crystal lattice, consisting of alternating double layers of corrugated double RTe layers and nominally square-planar Te sheets. The RTe layers are responsible for magnetism while the square Te sheets contribute to strongly anisotropic conduction. Interestingly, the two-dimensional conduction bands originating from Te-5px and 5py orbitals give rise to Fermi surface nesting, and this hidden one-dimensionality leads to a charge density waves (CDW) with a high transition temperature. Because of this intriguing feature, RTe3 attracts great interest in recent years. TbTe3 has localized 4f electrons, and exhibits successive phase transitions of magnetic origin, which all coexist with the CDW. Surprisingly, superconductivity emerges under high pressure. The question is whether the superconductivity coexists or competes with the other long range orders. Therefore, we investigated the long-range orders by the electrical resistivity and ac magnetic susceptibility at high pressures up to 8 GPa and at low temperatures down to 2 K using a constant-load cubic-anvil cell with single crystalline samples that were grown from Te-fluxes. We found that the antiferromagnetism survives even at the highest pressure of the present experiment. We also confirmed the emergence of the superconductivity at high pressures.

#### PD21

## Magnetic properties of $Ce_{3}Rh_{4}Sn_{13}$ and $Ce_{3}Co_{4}Sn_{13};$ a comparative study

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We report on the electronic structure (XPS) and basic thermodynamic properties of Ce<sub>3</sub>Rh<sub>4</sub>Sn<sub>13</sub> and Ce<sub>3</sub>Co<sub>4</sub>Sn<sub>13</sub> and the reference La-based compounds. The both compounds were classified as a heavy fermion materials with extremely large low temperature electronic specific heat C/T [1] of about 4 J/molK2. The data show that either Ce<sub>3</sub>Rh<sub>4</sub>Sn<sub>13</sub> or Ce<sub>3</sub>Co<sub>4</sub>Sn<sub>13</sub> are very sensitive to applied magnetic fields and exhibit a cross-over from magnetically correlated state to a single impurity state, when the applied magnetic field increases. We noted that the low temperature properties are dependent on the stoichiometry of the sample. We also present and discuss the low temperature properties of Ce<sub>3</sub>Co<sub>4</sub>Sn<sub>13</sub> doped with La.

1] U. Kohler et al. J. Phys.: Condens. Matter 19 (2007) 386207

#### PD22

## Thermoelectric properties of Kondo semiconductor CeRu<sub>4</sub>As<sub>12</sub> prepared under high pressure

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The filled skutterudite compound CeRu<sub>4</sub>As<sub>12</sub> exhibits a hybridization gap insulating state, with a small activation energy $\Delta$ /kB of 50K [1]. We report further results for CeRu<sub>4</sub>As<sub>12</sub> prepared under high pressure (4GPa). Thermoelectric power and thermal conductivity measurements have been performed on this material. The temperature dependence of thermoelectric power for CeRu<sub>4</sub>As<sub>12</sub> shows two peaks (around 80K, 200K). The phenomena could be related to Kondo behavior at high temperature and hybridization-gap-formation process at low temperature.

[1] C. Sekine et al., J. Mag. Mag. Mater., 310 (2007) 260.

#### PD23

## Single crystal growth and physical properties of $UT_2Al_{20}$ (T=transition metal)

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Many anomalous physical properties have been observed in  $RT_2X_{20}$  (R=Rare metal, T=transition metal, X=Al, Zn) system. For example, the quadrupolar order takes place in  $PrT_2X_{20}$  system, and the extreme heavy fermion state takes place in  $YbCo_2Zn_{20}$ . In the uranium analogue, for example,  $UT_2Al_{20}$  is worth investigating because the extended 5f wave function in uranium would lead to peculiar features due to strong hybridization. Therefore we have studied single crystal growth and physical properties of  $UT_2Al_{20}$ . The sample preparation is performed by using Al-self flux method. We have succeeded to prepare the single crystal of  $UT_2Al_{20}$  (T=Ti, Nb). We are going to present about the physical properties of the prepared samples in the Conference.

#### PD24

#### Anisotropic c-f hybridization in a Kondo semiconductor CeFe<sub>2</sub>Al<sub>10</sub>

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We report anisotropic properties of a Kondo semiconductor CeFe<sub>2</sub>Al<sub>10</sub>. This compound is the non-magnetic counterpart of Kondo semiconductors CeT<sub>2</sub>Al<sub>10</sub> (T=Ru and Os) that exhibit an unusual antiferromagnetic transition at about 28 K[1-5]. The magnetic susceptibility for B/a,  $\gamma_a(T)$ , shows a broad peak at 70 K due to the strong c-f hybridization. The anisotropy in  $\gamma$ ,  $\gamma_a \gamma_c \sim \gamma_b$ , is in common with those for T = Ru and Os. Electrical resistivity  $\rho(T)$  has a broad maximum at 45, 35 and 110 K for 1/a, 1/b and 1/c, respectively. Below 20 K,  $\rho$ 's along all axes show an uptum due to the opening of a pseudo gap. Thermopower S(T) also shows a maximum at 190, 150 and 200 K for a-, b- and c-axis, respectively. The temperatures at the maximum in  $\rho$  and S suggest the anisotropy of c-f hybridization strength Vc>Va>Vb in CeFe<sub>2</sub>Al<sub>10</sub> as reported in the study of optical conductivity[6]. Furthermore, only Sc(T) changes the sign from positive to negative at 100 K on cooling and shows a minimum at 50 K. The negative minimum vas also observed in Sc(T) for T=Ru and OS (S]7]. This behavior indicates the development of antiferromagnetic correlation along the c-axis in CeT2Al10.

[1] Y. Muro et al. J. Phys. Soc. Jpn. 78 (2009) 083707. [2] T. Nishioka et al. J. Phys. Soc. Jpn. 78 (2009) 123705. [3] Y. Muro et al. Phys. Rev. B 81 (2010) 214401. [4] D. D. Khalyavin et al. Phys. Rev. B 82 (2010) 100405R. [5] D. T. Adroja et al. Phys. Rev. B 82 (2010) 104405. [6] S. Kimura et al. J. Phys. Soc. Jpn. 80 (2011) 033702. [7] H. Tanida et al. J. Phys. Soc. Jpn. 79 (2010) 035709.

#### PD25

## Cu-NMR studies of heavy-Fermion compound CeCu<sub>6</sub> under high magnetic fields

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It is known that the intermetallic compound CeCu<sub>6</sub> is a heavy fermion system showing a huge specific heat coefficient  $\gamma \sim 1600$  mJ / K2 mol, without magnetic order or superconductivity at low temperatures down to 10 mK[1,2]. The electric resistivity shows characteristic temperature dependence of ? logT above 10 K, which is ascribed to the Kondo effect, and it has a maximum around 10 K. No detailed NMR results have been reported yet because of a drawback of the complicated Cu NMR spectrum due to a peculiar orthorhombic crystal structure. We have carried out Cu-NMR measurements to clarify the heavy fermion state of a single crystal CeCu<sub>6</sub> under several magnetic fields. We have measured Knight shift and the nuclear spin-lattice relaxation time T1 under magnetic fields up to 15 T along the a-axis and c-axis in the temperature range of 100-1.5 K. At low temperatures, the Cu-NMR Knight shifts are suppressed as magnetic field increases. Simultaneously, 1/T1 is suppressed. These features are understood by the concept that the heavy fermion state in CeCu<sub>6</sub> is formed through Kondo effect. We will report detailed field dependence of the heavy fermion state of CeCu<sub>6</sub> for the first time.

[1] Y. Onuki, et al.: JPSJ 53 (1984) 1210 [2] Y. Onuki, et al.: JMMM 63-64 (1987) 281

#### PE01

## The underscreened Anderson lattice : A model for uranium compounds.

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We present a model based on the Underscreened Anderson Lattice (UAL) with two 5f electrons per site to account for the Kondo-ferromagnetism co-existence observed in some actinide compounds. The UAL model allows a description of the decrease of the number of f-electrons and we present a calculation which explains the pressure dependence observed in some Uranium compounds and particularly in Uranium monochalcogenides. On the other hand, we have developed a novel type of phase transition within the UAL model and we suggest that this model might describe the "Hidden Order" transition in URu<sub>2</sub>Si<sub>2</sub>. The gaps that appear in the electronic dispersion relations of bands of different orbital character are in agreement with experimental photoemission results

#### PE02

## Transport properties of intermetallic compounds RCoGe<sub>2</sub> (R = Ce and La)

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The electronic and magnetic properties of the cerium-based ternary intermetallic compound compounds with the general formula CeMSi<sub>2</sub> (M = transition metals) has been of considerable interest recently. Among those compounds, CeCoGe2 is of particular interest due to its peculiar physical properties, such as it has a rather high Kondo temperature TK ~ 250 K and it is the first Kondo system clearly interpreted by the Cogblin-Schrieffer model with i = 5/2. In this study, measurements of the electrical resistivity. Seebeck coefficient, and thermal conductivity on CeCoGe<sub>3</sub> have been performed in the temperature range 10 - 300 K to investigate the electronic structure. For comparison the nonmagnetic counterpart LaCoGe, has also been studied. It is found that CeCoGe, exhibits a broad maximum in Seebeck coefficient at about 75 K, at which the sudden drop in electrical resistivity occurs. This is a typical behavior commonly seen in Ce-based Kondo lattice compounds. A theoretical attempt was made and found that the electrical ransport and thermoelectric properties of CeCoGe<sub>2</sub> can be well described by a two-band model with reliable physical parameters. On the other hand, both electrical resistivity and Seebeck coefficient show a typical metallic-like behavior for the nonmagnetic LaCoGe<sub>2</sub>

#### PE03

#### **The lattice Kondo effect - A fabric for superconducting correlations?** Oliver Bodensiek<sup>1</sup>\*, Thomas Pruschke<sup>1</sup> and Rok Zitko<sup>2</sup>

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Although the simple Kondo lattice model constitutes one of the paradigms for understanding the physics of heavy-fermion materials, a reliable theoretical investigation of superconductivity in this model is still lacking. We present results based on the dynamical mean-field approximation in combination with the numerical renormalization group as impurity solver. While superconducting order is commonly not expected for a correlated lattice model without additional bosonic degrees of freedom, we observe strong superconducting order in the plain Kondo lattice model away from half filling. The possible origin of this superconducting order is discussed in view of the frequency dependence of the order parameter. In addition, close to half filling we find this ordered phase to appear most pronounced in the vicinity of the antiferromagnetic quantum critical point. Since a characterization of quantum criticality within dynamical mean-field theory is not possible, we rather discuss questions of principle such as the actual nature of the ordered states, that is, local moment or heavyfermion physics. The latter issue is of considerable interest addressing the question whether the ordered states stay in direct competition or even cooperate with the coherent Kondo effect in the lattice.

#### **PE04**

## Lifshitz transition with interactions in high magnetic fields: application to CeIn<sub>3</sub>

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The Neel ordered state of CeIn<sub>3</sub> is suppressed by a magnetic field of 61 T at ambient pressure. There is a second transition at ~45 T, which has been associated with a Lifshitz transition. Skin depth measurements indicate that the transition is discontinuous as T >> 0 and that the transition has a weak pressure dependence until it merges with the Neel transition. We study the effects of Landau quantization and interaction among carriers on a Lifshitz transition. The Landau quantization leads to quasi-one-dimensional behavior for the direction parallel to the field. Repulsive Coulomb interactions give rise to a gas of strongly coupled carriers. The density correlation function is calculated for a special long-ranged potential [1]. For the lowest Landau level the problem can be mapped onto the interacting one-dimensional electron gas. It is concluded that in CeIn<sub>3</sub> (a) an electron or hole pocket is being emptied as a function of field and (b) in the ground state the pocket is emptied in a discontinuous fashion. This discontinuity is gradually smeared by the temperature in agreement with the skin depth experiments. Work supported by the Department of Energy under grant DF-FG02-98FB45707

[1] P. Schlottmann, Phys. Rev. B 83, 115133 (2011).

#### **PE05**

#### Dynamical mean-field theory of indirect exchange between magnetic adatoms on metallic surfaces

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Two magnetic impurities on a metallic substrate surface experience the Ruderman-Kittel-Kasuya-Yosida (RKKY) exchange interaction. This indirect non-local magnetic exchange competes with the local Kondo effect. While the latter is described in principle exactly within single-site dynamical mean-field theory, the effects of the RKKY coupling are taken into account approximately only. Here, this is demonstrated by comparing the DMFT results with numerically exact data obtained by the densitymatrix renormalization group for a one-dimensional model with two Anderson impurities. With the two-site DMFT, we also benchmark a simplified DMFT variant. Varying the inter-impurity distance d and the local exchange coupling J, different parameter regimes with dominating Kondo physics or with dominating RKKY interaction are studied. For the Kondo regime, our calculations show that DMFT is able to reliably describe magnetic nanostructures on metallic surfaces. The prospects of DMFT (and two-site DMFT) as a reliable approach to study nanomagnetism of more complex RKKY-interacting systems are discussed.

#### **PE06**

#### Kondo effect near the Van Hove singularity in biased bilayer graphene Stanislaw Lipinski\* and Damian Krychowski

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We study Kondo impurity in a bilayer graphene with Bernal stacking for adatom adsorbed on the outside of the bilayer or intercalated between the two graphene layers. The density of states (DOS) close to the gap diverges as the square root of energy - Van Hove singularity (VHS). The system is modeled by the generalized Anderson hamiltonian with bilayer described by Hubbard model. The many-body problem is studied within the mean field Kotliar-Ruckenstein slave boson approach for SU(2) and SU(4) symmetries. We examine the impact of tuning the Fermi level (EF) through the VHS. Increase of the DOS at EF increases hybridization of the impurity, what moves the system towards mixed valence regime or even completely suppresses the many body peak of DOS. Additionally a competitive influence of magnetic instability is analyzed. Crossing of the Fermi level with a Van Hove singularity leads to weak ferromagnetism of the bilayer graphene with strong spin polarization at EF. The spin symmetry is broken, what for the case of spin-orbital degeneracy results in a crossover from SU(4) Kondo effect to SU(2) and in the case of only spin degeneracy it leads to a strong suppression of SU(2) Kondo effect.

#### **PE07**

#### Indirect exchange between magnetic atoms on surfaces: From two impurities to diluted chains

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Magnetic atoms on metallic surfaces are subject to the competition between Kondo screening and indirect magnetic (RKKY) exchange. To study the crossover of the magnetic properties from the case of a single or two magnetic adatoms to the limit of a dense chain of indirectly coupled adatoms, we model the metal surface as a onedimensional tight-binding system and apply the density-matrix renormalization group (DMRG) to the resulting Kondo model. The case of three magnetic atoms turns out to be particularly interesting as finite-size gaps are found to qualitatively affect the balance between Kondo physics and RKKY. With the transition from the diluted case with a few magnetic atoms to the dense chain with all atoms RKKY-coupled ferroor antiferromagnetically, magnetic long-range order becomes possible. The resulting phase diagram is studied in the weak (RKKY) but also in the strong-coupling (Kondo) limit. Besides DMRG we employ real-space dynamical mean-field theory (R-DMFT) in addition and thereby extend the results to higher dimensions.

#### **PE08**

#### **ESR study of hybridization in some undoped Yb-based alloys** Vladimir Ivanshin\*, Tatyana Litvinova and Eduard Gataullin

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The hybridization between a wide conduction band and correlated f-electrons localized at lattice sites is investigated using method of electron spin resonance (ESR) in some dense heavy-fermion compounds. ESR signals arise here both from the local moment spins and conduction electrons (CE) due to the strong electronic correlations because of the hybridization between 4f-electrons and CE in the presence of ferromagnetic (FM) fluctuations [1]. We review very recent ESR experiments in several undoped Ybbased intermetallics [2, 3] and discuss the role of chemical composition, Kondo effect, anisotropy, and RKKY interactions for the ESR linewidth. It is shown that an increase of the Kondo interaction and FM correlations leads to a stronger hybridization and an essential ESR line narrowing. Possible hybridization mechanisms are proposed.

 P. Wolfle, E. Abrahams, Ann. Phys. (Berlin) 523 598 (2011) [2] V.A. Ivanshin et al., J. Phys.: Conf. Ser. 324 012019 (2011) [3] T. Gruner et al., Phys. Rev. B 85 035119 (2012)

#### **PE09**

## ESR study of influence of anionic and cationic substitutions in EuB6 on the magnetic phase separation

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We performed a comparative study of the features of the magnetic phase separation in europium hexaboride with anionic and cationic substitution with donor impurities. The ESR (electron spin resonance) were measured on the single crystal plates  $EuB_{-}[6x_1C_{-}[x]]$  and  $Eu_{-}[1-x_1Gd_{-}[x_3]B_{-}[6]$  with x < 0.07 in X-band in the temperature range T = 10-300. Both types substitution keep integrity of the magnetic sublattice and they add one electron per the impurity atom in the conduction band. of  $EuB_{\mu}$  The magnetic field was oriented perpendicular to the sample plane and along various crystal axis. In all case we observed the magnetic phase separation and phase transitions. However, the ratio of the magnetic phases and their behavior is slightly different for substitution boron in carbon and europium gadolinium. In the latter case, there was the coexistence of ferromagnetic and antiferromagnetic phases. The obtained ESR data are discussed in frame Falicov-Kimball model the hybridization of Kondo-like electrons with two band (valence and conduction band) timerant electrons and formation of Kondo-like electron states. Work was supported by Grants of RFBR and Presidium of RAS.

[1] Altshuler T S, Goryunov Y V, Shitsevalova N Y, Dukhnenko A, JETP Lett. 88, 779 (2008) [2] Kunes J, Picket W E, Phys.Rev. B 69, 1651 (2004) [3] Falicov L M, Kimball J C, Phys.Rev.Lett. 22 997 (1969) [4] Kugel K I, Rakhmanov A L, Sboychakov A O, Phys.Rev.Lett. 95, 267210 (2005)

#### **PE10**

High pressure synthesis of novel boron-cage compounds RB<sub>12</sub> (R=Gd, Sm)

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Rare-earth compounds with cage structures recently attract much attention because of possibility of new heavy fermion mechanism using the lattice vibration [1]. Rare-earth borides such as RB<sub>6</sub> and RB<sub>12</sub> also have such cage structures composed by boron. However, RB<sub>12</sub> series are known from R = Tb to Lu. Cannon has succeeded in high-pressure synthesis of GdB<sub>12</sub> in pressure of 6.6 GPa and 2,400 C, and reported lattice parameter a = 7.523 A in the UB<sub>2</sub>-type structure [2]. We have also succeeded in high-pressure synthesis of polycrystalline GdB<sub>12</sub> and SmB<sub>12</sub>. Magnetic susceptibility of GdB<sub>12</sub> shows a kink at 38 K indicating of antiferro-magnetic ordering. Highfield magnetization of GdB<sub>12</sub> shows the saturation moment is 7 µB. All the reported Neel ordering temperatures of RB<sub>1</sub> well agree to the values estimated from the trivalent rare-earth dependence of de Gennes factor. From the x-ray powder diffraction of SmB<sub>12</sub>, this compound also has the UB<sub>12</sub>-type structure with lattice parameter of a = 7.541 A. Magnetic susceptibility shows two anomalies at 9 K and 19 K different from other RB<sub>2</sub>.

[1] T. Hotta and K. Ueda: Phys. Rev. B 67 (2003) 104518. [2] J. F. Cannon, et al.: J. Less-Common Metals 56 (1977) 83.

#### PE11

## Magnetic anisotropy of tetragonal rare-earth compounds RRu<sub>2</sub>Al<sub>2</sub>B (R: rare-earth metals)

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Recently, we succeeded to synthesize polycrystalline samples of new tetragonal rare-earth compounds CeRu<sub>2</sub>Al<sub>2</sub>B and PrRu<sub>2</sub>Al<sub>2</sub>B and reported magnetic properties of these compounds [1]. Both compounds show successive phase transitions of an antiferromagnetic and a ferromagnetic ordering. The Neel and Curie temperatures for CeRu<sub>2</sub>Al<sub>2</sub>B are 14.3 K and 13 K, respectively, and those for PrRu<sub>2</sub>Al<sub>2</sub>B are 26 K and 11 K, respectively. The values of magnetization vary when the angle between sample direction and applied magnetic field is changed even in the polycrystalline samples. This fact implies that the grains of polycrystalline samples are oriented partially and the variation of magnetization values is due to the magnetic anisotropy. In this study, we will report on the results of magnetization and magnetic-susceptibility measurements of grain-oriented powder samples of RRu<sub>2</sub>Al<sub>2</sub>B (R: rare-earth metals) and will discuss the magnetic anisotropy.

[1] E. Matsuoka, Y. Tomiyama, H. Sugawara, T. Sakurai, and H. Ohta, J. Phys. Soc. Jpn. 81 (2012) 043704.

#### **PE12**

Collective magnetic resonance mode in CeB<sub>6</sub>

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<sup>2</sup> Institute for Problems of Materials Science NAS, Ukraine

The observation of the electron spin resonance (ESR) in low temperature phase (phase II) of CeB<sub>6</sub> has risen up new questions about the nature of its ground state [1]. In the present work we study the effects of La doping on the ESR parameters in Ce<sub>1</sub>,  $_xLa_xB_6$ . Experiments in the doping range 0 < x < 0.1 have shown that even small La concentrations  $x \sim 0.03$  cause strong broadening as well as suppression of the ESR line. Such behavior is accompanied by a tendency of g-factor increase from g $\approx 1.6$  in CeB<sub>6</sub> to higher values with La doping. The resonance line gets unobservable with further concentration increase up to  $x \sim 0.05$ . At the same time the phase border between the phase II and the paramagnetic phase I observed as a kink in microwave conductivity remains almost independent on the La concentration in allows suggesting the collective origin of the electron spin resonance mode.

[1] S. V. Demishev et. al., Phys. Rev. B 80, 245106 (2009)

#### **PE13**

## Transport and magnetic properties of CeFe<sub>4</sub>Sb<sub>12</sub> synthesized under high pressure

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The filled skutterudite compounds have attracted much attention, since they exhibit a variety of novel features such as the heavy fermion behavior, hybridization-gap semiconducting behavior, and unconventional superconductivity [1]. The filling fraction of rare-earth site plays the important role for the ground state in certain compounds, such as  $PxFe_4Sb_{12}$ , in which the ferromagnetic ground state for x-0.8 is changed to the non-magnetic one by increasing to x-1 [2]. In this work, we have succeeded in synthesizing the filled skutterudite  $CexFe_5Sb_{12}$  with the Ce-site filling fraction x-1 under high pressure, and found the semiconducting behavior in this compound.

[1] H. Sato, H. Sugawara. Y. Aoki and H. Harima: Handbook of Magnetic Materials, Vol. 18, "Magnetic Properties of Filled Skutterudites' ed. by K.H.J. Buschow. (Amsterdam, Elsevier), p.1-110, 2009. [2] K. Tanaka, Y. Kawahito, Y. Yonezawa, D. Kikuchi, H. Aoki, K. Kuwahara, M. Ichihara, H. Sugawara, Y. Aoki and H. Sato: J. Phys. Soc. Jpn. Vol. 76 (2007) 103704.

#### **PE14**

#### Interplay of Kondo effect and spin orbit coupling

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Physics near quantum critical point (QCP) has been an interesting unresolved problem for many decades now. While the ground states on either sides, away from QCP are reasonably well understood, the physics in the vicinity of QCP is a long standing puzzle. The question is, how the quantum critical fluctuations affects the electronic excitations and evolves as one goes across QCP. We studied the electronic structure of Ce<sub>2</sub>Rh<sub>1</sub>, Co<sub>x</sub>Si<sub>3</sub> with varying x using high resolution photoemission spectroscopy. Co substitution at the Rh sites in antiferromagnetic Ce<sub>2</sub>RhSi<sub>3</sub> leads to transition to a Komdo system, Ce<sub>2</sub>CoSi<sub>3</sub> via QCP at x = 0.6. High resolution spectra reveal signature of Kondo resonance feature (KRF) and its spin orbit split component (SOC) in the whole composition range indicating finite Kondo temperature scale at QCP and applicability of spin density wave picture at the approach to QCP. Interestingly, the intensity ratio of the Kondo resonance feature and its spin orbit split component, KRF/SOC gradually increases with the decrease in temperature suggesting importance of spin-orbit interactions in this regime.

#### PE15

Insulator-to-metal transition and magnetism of potassium metals loaded into regular cages of zeolite LSX

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Zeolite low-silica X (LSX) crystals have an aluminosilicate framework with the regular nanospace of supercages and beta-cages. They are connected by the double diamond structure. Many K-cations are distributing among negatively charged framework in K-type LSX. 4s-electrons of guest alkali atoms are shared with K-cations of zeolite. Loading density of guest K atoms per unit cage, n can be controlled from 0 to 9. At n < 2, simples are nonmagnetic and insulating[1]. When n is just above 2, number density of magnetic moments suddenly increases and electrical resistivity decreases as semiconductive materials. For n > 6, electrical resistivity shows metallic increase in resistivity is observed at low temperatures likely Kondo semiconductor. In materials with the strong electron-phonon interaction, excess electrons are self-trapped quantum mechanically at the deformation potential well[2]. These immobile electrons are small polarons with magnetic moments. If the electron-phonon interaction is much stronger than the repulsive interaction between two electrons, nonmagnetic small bipolarons, at babilized. Observed novel properties are explained by small bipolarons, small polarons, metallic large polarons, and Kondo-like mixing between metallic electrons in survercages and localized electrons, metallic large polarons,

[1] Y. Ikemoto, T. Nakano, M. Kuno, Y. Nozue and T. Ikeda, J. Mag. Mag. Mat. 226-230, 229 (2001).
 [2] Y. Toyozawa, Prog. Theor. Phys. 26, 29 (1961).

#### **PE16**

1/(N-1) expansion for a finite U Anderson model with an SU(N) symmetry

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We apply recently developed 1/(N-1) expansion to a particle-hole asymmetric SU(N) Anderson model with finite Coulomb interaction U. This approach is completely different from the conventional 1/N expansion, or the non-crossing approximation. In our approach the factor N-1 corresponds to the number of interacting orbitals, excluding the one prohibited by the Pauli principle. To leading order in 1/(N-1) it describes the Hartree-Fock random phase approximation (HF-RPA), and the higher-order corrections describe systematically the fluctuations beyond the HF-RPA. The low-energy local Fermi-liquid behavior is described properly with this approach. We show that the next-leading order results of the renormalized parameters agree closely with the numerical renormalization group results at N=4 in a wide range of the electron fillings, from the empty to fully occupied orbitals, specifically in the mixed-valence regions. This ensures the reliability of the next-leading order results for N > 4. Our expansion scheme uses the standard Feynman diagrams, and thus has wide potential applications. It can be applied to the Keldysh formalism, and to lattice systems such as the Hubbard model.

[1] A. Oguri, R. Sakano, and T. Fujii, Phys. Rev. B 84, 113301 (2011). [2] A. Oguri, arxiv:1112.6051.

#### PE17

## Formation of the Kondo resonance band in CeCoGe<sub>2</sub>: DFT+DMFT approach

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We have investigated the formation of the Kondo resonance (KR) band in the heavy fermion  $CeCoGe_2$  using the combined approach of the density functional theory (DFT) and the dynamical mean field theory (DMFT). The low temperature (T) spectral function shows the dispersive KR states in momentum space, which is well consistent with the experimental observation. The size of the induced hybridization gap is rather insensitive to T. During the evolution from the spdbands at highTto the dispersive KR bands at low T, whose topologies are different each other, we have found the existence of kinks in the spectral function near EF. We suggested that the kink can be intrinsically observed in measuring spdconduction bands during the formation of the fully coherent KR band.

#### PE18

**Charge ordering in the Kondo lattice model at quarter filling** Junki Yoshitake\*, Takahiro Misawa and Yukitoshi Motome Department of Applied Physics, The University of Tokyo, Japan

The Kondo lattice model is one of the fundamental models for heavy-fermion systems, where exchange interactions between itinerant electrons and localized spins play an important role. Among many different phases described by this model, emergence of charge-ordered state [1,2] is interesting possibility, because this model does not include any bare inter-site repulsions between electrons. The possibility was first pointed out by a perturbation expansion in the strong Kondo coupling limit at quarter filling in 1D [1], and recently examined by the dynamical mean-field theory in infinite dimensions [2]. However, it remains unclear whether the charge-ordered state appears in two and three dimensions and what type of magnetic order is accommodated in the chargeordered state. To clarify these issues, we investigate both the ground state and finitetemperature properties of the quarter-filled Kondo lattice model on a square lattice by complementarily using the variational Monte Carlo simulation and cluster dynamical mean-field theory. We found that charge-ordered state appears in a wider region than the previous result in infinite dimensions, and that the antiferromagnetic order between charge poor sites is accommodated in the charge-ordered state. We discuss the stabilization mechanism of charge ordering by carefully analyzing the electronic and magnetic properties.

[1] J. E. Hirsch, Phys. Rev. B 30, 5383 (1984). [2] J. Otsuki et al., J. Phys. Soc. Jpn. 78, 034719 (2009).

#### PE19

## Angular-dependent magnetoresistance of the filled skutterudite CeOs<sub>4</sub>As<sub>12</sub>

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The filled skutterudite CeOs<sub>4</sub>As<sub>12</sub> compound is a narrow-gap semiconductor whose low-temperature properties originate from a hybridization between 4f and conduction electrons. Nonmagnetic CeOs<sub>4</sub>As<sub>12</sub> is the cubic system showing a nodal structure of hybridization energy gap(s) as evidenced from low-temperature (T > 0.08 K) and high-magnetic field (B < 14 T) studies of a directional dependence of the electrical resistivity  $\rho$ (T): At T < 3 K and for j  $\parallel$ B, we found remarkable dissimilarities along the [001] and [111] directions, indicative of an anisotropic suppression of energy gap(s). Additionally, differences observed between the transverse and longitudinal magnetoresistivity cannot be ascribed to the Lorentz force and thus, provide a further evidence for magnetic-field-induced anisotropy of hybridization gap in CeOs\_4As\_{12}. High-quality single crystals of CeOs\_4As\_{12} enable studies of a change of Fermi surface via the angular-dependent magnetoresistivity both in the field-induced metallic as well as semiconducting state. Our experiments with the use of the commercial one-axis rotator (PPMS) provide an additional evidence for strongly correlated electron phenomena in CeOs\_4As\_{12}.

#### PF01

## Spin and charge correlation of electrons at variation of interorbital Coulomb interaction

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Of great interest for the researches is the ions with orbital degeneration and electron shell with filling more than half. The aim of our work is to evaluate spin and charge correlation at variation of interorbital Coulomb interaction and competition of electron hopping parameters between nearest sites on the same and different orbitales arising as a result of variation of orbitales filling. Variation of charge gap may be also cause a radical change in the electronic density and spin correlation functions. We calculated electron spectrum in the Hubbard model with electron filling n=1.5 and following parameters: energy level of ions, the matrix element of the hopping between the nearest sites and orbitales. U and V is the Coulomb repulsion of the electrons on site and between orbitales. We calculated the eigenvalue spectrum and the corresponding state vectors by exact diagonalization for small cluster, based on which we determined the spin correlation functions on the longitudinal and transverse spin components, correlator of electron density. As a results, critical parameters Vc of the interorbitales Coulomb interaction and hopping integral are found at which reveal change in the value and sign of exchange interaction and in the correlator of electron density.

#### PF02

## Superfluid state of repulsively interacting three-component fermionic atoms in optical lattices

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We investigate the superfluid state of repulsively interacting three-component (color) fermionic atoms in optical lattices. We first discuss the effective interactions using diagrammatic approaches. We then investigate the superfluid properties using complementary two methods: the dynamical mean-field theory with a perturbative approach and the self-energy functional approach. We show that when the anisotropy of the three repulsive interactions is strong, atoms of two of the three colors form Cooper pairs and atoms of the third color remain a Fermi liquid. An effective attractive interaction is induced by density fluctuations of the third-color atoms. This superfluid state is similar to the color superfluid of the attractively interacting three-component systems. We show from dynamical properties that the Cooper-pairing mechanism is basically similar to that of the conventional phonon-mediated superconductivity. The superfluid state is stable against changes in filling close to half filling. We determine the phase diagrams in terms of temperature, filling, and the anisotropy of the repulsive interactions. We also discuss the relation between the theoretical results and experiments.

#### PF03

## Magnetism in complex oxides: A challenge for advanced ab-initio methods

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The accurate description of magnetic properties in complex oxides by first-principles is tantamount with the possibility of achieving reliable theoretical design of oxide-based spintronic devices. Alas, standard first-principles approaches fails to properly describe strong correlated systems, and more sophisticated approximations are required. Recently the novel variational pseudo-self-interaction corrected density functional theory (VPSIC) has been introduced, showing great potential as theoretical tool for the study of complex oxides and strong-correlated systems in general, and promising to open the field to reliable designs of magnetic oxide-based devices. Here we apply the new approach to investigate two prototypical oxide families, i.e. transition metal monoxides (MnO, NiO) and perovskite titanates (LaTiO<sub>3</sub>, YTiO<sub>3</sub>). We give evidence that the new theory is capable to furnish an unprecedentedly accurate description of both ground-state and finite-temperature magnetic properties, both at equilibrium and under pressure. As highlights of the results obtained in the work, we mention the correct determination of critical temperature and the accurate description of orbital ordering in transition metal oxides.

#### PF04

## Resonating Hartree-Fock Studies for spin fluctuations in the Hubbard model on triangular lattice

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Electron correlation effects in geometrically frustrated systems have attracted much attention in condensed-matter physics. In fact, these systems show a variety of interesting and non-trivial ground states. On the other hand, theoretical description on the electron correlations in these systems is difficult, and origins of such non-trivial states are not yet clarified. In this research, the electron correlations are visualized from the viewpoint of quantum fluctuations. Such visualization is important for the fundamental understanding and applications of the frustrated system. A resonating Hartree-Fock method constructs a many-body wave-function by superposition of non-orthogonal Slater determinants, where superposition coefficients and orbitals in all the Slater determinants are simultaneously optimized. We can visualize quantum fluctuations from the structures of Slater determinants generating the wave-function. We have previously demonstrated that quantum fluctuations in the doped Hubbard model on a square lattice are described by vibrations and translations of polarons. [Phys. Rev. Lett. 103(2009)116401.] Now, the half-filled Hubbard model on a uniform triangular lattice are considered. The ground state is metallic and non-magnetic when U is weak. We show that non-magnetic ground state is realized by the large quantum fluctuations due to hybridization of two different co-linear spin-density waves and their modulations. [Phys. Rev. B84(2011)245101.]

## POSTER PRESENTATION

#### **PF05**

#### Theory of momentum-dependent variational ansatz to strongly correlated electron system

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We have recently developed Momentum dependent Local-Ansatz wavefunction (MLA) approach to describe the correlated electron system in the weak and intermediate Coulomb interaction (U) regime. The theory did not work best in the strong U regime because we started from the Hartree-Fock (HF) wavefunction. To overcome the difficulty, we propose here a new hybrid wavefunction in which the starting wavefunction can vary via a new variational parameter w from the HF wavefunction (w=0) suitable for the weak interaction regime to the alloy-analogy (AA) wavefunction (w=1) suitable in the strong interaction regime. The extended method overcomes the standard variational method such as the Gutzwiller wavefunction approach (GA) irrespective of U. We have demonstrate that the energy for the MLA-AA (w=1) is lower than those of the MLA-HF (w=0) in the strong U region for the half-filled band Hubbard model for the hypercubic lattice in infinite dimensions. Moreover, the energy of the MLA (w=0 and 1) is always lower than the GA. We have also verified that the double occupation number, quasi-particle weight and momentum distribution function are much improved by the new MLA.

#### **PF06**

#### Metal-insulator transition in orthorhombic Perovksite PbRuO<sub>3</sub> Young-joon Song1 and Kwan-woo Lee2\*

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Transition metal oxide perovskites show abundant physical phenomena such as superconductivity, CMR, ferroelectricity, and unusual magnetic behavior. Recently, the orthorhombic perovskite PbRuO3, which was first synthesized 40 years ago, has been revisited by two experimental groups. Commonly, the measurements of temperaturedependent magnetic susceptibility show a kink at 90 K, where a structural transition occurs, but no magnetic ordering has been observed even at very low T. Remarkably, the resistivity measurements of Kimber et al. indicate a metal-insulator transition, but ones of Cheng et al. shows a metal-metal transition though a kink appears at 90 K. To understand this discrepancy, we have carried out first principles calculations using various approaches: correlated band theory (LDA+U), LDA+U+SOC (spin-orbital coupling) modified BJ and fixed spin moment calculations. In this presentation, we will address the electronic and magnetic structures on  $\ensuremath{\text{PbRuO}_3}$  in more detail and a scenario on the discrepancy.

#### **PF07**

Ferromagnetic semiconductor-metal transition in heterostructures of europium monoxide Tobias Stollenwerk\* and Johann Kroha

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# WITHDRAWN

#### **PF08**

#### The temperature dependence of the staggered magnetisation in itinerant weak antiferromagnets

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Effect of spin fluctuations on the staggered magnetisation in itinerant weak antiferromagnets is investigated below the Neel temperature  $T_N$ . It is known that the discontinuous change of the magnetisation occurs at the critical temperature in the rotationally invariant treatment in the self-consistent renormalization spin fluctuation theory. The difficulty is resolved in itinerant weak ferromagnets by Takahashi []. Phys. Condens. Matter 13 (2001) 6323. ]. In this paper, we extend his theory to the case of weak antiferromagnets. We have succeeded in deriving the temperature dependence of the staggered magnetisation below T<sub>N</sub>, by taking account the effects of spin-waves excitations into the transverse component of the dynamical magnetic susceptibility within the small area of wave-vector space around the staggered wave-vector. We found T3/2-linear behaviour of the staggered magnetisation at low temperatures compared with T<sub>N</sub>.

#### **PF09**

#### Theory of excitonic insulator in the two orbital Hubbard model: Variational cluster approach

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Motivated by the recent experimental and theoretical studies of excitonic insulators, we investigate the spontaneous symmetry breaking of the excitonic insulator state in the two-orbital Hubbard model defined on the two-dimensional square lattice. Using the variational cluster approximation (VCA) and Hartree-Fock approximation (HFA), we evaluate the number of particles on each orbital, staggered magnetization in the Mott insulator phase, the order parameter indicating the coherence between electrons and holes in the excitonic insulator phase, and the ground-state phase diagram as functions of intra-orbital and inter-orbital Coulomb interactions. We also calculate the singleparticle excitation spectra of the band insulator. Mott insulator, and excitonic insulator phases. We thereby argue that the normal metallic phase is unstable to the formation of the excitonic insulator phase, which should appear in a wide parameter region between the band insulator and Mott insulator phases.

#### **PF10**

#### Layered chalcogenide Ta2NiSe5 as a candidate for excitonic insulators theoretical aspects

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The electronic structure of layered chalcogenide Ta2NiSe5, a candidate for excitonic Insulators, is calculated using the generalized gradient approximation in the density functional theory, where the Hubbard-type repulsive interaction is taken into account. We find for the low-temperature monoclinic phase that the conduction band has a cosinelike quasi-one-dimensional (1D) band dispersion coming mainly from the 5dxy orbitals of Ta ions arranged along the a-direction of the lattice, whereas the top of the valence band has a quasi-1D dispersion coming mainly from the Ni 3dxz+3dyz and Se 4px+4py orbitals arranged alternately along the a-direction. The electronic state near the Fermi level is thus described by the quasi-1D structural unit consisting of these orbitals. The electronic structure of Ta<sub>2</sub>NiS<sub>5</sub>, which has the orthorhombic crystal structure is also calculated to clarify the role of the lattice distortion. Three scenarios are suggested to explain the characteristic temperature dependence of the band structure observed in the angle-resolved photoemission; a scenario of the band deformation due to the monoclinic lattice distortion a scenario of the formation of the excitonic insulator state and a scenario made from the combination of the former two where the monoclinic lattice distortion and formation of the excitonic insulator state occur cooperatively.

#### **PF11**

#### Electric dipolar susceptibility of the Anderson-Holstein model Takahiro Fuse\* and Takashi Hotta

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Recently cage structure materials have attracted much attention due to intriguing magnetic phenomena induced by the coupling effect between electrons and rattling, i.e. anharmonic local oscillation of guest ion in a cage. When the guest ion oscillates in the cage, there should occur electric dipole moment P, given by P=Zex, where Z denotes the valence of the guest ion, e indicates electron charge, and x is ion displacement. In order to investigate properties of electron-rattling system, we evaluate electric dipolar susceptibility xp on the basis of the Anderson-Holstein model, which consists of conduction electron term, hybridization between conduction and localized electrons, Coulomb interaction between localized electrons, and coupling between ion oscillation and local charge density. We analyze the model by using the numerical renormalization group (NRG) method. Note that yp is directly related to the phonon Green's function D. For the evaluation of D, we pay due attention to the phonon excited states kept in the renormalization process. We carefully evaluate D with the use of NRG and compare it with that obtained from charge susceptibility in the combination with the Dyson equation. Then, we discuss the effects of anharmonicity and Coulomb interaction on χp.

#### **PF12**

#### Insulator version of the double-exchange ferromagnetism

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The double-exchange ferromagnetism is known to appear when the coherent motion of electrons occurs in the conduction band that is coupled with the localized spins via the ferromagnetic Heisenberg-type exchange interaction, i.e., the Hund's rule coupling. A natural and realistic issue that arises is then what happens with the ferromagnetism when a gap opens in the conduction band and coherent motion of electrons ceases. To answer this question quantitatively, we apply the density-matrix renormalization group method to the study of the one-dimensional double-exchange model, where we introduce the Peierls-type lattice distortion in the conduction band to make the system an insulator with the band gap. We find from the analysis that the doubleexchange ferromagnetism persists strongly unless the strength of the lattice distortion exceeds a certain critical limit; we thus call this an insulator version of the doubleexchange ferromagnetism. We note that this type of ferromagnetism actually occurs in a chromium hollandite K2Cr8O16, where a half-metallic ferromagnet due to the doubleexchange mechanism undergoes a metal-insulator transition, resulting in the realization of a ferromagnetic insulator [1].

[11 Torivama et al., Phys. Rev. Lett. 107, 266402 (2011)

#### **PF13**

#### Metallic ferromagnetism in the 3D Hubbard model at finite temperature

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The dynamical mean-field approximation (DMFA), when used to study strongly correlated electron systems (SCES), allows the obtainment of thermodynamic properties in a nonperturbative regime. The Hubbard model is the simplest model capable of describing the essential physics of SCES. Thus, we investigated the existence of ferromagnetism in the Hubbard model on fcc lattices, at finite temperatures, using the DMFA with effective parameters solved by exact numerical diagonalization. We calculated the magnetization for 0.3 n 0.9 (where n is the number of electrons per site), with temperature of 0.04t/kB (where t is the hopping integral between nearest-neighbor sites) and the value of the on-site Coulombian interaction U equal to 3W (where W denotes the noninteracting energy bandwidth). Our results reveal a significant magnetization for n = 0.6, 0.7 and 0.8. The densities of states indicate that these systems are metallic. In the particular case for n = 0.6, magnetization, internal energy and specific-heat curves versus temperature were calculated to that value of U. This work provides an affirmative answer to the question whether the Hubbard model with only nearest-neighbor hopping, on a three-dimensional lattice, and finite both temperature and on-site Coulombian interaction, exhibits ferromagnetism when solved using the DMFA.

#### **PF14**

#### BCS-BEC crossover in the extended Falicov-Kimball model: Variational cluster approach

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We study the spontaneous symmetry breaking of the excitonic insulator state induced by the Coulomb interaction U in the two-dimensional extended Falicov-Kimball model. Using the variational cluster approximation (VCA) and Hartree-Fock approximation (HFA), we evaluate the order parameter, single-particle excitation gap, momentum distribution functions, coherence length of excitons, and single-particle and anomalous excitation spectra as functions of U at zero temperature. We find that in the weak-tointermediate coupling regime, the Fermi surface plays an essential role and calculated results can be understood in close correspondence with the BCS theory, whereas in the strong-coupling regime, the Fermi surface plays no role and results are consistent with the picture of a Bose-Einstein condensate (BEC). Moreover, we find that HFA works well in both the weak and strong-coupling regimes, and that the difference between the results of VCA and HFA mostly appears in the intermediate-coupling regime. The reason for this is discussed from a viewpoint of the self-energy. Details are reported in Ref. [1].

[1] K. Seki, R. Eder, and Y. Ohta, Phys. Rev. B 84, 245106 (2011).

#### **PF15**

#### Theory of the metal-insulator transition and charge/orbital states in V(012

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The Wadsley-phase vanadium oxide V6O13 has the crystal structure composed of the two-dimensional single and double trellis lattices, where the V ions are in the mixed valent state with an average valence of V4.33+ (3d0.66); i.e., there are formally V5+ (3d0) and V4+ (3d1) in a 1 : 2 ratio. This material undergoes a metal-insulator transition (MIT) at 150K, which is accompanied by the structural distortion, and shows an antiferromagnetic order below 55K. To clarify the electronic structure of this material, we make the density-functional-theory-based electronic structure calculations and discuss the mechanism of the MIT, as well as its charge and orbital states above and below the MIT. We find that the band structure obtained is in reasonable agreement with the observed angle-resolved photoemission spectra but that the characteristic distribution of valence electrons is not necessarily compatible with the results of a recent NMR experiment. We also find that, if we assume the observed low-temperature crystal structure in the calculations, we obtain the insulating solution with an antiferromagnetic order, which is in agreement with experiment.

#### **PF16**

#### Statistical dynamical mean field study of correlated fermions with disorder

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Disordered systems with strong correlation effects have been studied in solid state physics, and even now have attracted much interest. Various experiments have been performed extensively, and in particular a disordered system of strongly interacting cold atoms has been experimentally realized recently[1]. Statistical dynamical meanfield theory[2,3] to treat disorders and strong correlation effects on equal footing and non-perturbatively makes it possible to capture essential features of the Anderson transition beyond the coherent-potential approximation. In this work, we analyze the tight-binding fermion model on the Bethe lattice in which the box disorder manifests itself by random on-site energies. We systematically investigate characteristic behaviors of the density of states and the conductivity to estimate competitions between disorders and correlation effects in this system.

[1] L. Fallani et al., Phys. Rev. Lett. 98, 130404 (2007) [2] V. Dobrosavljevic et al., Phys. Rev. Lett. 78, 3943 (1997) [3] D. Semmler et al., Phys. Rev. B 84, 115113 (2011)

# **PF17**

# Analysis of many-body effects on anisotropic magnetic properties of $\ensuremath{YbB_{12}}\xspace$

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 $YbB_{12}$  is one of the typical Kondo insulator and many experiments are already reported. Although it has cubic structure, critical magnetic fields of insulator-to-metal transition are anisotropic[1]. Moreover anisotropic magnetization is observed above the critical field[2]. One of the authors studied the magnetic properties by means of tight binding model, composed of Yb 5de and 4f  $\Gamma_8$  orbitals[3]. The results were qualitatively in agreement with the experiments but a quantitative disagreement was left, because many-body effects were not taken into account. Therefore, we introduce coulomb repulsion between f-electrons in the tight binding model. We apply the dynamical mean field theory, in which we solve the effective impurity problem by using the continuous-time quantum Monte Carlo method. We compare obtained result with the experimental one's quantitatively, and discuss the many-body effects in the magnetic properties of  $YbB_{12}$ .

[1] F. Iga et al. : Jpn. J.Appl. Phys. Series 11(1999)88 [2] S. Hiura, et al.: Physica B 281&282(2000)271 [3] T. Izumi, Y. Imai, T. Saso : J. Phys. Soc. Jpn. 76(2007)084715

# **PF18**

# The ground state energies of spinless free fermions and hard-core bosons in 2D square lattices

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We compare the ground state energies of hard-core bosons and spinless free fermions on the same lattice. Ground state energy of bosons is usually expected to be lower than that of fermions, because bosons can condense into the lowest energy state. However, in the presence of hard-core interaction among bosons, the condensation is not perfect and comparison is nontrivial. We find that, in the absence of frustration, the ground state energy of hard-core bosons is indeed lower than that of free fermions. In general, however, the hard-core bosons could have higher ground state energy than the fermions on the same lattice if there is a frustration. In fact, there is an exactly solvable one-dimensional model, in which the ground state energy of bosons is proved to be higher. However, the difference vanishes in the thermodynamic limit. We also studied the two-dimensional square lattice with uniform magnetic flux perpendicular to the lattice, by numerical diagonalization. As the magnetic flux introduces frustration, we actually find the ground state energy of hard-core bosons is higher for a range of flux and particle density. In the two dimensional case, this difference seems to persist in the thermodynamic limit, in contrast to the one-dimensional case.

[1]Phys. Rev. B 14, 2239 (1976) [2]Phys. Rev. Lett. 83, 2246(1999) [3]Phys. Rev. Lett. 96, 036406 (2006)

### **PF19**

**Correlation effect in ferromagnetic 3d transition metals** Muneyuki Nishishita<sup>1</sup>, Sudhakar Pandey<sup>2</sup> and Dai Hirashima<sup>1</sup>\* <sup>1</sup> Nagoya University, Japan <sup>2</sup> APTPC, Korea

Selfconsistent perturbation theory is applied to ferromagnetic 3d transition metals, Fe and Ni. Realistic tight-binding model is used to describe the electronic states of the system. It is found that a reasonable value of the Curie temperature for Fe is obtained for realistic strength of the Coulomb interaction. Dynamical transverse susceptibility is also calculated. Here, vertex corrections are considered so that the spin rotational symmetry is preserved. Then, the spin-wave energy correctly vanishes in the longwavelength limit. The effect of the electron correlation on the spin stiffness is clarified. We also compare the spin fluctuation spectra below and above the Curie temperature and discuss the possible short-range order above the Curie temperature.

# PG01

Ferrimagnetic compensation in a Fe<sub>64</sub>Er<sub>19</sub>B<sub>17</sub> glass, - the head of a dandelion, or the spokes of a wheel?

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Magnetization and neutron scattering measurements with polarization analysis were made on a  $Fe_{od}Er_{i9}B_{17}$  glass to investigate its ferrimagnetic compensation at T(comp) =  $112 \pm 2K$ . The measurements prove that the glass has a non-collinear magnetic structure. The two non-spin flip cross-sections interchange between 100K and 125K, proving the magnetic structure flips at T(comp). They are identical at 112K where the forward limit of the spin flip cross-sections was greatest. The mean collinear components of the moments therefore go to zero at T(comp). If the moment vectors were mapped to originate at a point, the data correspond either to the moments being randomly oriented like the head of a dandelion, or with a collapse of the moments into a plane normal to the field, like the spokes of a wheel. The data will be presented and the possible driving mechanisms for the ferrimagnetic compensation will be discussed.

#### PG02

# Evolution of reverse magnetized seed in monodomain uniaxial garnet film elements

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Evolution of cylindrical domain seeded in monodomain film and film circular element was investigated numerically and experimentally. It is found that in continuous film in zero external magnetic field cylindrical domain may expand unlimitedly or collapse independently on initial diameter of the seed. Single seed expansion followed by film magnetization reversal is observed if l/h < 1.38 (where l is garnet characteristic length, h is film thickness). Constant positive external magnetic field antiparallel to magnetization inside seed supports stable domain diameter which is inverse proportional to field intensity H. Cylindrical domain inside disk is found stable in zero and negative external magnetic field due radial gradient of inhomogeneous stray field of the disk boundary. If l/h > 1.38 reverse magnetized seed collapses in zero external field. Stable domain doesn't exist in constant negative external field parallel to seed magnetization this case. Unlimited domain expansion takes place if modulus of negative H exceeds critical value which is inverse proportional to initial seed diameter. Domain is stable if domain wall embraces circular hole in garnet film. This case garnet boundary stray field prevents from domain contraction in zero and positive external field. Two-domain ring-shape elements can be utilized as fast magnetooptical modulators

### **PG03**

#### Re-entrant structure and critical behaviour of Fe-Cr and Fe-V sigmaphase allovs

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The  $\sigma$ -phase (tetragonal unit cell, space group D144h-P42/mnm) belongs to a family of a complex Frank-Kasper phases. It can be formed in transition-metal alloys. Only two alloys, among about 50 examples of  $\sigma$  known to occur in binary systems, viz. Fe-Cr and Fe-V have well-evidenced magnetic properties. The magnetic ordering of  $\sigma$  has been regarded as ferromagnetic (FM) with the Curie temperatures, Tc, below ~40 K for  $\sigma$ -FeCr and below ~315 K for  $\sigma$ -FeV. Here we report a clear evidence that the ground state of the  $\sigma$ -phase Fe<sub>0.55</sub>Cr<sub>0.47</sub> and Fe<sub>0.52</sub>V<sub>0.48</sub> alloys is a re-entrant spin-glass (RSG). The evidence was found from a study of magnetization versus temperature, T, and magnetic field, H. Based on the field-cooled and zero-field-cooled magnetization curves recorded in different fields, H, phase diagrams in the H-T plane could have been drawn. They display the location of the paramagnetic (P), FM and RSG states, and are in a qualitative accord with predictions of the mean-field theory. Critical behaviour near Tc was also studied with both conventional and extended protocols, and the critical exponents  $\beta$ ,  $\gamma$ ,  $\delta$  were determined.

# PG04

# Pressure effects in the ferromagnetic shape memory alloys $Ni_2Mn_{1-x}Cu_xGa$

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The Heusler alloy Ni<sub>2</sub>MnGa attracts attention as an actuator material because the magnetic distortion of about 10% is observed just below the martensitic transition temperature (TM). The pressure dependence coefficient of the Curie temperature of Ni<sub>2</sub>MnGa is +1.0×10-8 K/Pa.[1] In the Ni<sub>2</sub>Mn<sub>1-x</sub>Cu<sub>x</sub>Ga alloy system, TC and TM were reported by Kataoka et al.[2] that the samples with the range of x < 0.23, 0.23 < x < 0.30 and x > 0.30 are characterized TC > TM, TC = TM and TC < TM, respectively. We carried out the initial permeability measurements under pressure up to 1 GPa for Heusler Ni<sub>2</sub>Mn<sub>1-x</sub>Cu<sub>x</sub>Ga system. TM and TC were determined from the change of the initial permeability under each pressure. It was found that TM was considered to coincide with TC in 0.23 < x < 0.30. As a result, the temperature-pressure magnetic phase diagram was obtained. We'd like to discuss the mechanism of the martensitic transformation in Heusler alloy Ni<sub>2</sub>Mn<sub>1</sub>. Cu<sub>x</sub>Ga system.

[1] T.Kanomata et al.: J. Magn. Magn. Mater 65 (1987) 76. [2] M. Kataoka et al.: Phys. Rev B 82 (2010) 214423.

#### PG05

# Electronic state of Cr and collapse-like decrease of Fe magnetic moment in amorphous (Fe-Cr)B alloys

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Magnetic properties in amorphous alloys have been extensively investigated in rare earth (RE)-transition metal (TM) systems and transition metal (TM)-mettaloids systems. Many interesting and important problems have been investigated and solved, however, some of the important problems seem to be hard nuts to crack and remain unsolved. One of the unsolved and interesting problems is an anomaly in Slater-Pauling curve discovered by Mizoguchi et al. and the Fe magnetic moment decreases abruptly, for example in (Fe-Cr)<sub>30</sub>B<sub>20</sub> amorphous alloys. For instance, the substitution of Fe by only 10 at% of Cr causes about 50 %decrease of Fe magnetic moment. Mizoguchi et al. suggested that the Cr retained magnetic moment and coupled anti-parallel to that of Fe, however there have been no decisive experimental results. In order to detect and examine the Cr magnetic moment, one of the powerful tools for this aim, is the magnetic Compton scattering profile (MCP) method which can obtain the informations of 3d electrons and the spin-polarized outer 4s, 4p electrons, separately. The MCP measurement was carried out for (Fe-Cr)<sub>30</sub>B<sub>20</sub> (Cr=0.1, 0.2, 0.3) and the result suggest the existence of Cr magnetic moment coupling antiparallel to that of Fe. Macroscopic magnetic properties were also measured and analyzed.

#### PG06

#### Micromagnetic simulation of CNT-MFM probes under magnetic field Takashi Manago<sup>1</sup>\* Hironori Asada<sup>2</sup> and Hiromi Kuramochi<sup>3</sup>

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Ferromagnetic-film-coated carbon nanotube in a magnetic force microscope (CNT-MFM) demonstrated the high spacial resolution of about 10 nm and low perturbation. The magnetic structures of CNT-MFM probes were investigated using threedimensional micromagnetic simulation and it showed the advantages of the CNT-MFM probes [1]. In this paper, the stability of the magnetic structure in external magnetic field was investigated to confirm magnetic robustness using micromagnetic simulation. At a remanent state, the magnetic moments of the CNT-MFM probe align almost along the longitudinal direction of the CNT lod due to the shape anisotropy. When the opposite magnetic field applied, the magnetic moments of the probes reversed all at once at about 200 mT. In the pyramidal probe, the vortex core was appeared around the tip at a remanent state. When the opposite magnetic field applied, the magnetization of the side wall gradually tilted, and at last the vortex core was reversed. Thus magnetization reversal process is quite different because of large difference in magnetic tructures between the CNT-MFM and pyramidal probes. We found that the CNT-MFM probes have better magnetic robustness than the conventional pyramidal probes.

[1] T. Manago et al., Nanotechnology, 23, 035501 (2012).

# **PG07**

### Pressure effect of metamagnetic shape memory alloy Pd<sub>2</sub>Mn<sub>1+x</sub>Sn<sub>1-x</sub>

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Recently, it has been reported that NiMnX (X = In, Sn and Sb) Heusler alloys show metamagnetic shape memory effect. In these alloys, the martensitic transformation temperature decreases with increasing magnetic field, and magnetic field-induced reverse martensitic transformation associated with metamagnetic transition occurs below the martensitic transformation temperature. Since the magnetic field-induced transformation leads an almost perfect shape memory effect, these alloys attract attention as new type of ferromagnetic shape memory alloys. Very recently, similarly to NiMnX Heusler alloys, it has been found that PdMnSn Heusler alloys show metamagnetic shape memory effects and a drastic change with martensitic transformation in magnetic and transport properties. In this study, in order to investigate pressure effect on martensitic transformation in PdMnSn Heusler alloys, we performed electrical resistivity measurements under high pressure. As a result, it was found that the pressure-temperature phase diagram is similar to the temperature-Mn content phase diagram. These results suggest that the atomic distance plays an important role for martensitic transformation in PdMnSn Heusler alloys.

### **PG08**

# Engineering of Co atomic 1-D chains on Ag(111) with tailored magnetic ground state

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Certain magnetic adatoms lying on a metal surface give rise to a Kondo resonance. We used the model system of Co atoms on Ag(111) to investigate their magnetic interactions through the Kondo resonance [1,2]. The Co atoms architecture and the substrate LDOS landscape are artificially arranged by means of lateral manipulation with a scanning tunneling microscope (STM) tip at 4.8 K. The resonance is subsequently characterized by means of STM spectroscopy. Individual Co atoms in natural equilibrium positions at minima of the surface LDOS exhibit Kondo temperature TK=90 K in agreement with earlier studies [3]. TK can be gradually varied from 60 to 145 K by positioning the same atom in different Ag LDOS values of the interference patterns. The coordination with neighboring Co atoms increases TK when their distance decreases below four Ag(111) lattice parameters (2.9 A). In Co chains with N atoms (N=2,3,...,19), both the amplitude and the spatial extent of the many-body wave function of the Kondo resonance are reduced with increasing coordination. In the frame of the two impurity Kondo theory, our data unveil the nature of direct Co-Co and substrate mediated magnetic interactions as a function of the local electronic and structural environment.

[1] A. Otte et al., Nature Phys. 4, 847 (2008) [2] J. Bork et al., Nature Phys. 7, 901 (2011) [3] M.A. Schneider et al., Phys. Rev. B 65, 121406R (2002)

### **PG09**

# Pressure-induced suppression of magnetic ordering in a chiral magnet $Cr_{10}NbS_2$

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 $Cr_{1/3}NbS_2$  is a chiral magnet without structural inversion center. The magnetic property is attributed to the localized moment of  $Cr_{t}$  and metallic conductivity is attributed to conduction electrons of Nb. We expect that  $Cr_{1/3}NbS_2$  would become a noncentrosymmetric superconductor after disappearance of magnetic ordering at high pressure. Thus we carried out the AC susceptibility measurements at pressures up to 10 GPa and powder X-ray diffraction experiments at room temperature up to 7 GPa for the powder sample of  $Cr_{1/3}NbS_2$ . In the pressure region of below 3 GPa, the anomaly of magnetic susceptibility due to the magnetic ordering shifted toward low temperature side and the magnitude was gradually suppressed. At present, the Meissner signal has not yet been observed in the temperature region down to 2 K and the pressure region up to 10 GPa. The lattice parameters exhibited the change of contraction at around 3 GPa, where it became difficult to evaluate the magnetic ordering temperature.

# **PG10**

#### Multiple ESR spectra in a chiral molecule-based magnet [Cr(CN)<sub>6</sub>] [Mn(R )-pnH(H<sub>2</sub>O)](H<sub>2</sub>O)

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A molecule-based magnet [Cr(CN)<sub>0</sub>][Mn(R)-pnH(H<sub>2</sub>O)](H<sub>2</sub>O) (abbreviated as R-GN) is a chiral magnet without inversion center and mirror reflection. The magnetic network is constructed with the help of a chiral ligand diaminopropane. In R-GN, multiple spectra of ESR were observed below magnetic order temperature by Morgunov et al. [1]. They concluded that the phenomenon was caused by the formation of a chiral soliton lattice, which originated from incommensurate magnetic structure. However, we have a suspicion against their interpretation, because R-GN indeed exhibits no incommensurate magnetic structure. Now, we conducted the experiment of X-band ESR measurement in the condition similar to the experiment by Morgunov et al., and performed the detailed spectrum analyses: In order to reproduce the main ESR spectra, at least two Lorentz spectra are required, and the multiplicity was explained by assuming the existence of two characteristic modes when each periodicity is defined as f or f' < f, there is a relation of f = 2f'. We consider that the magnetic sites owing to the diaminopropane with electric dinole moment.

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# PG11

# Valence and spin structures of TFe2O4 spinel oxides (T=Mn, Co, Ni, Cu) investigated by using synchrotron radiation

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AB<sub>2</sub>O<sub>4</sub>-type spinel oxides have attracted much attention because of the interesting physical phenomena, such as Jahn-Teller (JT) effects, multiferroic phenomena, and phase separations [1]. In normal spinels of AB<sub>2</sub>O<sub>4</sub>, A and B ions occupy the tetrahedral (Td) and octahedral (Oh) sites, respectively. When both A and B ions are magnetic ions, a ferrimagnetic (FiM) ordering is often observed because the A-B interaction is mainly antiferromagnetic. In this work, we have investigated the electronic structures of FiM spinels of TFe<sub>2</sub>O<sub>4</sub> (T=Mn, Co, Ni, Cu) by employing soft x-ray absorption spectroscopy (XAS) and soft x-ray magnetic circular dichroism (XMCD). It is found that the valence states of Fe ions are nearly trivalent (-Fe<sup>3+</sup>) and those of T ions are nearly divalent (-T<sup>2+</sup>). The 2p XMCD spectra of MnFe<sub>2</sub>O<sub>4</sub> show that the magnetic moments of Mn<sup>2+</sup> and Fe<sup>3+</sup> ions are antiparallel to each other. On the other hand, the magnetic moments of Co<sup>2+</sup>, Ni<sup>2+</sup>, and Cu<sup>2+</sup> ions are parallel to those of B(Oh)site Fe<sup>3+</sup> ions, but are antiparallel to those of A(Td)-site Fe<sup>3+</sup> ions, implying that Fe<sup>3+</sup> ions occupy both A and B sites while Co, Ni, Cu uby B sites. \*Corresponding Author: kanej@(actholic.ac.kr

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#### PG12

# Probing the distance dependence of the magnetic exchange interaction with atomic resolution

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Quantifying strength and distance dependence of magnetic interactions between atoms is crucial to gain a deeper understanding of magnetic phenomena on the nanoscale. Magnetic exchange force microscopy (MExFM) [1] is a new and very versatile tool to study magnetism with atomic resolution in real space. It utilizes an atomically sharp magnetic tip at the free end of a cantilever to detect the short-ranged electron-mediated magnetic exchange interaction between tip apex and surface atoms. This force microscopy based technique can be employed to map spin structures of insulating [1] as well as conducting [2] magnetic samples with atomic resolution. Here we demonstrate that its spectroscopic mode, i.e., magnetic exchange force spectroscopy (MExFS), is able to measure the distance dependence of the magnetic exchange interaction in a very direct and elegant fashion [3]. Results obtained on the antiferromagnetic Fe monolayer on W(001) are compared with theoretical calculations based on density functional theory. They were performed for realistic multiatom tips and agree very well with data acquired using non-dissipative stable tips.

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# PG13

#### Electron correlation in a mixed valence perovskite system of Sr<sub>1</sub>. ,Ca,Ru<sub>0</sub>,Mn<sub>0</sub>,O<sub>1</sub>

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Mixed valence systems are interesting for their fundamental and practical properties. Perovskite-type oxides  $Sr_{i,x}Ca_{k}Ru_{i,y}Mn_{i}O_{j}$  is a typical mixed valence system.<sup>1</sup> The host titnerant system  $Sr_{i,x}Ca_{k}RuO_{j}$  shows various interesting properties such as ferromagnetism and paramagnetism.<sup>2</sup>,<sup>3</sup> Mn-substitution induces a mixed valence state of Ru<sup>+</sup>, Ru<sup>+</sup>, Mn<sup>+</sup> and Mn<sup>+</sup> ions and localized ferromagnetism.<sup>1</sup>,<sup>4</sup> However, understanding of the electron correlation in the mixed valence systems is not yet enough. In this work, we studied the relationship between electron correlation and mixed valence state through the experiments for  $Sr_{i,x}Ca_{k}Ru_{0,x}Mn_{0,x}O_{1}$  using the traditional experimental methods and the synchrotron-based magnetic Compton scattering.  $Sr_{i,x}Ca_{k}Ru_{0,x}Mn_{0,x}O_{1}$  is a good system since the Ca-substitution will modify crystal symmetry and change the electron correlation. The following results have been obtained; 1. Cell volume decrease for Ca substitution 2. insulator-metal transition occurs for Ca substitution. 3. Ferromagnetism appears at x - 0.3. 4. Antiferromagnetic coupling between Ru and Mn. These facts can be explained as follows. The decrease in the cell volume enhances the hybridization between Ru and O orbitals inducing tinerant electrons. The itinerant electrons enhance the formation of a ferromagnetic order of Ru ions. Mn ions couples antiferromagnetically with Ru ions by a superexchange interaction via O 2p orbitals.

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#### PG14

# An x-ray scattering study of magnetic order in a pyrochlore iridate Eu<sub>2</sub>Ir<sub>2</sub>O<sub>7</sub>

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The pyrochlore iridate Eu<sub>2</sub>Ir<sub>2</sub>O<sub>7</sub> undergoes a metal-insulator transition at  $T_{\rm +}$ [MI] = 120 K [1]. Since the magnetic Ir<sup>4+</sup> ions occupy the pyrochlore B-site, strong geometrical spin frustration is predicted. Nevertheless, a muon spin rotation study [2] indicated long-range orderings spins of Ir<sup>4+</sup>. To determine the spin configuration, we have performed the resonant X-ray diffraction measurement at the LIII edge of Ir. A clear magnetic reflection was found at (10,0,0) below T\_{MI}. The result proposes the all-in-all-out-type magnetic structure.

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#### PG15

# Magnetic structure analyses by small-angle electron diffraction

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Electrons are deflected at small angles by Lorentz force in magnetic materials. In magnetic elements in functional materials and spintronic devices, observation of Lorentz deflection of electrons at the small angle turns to be of significance because Lorentz deflection of electrons becomes smaller in modern miniaturized magnetic devices composed of thinner magnetic films. In the present work, electron optical system is constructed in order to obtain small angle diffraction and Lorentz deflection of electrons at the order of down to 10-6 radian in the reciprocal space[1]. Long-distance camera length up to 3000 m is achieved in a conventional transmission electron microscope with LaB6 thermal emission type. The diffraction pattern at  $5 \times 10-6$  radian is presented in a carbon replica grating with 500 nm lattice spacing while the magnetic deflection pattern at  $2 \times 10-5$  radian is exhibited in Permalloy elements. A simultaneous recording of electron diffraction and Lorentz deflection is also demonstrated in 180 degree striped magnetic domains of La<sub>0125</sub>Sr<sub>0.175</sub>MnO<sub>3</sub>[2]. We believe that small-angle Bragg diffraction and Lorentz deflection analyses provide a quantitative way to analyze structural and magnetic properties in the reciprocal space. This study is partly supported by SCOPE project from the MIC in Japan.

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# PG16

### Raman scattering of metal-insulator transition in Cd<sub>2</sub>Os<sub>2</sub>O<sub>7</sub>

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 $Cd_2Os_2O_7$  is a pyrochlore oxide showing metallic behavior at room temperature, and it shows metal-insulator transition at 225 K [1]. It has been suggested that this metal-insulator transition is caused by magnetic order [1,2], because any superlattice reflection has not been observed by x-ray and neutron [3] diffraction measurements. However, to conclude no structural distortion precisely below the transition temperature, it is necessary to perform symmetry sensitive measurements rather than diffractions. Raman scattering method is a sensitive tool to observe small lattice distortion with symmetry breaking. We have measured Raman spectra of  $Cd_2Os_2O_7$ from 4 K to 300 K to investigate possibility of structural distortion in the insulator phase. Below the transition temperature, we have not observed any new peak indicating symmetry breaking from the pyrochlore structure with the space group Fd3m.

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# PG17

#### Effects of impurities in the chromic compound CuMoO<sub>4</sub>

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Recently interest in "Chromism" phenomena for compounds may be attributed to the fact that a reversible color change of organic and inorganic materials can occur as the compound experiences various changes in its environmental conditions. Copper molybdate, CuMoO4, has demonstrated attractive properties of both thermochromism and piezochromism that are induced by a structural phase transition. We report a tunable structural phase transition and magnetic phenomena using substitutional effects in  $Cu_{1-x}Zn_xMoO_4(0 \le x \le 0.1)$ ,  $CuMo_{1-y}W_yO4(0 \le y \le 0.1)$ , and  $CuMo_{1-z}Cr_zO4(0 \le z \le 0.1)$  over wide ranges of temperatures and magnetic fields. The substitution of magnetic Cu<sup>2</sup> ions for nonmagnetic Zn2+ ions causes the hysteresis in the structural phase transition to shift to remarkably lower temperatures as x increases by means of magnetic susceptibility measurement. In contrast, replacing Mo6+ ions with W6+ ions results in a shift of the hysteresis and the structural phase transition to higher temperatures as y increases. However, in the case of replacing the Mo<sup>6+</sup> ions with Cr<sup>6+</sup> ions, the structural phase transition around 200K is completely disappeared and a ferromagnetic component is induced below 150K as z increases. We are going to discuss the reason of the shifting structural phase transitions and also show the magnetic properties under extreme conditions

### PG18

#### Effective exchange interactions in 5d transition metal oxides

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5d transition metal oxides. Sr-IrO, and Ba-IrO, have attracted much attention because of their unique properties caused by a strong spin-orbit coupling for 5d transition element. Both of these compounds behave as Jeff=1/2 Mott insulator at low temperature [1,2,3]. Recently, these compounds are expected to be a superconductor as an analogous system to high-Tc cuprates. According to recent experimental results on magnetic susceptibility and uSR studies, the effective exchange interactions in Sr-IrO, and Ba<sub>2</sub>IrO, are expected to be very large, and comparable to the ones of high-Tc cuprates [3]. Motivated by these experiments, we have studied theoretically the electronic and magnetic properties of Sr<sub>2</sub>IrO<sub>4</sub> and Ba<sub>2</sub>IrO<sub>4</sub> using a three-band Hubbard model with the spin-orbit coupling [4]. We first obtain the phase diagram of 8-site cluster in a wide parameter region using exact diagonalization technique. We find that there are several phases including the Jeff=1/2 Mott insulating phase which is realized in Sr. IrO, and Ba-IrO. We then calculate the dynamical spin structure factor composed by the Jeff=1/2 state in the Jeff=1/2 Mott insulating phase. We then estimate the values of effective exchange interactions by comparing the results with the ones of an effective spin model.

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# PG19

# Magnetic and electrical transport behavior of Ir substituted NiTi shape memory alloy

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NiTi shape memory alloys (SMAs) show attractive functional properties for a number of engineering and medical applications [1]. Recently. Somsen et.al [2] have synthesized the novel Ni-Ti-Ir alloys and have studied its structural properties. This motivates us to study the other properties of the alloy. In the present work, we have systematically studied the structural, kinetics, magnetic and transport properties of the Ni-Ti-Ir alloy. The alloy is formed in the B2- type phase with a small amount of Ti<sub>2</sub>NiIr as an additional phase. In the DSC curve, two distinct peaks are observed on heating and cooling. These peaks correspond to the Martensitic transformations (MT) from the B2 to R-phase and from R to B19'-phase on cooling and, a reverse transformation from B19'-R-B2 phase on heating, respectively. From the temperature dependence of resistivity curve, the Martensitic Start (MS), Martensitic Finish (Mf), Austenite Start (AS) and Austenite Finish (Af) temperatures are found to be 193 K, 153 K, 267 K and 293 K respectively. In the Magnetization versus temperature (M-T) curve, no evidence of the ferromagnetic ordering temperature is observed. However, the Magnetization versus Field (M-H) curve at 5 K and 300 K shows existence of ferromagnetic order in the alloy.

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### PG20

### X-ray diffuse scattering of pyrochlore niobium oxides $R_2 N b_2 O_7$

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It was reported that  $Y_2Nb_2O_7$  is a non-magnetic insulator, although the electron configuration of Nb<sup>4+</sup> is 4d1 with S=1/2. To reveal the origin, we have performed X-ray scattering measurement of single crystals of  $R_2Nb_2O_7$  (R=Y,Sm,Dy,and Ho). All the compounds show diffuse scattering, indicating that there exists short-range order. Analysis of the average structure suggests that Nb ions are displaced along the (111) direction. The formation of spin-singlet tetramers cannot explain the observed diffuse pattern. The short-range correlation implies the formation of spin-singret dimmers, because the configuration of dimmers in a pyrochlore lattice can be mapped on the icerule problem. In addition, the diffuse pattern changes with R. The diffuse of  $Y_2Nb_2O_7$  is the broadest, while that of  $Dy_3Nb_2O_7$  is narrowest and almost spot.

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### PG21

# Freezing of local lattice strains in the magnetic martensitic/ferroelastic material system

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Glass state is a frozen disordered state, in which the long range order is destroyed but short range order persists. Recently, a new glass phenomenon called strain glass was found in the martensitic/ferroelastic material system [1]. The strain glass state is a frozen disordered state of local lattice strains (nanosized martensitic/ferroelastic domains). Experimental measurements show that strain glass possesses obvious glassy features including dynamic freezing and breaking down of ergodicity in its mechanical properties [1, 2], being similar to the glassy features of the cluster-spin glass with froze local spin order. Strain glass can be obtained by doping point defects into normal martensitic/ferroelastic system [1]. However, whether strain glass can be achieved in the defect doped magnetic martensitic/ferroelastic system is unknown so far. In this study, we report that a strain glass transition with freezing of local lattice strains can exist in the Co doped magnetic martensitic system Ni-Mn-Ga by using dynamic mechanical analysis. Moreover, our magnetic measurements show that the strain glass in the Co doped Ni-Mn-Ga system is ferromagnetic, which is a new form of strain glass. The finding of the strain glass in magnetic martensitic/ferroelastic system may lead to novel applications.

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# PG22

#### Magnetic structure of the new chiral compound [Cr(CN)<sub>6</sub>][Mn(S)pnH(DFM)](H<sub>2</sub>O)

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Research on magnetic molecular materials has received strong attention in the last years for several reasons, among them, the possibility of these materials to be multifunctional, thus exhibiting more than one functional property at the same time. A very interesting possibility in a molecular material is magneto-chirality. Despite this interest, very few molecular candidates with likely possibilities have been synthesized, because chirality must be controlled in the molecular structure and at the same time the crystal must exhibit nuclear chirality. In addition, the arrangement of the magnetic moments must also be chiral. Nowadays, single crystals of the cyano-bridged bimetallic complexes have been growth, a 2D chiral ferrimagnet [Cr(CN)<sub>6</sub>] [Mn(R)-pnH(DMF)]•(2H<sub>2</sub>O) (1) and a racemic one [Cr(CN)<sub>6</sub>][Mn(rac)-pnH(DMF)]•(2H<sub>2</sub>O) (2), where DFM is N<sub>1</sub>N-dimethylformamide. Compound 1 consists of two-dimensional bimetallic sheets, with the same space group P212121, whereas compound 2 crystallizes under Pnma [1]. These compounds have been studied employing neutron diffraction techniques at the Institute Laue Langevin. In this communication, we present the nuclear and magnetic structures for (1) and (2). The coexistence of nuclear and magnetic chirality has been demonstrated.

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# PG23

Origin of spin scalar chiral order in frustrated Kondo lattice model - higher-order Kohn anomaly and hidden positive biquadratic interaction -

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Recently, noncoplanar spin configurations with scalar chirality have drawn considerable attention as an origin of the anomalous Hall effect. In this mechanism, itinerant electrons acquire an internal magnetic field through the Berry phase according to the solid angle spanned by three spins, which can result in the anomalous Hall effect. In order to clarify how such nontrivial magnetic states appear in spin-charge coupled systems, we investigate a ferromagnetic Kondo lattice model on a triangular lattice by variational and perturbative calculations. As a result, we find that a noncoplanar four-sublattice spin ordering emerges near 1/4-filling [1], in a wider parameter region compared to the nesting-driven 3/4-filling phase predicted in the previous study. We unveil that a kinetic-driven positive biqaudratic interaction is critically enhanced and plays a crucial role on stabilizing a chiral ordering near 1/4-filling. The origin of large positive biquadratic interaction is ascribed to the Fermi surface connection by the four-sublattice ordering wave vectors, which we call the higher-order Kohn anomaly [2]. Similar hidden interactions and resultant noncoplanar states are also found in other frustrated lattices, such as fcc and pyrochlore lattices. Our results usggest a universal mechanism underlying nontrivial spin configurations in frustrated spin-charge coupled systems.

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### PG24

#### Doping effects on the metal-insulator transition of Li<sub>2</sub>RuO<sub>3</sub>

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Li<sub>2</sub>RuO<sub>3</sub> has a layered rocksalt-type structure and has attracted much attention because of its interesting physical properties. In particular, it exhibits a metal-insulator transition at a very high temperature around 540 K. At the same temperature, the magnetic susceptibility shows a big drop implying the formation of a spin gap. It is also interesting to note that all of these physical properties occur on the unusual honeycomb lattice of Ru ions. In order to understand the origin of the metal-insulator transition and its related physical properties better, we have carried some extensive studies ranging from bulk properties to neutron scattering experiments. First, we have investigated the doping effects on the metal-insulator transition by varying Li content. Second, we have examined the field dependence of the physical properties by measuring resistivity and magnetization over a wide temperature range. In this report, we are going to summarize our experimental results and its importance with respect to the reported metal-insulator transition.

# PG25

#### Ferromagnetism in hydrothermally treated glassy carbon Hyun Jin Cho, Kyu Won Lee and Cheol Eui Lee\*

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We have investigated the magnetic properties of pristine and hydrothermal-treated glassy carbon by means of superconducting quantum interference device (SQUID) and electron spin resonance (ESR) measurements. The magnetic properties of the hydrothermally treated glassy carbon showed a sensitive dependence on the solvents and the treatment time. In particular, prominent ferromagnetism was observed in samples hydrothermally treated in ethanol for a properly chosen time, with the saturation moment and coercive force being comparable to those of iron. The ferromagnetism in the hydrothermally treated glassy carbon may have to do with the relative fractions of the sp3- and sp2-bonds.

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### PG26

**On the structure and symmetry of the spin glass state (SGS)** Jerzy Warczewski\*, Pawel Gusin and Daniel Wojcieszyk

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The explanation of uniqueness of SGS has roots in the appearance of a certain probability function in the second term of the assumed Hamiltonian [1, 2]. This term describes the random distribution of either dopants or defects in the ferromagnetic matrix under the percolation threshold. The Gaussian type randomness was derived for both the global magnetic coupling constant and the magnetization vector M, the latter effect bringing to the statistical features of the magnetic structure and the magnetic symmetry group of SGS. In both cases the use has been made of the central limit theorem of the theory of probability (the Lyapunov theorem) [3]. The authors have earlier proved that a certain spontaneous minimum magnetic field H is necessary for the stability of SGS [4]. The structure of SGS can then be described in such a way that M is situated along a generatrix of a given cone whose axis coincides with the direction of H [4]. Let us call  $\varphi$  the angle between M and H. Any precession of M around the direction of H at  $\varphi = \text{const}$  (at constant energy) makes the symmetry operation of SGS.

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### PG27

Magnetic study of Fe/MgO/Fe and Fe/MgO/Fe/Co multilayer systems Jitendra Pal Singh<sup>1</sup>, Sanjeev Gautan<sup>2</sup>, K Asokan<sup>1</sup>, D. Kabiraj<sup>1</sup>, D. Kanjilal<sup>1</sup>, M Raju<sup>3</sup>, Braj Bhusan Singh<sup>4</sup>, S. Chaudhary<sup>3</sup>, R. Kotnala<sup>4</sup> and Keun Hwa Chae<sup>2</sup>\* <sup>1</sup> Inter University Accelerator Centre, Arnun Asaf Ali Marg-110067, New Delhi, India <sup>2</sup> Advanced Analysis Center, Korea Institute of Science and Technology (KIST), Seoul 136-791, Korea <sup>3</sup> Department of Physics, Indian Institute of Technology, New Delhi - 110016, India <sup>4</sup> National Physical Laboratory, Dr. K. A. Krishanan Marg, New Delhi - 110012, India

Theoretical calculations [1-3] show high TMR up to 1000% for Fe/MgO/Fe MTJ's however TMR ratios of 200% at room temperature could be realized in these systems [4]. The discrepancy between theoretical and experimental values of TMR is probably due to the interface properties of the heterostructures, as spin dependent tunneling is sensitive to the interface [3]. In present work, we investigated the magnetic properties of Fe/MgO/Fe (MFe/O) multilayers by vibrating sample magnetize properties of Fe/MgO/Fe (MFe/O) multilayers by vibrating sample magnetize anisotropy (PMA) are observed in MFe multilayer. The PMA arises due to Fe(3d) and O (2p) hybridization at metal-oxide layer interface and generally observed in MTJs [6]. The asymmetric coercivity observed in this structure is due to the presence of exchange bias which arises probably due to oxidation of Fe at surface; resulting in antiferromagnetic Fe<sub>0</sub>. The positive vertical shift along magnetization axis may be due to the pinned interfacial uncompensated spins of multilayers. The absence of these is expected almost 3 times increase of Fe layers in MFeCo system [5].

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# PH01

### Ionic size effect on the spin gap nature of SrCu<sub>2</sub>(TeO<sub>3</sub>)<sub>2</sub>Cl<sub>2</sub>

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We report the effects of partial substitution of Br onto the Cl sites of SrCu<sub>2</sub>(TeO<sub>3</sub>)<sub>2</sub>Cl<sub>2</sub> by means of magnetic susceptibility measurements. This material has been a subject of current interest due to indications of spin gap behavior. However, the magnetic spin couplings responsible for the spin gap nature remains unclear. For each composition of SrCu<sub>2</sub>(TeO<sub>3</sub>)<sub>2</sub>(Cl<sub>11,3</sub>Br<sub>x</sub>)<sub>2</sub>, the temperature-dependent susceptibility exhibits a character of low-dimensional magnetism with a broad maximum at Tmax. The magnitude of Tmax gradually shifts to higher temperatures as increasing the Br concentration. The experimental data can be well fitted to a coupled spin dimer model, yielding an increasing trend for the spin gap size from 170 K to 190 K. It thus points out that the substitution of Br ions has an effect to reduce the distance between Cu(1) and Cu(2), leading to a stronger antiferromagnetic spin interaction. On the other hand, the variation of the interdimer coupling is marginal, indicative of weak effects on the distances of Cu(1)-Cu(1) and Cu(2)-Cu(2) in SrCu<sub>2</sub>(TeO<sub>3</sub>)<sub>2</sub>Cl<sub>2</sub>.

#### PH02

#### Magnetism of $SrM_3P_4O_{14}$ ( $M^{2+}$ = 3d ions) investigated using neutronscattering measurements

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One of intriguing phenomena in magnetism is appearance of qualitatively different magnetism in isostructural substances. In AM<sub>3</sub>P<sub>4</sub>O<sub>14</sub> (A = Ca, Sr, Ba or Pb, M<sup>2+</sup> = Ni<sup>2+</sup>: spin-1, Co<sup>2+</sup>: spin-3/2, Fe<sup>2+</sup>: spin-5/2), magnetism in AMn<sub>3</sub>P<sub>4</sub>O<sub>14</sub> is qualitatively different from magnetism in the other substances. A 1/3 quantum magnetization plateau was observed only in AMn<sub>3</sub>P<sub>4</sub>O<sub>14</sub>. An antiferromagnetic (AFM) long-range order (LRO) without a spontaneous magnetization appears in AMn<sub>3</sub>P<sub>4</sub>O<sub>14</sub>, while a canted AFM LRO appears in the other substances. We studied magnetism of SrM<sub>3</sub>P<sub>4</sub>O<sub>14</sub> using neutron-scattering measurements. In SrMn<sub>3</sub>P<sub>4</sub>O<sub>14</sub>, we determined the coplanar spiral magnetic structure that had no spontaneous magnetization. We confirm that the spin system is a frustrated J1-J1-J2 trimerized spin chain with dominant AFM J1, small AFM J2, and small next-nearest-neighbor exchange interactions. The magnetic structure (the irreducible representation tau) in the space group P21/c). We consider that the spin system consists of ferromagnetic dimers formed by the dominant J2 interaction and monomers. The magnetization plateau is impossible.

### PH03

# Spin-Peierls-like lattice distortion and incommensurate magnetic structure of geometrically frustrated spinel CdCr<sub>2</sub>O<sub>4</sub>

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ACr<sub>2</sub>O<sub>4</sub> spinels (A = Zn, Cd, Hg) with antiferromagnetic Cr<sup>3+</sup> (S = 3/2) on B sites are well known for strong geometrical frustration preventing long-range order. The frustrations in their cubic phases are finally lifted via spin-Peierls-like lattice distortions at low temperatures. In spite of their similarities in the cubic phases, each of them undergo distinctive lattice distortions and exhibit unique magnetic order. The nature of their magnetic ground states is thus of great importance in order to understand these novel transitions. Nevertheless, the exact spin structure of CdCr<sub>2</sub>O<sub>4</sub> has so far never been reported. We revisited this issue by performing neutron powder diffraction measurement of isotope-enriched sample for reduced neutron absorption. We observed the cubic-to-tetragonal distortion and the subsequent incommensurate magnetic order that are consistent with previous reports. Using representation analysis, we found that the spins on different sublatices maintain nearly collinear arrangements within a single tetrahedron. Whereas the obtained spin structure is different from the one previously conjectured based on spin wave frequencies, it is consistent with the idea that quantum fluctuation plays an important role in the magnetic ground state of CdCr<sub>2</sub>O<sub>4</sub>.

### **PH04**

# Interplay of magnetic order and structural distortions in multiferroic $GdMnO_3$ single crystals

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Multiferroics are candidates for magneto-electrical applications. Magnetic and dielectric phenomena in orthorhombic GdMnO<sub>3</sub> are strongly intertwined. In addition, electrically induced magnons exist. Magnetostriction measurements by high-resolution capacitive dilatometry, performed in the different AFM phases below T<sub>N</sub> = 42 K, show lattice distortions with striction of about 10<sup>4</sup>. All longitudinal and transversal components of the magnetostriction tensor were measured for the first time. Although no changes of the lattice symmetry exist, the measurements reveal magneto-structural phenomena, especially in the canted-AFM phase. A strong anisotropy of the magnetolastic properties was found, in good agreement with the type and propagation of the magnetic structure. As the capacitive method is sensitive to expansion effects and changes of permittivity, dielectric anomalies could be detected and compared to the critical values of H and T of the magnetic transitions. The anomalies especially exist in the range 10....22 K in which the structure is characterized by magnetic order on both the Gd- and Mn-sites. This demonstrates the strong interplay of magnetic, charge and lattice degrees of freedom. First steps in modelling the magnetization of GdMnO<sub>3</sub> by the simulation program McPhase are presented and exchange mechanisms are discussed. These results are examined in connection with the magnetostriction data.

### PH05

# Uniaxial pressure effect on magnetic ordering in a frustrated ising antiferromagnet CoNb<sub>2</sub>O<sub>6</sub>

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Uniaxial pressure effect on magnetic ordering has been investigated by magnetization measurements for a frustrated Ising antiferromagnet CoNb<sub>2</sub>O<sub>6</sub> where quasi-one-dimensional ferromagnetic chains along the c axis form a frustrated isosceles triangular lattice on the a-b plane. Under uniaxial pressure and magnetic field along the a axis, a critical field of magnetic phase transitions from antiferromagnetic to incommensurate magnetic (ICM) phase as well as from ICM to ferromagnetic interchain exchange interactions. Above 700 MPa, a step-like behavior of the magnetization curve shows up at around the half of the saturation magnetization. This behavior implies the appearance of a pressure-induced magnetization (0,0.25,0). These observations were explained within mean field calculations with nearest-neighbor and next-nearest-neighbor interchain exchange interactions.

### **PH06**

#### Thermal and magnetic properties of LiNiPO<sub>4</sub> olivine

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On the contrary to LiMPO<sub>4</sub> olivines with M = Mn, Fe, Co, in LiNiPO<sub>4</sub> two step transformation between the paramagnetic and a uniform in space antiferromagnetic phase occurs. At first, a long-range magnetic order develops in the form of an incommensurate antiferromagnetic phase (via a second-order transition at 21.8 K) and then a first-order transition to the antiferromagnetic phase occurs at 20.9 K. To study thermal properties of LiNiPO<sub>4</sub>, and an influence of magnetic field, B, on the magnetic structure, specific heat measurements of a single crystal were carried out from 3 K to 60 K, in B=0 and in B (up to 9 T) applied along three main crystallographic directions. It was found that B applied along a and b directions has no effect on the transition superimentaries depending parabolically on B). To elucidate these effects, supplementary studies of magnetic torque were performed for several B values, for B rotating within a-c and b-c planes, for a series of fixed temperatures. This work was partly supported by the European Regional Development Fund, through the Innovative Economy Grants: POIG.01.01.02-00-108.09 and POIG.02.02.00-00-025.09.

# **PH07**

# Cr- and mo-doping effects on structural and orbital order phase transition in spinel-type $MnV_2O_4$

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MnV<sub>2</sub>O<sub>4</sub> has orbital degrees of freedom of t2g electrons in V<sup>3+</sup> (3d2), and undergoes magnetic and orbital order. In this study, we investigated the doping effects of Cr<sup>3+</sup> (3d3) and Mo<sup>3+</sup> (4d3), which have no orbital degrees of freedom, for the V<sup>3+</sup> site of MnV<sub>2</sub>O<sub>4</sub>, in order to clarify the Crand Mo-doping effect on the electronic properties in this system. In MnV<sub>2</sub>O<sub>4</sub>, the paramagnetic to collinear ferrimagnetic phase transition and the collinear to noncollinear ferrimagnetic phase transition occur at TN=59K and TOO=54K, respectively. The structural phase transition from cubic to tetragonal lattices takes place at TOO concomitantly with the orbital order. From the magnetic susceptibility and powder X-ray diffraction measurements, it is clarified that the orbital order is suppressed by Cr- and Mo-doping and disappears above x=0.12 (x=0.08) in Cr (Mo)doped samples. The resistivity and optical reflectivity measurements have clarified that the system becomes almost metallic with increasing Mo content, while the Cr-doped samples are still Mott insulators. These results indicate that the doped Cr and Mo with no orbital egrees of freedom suppress the orbital order in MnV<sub>2</sub>O<sub>4</sub>. The introduced 3d electrons remain localized in the Crdoped system, while 4d electrons become nearly titnerant in the Mo-doped one.

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#### **PH08**

# Dopant-dependence on charge/orbital order in impurity doped layered manganites

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A layered manganite  $La_{0.5}Sr_{1.5}MnO_4$  shows charge/orbital order and CE-type antiferromagnetic structure below 240 K and 110 K, respectively. We have studied how the ordering state is changed by the substitution of Cr, Fe and Ga ions for Mn ions using a resonant x-ray scattering (RXS) technique and magnetization measurement. In trivalent ions these impurity ions have no orbital degree of freedom and different spin values, in fact Cr<sup>3+</sup>, Fe<sup>3+</sup>, and Ga<sup>3+</sup> have S = 3/2, S = 5/2, and S = 0, respectively. As a result, the magnetization for Ga doped compound shows the same temperature dependence as pure compound, but in Cr and Fe doped compounds the magnetizations increase below 90 K and 70 K, respectively. On the other hand, in all doped compounds the RXS intensities which reflect charge/orbital order are smaller than that in pure compound, but degrees of decrease are different according to the dopants. Namely, RXS intensity of Ga-doped compound has the same temperature dependence as pure compound, while RXS intensities of Cr- and Fe-doped compounds drastically decrease below 90 K and 70 K, respectively.

#### **PH09**

# Magnetic and nonmagnetic impurity effect on magnetic orderings of isosceles triangular lattice antiferromagnet CuMnO<sub>2</sub>

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In recent years, chemical substitution effect on magnetic and dielectric properties in frustrated magnetic systems has been a vast field of challenging new physics. Good illustrations are the drastic and various effects of a small level of chemical substitution upon the magnetic and dielectric properties on frustrated magnetic systems such as CuFeO<sub>2</sub> and MnWO<sub>4</sub>. These magnetic ground states stabilized with strongly competing exchange interactions should be easily changed by slight chemical disturbance. In the present study, we have investigated magnetic and normagnetic impurity substitution effects on magnetic neutron diffraction measurements on Cu(Mn<sub>1x</sub>Cu<sub>x</sub>)O<sub>2</sub> and Cu(Mn<sub>1y</sub>Ga<sub>y</sub>)O<sub>2</sub>. While the magnetic wave vector of Cu(Mn<sub>1x</sub>Cu<sub>x</sub>)O<sub>2</sub> the magnetic ground state of "magnetically substituted with x=0.04 is Q=(-1/2 1/2 0)[1-3], that in the "nonmagnetically substituted Cu(Mn<sub>1y</sub>Ga<sub>y</sub>)O<sub>2</sub> with y=0.05, 0.15 and 0.20 is Q=(-1/2 1/2 1/2 1)(21) this is identical to that of pure CuMnO<sub>2</sub>. We discuss an influence of magnetic rononmagnetic impurity substitution on the magnetic ordering of Cu(Mn<sub>1x</sub>Ga<sub>y</sub>)O<sub>2</sub> with Cu(Mn<sub>1x</sub>Gu<sub>y</sub>O<sub>2</sub>.

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### PH10

# Resonant soft X-ray scattering studies on half-doped manganite LaSr,Mn,O<sub>7</sub>

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Recently, a careful doping control study, however, reported that the La<sub>2-2x</sub>Sr<sub>1+2x</sub>Mn<sub>2</sub>O<sub>7</sub> system has an exotic spin phase diagram very near the half doping with extremely narrow antiferromagnetic (AFM) phase boundaries at x ~ 0.5 +/- 0.005: the CE-type within the boundaries but the A-type outside. To elucidate complication of this spin phase diagram, we investigated the orbital and AFM ordering behaviors of the half-doped bilayer manganite LaSr<sub>2</sub>Mn<sub>2</sub>O<sub>7</sub> by using Mn L-edges resonant soft x-ray scattering. We confirmed the predicted CE-type AFM order for the true half-doped (x = 0.5) case. Furthermore, we found that such a narrow phase boundary is due to the close competition of the two AFM phases via 3d Mn eg orbital instability.

#### PH11

On the non-idle-spin behavior and the field-induced magnetic transitions of the trimer chain magnet Cu<sub>3</sub>(OH)<sub>4</sub>SO<sub>4</sub> Hyun-joo Koo<sup>1</sup>, Reinhard K Kremer<sup>2</sup> and Myung-hwan Whangbo<sup>3</sup>

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chemistry, Horni Caronna State Oniversity, OSH

The magnetic properties of the Cu(2)-Cu(1)-Cu(2) trimer chain magnet Cu<sub>3</sub>(OH)<sub>4</sub>SO<sub>4</sub> were examined by evaluating the spin exchange constants, the Dzyaloshinskii-Moriya (DM) vectors, and the magnetic anisotropy energies of the Cu(1) and Cu(2) sites on the basis of density functional calculations. Cu<sub>3</sub>(OH)<sub>4</sub>SO<sub>4</sub> is not a idle-spin magnet, the Cu(1) spins have a strong antiferromagnetic coupling in the Cu(1) chain, and the Cu(2) spin arrangements of a Cu(2)-Cu(1)-Cu(2) trimer chain result from the presence of two non-equivalent sets of AFM interactions between the Cu(1) and Cu(2) chains. We propose the probable spin arrangements of a trimer chain responsible for the field-induced successive magnetic transition, and suggest that the nature of the magnetic transition depends on the field-direction because of the DM interactions associated with two strong spin exchange paths.

#### PH12

# Spin state of LaCoO<sub>3</sub> investigated from non-magnetic-ion substitution effect of Co sites

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The perovskite oxide LaCoO<sub>3</sub> is a fascinating material because of a dramatic change of its spin state with temperature variation. In low-temperature range, LaCoO<sub>3</sub> is non-magnetic because the d electrons of Co<sup>3+</sup> (3d6) fully occupy the lower t2g orbitals (LS). On the other hand, the nature of its excited state is controversial. The spin state of this system is also sensitive to the chemical substitutions. Most interestingly, non-magnetic Rh<sup>3+</sup> (4d6: LS) substitution increases the magnetization at low temperature, which is different from Ga<sup>3+</sup> (3d10) substitution effect. We have measured the magnetization and x-ray diffraction of LaCo<sub>1-x</sub>Rh<sub>x</sub>O<sub>3</sub> [1]. We have found a weak ferromagnetic ordering for 0.1 < x < 0.4, indicating that the Rh substitution stabilizes the magnetic state of Co<sup>3+</sup> at low temperatures. The effective magnetic moment of LaCo<sub>1-x</sub>Rh<sub>x</sub>O<sub>3</sub> evaluated at room temperature is independent of x for 0 < x < 0.5, and rapidly decreases above x = 0.5. It indicates that the excited state of Co<sup>3+</sup> is well described as a mixture of the low-and the high-spin Co<sup>3+</sup> ion, and that a Rh<sup>3+</sup> ion favors to substitute for a low-spin

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# PH13

# First-Principles Calculation of the A-Site Ordered Perovskite CaCu, Fe<sub>2</sub>O<sub>12</sub>

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#### PH14

# Soft x-ray absorption spectroscopy study of Prussian blue analogue MCo[Fe(CN)<sub>6</sub>]H<sub>2</sub>O Nano-particles (M=Na, K, Rb)

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Molecular-based magnetic materials have been studied for their possible applications as nano-magnetic devices as well as for the understanding of the fundamental physics of their magnetic phenomena. Molecular magnetic materials of Prussian Blue (PB) analogues, which are represented by the general formula,  $An[B(CN)_{ij}]mxH_{i}O$ , where A and B are transition-metal ions [1], show various characteristic magnetic properties depending on their transition metal ions. PB crystallizes in the fcc structure with space group Fm3m [2]. Of special interest is the photo-induced spin transition. However, the origin of the photo-induced spin transition has not been understood well yet. In this work, we have studied the valence states of PB analogue MCo[Fe(CN)\_i]H<sub>2</sub>O (M=Na, K, Rb) nano-size particles by employing synchrotron-radiation soft x-ray absorption spectroscopy (XAS). The measured Fe and Co 2p XAS spectra of MCo[Fe(CN)\_i]H<sub>2</sub>O (M=Na, K, Rb) reveal the systematic changes as the size of M ion changes, from the Fe<sup>2+</sup> Fe<sup>3+</sup> mixed-valent states and the nearly divalent Co<sup>2+</sup> states to the nearly trivalent Fe<sup>1+</sup> states and the Co<sup>2+</sup>-Co<sup>3+</sup> mixed-valent states. This finding suggests that the electronic structures of PB nano-particles play an important role in determining their magnetic properties.

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### PH15

#### Long-time variation of magnetic structure in multistep metamagnets Ca<sub>3</sub>(Co-M)<sub>2</sub>O<sub>6</sub>: Effect of disorder

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We have reported that CeIr<sub>5</sub>Si<sub>2</sub>, Ca<sub>3</sub>Co<sub>2</sub>O<sub>6</sub> and some other materials show long-time variation of magnetic structure [1-4]. These materials are uniform compounds including no appreciable randomness or imperfections. All of these compounds show successive magnetic transitions and metamagnetic transitions at low magnetic field, which indicate the presence of competing magnetic interactions. Therefore, we think that the long-time variation of magnetic structure in these materials is caused by the competing magnetic interactions. To ensure that randomness or imperfections are not the main factor causing the long-time variation of magnetic structure, we have investigated the element substitution effects. In this presentation, we show the time dependence of the magnetization in Ca<sub>3</sub>(Co<sub>1,x</sub>M<sub>x</sub>)<sub>2</sub>O<sub>6</sub> (M=Ga, Mg, Al and Ti). The time variations of magnetization were measured after the sample was cooled to various target temperatures from T=30 K (T>Tc1) at B=5 mT. Characteristics of the time variation behavior were not much modified by the disorder produced by the substitution of non-magnetic elements. These results show that the long-time variation of the magnetic structure in Ca<sub>3</sub>Co<sub>2</sub>O<sub>6</sub>

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# **PH16**

Study of magnetic properties of NdCo<sub>1-x</sub>Ni<sub>x</sub>O<sub>3</sub> (x = 0, 0.2, 0.4)

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We report here the magnetization studies of the NdCo<sub>1.4</sub>Ni<sub>x</sub>O<sub>3</sub> (x = 0, 0.2, 0.4) samples synthesized using conventional solid state reaction method. X-ray diffraction followed by Rietveld refinement, confirms the single phase orthorhombic structure with Phom space group for all the samples. X-ray absorption spectroscopy studies revealed that both Co and Ni exist in trivalent state. The zero field cooled and field cooled magnetization measurements show that below room temperature the magnetization has strong influence of paramagnetic rare earth ion inhibiting magnetic ordering of cobalt ions. Below a certain low temperature (T< 50K), samples with x = 0, 0.2, show paramagnetic to anti ferromagnetic transition but the sample with x = 0.4 shows paramagnetic to Sing lass transition. Isothermal magnetize behavior has on the basis of various magnetic interactions between Nd-sublattice and those of transition metal networks.

# PH17

# Fluctuations of charge, orbital, and spin order in single layer manganite $Pr_{0.5}Ca_{1.5}MnO_4$

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The magnetic and transport properties of doped manganite oxides are intimately related to the dimensionality of the MnO<sub>2</sub> network. A three dimensional system tends to lead to metallic-ferromagnetic properties, whereas the two dimensional (2D) systems is dominated by charge-orbital order (COO) leading to insulating-antiferromagnetic properties. Pr<sub>0.5</sub>Ca<sub>1.3</sub>MnO<sub>4</sub> is one of the more interesting members of the 2D manganite family showing a charge and orbital ordering transition above room temperature (TCOO = 320 K) with distinct changes in the magnetic, transport, and structural properties, evidencing a strong coupling between the orbital structural, charge, and spin degrees of freedom. We study the structural, magnetic susceptibility, transport, specific heat, and Raman spectroscopy of this material. The metal insulator transition observed at TCOO is found to be more pronounced in comparison to the well studied isostructural La<sub>0.5</sub>Sr<sub>1.3</sub>MnO<sub>4</sub> compound. The same holds for the changes in the magnetic susceptibility at the antiferromagnetic transition (TN=120 K). Even though the COO is accompanied by rather abrupt changes in the physical properties, our data shows that no full ordering occurs at the transition temperature. The system exhibits orbital fluctuations down to a temperature of about 200K, leading to an unusual temperature dependence of the vibrational, magnetization, and transport properties.

# PH18

Magnetic properties of spinel oxide  $\rm CuCr_2O_4$  investigated by NMR

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Spinel oxide CuCr<sub>2</sub>O<sub>4</sub> is known to have non-collinear spin structure and orbital ordering in its ground state. NMR spectrum for the magnetic ions Cr<sup>3+</sup> and Cu<sup>2+</sup> was measured at liquid helium temperature to investigate magnetic properties of CuCr<sub>2</sub>O<sub>4</sub>. Temperature dependence of the Cr NMR frequency shows that there exist initial energy gap in the dispersion relation of the spin wave. This energy gap is larger than that of the Cu ion, implying that spin-orbit coupling of the Cr ion is stronger than that of the Cu ion. This is consistent with theory that predicts orbital order for the Cu ion in CuCr<sub>2</sub>O<sub>4</sub>. The canting angle of the Cr<sup>3+</sup> spin was estimated to be 80° by external magnetic field dependence of the Cr NMR frequency. Also the line width of the Cr NMR spectrum was measured in 7T external magnetic field. The broadening of the spectrum is induced by the anisotropy field. The broadening of the Cu NMR spectrum, implying that orbital ordering, rather than anisotropy field, is the main reason of the Cu NMR spectrum broadening.

# PH19

# Magnetic frustration effects in the new colossal magnetoresistance oxide $NaCr_{2}O_{4}$

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In itinerant electron systems on a geometrically frustrated lattice, the magnetic frustration is considered to play a crucial role for their transport properties as seen in LiV<sub>2</sub>O<sub>4</sub> which shows the heavy fermion-like behavior [1]. Recently, a new chromium oxide NaCr<sub>2</sub>O<sub>4</sub> with geometrically frustrated double chains was reported to show the unconventional colossal magnetoresistance (CMR) effect [2]. Unlike the conventional CMR materials represented by the ferromagnetic manganites, NaCr<sub>2</sub>O<sub>4</sub> is an antiferromagnet with the Neel temperature TN = 125 K. The origin of the CMR effect in NaCr<sub>2</sub>O<sub>4</sub> to investigate the spin structure in the antiferromagnetic phase and the magnetic correlations. We found the presence of a ferromagnetic interaction, which competes with an antiferromagnetic one, from the temperature dependence of the nuclear spin-lattice relaxation rate 1/T1. Below TN, 23Na NMR spectra, which change with increasing the external magnetic field, cannot be explained by a conventional spin structure. The CMR effect may be related to the spin structure in the structure for the spin structure in the structure.

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# PH20

# Site-dependent metal-insulator transition and orbital order in quasi-one-dimensional $V_6O_{13}$

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Vanadates are prototypes of metal-insulator transitions with orbital degrees of freedom. In a quasi-one-dimensional compound  $V_6O_{13}$ , the mixed valence gives non-uniform spin-charge distributions on inequivalent V sites even for the metallic state. Our sophisticated 51V NMR measurements using a single crystal distinguished the site dependent orbital order and the local spin susceptibilities on two inequivalent chains. One chain behaves as a metallic one with the moderately depressed susceptibility, whereas another behaves as an insulating one with the magnetically active local moment. The result arises from the site dependent electron correlation, giving the net duality in the conducting and magnetic properties. A metal-insulator transition occurring at -150 K accompanies with dramatic changes in the charge and orbital ordering textures, as we observed in the 51V nuclear quadrupole splitting anisotropy. Such ordering separates the metallic chain into magnetic and nonmagnetic sites, giving a complex antiferromagnetic spin structure below -50 K.

#### PH21

# Synthesis and characterization of $\rm CoMn_2O_4$ nanopowders by a reverse Micelle processing

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 $CoMn_2O_4$  nanoparticles was synthesized by reverse micelle processing the mixed precursor (consisting of  $Co(NO_3)_3$  and  $MnCl_3$ ·4H\_2O). The  $CoMn_2O_4$  was prepared by mixing the aqueous solution at a molar ratio of Co: Mn = 1: 2. The synthesized powders were calcined at 600°C for 2h. The average size and distribution of synthesized powders were in the range of 10-20nm and narrow, respectively. The average size of the synthesized powders increased with increasing water to surfactant molar ratio. The XRD diffraction patterns show that the phase of  $CoMn_2O_4$  was spinel (JCPDS no.77-0471). The synthesized and calcined powders were characterized by thermo gravimetry- differential scanning calorimeter (TG-DSC), X-ray diffraction analysis (XRD) and transmission electron microscopy (TEM). The magnetic property of the powder was measured by Vibrating Sample Magneto-meter (VSM) at 298K. The effect of synthesis parameter, such as the molar ratio of water to surfactant, is discussed.

# PH22

# The patterning with a circular magnet array, its observation and domain switching in ferromagnetic ZnCoO:H

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Today's spintronic technologies are considered to be the most progressive solution for overcoming the typical semiconductor electronics. Many researchers have extensively studied the effect of size, shape, structure on single-domain magnet and its wall motion with permanent magnetic metal for high-density magnetic memory and logic devices [1,2]. However, the fact that permanent magnetic metals could not reprogram after the first design for logic devices is still remaining as an important issue. Thus, hydrogen-mediated Co-doped ZnO(ZnCoO/H) having reversible room-temperature ferromagnetism has been considered as a strong candidate material for spintronic devices[3]. In this study, we observed domain switching in a ZnCoO/H in real-space and simulate it using object-oriented micromagnetic framework(OOMMF). ZnCoO circular patterns are fabricated by rf-sputtering and various lithography processes, and hydrogen is injected by plasma and hot isostatic pressing(HIP) methods. The reversible magnetic single-domain behavior is activated and deactivated by hydrogen manipulation, which are investigated with superconducting quantum interference device(SQUID) and magnetic force microscopy(MFM) measurements, and the results are compared with the OOMMF simulation. This reversible room-temperature magnetic switching of the oxide materials could greatly contribute to the development of reprogrammable spin logic devices.

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#### PH23

# Orbitally induced molecule formations in itinerant triangular vanadates

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An intriguing route that eliminates the geometrical frustration is the orbitally-induced symmetry breaking usually described by the Kugel-Khomskii model. In itinerant systems, however, the model for localized electrons may not be guaranteed, and the degenerated d orbitals can be partially occupied. We show that the orbital occupation continuously changes into a valence-bond-solid ground state in an itinerant vanadate BaV<sub>10</sub>O<sub>15</sub> with a triangular lattice, based on the orbital-resolved nuclear magnetic resonance technique. Owing to the odd number of d electrons, i. e., eleven electrons for a five-V unit, all the V sites cannot form the valence bonds. Instead, two of the five V sites, sharing five electrons, remain magnetic with the ferro-orbital ordering [1]. In clear contrast, no commensurate spin order and orbital order are stabilized in SrV<sub>10</sub>O<sub>15</sub>. The result suggests fractional and incommensurate spin-orbital orders that cannot fully remove the geometrical frustration.

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#### PI01

Investigation of magnetocaloric effect in La<sub>0.45</sub>Pr<sub>0.25</sub>Ca<sub>0.3</sub>MnO<sub>3</sub> by differential scanning calorimetry and thermal analysis Aparnadevi M, Sujit Kumar Barik and Ramanathan Mahendiran\* Department of Physics, National university of Singapore, Singapore

We investigated magnetocaloric effect in La<sub>0.45</sub>Pr<sub>0.25</sub>Ca<sub>0.3</sub>MnO<sub>3</sub> by direct (changes in temperature and latent heat) and indirect methods (magnetization isotherms). This compound undergoes a first-order paramagnetic to ferromagnetic transition with Tc = 200 K upon cooling. The paramagnetic phase becomes unstable and it transforms into a ferromagnetic transition (FIMT). The FIMT is accompanied by release of latent heat and change in temperature of the sample as evidenced from differential scanning calorimetry and thermal analysis experiments. A large magnetic entropy change of  $\Delta$ Sm = -7.3 Jkg-1K-1 at Tc = 212.5 K and refrigeration capacity of 228 J/ kg are observed for a field change of  $\Delta$ H = 5 T. It is suggested that the nanometer sized ferromagnetic clusters co-exist with short-range charge-ordered clusters in the paramagnetic state, resulting in a large value of  $\Delta$ Sm above the Tc.

# PI02

### Specific heat and magnetic properties of spinel compound $FeV_2O_4$

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The spinel compound FeV<sub>2</sub>O<sub>4</sub> has the orbital degrees of freedom at both eg (Fe<sup>2+</sup>) and t2g (V<sup>3+</sup>) orbitals. This compound shows successive structural phase transitions at about 140 K, 110 K accompanied by ferrimagnetic ordering, 70 K and 35 K. However, the details of the phase transitions below 110 K are still not understood. In order to clarify the origin of these phase transitions, we measured magnetization and specific heat capacity for polycrystalline and single crystal of FeV<sub>2</sub>O<sub>4</sub>. We first observed three anomalies at 140 K, 109 K and 68 K with decreasing temperature in the specific heat measurements using the polycrystalline sample. The  $\lambda$ -type anomaly at 109 K indicates the second order phase transition, which is related to the ferrimagnetic ordering. Moreover, the anomaly at 68 K has a sharp peak, which conforms to the anomaly on the temperature dependence of magnetic susceptibility. The electronic specific coefficient  $\gamma$  =1.8 mJ/molK2 of FeV<sub>2</sub>O<sub>4</sub> is smaller than that of of MNv<sub>2</sub>O<sub>4</sub> ( $\gamma$  = 12 mJ/molK2)[2] because the Jahn-Teller effect of Fe<sup>2+</sup> for FeV<sub>2</sub>O<sub>4</sub> may weaken the electron electron interactions. The details of the magnetic structures are also discussed.

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#### PI03

# The crossover to checkerboard charge order: Magnetic excitations of charge-stripe ordered of La<sub>2x</sub>Sr<sub>x</sub>NiO<sub>4</sub>

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Are stripes ubiquitous to the cuprates, and what is the role of stripes in high temperature superconductivity, are two questions that have been intensively debated since the discovery of charge stripe order in the cuprates[1]. From studying the magnetism of charge-stripe ordered La<sub>6.3</sub> Sr,NiO<sub>4</sub> (LSNO) the magnetic interactions of a charge-stripe ordered phase have been quantitatively described[2], although our knowledge is incomplete[3]. For x = 1/2 the ground state order of LSNO is part stripe and part checkerboard like in character, with both the magnetic excitations from the ordered moments and charge-stripe periodicity, discommensurations, create additional magnetic excitations from the ordered moments in the x = 1/2 compared to damping the excitations at x ~ 1/3. While the magnetic excitations from the charge stripe electrons have a different wavevector centring in the x = 1/2 [3,4]. We report a neutron scattering study of the magnetic excitations in LSNO x = 0.4 and 0.45. Between these two doping levels the magnetic excitations show a dramatic crossover in character. We discuss this crossover behaviour with regards to charge-stripe materials with larger charge stripe gacing.

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#### PI04

#### Magnetic properties of low-dimensional $\alpha$ and $\gamma$ CoV<sub>2</sub>O<sub>6</sub>

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 $CoV_2O_6$  is a low dimensional oxide existing in two allotropic phases, generally called  $\alpha$  and  $\gamma$ , and showing monoclinic and triclinic structures, respectively. We synthesized this oxide by solid state reaction and performed neutrons diffraction and magnetic measurements in order to understand the magnetic behavior in correlation with the low dimensional structure. The magnetic properties were supported by ab initio calculations. Both phases are constituted of parallel 1D Co chains organized in planes that are separated by vanadium oxide thin layers and are antiferromagnetic in the ground state. The magnetization curves recorded at 5 K show a stepped variation with sharp field-induced magnetic transitions and a magnetization plateau at one-third of the saturation magnetization. In  $\alpha$ -CoV<sub>2</sub>O<sub>6</sub>, additional steps are evidenced at 1.8 K. The magnetic moment per Co ion is larger in the  $\alpha$  phase (4.5 µB) than in the  $\gamma$  phase (3 µB), although Co has the same octahedral environment in both cases. Ab initio calculations show that  $\alpha$ -CoV<sub>2</sub>O<sub>6</sub> is an antiferromagnetic semiconductor with a gap of 1.1 eV in which the magnetic moment is essentially carried by Co and can reach a maximum value of about 4.2 µB for a meta-stable solution.

M. Lenertz, J. Alaria, D. Stoeffler, S Colis and A. Dinia, J. Phys. Chem. C, 2011, 115, 17190-17196

### PI05

Structural and magnetic properties of the parent compound  $T^{*}\mbox{La}_2 CuO_4$  of electron-doped cuprates

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Recently, we found that the newly synthesized metastable T-La<sub>2</sub>CuO<sub>4</sub> exhibits much lower Neel temperature (TN) and increased magnetic fluctuations compared to other mother compounds of electron-doped superconductors like Nd<sub>2</sub>CuO<sub>4</sub> or Pr<sub>5</sub>CuO<sub>4</sub>. Muon spin rotation (USR) neveals a gradual slowing down of dynamic magnetic fluctuations below TN1=220K and static magnetic order below TN2=115K in contrast to T-La<sub>2</sub>CuO<sub>4</sub> where TN2=300K. In comparison to our T'-La<sub>2</sub>CuO<sub>4</sub>, measurements were done on T'-Pr<sub>5</sub>CuO<sub>4</sub> powso scillations at high temperatures up to 220K in sharp contrast to the predominantly dynamic character of the depolarization in T'-La<sub>2</sub>CuO<sub>4</sub> in the same temperature range. This clearly demonstrates that the magnetism of T'-La<sub>2</sub>CuO<sub>4</sub> in the same temperature range. This clearly demonstrates that the magnetism of T'-La<sub>2</sub>CuO<sub>4</sub> in the same temperature and the T'-Pr<sub>5</sub>CuO<sub>4</sub> compounds consistent with a reorganization of the magnetic structure at this temperature. In addition, µSR and neutron scattering results on T'-Pr<sub>5</sub>CuO<sub>4</sub> indicate a similar behaviour previously observed for Nd<sub>4</sub>CuO<sub>4</sub> pointing to, up-to-now undetected spin reorientation transitions.

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### **PI06**

# Microscopic magnetic nature of the quasi-one-dimensional antiferromagnet $BaCo_2V_2O_8$

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BaCo<sub>2</sub>V<sub>2</sub>O<sub>8</sub>, belongs to a wide group of quasi-1D antiferromagnets (AF). The Q1D compounds display a variety of fascinating ground states governed by the strong spin-spin coupling along the 1D direction and a much weaker coupling along other directions. BaCo<sub>2</sub>V<sub>2</sub>O<sub>8</sub> display a long-range AF order below TN=5 K and possibly short-range order all the way up to 30 K. Further, a novel type of field induced magnetic order has been found for T<1.8 K and Hc>3.9 T[1.2]. It was determined to be an incommensurate spin structure caused by quantum fluctuations, fitting well to theoretical predictions for a so-called Tomonaga-Luttinger liquid (TLL). To the best of our knowledge, we present here the first  $\mu$ SR investigation of the microscopic magnetic nature of single crystalline BaCo<sub>2</sub>V<sub>2</sub>O<sub>8</sub> samples. Our data reveal several clear muon frequencies below TN indicating the onset of a long-range order. Above 5 K, the  $\mu$ SR spectra are well fitted to a simple power-exponential relaxing function. The temperature dependence of the relaxation-rate as well as the power display a clear anomaly around T=40 K, indicating the onset of short-range 1D correllations. Finally we also present initial and intriguing field dependent data.

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# **PI07**

#### Detection of orbital wave in YVO<sub>3</sub> using inelastic neutron scattering

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We focus on an orbital excitation in YVO<sub>3</sub>, which shows the complex magnetic and orbital orderings in the low temperatures [1]. In the phase where the G-type orbital order accompanied with the C-type spin order (G-OO/C-SO) appears, the existence of the large orbital fluctuation is suggested [2]. Additionally, a large dispersive orbital-wave along the c-axis, due to the strong one-dimensional spin-orbital correlation, is calculated [3]. In consideration of the neutron scattering cross-section obtained from the correlation function for the orbital angular moment, we have attempted to detect the orbital excitation using inelastic neutron scattering spectrometers. We observed a magnetic excitation up to 35 meV. This energy range is higher than that reported previously [3]. On the other hand, weak excitations are observed in the range of 40-70 meV. In this Presentation, we discuss whether these are originated from orbital excitation or not.

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# **PI08**

#### Orbital occupation and magnetism of tetrahedrally coordinated Fe in CaBaFe.O.

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 $CaBaFe_4O_7$  is a new mixed-valent transition metal oxide from the class of Swedenborgites [1,2] having both Fe2+ and Fe3+ ions in tetrahedral coordination situated in a Kagome lattice. This class of materials could be an interesting starting point to study orbital physics in tetrahedral coordination. Here we characterize its magnetic properties by magnetization measurements and investigate its local electronic structure using soft x-ray absorption spectroscopy at the Fe L2,3 edges, in combination with multiplet cluster simulations and band structure calculations [3]. We found that the Fe<sup>2+</sup> ion in the unusual tetrahedral coordination is Jahn-Teller active having a minority-spin electron

[1] B. Raveau et al., Chem. Mat. 20, 6295 (2008) [2] B. Raveau et al., Z. Anorg. Allg. Chem. 635, 1869 (2009) [3] N. Hollmann et al., Phys. Rev. B 83, 180405 (2011)

by multiplet interactions, thereby explaining the observed magnetic anisotropy.

with an x2-y2 character. We deduce that there is an appreciable orbital moment caused

# PI09

#### Search for topological spin order in the multiferroic insulator Cu<sub>2</sub>OSeO<sub>3</sub>

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We report a comprehensive study of the magnetic properties of single crystals of the multiferroic insulator Cu2OSeO3. Our study was motivated by the recent discovery of a skyrmion lattice phase in the B20 transition metal Si and Ge transition metal compounds [1], which are isostructural siblings of Cu2OSeO3. For our study large high-quality single crystals were grown by vapour transport. The magnetisation, ac susceptibility and specific heat were measured down to 2K under magnetic fields up to 9T. The magnetic properties were studied with small angle neutron scattering. We find a paramagnetic to helimagnetic transition at T\_c - 58K. The magnetic phase diagram of Cu<sub>2</sub>OSeO<sub>3</sub> shares remarkable similarities with the B20 transition metal silicides, including the existence of a small phase pocket below T c. The putative evidence for topological spin order is compared with the evidence for a skyrmion lattice in B20 transition metal Si and Ge compounds.

[1] S. Muehlbauer et al., Science 323, 915 (2009)

#### **PI10**

#### Studies of neutron scattering and bulk properties of honeycomb lattice Li<sub>2</sub>MnO<sub>3</sub>

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Li<sub>2</sub>XO<sub>2</sub> with transition metal ions occupying the X position forms an interesting honeycomb lattice. Among Li<sub>2</sub>XO<sub>3</sub>, Li<sub>2</sub>RuO<sub>3</sub> was recently found to have an unusual structural phase transition involving probably Ru dimerization. For example, it was reported that the crystal structure changes from P21/m to C2/m near 540 K, accompanied by a concomitant metal-insulator transition and an abrupt change in the magnetic susceptibility. Li<sub>2</sub>MnO<sub>2</sub> is another example having the C2/m crystal structure with the honeycomb lattice. In order to investigate the magnetic properties of Li<sub>2</sub>MnO<sub>3</sub>, we have successfully grown single crystal Li<sub>2</sub>MnO<sub>3</sub> by using a flux method. Using these single crystals, we have carried out both bulk measurements and neutron diffraction studies with high resolution powder diffractometer (HRPD) and single crystal diffractometer (FCD) at HANARO, Korea. We also measured the magnetization and heat capacity of the single crystals

# **PI11**

#### Magnetic and calorimetric properties of Mn<sub>2</sub>GeO<sub>4</sub> single crystals Natalia Mihashenok\*, Nikita Volkov, Klara Sablina, Alexander Balaev, Maxim

Molokeev, Sergei Popkov and Dmitriy Velikanov L.V. Kirensky Institute of Physics SB RAS, Russia

Within the framework of our research of manganese oxides the single crystals  $Mn_2GeO_4$  have been grown by the flux method using the original technique [1]. The crystal structure an isomorph of olivine was determined on single crystal by X-ray diffraction method: cell parameters, a=10.7401 A, b=6.3116 A, c=5.0766 A, and the space group Pnma. Magnetic measurements were performed in the temperature range from 2 to 300 K and in magnetic fields up to 50 kOe with the exact orientation of the applied magnetic field relative to the crystallographic directions of the crystal. There are three sharp magnetic phase transition at T1=47 K T2=17.5 K and T3=5.5 K with phase transition at T2 depended on applied magnetic field. Also the specific heat measurements of single crystals were carried out. The heat capacity of Mn2GeO4 exhibits three sharp maxima, which exactly correlate with the magnetic phase transition

[1] Sapronova N.V., Volkov N.V., Sablina K.A., Petrakovskii G.A., Bayukov O.A., Vorotynov A. M., Velikanov D.N., Bovina A.F., Vasilyev A.D., Bondarenko G.V. Phys. Stat. Sol. B., 246 (2009) 206

#### **PI12**

#### Effect of doping on the magnetic structure of $YMn_{1,x}M_xO_3$ (M = Ga, Ti, x < 0.1)

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The effect of doping at the Mn site of multiferroic manganites, YMn<sub>1-x</sub>M<sub>x</sub>O<sub>3</sub> (M = Ga, Ti, x  $\leq$  0.1) have been studied by neutron powder diffraction and magnetic measurements. These compounds have been prepared by solid state synthesis techniques. All the compounds studied are isostructural and crystallizes with hexagonal structure in P63cm space group. We find that both  $Ga^{3+}(d10)$  and  $Ti^{4+}(d0)$  doping leads to significant reduction in the magnetic transition temperature (TN) while they influence the magnetic structure of YMnO3 differently. The magnetic structure of YMnO<sub>3</sub> is described by the irreducible representation  $\Gamma1$  (or  $\Gamma3$ ) with moment on Mn~3µB at 5K, in agreement with previously reported values on this compound [1]. We find that the representation remains the same on doping with Ga in YMnO3 albeit with a decrease in the value of the moment to 2.5µB at 6K. However, the irreducible representation changes on doping with Ti, and this modified structure is given by the basis vectors of the irreducible representation Γ2 with moment 2.1µB at 6K. The geometrical frustration parameter,  $f = \theta p/Tn$ , ~ 7 and reduces with doping. These observations are different from the earlier reported doping studies carried out in this compound [2].

[1]. J.A. Alonso et al. Phys. Rev. B 62 9498 (2000) [2] J. Park et al., Phys. Rev. B 79 064417 (2009)

### PI13

Neutron diffraction and magnetic properties of Ba<sub>2</sub>Co<sub>2</sub>Fe<sub>12</sub>O<sub>22</sub>: Co<sub>2</sub>Y Chan Hyuk Rhee<sup>1</sup>, Jung Tae Lim<sup>1</sup>, Sung Wook Yoon<sup>1</sup>, Kwang Lae Cho<sup>1</sup>, Sung Baek

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Y-type hexa-ferrites with non-collinear magnetic structure have been studied for magnetoelectric effect, showing induced electric polarization under external magnetic field. The helical magnetic structure of these ferrites depends on the temperature and magnetic field [1]. Here, polycrystalline Y-type barium cobalt ferrite (Ba2Co2Fe12O22; Co<sub>2</sub>Y) was synthesized by conventional ceramic method in O<sub>2</sub> gas atmosphere to reduce the oxygen defect. The temperature and magnetic field dependence of magnetic structure was investigated by neutron diffraction and vibration sample magnetometer. At temperature below TN = 615 K, the crystal structure of Co<sub>2</sub>Y was determined to be rhombohedral with the space group R-3m. It showed soft ferrimagnetic behavior with Hc of 145 Oe at 297 K. Most of super-lattice peaks from spin structure decreased with increasing temperature. However, the peak at 21.10 increased with increasing temperature at temperatures above 205 K. In addition, the slope of zero-field cooled magnetization under low field of 0.01 T changed at 205 K. This magnetic structure transition temperature of Co2Y is higher than that of Ba2Mg2Fe12O22 [1] due to the presence of cobalt ion with strong magnetic anisotropy

[1] S. Ishiwata, Y. Taguchi, H. Murakawa, Y. Onose, and Y. Tokura, Science 319, 1643 (2008)

# **PI14**

# Mössbauer studies of olivine Fe1-yMnyPO4

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The olivine structured  $Fe_{1-y}Mn_yPO_4$  (y = 0.1, 0.3), a possible cathode material for lithium ion secondary battery, have been studied by x-ray diffraction and Mössbauer spectroscopy. These Fe1. Mn PO4 samples were prepared by chemical deintercalation lithium from the LiFe, Mn.PO, The crystal structures of the Fe, Mn.PO, samples were determined to be orthorhombic (space group Pnma) at 295 K by Rietveld refinement method. The lattice constants of the Fe. Mn PO, samples increased from a0 = 9.833 Å, b0 = 5.811 Å and c0 = 4.786 Å for y = 0.1 to a0 = 9.979 Å, b0 = 5.895Å and c0 = 4.799 Å for y = 0.3 with the increasing Mn concentration. From the Mössbauer spectra at 295 K, the electric quadrupole splitting (AEQ) and isomer shift ( $\delta$ ) values of the Fe1-vMnvPO4 were determined to be  $\Delta$ EO = 1.50 mm/s  $\delta$  = 0.31 mm/s for y = 0.1 and  $\Delta EQ$  = 1.35 mm/s,  $\delta$  = 0.31 mm/s for y = 0.3. The decrease in ΔEO of Fe. Mn PO, samples can be explained by the exchange interaction due to the asymmetry in FeO<sub>4</sub> octahedral sites depending on the Mn concentrations

[1] B. C. Melot, G. Rousse, J.-N. Chotard, M. Ati, J. Rodriguez-Carvaial, M. C. Kemei, and J.-M. Tarascon, Chem. Mater. 23, 2922 (2011). [2] W. Kim, C. H. Rhee, H. J. Kim, S. J. Moon, and C. S. Kim, Appl. Phys. Lett. 96, 242505 (2010).

#### PI15

#### Three- dimensional electronic structure of Na<sub>0.85</sub>CoO<sub>2</sub>

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### **PI16**

Formation of stoichiometric of FeO synthesized under high pressure Yasushi Kanke<sup>1</sup>, Takuro Yoshikawa<sup>2</sup>, Hideto Yanagihara<sup>2</sup>, Eiji Kita<sup>2</sup>, Yorihiko Tsunoda<sup>3</sup> Kiiti Siratori<sup>2</sup> and Kay Kohn<sup>3</sup> National Institute of Materials Science, Japan Institute of Applied Physics, University of Tsukuba, Japan <sup>3</sup> School of Science and Engineering, Waseda University, Japan

FeO has the simplest NaCl crystalline structure and one electron exists in the degenerated de orbitals due to the regular octahedral ligands. Even such simple structure, the partial quenching of the orbital angular momentum (S = 2 and L = 1) bring the electronic structure complex.[1] Difficulty in the sample preparation, only the samples Fe. O with  $X \le 0.96$  have been obtained makes the experimental investigation of FeO hard. We have prepared the stoichiometric FeO with the high pressure synthesis technique.[2] The starting material of the mixture of α-Fe and Fe<sub>2</sub>O<sub>2</sub> powder was sealed in a Pt cansule and was heated up to 1000 °C and 1200 °C under 5.5GPa for 1 hour. The residual ferromagnetic parts were estimated to be less than 0.69 at% Higher reaction temperature may bring lower concentration of ferromagnetic components From Mossbauer study, it is found that the Koch-Cohen cluster does exist especially when the reaction temperature is high. We present the influence of post annealing with various conditions for the improvement of degree of stoichiometry

[1] J. Kanamori: Prog. Theo. Phys. 17, 177 (1957) [2] T. Katsura, B. Iwasaki, S. Kimura, S. Akimoto : J. Chem. Phys. 47, 4559 (1967)

# PI17

#### Charge and spin ordering in Sr<sub>4</sub>Fe<sub>4</sub>O<sub>11</sub> system

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Ordering and disordering of charge, spin, lattice and orbital degrees of freedom play an important role in the physics of strongly correlated electron systems. In the mixed valence (Fe<sup>3+</sup>/Fe<sup>4+</sup>) compound Sr<sub>4</sub>Fe<sub>4</sub>O<sub>11</sub>, the spin and charge ordering have recently generated a debate arising from its two crystallographic (pyramidal and octahedral) sites and the distribution of Fe<sup>3+</sup> and Fe4+ species [1-5]. Due to symmetry reasons, either the Fe4+ sublattice is frustrated while Fe<sup>3+</sup> ions order antiferromagnetically or the opposite. We are trying to explore this puzzle by X-ray Absorption Spectroscopy, Photoemission Spectroscopy, muSR, Neutron, Synchrontron X-ray diffraction and MOssbauer measurement on precisely adjusted oxygen stoichiometric polycrystalline samples and single crystals.

1 J. P. Hodges, S. Short, J. D. Jorgensen, X. Xiong, B. Dabrowski, S. M. Mini, and C. W. Kimball. J. Solid. State. Chem. 151,190-209 (2000) 2 M Schmidt, M Hofmann and S J Campbell. J. Phys.: Condens.Matter 15, 8691-8701 (2003) 3 P Adler, A Lebon, V Damljanovic, C Ulrich, C Bernhard, AV Boris, A Maljuk, CT Lin, and B Keimer. Phys. Rev. B 73, 094451(2006) 4 R. l'idya, P. Ravindran, H. Fjellvag, and A. Kjekshus. Phys. Rev. B 74, 054422 (2006) 5 P.Adler Comment Phys. Rev. B 77, 136401 (2008) and R. Vidya Reply Phys. Rev. B 77, 136402 (2008)

### **PI18**

#### Catalyst structure determination from magnetic properties

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SbVO<sub>4</sub> series of compounds could be used as oxidation catalysts of industrial interest yielding 20% cheaper acrylonitrile (production 8 million ton/year) by the ammoxidation of propane compared to current method based on propylene as starting reactant. SbVO4 structure was just known to be related to tetragonal rutile. Nevertheless we have shown that actually there is a non-stoichiometric flexible series described as  $Sb_{0.9}V_{0.9+x}\square_{0.2-x}O_4$  (0<x<0.2), where cation vacancies ( $\square$ ) are introduced in the basic rutile type-structure following the mechanism:  $4V(3+) \rightarrow 3V(4+)+\Box$ , while antimony remains as Sb(5+). In this sense, for the compounds synthesized in oxidizing conditions the presence of vacancies has been confirmed, however no cation vacancies have been observed in the reduced phase but alternating Sb/V cation order along the crystallographic c-axis. On the other hand, our magnetic susceptibility studies could distinguish between different phases including possible existence of magnetic order. Using neutron diffraction we determined the magnetic structure of reduced SbVO4 coming from the ordering of vanadium magnetic moments, taking place at TN~50K, and we studied how the substitution of  $V^{3+}$  (S=1) by  $V^{4+}$  (S =  $\frac{1}{2}$ ) while the synthesis conditions become more oxidizing, affects the spin arrangement in SbVO<sub>4</sub>. The structure-property relationship could be evaluated thanks to the magnetic behaviour.

### **PI19**

#### Synthesis and characterization of the mixed perovskite Ba<sub>1</sub>. <sub>x</sub>La<sub>x</sub>Ti<sub>1/2</sub>Mn<sub>1/2</sub>O<sub>3</sub> as a function of La-doping

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Oxides materials with perovskite structure are outstanding examples of materials used in technological applications, as well as in fundamental studies in condensed matter physics due to the great number of observed ground states such as: multiferroicity, high-temperature superconductivity, colossal magnetoresistive effects, multiferroicity, and many others. One possible way to work and study on the coexistence of electric and magnetic orders within metal oxides perovskites may be making "mixed" perovskites with d^0 and d^n ions. In this work we present the results of the synthesis and structural characterization (x-ray powder diffraction and Raman spectroscopy analysis) within the family of transition metal oxides Ba1.  $La_{10}Mn_{10}O_3$  (x = 0, 0.1, 0.2, 0.3, 0.4, 0.5, 0.6) as we dope the 12R-type perovskite BaTi<sub>10</sub>Mn<sub>10</sub>O<sub>3</sub> with the non-magnetic ion La3+. Our results show the evolution of the structure symmetry from the hexagonal one for the non-doped (x=0) sample [1] to a tetragonal one for the highest La concentration as well as the mixture of these two phases for intermediate doping. These results are compared with the appearance of ferromagnetic and antiferromagnetic orders as a function of La-content, temperature and magnetic field.

[1] G. M. Keith, C. A. Kirk, K. Sarma et al., Chem. Mater. 16, 2007 (2004).

# **PI20**

Theoretical modeling of the magnetic properties and magnetocaloric effect in La<sub>01</sub>Ca<sub>00</sub>MnO<sub>3</sub> manganite by Monte Carlo study Oksana Pavlukhina\*, Vacily Buchelnikov and Vladimir Sokolovskiy Chelvahinsk State University, Russia

Manganites have a special place among various materials which have a magnetocaloric effect. This materials, first, allow varying temperature of phase transitions in wide region of temperatures, and secondly, they are cheap and ecological. Experimental studies have shown that the manganites are also attractive for the application in magnetic refrigeration [1]. In this work, using the Heisenberg's Hamiltonian, we investigated by Monte Carlo method magnetic and magnetocaloric properties of La<sub>0.1</sub>Ca<sub>0.9</sub>MnO<sub>3</sub>. In the simulation, magnetic Mn<sup>4+</sup> and trivalent Mn<sup>3+</sup> ions are described by classical Heisenberg spins, while oxygen and calcium ions are considered as non-magnetic. Mn magnetic ions are distributed on a lattice according to the perovskite structure of the manganite. For the modelling lattice samples of size L3 with L=15 were used. Around 5x105 Monte Carlo steps per spin were considered in order to compute equilibrium averages. Curie temperature obtained during the theoretical simulations (TC ~150 K) agrees well with theoretical result for this compound (TC ~150 K) [2]. The Curie temperature and magnetocaloric effect obtained during the theoretical simulations agree well with experimental data. Support by RFBR grants 10-02-96020-r-ural, 11-02-00601 and Federal Target Program № 14.740.11.1442 (03.11.2011) is acknowledged.

1 Phan, Manh-Huong. 2007, Review of the magnetocaloric effect in manganite materials, J. Magn. Magn. Mater. 308: 325-340. 2 Restrepo-Parra E. 2011, Magnetic phase diagram simulation of La1-xCaxMnO3 system by using Monte Carlo, Metropolis algorithm and Heisenberg model, J. Magn. Magn. Mater, 323:1477-1483.

# **PI21**

Pressure effects on magnetic ordering transitions of bilayer manganites Pr(Sr<sub>1-x</sub>Ca<sub>x</sub>)<sub>2</sub>Mn<sub>2</sub>O<sub>7</sub> (x=0,1) by neutron diffraction and muon spin rasonance

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<sup>3</sup> Laboratory for Developments and Methods, Paul Sche

# **PI22**

#### Magnetization reversal and chemical pressure effect in the electron doped manganite Ca(Mn<sub>1-x</sub>Sb<sub>x</sub>)O<sub>3</sub>

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In ref.1, we have demonstrated the effect of hydrostatic pressure of magnetic and transport properties, and thermal transport properties in the electron doped manganite  $Ca(Mn_1,Sb_2)O_2$ . The substitution of  $Sb^{5+}$  ion for Mn4+ site of the parent matrix causes one-electron doping since the X-ray photoemission spectroscopy of Sb substituted CaMnO<sub>2</sub> sample reveals that the valence of Sb ion is 5+. Anomalous magnetization reversals are observed at x=0.05 and 0.08 in the field cooled magnetization while the application of external pressure induces a suppression of the negative magnetization. The effect of chemical pressure on magnetization of  $Ca(Mn,Sb)O_2$  are examined through the substitution of Ca ion for Sr with larger ion radius and are compared with the effect of external pressure on that of the same compound. We believe that the magnetization reversal is not caused by ferrimagnet model but is associated with the presence of the Jahn-Teller type orbital lattice coupling between 3d eg electron and MnO<sub>4</sub> octahedron, giving stabilization of negative canted spin state through the Dzyaloshinsky-Moriya interaction.

Ref.1, Y.Murano et al., Phys.Rev.B83,054437 (2011).

# **PI23**

#### The effect of cu substitution on the structural, electrical and magnetic properties of LaMn1, Cu,O3 manganites

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The hole-doped manganites, La<sub>1-x</sub>A<sub>x</sub>MnO<sub>3</sub> (A = Sr, Ca, Ba or vacancies) with a Perovskite structure have attracted considerable attention due to the discovery of the phenomenon of colossal magnetoresistance (CMR) and its potential application [1-3]. In this study, the effect of Cu substitution on structural, electrical and magnetic properties of LaMn<sub>1-x</sub>Cu<sub>x</sub>O<sub>3+ $\delta$ </sub> (0 $\leq$  x  $\leq$  0.075) manganites are investigated by XRD, electrical resistivity and Ac susceptibility measurements. The XRD refinement result indicates that the samples are single phase and the lattice parameters and volumes increase by the increase of the Cu doping level. The resistivity measurement results show that by increasing Cu doping level, the resistivity decreases and the heavily doped samples show metal - insulator transition at low temperatures. The paramagnetic- ferromagnetic transition temperature, Tc increases for low level doped samples and decreases for heavily doped samples. The analysis of the frequency dependent Ac susceptibility measurements by critical slowing down model, indicates that the reentrant spin glass (RSG) state exists in x=0.025, 0.05 and 0.075 samples. The RSG state is mainly ascribed to the coexistence of ferromagnetic (FM) and antiferromagnetic (AFM) phases and increase of disorder in the FM matrix induced by the random Cu impurity.

#### **PI24**

#### Investigation of the structural and magnetic peroperties of La<sub>0.9</sub>Sr<sub>0.1</sub>MnO<sub>3</sub> nanoparticles

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In recent years, magnetic nanoparticles have been the subject of intense research not

only for their fundamental scientific interest such as superparamagnetism, but also for their potential applications in magnetic storage media, biosensor devices and medical applications, such as targeted drug delivery and hyperthermia [1-3]. In a large number of magnetic nanoparticles applications, it is important to know the effects of interaction between nanoparticles on physical properties of these systems. In this paper, we studied the structural and magnetic properties of La<sub>0.9</sub>Sr<sub>0.1</sub>MnO<sub>3</sub> nanoparticles. Magnetic nanoparticles of La<sub>0.9</sub>Sr<sub>0.1</sub>MnO<sub>3</sub> manganite with mean particle sizes of 20 nm were prepared by the Microwave synthesis method. The crystal structure of the samples were analyzed, using X-ray diffractometer with Cu-ka (0.154 nm) radiation. The morphology of the samples was characterized by a scanning electron microscope (FESEM). Magnetic dynamics of the samples was studied by the measurement of ac magnetic susceptibility versus temperature and frequency. The XRD pattern along with Rietveld analysis indicates that the samples are nearly single phase. By fitting the experimental data with Vogel-Fulcher law, the relaxation time, characteristic temperature, magnetic anisotropy energy and effective magnetic anisotropy constant have been estimated.

#### P.J01

#### Magnetotransport properties of anisotropic Co(tCo)/Au(tAu) multilavers

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Artificially tailored Co/Au multilayers have been the subject of intense research in search of novel properties for potential applications [1] [2]. Earlier, we reported various components of giant magnetoresistance (GMR) effects and correlated them with microstructure of the isotropic multilayers [2]. In this work, we have investigated the effects of incident angle of deposition and magnetic annealing on GMR and magnetic properties of anisotropic Co/Au multilayers as the layer thicknesses of Co and Au are changed. We have investigated physical mechanisms responsible for various MR effects and correlated them with the existing models. The multilayers deposited at an oblique incident angle of 45 degree exhibited strong magnetic anisotropy. The maximum MR ratio is found to be 2.1 %. The transverse MR effect is always found to be larger than the longitudinal MR effect. These multilayers have been considered potential in developing biomagnetic technology.

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# PJ02

### The new type of current and spin polarization oscillations

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In our work we found new type of spin polarization and current oscillations in conducting rings with inhomogeneous magnetic properties. In case of hydrodynamic electron transport (when frequency of electron-electron collisions is more than frequency of collisions that does not conserve momentum of electrons) we obtained spin hydrodynamics equations for non equilibrium electron spins when the electron spectrum was spatially inhomogeneous. We showed that in case of close conductor these equations had non zero solution in form of "spin pendulum" oscillations. i e oscillations of full current and spin polarization with frequency determined by characteristics of magnetic inhomogeneity. We found own oscillations of electron system in conducting ring consisting of two parts with different magnetic properties. For example it could be realized by connecting DMS and NMS conductors into a ring and placing it in magnetic field. We used equations similar to well-known two-current model and we showed that conductance of this system as a function of external EMF frequency had one maximum in hydrodynamic regime and this case was conform to 'spin pendulum', and conductance had many maximums in a case of ballistic regime. These maximums vanished when magnetic properties of parts of the ring were identical

#### PJ03

#### Pressure-enhanced giant magnetoresistance in Fe/Cr magnetic multilayers Gendo Oomi\*

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Pressure-enhanced Giant Magnetoresistance in Fe/Cr Magnetic Multilayers G Oomi1), S.Higashihara2), K. Suenaga2), K.Saito3), S.Mitani3) and K.Takanashi4) 1)Kurume Institute of Technology,Kurume,Fukuoka,Japan 2)Department of Physics,Kyushu University,Fukuoka,Japan,3)Institute of Materials Research,Tohoku University, Sendai, Japan The giant magnetoresistance(GMR) of Fe/Cr magnetic multilayers has been measured under high pressure at low temperatures. It is found that the GMR on the second peak samples is enhanced strongly by applying pressure:the magnitude of GMR becomes twice at 2 GPa compared with the value at ambient pressure. The saturation magnetic fields of these samples decreases with increasing pressure. However the GMR of the first peak samples decreases with increasing pressure and the saturationfields increases with pressure. These sharp differences in the behaviour against pressure will be discussed phenomenologically.

#### PJ04

#### Spin current-induced by sound wave Igor Ivanovich Lvapilin Institut Metal Physics RAS, Russia

In recent years, increased interest in studying the effects that arise in response to spin degrees of freedom of electrons, when the external perturbation acts on the kinetic degrees of freedom. As rule, the external perturbation (electric field) directly affects only the kinetic degrees of freedom of electrons and through spin-orbital interaction is transferred to the spin subsystem. There are other ways of influencing the system of conduction electrons, which are also reflected in the response of the spin degrees of freedom. So, in the ferromagnetic metals temperature gradient leads to the spin Seebeck effect. It is interesting to consider the mechanisms of interaction with external fields, in which the energy of the external field is transmitted simultaneously in both electronic subsystems (kinetic and spin). An example of such interaction is the interaction of conduction electrons with the field of sound waves. We studied the evolution of the electronic system for their interaction with the field of sound wave analysis of the conditions of occurrence and the the response of the spin subsystem in a constant magnetic field. It is shown that in this case the longitudinal sound wave propagation arises transverse spin current, which has a resonant character.

# PJ05

#### Enhancement of magnetocaloric width in La<sub>2/3</sub>Ca<sub>1/3</sub>MnO<sub>3</sub> compounds with remain the composition ratio between Mn3+ and Mn4+ Oing Ji and Xiaoshan Wu\*

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Large magnetic entropy in a relative narrow transition temperature in manganites has been reported. To explore the enhancement of transition temperature width with large magnetocaloric properties, we here designed the samples with the same composition ratio of Mn<sup>3+</sup> and Mn<sup>4+</sup>, and the same averaged radius at each lattice sites, which we previously used to find the roles of the local Jahn-Teller distortion on the magnetic properties in manganites. Here, Y and Sr are used together to replace for La and Ca in La<sub>22</sub>Ca<sub>22</sub>MnO<sub>2</sub> (i.e., ABO3) to remain the average radius and the average valance at A-site, Results show that the averaged unit cell lattice parameters are unchanged. The ferromagnetic-to-paramagnetic transition temperature Tc increases abruptly owing to the increase of the competition between  $\sigma^2$  and double-exchange interaction resulting from Mn-O-Mn bond angle. The disordered distribution of cations at A-site broaden the phase transition temperature range (~T~10K or more) with large magnetic entropy of  $\Delta$ SM=1.3 J/Kg K, which may be used as a promising candidate for magnetic refrigerant near room temperature.

### PJ06

#### Output voltage calculations in double barrier magnetic tunnel junctions with asymmetric voltage behavior

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We studied the asymmetric voltage behavior of the tunnel magnetoresistance (TMR) for single and double barrier magnetic tunnel junctions (S&DMTJ) in range of a quasi-classical tunneling model. Numerical calculations of the TMR-V curves, output voltages and I-V characteristics for the negative and positive applied voltages were carried out using MTJs with CoFeB/MgO interfaces. Asymmetry is explained by different values of the minority and majority Fermi wave vectors for the left and right sides of the tunnel barrier, which can arises due to different annealing regimes. Electron tunneling in DMTJs was simulated in two ways: (I) Coherent tunneling, where the DMTJ is modeled as one tunnel system and (II) consecutive tunneling, where the DMTJ is modeled by two SMTJ connected in series. We found that DMTJs, in range of model I, the output voltage peaks will be shifted into the region of low voltages. For the model II, we found that, in order to provide symmetric output voltage branches in DMTJ, the conditions of equal Fermi vectors ratios have to be fulfilled for the left and right magnetic layers while ratio in the middle layer can be different.

#### PJ07

#### Spin-current manipulation by domain wall motion in the non-local spin valve

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A spatially- and time-dependent magnetic spin texture exerts a spin-dependent force on conduction electrons[1]. In particular, spin-motive force generated by a domain wall motion has been theoretically investigated and experimentally observed by detection of electric voltage[2,3]. The spin-dependent field generated by spin-motive force is also expected to directly affect the non-equilibrium magnetization such as spin accumulation. In order to study the effect of the spin motive force on the spin accumulation in a lateral non-local spin-valve (NLSV) structure, we treat the NLSV device attached with a magnetic wire. The device consists of spin injector F1 and detector F2 strips bridged by an non-magnetic strip N and the magnetic wire is attached on N through the insulating layer in-between two magnetic stripes. A domain wall is introduced into F3 and driven by external magnetic fields. The spin-dependent field drives the pure spin current flowing in the N-channel. This effect can be detected as increase or decrease in the nonlocal spin-accumulation signals depending on the domain wall motion. We evaluated the spin-dependent field based on the LLG approach and calculate the spatial distribution of spin/charge current of this NLSV by applying a finite element method in three-dimensions

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# **PJ08**

# The study of microwave assisted magnetization reversal via spin pumping

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There has been a significant interest to switch high coercivity magnetic bits for the future magnetic media by lowering the coercivity during the writing process. MAMR is a popular method of reducing the coercivity. We present a novel method of studying MAMR via spin pumping and the inverse spin Hall effect (ISHE). In this method, a Pt layer was placed in contact with a ferromagnetic (Py) layer, and an rf signal was applied to the Py via a coplanar waveguide during loop measurements. Due to the presence of the rf power, magnetization dynamics in the Py layer pumps spin in the adjacent nonmagnetic Py layer, which creates a measurable dc voltage across the Pt due to the ISHE. The measured signal is of the order of several microvolts and thus does not affect the magnetization by any significant amount. It also effectively decouples the input rf input from the dc output. The coercivity can be easily detected by the field at which there is a change in the polarity of the measured voltage. The coercivity as a function of the applied rf power is studied.

#### P.J09

# Contributions of domain walls on large magnetoresistance effect in ultrathin TbFeCo wires

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Perpendicular anisotropy magnetic materials such as TbFeCo are expected to introduce large effectivity of spin-transfer torque and high domain wall (DW) motion speed as well as high magnetoresistance (MR) values. We investigated the size dependence of MR effect in TbFeCo wires. The 80-µm-length and different widths (1 to 80 µm) of Tb<sub>26</sub>Fe<sub>66.8</sub>Co<sub>72</sub> wire-patterns were designed by using electron beam lithography technique. The Tb<sub>20</sub>Fe<sub>66.8</sub>Co<sub>72</sub> films with different thicknesses of 6, 18 nm were grown by RF sputtering. It is found that MR values are strongly depend on injection current for all applied fields. Large MR values which up to 100% were observed at a low bias current and depend on DW structure inside the TbFeCo wires. On the other hand, the MR values rapidly reduced at large bias currents which induce magnon and phonon scattering effects. The introduction of DWs in a magnetic wire causes an increase in spin-dependent scattering inside the DWs and leads to an enhancement of resistivity expectedly introduce a new technique to obtain high and multi-level MR values aw la shigh potentials on future magnetic devices.

#### **PJ10**

# First-principles study of conductivity tensor in transition metals and alloys

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Spin-orbit interaction which is a relativistic correction to electron motion influences the transport properties, so the resistance anisotropy and Hall effect appear in conductivity tensor. In the present work, we calculate the conductivity tensor of transition metals and alloys in first-principles by the Fermi-surface term of Kubo-Streda formula. To calculate the conductivity tensor and electronic structure with spin-orbit interaction, we employ the tight-binding linear muffin-tin orbital method based on the local spin-density approximation [1]. The substitution disorder of alloys is treated using the coherent potential approximation. The obtained anomalous Hall conductivity, anisotropic magnetoresistance ratio, and also normal resistivity in ferromagnets such as Fe-Co and Fe-Ni as a function of the alloys composition are quantitatively consistent with the available experimental measurements in terms of their magnitude and sign. In particular we find that the Hall resistivities reveal the nontrivial behavior against the magnetization which can be varied by the modification of the composition ratio. Furthermore, the intrinsic and extrinsic contribution of the Hall conductivity is comparable in these alloys systems. Also we show the calculation results about the spin Hall conductivity of nonmagnetic compounds obtained in similar frameworks.

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#### PJ11

#### Magnetoresistance of CoFe/Pt nano-contacts Muftah Al-mahdawi\* and Masashi Sahashi

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Using an alumina-based Nano-Oxide Layer (AlOx NOL) between with many nanoholes between two ferromagnets [1] has gained more interest in recent years both in their fundamental physics of spin-scattering off the confined domain-walls and in their applicability for next-generation hard-disk-drive reading heads. Also, spin-torque driven oscillations of the confined domain-walls has been reported [2]. Following the report by [3] that Pt is magnetized at the interface with Co, we examined the insertion of AlO<sup>x</sup>NOL between the Co-rich Co<sub>90</sub>Fe<sub>10</sub> and Pt. Film stack was (in nm): Underlayer/Co90Fe10 3/AIOx NOL/Pt 2/cap. It was deposited using ion-beam sputtering and dc magnetron sputtering deposition. AlOx NOL was prepared using Ion-Assisted Oxidation (IAO) process[1]. IAO process forms direct nano-contacts (NC) in NOL between the CoFe and Pt (1~2 nm in diameter, with <1% areal density). One film was pin-annealed in a 10 kOe field at 270°C for 1.5 hours and the other was left as is Magnetization curve was measured using VSM at room-temperature and showed only the curve of CoFe. Current-In-Plane Tunneling (CIPT) at room-temperature showed that CPP RA product were around 0.24 Ωµm<sup>2</sup> for both films, and MR ratio was 3% (asdepo film) and 5% (pin-annealed film). Further discussion and data will be presented.

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#### PJ12

#### Inversion of magenetoresistance in La<sub>1-x</sub>Sr<sub>x</sub>MnO<sub>3</sub>/Nb-doped SrTiO<sub>3</sub>/ CoFe junctions

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Understanding of the spin transport thorough ferromagnet/semiconductor interfaces is a critical issue to realize novel spintronic devices such as spin transistors. Specially, spin transport in degenerated semiconductors, which becomes metallic by heavy doping, is important for avoiding conductivity mismatch problems. In this study, magnetic junctions using degenerated semiconductors, Nb-doped SrTiO<sub>3</sub> (Nb-STO), as intermediate layers were fabricated. XRD measurements indicate 15mol% Nb-STO and La<sub>0.5</sub>Sr<sub>0.9</sub>MnO<sub>3</sub> (LSMO) films were epitaxially grown on MgO substrates (MgO(001) [100]//Nb-STO(001)[100]/LSMO(001)[100]). In the hysteresis curves of the LSMO/Nb-STO/ CoFe at 4.2 K, two step magnetization correspond to the Hc of LSMO and CoFe were observed. The junctions using the trilayers showed magnetoresistance (MR) of ~5 % at 4.2 K. The sign and the magnitude of the MR were changed depending on the thickness of Nb-STO. Typically, positive MR of ~2% for 10 nm and negative MR of ~1% for 15 nm were observed. The origin of the MR is under investigation, however we consider these behaviours can be explained by regarding the trilayers as magnetic tunnel junctions with double tunnelling barriers (*FM/INMI/FM*), which show oscillation of MR [1]. The depletion layers (~2 nm) at the interface between Nb-STO and ferromagnets seem to be behaved as insulative barriers in the junctions.

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### PJ13

### Spin transfer torques in antiferromagnets

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Current-driven magnetization switching through spin transfer torque [1] has been widely studied in ferromagnetic (F) spin-valves [2,3]. Recently, Nunez et al. [4] theoretically showed that spin-valves based on antiferromagnetic (AF) electrodes also display outstanding spintronics properties such magnetoresistance and spin torque. Whereas the nature of current-driven torque has been theoretically examined in AF spin-valves [6,7], promising experimental results have been obtained on F/NF/AF spin-valves structures (F is a ferromagnet and N is a normal metal) [5]. However, the microscopic origin of the spin torque between F and AF layers has not been addressed. In the present study, we apply non-equilibrium Green's function formalism in the tight- binding approach to investigate the spin transport in spin-valves consisting in N/AF/N/AF/N and N/F/N/AF/N. Whereas the spin torque calculated in the AF spin-valve is consistent with previous results [4,6,7], it presents significant differences in the latter structure (angular dependence, spatial distribution and relative magnitude between the in-plane and out-of-plane torque). Interestingly, we find that the AF layer servers a torque on the F layer. Finally, we analyze the contribution of the implane torque and out-of-plane torques on the different sub-lattices of the AF layer and discuss the implications in terms of magnetization dynamics.

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# PJ14

# Domain wall configuration and magneto-transport properties in dual spin-valve with nanoconstriction

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In this work, we use a synthetic antiferromagnet-based dual spin-valve (SAF-DSV) structure (i.e. the SV structure is doubled symmetrically with respect to the FM) and present the effect of the direction of the applied magnetic field with respect to nanoconstriction shapes on the magneto-transport properties including domain wall (DW) configuration and reversal process. We can tune the DW configuration and its reversal process from a vortex to a transverse type by changing the direction of applied magnetic field respect to the nanoconstricted SAF-DSV, the perpendicular magnetic moments are developed due to the transverse magnetization reversal mode and result in a multistep switching process. This multi step switching process reflects the pinning and depinning of a DW at the nanoconstriction. Our results also show an asymmetric depinning field. We demonstrate, if nanoconstriction is asymmetric along its length, i.e. expansions from both sides of the neck into the two nanowires are not identical, and then an asymmetric energy barrier to domain wall propagation is formed. This is due to the difference in DW width, which leads to an asymmetry in the domain wall depinning forces.

### PJ15

Giant magnetocaloric effect of Mn<sub>0.91</sub>Ca<sub>0.09</sub>As thin film on Al<sub>2</sub>O<sub>3</sub> (0001) Duong Anh Tuan<sup>1</sup>, Cho Sunglae<sup>1</sup>\*, Dang Duc Dung<sup>2</sup>, Shin Yooleemi<sup>1</sup> and Jeon Seung Mok<sup>1</sup>

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The MnAs compound is promising to magnetocaloric effect (MCE) application [1]. It presents the first order phase transition from a ferromagnetic to paramagnetic ~ 318 K, which is accompanied by a structural transition from a hexagonal NiAs-type (a/MnAs) to orthorhombic MnP-type structure (B/MnAs) [2]. Recently, the room temperature giant magnetocaloric effect was obtained by substitution of Mn by Fe, Co, Cu or Cr in MnAs: Co-25 J/kgK at 4T, Cu-175 J/kgK at 5T, Cr and Fe-42 J/kgK at 5T [3-5]. However, there is less information about group II dopants such as Ca, Ba, and Sr in MnAs. In this work, we report the magnetism and transport properties of epitaxial Ca<sub>000</sub>Mn<sub>0.97</sub>As thin film on Al<sub>2</sub>O<sub>3</sub>(001) substrate. The RHEED and XRD results indicated that film were epitaxially grown on Al<sub>2</sub>O<sub>3</sub>(001). The temperature dependent resistivity showed metallic and the temperature dependent magnetoresistance at 7 kOe showed a peak around 325K. The negative anomalous Hall effect was observed in Hall measurement. The magnetic measurement indicates that the sample is ferromagnetic with TC was around 340K and display the giant magnetocaloric signal. The magnetization hysteresis loops exhibited the saturation magnetization (MS) was around 300 (emu/cm3) at 10K and 180 (emu/cm3) at 300 K.

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#### PJ16

#### Anomalous Hall effect of [Amorphous CoSiB/Pt] multilayer films Hana Lee, Insung Park, Hyungjun Kim, Sungyong Kim, Youngkwang Kim, Hwayong

Noh and Taewan Kim\* Seiong university, Korea

We have quantitatively investigated the Anomalous Hall effect (AHE) in [CoSiB/Pt] multilayer films. The [CoSiB/Pt] multilayers exhibited large Hall effect. In this study, we compared with the effects of the before and after of patterning. The [amorphous CoSiB 6A/Pt 14 A]n multilayers were deposited on SiOx substrate with a 30 A Pt and 50 A Ta buffer layer at the room temperature using a DC magnetron sputtering system. Thereafter the sample was patterned into a hall bar with the line width of 10  $\mu$ m. Hall effect measurements at room temperature were made up to 3 kOe with the applied magnetic field perpendicular to the film plane. The samples for the Hall effect measurement were prepared with a square geometry for the easy application of the Van der Pauw method to measure the Hall voltage was in the range of several mV.

# PJ17

# Spin transport phenomena through MgO/CuPc hybrid barrier in magnetic tunnel junctions

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For realistic applications of organic spintronic devices, the investigation of the growth mechanism and electrical properties is of great importance to understand the spin transport phenomena in metal-organic hybrid systems. As the first step, we focused on surface morphology, growth mode, and thermal stability of thin Cu-phthalocyanine (CuPc) organic layers grown on metal surfaces covered with a few monolayer thick insulating films, such as MgO(001). Recent reports of achieving spin polarized tunneling across organic-inorganic hybrid barrier with significant magnetoresistance (MR) values at low temperature (< 10 K) led us to study the organics semiconductor (OSC) thin films for substituting highly performing inorganic barriers. Therefore, we investigated systematically the spin transport phenomena through the OSC CuPc barrier with and without MgO(001) layer as a function of CuPc thickness. An epitaxial Fe(001)/MgO(001) layer was utilized as a spin injection layer, while a polycrystalline Co film as a spin detector. Here, we observed, for the MgO/CuPc hybrid barrier with a total thickness of  $\sim 3$  nm, the high MR (> 200 %) values at 77 K and  $\sim 10\%$  MR at 300 K in the Si(001)/MgO(001)/CuPc/Co hybrid tunnel junctions.

### PJ18

#### Theoretical approach to spin-current absorption at an interface

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In the development of spintronic devices, techniques to detect spin currents efficiently are indispensable as well as the generation of spin currents. In general, the spin-current detection is demonstrated via the inverse spin Hall effect and the absorption of the spin current [1]. In this presentation, we theoretically examine the spin current induced by an inhomogeneous spin-orbit coupling due to impurities [2]. Using the Keldysh Green's function formalism, we propose analytical expressions of the diffusive spin currents under an external electric field and find a spin-current generation in the presence of both gradient of the spin-orbit strength and that of the external field. The resulting diffusive spin current indicates the absorption of the spin current at the interface between materials with different spin-orbit strengths, which is exemplified with an experimentally relevant setup.

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 K. Hosono, A. Yamaguchi, Y. Nozaki, and G. Tatara, Phys. Rev. B 83, 144428 (2011).

### PJ19

#### Magneto-transport properties of Al<sub>2</sub>O<sub>3</sub>-doped Mn-Zn ferrites

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We report the magneto-transport properties of Al<sub>2</sub>O<sub>2</sub>-doped Mn-Zn ferrites prepared by the conventional solid state reaction. Various amounts of Al<sub>2</sub>O<sub>2</sub> were doned to a pure Mn-Zn ferrite to have the nominal composition of  $(Mn_{0.8}, Zn_{0.2})_{1-x}Al_xFe_2O_4$  $(x = 0.03 \sim 0.1)$ . Field dependency of magnetization was measured with a SOUID (superconducting quantum interference device) magnetometer, and low field magnetoresistance (LFMR) was also measured with the SQUID magnetometer using an external current source and voltmeter. The X-ray diffraction analyses revealed that all Al<sub>2</sub>O<sub>2</sub>-doped Mn-Zn ferrites of the spinel structure were a pure phase without the second phase. From the transport measurements of our samples in the temperature region of 100~300 K it was found that with the addition of the Al<sub>2</sub>O<sub>2</sub> dopant the resistivity level of the pure Mn-Zn ferrite was abruptly dropped over four orders of magnitude and further the LFMR ratio was greatly improved. High LFMR ratio over 1.7% at 300 K in 0.5 kOe was achievable from 3mol% Al<sub>2</sub>O<sub>2</sub>-doped Mn-Zn ferrite without an appreciable increase in its resistivity. Detailed effects of the Al<sub>2</sub>O<sub>2</sub> dopant on the microstructures, magnetic and magnetotransport properties of the Mn-Zn ferrite will be presented for a discussion

# **PJ20**

Current induced fluctuation of switching fields in Co/Pd nanowires Mahdi Jamali, Xuepeng Qiu, Kulothungasagaran Narayanapillai and Hyunsoo Yang\* Electrical and Computer Engineering, National University of Singapore, Singapore

Current induced effective magnetic fields in a thin metallic film of a ferromagnetic material sandwiched by two nonmagnetic materials have been studied intensively due to its potential applications for the magnetic random access memory [1, 2]. There is a controversy about the origin of the effective field such as the Rashba effect [3] or spin torque induced by the spin Hall effect [4]. We have recently observed current induced effective fields in nanowires made of Co/Pd multilayer [5]. The further investigation of the switching fields of nanowires reveals fluctuation in the switching fields upon an injection of different current densities. We have applied current pulses with a width of 2 ms and repetition of less than 1 Hz to minimize the Joule heating effect. At low current densities (J  $\leq 10^{11}$  A/m<sup>2</sup>) there is a fluctuation in the switching field of the nanowire. For a certain value of the input current, the switching field of the Co/Pd nanowire increases(>6%). Furthermore, the switching fields in the positive and negative fields are not the same. The relative angle between the sample and magnetic field has changed and the effect of the angle on this fluctuation has been studied.

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# PJ21

#### Role of structural inversion asymmetry on current-induced effective field in perpendicular magnetized trilayers

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The fact that a current pulse can switch the magnetic moment of a perpendicular magnetized nanodot has attracted much attention for the next new generation spintronic devices [1]. However, the physical origin of this effect is still under debate. Here we report the pulse delta measurements in perpendicular magnetized substrate/2 MgO<sub>2</sub> Pt<sub>0.8</sub> CoFeB<sub>2</sub> MgO (sample A) and substrate/2 MgO<sub>08</sub> CoFeB<sub>2</sub> Pt (in nm) micro-sized hall bar samples with opposite structural inversion asymmetry (SIA). For both structures, when the magnetic field applied along the current direction tilted off-plane by 2°, positive (negative) current induces a downward (upward) effective field with positive magnetic fields and the effective field direction switches with negative magnetic fields. The effective field is around 2,388 Oe per 1012 A/m2 in the sample A. For the other structure, the effective field shows a nonlinear correlation with the intensity of current. With small in-plane external fields, the magnetization can be switched by pulse current and the switching hysteresis was the same in the two structures. Our results suggest that the main contribution of the current induced effective field arises from the spin Hall effect rather than the Rashba or other SIA related effects

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### PJ22

#### Characterization of mechanically milled Cu-Co powder by 3D-FIB and atom probe tomography : effect of oxidation on the magnetoresistance

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The Giant MagnetoResistance (GMR) effect discovered in 1988 [1] had inspired many researches on magnetoresistive materials. The resistance of the device is reduced with the application of a magnetic field and it is due to the spin dependent scattering effect. However, an inverse effect has been observed after an oxide phase introduction by Khan et al. [2]. This effect appears at low magnetic field and could be interesting for improving the sensitivity of magnetoresitive devices. In our study, a Cu<sub>80</sub>Co<sub>20</sub> powder was elaborated by mechanical alloying under ambient atmosphere to investigate the oxide formation and its influence on magnetoresistive properties. 3D-Focused Ion Beam tomography (3D-FIB) and Atom Probe Tomography (ATP) [3] were used to characterize our samples. It was shown with 3D-FIB that after 1h30 of milling, the powder exhibits a lamellar structure with micro-sized Co rich regions. The APT experiments have revealed that after 20h of milling 1/3 of Co atoms is oxidized leading to the formation of nano-sized CoO clusters embedded in a Cu<sub>85</sub>Co15 matrix. These results bring crucial informations to correlate the structural properties with magnetoresistive properties of the Cu-Co-O powder and to understand the influence of the nano-sized oxide clusters on the inverse magnetoresista

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# **PJ23**

#### Effect of pressure on magnetotransport properties in Fe/MgO granular films

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A special example of granular thin films is the discontinuous metallic and insulator multilavers (DMIMs) consisting of metallic layers with different degrees of discontinuity intercalated between insulating layers. However the studies about the effect on the magneto-transport properties of the hydrostatic pressure are scarce[1] In this work we report the effect of hydrostatic pressure on the resistance and magnetoresistance of a DMIM of nominal composition [Fe(t=0.7nm)/MgO(t=3 nm)]15 deposited on a substrate of coming glass. The DMIM has been characterized previously showing a superparamagnetic behaviour [2]. The resistivity has been measured at room temperatures with hydrostatic pressure (Ph) up to 7 Kbar. The resistivity decreases linearly with pressure indicating an increase in conduction via tunneling. The value of the coefficient (1/o0)do/dPh =-3.9 10-2 kbar-1 is higher than the value obtained in other granular films. This result implies that the electronic state of the DMIM is near the percolation threshold. Magnetoresistance (MR) isotherms have been measured at room temperatures for Ph=1bar and Ph=7Kbar. At the maximum applied magnetic field (3KOe) the MR increases from 0.6% at ambient pressure to 1.1% at 7 Kbar. This result will be explained due to the reduction of the tunnel barrier induced by the hydrostatic pressure.

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#### PJ24

#### Measurement of anomalous nernst effect in [CoSiB/Pt] multilayer films

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Recently, a new research field taking advantage of both spintronics and thermoelectronics, called spin caloritronics, has started to attract a great deal of attention [1]. In contrast to the Anomalous Hall Effect driven by electrical currents, Anomalous Nernst Effect includes combination of thermal and spin transport driven by a temperature gradient in ferromagnets. We have investigated Anomalous Nernst Effect in [CoSiB/Pt] multilayer films with perpendicular magnetic anisotropy. The layer structure of samples was SiO<sub>2</sub>Ta<sub>50A</sub>Pt<sub>30A</sub>(CoSiB<sub>6A</sub>Pt<sub>14A</sub>)n. Standard Hall bars were patterned on the samples. Additionally, heating element was also patterned on the samples in order to apply temperature gradient. Nernst effect measurements were made up to 1.5 kGauss with the applied magnetic field perpendicular to the film plane. We obtained significant Nernst voltages (VN) by applying temperature gradient  $\bigtriangledown T$ to the samples. Nernst voltage exhibited similar characteristic as AHE voltage due to the spontaneous magnetization. Samples with different number of [CoSiB/Pt] multilayers were measured and compared. VN increased with increasing number of layers. Dependence of VN on the heating power (P) and the distance from the heating element (d) have been investigated. We have extracted [VN] a [P] and [VN] a [1/d] dependencies from the measured samples.

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# PJ25

#### GMR effect in Co-Cu microwires

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Granular materials exhibiting giant magnetoresistance (GMR) effect attracted considerable attention. For obtaining of the granular structures, melt spinning technique allowing rapid quenching from the melt is quite suitable. During last years considerable attention has been paid to the studies of glass-coated microwires produced by Taylor-Ulitovski technique allowing to achieve high quenching rate and producing amorphous, nanocrystalline, microcrystalline or even granular thin microwires. In this paper, we studied magnetic properties and GMR effect of  $Co^{x}Cu_{wa}$  (10<x<30 at%) microwires and observed considerable magnetoresistance (MR) AR/R For Co., Cu90 microwires AR/R increases when the temperature decreases, exhibiting negative MR, tending to saturate in high magnetic fields, H. On the other hand, for Co<sub>10</sub>Cu<sub>70</sub> samples AR/R(H) dependences showed non-monotonic behavior exhibiting AR/R increase with H at low H-values (up to 5 kOe) and also considerable GMR effect. X-ray diffraction (XRD) results reveal that the structure of the metallic core is granular with two phases: the main one, fcc Cu (lattice parameter 3.61 Å) and fcc α-Co (lattice parameter 3.54 Å) which is present in microwires with higher Co content. In the case of low Co content, XRD indicates that Co atoms are distributed within the Cu crystals.

# PJ26

#### Domain wall quantum interferometer John Eves, N. Grisewood and H. B. Braun School of Physics, UCD, Ireland

Ferromagnetic nanosized toroidal rings have recently attracted much interest as potential storage elements for conventional bit storage. Here we consider quantum transport properties of such rings, and their dependence on the intrinsic domain structure of the ferromagnetic materials. Considering samples with realistic anisotropies we study the propagation of conduction electrons in the background of nontrivial domain configurations. Due to sd-coupling, the propagating electrons experience a domain dependent phase shift due to Berry's phase. Using exact solutions for domain configurations on a ring we use a path integral approach to determine exactly how this phase depends on the domain configuration in the ring. We find that the transmission through the ring can be controlled by the relative position of the domain walls, giving rise to a novel domain wall interferometer.

#### PJ27

#### Synthesis and magnetic properties of trilayer NiFe/Bi/NiFe films

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Earlier, it was found [1] that the bismuth spacer formation influences essentially the system magnetization. The shape of m(H) curve is changed with bismuth spacer thickness increasing. In particular, the test film with tBi = 0 has the narrow hysteresis loop and magnetization curve of a ferromagnetic type. For the films with tBi  $\neq$  0 the magnetization curves are typical for films possessing either by <strong> intra-layer anisotropy or by antiferromagnetic interlayer coupling. Since anisotropy is not experimentally observed, we attribute such a behavior to the presence of interlayer antiferromagnetic exchange. The investigations of electron magnetic resonance were made in NiFe/Bi/NiFe films. It was established that the magnetic resonance spectrum consists of solitary line for films with tBi = 0 nm and tBi = 15 nm, whereas for films with intermediate value of semi metal spacer the spectrum is superposition of two lines. Theoretical treatment of temperature dependences of resonance fields gives conclusion that the interlayer coupling depends on both bismuth interlayer thickness and temperature. Also, the magnetoresistivity of order of percent units and dependence of its value on thickness of bismuth spacer were found in these films at helium temperatures [2].

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### P.J28

#### A study on the pulesd laser deposited metallic spin valve structures Savak Ghoshal1\* and P S Anil Kumar

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It is envisaged that the conventional Spin Valve (SV) structures will be more advantageous over Tunneling Magneto-resistance (TMR) for the perpendicular magnetic recording due to the suitable RA product. Here Pulsed Laser Deposited (PLD) Spin Valve (SV) and Pseudo Spin Valve (PSV) samples are grown at room temperature with moderately high MR values using simple NiFe/Au/Co/FeMn structure. Although PLD is not a popular technique to grow metallic SVs because of its expected large intermixing of the interfaces, particulate formation etc., still by suitably adjusting the deposition parameters we could get exchange bias as well as more than 3% MR in the Current In Plane (CIP) geometry. Different SV and PSV samples are grown to study anisotropy effects, exchange coupling and exchange bias with anti-ferromagnet. Angular variation of the MR is also obtained which shows that the hard layer (Co) has a four-fold anisotropy in the NiFe/Co/Au/Co structures which becomes two-fold in presence of anti-ferromagnet. In this presentation, properties of these PLD grown SVs will be discussed in terms of the angular variation of MR, intermixing at the interfaces as well as details of exchange bias.

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### P.J29

#### Spin torque in a finite two-dimensional ferromagnet with Rashba interaction

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The interaction between a ferromagnetic exchange coupling and a Rashba spin-orbit interaction has shown its potential in magnetization switching [1, 2]. For such a system, a coherent theoretical framework has been developed to study the coupled spincharge diffusive dynamics [3]. Using finite element method, we numerically solve, in a two-dimensional (2D) electron system with appropriate boundary conditions, the generalized equations for the diffusive charge and spin dynamics. This numerical approach allows for an accurate description of the spin and charge dynamics (i.e. nonequilibrium accumulation and current) for a wide range of parameters beyond the limit of weak Rashba interaction [1,3]. The spatial profile of the spin accumulation exhibits a strong presence of a spin precession arising from both the s-d exchange and the effective Rashba field. Spin relaxations due to random magnetic impurities as well as the D'yakonov-Perel mechanism are embedded in the solution. We also discuss the diffusive dynamics and the spin torques in different regimes with respect to the magnitude of the Rashba spin-orbit interaction.

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### PJ30

#### Spin transport of Py/Au/Py spin valves with different Au channel widths

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Most previous studies on spin injection and transport were dealt with one dimensional diffusive model and two-dimensional consideration has been neglected for simplicity [1-2]. However, all of the lateral spin valve devices experimentally demonstrated are fabricated based on thin film technology so that it is indispensable to have spatial distribution of spin current in a non-magnetic channel [1-2]. Actually, some of the experimental data can't be understood in terms of one-dimensional conduction model. Injected spin current usually produce inhomogeneous spin accumulation which appears to a different magnitude of spin signal ( $\Delta R$ ) depending on the position of voltage probes on the detector [3-4]. In the study, nonlocal spin valve (NLSV) measurements for lateral Permalloy(Py)/Au/Py spin valve devices with different Au channel widths were carried out in two viable voltage probe configurations. The larger channel width is at a fixed center to center channel length, the lower the magnitude of  $\Delta R$  is in the NLSV, which is attributed to the increase in junction area and enhanced spatial distribution of spin accumulation. The voltage probe configuration strongly affects the spin valve signal owing to the non-uniform spin current distribution in lateral Au channel.

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# PJ31

#### The effect of magnetic impurities in magnetic tunnel junction

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Magnetic tunnel junctions (MTJs) have been extensively investigated for more than a decade, not only because of potential applications such as next generation memory devices or high frequency spin oscillators, but also due to scientific interest. The physical properties of MTJs are very sensitive to the conditions of interface formation and, therefore, it is essential to understand the various influencing factors to control interface quality. The effect of oxidizing the ferromagnetic layer on tunneling magnetoresistance (TMR) is one of the most interesting problems in MTJs. The ferromagnetic layer in MTJs was oxidized with varying O2 concentrations and the corresponding effect on spin-dependent transport was studied. As expected from our previous results for MTJs with an over-oxidized AlOx tunnel barrier, a partially oxidized ferromagnetic layer plays an important role in spin-dependent transport. As the temperature is lowered, the junction resistance increases dramatically and the TMR is strongly suppressed. Increasing the O2 concentration enhances the increase of resistance and suppression of TMR. This work supports our previous conclusion that oxidizing the ferromagnetic layer generates localized magnetic moments, which act as a scattering center for spin-polarized tunneling electrons.

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# PJ32

# Effect of cobalt layer thickness on the magnetic and magnetoresistance properties of asymmetrical Co/Cu multilayers

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A series of asymmetrical multilayer thin films of Co/Cu with a systematically varying cobalt layer thickness tCo= 10, 15, 20, 30, 40 and 50 A was prepared by DC magnetron sputtering technique to study the effect of tCo on its magnetization and magnetoresistance properties. The XRD pattern of the multilayers showed the polycrystalline nature with texture (111). The magnetization measurements at different temperatures of the series showed that change in coercive field (HC). The magnetoresistance measurements at different temperature of the series of multilayer showed a temperature dependent giant magnetoresistance (GMR) due to the inelastic scattering of electrons by phonons and magnons. Moreover, an asymmetry in GMR curves at low temperature was observed because of the asymmetrical (pseudo spin valve) structure of the multilayer.

#### PJ33

# Current-induced motion of a transverse magnetic domain wall in the presence of spin Hall effect

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The electric manipulation of a domain wall (DW) in a magnetic nanowire is subject to the spin-transfer torque (STT) effect due to the coupling between local magnetic moments of the DW and spin-polarized currents [1]. Numerous studies on this subject have been performed to understand its fundamental physics [2] and to extend its potential to applications such as data storage and logic devices [3]. In this work, we study theoretically current-induced dynamics of a transverse magnetic domain wall in bi-layer nanowires consisting of a ferromagnet on top of a nonmagnet having strong spin-orbit coupling. Domain wall dynamics is characterized by two threshold current densities, J(WB) and J(REV), where J(WB) is a threshold for the chirality switching of the domain wall and J(REV) is another threshold for the reversed domain wall motion caused by spin Hall effect. The domain wall with a certain chirality may move opposite to the electron-flow direction with high speed in the current range J(WB) < J < J(REV) for the system designed to satisfy the conditions J(WB) > J(REV) and  $\alpha > \beta$  where  $\alpha$  is the Gilbert damping constant and  $\beta$  is the nonadiabaticity of spin torque. Micromagnetic simulations confirm the validity of analytical results.

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### PJ34

#### Universal spin-hall effect in metallic thin films

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Spin-orbit interaction fills the academic and industrial pursuit for an electrical manipulation of spins. Fabricated using the state-of-art nanotechnology, semiconductor and metallic thin films embedded between two asymmetric interfaces in the presence of spin-orbit interaction have attracted intensive theoretical and experimental interest in recent years. Spin Hall effect [1], voltage-controlled spin precession [2] as well as the recently observed spin-orbit torque [3] all stem from a careful design of spin-orbit interactions in an ultrathin normal metal film that is sandwiched by two insulators. While providing a confinement, the insulator/metal interfaces also serve as a source to the spin-orbit interaction due to inversion symmetry breaking. More interestingly, the interfacial spin-orbit interaction and is predicted to be much larger than that induced by bulk impurities.

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# PK01

#### Controlling magnetic isolation and moment reversal of FePt (001) films by Cu capping nanoislands

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We aimed to control the magnetic interaction of FePt(001) films by introducing Cu capping nanoislands for applications in magnetic recording and permanent magnets. Fe/Pt multilayers were prepared by e-beam evaporation onto MgO substrates, and the thickness of Cu capping layers was varied from 1 to 6 nm. Out-of-plane coercivity increases with increasing growth temperature, and the reduction of grain/domain size and intergrain interaction of FePt can be achieved. Cu could penetrate into FePt nanostructures along the grain boundaries and create a strain-energy modulation at the interface due to its lower surface energy. Studies of angular dependent coercivity show a tendency of a domain-wall motion shift toward rotation of reverse-domain type with Cu capping nanoislands. The intergrain interaction was confirmed from the Kelly-Henkel plot that indicated the strong exchange coupling between neighboring grains in pure FePt films, and the negative  $\delta M$  value was obtained while FePt films capped with Cu nanoislands

#### PK02

# Perpendicular magnetic anisotropy for annealed Co/Ir(111) ultrathin films

Wen-yuan Chan, Du-cheng Tsai, Cheng-hsun Chang and Jyh-shen Tsay\* National Taiwan Normal University, Taiwan

Compositions and magnetic properties of Co/Ir(111) ultrathin films thinner than 4 monolayers (ML) have been investigated. As the Co thickness increases to above 2 ML, magnetic hysteresis appears in both the longitudinal and polar configurations as revealed using surface magneto-optic Kerr effect (SMOKE) technique. After annealing treatments, the observations of the attenuation of Co LMM Auger signal and chemical shifts to lower kinetic energy show that Co atoms diffuse into the Ir(111) substrate. At higher temperatures, the Kerr intensities in both the longitudinal and polar configurations decrease due to the reduction of Co overlayer. At a temperature interval, magnetic hysteresis appears only in the polar configuration. This shows that the magnetic easy axis of Co/Ir(111) could be stabilized in the direction of the surface normal by thermal annealing treatments. By systematical investigations for Co/Ir(111) ultrathin films thinner than 4 ML, a magnetic phase diagram has been established. Several phases are observed, including a nonmagnetic phase, canted out-of-plane anisotropy and perpendicular magnetic anisotropy. On the top of the Co/Ir(111), the introduction of Ag overlayer shows an oscillatory behavior of the Kerr intensities. Possible mechanisms of the charge transferring, quantum well effects and completeness of the Ag layer are discussed.

#### PK03

# Anomalous easy-plane magnetocrystalline anisotropy of compressive strained (La,Ba)MnO, films

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Magnetic anisotropy of manganite films is sensitive to the epitaxial strain. So far, in Caand Sr-doped manganite films, many authors reported that a lattice expansion induces easy magnetization in the corresponding direction. In this work, magnetocrystalline anisotropy of La067Ba033MnO3 epitaxial films suffering large compressive strain on LaAlO<sub>3</sub> substrates (LBMO/LAO) has been investigated by using ferromagnetic resonance (FMR) technique. By mapping out the dependence of the FMR position on the angle between the applied magnetic field and crystallographic axes of the films, a large easy-plane anisotropy has been found in such compressive strained LBMO films, which is in contrast to the common tendency reported before. The excellent epitaxial crystallographic orientation gives rise to an observable fourfold symmetry of the in-plane anisotropy. The results are discussed by taking large tolerance factor and giant magnetostriction of Ba-doped manganites into account. In the phase diagram of temperature vs. tolerance factor for doped manganites, we noticed that LBMO has a larger tolerance factor but a lower Curie temperature than Sr-doped manganite, meaning that shrinkage of Mn-O-Mn bond length in LBMO is in favor of spontaneous magnetization, which in turn results in the anomalous easy-plane anisotropy in the compressive strained LBMO films.

# PK04

# Magnetic reversal of Co/Pt multilayer depending on Co thickness and annealing temperature

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Magnetic reversal behaviors in ferromagnetic thin films are continuously studied to achieve high performance of technological applications. Recently, CoPt multilayers and alloys have been an important issue in perpendicular magnetic anisotropy (PMA) applications. However, fundamental studies of the magnetic properties were not much presented. In this work, we present a magnetic reversal behavior of Ta(3.0 nm)/ [Co(x nm)/Pt(1.0 nm)]x3/Ta(1 nm) multilayers, when the Co thickness varied from 0.3 nm to 2.0 nm. The multilayers were deposited using DC magnetron sputter on Si/ SiO2 substrate, and annealed with various temperature after deposition. The magnetic and structural properties of the multilayers were measured using vibrating sample magnetometer (VSM) and x-ray diffraction (XRD) at room temperature. In addition, we observed the magnetic reversal behavior with magneto-optic Kerr effect (MOKE) microscope system, capable time-resolved domain observation.

#### PK05

#### Annealing effect on the microstructures and magnetic properties of [Fe/ Pt]16 multilayers on MgO (001) substrate

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[Fe/Pt]16 multilayers are deposited on MgO (001) single crystal substrate by magnetic sputtering layer by layer with varying the substrate temperature from 150-500oC to explore the annealing effect on the microstructures and magnetic properties of the films. Results show that FePt alloy with L<sub>10</sub> phase forms at about 300oC, and reaches a completely ordered state at about 400 oC. FePt film depositing at above 300oC has (001)-preferred texture which shows that L<sub>10</sub>-FePt film has perpendicular magnetic anisotropy. The out-of-plane coercivity increases with sample transforming from [Fe/Pt]16 multilayers to L<sub>10</sub>-FePt alloy. The out-of-plane coercivity, which is confirmed to have the effects on the out-of-phane coercivity, with the continuing surface has an order higher than that for the granular surface.

# PK06

#### Preparation and magnetic studies of room-temperature sputtered [Co/ Pt] multilayer films on glass substrates with perpendicular magnetic anisotropy

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Multilayer  $[Co_{tul}/Pt_7_3]n$  films were deposited at room temperature (RT) on glass substrates with and without a Pt(111) underlayer. The thicknesses of each Co and Pt layer were set at 1 and 0.75 nm respectively. The stacking period, n, was varied from 1 to 6, leading to the total thickness of [Co10/Pt7.5]n films from 1.75 to 10.5 nm. The results of magnetic studies indicate that the perpendicular magnetic anisotropy (PMA) of  $[Co_{tu}/Pt_{7.3}]n$  films is strongly influenced by inserting a Pt(111) underlayer and the number of n. Surprisingly, [CoPt]n films exhibit PMA behavior with a Pt(111) underlayer when n exceeds 3 even the films were deposited at room temperature. However without a Pt(111) underlayer, the films will not exhibit PMA in all cases. The appearance of L11 phase is found related to the appearance of PMA in  $[Co_{tul}/Pt_{7.3}]nPt(111)$  films. The [CoPt]n film is grown along Pt(111) resulting in the L11(222) orientation, which causes the easy axis <111> along the plane normal. In this study, the microstructure at  $[Co_{tul}/Pt_{7.3}]nPt(111)$  interface is also studied and will be presented, which is helpful for understanding the formation mechanism of PMA behavior in multilayer [Co<sub>ut</sub>/Pt<sub>7.3</sub>]n structure prepared at room temperature.

# PK07

# Magnetic anisotropy in FeCo thin films

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Sputtered FeCo films normally show magnetic anisotropy inside the film plane. It is believed that the magnetic anisotropy are due to the stray field of the magnet of the target. However, the details of the induced magnetic anisotropy are yet to be fully understood. In the present work, we have prepared FeCo films by both DC magnetron sputtering system and facing targets sputtering system. The magnetic properties, crystallographic properties, and microstructure of the films have been investigated. Films deposited by DC magnetron sputtering show isotropic magnetic properties in the film plane. However, films deposited by facing targets sputtering show strong magnetic anisotropy in the film plane. The magnetic anisotropy field increases from 40 Oe to 540 Oe with the increase of the film thickness from 10 nm to 150 nm for films deposited by facing targets sputtering system. XRD results show that the (110) lattice spacing of FeCo films increases with the increase of the film thickness. Of particular interest is thatilnternal stress of the films increased with the decrease of the film thickness.

#### **PK08**

#### Magnetic domain wall energy of Co/Ni superlattice with perpendicular magnetic anisotropy

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Co/Ni superlattice is a promising material for spintronics devices, in which the magnetic domain walls carry information stored, owing to a perpendicular magnetic anisotropy and a high saturation magnetization. Despite the numerous studies on the magnetic domain wall motion of the Co/Ni superlattice, the magnetic domain wall energy density has not been reported. In this work, we present the experimental determination of the magnetic domain wall energy density of the Co/Ni superlattice with perpendicular magnetic anisotropy (PMA). We confirmed that the epitaxial growth of Co(111)/Ni(111) superlattice on the Ag(111)-buffered Si(111) substrate. Several parameters were determined to obtain the magnetic domain wall energy  $\sigma W$  [1], i.e. the saturation magnetiz domain period on the ferromagnetic layer thickness, the dipole length, D0 ( $\equiv \sigma W/2\pi MS2$ ), is estimated as 22.5±2.5 nm. Using the experimentally obtained MS ( $\approx 660 \text{ emu/cc}$ ), the magnetic domain wall energy density is estimated as  $\sigma W = 6.4\pm 1.4$  erg/cm2. The estimated  $\sigma W$  is about three-times higher than those for L12-CrPt3[2], and about one-third of that for L10-FePt[3]. From the obtained values, the exchange stiffness constant is calculated as about 4.8×10-6 erg/cm, which is comparable to that of bulk-Co.

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# PK09

#### Anisotropic Magnetoresistance of Co/Ni multilayers

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Co/Ni multilayers exhibit perpendicular magnetic anisotropy, relatively low magnetic damping and unusual magneto-transport characteristic that make them useful for spintronic devices. Anisotropic magnetoresistance was measured for sputtered Ni/ Co films: SiO2Ta40APd40A(CO2ANiY)x10Ta40A, with Ni thicknesses from Y=2-14Å. Three AMR configurations were studied. Two configurations varying the applied field vs. current angle, the applied field always in the plane of the sample, AMR1, and varying the applied field from in plane to out of plane, AMR2, provide similar information. The %AMR1 or 2 increases monotonically as the fraction of Ni increases and shows the expected  $\cos\theta^2$  dependence. Rotating the sample with the field always perpendicular to the current, AMR3, we see anomalous angle-dependence for all samples, at all temperatures. At the sample orientation where a peak in the AMR is expected, there is a dip. The relative magnitudes of these contributions vary with temperature and Ni thickness. The magnitude of the AMR3 is approximately half the AMR2. At 5K, the magnitudes of all AMRs increase with Ni thickness except the AMR3 for the sample with 3Å of Ni is anomalously low. The temperature variations for AMR2 and AMR3 are quantitatively different but each varies by approximately a factor of 3 from 5K to 305K

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# **PK10**

#### Interface perpendicular magnetic anisotropy in thick amorphous CoSiB film by Pt laver

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Recently, magnetic tunnel junctions (MTJs) with perpendicular magnetic anisotropy (PMA) have attracted a lot of research interest, because they have a potential for spin transfer torque magnetic random access memories (MRAMs). Devices with perpendicular magnetic anisotropy have lower switching currents than devices with inplane magnetic anisotropy. We have investigated the interface perpendicular magnetic anisotropy (IPMA) of amorphous CoSiB thin film. CoSiB/Pt exhibited perpendicular magnetic anisotropy. The structure of samples was SiO<sub>2</sub>Ta<sub>504</sub>Pt<sub>204</sub>CoSiBx/Pt<sub>204</sub> and SiO<sub>2</sub>Ta<sub>601</sub>Pt<sub>201</sub>CoSiBx. The thickness of CoSiB were varied in the range of from 2 to 40Å. The sample without Pt capping layer exhibited perpendicular magnetic anisotropy at tCoSiB=10Å but the sample with Pt capping layer exhibited perpendicular anisotropy at lower thickness tCoSiB=5Å. Therefore, it was expected that IPMA properties of amorphous CoSiB on Pt layer was suitable for STT-MRAM.

### PK11

#### Perpendicular magnetic anisotropy and superparamagnetism in Ta/ CoFeB/MgO structures

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Collapses of tunnel magnetoresistance in perpendicular magnetic tunnel junctions have been reported recently[1-2]. These observations indicate that the coercivity of the reference layer and sensing layer behaves differently at low temperature. We extend to study the temperature dependence of a series of films with SiO2 substrate/Ta<sub>5 nm</sub>CoFeB<sub>x nm</sub>MgO<sub>1 nm</sub>, 1.5<x<1.9, structures by sputtering. The perpendicular magnetic anisotropy (PMA) is revealed as the thickness of CoFeB < 1.7 nm after a post annealing process. In this report, we focus on the magnetic response of Ta/CoFeB17nm/MgO which has a critical thickness of the transition from perpendicular to in-plane. Although the magnetic hysteresis shows weak in-plane anisotropy at room temperature the magnetic anisotropy changes to perpendicular at low temperature. The magnetization vs. temperature further indicates a superparamagnetic behavior with blocking temperature ~ 150K. The characteristic of a super-paparamagnetism suggests a pancake shape nanostructure with diameter ~ 40 nm exists at the interface. The present study demonstrates that small ferromagnetic particles may exist even after a post annealing process and the temperature-dependent coercivity is associated with an "intrinsic" structure of the interface in Ta/CoFeB/MgO structures.

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### PK12

#### Observation of symmetry of wavefunction in interface controlled Co/ Pd multilayer using magnetic compton profile

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Co/Pd multilayers have attracted attention as a potential application in high-density magnetic recording media because of their strong perpendicular magnetic anisotropy (PMA)[1,2]. It has been reported that thin Co layer and thicker Pd layer enhance the PMA enegy[2,3]. In this paper two Co/Pd multilayers which have smooth and rough interfaces, respectively, were studied. A Co1.5mmPd26mm multilayer with smooth interfaces was fabricated by the MBE technique. A Co1.6mmPd40mm multilayer with rough interfaces was fabricated by the sputter technique[4]. Both multilayers had almost the same PMA energy, 1.15Merg/ cc for the Co1.5mpPd26mm multilayer and 1.20Merg/cc for the Co1.6mpPd40mm multilayer, respectively. The symmetry of the wavefunction was measured using the magnetic Compton profile measurement. The symmetry of wavefunction of both the multilayers is almost the same. Population analysis of Co-3d shows that the population of magnetic quantum number |m|=2, 1 and 0 is 50%, 36% and 14%, respectively. This suggests that the smooth interface enhances the population of |m|=2 symmetry and hence enhances the PMA energy[5-7]. In conclusion smooth interface enhances |m|=2 symmetry and hence enhances the PMA energy even if a Co/Pd multilayer has a thinner Pd layer.

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# **PK13**

#### Tunable spin configuration in [Co/Ni] adjoined NiFe exchange spring structures

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Future spintronic devices based on spin-transfer torque require a delicate engineering of constituent magnetic layers. More specifically, realization of a tunable spin configuration is expected to allow for operation in zero magnetic field, as well as increased output efficiency and quality, e.g. linewidth, of the device.[1] Here, we report on a tunable spin configuration in exchange coupled NiFe, with in-plane magnetic anisotropy (IMA), to a [Co/Ni]×4 multilayer, with perpendicular magnetic anisotropy (PMA). A set of sputter deposited [Co/Ni]×4/NiFe samples (series 1) is compared with the reverse structure. NiFe/[Co/Ni]×4 (series 2), with various NiFe thicknesses (tNiFe) of 0.5, 1, 2 and 3 nm. Both series show the same eventual transition from dominant PMA to IMA as tNiFe is increased. However, this transition appears for smaller tNiFe in series 2, most likely due to the initially sputtered NiFe modifying the PMA of the [Co/Ni] multilayer. Most importantly, due to the competition between PMA and IMA, a tilted spin configuration can be stabilized and tuned for intermediate tNiFe

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#### **PK14**

#### Large perpendicular magnetic anisotropy in the MgO/CoFeB/Ta with thick Ta laver Tao Zhu

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It is widely recognized that current-induced spin transfer torque (STT) plays an important role in spintronic devices [1]. The magnetic tunnel junction with a perpendicular magnetic anisotropy (PMA) is a candidate for such STT devices because the PMA enables a small critical current density for current induced magnetization switching. The perpendicular magnetic anisotropy (PMA) in the Ta capping film, MgO/CoFeB/Ta, has been investigated by using anomalous Hall effect. Ta capping structures, MgO11CoFeB122Ta05544 (numbers are the nominal thicknesses in nanometers, Ta=300°C), were prepared. The out-of-plane squareness increases with the Ta cap layer thickness from 0.55 to 0.94 nm and reaches unity at Ta thickness of 0.94 nm. Different from the previous report [2], the perpendicular anisotropy of CoFeB retains at which Ta thickness is larger than 0.94 nm. Unity squareness can be observed at the Ta cap layer thickness up to 4.4 nm. It is clear that there is no strong relationship between PMA and the Ta cap layer thickness. Moreover, A large PMA in the Ta capping film is due to the large effective anisotropy.

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#### **PK15**

#### Strain induced overlayer effect on perpendicular magnetic anisotropy in Ta/CoFeB/MgO structures

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The perpendicular magnetic anisotropy (PMA) in ultrathin CoFeB-MgO based structure is an important feature for the applications on magnetic random access memory1. However, the interface structure is sensitive to seed (cap) metal layer and the annealing treatment2. In this study, we fabricated a series of structures of Sub/Ta5nmCoFeB12nmMgO(0<x<5 nm)/Ta(0<y<5 nm) and found that the PMA not only depends upon the thickness of MgO but also depends upon the Ta overlayer on top of MgO. In Ta/ CoFeB/MgO(0<x<5 nm)/Ta(5 nm), PMA is observed as the thickness of MgO is over 0.75 nm (~3 atomic monolayer) and a maximum of Hk ~ 5 kOe at MgO thickness ~1nm is identified. In Ta/CoFeB/ MgO1 nm Ta(0<y<5 nm), we found both Hk and coercivity increases from 2.5 kOe to 5 kOe and 15 Oe to 30 Oe, respectively as the thickness of Ta varies from 1 to 5 nm. These results may conclude that the match of the lattice parameter between CoFe and MgO is crucial for the PMA in the CoFeB-MgO based structure after a post anneal process. The strain induced by MgO layer and by Ta overlayer may substantially improve the (001) texture of CoFe and thus improves the PMA in these structures.

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# PK16

#### Ultrathin Co/Pt films with high thermal stability and large effective magnetic moment

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### PK17

Magnetic characteristics of amorphous [CoSiB/Pt]n multilavers Woosuk Yoo1, Kyujoon Lee1, Myung-hwa Jung1\*, Insung Park2, Taewan Kim2, E.h.m. Van Der Heijden3 and H.j.m. Swagten <sup>1</sup> Dept of Physic, Sogang University Seoul, Korea <sup>2</sup> Dept of Physic, Sejong University Seoul, Korea <sup>3</sup> Department of Applied Physics, Eindhoven University of Technology, Netherlands

The development of electronics demands an increasing effort in attaining more efficient electronic memories. The materials with perpendicular magnetic anisotropy (PMA) are important in applications of electronic memories. One promising new material is amorphous PMA multilayers [CoSiB/Pt]n with n the number of repeated bilayers. By using a VSM-SQUID measurement setup, the magnetic characteristics of the [CoSiB/ Pt]n are gained. The anisotropy constants k1eff, K1, K2 are obtained by fitting the hard-axis magnetization curve with GST equation. The hysteresis loops appear slanted increasing n. Furthermore magneto- optical Kerr effect (MOKE) imaging confirms magnetic domain wall motion. The coercive field is first increasing with n and after a maximum, decreasing with n. This result is caused by the RKKY coupling between the magnetic layers and the balance between the domain wall energy scaling with n and the magnetostatic energy scaling with n2. In addition the reversal process is not depending on the strength of the initial magnetization field.

### PK18

#### Effects of phase distribution and grain size on the effective anisotropy and coercivity of nanocomposite PtCo permanent alloy

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Phase distribution and grain size are the key factors controlling the coercivity of nanocomposite permanent magnets. Feng et al. have investigated the effective anisotropy and coercivity of nanocomposite Nd2Fe14B/a-Fe magnets. [1] The latest research shows that the PtCo alloy is composed of nano-structured soft and hard magnetic phases. [2] This system can be used as an ideal model system for studying exchange coupling between the magnetically hard and soft phases. There exist three types of grain interfaces with different exchange-coupling strengths in nanocomposite PtCo permanent magnets. It is necessary to describe the proportions of the three kinds of grain interfaces for analyzing the influences of the exchange-coupling interactions on the magnetic properties. We assume that soft and hard phases distribute randomly in nanocomposite permanent magnets and that all grains have cubic shapes. We caculate the dependence of the coercivity Hc of a nanocomposite on the hard-grain size Dh. For a given volume fraction vh and soft-grain size Ds, Hc shows a maximum value as a function of the hard-grain size Dh, and the maximum position shifts to a smaller Dh with the decreasing volume fraction of the hard phase. We shows the comparisons between the calculated coercivity Hc and experimental results

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# **PK19**

#### Perpendicular magnetic anisotropy and metal layer effect in MgO/ CoFeB/cap (cap = Ta, Ru, and Nb)

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Perpendicular magnetic anisotropy (PMA) has been observed in MgO/CoFeB/Ta but not in MgO/CoFeB/Ru, indicating that the metal cap layer is crucial<sup>1-2</sup>. In this report, we extend the study to MgO/CoFeB12 nmNb(1 - 5 nm) by sputtering and find again that PMA can be induced by Nb cap layer. Our experimental results also show that Nb and Ta have strong interface effect but the PMA occurs in a small window (~1 - 3 nm) while Ru shows much weaker effect and the induced interface anisotropy constant (Ki) is basically not high enough to overcome the bulk anisotropy contribution

#### **PK20**

#### Dependence of perpendicular magnetic anisotropy of CoFeB thin films on thickness of MgO overlayer

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The CoFeB/MgO interface is known to favor perpendicular magnetic anisotropy (PMA) which is already found application in the efficient magnetic tunnel junctions and electric field control the magnetization. Recently, the PMA of CoFeB/MgO structure has been achieved in both experiment [1,2] and theory [3], however the origin of PMA is still ambiguous. In this scope, the structure Ta<sub>5nm</sub>Co<sub>10</sub>Fe<sub>70</sub>B<sub>2009-201nm</sub>MgO<sub>(0.601nm</sub>Ta<sub>5nm</sub> (unit in nm) was investigated. All samples were fabricated by sputtering method on thermally oxidized Si substrates. We studied the dependence of the total magnetic moment and the effective magnetic anisotropy energy of the Co10Fe20B20 layer. Magnetic anisotropy is found to be very sensitive to MgO thickness, favouring to perpendicular direction with thin MgO film (the maximum interface anisotropy was 1.74 erg/cm2 with the CoFeB/MgO<sub>1mm</sub> structrure. Although thin MgO samples have thick magnetic dead layer, resulting in the reduction of total magnetic moment, only the existence of the magnetic dead layer cannot explain the change of the PMA. We found that the thickness of MgO overlayer plays the very important role for the direction of magnetic anisotropy easy axis. This work is supported by CREST, Japan Science Technology Agency, Japan.

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### PK21

#### The study of MgO based surface anisotropy of CoFeB layer

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STT-MRAM (Spin transfer torque-magnetoresistive random access memory)[1] has considered as an outstanding candidate for the next generation memory device. Higher anisotropy constant K and low saturation magnetization Ms are required for the perpendicular magnetic layers to satisfy the high thermal stability and low current density for current induced magnetic switching. The PMA(perpendicular magnetic anisotropy) have been reported in Fe-Pt alloy[2] ordered by L10, Co/Pd multi-layer[3] and CoFeB-MgO[4,5]. We report the PMA of CoFeB based with MgO using amorphous paramagnetic FeZr buffer layer. Samples are deposited on SiO<sub>2</sub> substrate by RF magnetron sputtering machine at room temperature. The structures are  $SiO_2FeZr_{2nm}CoFeB(tCoFeB:0.5-2)/MgO_{2.5nm}Ta_{1nm} \ (in \ nm) \ and \ annealing \ for \ 1 \ hour \ at \ 400^\circ C$ Figure 1 shows Hall resistance of PMA of CoFeB after annealing. During the annealing treatment, the lattice of CoFeB became BCC lattice structure along with crystal structure of upper MgO. Therefore, the interface anisotropy is formed between CoFeB and MgO layers in rather thin thickness. In summary, we observed the PMA in CoFe/MgO using the FeZr buffer layer We obtain a high anisotropic constant K with a various conditions. This method can be an excellent candiate for PMA- MTI with low critical current density

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# **PL01**

#### BiFeO<sub>3</sub> thickness dependence of exchange bias in polycrystalline **BiFeO3/FeNi bilavers**

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Due to both ferroelectric and antiferromagnetic transition temperatures above room temperature, multiferroic BiFeO3 has attracted much attention for uses in the electrical control of magnetic random assess memory. Exchange bias in ferromagnet (FM)/ antiferromagnet (AFM) bilayers has also attracted much attention because of its importance in developing magneto-electronic devices. Although there have been some reports on the exchange bias employing BiFeO, as antiferromagnetic laver, most of them the BiFeO<sub>3</sub> layers were grown epitaxially. In this presented work, single phase polycrystalline BiFeO<sub>3</sub> films were successfully grown on Si substrate with LaNiO<sub>3</sub> buffer laver. Clear exchange bias were observed at room temperature in a series of BiFeO<sub>10</sub>FeNi(3.6nm) bilayers, in which t varies from 8 nm to 240 nm. With increasing t, both the exchange bias field and the coercivity increase sharply and approach maxima when t is about 40 nm, which is close to one half of the spin cycloidal modulation period (64 nm) of the bulk BiFeO<sub>3</sub> material. Our results suggest that the cycloidal spin structure may still exist in the polycrystalline BiFeO<sub>2</sub> film and the interfacial uncompensated pinned spins may be responsible for exchange bias in the BiFeO3/FeNi bilayer films.

#### Temperature-dependent magnetic anisotropies in epitaxial Fe/CoO/ MgO(001) system studied by the planar Hall effect

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The exchange coupling between ferromagnet (FM) and antiferromagnet (AFM) has been intensively explored due to fundamental interest and its technological applications. The effect of the magnetic anisotropy of the AFM layer on the exchange coupling has not been well studied, which can only be performed in single-crystalline FM/AFM system. In this contribution, we investigated the exchange-induced inplane magnetic anisotropies in a single-crystalline Fe/CoO/MgO(001) system in the temperature range of 10-400K. The temperature-dependent magnetic anisotropies of the Fe film were quantitatively measured with the torque technique utilizing the planar Hall effect. We found the field cooling can induce a very strong uniaxial anisotropy up to 0.196erg/cm2 in Fe film with the easy axis(EA) // CoO<110>, and also a 4-fold anisotropy with the EA//CoO<100>. Such EA of the exchange-induced anisotropy was found insensitive to the cooling field orientation and only related to the CoO crystalline axis. The CoO thickness dependent exchange-induced anisotropy will be discussed in this contribution

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### PL03

#### Structural changes and magnetic properties of ultrathin Fe/Pt(111) films influenced by oxygen exposure

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Magnetic properties of Fe/Pt(111) films have received much attention in the past decade [1-2]. In this contribution, we focus on the structures and magnetic properties of ultrathin Fe/Pt(111) films influenced by oxygen exposure. Fe atoms in the submonolayer region adapt a smaller layer distance as compared to the BCC structure of a Fe layer above one monolayer (ML). At saturation conditions for absorbed oxygen on Fe/Pt(111), the increase of oxygen in the films coincides with the Fe thickness while the amount of adsorbed oxygen in the submonolayer range is rather small. In addition, the amount of oxygen exposure to achieve a saturation condition increases to a maximum value followed by a reduction for Fe/Pt(111) thicker than 1 ML. This shows the structure related oxidation efficiency of Fe/Pt(111). The growth of Fermi edge at 620 K shows the segregation of Pt atom to the top layers. Only the Fe/Pt(111) thicker than 3 ML shows the enhanced coercive force after thermal annealing. This could be explained by the blocking of the formation of FePt alloy by the iron oxides

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# **PL04**

#### Initial growth of bcc Co films on Au(001) studied by STM/BH imaging T. Kawagoe\*, E. Wakabayashi, Y. Murasawa and T. Sakata Osaka Kyoiku University, Japan

We have studied the surface structure and magnetism of metastable bcc Co films (>1 ML) on Au(001) [1]. Here we report the initial growth of submonolayer Co films by means of barrier height (BH) imaging using STM, which has capability of element discrimination. The experiments were performed in a UHV-MBE system equipped with an STM unit. After the clean Au(001) reconstructed surface was confirmed by STM, a submonolayer Co was deposited on it at RT. Spatially resolved BH maps, together with the topographic images were obtained at RT using a tip-sample distance modulation technique. We have observed numerous islands with different sizes and heights after 0.15 ML Co coverage and successfully obtained, from the BH imaging, an element-specific contrast, i.e. recognizing aggregated Au islands and Co islands, and information about inhomogenities of BH with proper consideration of the artifacts near the step edges. BH of metastable bcc Co film was observed for the first time, and that showed a large BH value (~6 eV), whereas the observed BH of the Au(001) surface (~3.5 eV) was consistent with the previous results.

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#### PL05

#### Magnetic field induced "switching" of the nanodomain state of ferromagnet - antiferromagnet frustrated system Alexander I Morosov and Alexander S Sigov Electronics Department, MSTU MIREA, Russia

The presence of atomic steps on the interface between the nanolavers of ferromagnet and antiferromagnet with uncompensated atomic planes parallel to the interface causes frustrations of the interlayer exchange interaction. If the exchange integral for the ferromagnet is less than that one for the antiferromagnet, then, for the layers of equal thickness, and in the case where the characteristic distance R between the neighboring edges of the steps on the interface exceeds the layer thickness, the ferromagnetic layer appears to be broken down into 180° domains. The domain walls separating such domains generated by frustrations arise at the atomic step edges and penetrate the ferromagnetic layer, the domain wall thickness increases with increasing distance from the interface. An external magnetic field applied parallel to the magnetizations of the domains, being in excess of some critical value B\*, makes the monodomain state of the ferromagnetic layer more favorable. If so, the domain walls with more complicated structure depending on the field magnitude arise between 180° domains in the antiferromagnet layer. The B\* value can be obtained from the equality of the gain in the Zeeman energy and the difference between the energies of the domain walls in different lavers.

#### **PL06**

#### Stable structural and electronic properties of adsorption Tl on the clean Tl/Si(111) surface

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The structural inversion asymmetry shows up surfaces or interfaces. It leads to spin split electronic state (Rashba-Bychkov effect) in the material. In particular, heavy metals on the light elements like Si show strong spin-orbit splitting [1] and such property in Tl/Si(111) has also been confirmed both in simulation and experiment [2]. This surface has been observed as a clean one, indicating the stability due to the insulating property verified by the experiments and theoretical approaches [3,4]. By donating (or accepting) electrons, one can introduce carriers with spin polarization at the part of states, implying an interesting basis for spintronics devices. In this work, as a method of donating electrons, Tl atom which is the same element as the top surface of Tl/Si(111) is considered within the density functional first-principles approach of fully relativistic pseudo-potential and planewave basis. First, the stable adsorption site has been investigated with using the  $\sqrt{3}\times\sqrt{3}$  superlattice of Tl/Si(111). It was found that one of hollow site models is most stable in the adsorption models investigated. In the presentation, we report details of the stable structure and, for typical systems, we also present electronic structures and spin splitting in the band dispersion relations.

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# PL07

### Exchange bias effect in Ni(Zn)O film

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Exchange bias (EB) has received significant attentions for fundamental interest and potential applications in magnetic memory and read heads. Recently, EB is observed in various alloys and compounds [1]. Magnetic nanoparticles have been found promising for significant EB effect where surface spins were found to have different magnetic anisotropy than core spins, causing sizable EB effect. Compared to the studies of EB in nanoparticles, these studies in film geometry having chemically single phase are rarely investigated, which is rather promising for technological applications. Here, we report EB in Ni1xZnxO films where Zn substitution in NiO significantly modifies magnetic properties and EB effect. Atomic Force Microcopy exhibits formation of granular films of average grain size ≈25 nm. We note that EB field (HE) and coercivity (HC) decrease considerably with Zn substitution at x = 0.3 and HE vanishes at x = 0.5. Appearance EB in NiO film is suggested due to significant difference in anisotropy between core and surface spins, as it has been conjectured for NiO nanoparticles. The Zn substitution appreciably decreases the difference between core and surface anisotropy causing the reduction of HE as well as HC.

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#### **PL08**

### Electronic properties of transition metal oxides interface between GdTiO<sub>3</sub>SrTiO<sub>3</sub>

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The creation of a two dimensional electron gas (2DEG) at the interface between two insulating oxides has been a subject of great interest, and much experimental activity has centered on the LaAlO<sub>3</sub>SrTiO<sub>3</sub> interface. The (001) GdTiO<sub>3</sub>SrTiO<sub>3</sub> interface has the same naive polarization discontinuity as the former system, but unlikely it actually exhibits a 2DEG with the expected e/2 charge per planar unit cell.[1] Motivated by recent experiments, which show that the 2DEG primarily lives in the SrTiO<sub>3</sub>, we construct a theoretical tight-binding model and discuss the electronic properties in comparison with experiment, and consider possible magnetism in this system at low temperature

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#### PL09

#### Parallel ferromagnetic resonance and spin wave excitation in exchange-biased NiFe/IrMn bilayers

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Ferromagnetic Resonance (FMR) investigation of magnetron sputtered exchange-biased Si111Ru7mmNiFe1IrMn6mmRu5mm bilayers reveals that spin wave and uniform resonance modes are excited by the microwave field. The samples studied were produced in the presence of a 400 Oe magnetic field to set the unidirectional anisotropy and the thickness t of the NiFe layer varied in the range 55-120 nm. The FMR experiments were done at room temperature using a commercial spectrometer with swept static magnetic field and the usual detection techniques. In-plane angular variation of resonance fields of spin wave and uniform modes reveal for both modes the dependence on  $\cos\!\theta,$  where  $\theta$  is the field angle with respect to the anisotropy axis. The unidirectional anisotropy field for the spin wave mode is twice larger than that for the uniform mode; they are equal to 40 and 20 Oe, respectively, for the Si111Ru7nmNiFe65nmIrMn6nmRu5nm sample. Spin wave resonance theory give for the NiFe layer of this sample the exchange constant of  $1.2 \times 10^{\circ}$ (-6) erg/cm. The parallel FMR results at microwave frequencies of 9.45 and 34.1 GHz agree with the values given by the resonance fields of spin wave and uniform modes also excited in perpendicular FMR at a frequency of 9.45 GHz.

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PL02

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# **PL10**

#### Morphological consideration for post-annealed Cr2O3 surface and exchange bias in bilayer system

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Magnetoelectric Cr<sub>2</sub>O<sub>2</sub> has been received renewal of attentions as a promising material for the low energy consumption electric controlled spintronic devices such as MERAM [1]. High exchange bias field (Hex) between ferromagnetic layer and Cr2O3 thin film layer, and low coercivity (Hc) are essential for the applications. Previously we reported the high Hex (~500 Oe) at the wide temperature range from 10 K to 250 K and the low Hc (~30 Oe) at 250 K for post-annealed Cr<sub>2</sub>O<sub>3</sub> thin film system [2]. In this study we compared the morphology and magnetic properties of the post-annealed sample with that of samples fabricated by other manners and examined the origin of the high Hex and low Hc of the post-annealed sample. The post-annealed sample was identified to be polycrystalline, which would contribute to the low Hc. From TEM and AFM images. flat (1-102) planes (r-planes) were observed diagonally to the film surface which are formed during the post annealing by grain growth due to the low surface energy [3]. From the fact that Cr2O3 has uncompensated spin not only c-plane but also r-plane, we conclude that the origin of the high Hex is the uncompensated r-plane.

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### PL11

#### Magnetic properties of [NiFe/NiFeCuMo/NiFe]/FeMn¬multilayers depending on a thickness of NiFeCuMo layer

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Variations in the exchange bias field (HEX) and coercivity (HC) of the intermediately super-soft magnetic NiFeCuMo layer were investigated for different thicknesses of the bottom NiFe layer. HEX of the triple pinned NiFe<sub>4 nm</sub>NiFeCuMo (tNiFeCuMo = 1 nm)/ NiFe4 nmFeMn multilayer has a maximum value that is more less than that of a single pinned NiFes me FeMn layer. If a NiFeCuMo layer is inserted between the pinned and free NiFe layers, this multilayer system can be used as a giant magnetoresistive spinvalve (GMR-SV) device for biosensors with improved magnetic sensitivity.

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#### PL12

#### Structural and magnetic properties of antiferromagnetic Heusler Ru<sub>2</sub>MnGe epitaxial thin films

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Recently, it is suggested that spin transport phenomena observed in conventional spintronics ought to occur in systems where all FM materials are replaced by antiferromagnetic (AFM) materials. Although various studies have investigated FM Heusler alloys, very few studies have been made on AFM Heusler alloys in terms of spintronic applications. The AFM transition temperatures (TN) of Heusler alloy such as Fe<sub>2</sub>VSi can be controlled by inducing biaxial strain[1]. Thus we investigate the correlation between structural and magnetic properties of epitaxial thin films of AFM Ru<sub>2</sub>MnGe. Structural characterization revealed that Ru<sub>2</sub>MnGe films on MgO were subjected to compressive strain (c/a = 1.0042). The strained Ru<sub>2</sub>MnGe exhibited markedly enhanced TN of 353 K compared to the bulk value (316 K)[2]. The in-plane Mn-Mn distance for the strained Ru<sub>2</sub>MnGe is slightly smaller than that of bulk Ru<sub>2</sub>MnGe, which may contribute to reducing the total J2 values, resulting in the enhancement of TN. We also observed the exchange coupling between Heusler-type FM half-metal Fe<sub>2</sub>CrSi and AFM metal Ru<sub>2</sub>MnGe. The present AFM Heusler alloy with relatively high TN is useful to fabricate high-quality all Heusler-type half-metal FM/AFM junctions and is a promising material for the emerging field of AFM spintronics.

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# PL13

Exchange-spring phenomenon of ultrathin Fe/CoPt bilayer

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A large enhancement in the polar magneto-optic Kerr signal was observed after one monolayer (ML) of Co adatoms on Pt(111) was annealed above 600 K to form a Co-Pt alloy at the interface [1]. In this contribution, the composition, surface structure and magnetic properties of ultrathin Fe films on the top of a Co-Pt(111) surface alloy has been investigated. As the Fe coverage increases on the top of the Co-Pt alloy, the polar Kerr rotation increases. By way of systematically changing the Fe thickness, the magnetic properties of Fe/CoPt in both the polar and longitudinal configurations are investigated for Fe thinner than 7 ML. At an optimal condition for both the Fe thickness and the annealing temperature for preparing the Co-Pt alloy, the polar coercive force remains the same as the Fe coverage increases. The Fe/CoPt system exhibits magnetic exchange-spring behavior in the polar configuration. The interfacial conditions for the Fe/COPt system are demonstrated to play an important role on the exchange-spring phenomenon and will be discussed in this contribution.

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# PL14

Neutron magnetic scattering study in manganite thin film system H. Nakao<sup>1</sup>\*, H. Yamada<sup>2</sup>, K. Iwasa<sup>3</sup>, J. Okamoto<sup>1</sup>, Y. Yamasaki<sup>1</sup>, Y. Murakami<sup>1</sup> and A.

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Perovskite manganites show various interesting phenomena including colossal magnetoresistance (MR) due to a close interplay among charge, orbital, spin, and lattice degrees of freedom. The superlattice  $[(LaMnO_3)_m(SrMnO_3)_m]_n$  was investigated as a stage to control the Mn valence artificially. The physical properties were reported to strongly depend on the periodicity m, and the interface state between LaMnO<sub>3</sub> and SrMnO<sub>3</sub> layer was noted.[1] The stacking structure of Mn valence state by resonant x-ray scattering technique has been studied.[2] It elucidated that the physical properties just depend on the quality of the film. Recently, however, new large MR effect was discovered in the high quality superlattice system.[3] To understand the MR effect microscopically, the Mn valence state and the magnetic field effect were studied. Here the magnetism is also important, and then neutron magnetic scattering experiment was started using the neutron spectrometer TOPAN in JRR-3. As a first step, SrMnO<sub>3</sub> film with 80 nm thickness has been investigated. A signal with about 2 cps at (0.5, 0.5, 0.5) corresponding G-type AFM was certainly discovered. The temperature dependence of the signal intensity was also measured. Consequently we have succeded in observing magnetic scattering from thin film system.

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### PL15

#### Direction and temperature dependences of exchange bias and coercivity of NiFe/Cr<sub>2(1-x)</sub>Fe<sub>2x</sub>O<sub>3</sub>(x= 0.25, 0.4) bilayers Sanghoon Ki, Byeong-geon Kim and Joonghoe Dho\* *Kyungpook National University, Korea*

The antiferromagnets  $\alpha$ -Fe<sub>2</sub>O<sub>3</sub> and  $\alpha$ -Cr<sub>2</sub>O<sub>3</sub> have the same crystal structure but different magnetic orderings. The AF spin structure for a-Cr<sub>2</sub>O<sub>3</sub>(TN~307 K) is represented by the up-down-up-down spin ordering along c-axis, while that for a-Fe<sub>2</sub>O<sub>3</sub>(TN~953 K) is done by the up-down-down-up one. As reported in previous neutron diffraction study for the solid solution Cr<sub>2(1-x)</sub>Fe<sub>2x</sub>O<sub>3</sub>(CFO), such contradictory spin structures may result in a new type AF structure such as spiral order or cycloidal order. Here, we report exchange bias in NiFe/Cr<sub>2(1-x)</sub>Fe<sub>2x</sub>O<sub>3</sub>(x= 0.25, 0.4) bilayers grown on c-plane and r-plane Al<sub>2</sub>O<sub>3</sub>. The direction and temperature dependences of magnetic property of NiFe/CFO layer substrates were checked by a surface magneto-optical Kerr effect (SMOKE) setup. At room temperature, the M(H) loop of the NiFe layer exhibited enhanced coercivity and exchange bias depending on the measurement direction and the temperature. In order to see the direction dependence of magnetic property of NiFe layer, the magnetic field was applied to two orthogonal directions within the plane. Interestingly, the temperature dependence of M(H) loop exhibited some features of spin reorientation of CFO layer below antiferromagnetic transition temperature.

# PL16

# Structure and magnetic properties of epitaxial Fe/MgO/Si (001) heterostructures

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The ferromagnetic metal/semiconductor structures with inserting a thin MgO tunneling barrier are expected to play a key role in spin-based electronic applications to enhance spin injection efficiency. Recently, spin-based devices using silicon channel enable seamless integration with conventional MOS-FET technology. However, there are only a few reports on ferromagnetic metal/MgO/Si system due to large lattice mismatch (3.5%) between Si and MgO. The study on strain induced microstructure of the Fe and MgO layers is essential because the lattice mismatch between Si and MgO results in distortion in both MgO and Fe layers. In addition, the magnetic property of the ferromagnetic metal with respect to its microstructure evolution is also of importance in terms of spin injection. In this study, Fe thin films of thinner than 2 nm, were prepared on Si (001) substrates with 4 nm thick MgO buffer layer by a molecular beam epitaxy system at room temperature and 200 °C, respectively. X-ray diffraction and reflection high-energy electron diffraction results show epitaxial relationship of the system. In-plane four-fold magnetic anisotropy was observed by vibrating sample magnetometer, suggesting that elevated growth temperature of the Fe films leads to the strain relaxation and the suppression of two-dimensional Fe layer formation.

#### PL17

# Uncompensated magnetic moment around Mn<sub>3</sub>Ir / Fe-Co-Ni bilayer interface

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We study the microscopic origin of the uncompensated moment in Mn<sub>3</sub>Ir / Fe-Co-Ni bilayer, which is practically utilized system due to the excellent property of exchange anisotropy, by the theoretical and experimental approach. In order to induce exchange anisotropy, the magnetic structure in the ferromagnetic / antiferromagnetic bilayer needs to have some asymmetry against the field reversal, and the uncompensated moments around the interface is a candidate for the spin asymmetry. We previously reported that there is a linear correlation between the exchange anisotropy and uncompensated moments measured by the x-ray magnetic circular dichroism (XMCD) experiment in Mn-Ir / Co-Fe bilayer [1]. Therefore, in the present work, we investigate the composition dependence of the uncompensated moment in Mn<sub>3</sub>Ir / Fe-Co-Ni by performing the first-principles calculation and the XMCD measurement to clarify the origin of the uncompensated moment. In the comparison of the calculation and experimental results, the obtained uncompensated moments are in good agreement for their sign and relative magnitude [2]. Furthermore, the calculated result reveals that the uncompensated moment originates from the local magnetic correlation near the interface. So we find that the reorientation of the magnetic moments of Mn occur at the very interface with the ferromagnetic layer.

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### PL18

#### Unusual exchange bias effects in NiFe/Mn thin films induced via ionbeam bombardment

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When spins at the interface between a ferromagnet (FM) and an antiferromagnet (AF) couple, a unidirectional anisotropy occurs, resulting in exchange bias [1,2]. In this study, we wish to identify the dependence of the exchange bias on the interface microstructure in NiFe/Mn thin films. A series of [NiFe/Mn] multilayer thin films were fabricated using a dual ion-beam deposition technique [3]. Different Ar-ion deposition voltages created striking differences in magnetic properties. The [NiFe/Mn] (VEH= 0 V, i.e. unbombarded) thin film (~200 nm) at 5 K when 20 kOe field-cooled (FC), exhibited an enhanced coercivity (Hc~935 Oe) and positive exchange bias field, Hex~+130 Oe. By contrast, the [NiFe/Mn] (VEH= 150 V) Ar-ion hombarded thin film (~20 nm) at 5 K when 20 kOe field cooled showed a larger positive Hex (~+160 Oe) with smaller Hc (~610 Oe) accompanied by an asymmetric hysteresis loop. The observed enhanced positive Hex indicated that the transition from FM to AF coupling was made possible by changing the ion-beam bombardment voltages. Further, the role of Ar ion-bombardment seems to change both the spin orientations in FM NiFe and AF Mn and thus gives rise to the positive enhanced Hex.

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# PL19

#### Exchange bias effect in BiFeO<sub>3</sub> thin films Kil-dong Sung and Jonghoon Jung\* Physics, Inha University, Korea

We fabricated BiFeO3 thin films with Bi<sub>2</sub>O<sub>3</sub> or  $\gamma$ -Fe<sub>2</sub>O<sub>3</sub> impurities by changing oxygen partial pressures and characterized their magnetic properties. Since Bi<sub>2</sub>O<sub>3</sub> impurity is a non-magnetic material, only antiferromagnetic spins of BiFeO<sub>3</sub> will contribute to whole magnetic properties. But  $\gamma$ -Fe<sub>2</sub>O<sub>3</sub> impurity is a well known ferrimagnetic material, so there is a possibility to show exchange bias effect with antiferromagnetic spins of BiFeO<sub>3</sub>. We measured M(H) loop at various temperatures after zero- and  $\pm$ 5 T-field cooling. Almost no but very small amount of M(H) hysteresis loop was appeared in BiFeO<sub>3</sub> with Bi<sub>2</sub>O<sub>3</sub> impurity. Which indicate that there is uncompensated spins in BiFeO<sub>3</sub> which generate exchange bias effect with  $\gamma$ -Fe<sub>2</sub>O<sub>3</sub>. Significant hysteresis loop was appeared in BiFeO<sub>3</sub> with  $\gamma$ -Fe<sub>2</sub>O<sub>3</sub> impurity. Depending on the field cooling direction, the M(H) loops also shifted to the opposite directions. But unlike other conventional exchange bias systems, the interaction between ferrimagnetic spins of  $\gamma$ -Fe<sub>2</sub>O<sub>3</sub> and antiferromagnetic spins of BiFeO<sub>3</sub> is not so strong as shown in the training effect result.

#### PL20

# Experiment evidence to the existence of interaction between antiferromagnetic domains in IrMn/Pt/Co/Pt multilayers

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Exchange bias phenomenon has become an integral part of the modern magnetism, with implications in basic research and applications in magnetoelectronic devices. In the numerous experimental and theoretical studies, the surface layer of antiferromagnetic (AFM) material is often regarded as an assembly of uniaxial domains, which are independence from each other. This basic assumption is really suspect. In our presented work, experiment evidence of interaction between AFM domains has been found. The IrMn/Pt/Co/Pt perpendicular exchange bias systems were fabricated and treated by AC demagnetization (ACDM) parallel to the film plane at high temperature. The purpose of ACDM is to get the pinned demagnetization state in ferromagnetic layer. In the corresponding hysteresis loop, two minor-loops shift oppositely along the magnetic field axis and locate totally on one side of the vertical axis. Compare the training effects of the minor-loop and full-loop, we found hat one minor-loop measurement gives influence to the other. It is reasonable that the influence is attributed to the interaction between the IrMn surface domains. Our experiment facts can not be explained by the past theory models, new energy item of interaction between AFM domains should be considered.

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### PL21

# Nanoscale investigation of the Cr/Fe(001) interface grown by oxygen assisted epitaxy

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Cr deposition on Fe(100) has been the subject of intense investigations in the past, being an excellent model system for thin film magnetism, and also because Fe/Cr multlayers played a fundamental role in the development of giant magneto resistive devices [11]. The magnetic properties of such systems can be affected by parameters like interfacial mixing, strain and defects. For instance, previous investigations have shown that, within the atom distribution resulting from Cr incorporation into the Fe substrate, nearest neighbour Cr-Cr coupling is energetically less favorable than next nearest neighbor coupling, due to magnetic firstrations [2]. Here we employ experimental and theoretical methods to highlight and characterize deviations in the micromagnetic structure of the CrFe interface due to the presence of adsorbed O, which induces a layer by-layer Cr growth at room temperature, as we recently demonstrated [3]. Scanning tunnelling microscopy (STM) is used to investigate the deposition of Cr on the Fe-p(1x1/O) surface since the first stages of interface formation. Atomic resolution STM images reveal the occurrence of different magnetic configurations, including first-neighboring Cr atoms in the Fe matrix. The observations are discussed in the light of density-functional simulations with the generalized gradient approximation and plane waves basis set [4].

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# PM01

# Correlation of soft magnetic properties with free volume and medium range ordering in metallic glasses

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The effect of cooling rate on thickness and soft magnetic properties of the metallic glass of alloy composition  $Co_{64,3}Fe_{3,5}Si_{16}B_{14}Ni_2$  have been studied. The amorphous structure has been characterized by positron annihilation spectroscopy (PAS) and fluctuation electron microscopy (FEM) techniques. These studies suggest that the first lifetime component is associated with MRO present in the amorphous structure. Theoretical calculations showed that second component is associated with nanovoids having free volume equivalent to that of a vacancy defect consisting of 4 and more atom vacancy cluster. Coercivity of the samples showed small variations with wheel speed and it was found to decrease with the decrease in effective size of the defects present in the samples. This study reveals that, in case of amorphous metallic glass, the variation in coercivity is mainly controlled by the defects formed during processing rather than the small variations in MRO introduced due to different cooling rates.

# PM02

# Microscopic magnetic hysteresis measurement of amorphous Tb-TM (TM=Fe and Co) thin films by magnetic Compton scattering

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Rare earth transition metal (RE-TM) alloys have attracted considerable attention from the viewpoint of fundamental research and technology. For instance, they have been studied as candidate heat-assisted magnetic recording media for high-density magnetic storage. It is known that magnetic moment of RE and TM form ferrimagnetisms and have anisotropy dispersion along magnetic easy axes, so-called sperri-magnetism [1]. The magnetic switching process of sperri magnetism is considered not to be simple. In this report, we have study spin, orbital and also element specific magnetic switching process of amorphous Tb-TM (TM=Fe and Co) films by means of combination of magnetic Compton scattering [2, 3] and conventional macroscopic magnetization measurements. As specimen, amorphous Tb<sub>45</sub>Co<sub>57</sub> and Tb<sub>25</sub>Fe<sub>55</sub> (O<sub>13</sub>) thin films were fabricated by an RF sputtering onto a thin Al foil substrates. The macroscopic magnetization curves were measured by SQUID and VSM magnetometers. Magnetic Compton scattering neasurements were carried out at AR-NE1A of KEK [2] and at BL08W of SPring-8 [3], Japan. The present results suggest that magnetic switching can be different among the spin moment, the orbital moment, and even the states of electrons, e.g., the Tb 4f electron, TM 3d electron, and itinerant electron.

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### PM03

#### A significant reduction of hysteresis in MnFe(P,Si) compounds

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The magnetocaloric effects in  $Mn_{1.3}Fe_{0.7x}Co_{x}P_{0.46}Si_{0.54}$  compounds with x = 0, 0.025, 0.05 and 0.1 are studied systematically. X-ray diffraction shows that the compounds in the range from x = 0 to 0.1 crystallize in the Fe<sub>2</sub>P-type hexagonal structure with space group P-62m symmetry. Magnetic measurements show that the paramagnetic ferromagnetic transition temperatures range from 247 to 298 K. The maximum of magnetic entropy changes in  $Mn_{1.3}Fe_{0.7}P_{0.46}Si_{0.54}$  compound reaches 8.3 J/kgK in a field change of 0 to 1.5 T, and the thermal hysteresis of these compounds is less than 3 K. The maximum adiabatic temperature change is 2.2 K in  $Mn_{1.3}Fe_{0.7}P_{0.46}Si_{0.54}$  and  $Mn_{1.3}Fe_{0.65}Co_{0.05}P_{0.46}Si_{0.54}$  compounds for a field change from 0 to 1.48 T.

# **PM04**

#### Temperature dependence of creep-induced anisotropy in nanocrystalline FeCuNbSiB alloys

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We have investigated the creep induced anisotropy in nanocrystalline Fe<sub>75</sub>Cu<sub>1</sub>Nb<sub>2</sub>Si<sub>1</sub>B<sub>21</sub>, alloys as a function of the measuring temperature. The samples were produced by annealing originally amorphous ribbons for about 4 s between 500°C and 700°C with a tensile stress applied along the ribbon axis. For Si-contents equal or larger than 9 at% the annealed ribbons reveal a linear hysteresis loop indicative of an easy magnetic plane perpendicular to the stress axis. For lower Si-contents stress annealing results in a square loop indicative of an easy magnetic axis parallel to the stress axis. This transition from an easy plane to an easy axis as a function of the Sicontent is well-known from previous investigations and is related to the dependence of local magnetostriction of the nanocrystals on the Si-content. However, the investigations so far have been only carried out at room temperature. The present work, for the first time, reveals that the Si-contents where the easy axis changes its direction shifts to higher concentrations as the temperature increases. This is also nicely reflected in a reversible change from a linear hysteresis loop with low remanence to a square loop as a function of the measuring temperature.

#### **PM05**

Magnetocaloric properties of Ni-Co-Mn-In ribbon

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Temperature dependences of two parameters characterizing magnetocaloric effect (isothermal change in entropy and adiabatic change in temperature) were determined for a Ni<sub>47</sub>Co<sub>3</sub>Mn<sub>355</sub>In<sub>145</sub> ribbon. Two experimental methods were used, i.e., measuring temperature dependences of a specific heat (from 3 to 380 K) in zero magnetic field and in the field B=1 T, and measuring dependences of magnetization on temperature and on B. In the former method, both parameters were determined by analyzing the system entropy (determined by integrating the measured specific heat). In the latter method, the Maxwell relation was applied to get the isothermal change in entropy. The magnetocaloric effect shows two extrema of opposite sign, one near the transition between martensitic and austenitic phases. and the other one at the Curie temperature. The parameters characterizing the magnetocaloric effect near room temperature in relatively low B (1 T) make this material promising for application in magnetic refrigerators. The magnetocaloric effect was found to be highly anisotropic, being much larger for B applied in the ribbon plane. This property is particularly useful for designs, in which the active medium rotates in the field. This work was supported by the European Regional Development Fund. Innovative Economy Grants POIG 01 01 02-00-108/09 POIG 01 03 01-00-058/08 and POIG 02 02 00-00-025/09

#### PM06

#### Influence of bismuth substitution on the magnetocaloric properties of Gd<sub>5</sub>Si<sub>2</sub>Ge<sub>2</sub> compound

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In order to study the influence of relative small Bi addition on the magnetocaloric effect of Gd<sub>5</sub>Si<sub>2</sub>Ge<sub>2</sub>, a series of Gd,Si<sub>2</sub>, Ge<sub>2</sub>,Bi<sub>2</sub>, (2x=0; 0.02; 0.04; 0.06 and 0.08) alloys was prepared by arc melting method. Investigations of Gd<sub>5</sub>Si<sub>2x</sub>Ge<sub>2x</sub>Bi<sub>2x</sub> properties where carried out using X-ray diffraction, thermomagnetic analysis, magnetization isotherms measurements and differential scanning calorimetry (DSC). After a heat treatment at 1250°C for 200 min., all the alloys have, at room temperature, a monoclinic structural phase of  $Gd_5Si_2Ge_2$  type (S.G.: P112\_1/a) with a small amount of hexagonal structural phase of  $Gd_5Si_3$  type (S.G.: P112\_1/a) with a small amount of hexagonal structural phase of  $Gd_5Si_3$  type (S.G.: P112\_1/a) with a small amount of hexagonal structural phase of  $Gd_5Si_3$  type (S.G.: P112\_1/a) with a small amount of hexagonal structural phase of  $Gd_5Si_3$  type (S.G.: P112\_1/a) with a small amount of hexagonal structural phase of  $Gd_5Si_3$  type (S.G.: P112\_1/a) with a small amount of hexagonal structural phase of  $Gd_5Si_3$  type (S.G.: P112\_1/a) with a small amount of hexagonal structural phase of  $Gd_5Si_3$  type (S.G.: P112\_1/a) with a small amount of hexagonal structural phase of  $Gd_5Si_3$  type (S.G.: P112\_1/a) with a small amount of hexagonal structural phase of  $Gd_5Si_3$  type (S.G.: P112\_1/a) with a small amount of hexagonal structural phase of  $Gd_5Si_3$  type (S.G.: P112\_1/a) with a small amount of hexagonal structural phase of  $Gd_5Si_3$  type (S.G.: P112\_1/a) with a small amount of hexagonal structural phase of  $Gd_5Si_3$  type (S.G.: P112\_1/a) with a small amount of hexagonal structural phase of  $Gd_5Si_3$  type (S.G.: P112\_1/a) with a small amount of hexagonal structural phase of  $Gd_5Si_3$  type (S.G.: P112\_1/a) with a small amount of hexagonal structural phase of  $Gd_5Si_3$  type (S.G.: P112\_1/a) with a small amount of hexagonal structural phase of  $Gd_5Si_3$  type (S.G.: P112\_1/a) with a small amount of hexagonal structural phase of  $Gd_5Si_3$  type (S.G.: P112\_1/a) with a small amount of hexagonal structural phase of  $Gd_5Si_3$  type (S.G.: P112\_1/a) with a small amount of hexagonal structural phase of  $Gd_5Si_3$  type (S.G.: P112\_1/a) with a small amount of hexagonal structural phase of  $Gd_5Si_3$  type (S.G.: P112\_1/a) with a small amount of hexagonal structural phase of  $Gd_5Si_3$  type (S.G.: P112\_1/a) with a small amount of hexagonal structural phase of  $Gd_5Si_3$  type (S.G.: P112\_1/a) with a small amount of hexagonal structural phase of  $Gd_5Si_3$  type (S.G.: P112\_1/a) with a small amount of hexa P63/mcm). The Curie temperature is around ~ 285K (±2K). The results of the DSC measurements show a difference of ~10K in the peak of the heating and cooling curves associated with a first order structural phase transition. The peak slightly decreases with Bi content. This is expected to be related to the increasing of the Si/Ge bond distance in the monoclinic phase with Bi addition. The maximum isothermal magnetic entropy change decreases with Bi addition, however a giant magnetocaloric effect was obtained for all alloys. As an example, the maximum magnetic entropy change for the Gd<sub>5</sub>Si<sub>198</sub>Ge<sub>198</sub>Bi<sub>0.04</sub> was found to be 12.6 J/kg•K under an applied field of 5T.

# PM07

#### Ion incient energy effects on the giant magnetoimpedance enhancement

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The present experiment investigates the physical effects of ion irradiation on the magnetic properties of the amorphous ribbon. Samples were commercial amorphous ribbon with sizes of 2 mm x 40 mm x 20 um. The Ar ions were provided by an ion implanter at various dosages from 100 keV to 200 keV at the 1.0 x1018 ion/cm2. The GMI ratio was measured as a function of the DC magnetic field (Hmax=35 Oe) at different driving frequencies up to 10 MHz with a 5 mA amplitude applied to the sample irradiated with different energy. The irradiated samples show a clear change in the impedance response with applied field. When energetic ions with energies of several tens of keV penetrate into a thin film the collision cascade cause displacements of the target atoms that induce a uniaxial anisotropy in the sample. Ion irradiation induces uniaxial anisotropy along the transverse direction, which is accompanied by an effective transverse susceptibility. Therefore, the increase in the GMI ratio in an ionirradiated sample is attributed to the increase in associated uniaxial anisotropy created by ion bombardment. We found that controlling the ion energy can be one of the effective ways of controlling the GMI property for ion

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#### PM08

#### Investigation of electrodeposited FeNi film prepared from tartaric acid based bath

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FeNi films have been plated by using boric acid as pH buffer agent. However, the boric acid is needed to be replaced by another one which does not include boron, because boron is an environmentally hazardous substance. In order to achieve boronfree electroplating of FeNi, we applied the method proposed by Luisa et al for the CuNi electroplating [1] to that of FeNi, in which they successfully plated CuNi with good surface morphology by using tartaric acid. We electroplated FeNi films with Fe content of 17-45 at.% from the bath including NiSO4•6H2O, FeSO4•7H2O, NaCl, C7H4NNaO3S•2H2O and C4H6O6 (tartaric acid). The content of Fe was controlled by varying the amounts of tartaric acid between 0-200 g/L. The coercive force of the asdeposited films varied with the content of Fe, and showed a sharp bottom of 20 A/m at around 22 at.% of Fe. This result is consistent with the fact that FeNi alloys have nearly zero values of magnetocrystalline anisotropy and saturation magnetostriction at around 22 at.% of Fe. From the above result, we concluded that tartaric acid based bath film is one of hopeful plating bath for electroplating of FeNi films.

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#### PM09

#### Electrodeposited Fe-Co film prepared from citric acid-based plating bath

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Although a lot of the electrodeposited Fe-Co films with good soft magnetic properties were reported [1-3], boron, which is harmful and restricted by the regulation of environmental protection in Japan, was included in the plating bath as boric acid for controlling pH. Therefore, in order to remove boron, we focused on the citric acidbased plating bath and investigated effect of citric acid for the electrodeposited Fe-Co films. The plating bath was consisted of FeSO.•7H<sub>2</sub>O(100 g/L). CoSO.•7H<sub>2</sub>O(100 g/L) L), NaCl(50 g/L), C<sub>6</sub>H<sub>8</sub>O<sub>7</sub>•H<sub>2</sub>O(0 - 270 g/L).An Fe plate as anode and a Cu plate as cathode were used, and Fe-Co film was deposited on the Cu plate. Current density was kept at 2 mA/mm2 and bath temperature was set at 50 °C during the deposition. The coercivity of the deposited film was drastically reduced from 1700 to 200 A/m with increasing content of citric acid from 0 to 100 g/L, and then slightly increased The Fe-Co film prepared from the bath with citric acid of 100 g/L has high saturation magnetization of over 2.2 T and we can conclude that citric acid-based plating bath is one of hopeful plating ones to obtain good soft magnetic film with high saturation magnetization

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# **PM10**

#### Magnetic and magnetocaloric properties of polycrystalline Fe<sub>2</sub>P under hydrostatic pressure

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Magnetocaloric-based refrigeration presents itself as an environmentally-friendly energy-efficient technology capable of beneficially replacing the current gas compression-based technology. Some of the most promising magnetocaloric working materials found to date, such as (Fe,Mn)2(P,Si), are based on the Fe<sub>2</sub>P binary compound. Fe,P itself presents a first order magnetoelastic transition between FM and PM states at around 219 K and a low magnetic entropy change spanning a large temperature range. What makes Fe-P so unique is the coupling presented by its magnetic and crystal lattices which gives rise to the so-called mixed magnetism[1]. We believe that the mixed magnetism is responsible for the giant magnetic entropy change present in Fe<sub>2</sub>P-based compounds. To further advance our understanding of this coupling we performed magnetization experiments under hydrostatic pressures up to 8 kbar on pure Fe<sub>2</sub>P in its polycrystalline form. Pressure is found to decrease Tc at a rate of 6 K/kbar, about twice the rate reported in literature[2]. The Tc dependence on field is found to be non-linear and is only weakly affected by pressure. The magnetic entropy change neak is slightly decreased by increasing pressure whereas the total entropy increases by  $\sim 6\%$ 

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#### PM11

#### Influence of Er doping on magnetic and magnetocaloric properties of (NiCo),MnGa

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Recently, an attention has been paid to magnetocaloric properties of the Co-doped Heusler Ni2MnGa alloys with a paramagnetic gap in martensitic phase that can be tuned by their composition [1]. The strong influence of a rare-earth doping has been calculated and observed on the valence orbitals of Mn that makes the main contribution to the magnetic properties of these alloys [2]. We have prepared Co-doped Ni-MnGa allovs with a very small amount of Er. The polycrystalline sample of Ni429C07Mn31Ga19Er01 was checked by the EDX analysis and a segregation of Er was discovered on grain boundaries as it was expected in [2]. Saturated magnetization M(5K,0)=2.38µB/f.u. and the Curie temperatures TCM=318 K and TCA=440 K of martensite and austenite, respectively, were observed in Er free sample by a SQUID magnetometer. The inverse magnetocaloric effect at vicinity of martensitic transition, TM-A=379 K, was measured by direct method with ∆T=1.2K. In the Erdoped alloy, M(5K,0)=2.83µB/f.u. and paramagnetic gap was diminished in martensitic phase. TCA=450 K has been detected only and magnetic susceptibility obey nicely the Currie-Weis law with paramagnetic moment of mP=5.09 µB/f.u. above TCA. The effect of the Er-segregation on crystal texture and magnetic behaviour of the Co-doped Ni2MnGa alloys will be discussed.

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### PM12

# Structural and magnetic properties of FeMnAl nanocrystalline alloys

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Nanocrystalline magnetic alloys have been promising materials for magnetic device applications requiring high saturation magnetic flux density, high permeability, and low core loss. Among the magnetic alloys, Fe-Mn-Al system has been one of the most interesting tenary alloys with several magnetic phases. The magnetic properties were studied in one particular compound in previous works [1,2]. More systematic research of Fe-Mn-Al alloys would contribute to better understanding with Al contens. Nanocrystalline Fe90-xMn10Alx prepared via mechanical alloying using Fe, Mn, and Al powder with 48 hrs milling time. The structural and magnetic properties were studied. All peaks of XRD data are broader and shifted to smaller angle with increasing Al contents depending on the crystalline size and the lattice parametter, respectively. All the samples exhibit alloys behavior with an average crystallite size around 10 nm. The magnetic saturation, the coercivity, and the permeability, obtained from measurements by vibrating sample magnetometer showed strong dependence on the Al contents, and these dependences were related to the changes in structure and crystallite size. In addition Mossbauer experiments confirm tha Al diffuses into Fe bcc grains. Furthermore, these results suggest that, by adjusting the Al content, appropriate structural transformation and appropriate magnetization values can be obtained.

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# **PM13**

# Magnetocaloric effect in Fe doped La<sub>0.67</sub>Ba<sub>0.33</sub>MnO<sub>3</sub> system

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The perovskite-type oxides of  $La_{0.67}Ba_{0.33}Mn_{1\mbox{-}x}Fe_xO_3$  were synthesized by the conventional solid-state reaction method. The Curie temperature Tc and magnetic entropy change (MCE) in these samples are determined and compared to those of Gddoped systems [1,2] The Curie temperature Tc decreases from 350 K to 290 K with increasing Fe concentration x from 0 to 0.05. The relatively large MCE with a broad peak around Curie temperature compared with other oxides is observed in this system. These results suggest this system is a suitable candidate as magnetic refrigerator material working at room temperature.

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### **PM14**

Study of magnetic transition and magnetic entropy changes of La<sub>0.7</sub>Sr<sub>0.3</sub>MnO<sub>3</sub> and La<sub>0.7</sub>Ca<sub>0.1</sub>Sr<sub>0.2</sub>MnO<sub>3</sub>

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Magnetocaloric effect (MCE) and its most straightforward application, magnetic refrigeration, are becoming fields of increasing research interest. On the one hand, there are reasonable expectations that these subjects will give rise to energy-efficient, environmentally friendly technological applications. On the other hand, the study of some model materials gives some more insight into the physics underlying these phenomena. In this paper, complex magnetic materials La<sub>0.7</sub>Sr<sub>0.3</sub>MnO<sub>3</sub>, La<sub>0.7</sub>Ca<sub>0.3</sub>MnO<sub>3</sub> and La0,7Ca0.1Sr0.2MnO3, suitable for the magnetic refrigeration, has been investigated. We have found that the substitution of the Sr with Ca leads to an important decrease of the Curie temperature (TC) from 374 K to 345 K. The magnetocaloric study exposes a quite large value of the magnetic entropy changes for La<sub>0.7</sub>Ca<sub>0.3</sub>MnO<sub>3</sub> and La<sub>0.7</sub>Ca<sub>0.1</sub>Sr<sub>0.2</sub>MnO<sub>3</sub>. In cooling applications, another important parameter is relative cooling power (RCP), a good indicator of cooling efficiency of a magnetic refrigerant. For an applied magnetic field of 2 T, the RCP value is found to be 203 J/ kg for La07Ca01Sr02MnO3. As a result, the studied compounds could be considered as potential material for magnetic refrigeration near and above room temperature.

### PM15

#### Perpendicular magnetic anisotropy of amorphous ferromagnetic CoSiB multilaver

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Perpendicular magnetic anisotropy (PMA) is worthy of notice in both fields of science and industry due to its high integrate. In the previous study, we used CoSiB (75:15:10, atm%) and CoSiB/Pt multilaver showed the strong dependence on both the thicknesses of CoSiB and Pt. For the [CoSiB 3A/ Pt 14A]x5 multilaver, the maximum coercivity was ~ 224 Oe and the maximum PMA constant was obtained as ~  $2 \times 10^{6}$  erg/cm<sup>31</sup> In addition this film has the coercivity (Hc) and the saturation magnetization (Ms) with respect to 1.6 Oe and 407 emu/cm<sup>3</sup>. Considering this result, we suggest the new idea of CoSiB/Au/CoSiB sandwich structure and CoSiB/Pd multilaver. 1 mono-laverthick of Au and Pt are 144pm, 139pm and they are diamagnetic and ferromagnetic materials respectively CoSiB/Au/CoSiB sandwich structure and CoSiB/Pd multilaver is expected to show PMA and it can be developed to the magnetic random access memory (MRAM).

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# PM16

# Thickness dependence of magnetic induction in inhibitor-free 3% silicon steels

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Cold-rolled sheets of differing thicknesses were prepared through the conventional hot and cold rolling process. After annealing at 1200°C, final texture greatly varied with heating rate and sheet thickness. The 0.2 mm thick samples contained no or few {110} grains within the tested range of heating rates, yielding a poor magnetic induction. On the contrary, the 0.15 mm thick samples consisted entirely of {110} grains, and consequently resulted in a very high magnetic induction of over 1.95 Tesla. The 0.1 mm samples were also composed of {110} grains either at a very slow or at a very fast heating rate, but they unexpectedly exhibited a poor magnetic induction due to the deviated <001> direction of {110} grains from rolling direction. Orientation distribution functions of the samples taken from each cold rolling pass revealed that initial {111}<112> texture components started to move toward {111}<110> at a thickness of 0.13 mm. According to the nucleation and selective growth mechanism [1], primarily recrystallized texture takes after cold rolling texture, and hence the deviated final texture was obtained. This study indicates that there exists an optimum thickness for obtaining high magnetic induction and low iron loss.

[1] N. H. Heo et al., Acta Mater: 51 (2003), 4953-4964.

#### **PM17**

#### Effect of M (= Ge, Y, Hf) addition on soft magnetic properties of Fe-B-Si-M metallic glass alloys

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Fe-based metallic glass (MG) alloys exhibit good soft magnetic properties such as low coercivity and core loss but have the limits for industrial application because of low saturation flux density below 1.5T. Therefore, Fe-based MG alloys with high saturation flux density above 1.7T are very attractive soft magnetic alloy systems. In order to investigate the effect of M(= Ge, Y, Hf) addition on soft magnetic properties of Fe-B-Si-M metallic glass alloys, structural variations and magnetic properties of Fe $_{83,X}B_{10}Si_7M_X$  (X=1-5 at atomic percent) MG alloys were examined by X-ray diffraction(XRD) and vibration sample magnetometer(VSM) after fabricating Fe $_{83,X}B_{10}Si_7M_X$  MG ribbons by a single-roller melt-spinning technique under argon atmosphere. The thermal stability was measured by Differential Scanning Calorimetry(DSC) at a heating rate of 20K/min. Fe $_{83,X}B_{10}Si_7M_X$  MG alloys added Ge, Y or Hf exhibited a fully amorphous structure. Saturation flux density and coercive force of these alloys were Bs = 1.6~1.8T and Hc = 20~40 A/m, respectively. After annealing for 1 hr at 773K, Fe80B10Si7Ge3 MG alloy displayed excellent soft magnetic properties such as saturation flux density of 1.8T and corecivity of 20A/m.

#### PM18

#### Effects of film composition and substrate temperature on the structure and magnetic properties of FeCoB alloy films formed on MgO singlecrystal substrates

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Tri-layer films consisting of FeCoB and MgO layers have been investigated for MTJ applications. Theoretical studies show that the TMR ratio of MTJ prepared using epitaxial layers is one order of magnitude greater than that obtainable with amorphous layers [1,2]. The structure of bulk FeCoB alloy changes from crystalline to amorphous as the B content exceeds around 10 at %. However in a form of thin film, there is a possibility of crystallization for B rich compositions. In the present study,  $[Fe_{07}Co_{03}]_{100,x}B_x$  (x = 0-15 at %) films were sputter deposited on MgO substrates of (100), (110), and (111) orientations at different substrate temperatures to investigate the effects of film composition, temperature, and substrate orientation on the structure and magnetic properties. Epitaxial films were obtained for all compositions by adjusting substrate temperatures. For example,  $(FeCo)_{xg}B_x$  and  $(FeCo)_{xg}B_1$  epitaxial films were formed below these temperatures. The epitaxial films formed on MgO (100), (110), and (111) orientations consisted of bcc(100), bcc(211), and bcc(110) crystals, respectively. The magnetocrystalline anisotropy decreased with increasing the B content.

[1] W. H. Butler, X.-G. Zhang, T. C. Schulthess, and J. M. MacLaren: Phys. Rev. B., 63, 054416 (2001). [2] J Mathon and A Umerski: Phys. Rev. B., 63, 220403 (2001).

### PM19

# Arrangement of different magnetic alloy sheets for effective magnetic shielding

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The purpose of this study is to provide an effective shielding method for power frequency magnetic field emanating from a neutral ground reactor (NGR). The 0.35 mm thick commercial grade grain-oriented electrical steel (GO), non-oriented electrical steel (NGO), and permalloy (PC) were used as shielding materials. The area of 1200x1200 mm2 was blocked with the alloy sheets at a place 600 mm away from the NGR. In a weak magnetic field, PC is the best in shielding performance, whereas GO is the best in a strong magnetic field. NGO is worse than GO, but is better than PC in a very strong magnetic field. For multi-layered shields, it was observed that GO/PC pair (GO is close to source) was most effective, yielding a shielding factor less than 0.06 in the wide range of magnetic fields. These results are explained by variation of shielding effectively shields a strong magnetic field first, and then the outer PC sheet effectively shields the weakened field. Furthermore, this study demonstrates some spacing between laminations is essential to achieve best shielding performance in the 'shunt' mcchanism of magnetic shielding.

#### PM20

#### Fabrication of pariculated Fe-Mg thin films by selective oxidation and their magnetic properties Pyungwoo Jang\*

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It is impossible to manufacture Fe-Mg bulk alloys even in vacuum melting system because of very low melting temperature, high affinity with oxygen, and very high vapor pressure of Magnesium. However, it is very easy to fabricate Fe-Mg thin films by sputtering. Becasue of big difference of enthalpy of Fe-O and MgO, Mg can be selectively oxidized in a mixture atmosphere of hydrogen and water vapor. 10 - 200 nm thick sputtered Fe-Mg films were selectively oxidized in a damp atmosphere and annealed at 800-900oC for 10 - 200 min. With increasing oxidation time, saturatio magnetizatin of Fe-Mg films increased due to selective oxidation of Mg. Furthermore, Initial permeability as well as the slope of magnetizatio loops of the films decreased, which was attributed from the facts that morphology of the films were changed from continuous type to pariculated type. when the films were annealed in pure hyderogen atmosphere the rate of change is very slow. The reason for the expedited change is probably due to energy transfer to Fe atoms during migration of Mg atoms to film surface

#### PM21

# Optimum spacing of magnetic alloy strips in open-type magnetic shielding

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This study suggests an effective shielding method for power frequency magnetic field in open-type magnetic shielding, which secures the clear view of field sources by aligning magnetic alloy strips sparsely. A neutral ground reactor (NGR) was used as a field source. The 0.35 mm thick commercial grade grain-oriented electrical steel (GO) and permalloy (PC) were cut into a size of 50x1200 mm2 and then placed vertically 600 mm away from the center of the NGR. In a weak magnetic field, PC was the best in shieldng performance, whereas GO was the best in a strong magnetic field due to change in hierachy of material permeability with increasing magnetic field strength. In case of multi-layered shields, it was observed that combination of GO/PC/ GO was more effective than the triple-layered shield of the same material. However, closely conjoined GO/PC/GO shield showed worse shielding performance than the same layered shield seperated from each other, which demonstrates it acts as one body regardless of presence of thin surface insulation layers. This study reports that introducing air gap between laminations of each layered shield is important to obtain desired shieding performances and optimum spacing between alloy strips exits in opentype magnetic shielding

# PM22

#### Magnetocaloric effect in Ni2.27Mn0.73Ga Heusler alloy

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Drobosyuk Mikhail, Buchelnikov Vasiliy, Taskaev Sergey, Fayzullin Rafael Chelyabinsk State University, Chelyabinsk, Russia m.syuk@mail.ru type of presentation: poster There are many theoretical and experimental works devoted a possibility any magnetic materials to change its temperature after applying an external magnetic field. This effect is called the magnetocaloric effect (MCE). Recent experimental studies have shown that Ni-Mn-Ga Heusler alloys have unique properties such as the shape memory effect, the large magnetostriction, the large magnetoresistance and other magnetic properties. In this work we present the experimental studies of the MCE for Ni<sub>2,27</sub>Mn<sub>0.75</sub>Ga Heusler alloy. Polycrystalline ingot was prepared by a conventional arc-melting method. Samples for measurements were performed by the Magnetocaloric Measuring Setup produced by AMT&C. The results of measurements will be shown on the poster.

#### PM23

#### Soft magnetic properties of Fe-6.5wt.%Si alloy sheets fabricated by powder hot-rolling

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It is widely known that Fe-6.5 wt.% Si alloys have excellent soft magnetic properties. In this work, Fe-6.5 wt.% Si sheets were rolled by powder hot-rolling process without cracking and soft magnetic properties were investigated with microstructural defects. The cylindrical preforms were prepared by inert gas atomization and subsequent spark plasma sintering of Fe-6.5 wt.% Si powders at 1223K. The Fe-6.5 wt.% Si sheets were fabricated by hot rolling of copper canned preforms at 1173K with total reduction of 70%, 80% and 90%, respectively. The hot rolled sheets were annealed at 1073K in an argon atmosphere. Microstructures and magnetic properties (magnetic flux density, coercivity and core loss) of the hot rolled and annealed Fe-6.5 wt.% Si sheets were examined by scanning electron microscopy (SEM), vibration sample magnetometer (VSM) and B-H analyzer. The hot rolled Fe-6.5 wt.% Si sheet with reduction of 90% and subsequent annealing exhibited soft magnetic properties such as magnetic flux density of 1.8 T and coercivity of 20 A/m. This result is due to the elimination of microstructural defects and the increase of the average grain size and density.

### PM24

#### Structure and magnetic properties of nano/micro-sized Mn-Al alloy powders produced by plasma arc-discharge and gas atomization Junggoo Lee<sup>1\*</sup>, Younkyoung Baek<sup>1</sup>, Hwijun Kim<sup>2</sup> and Chuljin Choi<sup>1</sup> <sup>1</sup> Powder & Ceramics Division, Korea Institute of Materials Science, Korea

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The τ-phase in the Mn-AI system has attracted great attention due to the remarkable magnetic properties superior to the Alnicos and hard ferrites, low cost, as well as good machinability and corrosion resistance. Due to these reason, several methods have been employed to produce τ-phase Mn-Al alloy powder until now. On the other hand, plasma arc-discharge and gas-atomization method has been widely used for the production of nano-sized and micro-sized powders, respectively. To our best knowledge, however, there is no comparison study of Mn-Al alloy powder produced by these two methods. In the present study, Mn-Al alloy powders were prepared by the methods of plasma arc-discharge and gas-atomization and influence of process parameters on composition, particle size, and magnetic properties of the powders was systematically investigated. Mn-Al powder produced by plasma arc-discharge was smaller than its single magnetic domain and exhibited the coercivity of 5.6 kOe, which is the highest among the values reported until now. It was also confirmed that the magnetic properties of micro-sized powder produced by gas-atomization have been improved by subsequent treatments such as annealing and mcchanical milling Based upon these results, the magnetic properties of thus two types of Mn-Al alloy powders will be discussed.

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# PM25

# Ultra high speed pm type synchronous motor-generator with amorphous core for micro turbine

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# PM26

# Temperature dependence of magnetic domains in grain-oriented silicon steel

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Abstract: Silicon steels are essential for transformers and the reduction of loss in them is of utmost importance in this green age. Since the loss in steels are closely related to their magnetic domain structures, the magnetic domain structures in turn are of exceeding interest. We studied magnetic domains of 3% silicon steel single crystals with (110) [001] texture as a function of temperature via magnetooptical Kerr effect (MOKE) microscopy. Generally, the silicon steels (110) [001] have 180° main domains oriented along the [001] axis, and they also have complimentary Lancet domains in case of misorientation. [1, 2, 3] The MOKE images shown in Figure 1 clearly reveal that the domain patterns change as temperature is changed. The width of the 180° main domains, for instance, decreases with increasing temperature. In order to go one step further from mere qualitative observations, we analyzed the MOKE images quantitatively, and, in particular, investigated temperature dependence of the 180° domain width and the Lancet domain density. These temperature characteristics will be interpreted from the standpoint of total energy minimization.

 [1] A. Hubert, R.Schaefer, Magnetic domains(1998), Springer [2] Y. Ushigami, J. Magn. Magn. Mater. 254-255(2003), 307-314 [3] M. Imamura, Physica Scripta. T 24 (1988)

# PM27

# A tri-layer stress impedance sensor using amorphous magnetostrictive thin film

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Mechanical stress sensors have a wide area of applications in strain and torque measurements. Thin film stress impedance (SI) sensors have several advantages over other stress sensors such as high sensitivity in small stress range, miniaturization possibility as well as the integration capability with other devices. Stress sensors that utilize amorphous magnetostrictive thin films have been proposed previously. The permeability change of the magnetostrictive material upon the applied external stress was evaluated throughout an impedance measurement. In this work, a tri-layer thin film stress sensor is fabricated and analyzed. The sensor has a Fe40Ni38Mo4B18/ Ti/ Fe40Ni38Mo4B18 tri-layer structure with 900 nm thickness each and dimensions of 10×0.1 mm2. After the deposition, the sensor is annealed at 300 °C with a magnetic field of 500 Oe applied in the transverse direction to induce an in-plane transverse anisotropy. The magnetic behavior is studied with a vibrating sample magnetometer and the crystallographic structure is studied with using X-ray diffraction. The sensor is tested by applying compressive and tensile stresses and the sensor response was obtained from an impedance analyzer. The results show high sensitivity at strains between 0-80 ppm with a maximum impedance change of 16% at 50 MHz

Frommberger, M.; Glasmachers, S.; Schmutz, C.; McCord, J.; Quandt, E., IEEE Trans. Magn., vol. 41, pp.3691 - 3693, 2005

# PM28

Magnetoimpedance(GMI) effect in the NiFe shell/Cu core wires fabricated by electrodeposition

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Microwires with NiFe shell on Cu core of 90  $\mu m$  diameter were fabricated by electrodeposition method. The thickness of NiFe shell was well controlled by deposition time in the range of 4  $\mu m$  ~18  $\mu m$ . Giant magnetoimpedance(GMI) effect in the NiFe/Cu core shell wire was measured in frequency range of 10 kHz ~ 10 MHz and in the magnetic field range of -300 Oe ~ 300 Oe. The variations GMI ratio with frequency and NiFe shell thickness were analyzed in terms of magnetic relaxation and skin depth at ac field. We obtained maximum GMI ratio of 250 % at 300 kHz in the wire with NiFe shell of 11.5  $\mu m$  thickness.

# PN01

# SQUID, XRD and Raman studies of Mn implanted gallium nitride at elevated temperature

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The magnetic and structural properties of GaN samples implanted with 325 keV  $Mn^{2+}$  ions at 350 0C substrate temperature for various fluences varying from 1.75 x 1015 to 2.0 x 1016 ions cm-2 were studied using SQUID, XRD and Raman scattering techniques. Magnetic properties were found to vary with the ion fluence. The 5 MeV Si<sup>-</sup> ion irradiation with fluence of 1 x 1016 ions cm-2 was performed on the sample implanted for ion fluence 2 x 1016 cm-2. The Curie temperature estimated for as-implanted and after irradiated with 5 MeV Si<sup>++</sup> ion samples was found to 301 and 326 K respectively. XRD spectra showed GaN peak at (002) reflection, this peak shifted toward lower angle indicated the incorporation of Mn ion in GaN film.Raman spectra of the samples showed bands at ~300 and 670 cm-1 were attributed to the local vibration modes of gallium and nitrogen vacancies related defects respectively.

#### **PN02**

#### Oxidation of monovacancies in graphene by oxygen

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We study the oxidation of monovacancies in graphene by oxygen molecules using first principles calculations. In particular, we address the local magnetic moments which develop at monovacancies and show that they remain intact when a molecule is adsorbed such that the dangling carbon bonds are not fully saturated. The observed value of magnetic moment is 1.35 Bohr Magneton for monovacancy and it becomes 1.86 Bohr Magneton by oxygen adsorption on monovacancy in graphene. However, the lowest energy configuration does not maintain dangling bonds and is found to be semiconducting. Our data can explain the experimentally observed behavior of graphene under exposure to an oxygen plasma.

### PN03

# Structural , compositional and magnetic study of bulk Fe doped ZnO system and impurity phase formation

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Structural, compositional, optical, vibrational and magnetic properties have been studied for (ZnO), \_([Fe,O\_3), \le 0.10 series of bulk samples. The ZnO based ceramic samples with different doping percentage of Fe,O\_3 were prepared by ball milling process. The phase composition of the pellets was determined via X-ray diffraction analysis. Raman spectroscopy give information about the additional modes appear in the ZnO spectrum having Fe doping which assures the presence of Fe in the ZnO matrix. The de-convolution of XPS spectra of core peaks of different elements indicated the presence of different bond breakage due to Fe substitution in ZnO lattice. The appearance of shaking satellites in XPS spectra also confirmed the presence of different valance states of Fe ions. The red shift in energy band gap estimated from reflectance UV-VIS spectroscopy was observed for all bulk samples. The magnetic behavior of the samples was examined by using Vibrating Sample Magnetometer (VSM) indicating ferromagnetic behavior aroom temperature. The effective magnetic behavior is due to substitution of Fe ions in the ZnO detice.

#### PN04

# Enhanced magnetization and spin injection in Co/ZnO films by Al doping

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An important hurdle in the development of semiconductor spintronic devices is the inefficient injection of spin polarized currents from metallic ferromagnets into semiconductors due to the large conductivity mismatch. Inserting a magnetic semiconductor at the interface was regarded as an effectively route for spin injection into semiconductors. In this work, we investigate the magnetism and electrical spin injection into ZnO and Al-doped ZnO (Al 2at.%) matrix from Co particles in Co/ ZnO and Co/ZnAlO films, which were prepared by depositing ultrathin Co layers and semiconductor layers at room temperature. The films consist of Co particles dispersed in the semiconductor matrix. The graded magnetic semiconductors were formed at the transition region between metallic Co particles and semiconductor matrix with the substitution of Zn2+ with Co2+ ions. The increasing of magnetization of Co/ ZnAlO films compared to the corresponding Co/ZnO films is probably due to the enhanced magnetization of graded magnetic semiconductors with Al enhancing carrier concentration. The room temperature magnetoresistance of Co/ZnAlO film with 1.0 nm of ZnAlO layer reaches -12.3% which is higher than -8.4% of the corresponding film without Al, probably due to the large spin filter effect occurring in the graded magnetic semiconductors having a relative lager polarization.

#### PN05

# Structure, magnetic, and transport properties in $Cu_{1-x}Mn_xO$ compounds

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By annealing at air,  $O_2$ ,  $N_2$  atmosphere, the polycrystalline  $Cu_{1-x}Mn_xO$  powders, which are prepared by sol-gel method with the concentration range of  $0 \le x \le 0.2$ , are synthesized. The phase component and atomic structure of the main phase are obtained by X-ray diffraction (XRD) with Rietveld refinements. The bond length of Cu(Mn)-O decreases, and the bond-angle of Cu(Mn)-O-Cu(O) increases, with increasing the dopant. Magnetism measurements (SQUID) show that the ferromagnetic transition temperature, Tc, and the saturation moment, Ms, increase with the doped concentration of Mn, although all the samples show that the Tc is less than 100 K. All samples show the well defined thermal active mechanism with the thermal active energy varying with the annealing atmosphere. Results reveal that annealed in  $O_2$  may increase the carrier concentration of the prepared samples while annealed in N. the reverse.

# PN06

#### A probe into the structural, magnetic and dielectric properties of barium and lithium substituted pseudobrookites

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The versatile nature of the Fe-Ti oxides makes them attractive candidates in applications in which their coupled semiconductor, magnetic and dielectric properties can be exploited and they appear to be good candidates for emerging technologies such as spintronics magneto-electronics and radar electronics. In the present work pure pseudobrookite Fe-TiO, and its lithium and barium substituted varieties namely  $Ba_{2v} Fe_{2x} Li_{3x}Ti_{1-v}O_5$  and  $Ba_{3v/2} Fe_{2-x} Li_{4v}Ti_{1-v}O_5$  (x = y = 0.05) are synthesized by the solid-state reaction method. Powder x-ray diffraction studies show the single phase orthorhombic structure. Substitution has reduced the crystallite size but enhanced the grain size. A new intense plane (122) is observed in the barium-substituted sample. Lithium intercalation is also observed. Room temperature dielectric studies have confirmed the Maxwell- Wagner interfacial polarization and space charge is seen to have almost doubled with doping. The magnetic hysteresis loops show an asymmetric shift and constriction of loops. Variation of normalised susceptibility with absolute temperature is plotted. Presence of barium and lithium have also brought about a long range magnetic ordering in the pseudobrookite. Susceptibility measurements also show that the Curie temperature in the barium-rich pseudobrookite approaches room temperature making it a potential switch material.

#### PN07

#### Oxygen vacancy and magnetism of a room temperature ferromagnet Co-doped TiO<sub>2</sub>

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We report the magnetic and local structural study of a room temperature ferromagnetic Co-doped TiO<sub>2</sub> prepared by the solid state reaction. The ferromagnet with 5% Co shows a rutile-type structure and has a saturation magnetization of 0.005 JT-1kg-1 and a coercive force of 0.02 T. We have found that it has some oxygen vacancy sites around a Co atom only from the extended x-ray absorption fine structure (EXAFS) measurement. The charge compensation makes the oxygen vacancy around Co<sup>2+</sup> and/ or Co<sup>3+</sup> substituted with Ti<sup>4+</sup>.

### PN08

# Effect of annealing on the magnetic anisotropy of GaMnAs film with low Mn concentration

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We have investigated the effect of annealing on the magnetic anisotropy of a GaMnAs film with a low Mn composition of 2 %. Three pieces of samples were cleaved from the sample and two of them we are annealed at 200 and 300 °C, respectively, for 3 hours in air. In case of as-grown sample, abrupt transitions in the Hall resistance appeared with a negative coercive field, when field strength was swept at a fixed direction in the film plan. The phenomenon turned out to be related to the presence of magnetic domains with a vertical easy axis in the film, which was identified via the Hall measurement that was performed with a field direction that was normal to the plane. The portion of domain with dominant out-of-plane magnetic anisotropy in the GaMnAs film was systematically decreased with an increasing annealing temperature. Especially, the negative coercive field disappeared in the sample annealed at 300 °C and two abrupt positive switching fields we are observed in the field scan of the planar Hall resistance. This indicates that the annealing significantly affects the magnetic anisotropy of GaMnAs film by changing from out-of-plane dominant to in-plane dominant anisotropy.

# **PN09**

# Low field magnetization reversal behavior of ferromagnetic GaMnAs film

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We have investigated magnetization reversal process of a GaMnAs film by using planar Hall effect. The angle-dependent measurements of the planar Hall resistance (PHR) showed sinusoidal behavior without hysteresis between clockwise and counter clockwise rotation with strong field. However, the magnetization reversal behavior changes significantly by showing abrupt transition at the <110> directions as the field strength reduced to 200 Oe. Furthermore, when a smaller field is used, the PHR displays four stable values arising from the formation of four different multidomain states during the rotation of the external field direction over 360°. Two of these states showing maximum and minimum values of the PHR, correspond to the sample fully magnetized along one of the easy axes. Other two intermediate states are multi-domain states in which a fraction of the domains populates one easy [100] direction, and the remaining fraction populates an orthogonal easy [010] direction. We showed that the relative populations of the magnetic domains corresponding to the two orthogonal easy axes can be controlled by the value of the applied field during the process of magnetization reversal. This phenomenon was understood by considering the differences in domainpinning fields and their distributions required for crossing the two hard axes.

#### **PN10**

Half-metallic antiferromagnetism in the ordered  $Cr_{1\!\!-\!x}Ca_xSb$  alloy from first-principles calculations

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Compared to half-metallic ferromagnets, half-metallic antiferromagnets (precisely called half-metallic fully compensated ferrimagnets) are more promising candidates for spintronic applications since their zero magnetization leads to lower stray fields and thus tiny energy losses. Using the first-principles calculations, we have systematically investigated the electronic and magnetic properties of the ordered Cr<sub>1,x</sub>Ca<sub>x</sub>Sb alloy. It is found that Cr<sub>1,x</sub>Ca<sub>x</sub>Sb with x = 0.125, 0.25, 0.5 and 0.75 all are half-metals like zinc-blende CrSb and CaSb which have been fabricated experimentally or predicted theoretically. Interestingly, Cr<sub>0.25</sub>Ca<sub>0.75</sub>Sb is a half-metallic antiferromagnet with complete spin polarization, and the half-metallic antiferromagnetism is robust against the lattice compression and expansion and the choice of electronic exchange and correlation functional.

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#### PN11

#### Translation & rotation of diamagnetic material induced by a low field of a permanent magnet and terminal identification of a micron-sized narticle

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It has been generally believed that field-induced motion of an ordinary diamagnetic material can be induced only at high field above several Tesla. However, translations were recently observed on various diamagnetic crystals in a direction of monotonously decreasing field produced by a permanent magnet [1]; the crystals were released in diffused micro-gravity area. In a given field distribution, terminal velocity of translation is uniquely determined by intrinsic susceptibility of material; the velocity is independent to mass of particle. Rotational oscillation, caused by diamagnetic anisotropy energy, was also reported for the crystals [1]. When the above motions are observable, diamagnetic susceptibility & anisotropy of a small particle can be detected, irrespective of its size, since the measurement of sample. Material identification of particle is easily performed by comparing the obtained susceptibility (& anisotropy) with published values. A simple and nondestructive method to identify the material of a small particle is desired in various fields of nano-science. By introducing a compact NdFeB magnet, the measurement was realized in an ordinary laboratory using a chamber-type drop shaft.

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# **PN12**

Magnetoresistance effect in electron-injected p-type silicon Michael P. Delmo\*, Eiji Shikoh, Teruya Shinjo and Masashi Shiraishi Graduate School of Engineering Science, Osaka University, Japan

Magnetoresistance (MR) effect in silicon has gained a renewed interest in recent years because it is large and non-saturating [1-3]. It is known that the presence of "static" inhomogeneity, like phosphorus in low-doped silicon, generates large and non-saturating MR effect [2,4]. Inhomogeneity can also be introduced "dynamically" by injecting electrons or holes into a low-doped silicon [1,3]. In n-type silicon (n-Si), electron injection generates space-charge effect [1], whereas, the hole injection induces the formation of a hole-electron (p-n) boundary [3]. These dynamic inhomogeneities induce large MR effect even up to room temperature. Most of the silicon devices used in inhomogeneity-induced MR effect studies are n-type, and no study has ever been reported for p-type, so far. In this study, we investigated the MR effect in indium (In)/ p-type silicon (p-Si)/In devices. Current (I)-voltage (V) characteristic of the device shows an Ohmic and non-Ohmic I-V regime (I is proportional to the square root of V), which indicates electron space-charge injection into p-type silicon [5]. MR of the In/p-Si/In device is small (10%) compared to that of In/n-Si/In device (1,000%) for the same condition (3 T, 300 K). This suggests that the electron-hole compensation or recombination suppress inhomogeneity formation in the silicon device.

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# PN13

# First-principles investigation of the influence of adsorbed atom on the defect and impurity substitute graphene

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Graphene is well-known to be two-dimensional material made of carbon atoms. Graphene is a substance that attracts attention not only as parts of the nanocarbons but also for its own interesting electronic and mechanic properties. It is important for the application that charge and magnetic are controlled by defect and impurity in the graphene. Furthermore, influence of defect and the impurity are important in the growth of the compound on the graphene. The technological uses of controllable charge and magnetic carbon systems are extensive, and it is important that the mechanism is known. In previous study, we investigated the adsorption energies and magnetism in adsorption sites on graphene from atomic number 1 to 83, using the DFT. We discussed the charge transfer and the magnetism. When several adatoms is adsorbed, the defect graphene develops magnetic. In this study, we performed atom substitute in the graphene and atomic adsorption on the graphene with the defect. The aim of the research is discussion about the magnetic and stability when several adatoms adsorbed on defects and impurities substitue in graphene. In addition, we discuss the growth mechanism of the nitrided compounds. We used a band calculation technique based on the DFT calculation.

### PN14

# Room-temperature fabrication of highly transparent magnetic nanocomposite systems by aerosol deposition

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Nanocomposite systems, materials containing particles of nanometer dimensions have shown interesting properties related to its extremely small size. Some of their optical, magnetic, electronic, mechanical, and chemical properties are different from those exhibited by the same composition in bulk material[1]. We have proposed novel concept for materializing nanocomposite systems using aerosol deposition (AD) process[2-3]. Our concept can apply various sizes of nanoparticles, and achieves desirable distribution of nanoparticles in host matrix because the structure of deposited layer is similar to that of composite powder. In this report, we developed highly transparent magnetic nanocomposite thick films consisting of cobalt nanoparticles embedded in a host matrix of ferroelectric lead zirconate titanate (PZT) by AD process at room temperature. Transparent thick nanocomposite toolalt/PZT films display very dense without any pores as well as any cracks. As the wt% of nanocobalt increased, magneto-optic effect measured in the transmission mode, especially Faraday rotation effect of nanocomposite film acquired the Faraday rotation angle of approximately 0.2 degree. These results would strongly show potential of optical multiferroics on transparent nanocomposite systems.

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# PN15

#### Effect of transition metal (Co, Ni and Cu) doping on lattice volume, band gap, morphology and saturation magnetization of ZnO nanostructures

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The interest in transition-metal (TM)-doped ZnO has increased because of promising applications in the field of semiconductor spintronics, which seeks to extend the properties and applications of established electronic devices by using the spin of electrons in addition to their charge. In this work, the effect of different TM (Co, Ni and Cu) doping on structural, optical and magnetic properties of ZnO nanostructures have been studied. Zn, TM,O (TM=Co, Ni and Cu) nanostructures were prepared by a microwave assisted chemical route and characterized by X-ray diffraction (XRD), field emission scanning electron microscopy (FESEM), transmission electron microscopy (TEM), Raman spectroscopy, UV-Vis and magnetization measurements. The X-ray diffraction (XRD) and TEM analysis showed that the TM-doped nanostructures had single phase nature with the wurtzite structure. There is a strong correlation between changes in the lattice parameters, bandgap energy, morphology and the saturation magnetization of Zn, TM O nanostrutures. Lattice volume and bandgap determined from XRD and UV-Vis, respectively, were found to decrease as the atomic number of the dopant moved away from Co. Magnetic studies showed that all the TM-doped ZnO exhibit room temperature ferromagnetism and the decreasing trend of saturation magnetization was observed with the increase of 3d electrons number from Co to Cu

#### PN16

# Functionalized graphene as a room-temperature ferromagnetic semiconductor

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Low-temperature magneto-transport and vibrating sample magnetometry (VSM) combined with superconducting quantum interference device (SQUID) measurements indicated that graphene, when functionalized with aryl radicals, behaved as a room-temperature ferromagnet and therefore might form the basis for a new approach to semiconducting ferromagnetism. After radical functionalization, epitaxially grown graphene samples became semiconductors and at the same time displayed high-density magnetism, with a room temperature saturation magnetization of approximately 0.1 Bohr magnetons per carbon atom. The bulk measurements with VSM and SQUID were corroborated with local studies using atomic force microscopy (AFM), electrostatic force microscopy (EFM), magnetic force microscopy (MFM), scanning tunneling microscoy (STM), and Raman Spectra measurements. Both local and bulk measurements indicated a strong dependence on the degree of coverage of the graphene surface with functionalized sites. The measurements showed the presence of a superparamagnetic order in the graphene regions with a relatively low-density of functionalized sites in the entire temperature range from 2 to 300 K. In contrast, the regions with a relatively high-density of functionalized sites displayed a ferromagnetic order even at room temperature, in summary, the measurements indicated that anyl radical functionalied multi-layered graphene os Silicon Carbide could simultaneously be a semiconductor and a ferromagnet at room temperature.

#### PN17

#### Tuning of the curie temperature by varying reduction potential in electrochemically prepared thin films of Prussian blue analogue based molecular magnets

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The quest for molecular magnets with ordering temperature at or above room temperature is highly desirable for their practical applications at ambient conditions. In order to increase the ordering temperatures, many efforts have been devoted to synthesizing compounds with appropriate choice of organic/inorganic ligands, their valence modulation and spin centers, and alkali metal ion doping, etc., using organometallic and/or coordination chemical methods. However, we have shown that, the magnetic ordering temperature in PBAs (Prussian blue analogues) can be enhanced by varying reduction potential in electrodeposition method. In this regard, the electrochemically prepared films of PBAs, XjFeIIk[CrIII(CN)6]l.mH2O with varying reduction potential are investigated using dc magnetization measurements. The magnetization data for film deposited at lower reduction voltage shows ferromagnetic ordering with TC (Curie temperature) of ~ 21 K, similar to previously reported ordering temperature for this compound. However, an increase in TC up to ~50 K has been observed for film deposited at higher reduction potential. The observed variation in TC is mainly attributed to the inclusion of potassium ion (K+) at higher reduction potential, which changes stoichiometry of the film. The ability of tuning TC just by changing reduction potential could be useful for designing future device applications.

# PN18

# Evolution of multifunctional behavior in site specific cation substituted $Na_{0.5}Bi_{0.45}Gd_{0.05}Ti_{0.95}Mn_{0.05}O_3$ ceramics

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A novel multifunctional behavior was observed by site specific cation substitution in NBT ceramics prepared by conventional solid state reaction technique. X-ray diffraction studies revealed the co-existence of orthorhombic reflections in co-substituted NBT system along with the primary R3c phase. Raman spectroscopy of NBT-GM ceramics reveals a substantial change in all phonon modes in terms of suppression and shift indicates that the substitution affects nature of the cation displacement, octahedral tilt and hybridization. Reduced polarization value in NBT-GM ceramics with non-saturated hysteresis loops gives evidence for the suppression of ferroelectricity. Temperature dependent magnetization measurements from R.T to 2.5K in NBT-GM ceramics show the occurrence of spin-paramagnetic behavior. Susceptibility plots show some evidence for weak interaction between Mn4+ ions at low temperatures below 30K. UV-Vis spectroscopy measurements reveal the indirect transition allowed band gap nature for both NBT and NBT-GM ceramics with a slight increase in band gap (Eg) value for NBT-GM ceramics. Photoluminescence measurements show a strong blue emission at ~390nm (3.18 eV) in Gd-Mn substituted NBT system, which corresponds to radiative transition from  ${}^{4}T_{1e} \rightarrow {}^{4}A_{2e}$  energy levels of Mn<sup>4+</sup>-ions.

#### PN19

#### Magnetic doping effect on physical properties of PbPdO<sub>2</sub> Kyujoon Lee, Seongmin Choo and Myung-hwa Jung\*

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The exotic features of gapless semiconductors have attracted many researchers in solid state physics. Materials such as HgCdTe have been studied as a gapless semiconductor. Recently PbPdO<sub>2</sub> has been theoretically proposed to be a gapless semiconductor. In addition, by substituting Co for Pd ions the material becomes a spin gapless semiconductor which has full spin polarization. In this study we have experimentally shown the changes in the physical properties of PbPdO<sub>2</sub> by doping different magnetic ions such as Co, Mn, and Zn. PbPdO, shows a metal-insulator-like transition at TMI=100 K in the resistivity vs. temperature measurements. The magnetic properties show a diamagnetic behavior at high temperatures and a ferromagnetic behavior at low temperatures. The TMI increases to 150 K by Co doping and the diamagnetic behavior changes to paramagnetic behavior. By Mn doping, the TMI decreases to 73 K and the magnetic behavior changes drastically to show antiferromagnetic ordering at low temperatures. However no difference has been found by Zn doping. The common feature for all materials is that the charge carrier density increases with doping. The easily tuned properties of PbPdO2 show evidence that PbPdO2 is a gapless semiconductor and could be tuned to be a spin gapless semiconductor.

#### PN20

# An experimental approach using EPR and XMCD to explore hydrogen mediated ferromagnetism

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Since the ferromagnetism of III-VI based DMS materials such as GaMnAs was turned out to be based on carrier induced model, much progress has been made in developing its applicability but low operation temperature has limited the approach to functional devices[1]. Among other materials, Co doped ZnO has been regarded as strong candidate for room temperature DMS but the origin of ferromagnetism has been on controversy[2]. Previously, we reported experimental results on Co-H-Co responsible for ferromagnetism depending on hydrogen contents and position[3-5]. Here, in order to apply these phenomena for new functional devices, more convincing experimental results on Co-H-Co is necessary. In this study, we show the experimental approaches to hydrogen mediated ferromagnetic spin ordering in ZnCoO through electron paramagnetic resonance (EPR) and x-ray magnetic circular dichroism(XMCD). The signal intensity of ferromagnetic a ad spectrum of magnetic circular dichroism within UV-Vis range suggests the possibility for creating and manipulating spin-polarized carrier using hydrogenated ZnCoO. Through advanced theoretical calculation for Co-H-Co model, we are expecting to obtain the direct evidence for Co-H-Co. Through advanced theoretical calculation for Co-H-Co unit.

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# **PN21**

#### A role of mobility in hydrogen mediated ferromagnetism of ZnCoO

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Diluted magnetic semiconductors (DMSs) have recently been attracting much attention for their potential as spin manipulation layers [1]. Although there have been reported several authorized DMSs, such as Mn doped GaAs, their limited operating temperature requires new DMS materials with high Curie temperature. Up to now, many theoretical and experimental studies about hydrogen mediated ferromagnetism in ZnCoO have been introduced [2-3]; however, the interaction between such magnetic units as Co-H-Co complexes has not yet been confirmed. Delocalized electron in ZnCoO lattice is possibly considered to play a crucial role in the correlation [4]. In this work, we investigated the correlation between carrier mobility and ferromagnetic spin ordering in ZnCoO. Electron mobility was artificially controlled with the manipulation of crystallinity and hydrogen contents. Hydrogen treatment is carried out using hot isostatic pressing (HIP) system. The correlation between carrier mobility and ferromagnetic spin ordering was analyzed in terms of Co-H-Co magnetic units. Mobility-dependent ferromagnetic spin ordering is discussed in our experiment.

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### PN22

# Magnetic and magneto-optical properties of TiO<sub>2</sub>:V semiconductor oxide films with various resistivity

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We present the investigation results of semiconducting films TiO2.V with 3-18% at. of vanadium, which develop our previous research of doped itianium oxide [1,2]. The films have been grown by RF magnetron sputtering on either LaAIO3 or rutile TiO2 substrates and revealed the room temperature ferromagnetic ordering in the wide range of resistivity (10-3-106 Ωrcm). For degenerate semiconductors the values of magnetization were found to be lower. Magnetization of TiO2 doped with vanadium is maximum for all semiconductor oxides doped with transition metal impurities. Maximum values of magnetization were observed in the films with a small content of V impurity and corresponded to 4.8 Bohr magneton per V atom. The V impurity was found to be utterly in the oxidized state only in the films with a small V concentration. The magneto-optical response in the 1.5-3.0 eV energy range was observed in samples with a high magnetic moment only. The results are discussed within the model of charge-transfer ferromagnetism suggested by Coev et al. [3,4]. Supported by RFBR and Ministry of Science and Education No.16.513.11.3088.

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### PN23

# Magnetic properties and electrical conductivity on oxygen-deficient europium monoxide

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Pristine europium monoxide (EuO) is a ferromagnetic semiconductor with a Curie temperature (Tc) of 70 K. Below Tc, the conduction band is 100% spin-polarised [1] while at Tc it undergoes a metal-to-insulator transition [2]. These properties render it a promising material for future spintronic devices such as spin injection electrodes or filters [3]. Tunable control of magnetic properties is obtained by doping with oxygen vacancies. We present results on co-sputtered EuO<sub>14</sub> films with Tc as high as 140 K. The films were characterized by SQUID magnetometry, x-ray and polarized neutron reflectometry. The magnetic moment was found to increase monotonically with oxygen vacancy concentration. Density-functional theory calculations of EuO<sub>14</sub> show that oxygen vacancies act as n-type dopants and that the excess electrons preferentially populate the majority spin branch of the conduction band [4]. We also present thickness-dependent magnetic properties of EuO<sub>14</sub>, thin films, where surface-induced changes were observed in films with below 5 nm thickness [5]. Muon spin rotation measurements investigated microscopic magnetic behaviour and magnetic volume fraction as a function of themperature. Electrical transport measurements in the Hall geometry will also be presented.

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# PN24

# Magnetism and optical properties of diluted magnetic semiconductor superlattice GaGdAs/GaAs with GaGdAs nanograins

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We found ferromagnetic coupling in the rare earth elements Gd doped GaAs DMS, abbreviated as GaGdAs, which was grown by molecular beam epitaxy (MBE). However, GaGdAs is easy to have distortions or dislocations, because the atomic size of Gd is large compared to GaAs matrix, and also there is lattice mismatch between GaGdAs layer and substrate. Superlattice (SL) structure of GaGdAs/GaAs, in which GaGdAs layers are interleaved with a few ML thick GaAs layer, is considered to have a good crystallinity. We fabricate magnetic semiconductor superlattice GaGdAs/ GaAs and mono layer GaGdAs by MBE and analyze the crystal structure by means of Transmission Electron Microscopy (TEM) and measured the macroscopic magnetic properties by Alternating Gradient Magnetometer (AGM). In high resolution TEM images, there are GaGdAs grains of 2-3 nm, which has high Gd concentration and lattice-matching in GaAs Matrix. The result of magnetic measurements shows that superlattice structures have larger saturated magnetization and smaller coercive field than that of monolayers. We think that the GaGdAs grains, which are observed only in SL structures, possibly has main contribution to generate ferromagnetism in room temperature. We also measured Photoluminescence of GaGdAs samples and found the energy band gap is larger than GaAs.

### PN25

#### Ferromagnetism in hydrogenated fullerene

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We identify Stoner ferromagnetism in fcc  $C_{60}$ Hn (n=odd) by using a local density approximation in the framework of the density functional theory. Hydrogen chemisorption on fullerenes creates quasilocalized  $\pi$ -electrons on the fullerene surface, overlapping of their wave functions giving rise to a narrow half filled impurity band in the fcc  $C_{60}$ Hn. The Stoner-type ferromagnetic exchange between the itinerant electrons leads to spin-split impurity bands. The magnetic moment per  $C_{60}$ Hn molecule is 1  $\mu$ B (for n=odd) or 0 (for n=even, including zero), only one of the hydrogens contributing to the spin-split states. Direct overlapping of the quasilocalized  $\pi$ -electron orbitals is essential for the ferromagnetism.

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#### **PN26**

#### Enhancement of the magneto-optical effect by an addition of Co in pseudo-quaternary II-VI magnetic semiconductor CdMnCoTe films Masaaki Imamura\* and Keisuke Ninomiya

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Magneto-optical properties with pseudo-quaternary magnetic semiconductor CdMnCoTe films deposited on quartz glass substrates by using MBE equipment [1] have been studied at a visible wavelength region at room temperature. Deposition was made at an average rate of ~1 Å/sec for 8h using a 4N Co target sputtered by an electron gun having an emission-current of 20 mA, a 5N CdTe effusion cell heated at 490°C and a 4N MnTe cell at 1020°C. Using this system, ~2-µm-thick CdMnTe and CdMnCoTe films were prepared on 0.5-mm-thick QG substrates. These films exhibited the preferred (111) growth and paramagnetic characteristics. The magnetization intensity at 15 kOe at room temperature was approximately 2.5x10-4 emu for the  $Cd_{0.65}Mn_{0.35}Te$  film and 2.8x10-4 emu for the  $Cd_{0.56}Mn_{0.39}Co_{0.05}Te$  film, showing the values that were almost the same for two films. Faraday rotation measured at 620 nm in the CdMnCoTe film was -0.32 deg/cm-G and that in the CdMnTe film was -0.10 deg/cm-G. Thus the Faraday rotation in the CdMnCoTe films was largely enhanced compared to that in the CdMnTe films. A direct ac Faradav rotation observation under 1.27 kHz ac fields generated by a ring magnet [2] has successfully been carried out for the CdMnCoTe films

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# PN27

# Origin of ferromagnetism in Co-doped (La,Sr)TiO<sub>3</sub> diluted magnetic semiconductors

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Diluted magnetic semiconductors (DMSs) are one of the important supporting materials for the new-generation of spintronic devices. We prepared (La,Sr)TiO<sub>3</sub> (LSTO) and Codoped LSTO DMS nanocrystals through solvothemal technique combined with postannealing process. After that, the effects of concentration of Co dopant, the annealing atmosphere, temperature, duration and other conditions on the microstructures and magnetic properties were investigated. When the doping concentration is lower than 5 at.%, the observed room-temperature ferromagnetism (RTFM) of our samples is of intrinsic property. Meantime, it is found that the magnetism has the close relation with the concentration of Co dopant and the amorphous carbon which is remained after annealing. Through controlling the annealing atmosphere and duration, we can synthesize single-phase Co-doped LSTO nanocrystals. It is demonstrated that the RTFM is correlated with the structural defects induced in the Ar annealing process. It is also shown that Co-doped LSTO nanocrystals will have the CoO secondary phase when the Co doping concentration is higher than 5 at.%. According to the interesting exchange biasing phenomenon, the RTFM of our samples is ascribed to the collection of the CoO nanophase or amorphous Co metallic impurities.

#### **PN28**

# Two dimensional growth of Nb doped SrTiO<sub>3</sub> thin films and its superlattices

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Two dimensional growth modes were obtained for Nb doped SrTiO<sub>3</sub> thin films grown on TiO<sub>2</sub> terminated SrTiO<sub>3</sub> (001) substrates; as monitored by in-situ reflection high-energy electron diffraction (RHEED). At low substrate temperature ordinary RHEED oscillation was observed, indicating layer by layer growth. At high temperature, recovery of RHEED intensity after switching off laser pulses indicated the step flow growth mode. In both cases, films show metallic behavior with resistivity (p)~10-3  $\Omega$ -cm. Nb:SrTiO<sub>3</sub>/SrTiO<sub>3</sub> (NSTOm/STOn) superlattices were also being fabricated at low temperature containing five unit cell thick NSTO layers with various stacking sequences. Film surfaces were characterized by atomic force microscopy (AFM), showing high crystalline quality having steps of 3.9 A. At room temperature, resistivity was found to be 1~2  $\Omega$ -cm. The oxide superlattice growth approach might provide a

# PN29

# Anomalous hall effect in ferromagnetic nanocomposite FeGa/Fe<sub>3</sub>Ga thin films

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The anomalous Hall effect (AHE) is a central topic in the study of the ferromagnetic materials because it exhibit one of the strong evidence for spin polarization [1]. There are three main mechanisms for AHE including intrinsic deflection, side jump and skew scattering. The AHE in intrinsic homogeneous magnetic material is related to spin polarize due to difference state of spin up and spin down at the Fermi energy level and matched to the Berry phase theory [2]. In inhomogeneous magnetic material, AHE implies the presence of spin dependent scattering mechanism. The skew scattering dominates at low temperature, while the side jump becomes important at high temperature [3]. Recently, we reported the observation of weak AHE in FeGa epitaxial thin film on GaAs(001) substrate where FeGa displayed the A2 structure [4]. The FeGa were recently rapidly studied because of its giant magnetostrictive characteristics [5, 6]. In this work, we report the observation of AHE in nancocomposite FeGa/Fe<sub>2</sub>Ga epitaxial thin film on GaAs(001). Unlike FeGa (A2-structure), the strong AHE were shown up to 400 K. Interestingly, the nanocomposite FeGa/Fe<sub>2</sub>Ga film exhibited strong out-of-plane magnetization. We suggest that the origin of strong AHE are magnetization orientation of epitaxial Fe<sub>2</sub>Ga (DO3-structure).

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# PN30

# MnAs nanoclusters embedded in GaAs: Magnetism and transport properties

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The fabrication of magnetic nanocluster embedded in conventional semiconductor is a research trend because of their unique physical properties in electronics, optics and magnetism. The understanding of magnetic clusters effect on the transport and magnetic properties in semiconductor is timely demanded for the spintronic device application. In this work we embedded MnAs nanoclusters in GaAs(001) by growing MnAs/GaAs multilayer structure using molecular beam epitaxy. After growing 2000 A GaAs buffer layer on GaAs (100) substrate, the [MnAs (20 A)/GaAs ((GaAs)]4 multi-layer configurations were grown at 400 oC, finally followed by the growth of 200 A-thickGaAs capping layer to avoid the oxidation of inner layer. The tGaAs was selected as 100, 200 and 300 A. The reflection high energy diffraction (RHEED) exhibited spotty patterns, indicating the island-growth of MnAs nanocluster. The result of cross-sectional TEM measurement also showed that MnAs nanocluster are embedded in GaAs matrix successfully. The resistance was increased rapidly with decreasing temperature, indicating that the samples demonstrate semiconductor behavior. The magnetic properties were carried out by the superconductor quantum interface devices (SQUID) measurement, which showed the TC above room temperature. The detail sample growth, and magnetic and transport properties of MnAs nanocluster embedded on GaAs will be discussed.

#### **PN31**

# Half-metallic and ferromagnetic properties of carrier doping in Zn<sub>1</sub>.

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Diluted magnetic semiconductors (DMSs) have attracted a great deal of attention because of the possibility of incorporating magnetic degrees of freedom in traditional semiconductors. However, it is one of the primary challenges to create the ferromagnetic semiconductors due to the difficulty in the spin-injection into the semiconductors to form DMSs at room temperature or above room temperature. Since ZnO is a direct wide-band gap semiconductor which is piezoelectric, ZnO-based DMS would be useful for transparent thin film transistors, blue and UV light-emitting diodes and laser diodes. The precipitates of nonmagnetic Cu dopants do not contribute to the magnetism. In our work, the stability of ferromagnetic state in ZnO-based DMSs with a concentration of 2.77%, 5.55%, and 8.33% has been investigated by first-principles calculations. The magnetic moments are not affected by changing the doping concentrations, while the band gap is decreased as for increasing of Cu concentration. The band gaps in Cu-doped ZnO are 0.65, 0.22, and 0.08 eV for the Cu concentration of 2.77%, 5.55%, and 8.33%, respectively. The half-metallic character shows at low Cu concentration, while it is disappeared as for increasing of Cu concentration. The effects on the electronic state by the hole (Nitrogen) doping have been investigated.

#### PO01

# Synthesis and characterization of surface functionalized magnetic polymer microspheres with multi-shell structure

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We synthesize magnetic polymer microspheres with amino/ carboxyl groups on the surface. We also fabricate the microspheres by incorporating magnetite nanoparticles(MPs) into a silica shell around preformed polystyrene(PSI)/SiO<sub>2</sub> core/shell microspheres. Due to the silanization reaction with 3-aminopropyl triethoxysilane(APTES), the outermost silica shell is formed and linked with amino group at the same time. And we use PEG diacid to obtain carboxyl functionalized microspheres via the reaction between -NH2 and -COOH. Optical microscope and scanning electron microscopy(SEM) show that the microspheres are monodisperse with the average diameter about 2um; vibrating sample magnetometer(VSM) indicates that the synthesized sample with amino-group on the surface exhibits superparamagnetism with the saturation magnetization 7emu/g; thermogravimetry(TG) reveals that the microspheres are thermally stable when temperature is below 300°C and the Fe<sub>5</sub>O<sub>4</sub> content is 20.68%. Last, the fluorescence microscope shows that the microspheres obtained can be successfully conjugated with biomolecular, including streptavidin and antibody.

# PO02

# Production of $Fe_{3-x}Zn_xO_4$ nanoparticles for agents in hyperthermia treatment

Hiromasa Takeuchi\*, Akinobu Kurokawa, Takuya Yanoh, Shinya Yano and Yuko Ichiyanagi

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 $Fe_{5,x}Zn_xO_4(x=0.2,\ 0.4,\ 0.6,\ 0.8,\ 1)$  nanoparticle with average diameters 12nm were produced by our novel wet chemical method. The crystal structure and magnetic property of the obtained particle were investigated by X-ray diffraction and SQUID magnetometer. DC magnetization measurement showed that the coercive force Hc and saturation magnetization Ms decreased as the composite parameter x increased. This phenomenon suggests that  $Zn^{2^+}$  ions located on A-sites weaken the superexchange interaction between A and B sites. From the AC magnetic susceptibility, a sample with composition of x=0.4 is expected for heating by external field. Temperature increase depending on the magnetic field strength and frequency supported that samples with composition of x=0.2, 0.4 were appropriate for use as an agent in hyperthermia treatment.

### PO03

# Gene delivery using polyethylenimine-coated magnetic nanoparticles by static and oscillating magnetic field

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Recently gene delivery without viral vectors attracts rising attention. Low transfection efficiencies of nonviral gene vectors such as transfection reagent limit their utility in gene therapy. To overcome this disadvantage, we report a transfection method using cationic polymer polyethylenimine(PEI)-coated magnetic nanoparticles and an external magnetic field. It shows high transfection efficiency as well as low cytotoxicity. In this study, transfection efficiency of plasmid DNA in complex with PEI-coated magnetic nanoparticles was studied. HeLa cells were seeded the day before transfection. Immediately before transfection, the PEI-coated magnetic nanoparticles ( $\gamma$ -Fe<sub>2</sub>O<sub>3</sub>) were mixed with plasmid DNA. The complexes were incubated and added to HeLa cells on a NdFeB magnet and a coil. After 48 hours, transfection efficiency of plasmid DNA to HeLa cells was evaluated by the average percentage of the fluorescent cells imaged by a fluorescent microscope in triplicate. The optimal amount of magnetic nanoparticles was 3.0  $\mu$ /sample on a magnet for 4 h. The transfection efficiency was increased five-times compared without magnetic nanoparticles (or PE1). Higher transfection efficiency was obtained by applying an ac magnetic field. Further analysis on other parameters of an applied magnetic field and endocytosis are reported.

#### PO04

# Fabrication of QD-anchored magnetic nanocomposites for biomedical applications

Hangdeok Oh and Sang-wha Lee\*

Department of Chemical and Bio Engineering, Kyungwon University, Korea



# PO05

# Determination of biomolecule interaction in magnetic particle by voltammetry and ac impedance

Hoon Song and D. G. Park\*

korea atomic energy research institute, Korea

Magnetic biochip based immunoassay have potential in many fields, such as environmental immunoassay, diagnostic immunoassay, and biochemical studies and immunosensors [1]. In the field of immunoassays, cyclic voltammetry (CV) technique have become known independently as methods suitable for the detection of immunological reactions. CV measurements are the current signals based on the electrochemical species consumed and/or generated during a biological and chemical interaction process of a biologically active substance and substrate. To apply to the real biomarker and verify the sensitivity of CV-based sensor, we used the biomarker scheme using the applied the real protein on the gold surface. The capture antibodies (P21 mono antibody) are first immobilized on the positive gold sensor surface. The samples with biomarkers (P21 recombination protein) to give antibody and antigen reaction are applied on gold surface. After the non specific binding labels were successfully removed through the washing progress, only the chemically bounded bio molecules were remained on the Au surface. This enables to increase the sensitivity of the sensor for the binding of biomolecules on the gold surface. The difference in the CV represents the quantity of the antigen of interest and could be applied to the biosensor.

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#### **PO06**

# The effects of pulsed magnetic field stimulus on electromyographic activity

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Pulsed magnetic fields (PMF) conducting impulses toward the deep tissues as well as transcutaneous tissues are well known to be alternative non-invasive medical treatment for influencing human physiology as compared with acupuncture. In traditional Chinese medicine, acupuncture is used clinically to relieve myalgia and many studies have attempted to elucidate its mechanism of action [1]. The aim of the present study was to examine the immediate effects of PMF stimulus on the electromyographic (EMG) activity at acupoint HT2, Quingling, which is a trigger point as hyperiritable spots in biceps brachii muscle on the arm [2]. Six healthy volunteers were asked to perform four repetitions of withstanding maximal isometric contraction with 5-kg load and an elbow angle of 900 to cause fatigue in biceps brachii muscle on the arm. Analyzing power density spectrum from fast Fourier transformation of EMG signal, we observed that median power frequency (MDF) decreased and median power increased due to muscle fatigue during isometric contractions. It was proved that PMF stimulus on acupoint was effective in relieving muscle pain by means of increased MDF and amplitude probability distribution of RMS EMG shifted to high amplitude. Therefore, it may be concluded that appropriate PMF stimulus affects neuromuscular function.

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#### **PO07**

# Magnetic anisotropy of $Co_dFe_{3-d}O_4$ nanoparticles for applications in magnetic hyperthermia

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The magnetic behaviour of the Co<sub>4</sub>Fe<sub>3.4</sub>O<sub>4</sub> (d = 0 ~ 1.5) nanoparticles strongly depends on the concentration (d) of the cobalt ions within the spinelic ferrite structure; this is a superparamagnetic behaviour for low d, a soft magnetic behaviour for moderate d, and a hard magnetic behaviour for high d. The d parameter determines a certain value of the magnetocrystalline anisotropy [1]. The effective magnetic anisotropy in the case of nanoparticles is a very important intrinsic parameter (besides their diameter) to obtain the magnetic hyperthermia (MHT), with a great influence on the specific absorption rate (SAR), in the case of their usage in medicine to obtain tumour magnetic nanoparticles, having the diameters is to determine the magnetic anisotropy of the Co<sub>4</sub>Fe<sub>3</sub>.  $_{4}O_{4}$  nanoparticles, having the diameters of 2 ~ 17 nm, to identify its origin (magnetocystallen or the experimental results, the most suited nanoparticles (as anisotropy and diameter) in order to use them in intracellular MHT [4]. The obtained results are presented and discussed in detail in the paper.

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# PO08

Magnetic nanoemulsion as advanced drug delivery system applied to synergic procedures in the photodynamic therapy and hyperthermia trials using human mesenchymal stem cells as biological model Fernando Lucas Primo<sup>1</sup>, Daniela Regina Jardim<sup>1</sup>, Paulo Cesar Morais<sup>2</sup> and Antonio Claudio Tedesco<sup>1\*</sup> <sup>1</sup> Chemistry, Nanotechnology and Tissue Engineering Center, FFCLRP, Sao Paulo University, Ribeirao Preto, SP, Brazil <sup>2</sup> Physical, Brasilia University - UnB, Physical Institute, Brasilia-DF, 70910-900, Brazil

Tissue engineering and nanotechnology are important research fields in the future development of biological models useful for the understanding of wound healing processes, extracellular matrix activity and other biological skin activities [1]. This study reports on the synthesis and characterization of advanced drug delivery system with magnetic properties useful to Photodynamic Therapy (PDT) and Hyperthermia (HPT). Magnetic-Nanoemulsion (MNE) was obtained from spontaneous nanoemulsification to entrapment/controlled-release of lipophilic drugs as described by Siqueira-Moura[2]. The results demonstrated appropriated physical-chemistry stability to MNE biomaterial with size < 200 nm, exhibited a narrow size distribution (polydispersity index < 0.1) and zeta potential with modular value >1401 mV. Biological studies were carried out on human bone mesenchymal stem cells determination of in vitro biocompatibility after incubation with MNE which shown a safe concentration drug specific for each cellular type. Besides, the in vitro synergic effect of PDT and HPT were evaluated after simultaneous induction of photodynamic effects and magnetic field (40 Oe) application. These results confirmed that the use of nanotechnology associated with PDT nocedures led to in vitro tumor inactivation by advanced protocols that can be useful to future oncological treatment. Financial grant CNPq (DR.J.) and FAPESP projects 2009/15363-9 and 2008/537194 (FL.P.).

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#### PO09

# The morphological change of red blood cells in the hand exposed to the stimulus of strong pulse magnetic field

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The influence of magnetic field on human body has been studied to use as application of therapeutic equipment for a long time. Some studies have carefully suggested that vascular blood flow volume can be increased by an electromagnetic field stimulus [1-3]. In this study, live blood analysis (LBA) of blood collected in the hand stimulated by strong pulsed magnetic field was used in vitro to get the morphological change of red blood cells before and after magnetic stimulus. The LBA is a test methodology to approach the risk factors of illness, states of immunity, nutritional states of cells, degrees of hidden lesions, or treatments over time. The analysis can accomplish these goals functionally and preventively by observing morphological changes of ingredients constituting blood, which is alive without being chromated and collect through peripheral blood vessel. Our system was designed to generate a pulsed magnetic field that had a maximum intensity variation of 0.48 T and a transition time of 0.102 msec. LBA were done in the blood collected before and after the stimulation of 10 minutes in the hand. As a result, the hemadsorption (adhesion of red blood

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#### PO10

#### An analytical comparison in electoencephalography and electrocardiography during stimulus of pulsed magnetic field and acupuncture on acupoint PC9

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With recent increasing attention on the pulsed magnetic field stimulus as non-invasive medical treatment, diverse studies are being conducted to elucidate its effects on human physiology [1]. Among several bio-signals obtained from human body, electroencephalography and electrocardiography are known to have no side effect and provide real-time information on autonomic nervous activity with spectral analysis. The aim of the present study was to compare the changes of EEG and ECG during the PMF and acupuncture stimulus on acupoint, which is known to sedation and tonification point. In order to compare qualitatively the effect of PMF and acupuncture stimulus, the difference of alpha activities are calculated from EEG spectra compared with before, during and after stimulus, and the spectrum curves of ECG were analyzed in the frequency domain of heart rate variability. The increase of alpha activities after both stimuli could be explained that the impulse of stimulus on PC9 might pass through sensory nerve following meridian and approach the cerebral cortex, causing the CNS to be activated for pacifying emotion and calming the mind [2]. The decrease in index of sympatheougal activity after both stimuli indicates that parasympathetic nerves were activated and the sympathetic nerves were in constrained condition.

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# PO11

Reliability of a head movement compensation method based on minimum norm estimation for magnetoencephalographic recordings Sanghyun Lim and Kiwoong Kim\* Brain and Cognition Measurement Lab, KRISS, Korea

Non-invasive multi-channel recordings of neuronal activity of brain play an important role in brain research fields. Magnetoencephalography (MEG) which measures magnetic field generated by synchronous firing of group of neurons, has superior spatial and temporal resolutions with quick and easy recording procedures. However, a movement of subject's head during recordings makes significant data distortions, since the array of multi-channel sensors is not physically fixed to subject's head. One approach compensating these distortions is to utilize a singular-value-truncated minimum norm estimation (MNE) method to generate virtual MEG signals at a common reference head position. Since MNE is biased to minimize the total output power, this approach can be applied to signals which have a number of distinct sources generating event related fields (ERF), but not to spontaneous neuronal recordings. In this study, we calculated compensation errors of ERF recordings and spontaneous neuronal recordings, by comparing real sensor signals to virtually-generated signals, to see the reliability of the method for spontaneous recordings. MEG data were recorded using a 152-channel whole head helmet MEG system (KRISS MEG system) within a magnetically shielded room, and each channel was simulated and compared with each time trials.

### PO12

# The effect of small quantities of irradiation damage on the magnetic properties of Steel 316.

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The Steel 300 series is often used in applications where a non-magnetic steel is required. The materials take the austentite FCC crystal structure as opposed to the ferrite structure common in magnetic steels. It has previously been suggested that large doses of radiation (approximately 1022 ions/cm2) can cause a change in the local crystal structure of steel-316 to the ferrite structure and thus creating a magnetic signal [1]. We have performed x-ray absorption, x-ray dichroism and neutron scattering using WOMBAT at ANTSO where we also commissioned the new ANSTO Oxford Instruments cryomagnet. The samples of steel-316 were irradiated to only about 10<sup>15</sup> ions/cm2 using He<sup>+</sup> ions. This level of irradiation should not be enough to change the bulk crystal structure and yet we observe some magnetic signal. We present our results and explanation for our findings.

### PO13

#### Study of a hybrid magnet array for an electrodynamic maglev control Chan Ham<sup>1</sup>, Kurt Lin<sup>2</sup>, Younghoon Joo<sup>3</sup> and Wonsuk Ko<sup>4</sup>\*

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This paper introduces an innovative hybrid array consisted of both permanent and electro magnets that would enable us to develop an active control mechanism for underdamped EDS Maglev systems. The proposed scheme is based on the Halbach array configuration in order to take the major technical advantage from the original Halbach characteristics: a strongly concentrated magnetic field on a side of the array and a cancelled field on the opposite side. The magnetic force produced by the proposed hybrid array can also be actively controlled. This force controllability resulted from a variable magnetic field is instrumental to provide a dynamic damping force to compensate the instability of EDS Maglev system caused by external perturbations and guideway irregularities. In this study, the magnetic characteristics and capability of the proposed array is capable to produce an equivalent suspension force of the basic Halbach permanent magnet array but in a controlled mode. The effectiveness of the proposed array confirms that this study can be used as a technical framework to develop an active control mechanism for an EDS Maglev system.

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# PO14

#### FEM simulation of magnetic treatment of surface vessel

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Since most of surface vessels and submarines are constructed with ferromagnetic materials, mainly Fe-C, Magnetic treatment (Deperm) of surface vessels and submarines is required to camouflage them against magnetic detection from enemy marine force. We investigated whether we could analyze practically magnetic treatment process of a surface vessel with FEM simulation. Because a surface vessel has non-linear ferromagnetic property as well as complicated structure with very large dimension, it is impossible to analyze the magnetic treatment process through conventional FEM using scalar material constants. We tried to do the FEM analysis a hysteresis model. The magnetic property including hysteresis of the constructing material of the submarine was calculated with the hysteresis model, and compared with the material property measured with the ferrite yoke method. The alternating magnetic field applied to the submarine in FEM analysis had several tens steps, and the maximum field was 1000 A/m. The analysis results showed that overall magnetic remanence was almost the same with that of hysteresis model. However, it was found that the magnetization in the endpiece of the vessel was changed non-linearly. The simulated results could be analyzed with the ferromagnetic exchanging coupling of magnetic nodes in FEM model.

# PO15

Magnetically exchange-coupled nanoparticles as efficient heat inductor Seung Ho Moon, Jung-tak Jang, Seung-hyun Noh, Jae-hyun Lee and Jinwoo Cheon\* Chemistry, Yonsei University, Korea

The conversion of electromagnetic energy into heat by nanoparticles has the potential to be a powerful, non-invasive technique for biotechnology, but poor conversion efficiencies have hindered practical applications so far. We demonstrate a significant increase in the efficiency of magnetic thermal induction by nanoparticles. We take advantage of the exchange coupling between a magnetically hard core and magnetically soft shell to tune the magnetic properties of the nanoparticle and maximize the specific loss power (SLP), which is a gauge of the conversion efficiency. We simulated SLP as a function of magnetocrystalline anisotropy K, diameter of the nanoparticle D, and magnetization M. The dependency on K chiefly results from the internal magnetic spin fluctuation (Neel relaxation), and the dependency on D is due to both Neel and Brownian relaxation. Our simulation indicates the optimal range of K and D for nanoparticles with high SLP values. High magnetization M is beneficial too. The optimized core-shell magnetic nanoparticles have SLP values that are an order of magnitude larger than conventional iron-oxide nanoparticles. We also perform an antitumour study in mice, and find that the therapeutic efficacy of these nanoparticles is superior to that of a common anticancer drug.

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# PO16

#### Theranostic magnetic nanoparticles

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The property of magnetic materials can be defined using some parameters such as saturation magnetization(Ms), coercivity(Hc), and magnetocrystalline anisotropy(K). Understanding the interplay of these parameters is critical for optimizing magnetic characteristics we need for their effective use. We can control these parameters by changing the shape, size and composition of nanoparticles. Because there are various species of application of nanoparticles and magnetic properties needed to be optimized in each field of application are all different flexible tuning of magnetic parameters is very important. Important capabilities of magnetic nanoparticles are the external controllability of magnetic heat generation and magnetic attractive forces for the transportation and movement of biological objects. We show that these functions can be utilized not only for therapeutic hyperthermia of cancer but also for controlled release of drugs through the application of an external magnetic field. Additionally, the use of magnetic nanoparticles to drive mechanical forces is demonstrated to be useful for molecular-level cell signaling and for controlling the ultimate fate of the cell. The wide range of accessible features of magnetic nanoparticles underscores their potential as the most promising platform material available for theranostics.

# PO17

Ion-texturing & Dynamics in Layered Compounds: From Electric Automobiles to Frustrated Magnetism Martin Marsson<sup>1</sup>\*, Jun Sugiyama<sup>2</sup>, Kazuhiko Mukai<sup>2</sup>, Yutaka Ikedo<sup>3</sup>, Hiroshi Nozaki<sup>2</sup>, Kazuya Kamazawa<sup>2</sup> Masashi Harada<sup>2</sup>, Marsa Medarde<sup>4</sup>, Fanni Juranyi<sup>2</sup>, Jorge Gavilano<sup>5</sup>, James S. Lord<sup>6</sup>, Isao Watanabe<sup>2</sup>, Ekaterina Pomjakushina<sup>4</sup>, Kazimierz Conder<sup>4</sup>, Vladimir Pomjakushin<sup>5</sup>, Tsutomu Ohzuku<sup>8</sup> and Tsunehiro Takeuchi<sup>9</sup>

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The drastic change from an antiferromagnetic to a superconducting state in cuprates exemplifies how a slight change in carrier-density strongly governs magnetic and electronic properties. In cuprates the intermediate ion-layers are thought to only serve as 'passive' charge-reservoirs. However, it has become increasingly clear that also ion-order within these layers is of great importance [1]. Further, it was recently demonstrated that also ion-dynamics needs to be considered. In the Na<sub>0</sub> COO<sub>2</sub> compound it was shown that the low-T magnetic properties could be connected to Na dynamics a much the rag/cos/compound it was shown one and the order imginite properties could be connected to rad optimize a much higher temperatures [3,4]. From the context of chargeable Li-ion batteries for electrical cars, our group has recently presented a novel experimental method [5] to investigate microscopic ion-dynamics in lithium-transition-metal-oxides (Li-TMO). By using the µSR technique we have obtained results regarding ion-dynamics in a wide range of Li-TMO's [6-7]. We will here give an overview of the method/results as well as discuss their importance for the development of thin-film batteries for the future. Finally, we will also present our latest results regarding Na-dynamics in magnetically frustrated compound Na<sub>2</sub>CoO<sub>2</sub>. With the use of µSR, neutron quasi-elastic neutron scattering and diffraction, we show how solid state dynamics is imperative for tuning this compound's magnetic/electronic/thermoelectric properties

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# **PO18**

#### Large resistive switching phenomenon induced by magnetic field in nano conduction path formed in SiO<sub>2</sub>

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A large magnetoresistance (MR) exceeded 800% was observed in a nano conduction path (NCP) formed in SiO<sub>2</sub> layer sandwiched between Co film and Si substrate. Applied a high voltage of 20 MV / cm between the Co film and the backside of Si substrate, the NCP was formed by dielectric breakdown of SiO<sub>2</sub> layer. The Magnetoresistance of the NCP in in-plane magnetic field for NCP was measured. Using Ni layer instead of Co, the large MR also appeared, but could not be observed in the case of using non-magnetic materials. The structure of the device was similar to the resistive switching memory device, regarding Co layer and Si substrate as top and bottom electrode. Therefore we also measured I-V characteristic of the NCP device, and the resistive switching effect was observed. We haven't fully understood the mechanism of this large MR switching phenomenon yet. However, it is speculated that the atomic scale NCP consisted of Co atoms was formed due to dielectric breakdown of SiO<sub>2</sub>, and the nano gap in NCP was formed by magnetic field.

Jpn. J. Appl. Phys., 50 (2011) 06GG15.

### PO19

#### Size dependence simulation of saturated field in circular permalloy Xinghao Hu, Byunghwa Lim, Ilgyo Jeong and Cheolgi Kim\*

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Manipulation of magnetic heads for biomoleculars and cells separation, purification and sensing application is of great fundamental and practical interest [1]. On-chip magnets of soft magnetic structure of permallov is a classic manipulation systems for magnetic beads, which utilizes external magnetic fields for control of the magnetic beads [2]. The maximum of magnitic force generated by on-chip magnet in an applied field should be obtained for transportation of magnetic heads. We present a simulation method that is available to enable a maximum value of magnetic force under an advisable field. As shown in simulations, when the applied field is gradually increased to 130 Oe on the on-chip magnet of 5 um radius, the magnetization of the on-chip magnet is saturated that means a maximum value of magnetic force is obtained. Therefore, this simulation method can assist researchers to obtain the advisable field for maximum of magnitic force in designing variform on-chip magnets.

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# PO20

#### A new definition of magneto-mechatronics and applications Sung Hoon Kim\*, Jaewon Shin, Shuichiro Hashi and Kazushi Ishiyama

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Recently, magnetic wireless sensors and actuators have been developed for use in biomedical fields. In particular, their mechanisms depend on magnetics, such as magnetic material and physical phenomena. However, their research boundary has not been clear. Researchers talk of magnetic micro-robots, magnetic actuators and sensors, and so on [1-3]. Therefore, a new and correct definition is required. Magneticmechatronics is synergistic integration of magnetic engineering, mechanical engineering, and electronics/ computer engineering in the design, manufacturing, and control process. In particular, magneticmechanical mechanisms are directly controlled by external magnetic fields. Magneto-mechatronics is offered as a new definition of magnetic sensors and actuators. In this study, we introduce and investigate methods of magnetic field control and their suitable magnetic mechanisms. In general, the control methods and mechanisms are determined according to the types of magnetic field: gradient, alternating, and rotating magnetic fields. Because the magnetic field controls are suitable for use in the restricted spaces, magnetic devices based on magneto-mechatronics can be widely used medical applications. In magnetomechatronics, the combination of magnetic, mechanical, and computer engineering provides physical hardware, system modeling, and analysis, whereas the combination of magnetic, electrical, and computer engineering provides control methods and its algorithm, and an interface, and so on.

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#### **PO21**

### AC magnetic field frequency dependence of drug release characteristics for magnetic hyperthermia based polymer-drug encapsulate system for cancer treatment applications

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Ferromagnetic nickel, cobalt and gadolinium ferrite nanomaterials with Curie temperature values in the range of 40-50°C have been synthesized by wet chemical routes. These nanoparticles were then coated with Poly n-isopropyl Acrylamide thermo-sensitive polymer and doxorubicin drug forming a drug-polymer-nanoparticle conjugate. The conjugate was then tested for the drug release characteristics invitro by the application of alternating magnetic field to it. The study was conducted to understand the heating rate of the nanoparticles in the conjugate system and the quantitative drug release due to the heating of the nanoparticle. The quantitative study was conducted using a homemade magnetic hyperthermia setup and High Precision Liquid Chromatography (HPLC) with doxorubicin standards. The study was conducted in the range of 100 to 1000 kHz with 20 W power rating. The maximum heating rate and the drug release rate per minute was seen at around 900 kHz which was in concurrence with the theoretical calculations done using the hysteresis loop area of the material. These results will later be useful for testing on animal models to analyze invivo drug release characteristics of the conjugate.

#### PO22

#### Account of the image forces in the Bi<sub>1-x</sub>Sb<sub>x</sub>-insulator film structures. Konstantin Nicolaevich Kashirin

The Russian State Agrarian University-Moscow Timiryasev Agricultural Academy Kaluga Branch, Russia

Semiconducting alloy Bi1, Sb, (x=0,11) films of the different fixed thickness with a thin layer of insulating coating (SiO2) of variable thickness are investigated. The measuring of kinetic parameters was conducted in a superconducting solenoid gap (with magnetic induction to 9 T) and on the different temperatures (4,2÷200 K). The calculations of magnetic resistance, of the carries mobility and the Hall coefficient was done. The resistance oscillations is observed depending on the thickness of insulating coating for all films. The kinetic parameters oscillations was observed as function of coating thickness at different temperature and on magnetic field. So, the following should be noted by the results of experimental investigations: - the resistance oscillations of all investigated bismuth-antimony alloy films is observed depending on the thickness of dielectric coating and the period of oscillations is amounts 80-100 A average; - the magnetic resistance and the Hall coefficient oscillate with the variation of the dielectric coating thickness in antiphase with the analogous oscillation resistance; - the sphere of negative magnetic resistance is observed. One of the possible mechanisms accounting for the oscillations of the kinetic parameters of semiconducting films with insulator coating is taking into accounts the image forces.

Semiconducting alloy Bi1.Sb. (x=0,11) films of the different fixed thickness with a thin layer of insulating coating (SiO<sub>2</sub>) of variable thickness are investigated. The measuring of kinetic parameters was conducted in a superc solenoid ean and on the different temperatures. The calculations of magnetic resistance, of the carries mobility and the Hall coefficient was done. The assumptions that the one of the possible mechanisms accounting for the oscillations of the kinetic parameters of semiconducting films with dielectric coating is taking into accounts the reflection forces were done

# PO23

#### Time-resolved pump-probe measurement of polarization rotation in nano-structured chiral metamaterial

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In this study, we performed time resolved pump-probe measurement in two nano-structured metamaterials with opposite chirality which were fabricated by e-beam lithography on 30nm thick gold film. Transient absorption was measured to investigate the temporal behavior of the polarization rotation in metamaterial. It is well known that surface plasmonic resonances in metamaterial are originated from its structure rather than composition [1,2]. Surface plasmonic resonances in metamaterials which were used in this study were located in the visible spectral range with resonance peaks at 625nm and 670nm and an optical rotatory dispersion was observed. Time resolved pumpprobe measurements were carried out with 35fs Ti:sapphire laser operating at 800nm. BBO crystal was used to generate a 400nm pump pulses. For the probe, a sapphire crystal plate was used to generate a white light continuum. With linearly polarized pump beam, the electron dynamics in the metamaterial structure was studied. Also with right and left circularly polarized pump, we investigated the time dependent rotation of polarization of the probe beam in a polarization bridge configuration.

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### PO24

#### Structural and magnetic properties of glassy like carbon synthesized by pyrolysis of sucrose

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The structure and properties of glassy carbon (GC), a technologically important and widely used material, is still not well established. The glassy carbon samples were produced by a low-temperature pyrolysis of sucrose. The thermal destruction occurred in five stages to minimize carbon losses. At the first stage (175 °C, 3 h) a viscous melt was obtained due to detachment of volatile components and water. After the second and the third stages (300 °C, 1 h and 400 °C, 0.5 h, respectively) the material becomes solid. The next stage (525 °C, 1h) and the last stage (700 °C, 1 h) are performed to complete removal of all volatile impurities. The samples were then characterized using AFM and SEM to determine its structure and morphology. The magnetic properties were studied using vibrating sample magnetometer. It was found that the glassy carbon sample has an inherent fullerene like morphology whose concentration increases with increased heat treatment. Further, the samples also exhibited a distinct change in magnetic properties based on the pyrolysis. Paramagnetic and diamagnetic properties were found in the samples. The same has been explained based on the composition of the samples and structural properties.

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### PO25

#### Simulation of energy dispersive mode for RITA-type cold neutron triple axis spectrometer SIKA

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SIKA is a high flux cold-neutron triple axis spectrometer in Bragg Institute. Equipped with a multiblade analyser system (13 PG(002) blade) and position sensitive detector. SIKA possesses high flexibility to efficiently run in a traditional step-by-step mode or various mapping (or dispersive) modes by changing the configuration of analysers and detectors. In this study, the energy dispersive modes, namely the so-called E1 mode [1] at two different energy transfers is simulated using Monte Carlo ray-trace package SIMRES. [2] The simulated results show that SIKA could effectively work on E1 mode at low and intermediate energy transfer with reasonable energy and Q resolutions. The energy resolution is around 0.23meV at Ei=5meV while it increases to1.8meV at Ei=15meV. This simulation shows SIKA has the potential in operation flexibility and data-acquisition efficiency, and provides valuable references to the future inelastic neutron scattering experiments on SIKA.

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# PO26

# Efficiency of Energy base deperm protocol

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Magnetic treatment of surface vessels and submarines (Deperm) is a required to camouflage them against magnetic detection from enemy marine force. So far, deperm has been accomplished by applying an alternating magnetic field of which amplitude decreases linearly. However, the reduction of the residual flux density in the direction of magnetic field is not linear in the case of the linear protocol, since the ferromagnetic material used to construct a surface vessel, mainly Fe-C, shows a nonlinear behavior in an alternating magnetic field. This is one of main reasons to make an ordinary deperm protocol inefficient. In this paper, we propose an energy base deperm protocol which calculated by an algebraic hysteresis model, and present how to optimize the deperm protocol with the minor loops of magnetic hysteresis. We found out that step number could be reduced in the energy base deperm protocol compare with in the linear protocol, because the larger numbers of deperm steps are dedicated in the irreversible domain process region on the magnetic hysteresis.

### **PO27**

#### Effect of a magnetic field on mixed convection of a nanofluid in a square cavity

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<sup>3</sup> Department of Mathematics, Payame Noor University, Najafabad, Isfahan, Iran

The problem of mixed convection in a differentially heated lid-driven square cavity filled with Cu-water nanofluid under effect of a magnetic field is investigated numerically. The left and right walls of the cavity are kept at temperatures of Th and Tc respectively while the horizontal walls are adiabatic. The top wall of the cavity moves in own plane from left to right. The effects of some pertinent parameters such as Richardson number (ranging from 0.1 to 10), the volume fraction of the nanoparticles (ranging 0 to 0.1) and the Hartmann number (ranging from 0 to 100) on the fluid flow and temperature fields and the rate of heat transfer in the cavity are investigated. It must be noted that in all calculations the Prandtl number of water as the pure fluid is kept at 6.8, while the Grashof number is considered fixed at 10<sup>4</sup>. The obtained results show that the rate of heat transfer increases with an increase of the Rayleigh number, while but it decreases with increase in the Hartmann number. Moreover it is found that by increase in volume fraction of the nanoparticles the rate of heat transfer can be enhanced or deteriorated compared to the based fluid.

### **PO28**

#### Magnetohydrodynamic free convection in a square cavity heated from below and cooled from other walls

Ali Akbar Abbasian Arani\*, Mostafa Mahmoodi and Saeed Mazrouei Sebdani Mechanical Engineering, University of Kashan, Iran

Magnetohydrodynamic free convection fluid flow and heat transfer in a square cavity filled with an electric conductive fluid with Prandtl number of 0.7 has been investigated numerically. The horizontal bottom wall of the cavity was kept at Th while the side and top walls of the cavity were maintained at a constant temperature Tc with Th>Tc. The governing equations were solved numerically while the SIMPLER algorithm was used. A parametric study was performed, and the effects of the Rayleigh number and the Hartman number on the fluid flow and heat transfer were investigated. The results showed that temperature distribution and flow pattern depended on both strength of the magnetic field and Rayleigh number. For all cases two rotating eddies were formed. Using the longitudinal magnetic field results in a force (Lorentz force) opposite to the flow direction that tends to decrease the flow velocity. The magnetic field decreased the free convection and flow velocity. Also it was found that for higher Rayleigh numbers a relatively stronger magnetic field was needed to decrease the heat transfer. Moreover for low Rayleigh numbers, by increase in the Hartman number, free convection is suppressed and heat transfer occurs through conduction mainly.

# PO29

Interaction of a magnetic field and buoyancy force in a square cavity filled with a fluid with low Prandtl number

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In this paper, influences of Lorentz force due to a magnetic field on buoyancy driven heat transfer in a square cavity containing a fluid with low Prandtl number is investigated numerically. The left and right walls of the cavity are kept at temperatures of Tc and Th respectively, while the horizontal walls are kept adiabatic. Effect of existence of an adiabatic body in the middle of the cavity is considered too. The dimensionless governing equations associated with the boundary conditions are solved numerically using the finite volume method. Streamlines and isotherms are shown for Grashof number (ranging from 105 to 107), Hartman number (ranging from 0 to 100) and Prandtl number (ranging from 0.005 to 0.1). The results show that variations of magnetic field donot have noticeable effect on buoyancy force, while increase in magnetic field increases Lorentz force. Moreover buoyancy and Lorentz forces

# PO30

#### Numerical study of magnetic field effects on flow field and heat transfer in a cavity filled with porous materials

increase by increase in Grashof number in a constant Hartman and Prandtl numbers.

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Natural convection flows in a square cavity filled with a fluid-saturated porous medium under the effect of magnetic field has been studied numerically. A square hot body with length W is placed in the center of the cavity (with length L). The ratio of W/L is assumed to be 0.5. Continuity, momentum and energy equations are solved using the finite volume technique under the boundary conditions. The two vertical sides of the cavity are cooled and two horizontal bottom and top walls are adiabatic. The Prandtl number is 0.71, the Rayleigh number changes from 10<sup>4</sup> to 10<sup>6</sup>, the Darcy number varies from 10<sup>-5</sup> to 10<sup>-3</sup> and the Hartmann number alters from 0 to 100. Numerical results are presented in terms of stream functions, temperature profiles. The results say increasing the braking effect of magnetic field as increasing the Hartman number; lead to decreases the heat transfer. Also at low Darcy number, fluid circulation is weak due to high hydraulic resistance offered by porous medium for Darcy number down to 10<sup>-5</sup>. Average Nusselt number increases with enhancement of Rayleigh number and decreasing the Hartmann number.

#### PO31

#### Numerical simulation of magnetohydrodynamic Benard convection in a shallow enclosure

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In this paper the problem of magnetohydrodynamic Benard convection in a shallow rectangular enclosure with aspect ratio of 10 (width to height ratio) has been investigated numerically. The enclosure is filled with an electric conductive fluid with Prandtl number of 1. The bottom and the top walls of the cavity are maintained at a constant temperature Th and Tc respectively, with Th > Tc, while; the side walls of the cavity are kept insulated. The governing equations written in terms of the primitive variables are solved numerically using the finite volume method and the SIMPLER algorithm is employed to couple velocity and pressure fields. Using the developed code, a parametric study is performed, and the effects of the Rayleigh number and the Hartman number on the fluid flow and heat transfer inside the enclosure are investigated. The results show that temperature distribution and flow pattern inside the enclosure depend on both the strength of the magnetic field and Rayleigh number. Moreover it was found that at low Rayleigh number with increase in Hartman number the flow intensity decreases and fluid becomes stagnant. At high Rayleigh numbers number of the Benard cell decreases by increase in Hartman number.

# PO32

#### Near-electrode effcets of magnetic fields in electrochemistry Michael Coey and Peter Dunne

Physics, Trinity College dublin, Ireland

Remarkable influences of magnetic fields have been discovered on the diffusion layer and on the electrochemical double layer. The former are mainly due to magnetic pressure, which allows the patterning of deposits from both paramagnetic and diamagentic ions (using a nonelectroactive paramagnetic species such as Dy3+) The effects of small cylindrical [1] and linear magnet arrays are demonstrated. The magnetic patterning is estimated to be effective down to about the 100 micron scale The essential requirement for patterning of this type are orthogonal concentration gradients and magnetic field gradients. More surprising is the evidence we have found for an influence of a static magnetic field on the electrochemical double layer, a poorly understood region about 1 nm thick at the cathode. Using a nitrobenzene model redox system, we show that the double layer capacitance can be modified by a factor of two at an appropriate potential. The influence of Maxwell stress on the free radicals in solution is discussed as a possible explanation. More remarkable are

[1] P. Dunne, L. Mazza and J. M. D. Coey, Phys Rev Letters 107 042501 (2011)

#### **OA01**

#### Resonant magnetic x-ray scattering: Beamline P09 at PETRA III at DESY

Joerg Strempfer, Sonia Francoual, Dinesh K. Shukla and Arvid Skaugen DESY. Germany

Resonant and non-resonant magnetic scattering are complementary tools to magnetic neutron diffraction for the investigation of magnetic properties in solids. Resonant magnetic x-ray scattering opens the possibility to investigate magnetic order in single crystals and layered systems element selectively. Also the interference of magnetic order with orbital or charge order can be determined through the polarization properties of the x-ray and the special form of the scattering cross sections, describing the variation of the polarization states through the scattering process. The new beamline for Resonant Scattering and Diffraction (RSD), P09 at PETRA III, is designed especially to address the above mentioned issues. Low temperature cryostats and high magnetic fields of up to 14T allow the investigation of magnetic properties of materials as function of polarization, temperature and field. Two diffractometers are available to conduct experiments using the highly brilliant synchrotron beam available at PETRA III. The polarization of the incident x-ray beam can be varied arbitrarily through x-ray phase-plates and the polarization of the scattered beam is analyzed by a polarization analyzer. Examples for the performance of the beamline and experimental results on magnetic scattering from multiferroic systems will be presented.

### **OA02**

#### Magnetic and dielectric properties of FeTiO<sub>3</sub>

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The magnetic and dielectric properties of single crystal iron ilmenite FeTiO3 have been investigated. The temperature dependence of the magnetization for H parallel to the c axis (M//c) has a sharp peak at 58 K ( $T_N$ ) and abruptly decreases below  $T_N$ , while the magnetization for H perpendicular to the c axis (M<sup>⊥</sup>c) is almost independent of the temperature. These features indicate that the ground state is Neel state with the easy axis in the c direction. Above  $T_N$  the value of M//c is larger than that of M<sup>1</sup>/c due to the magnetic anisotropy. These results are consistent with previous report by Kato et al... The temperature dependence of the dielectric constant for E parallel to the c axis also shows a kink at  $T_N$ . Furthremore the dielectric polarization for E parallel to the c axis abruptly increases at  $T_N$ . Hence this kink is ascribed to ferroelectric phase transition. These results indicate that FeTiO3 is a strong candidate for a new type of multiferroic materials

# **OA03**

#### Theoretical study of tuning polarization and magnetism of BiCoO<sub>3</sub> Yu-jun Zhao and Xing-yuan Chen Department of Physics, South China University of Technology, China



#### **OA04**

### Soft x-ray synchrotron radiation spectroscopy study of Co<sub>0.6</sub>Fe<sub>0.9</sub>Mn<sub>1.5</sub>O<sub>4</sub> spinel with nano-checkerboard patterns

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In the AB2O4-type spinels, transition-metal ions occupy either the tetrahedral (Td) A site or the octahedral (Oh) B site. Recently, nano-checkerboard (CB) patterns have been observed in ZnMnGaO<sub>4</sub>, MgMn<sub>15</sub>Fe<sub>05</sub>O<sub>4</sub>, and Co<sub>0.6</sub>Fe<sub>0.9</sub>Mn<sub>15</sub>O<sub>4</sub> [1]. When properly annealed, these spinels show two different structures with different chemical phases, where one is cubic and the other is tetragonal or orthorhombic. This structural and chemical phase separation has been explained by the Jahn-Teller (JT) distortion. To confirm this idea, it is important to determine the valence states of the transition-metal ions that constitute these spinel oxides. In this work, we have studied the electronic structure of checkerboard Co<sub>0.6</sub>Fe<sub>0.9</sub>Mn<sub>1.5</sub>O<sub>4</sub> single crystal by employing soft x-ray absorption spectroscopy (XAS) and soft x-ray magnetic circular dichroism (XMCD). The measured T 2p XAS spectra (T=Co, Fe, Mn) of Co<sub>0.6</sub>Fe<sub>0.9</sub>Mn<sub>1.5</sub>O<sub>4</sub> reveal that the valence states of Co, Fe, and Mn ions are nearly divalent (Co<sup>2+</sup>), trivalent (Fe<sup>3+</sup>), and mixed-valent (Mn<sup>2+</sup>-Mn<sup>3+</sup>), respectively. Very weak XMCD signals are observed in checkerboard Co<sup>0.6</sup>Fe<sup>0.9</sup>Mn<sup>1.5</sup>O<sub>4</sub>. \*Corresponding Author : kangjs@catholic.ac.kr

[1] C.L. Zhang, et al., Appl. Phys. Lett. 91, 233110 (2007).

#### **OA05**

#### Zn-substitution effects in multiferroic Cu<sub>3</sub>Mo<sub>2</sub>O<sub>9</sub>

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We present the Zn substitution effects on the quasi-one dimensional frustrating spin system  $Cu_3Mo_2O_9$ . This material has S = 1/2 distorted tetrahedral spin chains and undergoes the multiferroic phase below 7.9 K without magnetic field [T. Hamasaki et al., Phys. Rev. B 77 (2008) 134419, H. Kuroe et al., J. Phys. Soc. Jpn. 80 (2011) 083705]. The S = 0 Zn<sup>2+</sup> ions substitute for the S = 1/2 Cu<sup>2+</sup> ions at the three crystallographically different sites, namely Cu1, Cu2, and Cu3. The spins at the Cu2 and Cu3 sites form nearly spin-singlet spin dimers at the bisectors of the nearest neighbor Cul sites in the quasi-one dimensional spin chain [H. Kuroe et al., J. Phys.: Conf. Ser 200 (2010) 022028, H. Kuroe et al., Phys. Rev. B 83 (2011) 184423]. In this study, we focus on the magnetic, electric, and thermodynamical properties of the single crystal of (Cu,Zn)3Mo2O9. Comparing our previous results in polycrystalline samples [M. Hase et al., J. Phys. Soc. Jpn. 77 (2008) 034706], we discuss the reduction of the antiferromagnetic correlation and the singlet-pair breaking effect on the Cu2-Cu3 spin dimers. The dielectric constant and the specific heat in (Cu,Zn)3Mo2O9 will also be discussed in this study.

# **OA06**

#### Effects of bismuth substitution on the magnetic properties of Bi,Co, MnO.

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Coexistence of magnetism and ferroelectricity, along with coupling between them, ensures many technological applications of materials named multiferroics. Transition metal and/or rare earth oxides with ABO3 perovskite and AB2O4 spinel structures are good candidates due to the substitution of magnetic and non-magnetic ions at the A and/or B sites [1]. Bi-substituted Bi,Co2, MnO4 multiferroic samples, well-characterized structurally by XRD and Rietveld refinement together with SEM-EDX techniques, were synthesized by a polymeric precursors method. A gradual expansion of the lattice parameter due to the substitution of  $Co^{3+}$  by  $Bi^{3+}$  is observed for x = 0.1, 0.2 and 0.3. Magnetic characterization was carried out by SQUID measurements. The inverse susceptibility exhibited ferrimagnetic behaviour, as also proved by the strong antiferromagnetic and ferromagnetic interactions shown by the ZFC and FC curves. The M(H) loops performed at various temperatures below  $T_c$  showed larger loop areas when decreasing the temperature, indicating an increasing ferromagnetism. The inclusion of Bi in Co2MnO4 creates fluctuating valence states of both Co and Mn ions  $(Co^{2+}/Co^{3+}, Mn^{2+}/Mn^{3+}/Mn^{4+})$  which occupy the tetrahedral and octahedral sites [2]. Magnetization measurements showed that  $T_c$  and  $M_s$  increase with increasing Bi<sup>3+</sup> concentration due to the redistribution of the magnetic cations, in particular the Mn3+ and Mn4+ ones.

[1] P. Barahona, et. al. J. Chil. Chem. Soc. 50. N 2 (2005) 495. [2] N. E. Rajeevan, et. al. Mater. Sci. Eng. B 163(2009) 48.

# OA07

### Cross-correlation effects in multiferroic Cu<sub>3</sub>Mo<sub>2</sub>O<sub>9</sub>

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We present the cross correlation effects in a multiferroic material Cu<sub>3</sub>Mo<sub>2</sub>O<sub>9</sub>. This material contains spin tetrahedra made from the Cu2+ ions which form a edge sharing spin chain. The quasi-one dimensional magnon dispersion curve, the ferroelectric properties caused by the geometrical magnetic frustration, and a strong magnetocapacitance effect have been reported [H. Kuroe et al., Phys. Rev. B 83 (2011) 184423, H. Kuroe et al., J. Phys. Soc. Jpn. 80 (2011) 083705]. In this study, we focus on the change of the longitudinal magnetization under the electric field. Using a SQUID magnetometor which has two electric leads, we measured the longitudical magnetization Ma along the a axis and the electric polarization Pc along the c axis of a plate-like single crystal of Cu<sub>3</sub>Mo<sub>2</sub>O<sub>9</sub> simultaneously. We measured the longitudinal magnetization along the a axis (Ma-Ha loop) under the static electric field Ec along the c axis, and the Ma-Ec loop under the static Ha. The details of our results including the change of the Ma-Ha loop under the static electric field and butterfly-loop in the Ma-Ec loop will be presented in this study.

# **OA08**

#### Nonlinear current-voltage characteristics of (La<sub>0.5</sub>Eu<sub>0.5</sub>)<sub>0.7</sub>Pb<sub>0.3</sub>MnO<sub>3</sub> Single crystals: Possible manifestation of the internal heating of charge carriers

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Current--voltage characteristics of the polycrystalline substituted lanthanum manganite  $La_{0.7}Ca_{0.3}MnO_3$  were experimentally studied at T = 77.4 K in magnetic fields up to 13 kOe. In these characteristics, a portion of negative differential resistivity was observed above a certain threshold value of critical current density j caused, in our opinion, by nonequilibrium heating of the electron gas due to low thermal conductivity of the manganite material. Because of the nonlinearity of the current--voltage characteristics, the field dependences of resistivity  $\rho(H)$  appear extremely sensitive to the value of a transport current. In this case, the  $\rho(H)$  dependences reveal both ordinary negative and positive magnetoresistance.

### **OA09**

#### High field phase diagram in multiferroic Cu<sub>3</sub>Mo<sub>2</sub>O<sub>9</sub>

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We present the multiferroic behavior in the quasi-one dimensional antiferromagnet Cu-Mo-O-This material has spin chains of the distorted spin tetrahedra made from  $S = 1/2 \text{ Cu}^{2+1}$  ions. In ICM2009, we have presented that the magnon dispersion relation in Cu<sub>3</sub>Mo<sub>2</sub>O<sub>9</sub>, which is well explained by the hybridization effects between a magnon excitation from the quasi-one dimensional spin system and that from the isolated spin dimers [H. Kuroe et al. J. Phys.: Conf. Ser. 200 (2010) 022028. H. Kuroe et al., Phys. Rev. B 83 (2011) 1844231. Recently, we found that this material shows a multiferroic behavior [H. Kuroe et al., J. Phys. Soc. Jpn. 80 (2011) 083705]. We proposed that the origin of the multiferroic behaviors in Cu<sub>3</sub>Mo<sub>2</sub>O<sub>9</sub> is the charge redistribution effect in a frustrated Mott insulator studied by Bulaevskii et al. [L. N. Bulaevskii et al., PRB 78 (2008) 024402.], which is presented by Khomskii as an invited talk in ICM2009 [D. I. Khomskii, J. Phys.: Condens. Matter 22 (2010) 164209]. In this study, we present the magnetic-field-temperature phase diagram up to 55 T obtained from the temperature and magnetic-field dependences of the magnetization the dielectric constants, the electric polarization the pyroelectric current and the specific heat.

**OA10** 

Magnetoelectric effect in Ca<sub>2</sub>FeAlO<sub>5</sub>

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Ca<sub>2</sub>FeAlO<sub>5</sub> has the same crystal structure as mineral Brownmillerite. The oxygen deficient perovskite type structure is constructed by the alternate stacking of (Fe,Al)O<sub>6</sub> octahedra and (Fe,Al)O4 tetrahedra. An early structure analysis showed that Ca2FeAlO5 belongs to a noncentrosymmetric space group Ibm2 [1], which originats from the in-phase stacking of (Fe,Al) O4 tetrahedra. This compound shows antiferromagnetic order below 350 K. Mossbauer spectroscopy measurements conclude that the magnetic moments of Fe<sup>3+</sup> align along the a-axis[2]. Because the crystal structure lacks the inversion symmetry, Ca<sub>2</sub>FeAlO<sub>5</sub> is a candidate material which has a magnetoelectric coupling. Here we show the measurements of magnetoelectric effect in single crystalline Ca<sub>2</sub>FeAlO<sub>5</sub>. The field dependence of the magnetization along the a-axis shows a metamagnetic transition caused by the spin flop of Fe<sup>3+</sup> moments. Moreover, the dielectric constant and electric polarization show an anomaly at around the phase transition. These results indicate that Ca<sub>2</sub>FeAlO<sub>5</sub> shows magnetoelectric effect in a magnetic field induced spin flop phase, which is one example of a new route to realize a magnetoelectric coupling above the room temperature.

[1] A. A. Colville and S. Geller, Acta Cryst, B27, 2311(1971) [2] R. W. Grant et.al., J. Appl. Phys. 39, 1122 (1968)

# **OA11**

#### Electronic and magnetic phase separation in the semimetallic ferromagnet EuB<sub>6</sub>

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EuB6 is a low-carrier semimetal, which undergoes a unique paramagnetic to ferromagnetic transition displaying two consecutive features at  $T_{c1}$ = 15.5 K and  $T_{c2}$  = 12.6 K in electrical transport, magnetization and specific heat. Its structural and electronic simplicity makes EuB6 an ideal model system for studying the effect the magnetic state of a system has on its electronic transport properties. In particular, the fundamentals of intrinsic electronic and magnetic phase separation can be studied, which play a critical role for the colossal magnetoresistance (CMR) in magnetic semiconductors and manganites. In explaining CMR in EuB<sub>6</sub>, being largest at T<sub>cl</sub>, the concept of magnetic polaron (MP) percolation has been invoked. We report resistance noise and weakly nonlinear transport measurements, being sensitive to the microgeometry of the electronic system. We find direct evidence for electronic phase separation when cooling through T<sub>c1</sub> and a diverging noise level at T<sub>c2</sub>, which we interpret as magnetically-driven percolation resulting from the overlap of spatial inhomogeneities in conductivity produced by MP. Interestingly, MP percolation can be induced also by applying magnetic fields in the paramagnetic region. We find that the density of MP can be described by a universal scaling function, being related to a single critical magnetization.

# **OA12**

#### Study of magnetic and magnetodielectric properties of perovskite YhCrO<sub>2</sub>

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Perovskite YbCrO<sub>3</sub> shows the G-type antiferromagnetic ground state for  $Cr^{3+}$  moments ( $T_N \sim 118$  K). Because of the Dzyaloshinskii-Moriya interaction, the G-type structure becomes canted, and then YbCrO3 exhibits weak ferromagnetism. When weak magnetic fields are applied ( $B \le 0.3 T$ ) temperature-induced magnetization reversal is observed. The magnetization reversal is probably associated with a keen competition between antiferromagnetically coupled Cr<sup>3+</sup> - Yb<sup>3+</sup> moments. Recently, a distinct magnetodielectric effect was observed in another perovskite SmMnO<sub>2</sub> which also shows temperature-induced magnetization reversal due to the competition between antiferromagnetically coupled Mn<sup>3+</sup> - Sm<sup>3+</sup> moments [1]. Thus, it is intriguing to study magnetodielectric properties of YbCrO3. In this study, we investigated magnetic and magnetodielectric properties for single crystals of YbCrO<sub>3</sub> grown by the flux method. The comparison of magnetodielectric properties between YbCrO<sub>3</sub> and SmMnO<sub>3</sub> provides useful information to understand the distinct magnetodielectric effect observed in SmMnO<sub>2</sub>.

[1] J.-S. Jung et al. Phys. Rev. B 82, 212403 (2010)

#### **OA13**

#### Interplay among spin, orbital, and lattice degrees of freedom in a frustrated spinel Mn<sub>3</sub>O<sub>4</sub>

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Spinel oxides, AB<sub>2</sub>O<sub>4</sub>, are one of the typical systems for studying the effects of geometrical frustration and the possible coupling among spin, orbital and lattice degrees of freedom. Mn<sub>3</sub>O<sub>4</sub> exhibits the cooperative Jahn-Teller (JT) transition from cubic to tetragonal structure at 1443 K with the 3z2-r2 ferro-orbital order at B-site Mn3+. The gigantic JT distortion (~16 %) would encourage the formation of one-dimensional (1D) antiferromagnetic spin chains at B-site Mn3+ normal to the c-axis. A macroscopic degeneracy of the arrangements of 1D-spin-chains still remains, even if one considers the magnetic interaction between the nearest neighbor 1D-chains. It was reported that Mn<sub>3</sub>O<sub>4</sub> exhibits successive magnetic phase transitions from paramagnetic to Yafet-Kittel (43 K), incommensurate (IC) (40 K) and commensurate (C) phase (34 K)[1], which is an indicative of the surviving magnetic frustration. Another phase transition is induced by applying magnetic field in the low-temperature C phase [2]. In order to elucidate the magnetic-field induced phase transition, we have measured synchrotron x-ray diffraction and neutron scattering in a magnetic field along [100] c. We propose a scenario of the magnetic-field induced phase transition, where the rearrangement of 1D-spin-chains is triggered by orbital hybridization at B-site Mn3+ through the spin-orbit coupling.

[1] R. Tackett et al., Phys. Rev. B 76, 024409 (2007). [2] T. Suzuki et al., Phys. Rev. B 77, 220402(R) (2008)

### **OA14**

#### Structural and dielectric study of hexagonal Y<sub>0.8</sub>Sr<sub>0.2</sub>MnO<sub>3</sub> Compound

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Strontium doped YMnO3 compound have been prepared in single-phase form by using high temperature solid sate reaction method. The structural and dielectric properties of the prepared sample have been carried out in the wide range of temperature and frequency. Temperature and frequency variations of dielectric permittivity measurements show that the present hole doped YMnO3 compound exhibits a competing interaction in between the ferromagnetic double - exchange and antiferromagnetic super-exchange interactions, further a weak signature of ferromagnetic ordering at lower temperature has been witnessed. Our resistivity measurement reveals that the material is highly insulating at lower temperature just below the reported AFM transition in this class of materials

# **OA15**

#### Controlling the superparamagnetic limit using the magnetoelectric effect

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MRAM (Magnetic Random Access Memory) is attracting considerable attention as a possible candidate for next generation of non-volatile memory devices. Major challenges facing scientists today are overcoming thermal instability and reducing the writing energy. As bit density increases, this requires smaller magnetic regions to be written magnetically at small scales and external energies become a dominant feature. For example in the small scale, thermal energy becomes larger than magnetic anisotropy, resulting in superparamagnetic behavior. For a given small scale size, the temperature above which superparamagnetism begins is referred to as the blocking temperature. Superparamagnetic behavior is undesirable due to limits it places on the size of a bit of information. Recently, researchers have begun to study magnetoelectric (ME) memory related to some issues with memory. Their focus has been on writing processes using electric field-induced strain rather than addressing the more challenging topic of controlling superparamagnetism. Here we report experimental results showing the introduction of a magnetoelectric anisotropy energy to alter the blocking temperature and thus control the superparamagnetic behavior. This is achieved via an electric field-induced strain in a ME composite. Our results show 50 K increase of the blocking temperature with the application of an electric field.

1. Y. Jun, J. Seo, J. Cheon, Accounts of chemical research 41, 179 (2008). 2. F. Zavaliche et al., Nano letters 7, 1586 (2007). 3. T. Wu et al., Applied Physics Letters 98, 012504 (2011). 4. M. Bibes, A. Barthelemy, Nature Materials 7, 425 (2008).

#### **OA16**

#### Magneto-Capacitance Effect and Electric Polarization in Spinel Co<sub>2</sub>MnO<sub>4</sub>

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Structural phase diagram of spinel  $Co_{3x}Mn_xO_4$  (1.0  $\le x \le 2.0$ ) was reexamined at the room temperature with a focus on the cubic-tetragonal phase boundary. A specific composition of Co2MnO4 (x=1.0) was selected from the phase diagram study under the requirement of both high ferrimagnetic transition temperature (Tc) and non-mixed structural phase condition. Experimental result of magnetization measurements and local distortions due to the Jahn-Teller Mn3+ ions support the stabilization of Yafet-Kittel type magnetic configuration in the cubic inverse spinel Co2MnO4. Correlation between the observed forced magnetization - a continual increment of magnetization above the saturation field - and magneto-capacitance effect (MCE) was clarified. MCE of up to 7 % at 8 T around Tc was observed without an appreciable hysteresis and nonlinearity in the field variation. Diffraction-spectroscopy techniques were also used to identify the local chemical and structural environments of both Co and Mn ions at two crystallographic sites of the spinel structure.

### **OA17**

#### Magnetically induced polarization in copper metaborate CuB<sub>2</sub>O<sub>4</sub> Khanh Duy Nguyen<sup>1</sup>, Nobuyuki Abe<sup>2</sup>, Masashi Tokunaga<sup>3</sup>, Mitsuru Saito<sup>1</sup> and Takahisa Arima<sup>2</sup>

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Intriguing electro-magneto optic effect of non-centrosymmetric magnet copper metaborate has attracted huge consideration during last time [1]. In here we report the magneto-electric behavior of CuB2O4 under external magnetic field. It is well known that CuB2O4 undergoes a phase transition from paramagnetic state to canted-antiferromagnet at 21 K, then followed by a second transition to spiral order around 10 K. Although the magnetic structure at low temperature below 10 K has not completely been made clear yet, some evidences of incommensurate helical order in this range of temperature have been observed [2]. Result of magnetization measurement points out the existence of an anomaly around 1, 3 T at 2 K, then disappeared when increasing temperature above 10 K, which may be attributed to the transition from commensurate to incommensurate phase. Furthermore, electric polarization was also observed by applying a magnetic field along the [110] and [1-10] axes without poling electric field in low temperature regime. The magnitude of polarization in static field was found to be compatible with the measurement in pulse magnetic field. Besides novel optical characteristic, these results show the ability of generation and control electric polarization by magnetic field in CuB2O4 as well as application to opto-electronic devices.

[1] M. Saito, K. Ishikawa, S. Konno, K. Taniguchi and T. Arima, Nature Mater. 8, 634 (2009) [2] M.Boehm, B. Roessli, J. Schefer, A. S. Wills, B. Ouladdiaf, E. Lelievre-Berna, U. Staub and G. A. Petrakovskii, Phys. Rev. B 68, 024405 (2003)

# **OA18**

### Impedance spectroscopy of ferromagnetic oxide: Pr<sub>0.6</sub>Sr<sub>0.4</sub>MnO<sub>3</sub>

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In recent years, coupling between electrical and magnetic polarizations in multiferroicmagnetoelectrics have become a topic of considerable interest due to novel applications and fundamental physics involved [1]. Impedance spectroscopy has been extensively used to investigate cross coupling between magnetic and electrical dipole ordering. Anomalies in lowfrequency (f = 1-10 kHz) dielectric constant at magnetic transition have been reported in YMnO<sub>2</sub> [2] BiMnO<sub>2</sub> [3] etc. However, measurement in MHz range has been hardly reported. In this work, we report temperature, frequency and magnetic field dependence of impedance in Pr. Sr. MnO. which shows ferromagnetic transition near room temperature (T = 303 K). It is found that with increasing frequency of the accurrent an anomaly in the form of a sharp peak appears at the Curie temperature. A second anomaly in the form of an abrupt decrease occurs at T= 100 K much below the Curie temperature. These anomalies are suppressed under external magnetic field. It is suggested that the observed anomalies are not necessarily caused by ordering of electrical dipoles. The observed features in impedance closely correlate with the magnetization. We suggest that impedance spectroscopy in these materials can probe magnetization dynamics, which is often neglected in studies of magnetocapacitance effect in multiferroic materials.

[1] S.W. Cheong and M. Mostovoy, Nat.Mater. 6 (2007) 13. [2] S. Lee, A. Pirogov, M. Kang, K.H. Jang, M. Yonemura, T. Kamiyama, S.W. Cheong, F. Gozzo, N. Shin, H. Kimura, Y. Noda, and J.G. Park, Nature (London) 451, (2008) 451. [3] R. Seshadri and N. A. Hill, Chem.Mater. 13, (2001) 2892.

### **OA19**

# Structures and magnetic properties of $Tm_{1-y}Y_yMn_{1-x}Co_xO_3$

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The structure and magnetic properties of Tm<sub>1.v</sub>Y<sub>v</sub>Mn<sub>1.v</sub>Co<sub>v</sub>O<sub>3</sub> with 0≤x≤0.5 and  $0 \le y \le 0.3$  were investigated by X-ray diffraction, specific heat and magnetization measurements. A thulium manganite TmMnO<sub>3</sub> prepared by the solid state synthesis method in ambient pressure is hexagonal and antiferromagnetic with Neel temperature  $T_N$ =86 K. The substitution of Y for Tm in TmMnO<sub>3</sub> does not largely affect the fundamental hexagonal structure. We can roughly understand the results of the magnetization and the specific heat of Tm1-vYvMnO3 in terms of a dilution effect of Tm by Y. On the other hand, the structure of TmMn<sub>1-x</sub>Co<sub>x</sub>O<sub>3</sub> gradually changes from hexagonal to orthorhombic with the substitution of Co for Mn and hexagonal and orthorhombic phases coexist in samples for x≤0.3, and TmMn<sub>0.6</sub>Co<sub>0.4</sub>O<sub>3</sub> become almost orthorhombic single phase. The magnetization of TmMn<sub>0.6</sub>Co<sub>0.4</sub>O<sub>3</sub> in a field of 250 Oe increases rapidly below about 60 K with decreasing temperature. The discrepancy of the zero-field-cooled (ZFC) and the field-cooled (FC) magnetization remarkablely expands below about 60 K. Moreover, temperature dependence of ZFC and FC magnetization show peaks at 40 K and 27 K, respectively. TmMn<sub>1-x</sub>Co<sub>x</sub>O<sub>3</sub> we obtained shows complicated magnetic properties.

### **OA20**

#### Complex magnetic and electric orders in multiferroic Co<sub>3</sub>TeO<sub>6</sub>

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Multiferroics, where both ferroelectric and magnetic order coexist, are quite uncommon, but are of particular interest both to understand the fundamental interactions between the two types of order as well as for the potential for practical applications. The novel metal tellurates M3TeO6, where M is a first-row transition metal, have been shown to be rich in crystalline chemistry. Ferroelectricity and magnetic-field-driven polarization have been observed. Here, we report on the results of neutron diffraction, magnetic susceptibility, specific heat, and dielectric constant measurements made on a single crystal Co3TeO6 to study the interplay between the ferroelectricity and magnetic order. A strong interplay between the order parameters of ferroelectricity and both commensurate and incommensurate magnetic order are observed. Long range incommensurate magnetic order develops below TM1=26 K, which is followed by three additional zero-field phase transitions at TM2=19.5 K. TM3=18 K. and TM4=16 K where the incommensurate order changes and commensurate order develops. In magnetic fields up to 14 T we find that the magnetic intensities and incommensurate wave vector are dramatically altered as ferroelectricity develops, with a fifth abrupt transition around 10 T. The overall behavior characterizes Co3TeO6 as a type-II multiferroic.

# **OA21**

Annealing induced colossal magnetocapacitance and colossal magnetoresistance in in-doped CdCr<sub>2</sub>S<sub>4</sub>

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Colossal magnetocapacitance (CMC) and colossal magnetoresistance (CMR) induced by annealing polycrystalline Cd<sub>0.97</sub>In<sub>0.03</sub>Cr<sub>2</sub>S<sub>4</sub> are reported. In agreement with the appearance of CMR near the Curie temperature Tc, a insulator-metal transition is observed with decreasing temperature at zero magnetic field. On the contrary, after the same annealing treatment, CdCr<sub>2</sub>S<sub>4</sub> displays typical semiconductor behavior and does not show magnetic field dependent dielectric and electric transport properties. Accordingly, the simultaneous occurrence of CMC and CMR effects implies that the CMC in the annealed  $Cd_{97}In_{0.03}Cr_2S_4$  could be explained qualitatively by a combination of CMR and Maxwell-Wagner effect

#### Critical dynamics in LiCuVO<sub>4</sub>

**OA22** 

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### **OA23**

#### High quality crystal growth and Low temperature diffuse scattering studies

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The 120 K Verwey-transition [1] in magnetite Fe<sub>3</sub>O<sub>4</sub> is the classical example for charge or- dering, but the complex low-temperature structure was unresolved during decades. Early experimental studies and recent calculations suggest ferroelectricity (FE) due to charge ordering, which would be significant if confirmed. Very recent structural refinement [2] supports the necessary polar structure. Recently, signatures of relaxor FE were observed below 40 K [3] . Specific diffuse scattering would be expected in such a case [4], best re- solved on high quality single crystals. The Verwey transition depends sensitively oxygen stoichiometry [5] and the high quality crystals are obtained by the direct synthesis in an appropriate CO/CO2 flow [6]. We have grown high quality single crystals of magnetite under the tailored growth conditions by optical floating zone method and characterized them primarily by specific heat measurement. We present the results of high energy x-rays and neutrons diffuse scattering experiments, performed to test the proposed relaxor FE. The presence of weak diffuse scattering at very low temperature is observed only by neutron scattering, and thus likely magnetic, not supporting relaxor FE. Proof of the latter would require time resolved scattering.

1.E.J.W.Verwey, Nature 144, 327 (1939). 2.Mark S.Senn et al, Nature Letter (2011). 3.F.Schrettle et al, Phys. Rev. B. 83 195109 (2011). 4.Grace Yong et al., Phys. Rev. B. 62, 14763 (2000). 5.P.Shepherd et al., Phys. Rev. B 43, 8461 (1991). 6.R.Aragon et al., J. crystal growth, 61, 221 (1983)

# **OA24**

#### Anomalous magnetodielectric and magnetostrictive effect via spin reorientation in terbium iron garnet

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We have investigated the relationship among the magnetic, dielectric and magnetostrictive properties of terbium iron garnet (TbIG). In this presentation, we report experimental observations which are extraordinary magnetostriction and magnetodielectric effects in TbIG with a non-collinear spin structure in a single unit cell. The distinct effects of magnetism on the lattice are also demonstrated by the unprecedented magnetostriction with a negative poisson ratio. We attribute the observed magnetodielectric effects and the huge magnetostriction to the magnetic field induced spin reorientation from non-collinear to collinear. In this study, we analyzed crystallographic and spin structure of TbIG in various temperatures by the Fullprof program. We investigated the effect of the non-collinear spin structure on the crystal structure of TbIG to elucidate the relation between them through analyzing neutron diffraction data.

### OA25

#### Template based synthesis of multiferroic BiMnO3 nanotubes and shape dependent study of its magnetic properties

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The synthesis of Bismuth manganite nanotubes was done by templating precursors of bismuth and manganese on titanium dioxide nanotubes. The yield was post treated to remove the template material and sintered at 400c. The resulting powders were studied using X-ray Diffractometer and Scanning Electron Microscope, which revealed the formation of perovskite nanotubes . The BiMnO3 nanotubes were characterized for its magnetic properties by Vibrating sample magnetometer that revealed the intrinsic ferromagnetic nature with increase in its magnetic saturation that could have been due to the shape anisotropy of the nanotubes. The magnetization under field cooled and zero-field cooled was also found to differ significantly with applied magnetic field. These results indicate the co-existence of antiferromagnetic behavior at higher temperatures along with its intrinsic ferromagnetic nature which deviate from its bulk counterpart that only exhibits ferromagnetic nature.

1. Srikala, D., Singh, V.N., Baneriee, A., Mehta, B.R.Effect of induced shape anisotromy on magnetic properties of ferromagnetic cobalt nanocubes (2010) Journal of Nanoscience and Nanotechnology, 10 (12), pp. 8088-8094. 2. Tajiri, T., Harazono, M., Deguchi, H., Mito, M., Kohno, A., Kohiki, S. Synthesis and magnetic property of multiferroic BiMnO3 nanoparticles in the pores of mesoporous silica (2010) Japanese Journal of Applied Physics, 49 (6 PART 2), pp. 06GH041-06GH044. 3. Xu, E., Qian, T., Zhang, G., Zhang, T., Li, G., Wang, W., Li, X. Fabrication and magnetic properties of multiferroic BiFeO3 nanotube arrays (2007) Chemistry Letters, 36 (1), pp. 112-113.

#### **OA26**

#### Solitary reentrant superconductivity prediction in asymmetrical ferromagnet-superconductor-ferromagnet trilayer

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The theory of proximity effect for thin bilayer FS, trilayer FSF and fourlayered system FSFS, where F is a ferromagnetic metal, and S is superconductor, is proposed on the base of new boundary-value problem for the Eilenberger function. For all systems the dependencies of critical temperature on an exchange field of the F metal, electronic correlations in the S and F metals, and thicknesses of layers F and S are derived. It is shown that the possibility of the Fulde-Ferrell-Larkin-Ovchinnikov (FFLO) state observation is especially increased in the asymmetrical systems (FSF' and FSF'S') for which solitary reentrant superconductivity is predicted. On the basis of a proximity effect we propose new method of probe of electronic parameters of contacting materials. If well known BCS superconductor S is used as a probe, one can determine the exchange field, the electron-electron constant in various magnetics F for the FS structures. It allows us to predict the sign and value of the constant of electron-electron interaction in gadolinium and to explain a surprisingly high critical temperature Tc ~ 5K in the short-periodic Gd/La superlattice [1].

P.P. Deen, et al. J. Phys.: Cond. Matt. 17, 3305 (2005)

# **OA27**

#### Tuning magnetic order, electromagnons and exchange bias by epitaxial strain in BiFeO3 thin films

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<sup>1</sup> Unite Mixte de Physique CNRS/Thales, France <sup>2</sup>Universite de Rouen, France Universite Paris Diderot, France <sup>4</sup> Ecole Centrale Paris, France University of Arkansas, USA 6 Universiv of South Florida, USA

With a ferroelectric Curie temperature of 1100 K and a Neel temperature of 640K, BiFeO3 is one the very few room-temperature multiferroics [1]. Its ferroelectric polarization is as high as 100  $\mu$ C/ cm<sup>2</sup> and in the bulk it is a G-type antiferromagnet, with a superimposed cycloidal modulation. In addition to showing magnetoelectric coupling between its ferroic orders, BiFeO3 also exhibits coupling between dynamic lattice and spin excitations (electromagnons), enabling an electrical control of spin waves [2]. We have explored the influence of epitaxial strain [3] on the antiferromagnetic properties of BiFeO3 thin films. Combining Mossbauer spectroscopy and Raman scattering, we have found that the harmonic cycloidal order is stable at low strain but is replaced at higher strain by conventional antiferromagnetism. This is accompanied by strong changes in the electromagnon spectra and on their dependence on external stimuli. In addition, the mean antiferromagnetic vector progressively rotates as strain varies, offering novel strategies to control exchange bias with ferromagnetic overlayers. [1] G. Catalan et al, Adv. Mater. 21, 2463 (2009) ; [2] P. Rovillain et al, Nature Mater. 9 975 (2010) ; [3] I.C. Infante et al. Phys. Rev. Lett. 105, 057601 (2010)

#### **OB01**

#### Magnetic-enhanced electron-phonon coupling and vacancy effect in '111' type iron pnictides from first-principles calculations Mei Liu<sup>1</sup> and Bin Li<sup>2</sup>

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We investigate the lattice dynamics in both nonmagnetic and striped antiferromagnetic states of '111' type iron pnictides using the density-functional perturbation theory. The anisotropic spin-phonon effects lead to the softening of phonon frequencies and the renormalization of electron-phonon matrix elements, so as to increase the electron-phonon couplings by 68%, 65%, and 0.9% for LiFeAs, NaFeAs, and LiFeP, respectively. Taking into account the vacancy effects in the three iron pnictides, we find that the electron-phonon couplings are further enhanced by hole doping, yielding superconducting transition temperatures close to their experimental values.

#### **OB02**

#### Penetration depth and knight shift in iron-based superconductor Ba1. <sub>x</sub>K<sub>x</sub>Fe<sub>2</sub>As<sub>2</sub>

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<sup>4</sup> Department of Physics, Chiba University, Japan

<sup>5</sup> AIST, Japan

We have observed temperature and magnetic field dependence of penetration depth  $\lambda$ in single crystalline samples of  $Ba_{1,x}K_xFe_2As_2$  (x = 0.27, 0.64, 0.7, 0.8, 0.94 and 1.0) in order to clarify the superconducting gap mechanism. We have found that T dependence of  $1/\lambda^2$  in KFe<sub>2</sub>As<sub>2</sub> with H//c and H<sup>⊥</sup>c are completely different. While the T dependence of 1/2/2 observed with H//c is well fitted by a two-full-gap model, T-linear behavior is clearly observed below ~T\_c/2 with H<sup>1</sup>c, suggesting the existence of line node. These results are well interpreted when we take into account the horizontal line node (line node in ab-plane) in KFe2As2. At the presentation, we will also report the T and H dependence of  $\lambda$  and Knight shift over a wide x region of Ba<sub>1-x</sub>K<sub>x</sub>Fe<sub>2</sub>As<sub>2</sub>.

# **OB03**

# Raman scattering study of the lattice dynamics in LiFeAs and Fe1+vTe1-

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Following the recent discovery of superconductivity in F doped LaFeAsO with Tc of 26 K, numerous families of the Fe-based superconductors such as REFeAs(O, F.) (RE = rare earth), MFe<sub>2</sub>As<sub>2</sub> (M = Ba, Ca, Sr, K, ...), LiFeAs/NaFeAs and Fe<sub>1</sub>, Te<sub>1</sub> "Se, have been found and investigated. Most of these compounds have closely related phase diagrams and present an antiferromagnetic parent compound at low temperature which turns superconductor upon doping. A noticeable exception is LiFeAs, in which superconductivity is observed already in the parent compound. Recent ARPES results suggest strong electron-phonon coupling in LiFeAs, but so far no experimental study of its lattice dynamics is available. Here, we present first Raman light scattering investigations of lattice dynamics on superconducting LiFeAs. Five of the six expected phonon modes are observed, and show a conventional anharmonic behavior. No clear evidence for coupling of electron with zone center phonon is observed. Then, we present a systematic study of the lattice dynamics in Fe<sub>1+v</sub>Te<sub>1-v</sub>Se<sub>v</sub> as a function of Se and excess Fe contents, with special emphasis on the Fe phonon. A peculiar doping dependence of the mode is observed, which suggests an original coupling of this phonon to the low energy spin excitations induced by excess iron.

### **OB04**

### Magnetism and Superconductivity in Rb<sub>x</sub>Fe<sub>2-v</sub>Se<sub>2</sub>

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The recent discovery of superconductivity in  $A_x Fe_{2-y}Se_2$  (A = alkali metal), with transition temperatures up to Tc ~ 32 K, has led to a renewed interest for iron based chalcogenide systems. A remarkable observation is that, beside the superconducting state, a strong antiferromagnetic (AFM) state with magnetic moments up to  $3.3 \mu B$ per Fe ion is observed below TN~550 K. The key issue in A.Fe, Se, is whether AFM order and superconductivity cohabit due to phase separation or share a microscopic coexistence. We have performed muon spin relaxation/rotation measurements in single crystalline samples of superconducting (SC) Rb<sub>0.8</sub>Fe<sub>1.6</sub>Se<sub>2</sub> (Tc = 29 K) and non-SC Rb<sub>0.8</sub>Fe<sub>1.6</sub>Se<sub>2</sub>, which doesn't show superconductivity, in order to elucidate whether there is difference of magnetic state between SC and non-SC samples. In SC sample, we found that SC volume fraction is only ~6% and the rest of amount is AFM ordered phase, suggesting SC and AFM ordered phases are separated. We found that muon spin depolarization increases with decreasing temperature in non-SC sample, while no change was observed in SC sample. It suggests the appearance of new magnetic state below 40 K in non-SC sample.

# **OB05**

#### Magnetic Resonant mode in the Spin-Excitation Spectrum of Superconducting Rb2Fe4Se5 Single Crystals Iitae Park

Max-Planck-Institute for Solid State Research, Germany



# **QB06**

# Magnetic field-induced superconductivity in the canted antiferromagnet $Eu(Fe_{0.81}Co_{0.10})$ , As,

Vinh Hung Tran<sup>1</sup>\*, T A Zaleski<sup>2</sup>, Z Bukowski<sup>1</sup>, L M Tran<sup>1</sup> and A J Zaleski<sup>1</sup> <sup>1</sup> Institute of Low Temperature and Structure Research, Polish Academy of Sciences, 50-950 Wroclaw, Poland <sup>2</sup> Institute of Low Temperature and Structure Research, Polish Academy of Sciences,

50-950, Poland

We present results of ac-magnetic susceptibility, dc-magnetization, specific heat and magnetotransport measurements on a single crystal of magnetic  $Eu(Fe_{0.81}Co_{0.19})_2As_2$  superconductor. The compound undergoes multiple phase transitions; from paramagnetic to a spin-density-wave (SDW) at TSDW = 80 K, to canted antiferromagnetic (C-AF) state at TN = 16.5 K and superconducting (SC) state at Tc = 5.15 K at zero field. Upon applying fields both the C-AF and SC states evolve in an unconventional manner. Magnetic field distinctly affects the spin canting, resulting in separation of the C-AF into two a new C-AF and ferromagnetic phases. The unusual behavior of the SC state deserves the observation that Tc can be increased up to 6.5 K with magnetic fields applied parallel to the ab-plane of crystals. The observed field-induced superconductivity is interpreted as a result of a weakening of the orbital pairbreaking effect. From the experimental data we propose the field-temperature phase diagrams for  $Eu(Fe_{0.81}Co_{0.19})_{AS_{2.5}}$ .

**OB07** 

# Superconductivity and spin fluctuations in $Ca_{1-x}Pr_xFe_2As_2$ superconductors studied by 75As NMR

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The observation of superconductivity in the rare earth doped CaFe<sub>2</sub>As<sub>2</sub> with Tc as high as 47 K is still not understood. Here we report the first NMR study on the newly discovered Ca<sub>1,x</sub>Pr<sub>x</sub>Fe<sub>2</sub>As<sub>2</sub> (x=0.075, 0.15) single crystals. The magnetization shows superconductivity at ~25 K with a small volume ratio for both samples. A first order structural transition from the tetragonal to the collapsed tetragonal phase in both crystals is shown by 75As spectra. Residual tetragonal phase is seen in the x = 0.075 sample with a finite volume, but not seen in the x = 0.15 sample. We also observed a sharp drop of Knight shift below Ts in both samples, suggesting a suppression of local moments in the collapsed phase. However, the 1/75T1 is enhanced below 50 K and shows a prominent peak at about 25 K. Such behaviors indicate strong low-energy antiferromagnetic spin fluctuations, with static or quasi-static magnetic ordering at low temperatures, which is probably caused by Pr<sup>3+</sup> doping. The absence of the residual tetragonal phase and the observation of antiferromagnetic spin fluctuations in the x=0.15 sample suggest that superconductivity in Ca<sub>1,x</sub>Pr<sub>x</sub>Fe<sub>2</sub>As<sub>2</sub> is not tied to the tetragonal phase and may have a magnetic origin.

#### **QB08**

# Importance and details of the spin excitation spectra in high-Tc pnictide superconductors

Tanmoy Das and A. V. Balatsky

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One common feature that underlies all high-Tc superconductor families is the development of a magnetic resonance peak at or around the antiferromagnetic 'hot-spot' vector. This feature has proven to give valuable information about the pairing symmetry, and bulk Fermi surface properties. We provide a random-phase approximation (RPA) based BCS susceptibility calculation using material specific band structure for 122 family of ironpnictide and layered iron selenide superconductors. (i) In iron-pnictide, we show that when two hole-pockets at Gamma point and the electron pockets at M point are present on the Fermi surface, the spin-excitation spectra split into two distinct resonance branches. [1] The energy separation between the two resonances is related to the differences in the superconducting gaps on two hole-pockets. Again, the momentum resolution between the two modes comes from the differences in Fermi surface areas of the two hole-pockets. (ii) We also show that when one hole-pocket disappears, one of the resonances disappears. Furthermore, when both hole pockets vanish, the magnetic scattering between the remaining electron-pockets lead to a change in the gap symmetry from  $s\pm$  to d-wave. The latter is observed in the newly discovered iron-selenide superconductors. [2] Finally, we will present the resulting spin-excitation spectra for both materials

[1] T. Das, and A. V. Balatsky, Phys. Rev. Lett. 106, 157004 (2011). [2] T. Das, and A. V. Balatsky, Phys. Rev. B 84, 014521 (2011).

# QB09

#### Angular dependence of the resistive upper critical field of an ironbased superconductor Fe(Te,Se) in high magnetic fields

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The 11-system FeCh (Ch=S, Se, Te) superconductors are of great importance in understanding the mechanism of superconductivity in iron-based superconductors owing to their simple structures. Previously, we have reported the temperature dependence of the resistive upper critical field (Hc2) on Fe<sub>105</sub>Te<sub>0.85</sub>Se<sub>0.15</sub>, which exhibits superconductivity at T = 14 K [1,2]. The Hc2 of this compound at low temperature is considerably smaller than that expected from the Werthamer-Helfand-Hohenberg model, manifesting the Pauli limiting behavior. It suggests that this compound shows the spin-singlet pairing in the superconducting state. The anisotropy coefficient of Hc2 decreases from 2.4 near Tc to ~1 (nearly isotropic) at low temperatures. In this study, we have investigated an angular dependence of Hc2 of a single crystal of Fe<sub>105</sub>Te<sub>0.85</sub>Se<sub>0.15</sub> in pulsed high magnetic fields up to 52 T. At T = 10 K, the Hc2 increases with increasing 0, which is the angle between the c-axis and applied magnetic fields direction. We have found that the angular dependence of Hc2 is well agreement with the Gintzbrug-Landau effective-mass prediction.

[1] T. Kida et al., J. Phys. Soc. Jpn. Vol.78 (2009) 113701. [2] T. Kida et al., J. Phys. Soc. Jpn. Vol.79 (2010) 074706.

#### **QB10**

# Ab initio evidence of strong correlation and large Mott proximity in iron-based superconductors

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Recently discovered iron-based superconductors have attracted much interest in their high superconducting critical temperatures (Tc). Although it is believed that electron correlations mediate the unconventional high-Tc superconductivity, their roles are not fully understood yet. To clarify electron correlation effects from a microscopic point of view, we study the ab initio low-energy effective model for iron-based superconductors by using multi-variable variational Monte Carlo method. From this ab initio calculations, we show that iron-based superconductors discovered near d6 configuration (5 Fe 3d orbitals filled by 6 electrons) is located on the foot of an unexpectedly large dome of correlated electron matter centered at the Mott insulator at d5 (namely, half filling). The d5 Mott proximity extends to subsequent emergence of incoherent metals, orbital differentiations due to the Mott physics and Hund's-rule coupling, followed by antiferromagnetic quantum criticality, in quantitative accordance with available experiments.

#### OB11

# As-NQR study of LaFeAsO<sub>1-x</sub>F<sub>x</sub>

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<sup>2</sup> Beijing National Laboratory for Condensed Matter Physics, Institute of Physics, Chinese Academy of Science, China

The discovery of superconductivity in LaFeAsO<sub>1</sub> F<sub>2</sub> at superconducting transition temperature Tc=26K, followed by that in ReFeAsO, F.(Re:Ce.Pr.Nd.Sm) with Tc as high as 55K, has gained much attention. We have performed systematic Asnuclear quadrupole resonance(NOR) measurements on LaFeAsO1, F, to elucidate its superconducting gap structure and mechanism of the Cooper pair formation For x=0.03 we find the spin-lattice relaxation rate(1/T1) exhibits a small unturn at TN=58K, below which the spectra become broadened due to the internal magnetic field attributed to an antiferromagnetic ordering. Above TN, 1/T1T increases with decreasing T due to the antiferromagnetic spin fluctuation(AFSF), which persists in the x=0.04,0.06, and 0.08 compounds. A dome-shaped x dependence of Tc is found with the highest Tc=27 K at x=0.06 which is realized under significant AFSF With increasing x further. AFSF decreases, and so does Tc. These features resemble the cuprates La2-xSrxCuO4, which suggests that AFSF is also important in producing high Tc superconductivity in LaFeAsO<sub>1</sub>, F. . In the superconducting state, we find that 1/T1 decreases exponentially below Tc down to 0.13Tc for x=0.06, which suggests that the energy gaps are fully opened. Meanwhile, 1/T1 below Tc decreases nonexponentially for x either smaller or larger than 0.06, which is accounted for by impurity scattering.

# QB12

Interplay between 3d- and 4f-electrons in ReFe<sub>1-x</sub>Co<sub>4</sub>AsO (Re = Ce,Gd) T. Shang<sup>1</sup>, L. Yang<sup>1</sup>, Y. Chen<sup>1</sup>, L. Jiao<sup>1</sup>, J. L. Zhang<sup>1</sup>, J. Chen<sup>1</sup>, H. Q. Yuan<sup>1\*</sup>, N. Cornell<sup>2</sup>, A. Howard<sup>2</sup>, A. Zakhidov<sup>2</sup>, M. B. Salamom<sup>2</sup>, F. Ronning<sup>3</sup>, E. D. Bauer<sup>3</sup> and J. D. Thompson<sup>3</sup>

<sup>1</sup> Department of Physics and Center for Correlated Matter, Zhejiang University, China
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In order to study the evolution from a high Tc superconductor to a strongly correlated system with largely enhanced effective mass in iron pnictides, we have synthesized a series of ReFe1. "Co"AsO (x=0-1) polycrystalline samples and characterized their physical properties by means of electrical resistivity, magnetization and specific heat. It is shown that, for CeFe1xCoxAsO, the SDW transition of the 3d-electrons is quickly suppressed upon substituting Fe with Co and superconductivity appears in a narrow doping range of 0.05<x<0.20. On the other hand, the 4f-electrons of Ce is antiferromagnetically ordered with a nearly unchanged TN~4K over a wide doping concentration (0≤x≤0.8). Remarkably, the 3d- and 4f-electrons become ferromagnetically ordered above x~0.75 and x~0.80, respectively. Similar results are derived for the GdFe1-xCoxAsO upon partially substituting Fe with Co. However, the 4f-electrons are antiferromagnetically ordered over the entire doping concentration and likely undergo a spin reorientation above x~0.6. On the Co-side (x<0.8), the 3d-electrons undergo a FM transition, followed by an AFM transition upon decreasing temperature. These results suggest that the interplay of 3d- and 4f-electrons may play an important role in ReFe1. xCoxAsO, giving rise to rich physical properties attributed to the the Kondo scattering and the polarization between the 3d-and 4f-electrons.

#### QB13

### Anisotropic Hc2 curves determined up to 92 T and two-band superconductivity in $Ca_{10}(Pt_4As_8)((Fe_{1-x}Pt_x)_2As_2)_5$ superconductor Eundeok Mun<sup>1</sup>, Ni Ni<sup>2</sup>, Jared M Allred<sup>2</sup>, Robert J Cava<sup>2</sup>, Oscar Ayala<sup>1</sup>, Ross D

Mcdonald<sup>1</sup>, Neil Harrison<sup>1</sup> and Vivien S Zapf<sup>4</sup> <sup>1</sup> Nationa High Magnetic Field Lab, Los Alamos National Lab, USA <sup>2</sup> Department of Chemistry, Princeton University, USA

The upper critical fields, Hc2(T), of single crystals of the novel superconductor  $Ca_{10}(Pt_4As_8)((Fe_{1,x}Pt_x)_2As_2)_5$  with x=0.02 were determined over a wide range of temperatures down to T = 1.42 K and magnetic fields up to H = 92 T. The measurements of anisotropic Hc2(T) curves are performed in pulsed magnetic fields using radio-frequency contactless penetration depth measurements for magnetic fields up to 92 T and temperatures down to 1.42 K, used in this study, enable access to the complete evolution of the Hc2(T) curves. Whereas a clear upward curvature in Hc2(T) along H][c is observed with decreasing temperature, the Hc2ab(T) along H][ab shows a flattening at low temperatures. The rapid increase of the Hc2(T) suggests that the superconductivity can be described by two dominating bands. The anisotropy parameter, Hc2ab(T) / Hc2c(T), is ~7 close to Tc and decreases considerably to ~1

#### QB14

# Effect of Ni-doping on superconductivity and magnetism in $Eu_{0.5}K_{0.5}Fe_2As_2$

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# QB15

75As NMR/NQR study of hole-doped superconductor  $Ba_{1\text{-}x}K_xFe_2As_2$ 

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We have performed 75As NMR and NQR on single crystalline Ba<sub>1-x</sub>K<sub>x</sub>Fe<sub>2</sub>As<sub>2</sub> for x = 0.27-1. The Knight shift K in normal state shows Pauli paramagnetic behavior with slight temperature T dependence. The value of K increases gradually with increasing x. By contrast, nuclear spin- lattice relaxation rate 1/T1T in normal state has a large T-dependence, which indicates existence of large antiferomagnetic (AF) spin fluctuations for all x. The T-dependence of 1/T1T shows a gap-like behavior below approximately 100 K for 0.6 < x < 0.9. These behaviors are well explained by the change of band structure with expansion of hole Fermi surfaces and shrink and disappearance of electron Fermi surfaces at Brillouin zone (BZ) with increasing x. The T dependence of 1/T1 in SC state indicates multiple-SC-gap feature. A simple two gap model analysis shows that the larger superconducting gap gradually decreases with increasing x from 0.27 to 1 and smaller gap decreases rapidly and nearly vanishes for x > 0.6 othere the electron pockets in BZ disappear.

### QB16

# Homogeneous coexistence and phase segregation in the 1111 iron-based pnictides studied via NMR

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In iron-based pnictides, one of the interesting topics is homogeneous coexistence or phase segregation between antiferromagnetic (AF) and superconducting (SC) states, because these phenomena are deeply related to the pairing symmetry. We addressed this problem on a microscopic level by means of 75As NMR measurements in LaFeASO<sub>14</sub>,  $F_4$  (La1111 series), and CaFe<sub>1x</sub>Co<sub>2</sub>As<sub>7</sub> (Ca1111 series) having an intermediate electronic phase diagram between Ba(Fe<sub>1x</sub>Co<sub>2</sub>)<sub>A</sub>As<sub>2</sub> and the La1111 series [1]. The spectra for 6%-Co Ca1111 samples exhibited 75As signals similar to the undoped samples even below Tc [1], whereas those for 2.6%-F La1111 samples exhibited signals arising from both AF and SC phases at 3.0 GPa. At ambient pressure, we observed signals arising from AF and paramagnetic metallic (PM) states. These results indicate that homogeneous coexistence between AF and SC states is realized in the Ca1111 series, whereas SC or PM domains are formed other than AF domains in the La1111 series. The separation of two phases at ambient pressure indicates that the SC phase is not directly related to the AF phase and some pairing mechanism other than spin fluctuation would be important.

[1] T. Nakano, S. Tsutsumi, N. Fujiwara, S. Matsuishi, and H. Hosono, Homogeneous coexistence of superconducting and spin-density-wave states in CaFe1-xCoxAsF as seen via nuclear magnetic resonance, Phys. Rev. B 83 (2011) 180508(R)

# QB17

#### NMR study of Fe-based superconductors K<sub>x</sub>Fe<sub>2-v</sub>Se<sub>2</sub>

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 $K_xFe_{2,y}Se_2$  with a transition temperature Tc of 32K is an unique example among Fe-based superconductors.[1] Angle-resolved photoemission spectroscopy studies suggest that the electronic structure is distinct from other Fe based superconductor in the absence of hole-like Fermi surface around  $\Gamma$  point. [2]On the other hand, the diffraction studies suggest a phase separation of an antiferromagnetic (AF) phase and a superconducting phase.[3] We measured nuclear spin-lattice relaxation time T1 and Knight shift using 77Se-NMR in a single crystal  $K_xFe_{2,y}Se_2$ . Their anisotropy for H// ab and H//c including the field dependence were measured in a wide temperature range down to the SC state. In the SC state, the field dependence of 1/T1 was weak, suggesting that the relaxation is dominated by the SC quasiparticle and the contribution from the AF phase and the phase boundary is weak. Knight shift in the normal state is almost isotropic, but it consists of the spin part and the orbital part. From a measurement at low temperatures, we found that the orbital part is anisotropic, and hence the anisotropic spin part. We will discuss the SC gap symmetry and the character of spin fluctuations in K\_Fe\_2Se.

J. Gio et al. Phys. Rev. B 82, 180520(R) (2010) [2] I.I.Marzin et,al Phys. Rev. B 84, 024529 [3]
 W.Bao et.al Chinese Phys. Lett. 28, 086104

# **OB18**

Spectroscopy and anisotropies in the magnetic state of iron pnictides Belen Valenzuela, Maria Jose Calderon, Gladys Leon, Noel A. Garcia and Elena Bascones\*

Theory and Simulations of Materials, Instituto de Ciencia de Materiales de Madrid, Snair



Magnetic interactions in iron pnictides

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# WITHDRAWN

# **OB20**

# Te-doped K<sub>0.80</sub>Fe<sub>1.81</sub>Se<sub>2-x</sub>Te<sub>x</sub> single crystals

CTLin\*

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We report the growth a serious of Te-substituted K<sub>0.80</sub>Fe<sub>1.81</sub>Se<sub>2-x</sub>Te<sub>x</sub> (x~0, 0.09, 0.16, 0.34 and 0.55) single crystals using optical floating-zone technique under application of 8 bar of argon pressure. Study of XRD indicated that the undoped single crystals contain a significant intergrowth of two sets of the c-axis characterized by slightly different lattice constants, i.e., the phase separation phenomenon. Our results demonstrate that the partial substitution of Se by Te atoms can lead to the expansion of the c-lattice constant, whereas the phase separation phenomenon is suppressed continuously with the increase of substitution level and vanished at x~0.55. The magnetization data show the superconducting transition temperature is gradually depressed for Tc ~32, 27, 25, 15, and 0 K with the substitution level x~0, 0.09, 0.16, 0.34 and 0.55, respectively

# **OB21**

#### High-pressure resonant x-ray emission study of Fe1.01Se superconductors

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The Fe 3d valence states, electron structures and spin states of iron-chalcogenide Fe101Se superconductors under pressure up to 45 GPa were probed by resonant x-ray emission spectroscopy (RXES) and and x-ray absorption spectra. The Fe 1s3p-RXES spectra of Fe. Se at ambient pressure 6 GPa and 45 GPa were measured to deduce the variation of Fe 3d valence states. The 1s3p-RXES data obtained at the Fe K edge clearly reveal that unoccupied Fe 3d states exhibit a relatively delocalized character, stemming from hybridization of the Fe 3d states with the Fe 4p orbitals The pronounced pre-edge peak at ~7112.7 eV in Fe K-edge x-ray absorption spectra of Fe101Se predominantly originates from the dipole transition of a Fe 1s electron to unoccupied Fe 3d-Se 4p hybrid bands. The larger compression accompanied with a significant distortion around the Fe atoms along the c axis in Fe1.01Se upon applying pressure suppresses the Fe 3d-Se 4p and Fe 4p-Se 4d hybridization. The applied pressure suppresses the nearest-neighbor ferromagnetic superexchange interaction and enhances spin fluctuations on the Fe sites in Fe101Se. Fe1.01Se shows a small net magnetic moment of Fe2+ at ambient pressure.

#### **OB22**

Microscopic coexistence and competition of magnetism and superconductivity in Ba1-xKxFe2As2 : A structural, magnetic, and superconducting phase diagram

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It is widely believed that in contrast to its electron-doped counterparts, the hole-doped compound Ba1-xKxFe2As2 exhibits a mesoscopic phase separation of magnetism and superconductivity in the underdoped region of the phase diagram. Here, we report a combined high-resolution X-ray powder diffraction and volume-sensitive muon spin rotation (µSR) study of Ba1-xKxFe2As2 showing that this paradigm does not hold true in the underdoped region of the phase diagram ( $0 \le x \le 0.25$ ). Instead, we find a microscopic coexistence of the two forms of order. A competition of magnetism and superconductivity is evident from a significant reduction of the magnetic moment and a concomitant decrease of the magnetoelastically coupled orthorhombic lattice distortion below the superconducting phase transition. The doping dependence of structural, magnetic, and superconducting properties is completed in a phase diagram ( $0 \le x \le 1$ ) herein presented and discussed.

[1] E. Wiesenmaver, H. Luetkens et al., Phys. Rev. Lett. 107, 237001 (2011), [2] M. Rotter et al., Nature Physics 5, 141 (2009). [3] A. A. Aczel et al., Phys. Rev. B 78, 214503 (2008). [4] T. Goko et al., Phys. Rev. B 80, 024508 (20

### **OB23**

#### Anomalous superconducting phase in LaFeAsO<sub>1-x</sub>H<sub>x</sub> studied via 75As NMR

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LaFeAsO<sub>1-x</sub>F<sub>x</sub> ( La1111 series ) is a prototype of high-Tc pnictides and has been investigated extensively by means of various experimental techniques. The electronic phase diagram is unique compared to other high-Tc pinctides in that the antiferromagnetic (AF) phase hardly overlaps the superconducting (SC) phase. and the optimal doping level is located away from the AF phase boundary. The latter feature is common to the R1111 series ( R = Ce, Nd, etc. ). Recently, an exotic phase diagram has been obtained in LaFeAsO1, H, in which a high Tc is maintained even in an overdoping regime, and the second SC dome emerges upon excess H doping. To investigate the exotic electronic state on a microscopic level, we performed 75As NMR measurements for 20% and 40% doped samples. The former doping level corresponds to the boundary between two SC domes, and the latter corresponds to the optimal doping level of the second dome. The relaxation rates for 20% doped samples are comparable with those of LaFeAsO<sub>1-x</sub> $F_x$  (x = 14%), however rates for 40% doped samples enhanced at a fixed temperature implying that some spin fluctuations return and some magnetic ordering would occur upon excess H doping.

# **OB24**

# Multi-frequency ESR in EuFe<sub>2</sub>As<sub>2</sub>

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One kind of the iron based superconductors is called '122' system. Most parent materials of the '122' system exhibit a spin-density wave (SDW) transition that is accompanied by a structural phase transition [1]. EuFe2As2 shows the SDW transition at TSDW=190 K, and the large local moment Eu2+ (S =7/2, L =0) orders antiferomagnetically below TN=19 K. Above TSDW, it is believed that Eu<sup>2+</sup> localized magnetic moments interact with itinerant electrons. Below TSDW, the interlayer interaction turns to be weak and the system becomes anisotropic [2]. To study the above natures of EuFe2As2 in more details, we have performed multi-frequency (20-35 GHz) electron spin resonance (ESR) measurements. Above TN, the resonance signal can be fitted with the Dysonian function, and the temperature dependence of the linewidth obtained by this fitting obeys the Korringa relation above TSDW as observed at X-band ESR [2]. From this relation, we will discuss the exchange interaction between Eu<sup>2+</sup> moments and conduction electron spins.

[1]Y. Xiao et al., Phys. Rev. B 80 (2009) 174424. [2]E. Dengler et al., Phys. Rev. B 81 (2010)

### **OB25**

#### Analysis of the critical current density and flux pinning properties in iron-based Ba0.55K0.45Fe2As2 high Tc superconductor

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The critical current density and flux pinning properties of the Ba055K045Fe2AS2 high Tc superconductor with  $T_c = 37$ K are studied using field and time dependence of the magnetization. The temperature dependence of the critical current density  $J_c(T)$ have been found to follow the form  $J_c \propto (1 - T/T_c)^n$  and is consistent with  $\delta T_c$  type pinning due to the randomly distributed defects larger than the coherence length  $\xi$ . The temperature dependence of the magnetic relaxation rate S exhibits a peak at much lower temperature which corresponds to the smaller pinning energy in the Ba<sub>055</sub>K<sub>045</sub>Fe<sub>2</sub>As<sub>2</sub> superconductor. Moreover, the effective pinning potential U<sub>0</sub> increases with the increase in temperature this behavior is in agreement with those of high  $T_c$ cuprate superconductors.

### **OB26**

#### In-plane anisotropy of magnetic and electric properties of the Fe pnictide Ba(Fe<sub>1-x</sub>Co<sub>x</sub>)<sub>2</sub>As<sub>2</sub>

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The anisotropy of electronic and magnetic properties within FeAs-planes of Co(2%)doped BaFe2As2 was investigated by 75As-NMR measurements of the As sites As0 with all four nearest neighbor sites occupied by Fe. Interestingly, even in the tetragonal phase (T >100 K), two sets of NMR spectra with two fold symmetry of in-plane Knight shifts and electronic nuclear quadrupole frequencies were observed. They originate from the existence of two domains with their symmetry axes lying at rightangles to each other. The anisotropy of the electronic and magnetic properties become pronounced on cooling at ~140 K, below which the in-plane anisotropy was reported to appear for the electrical resistivity under the pressure along one of Fe-Fe directions [1]. 75As-NMR was also carried out at the As site As1 surrounded by one Co and three Fe. From the data, we can see that the spin susceptibility at As1 is ~1/3 of the value at As0 and that the antiferromagnetic spin fluctuation at As1 is fairly suppressed as compared with the value at As0. We discuss these behaviors of the FeAs layer in connection with the impurity-induced local orbital ordering suggested by the theoretical studies.

[1] J.-H. Chu, J. G. Analytis, K. D. Greve, P. L. McMahon, Z. Islam, Y. Yamamoto and I. R. Fisher :Science 329 (2010) 824.

190 | The 19<sup>th</sup> International Conference on Magnetism

**OB19** 

# **OB27**

#### Vortex tunneling spectra of iron-pnictide superconductors

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Symmetry and structure of the superconducting gap functions in iron-pnictides have been still under spirited debate. Quasi-particle states around surfaces/interfaces or impurities have played a crucial role for identifying the gap functions in several unconventional superconductors such as cuprates or Sr<sub>2</sub>RuO<sub>4</sub>. However, the gap symmetry of iron-pnictide is believed to be s-wave, and the point at issue is to identify whether the gap structure has inter-band sign reversal (s+-) or not (s++). Thus, it will be difficult to distinguish these gap structures based on surfaces/interfaces information. There exist several studies on impurity-induced states in iron-pnictides. However, because of the complex band-structures of such systems, theoreticalmodeling of impurities in iron-pnictide is not free of ambiguity. Vortex core states can be another "probe" to identify the gap symmetries and structures in unconventional superconductors. Since a vortex core is not a chemical substance, there are no ambiguities in theoretical modeling. Thus, in this work, we theoretically study the quasi-particle states induced around a vortex core in iron-pnictides using two- or fiveband effective two-dimensional lattice models [1][2] within the Bogoliubov-de Gennes theory. We will also examine the impurity-induced quasi-particle states, and discuss the relation between present results and recent STM/S experimental results.

[1]T. Kariyado and M. Ogata, J. Phys. Soc. Jpn. 79, 083704 (2010). [2]S. Raghu et al., Phys. Rev. B

# **OB28**

#### Direct observation of superconducting gaps and their anisotropies in Ba<sub>1-x</sub>K<sub>x</sub>Fe<sub>2</sub>As<sub>2</sub>

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<sup>0</sup> ISSP. JST-CREST. JST-TRIP. RIKEN. Japan

Among the iron-pnictides,  $(Ba,K)Fe_2As_2$  (BaK122) system is the most interesting: while the optimally-doped BaK122 has a full gap (no node in the superconducting(SC) gap), the extremely hole-doped KFe\_2As\_2 (K122), which has no electron pocket, has been suggested to have SCgap nodes from several experiments. From the previous studies of laser-excited angle-resolved photoemission spectroscopy (laser ARPES), the optimally doped BaK122 has no anisotropy and Fermi surface (FS) sheet dependence in SC-gap sizes on the three hole FS around the Brillouin zone center. On the other hand, we found that K122 has clear anisotropies and sheet dependence. In addition, we found that the middle hole FS has eight nodes and identified their Fermi momenta. It should be important to investigate doping dependence of anisotropies and FS-sheet dependence of SC-gap sizes, and node existence to clarify the paring mechanism at each doping level. We have directly observed superconducting gaps and their anisotropies in the overdoped BaK122 utilizing laser ARPES. We will discuss these results and contrast them with other experimental and theoretical results available in literature

# **OB29**

#### Contrasting superconducting property in Fe-based superconductors (Ca<sub>4</sub>Al<sub>2</sub>O<sub>6-v</sub>)(Fe<sub>2</sub>Pn<sub>2</sub>)[Pn=As and P]

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We report As-nuclear quadrupole resonance (NQR) study on (Ca<sub>4</sub>Al<sub>2</sub>O<sub>6-v</sub>)(Fe<sub>2</sub>As<sub>2</sub>) with Tc=27K, and P-nuclear magnetic resonance (NMR) study on (Ca<sub>4</sub>Al<sub>2</sub>O<sub>6-v</sub>)(Fe<sub>2</sub>As<sub>2</sub>) with Tc=17K. (Ca<sub>4</sub>Al<sub>2</sub>O<sub>6-y</sub>)(Fe<sub>2</sub>As<sub>2</sub>) possesses the characteristic structural parameters such as short a-axis length, high pnictgen height, narrow As-Fe-As angle, and thick perovskite-type blocking layer. A measurement of the nuclear spin relaxation rate 1/T1 revealed a significant evolution of antiferromagnetic (AFM) spin fluctuations in normal state, which originates from the possible well nested hole and electron Fermi surfaces. Below Tc, the 1/T1 decreases steeply upon cooling without any trace of Hebel-Slichter peak, which is consistently accounted for within the framework of nodeless s(+-)-wave multiple gap model. On the other hands, (Ca<sub>4</sub>Al<sub>2</sub>O<sub>6-v</sub>)(Fe<sub>2</sub>As<sub>2</sub>) is characterized by lower pnictgen height and wider As-Fe-As angle than that of (Ca<sub>4</sub>Al<sub>2</sub>O<sub>6+</sub>)(Fe<sub>2</sub>As<sub>2</sub>), where FeP\_4 forms nearly regular tetrahedron. The P-NMR- 1/T1 measurement revealed the presence of AFM spin fluctuations in normal state and gapless SC state. These facts indicate the close relationship between the local structural characters of FePn 4 and the SC properties.

H. Kinouchi, H. Mukuda, M. Yashima, Y. Kitaoka, P. M. Shirage, C-H. Lee, H. Eisaki, and A. Iyo

# **OB30**

#### Superconducting properties of FeTe<sub>1-x</sub>Se<sub>x</sub> single crystals: impact of disorder and hydrostatic pressure

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It was found that chemical and structural disorder in superconducting single crystals of Fe(Te,Se), originating from kinetics of crystal growth process, strongly influences superconducting properties. The sharpness of a transition to the superconducting state is evidently inversely correlated with crystallographic quality. Jons inhomogeneous spatial distribution seems to enhance the superconductivity and inhomogeneous distribution of host atoms might be an intrinsic feature of superconducting chalcogenides. Small disorder introduced into magnetic sublattice, by partial replacement of Fe ions by slight amount of nonmagnetic ions of Cu or by magnetic ions of Ni and Co with spin value different than that of Fe ion, completely suppresses superconductivity. Even if superconductivity is observed in the system containing magnetic ions it can not survive when the disorder in magnetic ions sublattice is introduced. We have found pressure-induced enhancement of all of the superconducting state properties of  $FeTe_{0.5}Se_{0.5}$  single crystals, which entails a growth of the density of superconducting carriers. We noticed more pronounced increase in superconducting carrier density under pressure than that in the critical temperature what may indicate an appearance of a mechanism limiting the increase of Tc with pressure. The critical current density increases under pressure by at least one order of magnitude.

# **OB31**

#### The influence of superconductivity with magnetism in superconductor/ magnetic heterostructures

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Interplay between superconductivity and magnetism continues to be one of the central themes in Condensed Matter Physics. We have investigated a variety of superconductor/ferromagnetic heterostructures, such as YBCO/LSMO, Nb/GCMO, and Nb/CeRu<sub>2</sub>Al<sub>2</sub>B using low temperature magnetic force microscopy, addressing mainly magnetic pinning and domain wall superconductivity. In YBCO/LSMO, we focused on the critical state and observed effects of magnetism on superconductivity. In Nb/GCMO, we found that magnetic domains provide weak pinning centers, and act as magnetic templates. A recently discovered CeRu<sub>2</sub>Al<sub>2</sub>B system shows a variety of magnetic phases such as antiferromagnetic, ferromagnetic, and ferrimagnetic domains in the field/ temperature plane. We will discuss how we can manipulate superconducting domains in superconducting film (Nb) via these magnetic phases in close proximity

### **OC01**

#### The influence of the magnetic moment on the atomic distance in amorphous CexRu<sub>100-x</sub>

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The extended x-ray absorption fine structure (EXAFS) measurement was performed in amorphous alloys CexRu<sub>100-x</sub> (x=9, 43 and 80). We found that the Ce-Ce interatomic distance increases with increasing the Ce content while those of Ce-Ru and Ru-Ru pairs are unchanged and that the Ce-Ce interatomic distance is exactly proportional to the Ce effective magnetic moment. We discuss the relation between the Ce-Ce atomic distance and the effective magnetic moment 1) from the point of view of the

1) Y. Amakai et al., J. Phys. : Conf. Ser. 150 (2009) 042004

# OC02

Theoretical studies of the superconductivity and antiferromagnetism coexistence and the divergence of effective electron mass near quantum critical point in CeRhIn-

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Competition and coexistence of superconductivity and antiferromagnetism have been studied in the framework of the periodic Anderson model taking into account the superexchange interaction between localized 4f electrons. It has been shown that pressure-induced energy change of the localized states leads to modifying antiferromagnetic and superconducting order parameters. This modification of the order parameters develops in such way that a strong competition is suggested between antiferromagnetism and superconductivity. It has been proposed that the energy increase is caused by pressure growth. In spite of the competition mentioned above, conditions have been found for the coexistence of superconductivity and antiferromagnetism in the model. Theoretical results of the pressure effects on the ground state character of the heavy-fermion systems are in good agreement with specific heat data for CeRhIn, [1, 2]. In the coexistence phase there occur two types of Cooper pairing due to the antiferromagnetic exchange field. The divergence of the effective electron mass produced by suppressing antiferromagnetic order has also been analyzed in the same framework. Further it has been shown that at quantum critical point there occurs the transition from small to large Fermi surface. Such features have been found in the de Haas-van Alphen experiments for CeRhIn; [3].

1. H. Hegger, C. Petrovich, E.G. Moshopoulou, M.F. Hundley, J.L. Sarrao, Z. Fisk, J.D. Thompson Phys. Rev. Lett. 84, 4986 (2000). 2. T. Park, J.D. Thompson, New Journal of Physics 11, 055062 (2009). 3. H. Shishido, R. Settai, H. Harima, Y. Onuki, J. Phys. Soc. Jap. 74, 1103 (2005).

#### **OC03**

#### Josephson effect between UPt<sub>3</sub> and Nb under pressure

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The heavy-fermion compound UPt3 is an odd-parity superconductor and possesses a complex field-temperature (H-T) phase diagram with three phases: A (low H and high T), B (low H and low T) and C (high H and high T) phases. We have investigated the Josephson effect between UPt<sub>1</sub> and Nb and observed an anisotropic temperature dependence of Josephson critical current Ic at the transition from A to B phases[1]. However, the Josephson effect, which is easily destroyed by a small magnetic field, has not been measured in C phase. Considering that C phase appears even in zero field by applying pressure[2], we have investigated the pressure dependence of Josephson effect together with resistance and Meissner effect of UPt. The superconducting transition occurs at lower temperature and becomes broader by applying pressure. As the pressure is increased, the Josephson effect appears at lower temperature than the zero-resistivity temperature of UPt, while Ic increases more rapidly; the increasing rate -dlc/dT is a linear function of pressure. Although the Josephson effect in C phase has been observed for the first time, an abrupt change in the temperature dependence of Ic caused by the transition to C phase has not been found yet.

[1] A. Sumiyama et al.: Phys. Rev. Lett. 81 (1998) 5213; A. Sumiyama et al.: Jpn. J. Appl. Phys. ries 11 (1999) 47. [2] M. Boukhny et al. : Phys. Rev. Lett. 73 (1994) 1707.

#### **OC04**

#### Hidden order in URu<sub>2</sub>Si<sub>2</sub> --- Analysis based on the first-principles approach

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The so-called 'hidden-order' phase transition at \$T {HO}=17.5\$K in the heavy-electron metal URu2Si2, specified by a clear jump in the specific heat measurement, is a long-standing mystery in the condensed-matter physics since the discovery in 1985. In spite of every effort, it remains still dark what kind of symmetry is broken below \$T {HO}\$. So far, there are no evidence for magnetic/structural and any low-rank multipole orders. Recent advanced angle-resolved photo-emission spectroscopy indicates that \$5f\$ electrons appears to be itinerant rather than localized. Thus, the detailed analysis of the complicated band structures including the spinorbit coupling has been strongly desired. We show here the first report of the complete set of itinerant multipole correlations up to the fifth rank within the random-phase approximation and bevond. The itinerant Hamiltonian is obtained by the advanced Wannier method from the firstprinciples calculations. From the most divergent correlations, we obtain nearly degenerate two staggered-ordered states of rank-1 dipole with \$A 2^-\$ and rank-5 dotriacontapole with \$E^-\$; the former corresponds to the antiferromagnetic state realized under high pressures, and the latter the hidden-order state. This novel high-rank electronic state provides natural explanations of the key features for the hidden order including spontaneous fourfold symmetry breaking.

# **OC05**

#### Sb NQR study of filled skutterudite CeFe<sub>4</sub>Sb<sub>12</sub> synthesized under high nressure

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Among filled skutterudites, the influence of the filling fraction of R ion on the physical properties has been argued, especially in Sb-based filled skutterudites with a large pnictogen cage. On RFe<sub>4</sub>Sb<sub>12</sub> system, the drastic effect of the R-site vacancy on the basic properties has been pointed out. Recently, it has been reported that CeFe<sub>4</sub>Sb<sub>12</sub>synthesized under high pressure (HP) in order to increase the Ce-site filling fraction shows the semiconducting transport property at low temperatures, in contrast that the sample synthesized under ambient pressure (AP) shows the metallic behavior. At the presentation, we report on the result of the Sb NQR measurements for the HP sample of CeFe<sub>4</sub>Sb<sub>12</sub>, and compare with that for the AP sample in order to understand the influence of the Ce-site filling fraction. In the high temperature range, the nuclear spin-lattice relaxation rate for the HP sample decreases on cooling, and follows the exponential decrease above 100 K with the energy gap of 270 K. The value of the gap for the HP sample is larger than that for the AP sample, which suggests that the hybridization is enhanced by increasing the Ce-site filling fraction. More detailed results and analyses will be presented at the conference.

#### **OC06**

#### High-mobility magnetotransport of the narrow-gap semiconductor FeSb,

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Recently, a huge Seebeck coefficient S (-45 mV/K at 10 K) was reported for singlecrystalline FeSb<sub>2</sub> samples. Such a huge value seems beyond conventional semiconductor physics, indicating a novel mechanism to enhance the Seebeck coefficient such as a strong electron correlation expected in so-called 'Kondo insulators'. We study the magnetotransport properties of FeSb<sub>2</sub> single crystals to investigate its electronic state at low temperatures[1]. The conductivity tensor is well analyzed using a two-carrier model, leading to proper evaluation of the mobility and carrier concentration. The mobility significantly increases with decreasing temperature and reaches 28000 cm2/Vs at 4 K. The carrier concentration decreases from 1 down to 10-4 ppm/unit cell with decreasing temperature from 30 down to 4 K, which is well described by thermal activation with a small gap of 6 meV. We also find that the magnetoresistive behavior dramatically changes from positive to negative at lower temperatures. The negative magnetoresistance is generally observed in doped semiconductors and derives from a weak localization of carriers, indicating a considerable impurity level in the electronic structure. Our magnetotransport measurements, as well as the impurity-effect study[2], strongly suggest that the low-temperature transport is well understood within an extrinsic semiconductor picture with ppm-level impurity.

[1] H. Takahashi et al. Phys. Rev. B 84, 205215 (2011). [2] H. Takahashi et al. J. Phys. Soc. Jpn. 80 054708 (2011)

### **OC07**

#### A possible ferromagnetic quantum critical point in CeFe<sub>1-x</sub> Ru<sub>x</sub>PO Tetsuro Nakmura<sup>1</sup>, Takashi Yamamoto<sup>2</sup>, Masanori Matoba<sup>1</sup>, Yasuaki Einaga<sup>2</sup> and Yoichi Kamihara

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Many materials have their own magnetic transition temperature, and it is controllable by the application of a pressure, field or through chemical substitution. In particular, phase transitions which take place at absolute zero are called quantum phase transitions, and the transition point is called a quantum critical point. Quantum critical points have attracted much attention since some materials exhibit superconductivity in proximity to a magnetic quantum critical point. [1] Both CeFePO and CeRuPO are heavy-fermion materials containing Ce 4f-electron's sublattice inducing Kondo effects. Although CeFePO exhibits no magnetic transition down to 0.4 K, CeRuPO is one of rare examples of a ferromagnetic Kondo lattice, whose transition temperature is 15 K [2]-[5]. In this study, we synthesized CeFe<sub>1-x</sub>Ru<sub>x</sub>PO by a solid reaction to quantify a chemical composition which is a boundary between ferromagnetic phase and paramagnetic phase indicating the quantum critical point. Electrical resistivity and magnetic susceptibility are measured at temperatures from 4 K to 300 K. Furthermore, 57Fe Mossbauer spectroscopy was applied to CeFe1-Ru,PO to confirm the ordering of the magnetic sublattice. Our results tentatively indicate that CeFe<sub>1-x</sub>Ru<sub>x</sub>PO is heavy-fermion metal, and the magnetic quantum critical point exists at x~0.15.

1) A. J. Millis: Phys. Rev. B 48 7183 (1993). 2) E. M. Bruning, C. Krellner, M. Baenitz, A. Jesche, F. Steglich, and C. Geibel: Phys. Rev. Lett. 101 117206 (2008). 3) S. Kitagawa, H. Ikeda, Y. Nakai, T. Hattori, K. Ishida, Y. Kamihara, M. Hirano, and H. Hosono: Phys. Rev. Lett. 107 277002 (2011). 4) Y. Kamihara, H. Hiramatsu, M. Hirano, H. Yanagi, T. Kamiya and H. Hosono: J. Phys. Chem. Solids. 69 2916 (2008). 5) C. Krellner, N. S. Kini, E. M. Bruning, K. Koch, H. Rosner, M. Nicklas, M.Baenitz, and C.Geibel: Phys. Rev. B 76 104418 (2007).

# **OC08**

#### Ru doping evolution of magnetic properties in Ce(Fe<sub>1-x</sub>Ru<sub>x</sub>)PO studied by 31P-NMR

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CeFePO and CeRuPO have the same lavered crystal structure as iron pnictide superconductors RFeAs(O, F)(R: rare earth). CeFePO is a paramagnetic down to 80 mK and possesses a large Sommerfeld coefficient gamma 700 mJ/mol K2 at low temperatures, indicating that CeFePO is a heavy fermion compounds originated from Ce 4f electrons[1]. On the other hand, CeRuPO is a ferromagnetic Kondo lattice system with TCurie = 15 K and coherent temperature TK ~ 10 K[2]. We have performed 31P-NMR on Ce(Fe1-Rux)PO system to investigate the evolution of magnetic properties at low temperatures. Knight shift and spin-lattice relaxation rate 1/T1 shows extremely large anisotropy. In CeFePO, we found a metamagnetic transition at around muOH = 4 T and below 5 K characterized by the sudden increase in Knight shift when fields are applied perpendicular to the c axis[3]. As far as we know, this is the first example of the metamagnetic behavior in a heavy fermion observed with XY spin anisotropy. Moreover,  $Ce(Fe_{125}Ru_{125})PO$  shows ferromagnetic ordering at ~ 3 K with strong two dimensional ferromagnetic fluctuations. We will show the phase diagram in Ce(Fe1, Ru,)PO and discuss how two dimensional ferromagnetic fluctuations are changed with the Ru doping

[1] E. M. Bruning et al., Phys. Rev. Lett. 101, 117206 (2008). [2] C. Krellner et al. J. Cryst. Growth 310, 1875 (2008). [3]S.Kitagawa et al., Phys. Rev. Lett. 107, 277002 (2011).

# OC09

# Determining the orientation of the 4f ground state orbital in CeCu-Si-

with vector q dependent non-resonant inelastic X-ray scattering (NRIXS). Thomas Willers<sup>1</sup>\*, Fabio Strigari<sup>1</sup>, Yong Cai<sup>2</sup>, Nozomu Hiraoka<sup>3</sup>, Ku-ding Tsuei<sup>3</sup>, Yen-fa Liao<sup>3</sup>, Maurits Wim Haverkort<sup>4</sup>, Silvia Seiro<sup>5</sup>, Christoph Geibel<sup>5</sup>, Frank Steglich<sup>5</sup>, Liu Hao Tjeng<sup>5</sup> and Andrea Severing<sup>1</sup> Institute of Physics II, University of Cologne, Germany

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Max Planck Institute for Chemical Physics of Solids, Dres

WITHDRAWN

### **OC10**

#### Possible superconducting fluctuation in pressure-induced heavyfermion superconductor CeRhSi<sub>3</sub>

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CeRhSi<sub>2</sub> is a heavy fermion antiferromagnet whose transition temperature is 1.6K. The itinerant f electron is thought to be responsible for the magnetism in this material. Application of pressure suppresses the magnetism and simultaneously induces the superconductivity. At pressures from 0.2GPa to 2.4GPa, the superconductivity coexists with the antiferromagnetism. In this pressure region, a broad drop of the electrical resistivity and weak shielding in the ac-susceptibility are observed at the superconducting transition temperature. We argue a possibility of the superconducting fluctuation realized in the coexisting region of the superconductivity and antiferromagnetism

# **OC11**

Thermoelectric power in a single-crystalline CeRhSi<sub>3</sub> Hidekazu Tanaka<sup>1</sup>\*, Naofumi Aso<sup>2</sup>, Yoshinao Takaesu<sup>1</sup>, Masato Hedo<sup>2</sup>, Takao

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CeRhSi3 is a non-centrosymmetric heavy fermion superconductor.[1] The a- and c-axis thermoelectric power S\_a and S\_c of a single-crystalline CeRhSi3 were measured under pressure up to around 2.7 GPa in the temperature range from 1.5 K to 300 K. Both S a and S c exhibit a large positive value up to around 60 µV/K, which is characteristic of heavy-fermion systems. S\_a and S\_c show a maximum in its temperature dependence due to the Kondo effect and its maximum temperature T\_Smax gradually moves to higher temperatures from about 90 K (at ambient pressure) to about 130 K (at 2.7 GPa) by applying the pressure. In both directions, such behaviors of thermoelectric power are very similar to those in the CeIrSi<sub>3</sub>[2] In the conference, we will discuss the anisotropy of the thermoelectric power in connection with the superconductivity in CeRhSi.

[1] N.Kimura et al., Phys. Rev. Lett. 95 (2005) 247004. [2] Y.Takaesu et al., J. Phys. Soc. Jpn. 80 (2011) SA068

# **OC12**

# Valence of CeM2Al10 (M=Ru, Os, and Fe) determined with hard X-ray

photo emission spectroscopy (HAXPES). Fabio Strigari<sup>1</sup>, Thomas Willers<sup>1</sup>, Ku Ding Tsuei<sup>2</sup>, Yen Fa Liao<sup>2</sup>, Arata Tanaka<sup>3</sup>, K. Yutani<sup>3</sup>, Y. Muro<sup>3</sup>, Toshiro Takabatake<sup>3</sup>, Liu Hao Tjeng<sup>4</sup> and Andrea Severing<sup>1</sup>\* University of Cologne, Germany

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# **OC13**

### Studies on novel tetragonal Ce2RhGa12 heavy fermion compound

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Single crystals of Ce<sub>2</sub>RhGa<sub>12</sub> have been synthesized using Ga flux and their structure was determined by single-crystal X-ray diffraction. The Ce<sub>2</sub>RhGa<sub>12</sub> crystallizes in the tetragonal P4/nbm space group (No. 125), which is isostructural to Ce<sub>2</sub>PdGa<sub>12</sub>, with Z = 2 and lattice parameters a 6.0405 A and c= 15.706 A. Data were collected at the Swiss Norwegian Beam Line at the European Synchrotron Facility, Grenoble France. Also the Laue diffraction was carried out to confirm the single crystal quality possessing tetragonal symmetry. The tetragonal structure of Ce<sub>2</sub>RhGa<sub>12</sub> is composed of Ce-Ga and PdGa8/2 layers similar to CePdGa<sub>6</sub>[1]. The structure is isomorphous to Ce2MGa12 M=(Ni, Cu[2], Pd). Along the c axis slabs of edge-sharing rectangular prism of RhGa8 alternates with Ga rich areas in which Ce atoms occupy cavities made of Ga atoms. The bulk anti ferromagnetic transition of Ce<sub>2</sub>RhGa<sub>12</sub> at 3.5 K is exhibited by a big jump up to 12 J/mol K. The value of  $\gamma$  above magnetic transition in specific heat measurements suggests that it is a heavy-fermion system with  $\gamma$  = 420 mJ mol-1 K-2. The magnetic measurements are in progress to confirm the consistency of the ordering temperature and the further studies are in progress.

I. Robin T. Macaluso, S. Nakatsuji, H. Lee, Z. Fisk, M. Moldovan, D.P. Young, and Julia Y. Chana Journal of Solid State Chemistry 174, 296-301(2003) 2. J.Y. Cho, J.N. Millican, C. Capan, D.A. Solokov, M. Moldovan, A.B. Karki, D.P.Young, M.C. Aronson, J.Y. Chan, Chem. Mater. 20, 6116-6123(2008)

# **OC14**

#### Electronic structures of plutonium compounds with the NaCl-type monochalcogenides structure

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Some of the uranium and the transuranium compounds are characterized as strongly correlated f electron systems and have attracted particular interest because of rich variety of anomalous physical phenomena such as the valence fluctuation, the Kondo lattice, exotic magnetism, unconventional superconductivity, etc. Among such compounds, the plutonium monochalcogenides display a variety of anomalous physical properties. For PuX (X=S, Se, Te) have been studied intensively both theoretically and experimentally[1,2]. To understand the electronic properties of these materials, the features of the ground state should be first clarified. The calculations for the energy band structure have been done by using the relativistic linear augmented-plane-wave (RLAPW) method. Note here that relativity should be taken into account, because of arge atomic numbers of the constituent atoms. In the presentation, we try to understand what the key issues are to construct the energy band structures around the Fermi energy for PuS. PuSe and PuTe.

[1] A. Hasegawa, and H. Yamagami, J. Magn. Magn. Mater. 104-107, 65 (1992). [2] P. M. Oppeneer, T. Kraft, and M. S. S. Brooks, Phys. Rev. B61, 12825 (2000).

### OC15

#### Electronic property of ThSn<sub>3</sub> in comparison with uranium and transuranium compounds

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Recent rapid expansion of the research frontier of condensed matter physics from uranium to transuranium compounds has stimulated renewed attention and much interest on exotic properties of actinides and related compounds. It is a challenge to modern electronic-structure theory to understand the appearance of a large number of exceptionally complicated crystal structures of actinide metal. It is important to clarify electronic structure of actinide compounds which exhibit exotic magnetism and unconventional superconductivity. By using a relativistic linear augmentedplane-wave (RLAPW) method with the one-electron potential in the local-density approximation, we investigate energy band structures and the Fermi surfaces of transuranium compounds ThSn<sub>3</sub>, USn<sub>3</sub>, NpSn<sub>3</sub> and PuSn<sub>3</sub>. It is found in common that the energy bands in the vicinity of the Fermi level are mainly due to the large hybridization between 5f and Sn 5p electrons. Thorium contains no occupied 5f states, thorium compounds provide good comparative systems for investigating the role of 5f electrons. In the presentation, we try to understand what the key issues are to construct the energy band structures around the Fermi energy for ThSn<sub>3</sub>, USn<sub>3</sub>, NpSn<sub>3</sub> and PuSn<sub>3</sub>, we attempt to unveil 5f electron properties purely originating from actinide atoms.

### **OC16**

#### Multipolar phase transition of the 4f2 nonmagnetic doublet in a caged compound PrRh<sub>2</sub>Zn<sub>20</sub>

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The caged compound  $PrIr_2Zn_{20}$  with the  $Pr^{3+}$  ion in the  $\Gamma 3$  doublet ground state exhibits a superconducting transition at TSC=0.05K in the presence of an antiferroquadrupolar (AFQ) order below TQ=0.11K [1, 2]. The entropy release up to TQ is only 20% of Rln2 which is expected for the  $\Gamma$ 3 doublet with the quadrupolar degrees of freedom. It suggests that quadrupolar degrees of freedom remain active even below TQ and quadrupole fluctuations play an important role in forming the Cooper pair in the superconducting state. In the present work, we have focused an isostructural and isoelectronic compound PrRh<sub>2</sub>Zn<sub>20</sub>. The electrical resistivity shows a hysteretic behavior between 140K and 470K, suggesting a firstorder structural transition. In spite of the possible structural transition, a Schottky-like peak at around 10K in the specific heat and anisotropic behavior of the isothermal magnetization at 1.8K indicate the crystalline electric field ground state to be the F3 doublet. The specific heat exhibits a peak at TQ=0.06K, where the entropy release is evaluated to be only 6% of Rln<sub>2</sub>. Because of the increment of TQ in the magnetic fields and the anisotropic B-T phase diagram, the phase transition at TO should be an AFO order

[1] T. Onimaru et al., J. Phys. Soc. Jpn. 79 033704 (2010). [2] T. Onimaru et al., Phys. Rev. Lett. 106 177001 (2011)

# **OC17**

#### Crystal growth and low temperature properties of non-centrosymetric heavy-fermion comound CeTAl<sub>3</sub> (T = Cu, Ag, Au)

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For a long time full inversion symmetry was thought to be a precondition for heavyfermion superconductivity. The discovery of heavy-fermion superconductivity in the antiferromagnetic state of CePt, Si at ambient pressure and CeRhSi, CeIrSi, and CeCoGe<sub>3</sub> at high pressure contradicted these paradigms. All compounds in this class crystallize in a P4mm or subclass crystal structure. Here we report on the single crystal growth of CeTAl<sub>2</sub> (T = Cu, Ag, Au) with crystal structure I4mm, being isostructural to CeRhSi<sub>3</sub> and CeIrSi<sub>3</sub>. High purity elements (>5N) have been used as starting materials. A special build image furnace was used for crystal growth. Crystal quality was determined by Laue x-ray scattering, x-ray powder diffraction and edx. Low temperature properties have been investigated by measurements of bulk properties including magnetisation, heat capacity and electrical resistivity.

# **OC18**

#### Enhancement of the hall coefficient under pressure in CeCu<sub>2</sub>Si<sub>2</sub>

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CeCu<sub>2</sub>Si<sub>2</sub> is the heavy fermion superconductor discovered in 1979. The unconventional superconductivity with Tc ~ 0.6 K at ambient pressure is believed to be mediated by antiferromagnetic spin fluctuation around the AFM quantum critical point. With increasing pressure, Tc increases to reach the maximum at 4.2 GPa, although antiferromagnetic spin fluctuation decreases. The existence of the critical valence fluctuation around the critical valence transition is proposed to explain the enhancement of Tc under pressure [1]. However, there is no direct evidence to reveal the existence of the valence fluctuation. In this study, we investigated the Hall effect on single crystal of CeCu<sub>2</sub>Si<sub>2</sub> under high pressure. The temperature dependence of the Hall coefficient at ambient pressure shows a peak at TmL = 4 K for H // [001] and [100]. In addition, the Hall coefficient shows the other peak at T1max for H // [100]. The magnitude of the Hall coefficient at TmL increases with increasing pressure and reaches the maximum at around 4 GPa, coincident with the enhancement of Tc. We will discuss the relation between the enhancement of the Hall coefficient and Tc

[1] K. Mivake et al.: J. Phys.: Condens. Matter 19 (2007) 125201.

# **OC19**

#### Saturation moment in the ferromagnetic state of EuB<sub>6</sub>

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Two consequtive phase transitions featuring the onset of ferromagnetism in europium hexaboride [14] are commonly associated with percolation in the system of linked magnetic polarons at  $T_M \sim 15.6$  K and with long range magnetic ordering below T\_C=12.6-13.9 K [1]. However, the recent study of high frequency electron spin resonance [5] showed neither magnetic phase separation nor MP formation features in the system of localized magnetic moments of Eu<sup>+</sup>. To shed more light on the nature of ferromagnetism we studied the magnetic and thermal properties of the EuB<sub>6</sub> single crystals in pulsed (up to 40 T) and steady (up to 9 T) magnetic fields at temperatures 2-40 K. The estimated saturation moment  $M_{sat}$ -7.4 $\mu$  B is found to exceed noticeably the magnetic moment of Eu<sup>2+</sup> ion. The difference is shown to be induced by an additional contribution to magnetization due to the formation of intermediate magnetic phase below T\_\_\_\_\_M. The observed renormalization of M {sat} is discussed in terms of the effects of the spin polarization in 5d-band and the short range magnetic ordering in the system of Eu<sup>2+</sup> ions. Support from RFBR 11-02-00623 project, Federal Programme "Scientific and Educational Human Resources of Innovative Russia" and the Methusalem Funding by the Flemish Government is acknowledged.

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 V.V.Glushkov et al., JETP v.105, p.32 (2007). [4] J.Kim et al., Phys Rev B, v.78, p.165120 (2008). [5] A.V.Semeno et al., Phys Rev B, v.79, p.014423 (2009).

# **OC20**

#### Magnetization of Tm<sub>1-x</sub>Yb<sub>x</sub>B<sub>12</sub> in strong pulsed and steady magnetic fields

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The magnetization of solid solutions Tm1-xYbxB12 (0<x\leq 0.8) crossing metal-insulator transition point at x\_{c}~0.3 have been studied at low temperatures 1.8-40 K in strong steady (up to 11T) and pulsed (up to 50T, pulse duration ~20-100 ms) magnetic fields. The analysis of the magnetization data allows extracting three contributions to the paramagnetic response. Itinerant electron components are attributed to heavy fermion and ferron type manybody states at EF with densities of states N(E<sub>F</sub>)= 1\div 4?10<sup>22</sup> cm<sup>-3</sup>meV<sup>1</sup> and 3\div 4?10<sup>21</sup> cm<sup>-3</sup> <sup>3</sup>meV<sup>1</sup>, correspondingly. The third contribution may be identified as the response from the localized magnetic moments of the rare earth ions (~0.8-3.7 µB per unit cell) which saturates at helium temperatures in strong magnetic fields. The reduced intensity of the manybody resonance at E<sub>F</sub> is shown to increase with Yb concentration reaching 2 electrons per unit cell for x>0.7. Comparison between the steady and pulsed field magnetization behavior provides additional arguments in favor of magnetic nanoclusters formation, short range ordering and electron phase separation effects in Tm1-xYbxB12.

#### OC21

#### Stabilization of ferromagnetism and existence of ferromagnetic quantum criticality in UCo1.,Ru,Ge system

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We have prepared polycrystalline samples of UCoGe doped by Ru on the Co-site about nominal stoichiometry UCo<sub>1,x</sub>Ru<sub>x</sub>Ge where x = 0.01-0.35 with step 0.01 to studied the magnetic state and ferromagnetic instability in UCoGe. The whole series of doped materials keeps the orthorhombic TiNiSi-type structure. The evolution of the lattice parameters well follows Vegard's law respecting the pure parent UCoGe and URuGe. While the UCoGe is well known ferromagnetic superconductor (TC = 2.5 K, TSC = 0.6 K) the URuGe is paramagnetic down the lowest temperatures. Nevertheless the ruthenium doping of UCoGe leads to surprising stabilization of the ferromagnetic state when the maximum TC = 8.5 K has been found in 12 % Ru doped sample. The further increase of the Ru content destabilizes previous ferromagnetic state and it is terminated by ferromagnetic quantum critical point at concentration of 30 % of Ru. We have prepared single crystal of the UCo<sub>0.88</sub>Ru<sub>0.12</sub>Ge and UCo<sub>0.70</sub>Ru<sub>0.30</sub>Ge for detail study of the ferromagnetic state, individual contributions of the U, Co and Ru components to total magnetic moment, magnetic anisotropy and destabilization of the ferromagnetism in UCoGe in relation to superconductivity. The results are discussed on the basis of the change of electronic structure.

# **OC22**

#### Thermal properties of RB<sub>6</sub> (R-La, Ce, Pr, Nd)

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Rare earth (RE) hexaborides RB 6 represent a class of compounds with variety of magnetic properties including antiferromagnetic (R-Ce, Pr, Nd, Gd, Dy, Tb, Ho), ferromagnetic (EuB 6), intermediate valence (SmB 6) and heavy fermion behavior (CeB 6). Here we investigate the heat canacity C P(T) on the high quality single crystals of light hexaborides (R-La Ce, Pr, Nd) in the wide range of temperatures 2-300K. The data obtained allow to estimate correctly (i) the electronic component with v(LaB 6)~24mJ/(mol•K). (ii) contribution from quasilocal vibrational mode of R^{3+} ion with O E~119 - 152K, (iii) the Debye-type term from rigid boron cage with O\_D~1160K. Our data also suggest the additional (iv) defect mode component which is provided by the contribution of the boron vacancies. The estimated values of residual concentration N {vac}~1-4% agrees with the results of X-ray studies of RB 6 [1], [2]. The analysis of magnetic contribution shows that magnetic entropy  $S=\Delta C \{mag\}/TdT$  reaches the value corresponding to the ground state of 4f-multiplet in the range 2.4T N Taking into account the results of previous thermal investigations [3] we conclude in favour of the spin-

polaron formation (see [4], [5] for details) in the paramagnetic vicinity of Neel temperature

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# QC23

#### Magnetic penetration depth and skin depth study of superconductivity and quantum criticality in Ce<sub>1.x</sub>R,CoIn<sub>5</sub> (R=La and Nd)

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A heavy fermion superconductor CeCoIn<sub>5</sub> shows different responses to Nd- or Lasubstitutions for Ce, with the former inducing long range magnetic order coexisting with superconductivity. To understand the origin of the differences, we studied the temperature and field dependent in-plane magnetic penetration depth, \$\lambda(T)\$, in single crystals of (Ce,R)CoIn<sub>5</sub> (R=La, Nd). Measurements were performed with a tunnel diode resonator down to 50-mK in a dilution refrigerator, in magnetic field up to 14 T parallel to the \$c\$-axis. These low-temperature and high field measurements allowed for the exploration for the full domain of superconductivity and quantum criticality in the  $T_{ii}$  phase diagram. Some previously unreported features were observed and will be discussed from the point of view of measured differential magnetic susceptibility. These measurements bring new insight into the interplay between superconductivity and magnetism as well as field-tuned quantum critical behavior of doped 115 systems.

# QC24

# Low energy spin excitations in single-crystalline $CeCu_2Ge_2$ in magnetic fields up to 10T

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CeCu<sub>2</sub>Ge<sub>2</sub> is Kondo lattice with a moderately enhanced Sommerfeld coefficient of the specific heat  $\gamma = 140$  mJ/molK2. It exhibits antiferromagnetic order below TN = 4.1 K. The incommensurate magnetic structure can be explained by the nesting properties of the calculated Fermi surface when taking into account an itinerant component of the Ce 4f-moments in addition to their local character. We performed low-energy inelastic neutron scattering experiments on a 2g single crystal of CeCu<sub>2</sub>Ge<sub>2</sub> using the cold three-axes spectrometer PANDA at FRM II. Data were taken at 1.5 K and in magnetic fields up to 10T applied perpendicular to the (110/001) scattering plane. At zero field the spin excitations show an energy gap of 0.5 meV at the magnetic zone center. Away the magnetic excitations become dispersive merging into a band of excitations around 1 meV. For increasing magnetic fields, the gap energy is lowered indicating opposite action of external and internal fields. The results will be discussed in the framework of local, crystal-field related, and non-local, spin-wave-like, magnetic excitations.

### QC25

#### Magnetization measurements under high pressure on incommensuratecommensurate phase transitions on UPd<sub>2</sub>Si<sub>2</sub>

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UPd<sub>2</sub>Si<sub>2</sub> has the ThCr<sub>5</sub>Si<sub>2</sub>-type body-centered tetragonal structure. This compound shows an incommensurate-antiferromagnetic (IC-AFM) ordering with a propagation vector Q = 0.73c\* at 135 K (= TNh), and successively undergoes a simple commensurate (C) type-1 AFM ordering at 108 K (=TNI). It is known that the phase transitions are of second order at TNh and of first order at TNI. In addition, an antiphase ordering state (Q = 2/3c\*) becomes stable in high magnetic field applied along the c-axis. The magnetism of UPd<sub>2</sub>Si<sub>2</sub> is phenomenologically well explained in terms of the one-dimensional Ising model assuming localized 5f spins, whereas Fermi liquid behavior with a moderately large electric-specific-heat coefficient ~ 30 mJ/molK2 is also seen at low temperature [1]. In the present study, we have measured magnetization of UPd<sub>2</sub>Si<sub>2</sub> under high pressure up to 5.5 GPa and high magnetic field up to 7 T in order to investigate the duality between the localized and the itinerant features of the 5f-electronic states in U compounds. From the results of the magnetization and previous resistivity measurements under high pressure [2], we will discuss about the 5f-electronic states of the compound.

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# QD01

# Fe doping and magnetic field effect in the valence fluctuating heavy fermion system α-YbAlB<sub>4</sub>

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The Yb-based heavy fermion system,  $\alpha$ -YbAlB<sub>4</sub>, has a noncentrosymmetric crystal structure, while the structure of the isostoichiometric system,  $\beta$ -YbAlB<sub>4</sub> is centrosymmetric [1].  $\beta$ -YbAlB4 is a unique example of the Yb based heavy fermion systems that shows pronounced non Fermi liquid behaviors and a superconducting transition at 80 mK under ambient pressure and field [2,3,4]. On the other hand,  $\alpha$ -YbAlB<sub>4</sub> at low temperatures is well fit to a Fermi liquid type description. Both  $\alpha$  and  $\beta$ -YbAlB<sub>4</sub> are valence fluctuating systems with the Yb valence of 2.73 for  $\alpha$ -YbAlB<sub>4</sub> and 2.75 for  $\beta$ -YbAlB<sub>4</sub> [5]. We succeeded in substituting Fe for Al in  $\alpha$ -YbAlB<sub>4</sub>, and found an antiferromagnetic order at 10 K. By a small substitution to  $\alpha$ -YbAlB<sub>4</sub>, we also succeeded in inducing non-Fermi liquid behaviors. The antiferromagnetic ordering temperature is suppressed by magnetic field and non-Fermi liquid behaviors are induced where the Neel temperature vanishes. In this presentation, we will discuss the difference of the two non-Fermi liquid behaviors in Fe doped  $\alpha$ -YbAlB<sub>4</sub> showing thermodynamic and transport properties.

[1] R. T. Macaluso et al, Chemistry of Materials 19, 1918 (2007). [2] S. Nakatsuji et al, Nature Physics 4, 603 (2008). [3] K. Kuga et al, Physical Review Letters 101, 137004 (2008). [4] Y. Matsumoto et al, Science 21, 316 (2011). [5] M. Okawa et al, Physical Review Letters 104, 247201 (2010).

#### QD02

#### Magnetic and electric properties of single crystal $\rm SmBaMn_2O_6$

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We have succeeded in growth of single crystals of an A-site ordered perovskite maganese oxide, SmBaMn<sub>2</sub>O<sub>6</sub>, and investigated the magnetic and electric properties. Since Sm and Ba ions alternately occupy the perovskite A-site, it is expected that the magnetic and electric properties should be anisotropic. Moreover, the anisotropy may be modified with the charge orbital ordering. We also found that a peak in magnetic susceptibility of polycrystals at about 250 K. The magnetic susceptibility of a single crystal shows an anomaly which is accompanied by a hysteresis between 170 K and 210 K. This anomaly is caused by a rearrangement of charge orbital ordering. A clear magnetic anisotropy is observed below this temperature indicating that this rearrangement of charge orbital ordering accompanies the simultaneous antiferromagnetic spin ordering. Temperature dependence of the dielectric constant along the c axis also shows an anomaly at the rearrangement temperature, while no clear anomaly is observed for E perpendicular to the c axis. This result is consistent with the rearrangement of the stacking of the charge- and orbital-ordered sheets along the c axis.

### QD03

# Phase diagram and Eu valence state in $EuPtP_{1-x}As_x$

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EuPtP is known to exhibit two valence order transitions at  $T_1 = 235$  K and  $T_2 = 190$  K [1]. At  $T > T_1$ , the Eu ions have a single valence state with the mean valence of 2.16. At  $T_2 < T < T_1$ , the Eu valence state is ordered with the sequence of  $-Eu^{2+}-Eu^{2$ 

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# QD04

# Investigation of crystal structure, magnetism and transport properties of $SrFe_{1,x}Ti_xO_{3,6}$ systems

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SrFeO<sub>3</sub> has attracted lot of attention in the recent past due to the exciting electric and magnetic properties that can be achieved by varying oxygen stoichiometry. In this investigation, partial replacement of Fe<sup>+</sup> with a non-magnetic on of Fi<sup>+</sup> in SrFeO<sub>34</sub> is carried out to reduce the magnetic complexity of the helical structure and determine the resulting chemical and magnetic structure using neutron powder diffraction and Mossbauer spectroscopy. Rietveld refinement of the neutron data indicated the presence of tetragonal symmetry with H/mmm space group in  $\delta$ -0.125. Our neutron and Mossbauer measurements confirm that Fe exists in three different sites both crystallographically as well as magnetically in all four compositions, oxygen deficient SrFe<sub>3</sub>, Ti<sub>1</sub>O<sub>35</sub> (x = 0 to 0.3) show an insulating behavior with a canted antiferromagnetic spin ordering, in comparison to a metallic, helical spin arrangement of x=0. Resistivity of x = 0 to 0.3 follows the Mott's 2-dimensional variable range hopping (VRH), p= p0 exp(TiOT)-1/3. The composition dependence of ac conductivity is carried out in the fragment presenting the grain contribution. The temperature coefficient fresistance (TCR) changes from negative to positive (PTCR) for x = 0.2.

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### QD05

# Evolution from a localized to an intermediate valence regime in Ce<sub>2</sub>Cu<sub>2x</sub>Ni,In

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Ce<sub>2</sub>Cu<sub>2</sub>In and Ce<sub>2</sub>Ni<sub>2</sub>In crystallize in a primitive tetragonal structure of the Mo<sub>2</sub>FeB<sub>2</sub> type. The former compound orders antiferromagnetically at the Neel temperature of 5.5 K, while the latter one is a system with fluctuating valence [1]. Here we report on low temperature physical properties of the solid solution Ce<sub>2</sub>Cu<sub>2</sub>, Ni<sub>4</sub>In studied by means of x-ray powder diffraction, magnetic susceptibility and electrical resistivity measurements, using polycrystalline specimens. We show that partial substitution of Cu by Ni results in a quasi-linear decrease of the lattice parameters and unit cell volume of the system. The lattice compression leads to an increase in the exchange integral and yields a reversal in the order of the magnetic 4f1 and nonmagnetic 4f0 states, being in line with the Doniach phase diagram. In the localized regime, where an interplay of the Kondo scattering and the crystalline electric field effect occurs, the rise in the hybridization strength is accompanied by a relative reduction in the scattering conduction electrons on excited crystal field levels.

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#### QD06

#### Substitution effect in CeFe<sub>2</sub>Al<sub>10</sub>

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CeFe<sup>2</sup>Al<sup>10</sup> is an important reference material as CeRu<sup>2</sup>Al<sup>10</sup> which indicates a mysterious phase transition. We have performed the electrical resistivity and magnetic susceptibility of substituted systems, i.e., (Ce,La)Fe<sup>2</sup>Al<sup>10</sup> and Ce(Fe,Co)Ru<sup>2</sup>Al<sup>10</sup>. The former system reveals the effect of periodicity and the latter one reveals that of electron number. A broad peak, which is attributed to Kondo effect, monotonically shifts to higher and lower temperatures with La- and Co-substitution, respectively. The electrical resistivity of CeFe<sup>2</sup>Al<sup>10</sup> shows a broad peak at 60 K and sharp increase below 15 K. A slight substitution of Ce by La and Fe by Co site remove the broad peak and the increase at low temperatures. These results indicate the broad peak is due to coherence effect and the increase is due to a semiconducting gap.

# **QD07**

Fermi surface and electronic correlations in a valence fluctuating Eusystem: an experimental and theoretical study of EuIr<sub>2</sub>Si<sub>2</sub>

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Electronic correlations can induce strong changes in electronic states, especially those at the Fermi level. These effects are presently the subject of intense theoretical and experimental investigations, especially on Ce- and U-based heavy fermion systems as well as on High Tc cuprates. In contrast valence fluctuating Eu-systems have yet not been investigated in this respect, although the valence fluctuations of Eu should have a strong impact. We have recently grown high quality single crystals of the valence fluctuating compound EuIr<sub>2</sub>Si<sub>2</sub> and started a thorough investigation of its electronic properties. From previous Moessbauer studies the valence v of Eu in EuIr<sub>2</sub>Si<sub>2</sub> was reported to decrease continuously with temperature, from v = 2.3 at 300 K to v = 2.8 below 20 K [1, 2]. Our T-dependent XRD, susceptibility, resistivity, and specific heat measurements on these single crystals confirm such a strong evolution of the valence and give insight into global effects of the strong correlations. These single crystals allowed the first observation of de Haas-van Alphen oscillations in a valence fluctuating Eu-system, as well as the first determination of the optical conductivity in such a system. Comparing these results with DFT-based calculations give a deeper insight into the correlation effects.

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# QD08

### Heat capacity and electrical resistivity of $CeNi_{5-x}Ge_x$

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Metallic Ce compounds are of great interest for the intriguing physical properties associated with the 4f electrons. One notable example of compound intensively studied for many years is CeNi<sub>5</sub> crystallizing in the hexagonal CaCu<sub>5</sub> type, which results to be a Stoner enhanced paramagnet characterized by a spin fluctuation contribution on its transport properties. Here the effect of the Ni/Ge substitution on the ground state of CeNi<sub>5</sub> is presented. Several ternary germanides CeNi<sub>5-x</sub> Gex (x=0.1, 0.2, 0.5, 0.8) were prepared by arc melting method. All the alloys crystallize in the CaCu<sub>5</sub> structure. Single phase of the polycrystalline samples was confirmed by X ray diffraction. The measurements of resistivity between 2-300 K and heat capacity between 0.4-300 K were provided and the results are presented. The resistivity measurements show typical metallic behaviour. No phase transition was observed in resistivity and heat capacity measurements.

# QD09

#### Antiferromagnetic order in Yb<sub>5</sub>Rh<sub>4</sub>Ge<sub>10</sub>

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The titled compound has three Yb sites in the tetragonal Sc<sub>5</sub>Co<sub>4</sub>Si<sub>10</sub>-type structure. The magnetization process shows strong anisotropy [1]. At 2 K, M (B/a) is saturated to 4.5 µB/f.u. in 5 T, whereas M (B/c) increases linearly up to 0.9 µB/f.u. at 5 T. The effective magnetic moment for B//a is 4.25 µB/Yb, which is close to the value of 4.53 µB/Yb expected for the Yb<sup>3+</sup> free-ion [1]. The continuous increase in the specific heat down to the lowest measured temperature 2 K suggests a magnetic order below 2 K [1]. Therefore, we have measured the temperature dependence of the magnetization M, specific heat C and resistivity down to 0.3 K. A cusp in M (B//a) and sharp peak in C(T) were observed at 0.5 K, both of which were suppressed by applying magnetic fields up to 0.8 T. These findings indicate an antiferromagnetic order at TN = 0.5 K. However, the magnetic entropy Sm at TN is only 60 % of Rln<sub>2</sub>. This observation suggests a heterogeneous mixed valence state where three of five Yb ions in formula unit are in trivalent state and the rest two are in divalent state.

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# QD10

Magnetic field dependence of the resistivity minimum of nanosized YbAl<sub>3</sub>

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Yb-based alloys follow the Doniach phase diagram where the competition between RKKY and Kondo interaction plays a predominant role. Among them, a simple cubic compound (Pm-3m) YbAl<sub>3</sub> is a fine example of strong electron correlations leading to intermediate valence behavior [1]. Recently, we have investigated the effect of crystallite size on the electronic and magnetic properties of the nanometric samples of this alloy, prepared by high-energy milling of arc-melted ingots. Milling yields massive quantities of 4f-alloys in nancorystalline (nc-) form with a relatively narrow size distribution [2]. In nc-YbAl<sub>3</sub>, the variation of valence is due to a modification of the fraction of surface atoms compared to the total number of particles. This affects both specific heat and susceptibility data [3]. Here, we present a magnetic field (H)-dependent study of the electrical resistivity, p, of annealed and pressed pellets of nc-YbAl<sub>3</sub>. The XRD Rietveld refinements yield particle sizes of maximum ~10 nm whereas slightly smaller values are detected by HRTEM. The results are striking in that the physical properties change drastically compared to the bulk [4] and an unexpected minimum in p(1) occurs around 15 K. With increasing H, the minimum becomes more pronounced and shifts to higher temperatures.

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# QD11

Electronic states of Eu<sub>4</sub>As<sub>3</sub> under high pressure

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The mixed valence Eu<sub>4</sub>As<sub>3</sub> compound with an anti-Th<sub>3</sub>P<sub>4</sub> structure shows the charge ordering at 345 K and ambient pressure accompanied by trigonal distortion, which is similar to those in Yb<sub>4</sub>As<sub>3</sub> and Sm<sub>4</sub>Bi<sub>3</sub> [1-3]. The hybridization of a narrow 4f band with broad conduction bands is one of the important parameters in these charge-ordering compounds, which is possible to be controlled by hydrostatic pressure. We observed the melting of the static charge order at about 10 GPa and room temperature without structural change using high-pressure X-ray diffraction [4]. We have carried out 151Eu nuclear forward scattering (NFS) measurements using synchrotron radiation under high pressures at low temperatures to investigate the magnetic properties of Eu<sub>4</sub>As<sub>3</sub>. Clear quantum beats due to magnetic hyperfine field were observed in 151Eu NFS spectra at 6 K under pressure up to 14.5 GPa. The results of temperature dependences of 151Eu NFS spectra reveal that a magnetic ordering temperature increases linearly with pressure. Consequently, the pressure dependence of the magnetic ordering temperature is related to the melting of the static charge order.

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### QD12

# Pressure effect on intermediate valence semiconductor SmB<sub>6</sub> : 11B-NMR

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The intermetallic compound SmB<sub>6</sub> is one of the most famous compounds in terms of the physics of intermediated valence state, heavy Fermion systems and Kondo insulators. Recent high pressure measurements of 149Sm nuclear forward scattering and resistivity revealed that a pressure-induced insulator-metal transition of SmB<sub>6</sub> is intimately connected with a long-range magnetic order, which has renewed interest in this compound. In order to investigate variations of the ground state and magnetic properties in SmB<sub>6</sub>, we have carried out 11B-NMR measurement under pressure. While our group had reported the results of 11B-NMR measurements up to 2.6 GPa with a piston-cylinder type pressure cell, the use of a Bridgman Anvil Cell has recently enabled us to perform NMR measurements at higher pressures than 4.5 GPa. Spin-lattice relaxation rate 1/T1 above 20 K roughly follows an exponential function of temperature, Exp(-Eg/2kBT), consistent with previous reports. 1/T1 between 20 and 100 K increases with increasing pressure, indicating a reduction of the insulating gap. We will discuss detailed pressure dependences of the ensy gap and the Sm valence.

### QD13

# Magnetic properties of ytterbium fluoride sulfide $Yb_3F_4S_2$

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The ternary ytterbium fluoride sulfide Yb<sub>3</sub>F<sub>4</sub>S<sub>2</sub> crystallizes in the tetragonal structure with space group I4/mmm. This structure can be described as an alternative stack of one layer of YbF<sub>2</sub> and two layers of YbFS along the c-axis. In a previous paper, it was suggested that there are two different ytterbium cations (Yb2<sup>+</sup>(Yb<sup>+</sup>)2(F)4(S<sup>2</sup>)2 present [1]. An anomaly with the long-range magnetic order has not been observed down to 0.4 K by magnetization measurements, and the magnetic susceptibility exhibits a Curie-Weiss behavior above 50 K. However, it has become clear that the estimated value of the effective magnetic moment cannot be explained by a model which assumes that Yb<sup>2+</sup> and Yb<sup>3+</sup> occupy the two different crystallographic sites. A Schottky anomaly is found in the temperature dependence of the specific heat around 1.0 K at zero magnetic field. This anomaly is shifted towards higher temperatures with increasing applied magnetic field, and a release of the residual entropy is observed simultaneously.

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#### QD14

# Physical properties of the layered oxypnictides (CeO)MnPn; Pn=As, Sh

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Recently, the layered iron oxypnictides (LaO)FePn ; Pn = P, As, Sb, which belong to the tetragonal system, have attracted much attention due to the discovery of the superconductors (LaO)FeP [1] and (LaO)FeAs [2]. Many efforts have been dedicated to explore the superconductor with the highest transition temperature. In this tidal stream, many new materials with different physical properties have been found. In case of selection of the cerium atom as a rare earth element, these materials of (CeO)FeP [3], (CeO)FeAs [4], (CeO)CoP [5], (CeO)CoAs [6] show the typical metallic properties and the Fermi liquid behaviors at low temperature. In these systems, the 3d electrons of the iron and cobalt have finite density of state at the Fermi energy, and the cerium atom provides the f electrons as carriers at the same time. Namely, it is difficult to know which effects are dominant. In our group, the semiconductor based materials (LaO)MnPn have been focused and the substitution effects of the rare earth elements on the physical properties have been investigated. (CeO)MnAs and (CeO)MnSb indicate metallic character and the Fermi liquid behaviors below 30K and 8 K, respectively due to the valence fluctuation of the Ce f electrons.

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### QD15

#### 31P-NMR study of valence fluctuating compound EuPtP

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We report the results of the first 31P-NMR measurements on the hexagonal layered compound EuPtP. This compound exhibits two valence transitions at T1= 235 K and T2= 190 K. Various experiments, including magnetization, Eu Mossbauer, and x-ray diffraction measurements [1-5], have been carried out in order to investigate the valence transitions so far, however the ordered structures below T1 and T2 have not been fully clarified. We observed 31P-NMR signals in all three phases. They show anomalies at T1 and T2, and the Knight shift of the 31P-NMR lines reach ~ 20 % at low temperatures due to large effective moments of magnetic Eu<sup>2+</sup> ions. Below T2, three resonance lines of 31P-NMR were observed, which should reflect local symmetry at the P-site. Based on these information obtained from the NMR measurements (spectrum and spin lattice relaxation time T1, and so on), we discuss possible ordered structures at low temperatures of EuPtP.

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# QD16

**Magnetic behavior of polycrystalline Eu<sub>5</sub>Si<sub>3</sub> compound** Sujata M. Patil<sup>1\*</sup>, P. L. Paulose<sup>2</sup> and E. V. Sampathkumaran<sup>2</sup> <sup>1</sup> Wilson College, Mumbai 400007, India

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The series  $R_s X_3$  (R= light rare earth and X = Ga,Si,Ge,Sn,Sb), form in  $Cr_5B_3$ -type tetragonal structure. There are two inequivalent 'R' sites in this structure, one of them showing a very short distance between rare-earth atoms. Single crystal studies of  $Ce_sSi_3$  show that metamagnetic transition temperature is different along 'a' axis and 'c' axis, which was attributed to the two inequivalent sites of Ce atoms. Studies of  $R_s X_3$  have been extensively reported for different rare earths with exception of  $Eu_sSi_3$ . Rietveld analysis of x-ray diffraction shows the formation  $Eu_sSi_3$  in the  $Cr_3Br_3$ -type tetragonal structure (space group 14/mcm). Magnetization display a prominent cusp around 32K corresponding to antiferromagnetic ordering. Zero field cooled and field cooled magnetization is different below 30K, thereby showing a complex magnetic structure. Magnetizention like that in  $Ce_sSi_3$ . Heat capacity shows a peak below 32K which shifts to lower temperatures upon application of magnetic fields, like in a typical antiferromagnet. Specific heat also shows a strong magnetic field dependence below 20K, revealing the complex nature of magnetic ordering.

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### QD17

# Valence fluctuation study by using X-ray absorption and emission spectroscopies at Yb L3-edge in YbNi<sub>3</sub>X<sub>9</sub> (X=Al and Ga)

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Quantum criticality of Ce- and Yb-based heavy fermion system has been recently discussed from a viewpoint of valence fluctuation [1], thus the estimation of valence guides in the solution for several anomalies observed near the quantum critical point (QCP). Recently, high-quality YbNi<sub>3</sub>X<sub>9</sub> (X = Al and Ga) is newly synthesized [2,3] and is desired to investigate the electronic states including the valence state as a heavy fermion system. In order to clarify the Yb valence state in YbNi<sub>3</sub>X<sub>9</sub>, X-ray absorption and emission spectroscopies at the Yb L3-edge have been performed at BL39XU of SPring-8. For YbNi<sub>3</sub>Ga<sub>9</sub>, the Yb valence of 2.71 at room temperature almost linearly decreases to 2.60 at 55 K with decreasing temperature, and then the valence maintains a constant value at 2-55 K, while the valence of YbNi<sub>3</sub>Al<sub>9</sub> is almost trivalent in a temperature range of 2-300 K. Moreover, the XES spectra for YbNi<sub>3</sub>X<sub>9</sub> give us detail information of Yb electronic states by comparison with the result in several Yb-based compounds.

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#### QD18

Angle-resolved photoemission spectroscopy of mixed-valence Sm<sub>1-x</sub>Y<sub>x</sub>S Keiichiro Imura<sup>1\*</sup>, Tetsuya Hajiri<sup>2</sup>, Masaki Kaneko<sup>2</sup>, Yusuke Nishi<sup>2</sup>, Hiroyuki S. Suzuki<sup>3</sup>, Noriaki K. Sato<sup>2</sup>, Takahiro Ito<sup>2</sup>, Masaharu Matsunami<sup>1</sup> and Shin-ichi Kimura<sup>1</sup> <sup>1</sup> UVSOR Facility. Institute for Molecular Science. Japan

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It is known that samarium mono-sulfide (SmS) exhibits an insulator-to-metal transition, so called Black-to-Golden phase transition, by the application of pressure or by the substitution of Sm ions [1,2]. Although extensive studies to investigate the mechanism of this transition as well as the ground state of golden phase have been made, the phase transition has net been perfectly clarified yet. In this study, we focus on Sm<sub>1-x</sub>Y<sub>x</sub>S system in order to investigate the electronic structure of these two phases by an angle-resolved photoemission spectroscopy (ARPES). Single crystals of Sm<sub>1-x</sub>Y<sub>x</sub>S (x = 0, 0.17, 0.33) were grown by the Bridgman technique using a high-frequency induction furnace installed in NIMS. Photoemission spectroscopy measurement on these crystals was performed at UVSOR-II BL5U end station. The obtained ARPES spectrum of pure SmS (x = 0) has an energy gap of about 0.4 eV which is consistent with a previous report [3]. This energy apd decreases with increasing x, and the top of the Sm<sup>2+</sup> 4f band touches the Fermi level in x = 0.33. This result suggests that the Golden phase induced by the ythrium substitution is a mixed-valence state.

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# **QE01**

# Critical behavior of a spin-1 triangular lattice Ising antiferromagnet

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We employ Monte Carlo simulations in order to investigate critical behavior of a frustrated spin-1 Ising antiferromagnet on a triangular lattice in the presence of a single-ion anisotropy. It has been previously found that no long-range order can exist in the system with spin 1/2 [1] but it can occur if the spin is larger than some critical value estimated as 11/2 [2]. We demonstrate that even small amount of the single-ion anisotropy can change this scenario. More specifically, we show that long-range order can exist in the low-temperature region even below this critical value, namely for spin 1, within a certain range of the single-ion anisotropy strength. At higher temperatures we identify another phase of the Berezinsky-Kosterlitz-Thouless type characterized by a power-law decay of the spin correlation function. We perform finite-size scaling analyses to calculate the correlation decay exponent. Subsequently, we estimate the critical boundaries of both phases and plot the phase diagram as a function of the single-ion anisotropy strength. This work was supported by the Scientific Grant Agency of Ministry of Education of Slovakia Republic (Grant No. 1/0234/12).

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### **QE02**

# Selectively diluted triangular lattice Ising antiferromagnet in an external magnetic field

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It is well known that the pure Ising model with nearest-neighbor antiferromagnetic interactions on a triangular lattice shows no long-range ordering due to a high degree of frustration. Dilution of one out of three sublattices with nonmagnetic impurities can relieve the geometrical frustration globally and, above a certain dilution threshold, the system starts showing long-range magnetic ordering characterized by the second-order phase transitions [1]. Within the effective-field theory with correlations we demonstrate that the presence of a finite external field in the selectively diluted system can dramatically change its critical behavior. In particular, above a certain field value the system starts displaying tricritical behavior, i.e., as the dilution increases the phase transition becomes initially of first order and only upon further dilution changes to the second order. We confirm the existence of the first-order transition in the region of moderate dilution and relatively high fields by Monte Carlo simulations. An Ising antiferromagnet on a honeycomb lattice is treated as a special case when one sublattice is completely removed by dilution. This work was supported by the Grant VEGA No. 1/0234/12.

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# QE03

# Possible magnetic transition observed in S=1/2 kagome antiferromagnet volborthite by high field ESR

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Due to the recent discovery of model substances, such as  $ZnCu_3(OH)_6C_{12}$  (herbertsmithite),  $Cu_3V_2O_3(OH)_{2'}2H_2O(volborthite)$  and  $BaCu_3V_2O_8(OH)_2(vesignieite)$ , S=1/2 kagome antiferromanet has attracted much interest experimentally, and we suggested a gapless spin liquid state in vesignieite [W. Zhang et al, J, Phys. Soc. Jpn. 79 (2010) 023708.]. In this paper we will report on high field ESR results of volborthite powder sample. For the spin dynamics volborthite shows the g-shift below around 20 K similar to vesignieite [H. Ohta et al., Phys. Soc. Jpn. 79 (and mode B with no gap. Existence of mode B may suggest the spin liquid state with no spin gap at 1.8 K. Moreover, two modes cross at around 4.5 T, which corresponds to magnetization anomaly at 4.3 T. Detailed frequency dependence measurement up to 30 T at 1.8 K shows also new anomaly at 14 T. Possible magnetic transitions at 4.3 and 14 T will be discussed in connection with recent NMR result on volborthite will be also presented.

# **OE04**

#### Field-induced staggered moment stabilization in frustrated quantum magnets

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For low-dimensional frustrated magnets, the dependence of the staggered moment on a magnetic field is nonmonotonic: For small and intermediate fields, quantum fluctuations are gradually suppressed, leading to an increase of the staggered moment as function of the field strength. For large applied magnetic fields, the classically expected field dependence is recovered, namely a monotonous decrease with increasing field strength. The staggered moment is eventually suppressed when reaching the fully polarized state at saturation. The quantitative analysis of this behavior is an excellent tool to determine the frustration parameter of magnetic compounds. We have developed a general finite-size scaling scheme for numerical exact-diagonalization data of low-dimensional frustrated magnets, which we apply to the recently measured field dependence of the magnetic neutron scattering intensity of Cu(pz)2(ClO4)2 in the framework of the S=1/2 two-dimensional J 1-J2 Heisenberg model. We also apply linear spin-wave theory to complement our numerical findings. Our results show that Cu(pz)<sub>2</sub>(ClO<sub>4</sub>)<sub>2</sub> is a quasi-2D antiferromagnet with intermediate frustration J\_2/ J\_1=0.2. A self-consistent RPA theory for the magnetic ordering temperature shows that with this ratio the observed reentrant behavior of the latter as function of the applied magnetic field can be understood as a consequence of the reduced quantum fluctuations as well

# **OE05**

### Structural and magnetic properties of single crystals of volborthite comprising a distorted spin-1/2 kagome lattice

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In geometrically frustrated magnets, exotic ground states such as a spin liquid are expected to realize as a result of competing magnetic interaction. Volborthite Cu<sub>3</sub>V<sub>2</sub>O<sub>7</sub>(OH)<sub>2</sub>·2H<sub>2</sub>O has drawn attention as a model compound for the distorted spin-1/2 kagome antiferromagnet. Previous studies using powder samples showed a peculiar magnetic transition at 1 K. Recently, H. Yoshida et al. synthesized a small single crystal, and found a structural phase transition at 310 K and two magnetic transitions at around 1 K, which have not been observed in powder samples. To clarify the physical properties of volborthite at low temperatures, we tried to synthesize a large and highquality single crystal and obtained large single crystals of 2.5×1.5×0.1 mm in size as a result of improving synthesis conditions. By means of single crystal XRD, we found a new crystal polymorph at room temperature which crystallizes in space group C2/ c, not in C2/m as reported before. There is a phase transition upon cooling at around 290 K to an I2/a phase, similar as in the C2/m phase. In this presentation, we show the differences of structure and magnetic properties between various polymorphs.

### **OE06**

#### High-field study of multiferroic Ni<sub>3</sub>V<sub>2</sub>O<sub>8</sub>

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# **OE07**

### Gapless spin excitation of the triangular-lattice antiferromagnet Hiroki Nakano1\* and Toru Sakai

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Two-dimensional frustrated magets including the triangular-lattice antiferromagnet (TLA) have attracted much attention. However, a useful method among unbaised computational methods is limited to being the numerical-diagonalization method that can treat only small systems. In spite of such situation, great efforts have been done for studies of the TLA: it is widely believed that the ground state of TLA has a long-range order of 120-degree structure. The existence of this order indicates that spin excitation is gapless; however, the system size (N) dependence of the singlet-triplet energy difference of finite-size clusters shows a peculiarity that a nonzero spin gap seems to survive in the extrapolation with respect to N[1]. Under circumstances, we calculate the singlet-triplet energy difference of possible all the finite-size clusters up to 39 sites by our MPI-parallelized code of the numerical-diagonalization method. The analysis used in the study of the spin-gap issue of the kagome-lattice antiferromagnet[2] is applied to the TLA. The present examination successfully confirms a gapless spin excitation, which is consistent with the existence of the long-range order.

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#### **OE08**

#### Magnetic properties of the spin-1/2 kagome antiferromagnets: vesignieite BaCu<sub>3</sub>V<sub>2</sub>O<sub>8</sub>(OH), and CdCu<sub>3</sub>(OH)<sub>6</sub>Br<sub>2</sub>

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An ideal model system for the spin-1/2 kagome antiferromagnet has not yet been obtained and the search for that is still going on. Here we report the magnetic properties of two Cu minerals, vesignieite BaCu<sub>3</sub>V<sub>2</sub>O<sub>8</sub>(OH)<sub>2</sub> and CdCu<sub>3</sub>(OH)<sub>6</sub>Br<sub>2</sub>, both of which comprise spin-1/2 kagome lattices. Vesignieite has a nearly isotropic kagome lattice and was found to show neither Neel order nor a spin-glass transition above 2 K [1]. However, the intrinsic ground state has remained unsolved, because previous samples contained many impurity spins caused by the low crystallinity. Recently, we succeeded in improving the sample quality by optimizing preparation parameters under hydrothermal conditions. The magnetic susceptibility of the new sample increases steeply below 15 K, followed by a sharp peak at 9 K that is indicative of long-range antiferromagnetic order. CdCu3(OH)6Br2 has an isotropic kagome lattice made of Cu2 ions. The compound is an antiferromagnet with a Weiss temperature of -48 K and shows long-range antiferromagnetic order at 7 K. Thus, both the compounds show long-range order at rather high temperatures, which may be due to significantly large interlayer magnetic interactions

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### **OE09**

### Spin dynamics of triangular spin tubes

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Triangular spin tubes belong to a new category of one-dimensional Heisenberg antiferromagnets along the tubes coexisting with geometrically frustrated spin systems in the triangular plane. Recently, we discovered that chromium fluoride CsCrF4 forms ideal equilateral triangular spin tubes with S=3/2[1]. All Cr ions are at equivalent sites, and each angle of Cr-Cr-Cr in the triangular plane is 60 deg.[2] Because superexchange interactions through the three Cr-F-Cr paths in each equilateral triangle may be kept in equilibrium at 0 K, we propose resonating spin-singlet pairs in each equilateral triangle. From DC magnetic susceptibility and adiabatic heat capacity data above 1.5 K, we concluded that there is no magnetic phase transition. Therefore, CsCrF4 consists of a gapless spin-liquid ground state encompassing resonating spin-singlet pairs not only in each equilateral triangle but also along the tubes[2]. To clarify spin dynamics of the spin-liquid ground state, we performed linear and nonlinear AC magnetic susceptibility measurements on CsCrF4. As a result, temperature dependence of a linear susceptibility exhibits a positive peak at around 4 K, while a nonlinear susceptibility exhibits a broad and negative peak at around 3.5 K. We found that a spin-glass like state was realized helow 4 K

[1] H. Manaka et al., J. Phys. Soc. Jpn. 78, 093701 (2009). [2] H. Manaka et al., J. Phys. Soc. Jpn. 80. 084714 (2011)

# **OE10**

#### Origin of field induced magnetic ordering in frustrated honeycomb lattice antiferromagnet

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The ground state of honeycomb lattice antiferromagnet (HLAF) have been considered theoretically to be the Neel state similar to the square lattice antiferromagnet because both antiferror are bipartite lattices. However, the honevcomb lattice has smaller coordination number 3 than the square lattice of 4 and it is close to the one dimensional chain. Therefore, the stronger spin fluctuation is expected in the HLAF. Actually, anomalous spin dynamics are reported in S=1/2 HLAF InCu<sub>3/2</sub>V<sub>1/3</sub>O<sub>3</sub> [1]. On the other hand, the spin liquid state is expected in S=3/2 HLAF Bi3Mn4O12(NO3) due to no long range order down to 0.4K despite of the large Weiss temper Although no spin gap has been observed in the magnetic susceptibility of Bi<sub>2</sub>Mn<sub>2</sub>O<sub>2</sub>(NO<sub>2</sub>). it shows the field-induced magnetic ordering (FIMO) above around 6 T. Neutron diffraction measurements confirmed FIMO state and suggested that Dzyaloshinski-Moriya (DM) interaction plays an important role in the FIMO state. In order to estimate the DM interaction high-field ESR arements have been performed. The D-term of DM interaction is estimated to be 1.3K from the analysis of observed antiferromagnetic resonance mode. Origin of FIMO will be discussed in connection with the DM interaction and the other magnetic anisotropy. [1] S. Okubo et al. J. Phys. Soc. Jpn. 80 (2011) 023703

# **OE11**

Semi-classical spin-liquid state as a low-energy excited state in frustrated quantum spin systems on triangle-based lattice system Makoto Isoda1, Hiroki Nakano2 and T^oru Sakai3 Kagawa University, Japan University of Hyogo, Japan <sup>3</sup> JAEA SPring-8, Japar

The magnetic and the thermal properties of the spin-1/2 Heisenberg antiferromagnets on the trianglebased lattice are summarized, focusing on the realization of the semi-classical spin-liquid state of the resonating doublet trimers formed on triangles, through the calculations for various systems such as the triangular, the kagome, the triangulated-kagome1), the Sierpinski-gasket like2) and the spatially anisotropic triangular lattices, which interpolates the triangular and the kagome lattices. The calculations have been performed by the numerical diagonalization method. On these systems, the peculiar behaviors have been found in the Heisenberg3,4) and the Ising5) models for spin-1/2 on those triangle-based lattices with nearest neighbor (nn) antiferromagnetic interaction. The peculiarities are the bending of the magnetic susceptibility as a function of temperature, the more than two peaks structure on the specific heat and the 1/3 magnetization plateau under magnetic field. Each of these peculiarities appears at the characteristic temperature or magnetic field corresponding to the magnitude of nn interaction, suggesting the predominance of the doublet trimer low energy excited state. The peculiarities reveal as an intrinsic character of the systems with the nn antiferromagnetic interaction depending on whether the lattice is formed by the equilateral triangle as a unit cell or not.

1) M.Gonzalez et al., Mol. Cry. Liq. Cry. 233 (1993) 317. 2) T. Stosic et al., Phys. Rev. E 49 (1994) 1009. 3) M.Isoda, H.Nakano and T.Sakai, Mod. Phys. Lett. B 25 (2011) 909. 4) M.Isoda, H.Nakano and T.Sakai, J. Phys. Soc. Jpn. 80 (2011) 084704, 5) M.Isoda, J. Phys.: Condens. Matter 20 (2008) 315202.

### **OE12**

#### High pressure and low-temperature 31P NMR study of the twodimensional frustrated square lattice compound BaCdVO(PO<sub>4</sub>)<sub>2</sub>

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We report first NMR study under pressures on BaCdVO(PO<sub>4</sub>)<sub>2</sub>, a S = 1/2 frustrated square-lattice (FSL) compound with a nearest neighbor exchange coupling J1  $\sim$  - 3.36 K and a next-nearest neighbor exchange coupling J2  $\sim 3.53$  K bearing J2/ J1  $\sim 1.05.$ Based on the J2/J1 ratio the system is known to be located close to the disordered ground state (known as 'nematic state') regime of the phase diagram. We have carried out 31P-NMR measurements under high pressure (up to  $\sim 2$  GPa) and at low temperatures (down to ~0.1 K) using a dilution refrigerator to investigate the pressure effects on the magnetic properties of the system from microscopic point of view. Under ambient pressure, we observed a sharp peak in 31P spin lattice relaxation rate (1/T1) at TN  $\sim 1.05$  K at H = 2.67 T, which corresponds to the antiferromagnetic ordering temperature. With increasing pressure, the peak position of 1/T1 shifts to lower temperature. This indicates that TN decreases with the application of pressure. We will discuss pressure effects on magnetic state based on the temperature dependence of the NMR spectra and of the 1/T1 under different magnetic fields and pressures.

# **OE13**

### Kasteleyn transitions in the spin ice Dy<sub>2</sub>Ti<sub>2</sub>O<sub>7</sub>

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The spin ice model in the pyrochlore lattice, e.g. Dy2Ti2O7, consisting of corner sharing tetrahedra has attracted much attention because of its intriguing frustration mechanism. The ground states of this system are macroscopically degenerate with the finite zero-temperature entropy, and the spins freeze in this manifold at low temperatures. The macroscopic degeneracy of the ground states is partly lifted under magnetic fields, where degrees of the degeneracy depend on direction and strength of the field. Interesting theoretical predictions, which have not been seriously investigated to date by experiments, are Kastelevn transitions [1] in three and two space dimensions, which may occur under magnetic fields along a [100] direction [2] and fields tilted slightly from a [111] direction [3]. These transitions are classified into topological phase transitions which are characterized by line (string) defects. We have investigated these possibilities in the spin ice compound Dy2Ti2O7 using magnetization and specific heat measurements

[1] P. W. Kasteleyn, J. Math. Phys. 4, 287 (1963). [2] L. D. C. Jaubert et al. Phys. Rev. Lett. 100, 067207 (2008), [3] R. Moessner, S. L. Sondhi, Phys. Rev. B 68, 064411 (2003).

#### **OE14**

Novel frustrated quantum antiferromagnets in the solid-solution Cs<sub>2</sub>CuCl<sub>4x</sub>Br<sub>x</sub> through site-selective halide substitution

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We report on the magnetic properties of tetragonal and orthorhombic single crystals of the solid solution  $Cs_2CuCl_{4x}Br_x$  ( $0 \le x \le 4$ ). Two phase regions, namely an A-type orthorhombic and a B-type tetragonal structure, were detected in this system depending on the growing conditions. The two known orthorhombic end-member compounds Cs2CuCl4 and Cs2CuBr4 are classified as quasi-two-dimensional quantum antiferromagnets with different degrees of magnetic frustration. By measurements of the magnetic susceptibility  $\chi(T)$  on orthorhombic single crystals with different Br concentrations, we found that the in-plane and out-of-plane magnetic correlations, probed by the position and height of a maximum in  $\chi(T)$ , respectively, do not show a smooth variation with x. Instead three distinct concentration regimes can be identified, which are separated by critical concentrations xc1 = 1 and xc2 = 2. This unusual magnetic behaviour can be explained by considering the structural peculiarities of the materials, which support a site-selective replacement of Cl- by Br- ions [1]. Consequently, the critical concentrations xc1 (xc2) mark particularly interesting systems, where one (two) halide-sublattice positions are fully occupied. The magnetic properties of the tetragonal phase of Cs2CuCl4xBrx differ strongly from those of the orthorhombic variants, reflecting changes in the Cu-coordination.

[1] P. T. Cong et al., PRB 83, 064425, (2011)

# **OE15**

Strong geometrical frustration in Fe oxychalcogenide Sungdae Ji1, K. Horigane2 and K. Yamada3

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After Fe-based superconductors have emerged, it has generated a lot of controversies regarding to low ordered moments in parent compounds which have been explained by itinerant magnetism with Fermi surface nesting vs. magnetic frustration and/ or fluctuation effects in a large-local-moment system [1]. Very recently, Feoxychalcogenide has drawn attention, because it has been identified as a layered insulating antiferromagnet with the geometrically frustrated, so-called, "double diagonal stripe" spin structure which is realized only in Iron telluride among the family of Fe-based superconductors [2,3]. Here, we report elastic and inelastic neutron scattering study on La2O2Fe2OSe2. Our elastic neutron measurement at 5 K shows that magnetic Bragg peaks are diffusive and a obtained magnetic moment is 2.2 µ B, which is smaller than a localized spin, S=5/2. Inelastic neutron scattering measurement shows that a strong magnetic fluctuation survives even at 300 K far above magnetic transition temperature, T\_N = 90 K. Moreover, observed dispersive excitation reveals that the 3rd nearest neighbor interaction is the key to explain a magnetic ground state on the geometrically frustrated lattice.

[1]. David C. Johnston, Adv. in Phys. 59, 803 (2010). [2]. Jian-Xin Zhu et al., Phys. Rev. Lett. 104, 216405 (2010). [3]. David G. Free and John S. O. Evans, Phys. Rev. B 81, 214433 (201

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CROSS, Japan

# **QE16**

Magnetic properties of the frustrated magnet  $Cu_5(PO_4)_3(OH)_4$  on a peculiar spin network composed of pentagons and triangles

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Orthorombic Cu<sub>3</sub>(PO<sub>4</sub>)<sub>2</sub>(OH)<sub>4</sub> (pseudomalachite) is a layered compound. Each Cu-O layer in this compound is well separated by PO<sub>4</sub> tetrahedrons. Spin network in beplane of pseudomlachite is composed of triangles and pentagons, which will induce the spin frustration effect. This network can be a new type of geometrically frustrated spin lattice. Magnetic susceptibility, specific heat, high field magnetization are measured to study magnetic properties of Cu<sub>3</sub>(PO<sub>4</sub>)<sub>2</sub>(OH)<sub>4</sub> using natural mineral samples. Magnetic ordering is observed at around 4.2 K. Spin entropy change at the transition temperature is about one fifth times smaller than that expected for a usual long range order. 1/5 magnetization plateau is observed in the high field magnetization curve. The plateau can be explained based on the simple model in which the spin network is assumed to be composed of chain, dimer and monomer units.

**OE17** 

uly 10 (Tue)

# Single crystal growth and magnetic properties of novel kagome compound KMn3Ge2O9

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We have succeeded in synthesizing single crystals of a novel manganese oxide KMn<sub>3</sub>Ge<sub>2</sub>O<sub>9</sub> by a hydrothermal method. This compound has a hexagonal unit cell with lattice constants a = 11.77 Å and c = 13.77 Å. In this structure, edge-shared MnO6 octahedra form layers separated by GeO4 tetrahedra. In each layer, Mn ions with S = 2 spin form a rippled kagome lattice. Therefore, we can expect a frustrated magnetism. We measured temperature dependence of the magnetic susceptibility on a single crystal. The result suggests the existence of a strong magnetic frustration. The Curie constant and Weiss temperature are estimated at 3.4 emu K/mol Mn and - 84 K, respectively. The negative Weiss temperature indicates that the nearest-neighbor exchange interaction among Mn ions is antiferromagnetic. At low temperatures, the susceptibility suggests the existence of multiple magnetic phases. For example, the susceptibility shows an antiferromagnetic like transition at 40 K and a spin-glass like transition at 20 K under 0.5 T of magnetic field perpendicular to the kagome lattice. The successive transitions are also observed in other conditions of magnetic field. We consider that these behaviors are driven from the strong magnetic frustration of the kagome lattice.

### **QE18**

# Specific heat study of geometrically frustrated magnet botallackite Cu<sub>2</sub>(OH)<sub>2</sub>Cl

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Specific heat measurements have been performed in the geometrically frustrated magnets botallackite,  $Cu_2(OH)_3Cl$  with the triangular structure of the  $Cu^{2+}$  ions. This is a different polymorphic form of clinoatacamite with the corner-sharing tetrahedral structure, which exhibits the coexistence of spin fluctuation and long range ordering below 6.2 K[1]. The specific heat of botallackite shows a sharp peak at TN=7.0K and approximately exhibits T3-dependence with lowering temperature in zero magnetic field. When a magnetic field is applied, TN shifts to lower temperatures and the peaks disappears at H=8 T. Moreover, no residual entropy is observed at lower temperatures.

[1] X.G. Zheng, et. al., Phys. Rev. Lett. 95 057201 (2005).

# QE19

**Order and excitations in the frustrated quantum spin ladder BiCu<sub>2</sub>PO<sub>6</sub>** P. Merchant<sup>1</sup>, S. Wang<sup>2</sup>, O. Zaharko<sup>3</sup>, Ch. Niedermayer<sup>3</sup>, L. P. Regnault<sup>4</sup>, M. Boehm<sup>4</sup>, M. Kenzelmann<sup>2</sup> and Ch. Ruegg<sup>3</sup>

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Quantum spin ladders are exceptional model systems showing non-trivial phenomena ranging from complex spin correlations in the ground state to rich phase diagrams with several quantum critical points. A fully quantitative understanding of their magnetic and thermodynamic properties may be achieved e.g. by powerful computational methods (DMRG, ED, bond operators) [1, 2] and comparison with high-precision experimental data, which may only be possible in such model systems [3]. BiCu<sub>2</sub>PO<sub>6</sub> is a new prototypical model material in which the spin ladders are furstrated by next-nearest neighbour exchange along the ladder legs, furthermore it can be doped with non-magnetic  $Zn^{2+}$  impurities. The results of comprehensive neutron scattering studies on single crystals are presented of the elementary excitations and spin Hamiltonian of this frustratel ladder compound, as well as of the Zn-induced magnetic order and phase diagrams. The physics in this system is dominated by incommensurate correlations on all energy scales.

[1] P. Bouillot et al., Phys. Rev. B 83, 054407 (2011). [2] B. Normand and Ch. Ruegg, Phys. Rev. B 83, 054415 (2011). [3] e.g. see B. Thielemann et al., Phys. Rev. Lett. 102, 107204 (2009).

#### **OE20**

# Melting of the spin ice state in $Dy_2(Ti_{1:x}Zr_x)_2O_7$ without dilution of rareearth ion

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The pyrochlore  $Dy_2Ti_2O_7$  is a spin ice material, in which the spins reside on the lattice of corner-sharing tetrahedra with single-ion anisotropy along each local <111> axis. The effective ferromagnetic interaction between spins leads to a highly degenerate ground state. Earlier researches [1][2] reported that the magnetic relaxation shows three successive processes; thermally activated Arrhenius process above 13 K, quantum tunneling process between 4 K ~ 13 K, and process with strong spin correlation below 4 K. The relaxation in the medium temperature regime is the Davidson-Cole type [3] characterized by a wide distribution of relaxation time  $\tau$  with a cut-off  $\tau$ . The dilution of Dy causes disappearance of the ac susceptibility 15 K peak [1] due to a speeding up of relaxation [2]. In this talk, we examined the magnetic relaxation of  $Dy_2(Ti_{1,2}T_2)_{2}O_7$  through ac magnetic susceptibility measurements in order to study B-site disorder effect. We found, for x above about 0.01, disappearance of the 15 K peak, small frequency dependence, and substantial change of the magnetic relaxation from the Davidson-Cole type into the Cole-Cole one having no cut-off  $\tau$ , indicating a major effect of B-site disorder on the spin ice state toward melting without dilution of Dy ion.

[1] J. Snyder et al., Phys. Rev. B 70 (2004) 184431. [2] G. Ehlers et al., J. Phys. Condens. Matter 15 (2003) L9. [3] K. Matsuhira et al., J. Phys. Condens. Matter 13 (2001) L737. [4] L.D.Jaubert and P.C.Holdsworth, J. Phys. Condens. Matter 23 (2011) 164222.

#### QE21

#### Unusual magnetic ordering of kagome lattice magnet [Cu<sub>3</sub>(CO<sub>3</sub>)<sub>2</sub>(bpc)<sub>3</sub>]•<sub>2</sub>ClO<sub>4</sub>

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Cu<sup>2+</sup> ions in a hexagonal compound [Cu<sub>3</sub>(CO<sub>3</sub>)<sub>2</sub>(bpe)<sub>3</sub>]•2ClO<sub>4</sub> (space group; P-6) form slightly distorted kagome lattice in its c-plane. Former study1) reported that the magnetic susceptibility increased rapidly below about 7 K indicating an occurrence of the ferromagnetic transition. Surprisingly, however, no anomaly in the specific heat was observed at the transition temperature in this compound. In order to scrutinize the magnetic properties, we measure the specific heat of this compound using single crystals under an applied magnetic field up to 7 T. 1H-NMR measurement is also carried out. NMR spectra-broadening and a divergent behavior of the spin-lattice relaxation rate are observed at around 7 K. Slight anomaly in the specific heat is observed and the anomaly temperature is found to depend on the magnetic field. These results suggest that entropy change associated with the magnetic ordering of this compound is vanishingly small.

1) P. Kanoo et al. Dalton Trans. 2009, 5062.

# QE22

#### Field-induced staggered moments in the spin-gapped antiferromagnet on a deformed kagome lattice, Rb<sub>2</sub>Cu<sub>3</sub>SnF<sub>1</sub>,

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Rb\_Cu<sub>5</sub>SnF<sub>12</sub> is a kagome lattice antiferromagnet in which Cu<sup>2+</sup> ions with spin 1/2 interact with four exchange interactions, J1-200K>J2>J3>J4-100K. The ground state of this compound is a valence bond solid owing to the quantum and frustration effects, so that no magnetic ordering occurs [1,2]. We have performed 19F-NMR experiments on a single crystal of Rb<sub>2</sub>Cu<sub>5</sub>SnF<sub>12</sub> in order to clarify the microscopic property of the ground state. One sharp 19F-NMR spectrum at high temperature splits to several peaks at low temperatures due to different magnetic sites of 19F. The shifts of these peaks increase steeply below around 70K, whereas the static susceptibility decreases exponentially below this temperature owing to the singlet ground state with the spin-gap. The shifts are proportional to the applied field. These results strongly suggest that the local fields are produced by the staggered moments induced by the external field through Dzyaloshinsky-Moriya interaction. The staggered moments become dominant below 70K, while the parallel moments are remarkably small because of the large exchange interaction. This characteristic temperature differs from the spin-gap of 20K obtained from the susceptibility, because the large density of state in the dispersion curve around 70K contributes to the generation of the stageered moments [3].

[1] K. Morita, M. Yano, T. Ono, H. Tanaka, K. Fujii, H. Uelasa, Y. Narumi, and K. Kindo, J. Phys. Soc. Jpn. 77, 043707 (2008). [2] K. Matan, T. Ono, Y. Fukumoto, T.J. Sato, J. Yamcura, M. Yano, K. Morita, and H. Tanaka, Nature Phys. 6, 865-869 (2010). [3] H. Tashiro, M. Nishiyama, A. Oyamada, T. Itou, S. Maegawa, M. Yano, T. Ono, and H. Tanaka, J. Phys.: Conf. Ser. 320, 012052 (2011).

### QE23

# Magnetic order in finite size domains of the honeycomb lattice compound $InCu_{23}V_{13}O_3$

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The results of high field electron spin resonance, nuclear magnetic resonance, nuclear quadrupole resonance and magnetization studies addressing the ground state of the quasi two-dimensional spin-1/2 honeycomb lattice compound InCu<sub>2/3</sub>V<sub>1/3</sub>O<sub>3</sub> are reported. Uncorrelated finite size structural domains occurring in the honeycomb planes are expected to inhibit long range magnetic order. Surprisingly, magnetic resonance data show the development of two collinear antiferromagnetic (AFM) sublattices below 35K and the presence of the staggered internal field. Magnetization data evidence a spin reorientation transition at 5.7T. Quantum Monte-Carlo calculations show that switching on the coupling between the honeycomb spin planes in a finite size cluster yields a Neel-like AFM spin structure with a substantial staggered magnetization at finite temperatures. ESR, NMR and NQR data allow to retrace the development of staggered magnetization during the crossover to ordered ground state.

A. Yehia, et al., Phys. Rev. B 81, 060414 (2010)

#### QE24

# $Z_2$ vortices in frustrated background of cuprates

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La<sub>2x</sub>Sr<sub>x</sub>CuO<sub>4</sub> (LSCO) in the underdoped regime (x=0.02 - 0.05) is a non-collinear quantum antiferromagnet. The relevant order parameter is a triad of orthonormal vectors. The topology of the order parameter space allows the existence of Z\_2 vortices. Randomly distributed Sr dopants introduce holes into the Cu-O planes which are the source of fluctuating dipolar fields resulting in frustrated spin background. We study Z\_2 vortices in the system in the framework of non-linear sigma-model coupled to SO(3) gauge fields which represent fluctuating frustrated background. In this model we describe explicitly not only spin degrees of freedom, but also the distribution of frustrations as opposed to pure spin systems. We have found explicitly a finite energy Z\_2 vortex solution. Frustrations are concentrated in the soft core of Z\_2 vortices and exponentially decrease at large distances. The chirality of spin system changes its sign in the vortex core as expected. Our vortex seems to be more relevant to the study of frustrated Heisenberg spin systems and the charge transport by vortices in cuprates than Z\_2 vortices in pure spin degrees of freedom.

# QE25

### First-principles study of S=1/2 Kagome antiferromagnet

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3-D systems with interacting spins usually develop static long-range order when they are cooled. However, quantum fluctuations are enhanced in low dimensionality system with low spin value and lead to suppress the long-range ordering, and can show some extremely subtle, complex, and sometimes even useful magnetic behavior in solids. Here, using first-principles calculations, we investigated the atomic structural and magnetic properties of S=1/2 Kagome antiferromagnet A<sub>x</sub>Cu<sub>+x</sub>(OH)<sub>6</sub>Cl<sub>2</sub> (A=Zn, alkali metal ion). Cu ions are indeed found to locate at the tetragonally elongated intralayer site and Zn ions favorably rest on the higher symmetry interlayer site, in support of experimental contention. The intra- and inter-layer exchange interactions are studied in details. The effect of Dzyaloshinskii-Moriya interaction on the magnetic structure is also discussed.

[1] M. P. Shores, E. A. Nytko, B. M. Bartlett, and D. G. Nocera, J. Am. Chem. Soc 127, 13462, (2005) [2] S. Chu, T.M McQueen, R. Chisnell, D.E. Freedman, P. Muller, Y.S. Lee, and D.G. Nocera, J. Am. Chem. Soc. (2010).

# **QE26**

### Exotic phases in the frustrated hexagonal lattice

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We study the phase diagram of the Heisenberg model on the Honeycomb lattice with antiferromagnetic interactions up to third neighbors along the line J2=J3, close to the point where it has macroscopic degeneracy at the classical level. Using the Schwinger boson technique followed by a mean field decoupling and exact diagonalization for small systems we find an intermediate phase with a spin gap and short range Neel correlations even in the strong quantum limit 1/2. We have also studied the magnetization curve for J3=0, close to J2/J1 = 1/2 where the classical ground state is also highly degenerate. The properties of the magnetization curve and phase transitions are studied by means of Montecarlo, which shows many interesting transitions.

# **QE27**

#### Light scattering in spin liquid systems

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When quantum spin systems are restricted in dimensionality and coordination they realize spin liquid states with enhanced quantum fluctuations and exotic correlation functions. We compare the experimentally determined excitation spectra of different spin liquid candidates such as the s=1/2 Heisenberg antiferromagnet on a kagome lattice and weakly coupled spin chain system in Herbertsmithite and (NO)Cu(NO<sub>3</sub>)<sub>3</sub> using Raman scattering. This technique is sensitive to fractional spinon excitations. The effect of the crystal's lattice structure and defects on the spin dynamics is investigated. The data is also compared with recent theoretical modelling. Work supported by DFG, B-IGSM and NTH School for Contacts in Nanosystems.

D. Wulferding, et al., Phys. Rev. B 82, 144412 (2010) V. Gnezdilov, et al., arXiv:1203.2818 (2012) D. Wulferding, et al., arXiv:1111.2167 (2012)

# **OF01**

#### Molecular nanomagnets as quantum simulators

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Quantum simulators (OSs) are controllable quantum systems that can be used to simulate other quantum systems. OSs can tackle problems intractable on classical computers. Here we focus on the dynamics of gubits encoded in chains of molecular nanomagnets. For istance, we have shown that Cr<sub>7</sub>Ni rings are good candidates qubits and can be linked to each other either directly or through magnetic complexes [1,2]. We theoretically show that the dynamics of such chains can be controlled by means of uniform magnetic pulsed fields and used to mimic the coherent time evolution of other quantum systems (e.g., spin-one chains) [3]. We propose two significant proof-ofprinciple experiments [3].

[1] G.A. Timco et al., Nature Nanotechnology 4, 173 (2009) [2] A. Candini et al., Phys. Rev. Lett. 104, 037203 (2010) [3] P. Santini, S. Carretta, F. Troiani, G. Amoretti, Phys. Rev. Lett. 107, 230502 (2011)

# **OF02**

#### ESR signature of the next-nearest-neighbor interactions in the S = 1/2chain compound (6MAP)CuCl<sub>3</sub>.

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Electron spin resonance (ESR) studies of the S = 1/2 chain material (6MAP)CuCl<sub>3</sub>  $[6MAP = C_6H_9N_2]$  are presented. Two modes with asymmetric with respect to a simple gapless Zeeman splitting resonance positions have been observed in the low-temperature ESR spectrum. This shows that the simple S = 1/2 Heisenberg antiferromagnet chain model, employed so far, is not sufficient for a complete description. The frequency-field diagram of magnetic excitations is interpreted in the frame of the recently developed theory for S = 1/2 chains with nearest- and nextnearest-neighbor interactions [A.A. Zvyagin, Phys. Rev. B 79, 064422 (2009)]. A good qualitative agreement with experiment has been achieved. The work was supported in part by DFG and EuroMagNET II (EU Contract No. 228043).

# **OF03**

#### Electron spin resonance in the spin-ladder compound BPCB

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Magnetic excitations in the spin-ladder material (C5<sub>2</sub>H<sub>12</sub>N<sub>2</sub>)CuBr<sub>4</sub> (known as BPCB) are probed by means of electron spin resonance (ESR) spectroscopy in magnetic fields up to 16 T. The gap between the ground state and the first excited state, 16.5 K (or 1.42 meV), is detected directly. In addition, our experiments provide clear evidence for a pronounced anisotropy (~ 5% of the dominant exchange interaction), responsible for a shift of the observed ESR line and a specific angular dependence of resonance absorptions. The results are explained in frame of the recently developed theory for ESR in spin-ladders (Furuya et al., arXiv: 1107.5965). The work was done in collaboration with B. Thielemann, K. W. Kramer, Ch. Ruegg, O. Piovesana, M. Klanjsek, M. Horvatic, and C. Berthier. The work was supported in part by DFG and EuroMagNET II (EU Contract No. 228043).

# **OF04**

#### Non-magnetic impurity effect of S=1/2 spin ladder system (pipdH)<sub>2</sub>Cu<sub>1-</sub> <sub>x</sub>Zn<sub>x</sub>Br<sub>4</sub>

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Low dimensional spin system shows fruitful magnetic properties such as the Bose-Einstein condensation of magnon or the spin liquid state. Non-magnetic impurity effect in a quantum spin chain, which arises as an appearance of magnetic moment with spin correlation at around the impurity, is one of typical quantum effects of correlated spins. (pipdH)2CuBr4 is known as a model substance of weakly coupled spin 1/2 Heisenberg antiferromagnetic ladder system. Recently, we succeeded in synthesizing Zn-doped (pipdH), Cu, Zn, Br, single crystal. Samples were confirmed to be in a single phase and have the same crystal structure as the pure system (pipdH), CuBr<sub>4</sub> by X-ray diffraction measurements. To check magnetic properties, the magnetic susceptibility and the magnetization measurements are performed by the SQUID magnetometer. The magnetic susceptibility shows a maximum related to the spin ladder and the Curie increase at low temperature related to the doping of Zn. The Curie term increases as the Zn concentration increase. Observed non-magnetic impurity effect of S=1/2 spin ladder system will be discussed in connection with AKLT model

#### **OF05**

#### High-field multi-frequency ESR in the S=2 Heisenberg antiferromagnetic chain compound MnCl<sub>3</sub>(bpy)

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# **OF06**

#### High field magnetization of bimetallic chain with alternating ising and Heisenberg spins

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3d-4f bimetallic chain [Dy(NO3)(DMSO)2Cu(opba)(DMSO)2] built from alternating dysprosium(3+) and copper(2+) ions has been synthesized and magnetically studied. Magnetic susceptibility as well as high-field magnetization of the polycrystalline sample was measured under the temperature down to 1.3 K and the magnetic field up to 52 T. After performing correction for the temperature independent paramagnetism, the magnetization curve recorded at 1.3 K shows a rapid increase with the magnetic field in the low-field range 0-3 T, followed up with a less steep increase in the field range 5-34 T before tending towards its saturation value at 34 T. The observed magnetization process can be interpreted in terms of the spin-1/2 chain model with alternating Ising and Heisenberg spins, which accounts for both parallel and perpendicular Zeeman's energy of Dy and Cu ions. The good fitting of the experimental data by using the above theoretical model suggest the effective antiferromagnetic exchange constant J/kB= -26 K, the highly anisotropic g-factor of Dy ions and almost isotropic g-factor of Cu ions. It was concluded that a new type of onedimensional Ising-Heisenberg spin chain was found.

# **OF07**

Thermal conductivity and magnetic susceptibility of the 4-leg spinladder system (La<sub>1,x</sub>Y<sub>x</sub>)<sub>2</sub>Cu<sub>2</sub>O<sub>5</sub> and the 5-leg spin-ladder system (La<sub>1</sub>  $_{x}Eu_{x})_{8}Cu_{7}O_{19}$ 

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We have measured the thermal conductivity and magnetic susceptibility of the 4-leg spin-ladder system  $(La_{1x}Y_x)_2Cu_2O_5$  (x = 0.01) and the 5-leg spin-ladder system  $(La_{1x}Eu_x)_8Cu_7O_{19}$  (x = 0, 0.01). It has been found that the temperature dependence of the thermal conductivity along the spin-ladders of (La099Y001),Cu2O5, La8Cu2O19 and (La099Eu001),Cu2O19 shows only one peak originating from the thermal conductivity due to phonons at ~15K. That is, the thermal conductivity due to spins, kspin, has not been observed in these compounds, though kspin was expected to be large owing to the exchange interaction along the spin-ladders as large as ~1000K. Since the antiferromagnetic transition temperature, TN, is as large as 137K and 103K in La2Cu2O5 [1] and La8Cu7O19 [2], respectively, it has been concluded that the comparatively large exchange interaction between spin-ladders disturbs the thermal conduction due to spins along the spin-ladders. Moreover, it has been found that the temperature dependence of the magnetic susceptibility below TN is clearly different between (La<sub>099</sub>Y<sub>001</sub>)<sub>2</sub>Cu<sub>2</sub>O<sub>5</sub> and (La099Eu001)&Cu7O19, which may be related to the difference of the ground state between the odd- and even-leg spin-ladder systems.

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#### **OF08**

#### Magnetic susceptibility of the quasi one-dimensional spin system Sr<sub>2</sub>V<sub>2</sub>O<sub>0</sub>

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According to the magnetic susceptibility of polycrystalline samples of the quasi onedimensional spin system Sr<sub>2</sub>V<sub>3</sub>O<sub>9</sub>, this compound undergoes a weak ferromagnetic transition due to the Dzyaloshinsky-Moriya (DM) interaction at TN ~ 5K [1]. We have grown single crystals of Sr<sub>2</sub>V<sub>3</sub>O<sub>9</sub> [2] and measured the magnetic susceptibility along the [101] direction, x[101], along the [10-1] direction, x[10-1], and along the b-axis,  $\chi$ b. It has been found that temperature dependences of  $\chi$ [101],  $\chi$ [10-1] and  $\chi$ b show Bonner-Fisher-type behavior and that the anisotropy is very small at high temperatures above TN, indicating that Sr<sub>2</sub>V<sub>3</sub>O<sub>9</sub> is regarded as an isotropic one-dimensional Heisenberg spin system. At low temperatures below TN, however, both  $\chi$ [10-1] and  $\chi$ b increase with decreasing temperature, while  $\chi$ [101] decreases. This anisotropic behavior is consistent with the canted spin state proposed from the ESR measurements [3].

[1] E. E. Kaul et al., Phys. Rev. B 67 (2003) 174417. [2] M. Uesaka et al., J. Phys.: Conf. Ser. 200 (2010) 022068, [3] V. A. Ivanshin et al., Phys. Rev. B 68 (2003) 064404.

#### **OF09**

#### Magnetic property of a single crystal of spin-1/2 triple-chain magnet Cu<sub>3</sub>(OH)<sub>4</sub>SO<sub>4</sub>

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A spin-1/2 triple-chain magnet Cu3(OH)4SO4 (antlerite) is a candidate 'idle-spin' system in which magnetic moments on central spin chain seem to disappear, while the rest exhibit a long-range order below TN=5.5 K [1]. Our previous proton NMR study with a powder sample almost supported this scenario, although it was difficult to determine the spin structure [2]. Recently, successive phase transitions and anisotropy of phase diagram were found from the results of magnetization measurements with high-quality single crystals of Cu<sub>3</sub>(OH)<sub>4</sub>SO<sub>4</sub> [3]. In order to study phase transitions in the present compound in detail, we have performed specific heat measurements under magnetic fields up to 7 T and in the temperature range down to 2 K and proton NMR experiments at several frequencies with single crystals. We found at least three successive phase transitions at zero field around TN. Complicated phase diagrams involving changes of spin structures were suggested.

[1] S. Vilminot et al., J. Solid State Chem., 170, (2003) 255. [2] Y. Fujii et al., J. Phys.: Conf. Ser. 145. (2009) 012061. [3] S. Hara et al., J. Phys. Soc. Jpn., 80, (2011) 043701

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# **OF10**

Thermal conductivity due to magnons in high-quality single crystals of the two-leg spin ladder system (Ca,Sr,La)14Cu24O41

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In order to enhance the thermal conductivity due to magnons, kmagnon, in the two-leg spin ladder system (Ca,Sr,La)14Cu24O41, we have grown single crystals using high-purity raw materials, improved the quality of the single crystals by annealing in O2 and measured the thermal conductivity along the spin ladders. In both Ca<sub>2</sub>La<sub>5</sub>Cu<sub>24</sub>O<sub>41</sub> and Sr<sub>2</sub>La<sub>5</sub>Cu<sub>24</sub>O<sub>41</sub>, the thermal conductivity shows a very large peak around 150K owing to the contribution of kmagnon along the spin ladders[1,2]. However, no enhancement of kmagnon has been observed in spite of the increase of the purity of the raw materials and the O2-annealing. In Sr14Cu24O41, on the other hand, a large enhancement of kmagnon has been observed through the O2-annealing. These results indicate that the mean free path of magnons, Imagnon, has already limited by the disorder of the atomic arrangement due to the substitution of Ca and La for Sr in Ca<sub>9</sub>La<sub>5</sub>Cu<sub>24</sub>O<sub>41</sub> and Sr<sub>9</sub>La<sub>5</sub>Cu<sub>24</sub>O<sub>41</sub>, while Imagnon is enhanced owing to the decrease of spin defects through the O2-annealing in nonsubstituted Sr14Cu24O41. Accoringly, it is necessary for the enhancement of kmagnon to remove the disorder of the atomic arrangement in the spin ladders in (Ca,Sr,La)14Cu24O41.

[1] K. Kudo, S. Ishikawa, T. Noji, T. adachi, Y. Koike, K. Maki, S. Tsuji, and K. Kumagai: J. Phys. Soc. Jpn. 70, 437 (2001). [2] C. Hess C. Baumann, U. Ammerahl, B. Buchner, F. Heidrich-Meisner, W. Brenig, and A. Revcolevschi: Phys. Rev. B 64, 184305 (2001).

#### **OF11**

Magnetic property of Ni<sup>2+</sup> antiferromagnetic perfect triangle cluster

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The magnetic properties of spin systems with frustration have attracted much attention for a long time. The fundamental structure causing the frustration is "triangle". The magnetic properties of the antiferromagnetic perfect triangular cluster (APTC) with half-integer spins have been already studied well. [1] By contrast, there are little experimental data on the APTC with integer spins. Such an APTC is expected to possess a non-magnetic, singlet ground state. Recently, Liu et al. have investigated magnetic properties of S=1 APTC,  $[Ni_3(\mu_{13}-N_3)_3(2,2'-bpy)_3](ClO_4)_3 \cdot 3H_2O$ . It is believed as one of model materials for the S=1 APTC. [2] We synthesized it and measured the magnetization process (M(H)) and the temperature dependence of the magnetic susceptibility  $(\chi(T))$  of the material. We calculated M(H) and  $\chi(T)$  in the simple ISOTROPIC S=1 APTC. However, the experimental results clearly cannot be accounted for by the calculated results. When the spin is greater than or equal to 1, generally a single ion anisotropy D can present. Thus we consider an anisotropic model Hamiltonian by including D. We will discuss the experimental data based on calculations using the new Hamiltonian with D.

[1] for example T. Yosida et al., J. Phys. Soc. Jpn. 57 (1988) 1428. [2]J.L. Liu et al., Aust. J. Chem. 63 (2010) 1111.

# **OF12**

#### Thermal conductivity of anisotropic spin ladder

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We have studied the thermal conductivity of anisotropic spin ladder model with antiferromagnetic coupling constants between spins on the both rung and ladder directions. Kubo formalism has been applied to study temperature dependency of thermal conductivity of this model hamilltonian. The bond operator formalism is used to transform the spin model to a hard core bosonic gas. We have used the green's function approach to obtain the temperature dependence of spin excitation spectrum. We have found the temperature dependence of the thermal conductivity for various exchange coupling constants and anisotropies in both coupling strengths. We have obtained the increase of coupling constant along ladder direction leads to decrease of thermal conductivity. Furthermore the effect of local anisotropy anisotropy on the thermal conductivity has more significant in comparison with other one.

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# QF13

# The magnetic properties of the newly synthesized trinuclear copper complex

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We present magnetic properties of a newly synthesized single molecule magnet (SMM) Cu<sub>3</sub>(µ<sub>3</sub>-OMe)(µ-OMe)(µ-O<sub>2</sub>CAr<sup>4F-Ph</sup>)<sub>2</sub>(O<sub>2</sub>CAr<sub>4</sub>F-Ph)<sub>2</sub>(HOMe)<sub>3</sub>. In this compound, three copper ions with s=1/2, connected by one µ<sub>3</sub>-bridged methoxide, form a spin triangle, whereby two Cu<sup>2+</sup> ions are 5 coordinated and the other Cu<sup>2+</sup> ion is 4 coordinated. By means of SQUID magnetometer, temperature dependence of magnetic susceptibility  $\chi$ (T) was measured in the range from 2 K to 300 K at the external field of 100 Oe and field dependence of magnetization M(H) was measured at 2 K and 300 K in the range from 0 to 7 T. From an analysis of the magnetization data of the polycrystalline powder sample using ITO (Irreducible Tensor Operator) method, we obtain the antiferromagnetic exchange coupling constants of J<sub>1</sub> = -143.65 K, J<sub>2</sub> = -143.78 K, J<sub>3</sub> = -143.61 K and the g-factor of g = 2.66. A further analysis with a model Hamiltonian including anisotropic exchange interactions is on-going.

**OF14** 

# Magnetization process of S=1/2 diamond chain compound Na,Cu,Ge,O<sub>1</sub>,

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Frustrated magnetism is one of the most fascinating systems because these systems have a great variety of quantum phenomena, and a lot of research activities are going on. Diamond chain in which spins exist on the Diamond-shaped lattice is one of examples of 1D-frustrated magnet. In this system, it is theoretically expected that not only the frustration but also low-dimensionality induces the quantum phenomena. Experimentally,  $Cu_3(CO_3)(OH)_2$  was reported as an S=1/2 Diamond chain compound, and its magnetic properties are studied actively [1]. Recently, Mo et al. reported the crystal structure of Na\_2Cu\_3Ge\_{Q\_{12}}, which has magnetic ions Cu<sup>2+</sup> (3d9:S=1/2) arranging in Diamond chain [2]. This compound has a shortrange order at T(SLO)–10K and a long-range order at T(N)=2K by means of magnetic susceptibility and heat capacity measurements [2][3]. We performed high field magnetization measurement using pulse magnetic field up to 5ST at T(N)<T-T(SLO). The magnetization curve increase convex downward up to one-third of saturation value, and clearly exhibit a plateau behavior with increasing magnetic field more than 30T. The convex downward curve indicates the components of S=1/2 1D antiferromagnetic nhenomena with a spin reduction. More detailed results will be discussed in the conference.

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# QF15

#### **Quantum criticality in a frustrated ising chain columbite** Kazuhiro Igarashi, Yasuhiro Shimizu\* and Masayuki Itoh

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A ferromagnetic Ising model in transverse magnetic field is a textbook example of the quantum phase transition. Nevertheless, few experiments have been accessible to the quantum criticality.  $CoNb_2O_6$  with a triangular lattice of ferromagnetic Ising chains is the highlight compound that has been recently found to exhibit the emergent symmetry in the excitation spectrum. Here we present the NMR observation of the low-energy quantum critical behavior in  $CoNb_2O_6$ . We find that the dynamical critical exponent toward the quantum critical point exhibits close to that of quantum antiferromagnets. Namely, the nuclear spin-lattice relaxation rate is independent of temperature at the critical transverse magnetic field of 5.4 T. The strong field dependence is manifested only in the dynamical susceptibility, whereas the uniform part is almost independent of magnetic field. Our results can be served as the ideal Ising spin dynamics on the quantum phase transition.

# QF16

#### Ligand-driven geometric and electronic structures of FeII spincrossover molecules

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Transition-metal complexes that exhibit spin-crossover (SCO) between low-spin (LS) state and high-spin (HS) state are now potential candidates for many applications such as display and memory devices. In these applications, thermal hysteresis in SCO of transition-metal complexes is required. Our previous study demonstrated that there is a correlation between thermal hysteresis in SCO and the electrostatic-energy difference ( $\Delta U$ ) of transition-metal complexes [1]. In this work, in order to explore more about the role of ligands in tailoring SCO behavior of transition-metal complexes, we study SCO of a series of FeII molecules with different ligand configurations [2-4]. Our results showed that the  $\Delta U$  of transition-metal complexes can be determined by pKa constant of ligands. These results should be helpful for designing new SCO complexes.

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### QF17

# Magnetic properties of S=1/2 zigzag antiferromagetic chain compounds, $VO(XO_4)(2,2'-bpy)$ (X=S, Mo; bpy = bipyradine)

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One dimensional Heisenberg antferromagnets (1D HAF) with the nearest J1 and second-nearest J2 interactions (zigzag chain) are the simplest frustrated magets. We measured magnetic susceptibility, specific heat and high field magnetization up to about 60 T of VO(XO4)(2,2'-bpy) (X=S, Mo; bpy = bipyradine), model compounds of the zigzag chain. V<sup>++</sup> ions carry S =1/2 in these compounds. No magnetic long range order is observed in both compounds, assuring low dimensionality of the present compounds. The exchange interactions are determined by comparing the experimental data with calculated results based on the exact diagonalization to the finite spin chain up to N=12. For VO(SO<sub>4</sub>)(2,2'-bpy), data agree well with the theoretical curve for the simple ID HAF, thus J2 is found to be zero. On the other hand, for VO(MoO<sub>4</sub>)(2,2'-bpy), a ratio of J2 to J1 is estimated to be 0.2-0.3.

### QF18

# Inter-chain coupling and anisotropy in the frustrated chain cuprate Li,CuO,

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Frustrated spin-chains with nearest neighbor ferromagnetic and next-nearest neighbor antiferromagnetic interactions have recently attracted lots of interest. In this class of materials, Li<sub>2</sub>CuO<sub>2</sub> represents a reference system with a particularly simple structure. Here we present a detailed experimental study of this quasi-one-dimensional magnet with focus on the properties at low temperature where long range order sets in due to the present inter-chain interactions. We discuss the magnetic phase diagram as obtained from magnetization, specific heat and lattice expansion measurements. Using inelastic neutron scattering we have further examined the magnetic excitations throughout a vast part of the Brillouin zone and in great detail, revealing an exceptional spectrum. In conclusion, we demonstrate that the low temperature dynamics of this S=1/2 chain material is well described within spin-wave theory and that the observed ground state is largely due to the exchange anisotropy.

# QF19

#### Spectral signatures of magnetic Bloch oscillations in 1D ferromagnets Sergey Shinkevich and Olav F. Syljuasen Department of Physics, University of Oslo, Norway

Domain walls in a one-dimensional gapped easy-axis ferromagnet can exhibit Bloch oscillations in an applied magnetic field. We investigate how exchange couplings modify this behavior within an approximation based on non-interacting domainwall bound states. In particular we obtain analytical results for the spectrum and the dynamic structure factor, and show where in momentum space to expect equidistant energy levels, the Wannier-Zeeman ladder, which is the spectral signature of magnetic Bloch oscillations. We compare our results to previous calculations employing a single domain-wall-approximation, and make predictions relevant for the material cobalt chloride dihydrate (CoCl<sub>2</sub>•2H<sub>2</sub>O).

 Jordan Kyriakidis and Daniel Loss, Phys. Rev. B 58, 5568 (1998). 2) W. Montfrooij, G. E. Granroth, D. G. Mandrus, and S. E. Nagler, Phys. Rev. B 64, 134426 (2001). 3) N.B. Christensen, K. Lefmann, I. Johannsen, and O. Jørgensen, Physica B 276, 784 (2000).

# QF20

# Quasi-one-dimensional magnetic phase as a competing ground state in a frustrated magnet

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Low-dimensional physics is interesting because it is driven by fluctuations and it may result in exotic phenomena, like the fractionalization of quantum numbers and spinons in magnetic chains. In magnetic systems, the reduction of dimensionality occurs due to anisotropic short-ranged interactions. Generally, transition metal oxides order at low temperatures. Here, in geometrically frustrated Ca<sub>3</sub>Co<sub>2</sub>O<sub>6</sub>, we show that by geometric frustration a stable quasi--one-dimensional (Q1D) magnetic phase emerges from a competition of degenerate ground states. This phase coexists and competes with the 3D antiferromagnetism in a large range below the 3D ordering temperature. We show that (i) constituent entities of this Q1D phase are shaped as rods in real space with alternating up and down magnetized sections (d=15A, l=320A at 15 K) along the chain axis where (ii) the sections belong to a degenerate ferrimagnetic class of ground states, (iii) these magnetic rods are arranged randomly in the frustrated triangular plane, (iv) they occupy a significant volume fraction (estimated 25% of moments at 15K) and (v) they are correlated with the magnetic response of this material. These results provide a new general possibility of a frustration-stabilized one-dimensional phase in physics.

# QF21

# Crossover of magnetic relaxation from 2D-spin ice like state to ordered state in layered single molecular magnet networks

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Layered 2D single molecular magnet (SMM) network [Mn(saltmen)]4[M(CN)<sub>k</sub>]ClO<sub>4</sub> nH<sub>2</sub>O (M=Mn, Fe[1]) is a strongly frustrated system. The each SMM unit [Mn(saltmen)] (ST =4) is ferromagnetically connected through the [M(CN)6] unit (S=1 for Mn, and S=1/2 for Fe) within a layer. The origin of the frustration [2] is a competition between the orthogonally oriented two uniaxial axes of the SMM unit and the ferromagnetic interaction, leading to a macroscopically degenerated ground state as the pyroclore 3D-spin ice. Our previous work [2] for M=Mn showed no transition down to 2 K through the Weiss temperature  $\Theta$ = 8.3 K and Davidson-Cole (DC) type magnetic relaxation as 3D-spin ice [3]. However, for M=Fe, a distinct anomaly of relaxation time is observed at 2.6 K ( $\Theta$  = 7.5 K), and the relaxation is the Cole-Cole (CC) type, suggesting a magnetic order. In this work, we measured ac susceptibility of M = Mn1-xFex (x = 0 ~ 1.0). We found that the relaxation is not the DC type for x>0.3 accompanying the distinct anomaly of relaxation time. The results suggest that 2D-spin ice. The set is ordered into a randomly frozen state above x ~ 0.1.

[1] H. Miyasaka et al., J. Am. Chem. Soc. 118 (1996) 981. [2] T. Saito et al. (2012), in preparation (partly presented in ICM 2009). [3] K. Matsuhira et al., J. Phys. Condens. Matter 13 (2001) L737.

# QF22

Magnetic properties of the novel low-dimensional spin-1/2 magnet  $\alpha$ -Cu<sub>2</sub>As<sub>2</sub>O<sub>7</sub>

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In copper pyroarsenate  $\alpha$ -Cu<sub>2</sub>As<sub>2</sub>O<sub>7</sub>, the quantum spins S = 1/2 localised at the Cu sites in the structurally well defined dimer chains interact predominantly in one spatial dimension (one-dimensional spin systems). Magnetic properties of this new material have been thoroughly investigated by means of magnetization, heat capacity, electron spin resonance, and nuclear magnetic resonance techniques, as well as by density functional theory (DFT) calculations and quantum Monte Carlo (QMC) simulations. The data reveal that the magnetic Cu-O chains in the crystal structure represent a realization of a quasi-one-dimensional (1-D) coupled alternating spin-1/2 Heisenberg chain model with relevant pathways through nonmagnetic AsO4 tetrahedra. Owing to residual 3-D interactions, antiferromagnetic long range ordering at TN ~10 K takes place. The experimental data suggest that substantial quantum spin fluctuations take place at low magnetic fields in the ordered state. DFT calculations and QMC fits enable quantitative evaluation of the exchange couplings. We conclude that the electronic state of the central ion in the nonmagnetic AsO<sub>4</sub>, side groups plays a crucial role for determining the relevant interchain pathways which makes ar-Cu<sub>4</sub>As<sub>4</sub>O<sub>7</sub>, distinct in its magnetic properties from other representatives of this class of compounds [1].

[1] Y.C. Arango, E. Vavilova, M. Abdel-Hafiez et al., Phys. Rev. B 84, 134430 (2011)

# QF23

#### Low temperature magnetic properties of the dilutable frustrated spinladder Bi(Cu<sub>1,x</sub>Zn<sub>x</sub>),PO<sub>6</sub>

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Quantum antiferromagnets with a spin-ladder magnetic network show an energy gap in the spin excitation spectrum and may possess an intrinsically disordered (spin liquid) ground state. BiCu<sub>2</sub>PO<sub>6</sub> (BCPO) shows a structure where Cu<sup>2+</sup> magnetic ions (S=1/2) form two-leg zigzag ladders along the b-axis, which closely corresponds to theoretical quasi-1D AF spin ladder model. We can reveal quantum nature of low temperature behavior by studying the effect of non-magnetic impurities (e.g. Zn<sup>2+</sup>, Ni<sup>+</sup>) introduced at magnetic Cu-sites. Recently, single crystals of Bi(Cu<sub>2</sub>, Zn<sub>2</sub>, PO<sub>6</sub> (O<Cu05) were successfully synthesized and studied by NMR for the first time [1,2]. We present here the specific heat Cp (H,T), ac susceptibility (f,T) and dc magnetization M(H,T) on single crystal sample of the undoped BCPO and 5% Zn-doped BCPO (Zn005BCPO) with field applied along main crystallographic axes focusing on low temperatures T < 10 K. We argue that the ground state of the Zn doped compound (Zn005BCPO) has features of both spin-glass and long range magnetic order, and that glassiness is lost under application of higher magnetic fields. We show the H-T magnetic phase diagram.

[1] SWang, E. Pomjalashina, I. Shinoka, G. Deng, N. Nikzeresht, Ch. Ruegg, H.M. Rømnow, and K. Conder. J.Crystal.Growth 313, 51-55 (2010) [2] F. Casola, T. Shiroka, S. Wang, K. Conder, E. Pomjakashina, J. Mesor and H.-R. Ott, Phys. Rev. Lett. 105, 067203 (2010)

# QF24

# 51V-NMR study of the quasi-one-dimensional antiferromagnet $BaCo_2V_2O_8$

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The double rectangle-like powder pattern has been observed in the 51V-NMR spectrum for the quasi-one dimensional antiferromagnet  $BaCo_2V_2O_8$  below TN = 5.4 K. The powder pattern indicates the existence of two V sites with different internal field, 2.1 and 3.8 kOe, although V atoms occupy only one magnetically symmetric site where the scalar type transferred field from the surrounding  $Co^{2+}$  magnetic moments is zero. The internal field at V site is explained by taking into account of the classical and the pseudo-dipolar fields from  $Co^{2+}$  magnetic moments. In the paramagnetic state, the nuclear spin-lattice relaxation is dominated by the antiferromagnetic spin fluctuation via the dipolar field, which is proved by the linear relation between the nuclear spin-lattice relaxation rate divided by temperature 1/T1T and the magnetic susceptibility. The change in the slope of 1/T1T against the magnetic spectrum. Below 60 K, 1/T1 shows the thermal activation-type temperature dependence signaling the formation of spin gap.

# **OF25**

#### Magnetic properties of one-dimensional chain of O<sub>2</sub> confined in nanospaces of MFI-zeolite

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We have studied the molecular arrangement and the magnetic properties in O2 adsorbed in nanoporous compounds [1-3]. We found that the O<sub>2</sub> adsorbed in MFI zeolite form a one-dimensional chain. The O<sub>2</sub> adsorption isotherm shows a plateau and the step, indicating a phase transition. It is found by means of the x-ray diffraction measurements that the zeolite deforms from monoclinic to orthorhombic corresponding to the step of adsorption isotherm. In the low-adsorption case, the adsorbed O2 aligns linearly in the straight and sinusoidal channels in MFI zeolite. In the high-field magnetization measurements, the magnetization process does not show the linear field dependence characteristic to one-dimensional Heisenberg antiferromagnet and presents the metamagnetic transition similar to the case of O2-O2 dimer reported previously [1, 2]. This result may suggest that the singlet ground state is realized at low fields and the fieldinduced orientational change of molecular axis occurs. In the high-adsorption case, the adsorbed O<sub>2</sub> seems to align in rows of two. The temperature dependence of susceptibility shows the first-order transition with large hysteresis. By XRD measurement, the structural phase transition from orthorhombic to monoclinic with decreasing temperature was found, which is not observed in N2 adsorbed compound.

[1] T. C. Kobayashi et al., Prog. Theor. Phys. Suppl. 159 (2005) 271. [2] A. Hori et al., J. Phys. 200 (2010) 0022018. [3] A. Hori et al., J. Low. Temp. Phys. 159 (2010) 122.

# **OF26**

Quantum spin transport in a Heisenberg spin chain Nan-hong Kuo<sup>1</sup>, Sujit Sarkar<sup>2</sup> and Chong Der Hu<sup>1</sup>\* <sup>1</sup> Physics NationalTaiwan University Taiwan <sup>2</sup> PoornaPrajna Institute of Scientific Research, India

# WITHDRAWN

### **OG01**

#### Magnetic transition of plastic deformed Si-doped Ni<sub>3</sub>Mn alloy

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Structural and magnetic properties of plastic deformed Si-doped Ni<sub>3</sub>Mn alloy were studied and compared with those of undeformed specimen investigated by powder neutron diffraction and diffuse purely magnetic scattering measurements in which coexistences of ferromagnetic and antiferromagnetic interactions were still existing even far above Neel temperature. Two specimens showed spin glass properties and disordered face-centered cubic structure but lattice parameters were a little different value of 3.5961 Å and 3.5890 Å. Magnetic properties showed very different Neel temperatures of 87.5 K and 62.5 K respectively. The paramagnetic Curie temperature appeared by -325 K and 125 K and therefore effective magnetic moments were calculated by 3.51 µB and 0.41 µB respectively. The reasons for different magnetic transition could be introduced by the elongation of ferromagnetic clusters and change of atomic environments effect after plastic deformation. This report shows two kinds of magnetic properties by plastic deformation

# **OG02**

#### Incommensurate-commensurate phase transition in TbNis induced by external magnetic field

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The rare earth intermetallic RNi5 compounds are attractive subjects for the study of the exchange and crystal field interactions due to their simple crystal and magnetic structures. The magnetic ordering in TbNi5 compound has been interpreted in terms of FAN-like incommensurate structure at temperatures between 23 K and 10 K. On cooling sample down to T = 10 K order-order type magnetic phase transition from the incommensurate phase to a "lock-in' phase takes place. In the present work we report the results of neutron diffraction experiment on TbNi5 single crystal sample in external magnetic field up to  $\mu$ 0H = 1T. Abrupt change of the intensity of (001) Bragg peak and satellites (001)± occurs within temperature interval 7 K -10 K due to magnetic phase transition from the incommensurate phase to a "lock-in' phase. It has been found that the application of external magnetic field up to  $\mu 0H = 0.1T$  at the temperature T = 11Kalong [100]-axis of TbNi5 single crystal recovers zero-field low temperature 'lockin' magnetic structure. The external magnetic field increase up to  $\mu$ 0H = 0.4T leads to transformation of the modulated magnetic structure to a ferromagnetic structure.

#### **OG03**

### First-principles dynamical CPA study of ferro- and antiferromagnetism of transition metals

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Quantitative explanation of the finite-temperature properties of 3d transition metals and alloys has been a long standing problem in metallic magnetism because their Coulomb interactions are comparable to the band widths and simple perturbation approach is not applicable to such systems. We present here the first-principles dynamical CPA combined with the LDA+U Hamiltonian toward quantitative calculations of magnetic properties at finite temperatures and explain quantitatively or semi-quantitatively the finite-temperature magnetism of 3d transition metals from V to Ni. We obtained the Curie-Weiss spin susceptibility with meff=1.8 muB (Expt. 2.1 muB) in V, the Pauli paramagnetic susceptibility in Cr being in agreement with the experimental data, as well as the Curie-Weiss susceptibilities in Fe, Co, and Ni whose meff quantitatively agree with the experimental data (3.2, 3.2, and 1.6 muB). Calculated ground-state magnetizations are 2.58, 1.72, and 0.64 muB for Fe, Co, and Ni, being in good agreement with the experimental data (2.2, 1.74, and 0.62 muB). On the other hand we obtained the Curie temperatures 1900K, 2100K, and 620K for Fe, Co, and Ni which are somewhat overestimated because of the single-site approximation. We also clarify the systematic change of the single-particle DOS in these systems.

#### **OG04**

#### High-coercive metastable ferromagnetic state induced in the Ising antiferromagnet Fe0.5 TiS2

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The transition metal (T) dichalcogenides TX2 (X = S, Se, Te) with a layer crystal structure being intercalated with 3d-transition (M) metals demonstrate different magnetic states depending on the type and concentration of M-atoms as well as on the parent TX2 matrix [1]. In the present work, the magnetic susceptibility, magnetization, electrical resistivity, magnetoresistance and powder neutron diffraction measurements have been performed for the Fe intercalated compound Fe05TiS2. It has been shown that this compound exhibits an antiferromagnetic (AF) ground state below the Neel temperature TN = 140 K. Application of the magnetic field at T < TN induces a metamagnetic phase transition to the ferromagnetic (F) state, which is accompanied by the large magnetoresistance effect (~ -27 %). The field-induced AF-F transition is found to be irreversible below 100 K. At low temperatures, the magnetization reversal in the metastable F state is accompanied by substantial hysteresis (up to 100 kOe) which is associated with the Ising character of Fe ions. This work was supported by the RFBR (project 12-02-00778) and by the program of the Ural Branch of RAS (project No 12-T-1012).

[1] Baranov N. V. et al. J. Phys. Condens. Matter. 17 (2005) 5255

# **OG05**

#### Magnetic and thermoelectric properties of the solid solutions Mn<sub>1</sub>. vNivS

Sergey Aplesnin<sup>1</sup>, Oksana Romanova<sup>1</sup>, Ludmila Ryabinkina<sup>1</sup>, Olga Demidenko<sup>2</sup>, Anatoly Galvas<sup>2</sup> and Kazimir Yanushkevich<sup>2</sup> <sup>1</sup> L.V. Kirensky Institute of Physics, Russia

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In connection with intensive development of a microelectronics and spintronics, interest of researchers to the magnetic semiconductor materials showing magnetoresistivity effect and metal-insulator transition has increased. New sulphide solid solutions Mn<sub>1</sub> Ni<sub>x</sub>S (0<X<0.1) are synthesized. The structural, magnetic and thermoelectric properties of the obtained materials have been studied at the temperatures 77-1000 K in magnetic fields up to 10 kOe. The X-ray diffraction analysis has shown that the samples synthesized Mn<sub>1,x</sub>Ni<sub>x</sub>S have a NaCl-type FCC lattice and are observed four reflexes of a weak intensity associated with  $\gamma$  - modifications of MnS. Solid solutions Mn<sub>1</sub> Ni<sub>v</sub>S are antiferromagnetic with the Neel temperature (TN = 180 K for X=0.05 and TN= 200 K for X=0.1). The temperature dependence of magnetization is described by the Curie-Weiss law at T> TN. The effective magnetic moment are found to be µef. f=5,04 µB for X=0,05 and µeff.=5,16 µB for X=0,1. Change in the conductivity type from the hole to the electronic at X=0.05 is revealed on the basis of measurements of thermoelectric power.

#### **OG06**

#### Field induced anisotropy in NiMn and NiMnPt alloys Yildirhan Oner

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We report the detailed characterization of the magnetic properties of polycrystalline disks of Ni<sub>72</sub>Mn<sub>28</sub> and Ni<sub>72-x</sub>Mn<sub>28</sub>Pt<sub>x</sub> (x=1.0, 4.0, and 10.0) at 4.2 K induced by field cooling (FC). It is found that the FC-induced anisotropy field ,HK and coercivity , HC are strongly enhanced by the addition of nonmagnetic Pt impurities. The remanent magnetization in the direction of the initial applied field (in the disk plane) for each samples can be rotated from 0° to 1800 and back to 0° in various stationary fields above and below HK and the parallel component of the rotated remanents are measured. From the analysis of the angular dependence of ML, we show that these results can be accounted for by the coexistence of Mn(Ni)-rich and /or Mn(Ni)-deficient nano scale regions coupled antiferromagnetically. It is found that the unidirectional anisotropy originates from interfacial exchange interactions between these regions. Up to some critical angle rotation ( $\Theta$ C) relative to H, the unidirectional anisotropy field turns rigidly with the sample, while above  $\Theta C$ , the coupled regions become unstable and magnetically rearrange such that a unidirectional anisotropy is induced along H.

#### **OG07**

#### X-ray-absorption near-edge structure and X-ray magnetic circular dichroism studies of a Lu<sub>2</sub>Fe<sub>165</sub>Ru<sub>0.5</sub> single crystal

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The magnetism of R2Fe17 with the smallest rare-earth R=Lu (hexagonal Th2Ni17 structure type) is essentially associated with its Fe sublattice behavior. Strong contribution of the local antiferromagnetic (AF) Fe-Fe interactions in the overall ferromagnetic (F) exchange considerably lowers the ordering temperature (TN=274 K) and balances the non-collinear antiferromagnetic phase above the Curie temperature TC=130 K. Generally, various doping elements (M) in Lu<sub>2</sub>Fe<sub>17-x</sub>M<sub>x</sub> suppress fully the AF states by strengthening the ferromagnetic ones accompanied by a dramatic TC increase. However, recently a unique case of antiferromagnetism stabilization in the whole range of magnetic ordering in Lu<sub>2</sub>Fe<sub>17-x</sub>Ru<sub>1x</sub> was reported [1], and this work is dealing with the local studies of intrinsic magnetic properties of Lu2Fe165Ru05. A single-crystalline sample of Lu2Fe165Ru05 has been investigated by X-ray-absorption (XANES) and X-ray magnetic circular dichroism (XMCD) techniques in the applied magnetic field up to ±3 T at 7 K. By probing the lutetium L2,3-, iron K- and ruthenium L2 absorption edges, information on magnetic states of the constituents of Lu<sub>2</sub>Fe<sub>165</sub>Ru<sub>05</sub> is obtained. Small but noticeable induced magnetic moments with an antiparallel orientation to Fe moments have been registered on Lu and Ru sites

[1] E.A. Tereshina et al., J. Appl. Phys. 105 (2009) Art. No. 07A747.

# OG08

#### Study of the metamagnetic behavior of Ni-Co-Mn-Sb alloy in high magnetic fields

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In order to investigate metamagnetic behavior and temperature dependence of equilibrium magnetic field of the martensitic transformation for Ni-Co-Mn-Sb, magnetization experiments up to 18 T static magnetic field and up to 40 T pulsed magnetic field were performed. Polycrystalline specimen of Ni, Co, Mn, Sb, was prepared by induction melting, and annealed at 1173 K for 1 day. Transformation temperatures, TM, and entropy change AS between the martensite phase and parent phase were examined with DSC. Magnetic measurements were performed with using a SQUID magentometer up to 5 T, by extracting method up to 18 T static magnetic fields, and also by induction method up to 40 T pulsed magnetic fields. In the thermomagnetization curves, it was observed that TM decreased with increasing the applying magnetic field, H, with a rate of dTM/dH = 3.8 K/T. An estimated value of  $\Delta$ S from the Clausius-Clapevron relation, is about 15.8 J/K-kg, being in good agreement with the value obtained by the DSC. For the isothermal magnetization curves, metamagnetic behavior associated with the magnetic field induced martensitic reverse transformation was observed. Arresting behavior of the martensitic transformation was confirmed in lower temperature, as similar to the Ni(Co)-Mn-Z (Z = In, Sn, Ga) alloys.

### **OG09**

### Magnetic Transition and Thermal Expansion in LaFe<sub>13-x-v</sub>Co<sub>v</sub>Si<sub>x</sub>

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Materials with a large magnetocaloric effect (MCE) continue to attract significant attention for cooling applications with particular interest in the isothermal entropy changes at first-order magnetic transitions [1]. The large entropy change in LaFe12, Si, is associated with negative lattice expansion and metamagnetic transition behaviour [2]. We have investigated LaFe<sub>13-x-y</sub>Co<sub>y</sub>Si<sub>x</sub> to clarify the influence of magneto-structural coupling on the MCE. The Curie temperatures of LaFe<sub>13</sub>, Si, increase with Si content from TC=219 K for x=1.6 to TC=250 K for x=2.6 with further enhancement to TC=281 K on substitution of Co to LaFe, CoSi, A pronounced positive spontaneous volume magnetostriction has been observed below TC and the anomalous thermal expansion attributed to the volume dependence of the magnetic energy. Our results show that the magnetic transition changes from first order for LaFe114Si16 to second order for LaFe104Si26 and LaFe04CoSi26. The different natures of the magnetic transitions in LaFe12, Co.Si, have been discussed in terms of the classical model for itinerant ferromagnets and the volume dependence of the magnetic energy

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# **OG10**

### The magnetovolume effect of Y<sub>2</sub>Fe<sub>17-x</sub>Ga<sub>x</sub>

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The unit cell volume and the Curie temperature of Y2Fe17 increases by substitution of Ga for Fe. This phenomenon is understood as a property of Invar materials, dTc/dV > 0, because the compounds R<sub>2</sub>Fe<sub>17</sub> have Invar properties. However, in the substitution of Si for Fe, dTc/dV is negative. Kamimori et. al. proposed damping of the Invar properties by these substitutions. In this work, the magnetovolume effect of  $Y_2Fe_{17}$ xGax was investigated by measuring temperature dependence of unit cell volume and forced volume magnetostriction. X-ray diffraction was used to measure unit cell volume at under R.T. and thermal expansion meter at above R.T.. The forced volume magnetostriction was measured by three terminals capacitacne method under magnetic field up to 1.5T. The forced volume magnetostriction at R.T. is positive. This shows Y2Fe17-xGax compounds have Invar property. However, the magnitude of the forced volume magnetostriction decreases violently by the substitution of Ga for Fe. This tells us the substitution of Ga for Fe in R2Fe17 decreases magnetovolume effect and attenuates Invar characters of the compounds

# **OG11**

#### Pressure effect on the electrical resistivity of $La_{1.09}(Fe_{0.845}Si_{0.155})_{13}$ compound

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**OG12** 

### One- and two-magnon and exciton raman scattering in antiferromagnetic CoF2: Experiment and theory

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Experimental data are reported for the temperature and polarization dependence of the one- and two-magnon Raman light scattering in the rutile-structure antiferromagnet CoF<sub>2</sub> (Neel temperature = 38 K). A detailed analysis of the one-magnon Stokes and anti-Stokes Raman spectra is presented resulting in comprehensive experimental data for the temperature variation of the one-magnon frequency, line width, and integrated intensity. A similar analysis was carried out for the two-magnon scattering. The lowlying excitons were also investigated at low temperatures and comparisons are made with results from earlier Raman, infrared, and neutron scattering work. The Green's function equation of motion method was employed to derive the excitation energies and spectral intensities over a broad range of temperatures for the magnetic excitations in a spin S = 3/2 anisotropic antiferromagnet with strong spin-orbit coupling. Results were obtained using RPA for the product of operators at different sites while the singleion anisotropy terms were treated exactly by generating a closed set of coupled Green's function equations. The numerical results are in good agreement with one-magnon Raman light scattering data. Results were deduced for some exchange and anisotropy parameters and the relative magnitudes of the magneto-optical coefficients in the interaction Hamiltonian

# **OG13**

#### Thermodynamic and transport properties of $Ru_{2,x}Fe_xCrSi$ (1.3 $\leq x \leq$ 1.8)

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Full-Heusler compounds with the generic chemical formula X<sub>2</sub>YZ (X and Y are transition elements, Z is sp element) attract much interests because this system shows a number of physical properties and some of them have a potential for the technological applications. The ferromagnetic Heusler compound Ru2-xFexCrSi is one of candidate materials expected a half metallicity1] and this system has relatively high Curie temperature TC~500 K for x > 1.4 2]. We have carried out electric resistivity  $\rho(T)$  and specific heat Cp(T) measurement of Fe-rich  $Ru_{2,x}Fe_xCrSi$  (x = 1.3, 1.6, and 1.8). For all samples,  $\rho(T)$  shows metallic behavior ( $\partial \rho/\partial T$ ) in 20 < T < 300 K, and upturn at around 20 K. In low-temperature range, Cp(T) can be described by Debye's law. Although a coefficient of lattice specific heat  $\beta$  is insensitive for Fe concentration x, a electronic specific heat part  $\gamma$  decreases slightly with increasing x. For each x sample, we estimated total density of states (DOS) at the Fermi level D(EF) from  $\gamma$ .

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# **OG14**

#### Magnetic structure and excitations of the one-dimensional quantum antiferromagnet RbCoCL

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RbCoCl<sub>3</sub> is a hexagonal perovskite in which the  $Co^{2+}$  ions have an effective spin S = 1/2 and form onedimensional (1D) chains perpendicular to the basal plane. Measurements of the magnetic susceptibility confirm its antiferromagnetic behaviour; the difference between the maximum of the susceptibility and the Neel temperature indicates 1D magnetic correlation. Isostructural materials, such as CsCoCl<sub>3</sub> and TlCoCl<sub>3</sub>, are quasi-1D Ising-like spin systems[1][2]. The lowest-lying magnetic modes in their spectrum consist of pairs of domain-wall excitations and are described by a strongly anisotropic Heisenberg Hamiltonian[3]. Here we report the first results of a comprehensive neutron scattering study of the magnetic structure and excitations in RbCoCl. A neutron diffraction experiment on HRPT (SINO.PSI.Switzerland) revealed the existence of two phase transitions at TN1=28K and TN2=14K. At TN1 there is a paramagnetic-to-antiferromagnetic transition, while at TN2 a low-temperature phase with a new propagation vector develops. Inelastic neutron scattering experiments were performed on LET (ISIS, RAL, UK) and IN22 (ILL, Grenoble, FR). Between TN1 and TN2 as well as above TN1, excitations characteristic for an Ising magnet are observed. Below TN2 an additional splitting of the modes is discovered. Results and analysis of the data from LET will be presented and discussed.

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#### **OG15**

# Antiferromagnetic transition in Ru<sub>2</sub>CrSi in magnetic fields

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Heusler compounds Ru<sub>2-x</sub>Fe<sub>x</sub>CrSi for Fe-rich region were recently found to be ferromagnetic and to be candidates for a half-metal, where conduction electrons are fully spin-polarized.[1] In the Heusler compound Ru2CrSi, on the other hand, an antiferromagnetic (AF) transition at TN =14 K was revealed by specific heat and magnetization measurements.[2] In this study electrical resistivity is measured in magnetic fields up to 14.5 T and the AF transition in magnetic fields is investigated. In the temperature dependence of the resistivity at zero field a clear kink is observed at TN =14 K. With decreasing temperature the resistivity turns upward from the kink at TN. This kink can be regarded as an indication of the AF transition. With increasing magnetic field the resistivity only slightly changes and even at 14.5 T the clear kink is observed. From the kink TN is found to decrease just a little in high magnetic fields. TN at 14.5 T is lower than TN at zero field by 0.5 K. These results demonstrate that the AF state in Ru2CrSi is unusually unaffected by strong magnetic field.

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### **OG16**

#### Assessment of Curie temperature by magnetization vs. temperature (M-T) scans in a granular magnetic Cu-Fe-Ni alloy

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The present study has investigated the Curie temperature of Cu-Fe-Ni ternary alloy which is known as a granular giant magneto-resistance (GMR) material with a high magneto-resistance (MR) ratio. At first, we surveyed the relationship between the microstructure and the aging temperature of precipitates formed in Cu-Fe-Ni alloys isothermally annealed at 873K and 1073K for various aging times, using a transmission electron microscopy (TEM). According to TEM observations, we have revealed that two or more nano-scale particles with a cubic shape and rod shape precipitates were formed at 873K and 1073K, respectively. The critical temperature of the drastic change of the microstructure coincides with the Curie temperature of magnetism of the alloy. And, the Curie temperature has been determined by the magnetization (M) vs. temperature (T) scans, using the SQUID magnetometer. The experimentally determined temperature of 1028K was as high as a bulk Fe-Ni alloy, even though nano-scale ferromagnetic particles are finely distributed in copper matrices. The high Curie temperature has, however, well explained the temporal evolutions in both microstructure and magnetic properties occurring during isothermal annealing.

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# **OG17**

#### Design strategy for strongly coupled diradicals: Systematic Approaches of intramolecular magnetic interactions

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A new series of neutral radicals (DP1-DP6) based on diazaphenalenyl were designed via heteroatomic modifications. As spin sources, the designed radicals were implemented in four different diradical model systems (Model I, II, III, and IV) by changing the reference radical and the linkage, and their magnetic interactions between the designed radicals and the reference radical were investigated by using density functional theory calculations. The trend in strength of magnetic interactions of radicals was found to be identical in different model systems. In particular, as a new family of spin source radical, DP3 could be a potential candidate in designing new organic magnetic materials due to its strong magnetic coupling and high stability in diradical systems.(1) We divided 8 monoradicals into  $\alpha$ -group and  $\beta$ -group according to Mulliken spin density values of the connected atoms. The overall trends in the strength of magnetic interactions of diradicals were found to be identical in three different model systems (Model I, II and V). NN-ethylene-PO was calculated to have the strongest magnetic coupling constant with ferromagnetism, and even stronger (more than twice) than NN-ethylene-NN which was reported to have strong antiferromagnetic interactions in a previous experiment.

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### **OG18**

#### Single crystal growth and physical properties of metalic antiferromagnet (Mn,Fe)<sub>3</sub>Si

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Heusler alloys Mn<sub>3-x</sub> Fe<sub>x</sub> Si(x=0~3) have a face centered cubic (DO3 type) structure in the whole x-region from the antiferromagnet Mn<sub>3</sub> Si(TN~23K) to the ferromagnet Fe<sub>3</sub> Si(TC~850K). As a characteristic point, Mn3Si shows anomalous magnetic excitations so-called "chimney-type" ones, which is difficult to understood by normal spin waves, while Fe3Si shows typical ferromagnetic spin wave excitations. Our final goal is to clarify how the spin dynamics qualitatively changes from antiferromagnet to ferromagnet in this (Mn,Fe)3Si system, so that a deep insight into the origin of magnetism would be gained. The purpose of current work is to grow high quality large single crystals of Mn<sub>2,x</sub> Fe, Si for inelastic neuron scattering experiments, which can reveal generalized magnetic susceptibility directly. Besides, for accurate measurement of bulk properties, such as electric resistivity  $\rho(T)$  and magnetic susceptibility  $\chi(T)$ , minimizing amount of impurities is important. So far, we have succeeded in growing large single crystals of x=0, 0.2, 0.6, 1.0, and 1.5 by the Bridgman method. At present, an impurity phase (several-% order) is still observable in X-ray powder diffraction patterns. In this work, we report the details of single crystal growth of  $Mn_{3-x}$  Fe<sub>x</sub> Si, assignment of impurity phase, and resultant  $\chi(T)$ , and  $\rho(T)$ .

### **OG19**

#### AC magnetic measurement of LiFeAs at pressures up to 5.2 GPa: Verification of the relation between Tc and structural parameters

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We have already conducted AC susceptibility measurements at pressures up to 1.5 GPa and X-ray diffraction experiment up to 20 GPa for a 111 type of iron-based superconductor, LiFeAs[1]. Now, we have enlarged the pressure range of the AC susceptibility measurement for the same sample up to 5.2 GPa. The superconducting transition temperature Tc decreased from 17 K at 0 GPa to 10 K at 5.2 GPa. Tc has been related to two structural parameters, an As-Fe-As bond angle  $\boldsymbol{\alpha}$  and the height of As from the Fe layers (hAs): Lee et al. have suggested that Tc has the highest value at the optimal  $\alpha$  of 109.47°[2]. Kuroki et al. have suggested that Tc increases with increasing hAs[3]. According to our previous study, α changes from 101.5° at 0 GPa to 97.8° at 17 GPa, and hAs changes from 0.154 nm at 0 GPa to 0.156 nm at 17 GPa. The present pressure-induced decrease in Tc is not explained by the change in hAs, but is qualitatively consistent with the change in a

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# **OG20**

#### Fundamental Magnetism of Fe-P Alloys and Fe3P Compounds: A density functional study

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Phosphorus (P) is quite harmful on the mechanical properties of steel. For instance, the decrease of the strength and plastic properties of G13L steel is nearly half when it contains 0.065 wt.% P compared to the case of 0.028 wt.% P steel. On the contrary, P is an important component in the making phosphor bronze which enhances wear and water corrosion resistant. On the purpose of solid solution strengthening. P containing interstitial free steel is very important in the automotive industry due to their excellent deep drawability. However, there is no systematic first-principles study done for understanding the Fe-P alloys and Fe3P compounds. Thus, we investigated fundamental magnetic and thermodynamic properties of ferrites with the various P concentrations  $(Fe_1, P_2, alloys, x = 0.037, 0.0625, and 0.125)$  and (L12, D03, and I-4 structured) Fe<sub>3</sub>P compounds using the highly precise all-electron full-potential linealized augmented plane wave method. As a result, the ferromagnetic state for I-4 structured Fe<sub>3</sub>P compound is more stable than the nonmangetic one by energy difference of about 809 meV. The calculated magnetic moments are rather overestimated compared to the neutron diffraction study. Detailed discussion on the magnetic and thermodynamic properties of Fe1-xPx alloys and Fe3P compounds will be given.

# **OG21**

### Analysis of spin-polaron formation in Hund lattices

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The ferromagnetic Kondo lattice model, also known as Hund lattice model, is a well-established model to study the interplay between charge and spin degrees of freedom in strongly correlated fermionic systems such as transition-metal oxides. Their electronic, magnetic and transport properties have displayed interesting physics such as the phenomenon of magnetoresistance [1]. Their properties depend as well strongly on the doping of the materials, which leads to formation of magnetically ordered regimes [2]. In our work, we study numerically the formation of spin-polarons in one-dimensional systems. We consider a Hund lattice model where a conducting band is coupled to a background of localized spins interacting antiferromagnetically with coupling constant J. Even though the zero-temperature of such lattice model has been intensively studied [3 .4], the physics of its phase diagram is still not exhausted. We investigated the ground state phase diagram as a function of the exchange coupling J and as a function of the band filling and found different evolving magnetic orderings. We report measurements of static properties such as spin structure factor. We study the quasi-particle formation and phase separation using the density-matrix renormalization group method, which is a highly efficient method to investigate quasi-one-dimensional strongly correlated systems[5].

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# **OG22**

#### Chiral magnetic orders in chiral helimagnet Cr<sub>1/3</sub>NbS<sub>2</sub>

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We directly present that chiral magnetic orders emerge in chiral helimagnet Cr. NbS, in small magnetic fields by means of low-temperature Lorenz transmission electron microscopy and small-angle electron diffraction technique. Based on precise analyses in both real and reciprocal space, we undoubtedly demonstrate that chiral magnetic soliton lattice (CSL) develops from chiral helimagnetic structure (CHM) with increasing the spatial period from 48 pm toward infinity in rising magnetic fields perpendicular to the helical axis. CSL and CHM do not exhibit any structural dislocation, indicating their high stability and robustness. This is because chiral magnetic orders are macroscopically induced by the uniaxial Dzyaloshinkii-Moriva (DM) exchange interactions that is allowed in Cr., NbS, hexagonal crystals belonging to noncentrosymmetric chiral space group. Theoretical sides of this talk will be given in this topic category by Prof. Jun-ichiro Kishine from the Open University of Janan

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# QG23

### ESR study of AFM - Ordering in the orthorhombic CuMnAs

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At the present time the Mn based compounds with elements of I or V group were studied intensively as advances materials for spintronic. One from such compounds is CuMnAs. CuMnAs is ordered antiferromagnetically above room temperature. We have studied electronic spin resonance (X-band) in this compound at temperature range 10 - 400 K At all temperature we observed symmetric Lorentz's shape resonance lines from Mn<sup>2+</sup> ions with g - factor 2.00. The lineswidth increase in 2 times at temperature decreasing to 10 K. The second line with a g-factor of 2.55 has appeared in the spectrum at temperatures below 60 K Non resonant absorption was observed for powder samples. We interpreted our results in terms Mn-Cu-Mn and Mn-As-Mn indirect superexchange interaction. Work was supported by Grants of RFBR and Presidium of the RAS.

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# QG24

Elastic properties and stability of Heusler compounds

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The physical properties of Heusler compounds which are promising materials for spintronics are important. Various Heusler compounds are investigated with respect to their malleability and stability. The elastic constants are calculated by applying isotropic strain in different ways to the cubic crystal. We used FP-LAPW implemented in Wien2K to calculate the elastic constants as well as further physical properties (bonding, hardness, velocity of sound, Debye temperature, etc.) that can be derived from the elastic constants. The result of the cubic elastic anisotropy can be used to decide the structural stability. Zener's ratios of the most stable compounds are in the range of 1.0 < Ae < 2.5. Ae < 0 is not stable and Ae = 0 is isotropic. We observed phase transitions in some of Heusler compounds with Ae < 0. It is shown that calculating elastic constants provides an estimate for studying phase transitions.

### QG25

# Electronic and magnetic properties of ferromagnet/dilute-magnetic semiconductor interfaces

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A. Continenza and G. Profeta Experimental findings [1-4] report exchange-bias effects at the Fe/GaMnAs together with interesting possible magnetic prossimity effects at the interface between the ferromagnetic (FM) material and the dilute semiconductor (DMS). In this work we perform ab-initio calculations of different FM/DMS geometries (superlattice systems and periodic slab) for Fe/Ge:Mn and Fe/GaAs:Mn interfaces. We examine different interfaces to assess the most stable structures, the magnetic coupling between the magnetic overlayer and the Mn dopant atoms and how the ferromagnetic overlayer affects the interaction between two Mn atoms dispersed in the semiconductor matrix. An antiferromagnetic coupling between the overlayer Fe and the Mn is found for Mn close to the interface layers, fading away with the distance of the Mn atom from the interface. A study of how the magnetic coupling between overlayer and dilute Mn and between dilute Mn's at different locations is performed and discussed. We find a rather large magnetic coupling enhancement in the case of Ge while the situation for GaAs looks to be much more complex due to chemical effects at the GaA/Fe interface. Results are discussed in terms of local densities of states, magnetic moments and total energies.

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# QG26

#### Mn-Sublattice of YbMn<sub>2</sub>Si<sub>2</sub>

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The instability of the f shell in Yb and the small energy difference between the Yb<sup>2+</sup> divalent and Yb<sup>1</sup> trivalent states, has led to the wide range of physical and magnetic properties exhibited by ytterbiumbased intermetallic compounds. The ranges of effects are linked with many physical states such as: mixed valence, Kondo effect, heavy fermions, spin wave excitations, superconductivity and magnetic quantum critical point [e.g. 1]. Trivalent YbMn<sub>2</sub>Si<sub>2</sub> exhibits a collinear antiferromagnetic AFil structure below TN1 = 526(4) K with a doubling of the magnetic c-lattice below TN2 = 30(5) K; this structure of the Mn-sublattice would therefore be expected to result in different environments for Yb atoms in the doubled unit cell [2]. However 170Yb Mossbauer studies led to the conclusion that the observed spectra are inconsistent with the cell-doubling [3]. To clarify this discrepancy, we have re-investigated the magnetic behaviour of YbMn2Si2 over the temperature range -0.4-50 K using the E6 neutron diffractometer, HZB. The high density of diffraction patterns confirm the cell doubled magnetic behaviour of the Mn-sublattice behaviour of the satellite reflections (propagation vector k = 00/2) described well by a power law.

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#### QG27

# Electronic structures and magnetic properties of full and half Fe-Mn-Ga Heusler alloys

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Electronic and magnetic structures of bulk and film Fe<sub>3</sub>MnGa and FeMnGa alloys near 2:1:1 and 1:1:1 stoichiometric compositions have been investigated theoretically and experimentally. It was experimentally found that unlike most of full stoichiometric Heusler compounds, bulk Fe<sub>36</sub>,4Mn<sub>34</sub>,2Ga<sub>25</sub>,4 alloy with the Curie temperature of 800 K and saturation magnetization of 4.86 µB/fLu crystallizes in the face-centered-cubic (FCC) ordered lattice of L12 (Cu<sub>3</sub>Au)-type with small admixture of B2-phase. According to results of our first-principle calculations of full electron energy feromagnetic metallic FCC phase (µFCC=6.11 µB/fLu) is energetically more preferable than ferrimagnetic half-metallic phase with L21 type of structure (µL21=2.04 µB/fLu) or ferrimagnetic metallic FCC phase (µFCC=0.48 µB/fLu). Metamagnetic transition in Fe<sub>36</sub>,4Mn<sub>34</sub>,2Ga<sub>35</sub>,4 alloy from antiferromagnetic to ferromagnetic phase was observed at low temperatures. Bulk Fe<sub>313</sub>,Mn<sub>320</sub>,Ga<sub>354</sub> alloy has hexagonal structure and saturation magnetization at T=15 K of µ=3.00 µB/fLu and the Curie temperature. Atomic disorder in amorphous state significantly reduces the film magnetization nearly to zero at room temperature. Depending on deposition conditions, unlike the bulk alloy, crystalline Fe<sub>3136</sub>,Mn<sub>323</sub>,Ga<sub>238</sub> alloy films exhibit BCC- or FCC-ordering that result in anomalous temperature dependencies of magnetization in low-temperature region.

#### **OG28**

#### Microscopic analysis of magnetic orders in MnP single crystals

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Manganese phosphide (MnP) has been extensively studied because of its richness of magnetically ordered phases in the magnetic phase diagram. In this study, as a first step to understand the mechanism of the magnetic phase transition from the ferromagnetic (FM) phase to the proper screw spiral phase or FAN phase, we have studied FM structures in the FM phase of MnP single crystals by means of low-temperature Lorentz transmission electron microscopy and small-angle electron diffraction[1]. In Lorentz Fresnel micrographs, striped FM domain structures were observed in external magnetic field less than 10 Oe in the specimens with the ab-plane in the plane of the specimen. From real- and reciprocal-space analyses, it was undoubtedly identified that striped FM domains oriented to the c-axis appear with Bloch-type domain walls in the b direction and arrange regularly along the a-axis with 64 nm separation. Moreover, the magnetic chirality reverses in alternate FM domain walls. These specific spin configuration of striped FM domains will affect the magnetic phase transition from the FM phase to the proper screw spiral phase at the low temperature or the FAN phase in the magnetic field in MnP.

[1] T. Koyama, S. Yano, Y. Togawa, Y. Kousaka, S. Mori, K. Inoue, J. Kishine, and J. Akimitsu, submitted.

# QH01

# Structural and magnetic properties of $Fe_2Mn_{0.5}Cu_{0.5}Al$ Nanocrystalline Alloys

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We synthesized  $Fe_2Mn_{0.5}Cu_{0.5}Al$  nanocrystalline alloys by mechanical alloying technique and studied as a function of milling time. The milling times were 1, 3, 6, 9, and 24 hrs. The effect of milling time on structure formation was investigated by X-ray Diffraction (XRD), and magnetic properties was investigated using vibration sample magnetometer (VSM). Alloy nature of  $Fe_2Mn_{0.5}Cu_{0.5}Al$  was confirmed after 24 hrs. The crystallite size after 24 hrs of milling time was estimated using Scherer formula, and the value was 24 nm. Magnetic saturation was decreased with increased milling time, which seemed be due to co-existence of Mn and Cu atoms. The magnetic saturation was 86.19 (emu/gr) and the coercivity was 376.59 Oe after 24 hrs of milling time. The coercivity (Hc) was increased with milling time presumably due to magnetic domain growth towards stable structure. Keywords: Heusler compound, nanocrytalline

References: [1] Tanja Graf, Claudia Felser, Stuart S.P. Parkin, Progress in Solid State Chemistry 39 (2011) 1-50 [2] Shane Stadler1, Mahmud Khan1, Joseph Mitchell1, Naushad Ali1, Angelo M. Gomes2, Igor Dubenko3, Armando Y. Takeuchi3, and Alberto P. Guimaraes3 Appl. Phys. Lett. 88, 192511 (2006)

### QH02

#### Antiferromagnetic resonance in the one-dimensional magnet IPACu(Cl<sub>1-x</sub>Br<sub>x</sub>)<sub>3</sub> (x=0.83)

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Low-dimensional antiferromagnets with a small spin value exhibit various fascinating properties. IPACuCl<sub>3</sub> and IPACuBr<sub>3</sub> (IPA: isopropyl-amine) having magnetic Cu<sup>2+</sup> (spin S=1/2) ions form the ladder chain in the b plane. The ground states of IPACuCl<sub>3</sub> and IPACuBr<sub>3</sub> are considered to be the Haldane state [1] and the singlet dimer state [2]. It was reported that IPACu(Cl<sub>3</sub>Br<sub>1,3</sub>)<sub>3</sub> (0.44 < x<0.87) undergo an antiferromagnetic longrang order (AFLRO) at low temperatures. However, the details of the AFLRO phase have not been clarified yet [3]. Therefore, we have performed electron spin resonance (ESR) measurement on single crystals of IPACu(Cl<sub>4</sub>Br<sub>1,3</sub>)<sub>2</sub> (x=0.83) in magnetic fields of up to 14 T. We have observed the typical antiferromagnetic resonance (AFMR) signals, and successfully explained all the observed resonance modes by the conventional AFMR modes with orthorhombic anisotropy.

 [1] T. Masuda et al., Phys, Rev. Lett. 96 (2006) 047210. [2] H. Manaka and I. Yamada, J. Phys. Soc. Jpn. 66 (1997) 1908. [3] H. Manaka et al., Phys, Rev. B 64 (2001) 104408.

#### QH03

#### Magnetization steps in Yb<sub>2</sub>Pt<sub>2</sub>Pb with the Shastry-Sutherland lattice Yasuyuki Shimura<sup>1</sup>\*, Toshiro Sakakibara<sup>1</sup>, Ken Iwakawa<sub>2</sub>, Kiyohiro Sugiyama<sub>2</sub> and Yoshichika Onuki.

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In the tetragonal compound Yb<sub>2</sub>Pt<sub>2</sub>Pb, Yb<sup>3+</sup> ions form the Shastry-Sutherland lattice [1]. The Yb moment shows a strong Ising type anisotropy with the easy axis lying along either [110] or [1-10] directions, and exhibit an antiferromagnetic long range order at TN=2.1 K [1]. Surprisingly, the recent H-T phase diagram investigation [2] has revealed that when a sufficiently strong magnetic field is applied along the [110] direction, the system exhibits a partially disordered state below TN; half of the Yb sites remain antiferromagnetically ordered whereas the rest become paramagnetic. This unusual behavior has been interpreted by a decoupling of two Yb sublattices in a [110] magnetic field due to the characteristic crystal structure and the competing Ising anisotropy [2]. In this paper, we measured the DC magnetization of Yb<sub>2</sub>Pt<sub>2</sub>Pb at very low temperatures down to 0.08 K for H/[100] and [110] in order to further explore its unusual phase diagrams. For both directions, a sequence of sharp magnetization steps is observed at T~0.5 K. Comparison of the differential susceptibility data for the two directions indicates the presence of an intersublattice coupling for H/[100].

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# **QH04**

#### Mossbauer study on Fe (Si,Ge) alloys

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Calculations of electronic structures in FeSi<sub>1-x</sub>Ge<sub>x</sub> with a cubic B20-type structure found that they were narrow-gap semiconductors in the non-magnetic state for x<0.5, however, the metallic state with the magnetic moment of about 1µB per Fe was shown to be more stable for x≥0.5 [1]. In this paper, we report the results of the magnetization measurements and 57Fe Mossbauer effect measurements on FeSi<sub>1-x</sub>Ge<sub>x</sub> compounds. The 57Fe hyperfine fields for x=0.6 and 0.8 estimated by fitting of Mossbauer spectra at 80K were 6.0T and 8.0T, respectively, while they were almost zero for 0≤x≤0.5. The Curie temperature and spontaneous magnetization at 4K were found to be 124K and 0.66µB/fu. for x = 0.6 and 171K and 0.71µB/f.u. for x = 0.7, respectively.

[1] H.Yamada, K.Terao and T.Goto, Physica B 344 (2004) 451.

#### QH05

# Ab initio and Monte Carlo investigations of the magnetic exchange and curie temperature of $Ni_2Mn_{1+x}Sn_{1-x}$ alloys

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The ab initio and Monte Carlo calculations of magnetic exchange constants, magnetic moment and the Curie temperatures of stoichiometric and non-stoichiometric Heusler  $N_{I_{00}}M_{D_{25rx}}S_{D_{25x}}$  alloys have been performed. For the ab initio calculations of magnetic exchange constants and magnetic moment the Munich SPR KKR package was used [1]. The calculations were made for order and disorder structures. We consider two types of disorder. First is a disorder between Ni and Mn atoms and second is a disorder between Mn and Sn atoms. The Curie temperatures have been calculated by mean field approximation and Monte Carlo method using the Heisenberg model. It is shown that the experimental Curie temperature of non-stoichiometric alloys can be explained theoretically by the assumption that in alloys the disorder between Mn and Sn atoms exists. Our calculations have shown that the magnetic moment and Curie temperatures for Ni-Mn-Sn alloys are in qualitative agreement with experimental magnetic moment and transition temperatures [2].

 H. Ebert, Electronic Structure and Physical Properties of Solids, ed. H. Dreysse, Springer, Lecture Notes in Physics, Berlin, 1999, Vol. 535, p. 191. 2. T. Kanomata, et al., Magnetic and crystallographic properties of shape-memory alloys Ni2Mn1+xSn1-x, Materials Science Forum 583 (2008) 119.

### QH06

# Structural, magnetic, magnetocaloric and magneto-transport properties in Ge doped Ni-Mn-Sb Heusler Alloys

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A. K. Nigama, Roshnee Sahoob and K. G. Sureshb\* aTata Institute of Fundamental Research, Homi Bhabha Road, Mumbai-400005, India bDepartment of Physics, Indian Institute of Technology Bombay, Mumbai-400076, India Abstract The effect of Ge substitution on the magnetic, magnetocaloric and transport properties of Ni<sub>45</sub>Co<sub>5</sub>Mn<sub>38</sub>Sb<sub>12-x</sub>Ge<sub>x</sub> (x=0-3) has been investigated. The decrease in the exchange interaction brought by Ge substitution can be seen from the reduction in the magnetization of austenite phase and the increase in the martensitic transition temperature. The hysteresis observed for x=0 between the heating and the cooling curves in the magnetization data decreases considerably for x=3. The reduction in the sharpness of the metamagnetic transition is also observed with Ge doping. From the resistivity data it is found that the disorder in the martensitic transitic rans be tuned over a large temperature interval with a small change in Sb/Ge ratio, giving rise to a large MCE and a large MR value over a broad temperature interval. This makes the system as a potential material system for various applications.

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## **QH07**

Effect of doped Mn ions in thermoelectric material of Mg<sub>3</sub>Sb<sub>2</sub>

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Spin caloritronics is about the physics that control of heat, charge and spin currents in structures, devices, and materials. Here we focus on the spin-related thermoelectric properties and study the effect of magnetic ion doping on the thermoelectric material of Mg<sub>3</sub>Sb<sub>2</sub>. We have synthesized single crystals of Mg<sub>3</sub>Sb<sub>2</sub> and Mn-doped Mg<sub>3</sub>Sb<sub>2</sub> by using Bridgeman method. The X-ray diffraction data show that the crystals are single-phased and c-axis oriented. In order to illustrate the thermoelectric performance, we have measured the seebeck coefficient (S), electrical conductivity ( $\sigma$ ), and thermal conductivity ( $\kappa$ ) as a function of temperature. The efficiency of thermoelectric performance is dominated by the figure of merit,  $ZT=S^2 \cdot \sigma \cdot T/\kappa$ . The obtained S,  $\sigma$ , and  $\kappa$  values at 300K are 600  $\mu V/K$ , 16 S/m, and 3.5 W/mK, respectively for Mg<sub>3</sub>Sb<sub>2</sub>. Because of the lower  $\sigma$ , the ZT value is estimated to be small. However, by doping Mn ions into Mg<sub>3</sub>Sb<sub>2</sub>, we find the electrical properties to be improved. We will discuss more about the optimal condition of Mn doping rate in Mg<sub>3</sub>Sb<sub>2</sub> to promote the thermoelectric performance.

**OH08** 

uly 10 (Tue)

# X-ray diffraction study on crystal structure of $Mn_{1.8}Co_{0.2}Sb$ under high magnetic fields

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Measurements of magnetization, electrical resistivity and high field X-ray powder diffraction (XRD) were carried out for polycrystalline  $Mn_{1.8}Co_{0.2}Sb$  in magnetic fields up to 5 T in 4.2-300 K temperature range, in order to investigate the structural properties affected by magnetic fields. In a zero field, the compound shows the first-order magnetic transition from the ferrimagnetic (FRI) to antiferromagnetic (AFM) states at Tt = 145 K with decreasing temperature. By applying magnetic fields of 5 T, Tt decreased to 60 K with thermal hysteresis of 35 K. From the XRD measurements, not only the AFM phase but also the residual FRI phase was confirmed even at 10 K in cooling process under 5 T. Results obtained indicate that the residual FRI phase is kinetically arrested [1].

[1] H. Orihashi et al. Mater. Trans. [Selected Papers from JIM Fall Meeting(2011)] 76 (2011) in press

## **QH09**

## Emergence of Ferromagnetism under pressure in (Mn<sub>1-x</sub>Fe<sub>x</sub>)<sub>3</sub>GaN

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The recent development of "metallic" antiperovskite nitrides Mn<sub>3</sub>AN, where A is a metal or a semiconducting element, has promoted research toward practical use. We found an intimate relationship between the large magneto-volume effect (MVE) in cubic Mn<sub>3</sub>AN and the  $\Gamma^{cg}$ -type antiferromagnetic (AF) ordering. As a tunable thermodynamic parameter, the external pressure may play an important role in this system, because it can affect not only the crystal structure but also the magnetic structure. We report the pressure effect on the stability of magnetic structure of (Mn<sub>1</sub>, Fe<sub>3</sub>, GaN.

In this system, it was reported that a few percent of Fe dopants at the Mn sites alters the AF spin structure of Mn<sub>3</sub>GaN and induces a first-order antiferromag-netic-ferromagnetic phase transition. We have measured the magnetization of (Mn<sub>14</sub>Fe<sub>3</sub>)<sub>5</sub>GaN with x=0,0.05 and 0.07 and Mn<sub>5</sub>Ga(N<sub>12</sub>, yC<sub>3</sub>) with y=0,0.05 and 0.1 under pressure up to 1.3 GPa. A magnetic transition from an AF to an ferromagnetic (FM) state with spontaneous magnetization is observed under pressure. Especially, a spontaneous magnetization in Mn<sup>3</sup>GaN is 0.24 µB per Mn under pressure. The application of pressure increases the magnetization in the FM phase for all these compounds.

# QH10

# First-principles-calculations of the magnetic anisotropy energy of FeCo and FeNi alloys

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The calculations of T.Burkert et at. predict an unusually large magnetic anisotropy energy (MAE) of ~700  $\mu$ eV/atom for the Fe<sub>0.4</sub>Co<sub>0.6</sub> random alloy at a ratio of 'c/a'[1]. This result suggests that other transition metal alloys are also candidate for emerging a high MAE with altering the 'c/a'. In order to seek new design of materials, we research MAE and electronic structures for Fe-Co and Fe-Ni alloys with using the first-principles-calculations. The calculations were carried out using the VASP code. At first, we optimized the length of lattice constant 'a' against the fixed 'c/a' at range of 1.10-1.35. The MAE were evaluated by subtracting the energy with polarized spin direction parallel to [100] from that to [001]. For the Fe-Co alloy, the estimated MAE has maximum at c/a-1.25 and Fe<sub>0.5</sub>Co<sub>0.55</sub>, which is roughly consistent with the previous result of Burkert et al. On the other hand, Fe-Ni system has the high MAE (~300  $\mu$ eV/atom) is found at the Fe<sub>0.75</sub>Ni<sub>0.25</sub>. We will present the origin of the large MAE from the electronic structures in these systems.

[1]T. Burkert et al. Phys. Rev. Lett. 93 027203 (2004)

### OH11

# Tuning the magnetic properties of amorphous FeZr thin by hydrogen implantation

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Amorphous FeZr alloys order ferromagnetically, with a maximum in the ferromagnetic transition temperature near 240K for Zr contents between 10-20 at.% . Interestingly, hydrogen can be introduced into the alloy, enhancing its Curie temperature, magnetic moment and soft magnetic properties. In order to understand the magnetism of the Fe(Zr) binary system, and the e ect of hydrogen on the magnetic interaction, we have investigated in details the magnetic properties of amorphous thin lms of Fe(Zr). Fe1 xZrx lms with compositions x=7.5, 9.5, 11.5, 12.5 and 13.5 at.% were grown by cosputtering on Si substrate in ultra high vacuum (UHV) conditions. The compositions, structural and magnetic properties of the lms were investigated by Rutherford Backscattering spectrometry (RBS), X-ray di raction, and SQUID magnetometery respectively before and after hydrogenation. The hydrogen implanted samples show an increased value of the Curie temperature with increasing hydrogen content. The magnetization curves at low temperature show a continuous increase with increasing eld also at quite high temperature con rming the presence of random anisotropy. This anisotropy is heavily suppressed by hydrogen implantation. Keywords: Amorphous Fe(Zr), Magnetization, H-implantation

# QH12

### Normal and intrinsic anomalous Hall effect in NbFe<sub>2</sub>

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Hall effect studies are a powerful probe of the Fermi surface of metals via the normal contribution whilst the anomalous contribution may give substantial insight into spinorbit interactions. Here, we present Hall effect measurements on selected samples of the band magnet system NbFe<sub>2</sub>. In the dilution series Nb<sub>1.7</sub>Fe<sub>2.47</sub> composition tuning allows to access various magnetic ground states including a complete suppression at a quantum critical point. We observe normal and anomalous contributions in the Hall effect. Consistent analysis of the anomalous contribution, however, is only possible for Fe-rich Nb<sub>0.983</sub>Fe<sub>2.015</sub> featuring a ferromagnetic ground state. Here, a positive normal Hall coefficient is found at all temperatures with a moderate maximum at the spin-density-wave transition. The anomalous Hall effect is covered by an intrinsic (Berry-phase) contribution which is constant below the ordering temperature TC and continuously vanishes above TC. For samples with antiferromagnetic ground state, i.e., stoichiometric NbFe<sub>2</sub> and Nb-rich Nb<sub>1.01</sub>Fe<sub>1.99</sub>, an additional contribution to the Hall resistivity impedes a complete analysis and indicates the need for more sophisticated models of the anomalous Hall effect in itinerant antiferromagnets.

# QH13

# Magnetovolume effect in the itinerant-electron frustrated magnet $Fe_{s}Mo_{s}N$

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The  $\eta$ -carbide-type compound Fe\_3Mo\_3N, including the frustrated stella-quadrangula Fe-sublattice, is a novel itinerant-electron frustrated magnet and exhibits intriguing features such as non-Fermi-liquid behavior without parameter-tuning and sharp metamagnetic transition [1]. We performed thermal expansion and magnetostriction measurements of Fe\_3Mo\_3N up to 15 T. A large volume expansion is found at the metamagnetic transition field and the magnetovolume coupling constant is estimated as about  $4 \times 10^{-3} \ \mu\text{B}^2$ , being the same of magnitude of those of Co Laves phase compounds [2]. It indicates that the metamagnetic transition of Fe\_3Mo\_3N as typical moment-induced-type in nonmagnetic itinerant electron systems. The nonmagnetic ground state of Fe\_3Mo\_3N at zero field is emergent due to the frustration effect built into the stella quadrangula lattice.

 T. Waki et al., J. Phys. Soc. Jpn. 79, 043701 (2010).; T. Waki et al., Europhys. Lett. 94, 37004 (2011). [2] Y. Muraoka et al., J. Phys. Soc. Jpn. 53, 331 (1984).

## QH14

## **NMR studies of magnetic properties in the spinel-type Cu(Cr<sub>1-x</sub>Hf<sub>x</sub>)<sub>2</sub>S<sub>4</sub>** Haruo Niki<sup>1</sup>\*, Morihito Oshiro<sup>1</sup>, Saori Nakamura<sup>1</sup>, Ayaka Uechi<sup>1</sup>, Mamoru Yogi<sup>1</sup>, Shuji Ebisu<sup>2</sup> and Shoichi Nagata<sup>2</sup>

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A chalcogenide tiospinel CuCr<sub>2</sub>S<sub>4</sub> is a metallic ferromagnet with Tc = 380 K. The increase of x for the mixed spinel Cu(Cr<sub>1,x</sub>Hf<sub>2</sub>)<sub>2</sub>S<sub>4</sub> shows a successive change from a ferromagnetic, a re-entrant spin-glass, a spin-glass, and finally to a paramagnetic state. A step-like anomaly in magnetization is also found around 160 K in the range of  $0.50 \le x \le 0.70$  due to a spin crossover phenomenon [1]. In order to investigate the magnetic properties of Hf rich side in Cu(Cr<sub>1,4</sub>Hf<sub>2</sub>)<sub>2</sub>S<sub>4</sub> from a microscopic point of view, 63Cu NMR have been carried out between 4.2 and 300 K. Temperature dependences of resonance line-width, Knight shift and spin-lattice relaxation time (T1) have been measured. Knight shift in the range of  $0.50 \le x \le 0.70$  gradually shifts to negative side below 130 K, but rapidly shifts to negative side below 130 K due to the core polarization of d conduction electrons. Temperature dependence of 1/T1T in the range of  $0.50 \le x \le 0.70$  obeys the Korringa relation (T1T = constant) down to 130 K, as expected in a Pauli-paramagnetic metallic state. However, the values of 1/T1T =constant abruptly change between 130 and 90 K, and keep constant below 90 K.

[1] F. Kariya, K. Ebina, K. Hasegawa, K. Koshimizu, B. Wuritunasitu, K. Hondou, S. Ebisu and S. Nagata, J. Solid State Chem. 182 (2009) 2018.

# QH15

### Relationship between microstructure and magnetic properties of nanoscale magnetic particles formed in a Cu-Ni-Co alloy

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The discovery of Giant Magnetic Resistance (GMR) has invoked strong interest to relationship between microstructure and magnetic properties of nano-scale particles. The magnetic properties of nano-scale particles are strongly dependent on their structures, particle sizes, shapes, and interparticle distances particularly in granular magnets. The present study aimed to investigate the relationship between the microstructure and the magnetic properties of granular Cu-Ni-Co alloys, using transmission electron microscope (TEM), a superconducting quantum interference device (SQUID) magnetometer and thermo gravimetric-differential thermal analysis (TG-DTA). The present TEM observations revealed a specific feature that spherical ferromagnetic precipitates are transformed into cubic ones, and subsequently two or more cubic precipitates are aligned along the <100> orientation of a Cu matrix, which feature was significantly different from previous recognition to conventional precipitation. SQUID measurements shows that the magnetic property of the specimen changes from superparamagnetism to ferromagnetic in character with increasing isothermal annealing period from M-T curves. The present TG/DTA curves showed that a ferromagnetic to paramagnetic transition occurred at low temperatures (~950K) but paramagnetic to ferromagnetic transition at high temperatures(600~700K). This suggests that the coalescence of magnetic atoms and accordingly the Curie point are sensitive to annealing temperature.

# QH16

### Magnetic study of mechanically deformed FeAlSi alloys

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The mechanical deformation of Fe-Al alloys produces order-disorder transition that induces an enlargement of the magnetism of the alloys. It has been demonstrated that this enlargement is related to atomic disorder and cell volume increase [1]. The behaviour of the saturation magnetization with milling and Si content in Fe<sub>27</sub>Al<sub>32</sub>, s<sub>1</sub>Si, and Fe<sub>37</sub>Al<sub>304</sub>, Si, alloys is similar to the behaviour shown by the mean magnetic hyperfine field. Their values increase with disordering in alloys with low Si content, but decrease in alloys with higher Si content. The behaviour observed indicates clearly that the introduction of Si in the binary Fe<sub>57</sub>Al<sub>325</sub> and Fe<sub>37</sub>Al<sub>345</sub> and Fe<sub>37</sub>Al<sub>345</sub> and Fe<sub>37</sub>Al<sub>355</sub> and Fe<sub>37</sub>Al<sub>355</sub>

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### QH17

# The propagation of electromagnetic waves in magnetic with ferromagnetic spiral

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Recently, spiral magnetic materials have attracted researchers' attention for their unusual physical properties. However, the spectrum and dynamic properties of magnets in a phase <</ferromagnetic spiral>> are not studied enough. In the present work the spectrum of the coupled spin and electromagnetic waves in spiral magnetic structure of type <</ferromagnetic spiral>> is investigated. Also the reflection of electromagnetic waves from a layer of magnetic having the ferromagnetic spiral and the Faradey effect are considered. We use approach L >> a, where L - the spiral period, a - the lattice constant. Studies showed that the spectrum of the coupled waves has band structure. The band gap depends on the angle of the ferromagnetic spiral, and hence on the external magnetic field. The frequency dependence of the reflection coefficient of electromagnetic waves from the plate of the magnetic with a ferromagnetic spiral at different angles of the helix has been calculated. As the angle increases the opacity region broadens and shifts to higher frequencies. Thickening of the plate also leads to broadening of the opacity region. The Faraday effect has been observed for different frequencies. The angle of plate-of-polarization rotation increases near the band gap. The Verdet constant has been calculated

## QH18

### Magnetism of sigma-phase Fe-Mo and Fe-Re systems

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Sigma phase exists in five binary systems of the Fe-X type viz. for X = Cr, V, Mo, Re and Tc. It is characterized by a lack of chemical order on all five sublattices of the tetrahedral unit cell. Until now, the magnetic properties have been known only for sigma-FeC and sigma-FeV systems. The newly discovered low temperature magnetic properties of sigma-FeMo and sigma-FeRe phases will be presented and analyzed in detail in the whole concentration range in which this phase occurs. In particular, in both cases, the observed increase in an average magnetic moment and magnetic ordering temperature with increasing of Fe content is evident. Differences in Zero-Field Cooling and Field Cooling magnetization measurements suggest the magnetic ordering of spin-glass type or a large magnetic anisotropy in the measured samples. Electronic structure calculations using the KKR method were carried out for more than a dozen semi-ordered unit cells. The results of these calculations show a linear dependence of magnetic moments and hyperfine fields of Fe atoms at different sites versus the number of Fe atoms being their nearest neighbors. They also suggest a noncollinear magnetic structure in some sublattices.

# QH19

# Second order magnetization effects in bcc Fe

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Second order effect in magnetization are becoming an important way of characterization of the magnetic samples. One of the examples of second order effect in magnetization are anisotropic magnetoresistance (AMR), Voigt effect, quadratic magneto-optical Kerr effect (QMOKE) or X-ray magnetic linear dichroism (XMLD). All those effects are originating from conductivity (or permittivity) tensor, which elements are varying with direction of the sample magnetization. This dependence is determined by symmetry arguments describing break of the symmetry by magnetization in the crystal. From first principles we have calculated the complete angular dependence of all elements of the permittivity tensor with respect to the change of magnetization axis in a wide spectral range (from dc to ultraviolet) for a body-centered cubic Fe. The correctness of calculations are verified by careful comparison with expected dependences provided by symmetry arguments. The calculations are based on density functional theorem utilizing one of the most accurate full potnetial augmented plane-wave code. The effect of the spin-orbit interaction for valence and semicore states is included as well.

# QH20

Magnetic structure of nearly equiatomic MnRh alloy Yuki Matsuoka<sup>1</sup> and Aya Takasaki<sup>2</sup>

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MnRh shows various kind of magnetic state. The magnetic structure and the transition temperature of MnRh are expected to be very sensitive to the atomic composition. At the equiatomic composition, this alloy shows antiferromagnetic - paramagnetic transition. By increasing Mn ratio, this alloy shows complicated magnetic property, ferromagnetic - weak ferromagnetic transition.[1] Yamada et.al. reported that antiferromagnetic structure is most stable in this alloy.[2] To determine the magnetic structure, we performed neutron diffraction measurements and observed some magnetic scattering peaks. The intensities of magnetic scattering peaks were almost same even just below the phase transition temperature. On the other hand, magnetization curve of same component alloy suggests magnetic phase transition. We attempt to determine the magnetic structure and to make magnetic phase diagram of nearly equiatomic MnRh.

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## QH21

### Magnetic and transport properties of Pd2Mn1+xIn1-x Heusler alloys

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Recently, it was reported that Ni-Mn-In and Ni-Mn-Sn Heusler alloys having a ferromangetically ordered L21 type structure show a magnetic-field-induced reverse martensitic transformation associated with a metamagnetic transition [1]. The interplay between the martensitic transformation and the magnetic transition gives rise to a magnetic-field-induced shape recovery effect and a magnetoresistance effect. Pd<sub>2</sub>MnIn has an antiferromagnetically ordered L21 type Heusler structure [2], which is different from the magnetic ordered state of Ni-MnIn. In this work in order to investigate the magnetic and transport properties of the antiferromagnetic Heusler alloys, we performed magnetization and electrical resistivity measurements for Pd2Mn1+xIn1-x-In Pd<sub>2</sub>Mn<sub>14x</sub>In<sub>1</sub>, with x > 0.2, a sudden increase of the resistivity was observed. The temperature at which the anomaly of the resistivity appears shifts to the higher temperature with increasing Mn content, and no anomaly was detected for the Mn content with x = 0.55. The results of powder X-ray diffraction experiment indicate that a martensitic phase appears for the Mn content with x = 0.55 at room temperature Therefore, it is considered that the anomalous behavior of the resistivity results from the martensitic transformation

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# QH22

# Structural and magnetic properties of (Co, Fe, Ni), (Fe, Mg, Zn) and (Fe, Mn, Zn) alloys deposited onto Al<sub>2</sub>O<sub>3</sub> and SiO<sub>2</sub>

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Structural and magnetic properties of (Co, Fe, Ni), (Fe, Mg, Zn) and (Fe, Mn, Zn) alloys deposited onto Al<sub>2</sub>O<sub>3</sub> and SiO<sub>2</sub> supports are reported. The samples were synthesised by ball milling and characterised using magnetic measurements (MM), Mossbauer spectroscopy (MS), XRD, SEM, TPR and (BET). Granular and irregular shaped particles were formed. MM, MS measurements and XRD data revealed the formation of peculiar solid solutions some of them having spinel structures. In all the samples small crystallites in a superparamagnetic (SPM) state at room temperature were found to coexist with larger ones, their fractions depending on the energy imparted during the ball milling process. That is, crystallites were prevalent at a relatively low ball-to-powder ratio. The ball milling process resulted also in low surface areas and collapse of pores. Further, the existence of strong metal support interactions (SMSI) was inferred in all the samples; this rendered reduction difficult and/or led to incomplete reduction. Fittings of the Curie-Weiss model to the experimental temperature dependences of AC magnetic susceptibility were used to determine the effective magnetic moments associated with magnetic cations and the nature of their interactions.

### QH23

### Structure disorder effect of L1<sub>0</sub> FePt within supercell method Kazuhiko Uebayashi

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L1<sub>0</sub> FePt alloy is famous for a good candidate of high density magnetic recording media. In my previous work[1], I have performed energy-band calculations for the L1<sub>0</sub> bulk alloy FePt by using the linear muffin-tin orbital method with atomic sphere approximations. This previous calculated study shows that the calculated lattice constants and magnetic moments are good agreement with observed ones. The experimental stable-magnetic state is ferro magnetic, but in the calculation one of the anti-ferro magnetic state is most stable, and the second is ferro magnetic above 1 mRy/ f.u. The difference between the experimental and calculated result for L1<sub>0</sub>FePt is known by another work[2]. To improve the difference, I carry out structure-disorder band calculation for FePt with supercell method. I assumed that the supercell is a 8 times large as a formula unit. In the supercell with only one substitution Ferro magnetic state is about  $10^3$ mRy. The result indicate that the structure disorder of L1<sub>0</sub> FePt in the supercell makes ferro magnetic stable with observed lattice constants.

[1] K. Uebayashi, JOURNAL OF APPLIED PHYSICS 107, 09A733 (2010) [2] G. Brown, PHYSICAL REVIEW B 68, 052405 (2003)

## QH24

## Itinerant antiferromagnetism in high-quality single crystal CrB<sub>2</sub>

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Transition metal and rare earth diborides with the hexagonal AlB<sub>2</sub> structure form a series of compounds with a remarkably wide range of different electronic ground states. We report the properties of isotopically enriched high-purity single crystals of CrB<sub>2</sub> grown by optical float-zoning. The temperature dependence of the electric resistivity and specific heat provide unambiguous evidence of spin density wave type order with a slightly anisotropic charge gap  $\Delta(001) = 230$  K and  $\Delta(100) = 212$  K, respectively. We also find that the Neel temperature of 88.5 K remains unchanged under magnetic fields up to 14 T, the highest field studied. A detailed neutron scattering study, made possible through the use of 99% isotopically enriched 11B, provides for the first time unambiguous evidence of antiferromagnetic order, and identifies CrB<sub>2</sub> as a textbook example of internant antiferromagnetism. In comparison to the properties of other high-quality single crystal transition metal and rare earth diborides we have prepared, our studies of CrB<sub>2</sub> shed an unexpected new light on the properties of this exciting series of compounds.

# QH25

### Magnetic orderings in Fe-intercalated TiX<sub>2</sub> (X=S, Se) Seung Ill Hyun and J.h. Shim\* Chemistry, POSTECH, Korea

Titanium dichalcogenides shares the same layered structure, except TiO<sub>2</sub>. Their physical properties, however, are different from each other. TiSe<sub>2</sub>, for example, is the only compound that bears an unconventional charge-density wave (CDW) transition under 200K, whereas the others show no distinct structural distortion.[1] Furthermore, fron-intercalated TiS<sub>2</sub> and TiSe<sub>2</sub> exhibit different types of magnetism. Fe<sub>0.5</sub>TiS<sub>2</sub> displays a ferromagnetic ordering under 80K, Fe<sub>0.5</sub>TiSe<sub>2</sub> shows an antiferromagnetic ordering under 135K.[2,3] In order to describe the magnetic interactions between Fe atoms, we apply the first principle calculation on Fe<sub>0.5</sub>TiX<sub>2</sub> (X=S, Se) and explain the experimental observations correctly. We also discuss the origin of different magnetic behaviors based on the structural & elemental effect.

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### QH26

### Non-Stoner Itinerant Ferromagnetism of a Transition Metal Monomer Hanhim Kang, Geunsik Lee, P. Dua, Ji-hoon Shim and Kwang S Kim\* Department of Chemistry, POSTECH, Korea

We investigated the one-dimensional monomer chain of 3d transition metals (Fe, Co, and Ni ) by ab-initio band structure calculation with quasiparticle GW method [1]. We also studied the role of pair hopping interaction on the magnetic properties using the extended Hubbard model [2][3]. GW calculation of transition metal monomer shows the spin-dependent band width changes as well as spin-dependent transport properties in ferromagnetic phase. Such behavior is well described by the extended Hubbard model, where the pair hopping enhances the metallic ferromagnetism in one-dimension. We suggest that the electron correlation effect such as pair hopping usually has important role in the metallic ferromagnetism of low dimensional system.

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## QH27

# Magnetic structure variation with structural phase transition in $\mbox{SrMnSb}_2$

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RMnSb<sub>2</sub> system (R = Alkali Earth Metal) is analyzed by first principle DFT calculation using full potential linearized augmented plane wave method while PBE-GGA is used to account for the Exchange-Correlation potential. SrMnSb<sub>2</sub> has zig-zag chains of Sb ions while BaMnSb<sub>2</sub> has square planes of it. The application of pressure on SrMnSb<sub>2</sub> leads to a phase transformation from chain like semi-metal phase to BaMnSb<sub>2</sub> like Square planner structure showing metallic properties. The density of states for AFM SrMnSb<sub>2</sub> changes with the structural phase change indicating a magnetic transition. While for a similar chain to square transition in BaMnSb2 accompanying an Insulator to Metal transition, no such variation in magnetic structure is observed.

# QH28

### Magnetoresistance effect of Heusler-type Fe<sub>2</sub>VAl single crystal

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The Heusler-type intermetallic compound Fe<sub>2</sub>VAl has attracted strong attention because of a semiconducting behavior of  $\rho$  and a large electronic specific heat coefficient  $\gamma$ , 14 mJ/(K2mol). The band calculation revealed that the Fe<sub>2</sub>VAl was a semimetal and nonmagnetic with a pseudogap around the Fermi level [2]. Endo et al. discovered the giant magnetoresistance (GMR) effect in the Heusler type alloys Fe<sub>2-x</sub>V<sub>1-x</sub>Al poly crystal of about 30 % for Fe<sub>21</sub>V<sub>0.9</sub>Al at 4.2 K in the field of 55 kOe [3]. Recently we have succeeded in fabricating single crystals of Fe<sub>21</sub>V<sub>0.9</sub>Al using Czochralski method and investigated their electric and magnetic properties. The temperature dependence of GMR is negative below the Curie temperature (TC). The maximum value of GMR is about 35 % around the TC. The value of GMR in Fe<sub>21</sub>V<sub>0.9</sub>Al single crystal is similar in that of poly crystal.

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## QH29

# Electronic structures and magnetic properties of antiferromagnetic BaFe<sub>2</sub>As<sub>2</sub>

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In low dimensional metals the electrons can be organized into regular patterns, which is known to be the charge-density wave(CDW) or spin-density wave(SDW). The parent compound of iron-based superconductors, BaFe<sub>2</sub>As<sub>2</sub>, undergoes both structural and magnetic phase transitions. The crystal structure changes from tetragonal to orthorhombic, and the magnetic ground state changes from paramagnetic to antiferromagnetic at ~140 K and neutron scattering measurements show the magnetic moment of 0.87µB per Fe in antiferromagnetic state. SDW anomaly at 140 K, which is believed to be an important prerequisite for high-Tc superconductivity in iron-based superconductors, was also observed. We performed the electronic-structure calculation for the BaFe<sub>2</sub>As<sub>2</sub> compound in various antiferromagnetic states, and the electronic structures and Fermi surfaces are carefully investigated. While there have been considerable debates as to whether the antiferromagnetism is due to the Fermi-surface nesting or the second-neighbor superexchange, via comparing our results with experimental and hypothetical cases, we can deduce the correlation between the SDW

## QI01

# Development of magnetic order in the TbNi(Al,In) series and magnetocrystalline anisotropy in TbTX compounds

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We present recent study of magnetocrystalline anisotropy in TbNi(Al,In) series, crystallizing in hexagonal ZrNiAl structure. The properties of this series were investigated by magnetic measurements, low-temperature x-ray diffraction in magnetic field and mainly by powder neutron diffraction. Parent compounds of our series were previously studied. TbNiAl orders antiferromagnetically with propagation (1/2,0,1/2) and at lower temperatures undegoes another magnetic phase transition when one third of magnetic moments is reduced to almost zero. TbNiIn shows complex magnetic ordering with propagation (0,0,0) and other propagation (1/2,0,1/2) occurs at lower temperatures. Our data reveal suppression of the antiferromagnetic structure of TbNiAl with substitution of 10% Al by In. Instead, dominant ferromagnetic corder appears, while the direction of Tb magnetic moments along the c-axis is retained. Additional weak antiferromagnetic component is arranged within the basal plane and described by ket (1/2,0,1/2) appear at low temperatures. This antiferromagnetic component is presented in the whole rest of series. Dominant component described by (0,0,0) propagation the undergoes change from c-axis ferromagnetic order to a more complex basal-plane structure around 40% of In. Anisotopy change in this series is related to strong change of the c/a ratio of lattice parameters. This behaviour can be generalized for other hexagonal Tb-based compounds.

# **OI02**

### Magnetic properties of Ho<sub>1-x</sub>Lu<sub>x</sub>B<sub>12</sub> solid solutions

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Rare earth dodecaborides (REB<sub>12</sub>) have attracted considerable attention in recent years, above all due to the wide variety of their physical properties. In our contribution we present the magnetic properties of the geometrically frustrated antiferromagnet  $HoB_{12}$  (with ordering temperature T\_N = 7.4 K) modified by substitution of magnetic Ho atoms through nonmagnetic Lu ones. In this case, in the so-called solid solution Ho<sub>1-x</sub>Lu<sub>x</sub>B<sub>12</sub>, both chemical pressure (as Lu<sup>3+</sup> ions and Ho<sup>3+</sup> ions have different radii) and magnetic dilution take place with the increasing content of Lu that change the properties of the system. The received results show strong indication for the existence of a quantum critical point close to x = 0.9. This critical point separates the region of magnetic order (starting with  $HoB_{12}$  for x = 0) and the nonmagnetic region (ending with superconducting  $LuB_{12}$  for x = 1). The work was supported by the Slovakia Scientific Agency VEGA (0148-10), the Slovakia Research and Development Agency APVV, and by the Center of Excellence of the Slovakia Academy of Sciences.

**OI03** 

High-field magnetization of Tm<sub>2</sub>Fe<sub>17</sub> and its deuteride

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A magnetization study of a Tm<sub>2</sub>Fe<sub>17</sub> single crystal and aligned powder of the deuteride Tm<sub>2</sub>Fe<sub>17</sub>D<sub>3</sub> has been carried out in steady (14 T) and pulsed (60 T) magnetic fields at 1.5-300 K. Tm<sub>2</sub>Fe<sub>17</sub> is a ferrimagnet with T<sub>c</sub>=295 K and a spontaneous moment of 21 µ<sub>B</sub>/f.u. Due to competing magnetic anisotropy of the Tm and Fe sublattices, a spin-reorientation occurs at 70-100 K, from the easy-axis to the easy-plane type via a cone of easy axes. Two metamagnetic transitions, at 41 and 53.5 T, are observed at 2 K in fields applied along the c axis. At 60 T the magnetization reaches 43 µ<sub>B</sub>/f.u., i.e. the sublattices are not yet parallel ( $M_{E_e}+M_{T_m}=49 \mu_B$ ). In the basal plane a wide S-shape magnetization curve is observed. The high-field behavior of Tm2Fe17 is compared with that of two other easyaxis R2T17 compounds, Tm2Co17 and Er2Co17. Hydrogenation (deuteration) stabilizes easyplane anisotropy in  $Tm_2Fe_{17}D_3$  over whole ferrimagnetic range (up to  $T_c$ =465 K). Despite different types of magnetic anisotropy, the high-field behavior of the deuteride and the parent compound is qualitatively similar. The transitions along the c axis occur in the same fields, being less pronounced in the deuteride.

## **OI04**

## RKKY interaction and magnetic properties in (Y-Gd)Ni compounds

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CeNi is well known as one of the most famous heavy fermion compounds and shows a peculiar behavior as an enhanced Pauli-paramagnet. That is, both Ce and Ni seem to be non-magnetic in CeNi. On the other hand, GdNi, which has the same structure as that of CeNi, is a ferri-magnet and not only Gd but also Ni are magnetic. It follows that the result that Ni is non-magnetic in CeNi becomes not self-evident and the mechanism of disappearance of Ni magnetic moment is important and interesting problem. Based on such scenario, we have been investigating a (Ce-Gd)Ni compound system. One of the results obtained is a linear relationship in M(T) in Gd-poor content (Gd<20 at%) and it is found that this result reflects the collapse-like decrease of Gd-Gd exchange interaction. In order to examine this collapse-like decrease of RKKY interaction, we have investigated the (Y-Gd)Ni compound sysytem where Ce is replaced by Y. In this system the exchange interaction between Gd and Gd does not collapse-likely decrese. and seems to revive even in Y=15 at%. Furthermore, YNi (Gd=0) seems to be itinerant. however, the substitution of Y by Gd causes localization and Curie-Weise law becomes valid

## **OI05**

Magntization curves in high magnetic fields of TbZn,

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Rare earth Zinc compounds RZn2 (R=Gd, Tb, Dy, Ho, Er) crystallize in CeCu2-type structure. In the case of GdZn<sub>2</sub>, its magnetism is ferromagnetic, in the other cases, antiferromagnetic. Debray et al. [1] was reported that TbZn2 is the antiferromagnet with two antiferromagnetic phases which are the high temperature phase AF(a) and the low one AF(b). Two transition points TN(a) = 75 K and TN(b) = 60 K are corresponding to those phases. The magnetic structure of AF(b) is a collinear type and AF(a) is a sinusoidal transverse wave type. The magnetic measurements were performed with a single crystalline sample by Kitai [2]. Two transition points TN and Tt were observed corresponding to TN(a) and TN(b). We performed the measurements of the magnetization curves in high magnetic fields with the single crystalline sample of TbZn2. The metamagnetic transitions were observed in the directions of all axes. Especially in the direction of c-axis, the metamagnetic behaviors were observed above TN. In this paper, we will discuss the origin of the metamagnetic transitions for TbZn<sub>2</sub>.

[1] D. Debray et al.: J. Chem. Phys., 56 (1972) 4325. [2] T. Kitai: J. Phys. Soc. Jpn., 64 (1995)

### **OI06**

### Deuteration induced ferromagnetic metallic properties in R<sub>7</sub>Rh<sub>3</sub>Dx (R = Tb, Dv)

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Heavy rare earth compounds  $R_7Rh_3$  (R = Gd to Er) show the antiferromagnetic character at low temperature and semimetallic behavior in which the temperature coefficient of electrical resistivity is negative in high temperature region [1]. We studied the magnetic properties of Tb<sub>7</sub>Rh<sub>3</sub> and Dy<sub>7</sub>Rh<sub>3</sub> hydrides and reported the hydrogenation induced ferromagnetic states in these compounds [2]. At this time, we have investigated the magnetic and electrical properties of the R-Rh3D, deuterides and found the ferromagnetic metallic state in these compounds. Deuteration of Tb7Rh3 and Dy<sub>7</sub>Rh<sub>3</sub> were carried out using the high pressure apparatus in the pressure range up to 1 GPa at 100 °C for 24 hours and Tb7Rh3D27 and Dy7Rh3D28 were obtained. Dy7Rh3 has an antiferromagnetic phase below TN = 59 K and a semimetallic properties. On the other hand,  $Dy_7Rh_3D_{28}$  shows a ferromagnetic state below TC = 11.5 K and a metallic property below 300 K. In Tb<sub>7</sub>Rh<sub>3</sub>D<sub>27</sub>, almost the same changes in magnetic and electrical properties were observed.

[1] T. Tsutaoka, Y. Nakamori, T. Tokunaga, T. Itoh, J. Phys. Soc. Jpn., 70 (2001) 199. [2] R. Sato, T. Tsutaoka, S. M. Filipel, J. Alloys Comp., 446-447 (2007) 610.

### **OI07**

# Magnetic properties of Nd7Pd3 single crystal

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Rare earth intermetallic compound Nd<sub>2</sub>Pd<sub>3</sub> shows relatively large magnetocaloric property and magnetoresistance [1]. This time, we have succeeded to make single crystals of Nd<sub>2</sub>Pd<sub>3</sub> and investigated the magnetic properties by measuring magnetization and magnetic susceptibility. Single crystals were grown by the Czochralski method from single-phase polycrystalline samples, which were made by arc-melting the constituting 99.9%-pure Nd and 99.95%-pure Pd elements, under highpurity argon atmosphere. The crystal orientation was determined by the back-reflection Laue method. Magnetization curve at 2 K indicates ferromagnetic property for both a- and c-axes: the easy magnetization axis locates in the c-plane. Metamagnetic transitions indicating an antiferromagnetic state were observed above TC = 34 K in the magnetization curves, which were reported so far. A first order magnetic transition from ferromagnetic to antiferromagnetic was observed at TC in the temperature dependence of magnetic susceptibility in both a- and c-axes. Magnetic susceptibility curves for both axes also indicate a small cusp at around TN = 40 K indicating the antiferromagnetic to paramagnetic transition.

[1] N. K. Singh, P. Kumar, Z. Mao, D. Paudyal, V. Neu, K. G. Suresh, V. K. Pecharsky, K. A. Gschneidner Jr., J. Phys. Cond. Mat., 21 (2009) 456004.

# **OI08**

### Giant magnetoresistance and field-induced phase transitions in Tb<sub>2</sub>Rh<sub>3</sub> single crystal

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The rare earth intermetallic compound Tb7Rh3 crystallizes in the Th7Fe3 type hexagonal structure with the space group P63mc in which Tb occupies three non-equivalent sites.  $Tb_7Rh_3$  enters an antiferromagnetic (AF) state below TN = 91 K and shows further magnetic transformations with decreasing temperature below ~ 73 K [1]. Unusual negative temperature coefficient of the resistivity in a wide temperature range above TN [1] and the large magnetoresistance even in paramagnetic state [2] were observed for this compound. In this study, we have performed high field magnetization (up to 350 kOe) and magnetoresistance (up to 120 kOe) measurements on a Tb-Rh- single crystal. The metamagnetic transition along the b-axis is observed to be accompanied by the giant magnetoresistance effect (~ 65 %) in the field interval (70 -110) kOe. The magnetization measurements along the c-axis have revealed multiple phase transitions at fields above 200 kOe. The magnetization curves measured in the paramagnetic state along the c-axis show an inflection point up to 140 K, which may be indicative of the persistence of antiferromagnetic short-range magnetic order in the c-plane.

[1] Y. Nakamori et al. Physica B, 304 (2001) 1. [2] K. Sengupta et al. Solid State Comm., 39 (2006)

### **OI09**

# Magnetic properties of a GdFe<sub>5</sub>Al<sub>7</sub> single crystal

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A magnetization study of a GdFe<sub>5</sub>Al<sub>7</sub> single crystal has been carried out in static magnetic fields up to 14 T at temperatures between 2 and 300 K. The compound is found to order ferrimagnetically at T<sub>c</sub> = 266 K. The spontaneous magnetic moment is 0.56  $\mu_B$  at 2 K. There is no compensation point in GdFe\_5Al\_7, as distinct from isomorphic compounds DyFe<sub>5</sub>Al<sub>7</sub> and HoFe<sub>5</sub>Al<sub>7</sub>. GdFe<sub>5</sub>Al<sub>7</sub> is endowed with an appreciable magnetic anisotropy, the 4-fold axis being a hard magnetization direction. The easy direction is [110] and the anisotropy within the basal plane is also rather strong. The latter is quite remarkable for a system devoid of significant orbital moments. The possible cause of the in-plane anisotropy in GdFe<sub>5</sub>Al<sub>7</sub> could be the unusually large crystal field parameter B<sub>44</sub>, as observed in other compounds with the ThMn<sub>12</sub>-type structure, e.g., in DyFe<sub>11</sub>Ti. The crystal field effect could be mediated by a small admixture of the  $4f^8$  configuration to the ground state  $4f^7$ . The exchange field on the Gd sublattice is determined from the high-field slope of the basal-plane magnetization curves to equal 62 T.

### **OI10**

### Study of multipole ordering in CePd<sub>3</sub>S<sub>4</sub> by resonant X-ray diffraction with full polarization analysis

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In rare-earth palladium bronzes RPd<sub>3</sub>S<sub>4</sub> (R = La-Yb) compounds with a cubic structure, it is expected that the multipole interactions dominate the low temperature properties because the degeneracy of the orbital degree of freedom remains in the crystalline electric field ground states of  $R^{3+}$  in RPd<sub>3</sub>S<sub>4</sub> except for R = Tm.[1] CePd<sub>3</sub>S<sub>4</sub> orders ferromagnetically below the transition temperature TC = 6.3 K.[2] Recently we revealed that the O20-type quadrupole moments order below TC, and infered that the magnetic structure is a ferrimagnetic structure by resonant x-ray diffraction (RXD) experiments.[3] Since the symmetry of the magnetic and O20-type quadrupole moments differs from each other, simultaneous order of them below TC suggests the involvement of the octupole interaction. We performed full polarization analysis of RXD experiments in magnetic fields to study the contribution of an octupole interaction. The experiments were performed at the Ce-L<sub>3</sub> absorption edge using a single crystal at BL22XU in SPring-8. In our presentation, we will discuss the contribution of multipole interactions for the low temperature properties in detail.

[1] K. abe et al.: Phys. Rev. Lett. 76 (1999) 5366. [2] E. Matsuoka et al.: J. Phys. Soc. Jpn. 77 (2008) 114706. [3] S. Michimura et al., submitted to J. Phys. Soc. Jpn

# **OI11**

### Effect of temperature and magnetic field history on magnetization behavior of NdVO<sub>3</sub> polycrystalline

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The magnetic properties of NdVO<sub>3</sub> polycrystalline have been studied in the temperature range from 300 to 1.8K and in magnetic fields up to 16 T. It is found that in an applied field less than 1 T remarkable magnetization reversals occurred in field cooling mode: a diamagnetic behavior was observed below a compensation temperature Tcomp. Most interestingly, the diamagnetic behavior was dependent to both the temperature and applied magnetic field, where the temperature window is 50-100K, and magnetic field below 1T. Beyond the temperature window and above the magnetic field, the diamagnetic behavior disappeared. That is the magnetization always switches to the opposite direction if the sample is cooled through the temperature window in a field less than 1 T. However, the temperature dependence of the resistance has an almost independent to the applied magnetic field in the range from 0 to 16T. A possible mechanismbased on the competition of the single-ion magnetic anisotropy and the antisymmetric Dzyaloshinsky-Moriya interaction accompanied by a change of orbital ordering is discussed.

# **OI12**

### Phase diagram and transport properties of Y<sub>1-x</sub>NdxCo<sub>2</sub> pseudo-binary allovs

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Electrical resistivity p and thermopower S of pseudo-binary Y1, Nd, Co2 alloys have been measured at temperatures from 2 K to 300 K, under pressures up to 3.5 GPa. NdCo2 is a ferromagnet with Curie temperature TC of 99 K. In ferromagnetic phase it undergoes 4f spin-reorientation transition at TR = 40 K. At this transition the easy direction of 4f magnetization is changed from [100] above TR to [110] below TR. The system phase diagram TC(x,P) and TR(x,P) is inferred from the temperature dependencies of the electrical resistivity and thermopower. The Curie temperature of the alloys, TC decreases with decreasing Nd concentration x and vanishes around xc=0.3. Above xc the Curie temperature decreases with pressure. The spin-reorientation temperature is weakly dependent on composition and on pressure until it merges with TC. There is a large region of the phase diagram, partly overlapping with the ferromagnetic phase, where cobalt 3d electron system is non-uniformly magnetized due to spatial fluctuations of the 4f-3d exchange field, related to random distribution of Nd 4f magnetic moments over R-sublattice. In this region, spanning the range 0 < x < xm, large static 3d magnetic fluctuations govern the electronic transport.

## **OI13**

### Effect of partial magnetic order on resistivity and thermopower of Ho(Co<sub>1-x</sub>Al<sub>x</sub>)<sub>2</sub> alloys

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The electrical resistivity  $\rho$  and thermopower S of Ho(Co<sub>1-x</sub>Al<sub>x</sub>)<sub>2</sub> alloys (x = 0, 0.15 and 0.2) were measured from 2 K to 300 K in magnetic fields up to 10 T and under pressures up to 3 GPa. While  $\rho(T)$  and S(T) of the pure HoCo<sub>2</sub> reveal abrupt changes at the magnetic ordering temperature TC, indicating the first order transition, the temperature variation of both properties across the ordering temperature in Ho(Co1.xAlx)2 alloys clearly show that the type of the magnetic transition changes under substitution of Al for Co from the first order to the second order.  $\rho(T)$  of the Al substituted samples (x = 0.15 and 0.2) has a very unusual variation below TC. Usually, the resistivity of an itinerant ferromagnet decreases below TC with decreasing temperature at higher rate than above TC, its temperature derivative having positive "lambda" anomaly at the ordering point. In contrast, p(T) of Ho(Co1-xAlx)2 alloys (x=0.15 and x=0.2) increases below TC. Moreover, the magnetoresistance of these allovs is positive around TC, this is also in a sharp contrast to usually observed negative magnetoresistance.

# QI14

### Breakdown of Hund's third rule for intrinsic magnetic moments

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Hund's rules were originally developed in an atomistic picture, and it is therefore surprising that they also hold, almost universally, for molecules and solids. Here we report the first experimental example of a breakdown of Hund's third rule in an intrinsic magnetic moment: the Co 3d moment of the unorthodox ferrimagnet ErCo<sub>2</sub>. The application of magneto-optical sum rules to X-ray magnetic circular dichroic spectra allows for the determination of the Co orbital and spin magnetic moments both in the ferrimagnetic and paramagnetic phases of ErCo<sub>2</sub>. In agreement with the predictions of Hund's third rule, Co orbital and spin moments are parallel throughout the overall temperature range, except in a narrow temperature window within the paramagnetic phase, in which they are overall coupled antiparallel. A qualitative consideration of the balance between the interatomic and intra-atomic interactions acting on the cobalt atoms explains the phenomenon. This situation also leads to the observation of other exotic forms of magnetism in solids, such as a "pure" spin or a "pure" orbital moment on the same compound at selected temperature and fields.

OI15

Magnetic properties of Y substituted TbB<sub>4</sub>

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We synthesized the single crystals of  $Tb_{1,x}Y_xB_4$  (x=0.0, 0.1, 0.2, 0.3, 0.35, 0.4, 0.5, and 0.6)using metallic flux method. The lattice parameters of these crystals are confirmed by powder x-ray diffraction using Fullprof program. The lattice parameters are decreased with increasing substituted Y ratio, according to Vegard's law. The measurements of temperature-dependent magnetization M(T) is carried out for the single crystals of  $Tb_{1,x}Y_xB_4$  (x=0.0, 0.1, 0.2, 0.3, 0.35, 0.4, 0.5, and 0.6) in 1 Tesla. The compounds reveal antiferromagnetic ordering and its Neel temperatures are decreased with increasing Y ratio. We confirm the spin direction of antiferromagnetic contribution in  $Tb_{0.65}Y_{0.35}B_4$  using single crystal neutron scattering. The direction of Tb spin magnetic moment is tilted from a axis.

## QI16

# High pressure effects on antiferroquadrupolar orders in $RB_2C_2$ (R = Dy and Ho)

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Single-crystal neutron diffraction experiments under 10 GPa class pressure were performed on tetragonal DyB<sub>2</sub>C<sub>2</sub> and HoB<sub>2</sub>C<sub>2</sub> compounds showing antiferroquadrupolar (AFQ) and antiferromagnetic (AFM) orders. Both magnetic structures in the ground state phases are basically same and described mainly by the propagation vectors k1 = (1, 0, 0) and k2 = (0, 1, 1/2) for an AFM alignment in the c-plane and an orthogonal arrangement along the c-axis. This characteristic zigzag structure is attributed to the competitive coexistence of the underlying AFQ order with the AFM order. Our purpose is to investigate pressure effects on the AFQ and AFM orders through the alteration of the magnetic structures with k1 and k2. The high pressure experiments revealed that applying pressure enhanced the AFM interactions and simultaneously suppressed the AFQ orders in the both compounds. In addition, we found that applying pressure shranks the crystal volumes more greatly in the direction of the c-axis even under hydrostatic pressure. Therefore, we presume that the volume shrinkage by applying pressure constrains the buckling accompanied by the AFQ orders, which means a local atomic displacement of (B-C) layers along the c-axis, and leads to the suppression of the AFQ orders in the both compounds.

## QI17

## Antiferromagnetism of TbPd2Ge2 Single Crystals

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The compound TbPd<sub>2</sub>Ge<sub>2</sub>, having the tetragonal ThCr<sub>2</sub>Si<sub>2</sub>-type crystal structure, exhibits an antiferromagnetic order below TN=4 K [1]. The details are not clear. In order to clarify magnetic behavior, we have grown TbPd<sub>2</sub>Ge<sub>2</sub> single crystals, and carried out magnetic and specific heat measurements using the single crystals. The magnetic susceptibility shows an anomaly at 13 K in susceptibility vs. T curves, indicating an antiferromagnetic order. The appearance of anomaly at 13K is confirmed in specific heat. This transition temperature is fairly higher than one reported previously [1]. The easy magnetization direction is the [100] direction in the basal plane within this experimental magnetic field up to 7 T. Three metamagnetic transitions appear in the process at low temperatures. In the [110] magnetization process, two metamagnetic transitions are observed. In magnetization curves along the all directions, downward curvatures can be seen, suggesting a precursor of further transition. From these results, the magnetic behavior of this compound is discussed in detail.

[1] A.Szytula et al., Solid State Commun.66 (1988) 309.

### QI18

Multi-step metamagnetic processes in PrPd<sub>2</sub>Si<sub>2</sub> single crystal

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Magnetic studies have been performed in detail on a PrPr<sub>2</sub>Si<sub>2</sub> single crystal which is crystallized in the tetragonal ThCr<sub>2</sub>Si<sub>2</sub>-type structure. The temperature dependence of magnetic susceptibility indicate the compound orders antiferromagnetically below TN=3.0 K. The transition has been confirmed by specific heat. Magnetization measurements up to 18 T have been performed and obtained following results. The easy magnetization direction is the [100] direction in the basal plane. In the magnetization process, four metamagnetic transitions appear; it is a four-step metamagnetic one. The [110] magnetization process is a three-step one. A strong magnetic anisotropy between the [100] and [110] directions is observed within the basal plane in high magnetic fields. Along the hard magnetization direction of the c-axis, a metamagnetic transition appears. A magnetic anisotropy between the [100] and [001] directions is solven. These behaviors are discussed from an analysis of crystal field effects.

### QI19

### Anisotropic properties of Tb<sub>2</sub>Pd<sub>2</sub>In single crystal

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Tb<sub>2</sub>Pd<sub>2</sub>In (tetragonal Mo<sub>2</sub>FeB<sub>2</sub> structure type) was reported to be an antiferromagnet with the Neel temperature TN = 32 K and metamagnetic transition at  $\mu$ 0Hc = 4 T [1]. On the contrary, our magnetization measurements on polycrystalline sample indicate that the magnetization curve of Tb<sub>2</sub>Pd<sub>2</sub>In has 4 metamagnetic features in the field range to 14 T. We have succeeded to prepare singe-crystalline sample of Tb<sub>2</sub>Pd<sub>2</sub>In using Czochralski method. The magnetic susceptibility measured on single-crystal in various crystallographic directions shows that magnetic properties of Tb<sub>2</sub>Pd<sub>2</sub>In are highly anisotropic. c-axis was found to be an easy magnetization direction. The magnetization curve along the c-direction at T = 2 K shows two metamagnetic transitions ( $\mu$ 0Hc = 5 and 8 T). The magnetization is completely saturated at 14 T (Ms = 19.1  $\mu$ B/ f.u.). We have also observed another metamagnetic transition along the a-direction. The magnetization in the a-direction does not show any tendency to saturation in the highest available fields.

[1] M. Giovannini et al., J. Alloys Compd. 280, 26 (1998).

# QJ01

## Successive magnetic transitions of PrRh<sub>2</sub>Ge<sub>2</sub> single crystal

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The ternary compound PrRh<sub>2</sub>Ge<sub>2</sub>, having the tetragonal ThCr<sub>2</sub>Si<sub>2</sub>-type crystal structure, shows an interesting magnetic behavior; it orders antiferromagnetically below TN=45 K and exhibits an additional magnetic transition at Tt=30 K [1]. The details have been, however, unknown yet. In order to elucidate its magnetic behavior, measurements of magnetization and specific heat have been performed on a PrRh<sub>2</sub>Ge<sub>2</sub> single crystal. A much large magnetic anisotropy with the easy magnetization direction of c-axis is observed. Four anomalies appear in the temperature dependence of magnetic susceptibility along the c-axis; a cusp at TN=52 K, a high peak at T3=38 K, a shoulder at T2=33 K and a small peak at T1=4 K are observed. Specific heat also shows four anomalies at corresponding temperatures where magnetic anomalies appear, confirming an occurrence of successive magnetization direction of the c-axis at low temperatures. The magnetization in the first step is almost flat and is corresponding to (1/5) Ms, indicating this field-induce phase has a long period magnetic structure. This behavior is discussed as one of characteristic features for lising-like system with competing exchange interactions.

[1] A. Szytula et al., Intermetallics 18 (2010) 1766.

## QJ02

# Phase diagram of the magnetically frustrated system $SmPd_2Al_3$ studied by neutron diffraction

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SmPd<sub>2</sub>Al<sub>1</sub> represents a distinctive example of Sm magnetism exhibiting a complex magnetic behavior [1]. The complexity of magnetism is caused by specific features of the Sm<sup>3+</sup> ion, namely by nearness of the ground-state multiplet J = 5/2 and the first excited multiplet J = 7/2 in conjunction with strong crystal field. We have established a magnetic phase diagram based on magnetization data. Detailed information regarding the magnetic trutter of each magnetic phase was lacking and neutron diffraction was highly desirable. To reduce the neutron absorption by the Sm we have performed an experiment at the ILL Grenoble D9 diffractometer with the short wavelength neutrons which are much less absorbed by the natural Sm. The four successive magnetic phases have been detected within the neutron diffraction experiment providing the values of the critical temperatures 3.5, 4, 4.5 and 12.4. Despite the still strong Sm neutron absorption and the low effective magnetic interval 12.4.4.5 K which classifies the SmPd<sub>2</sub>Al<sub>3</sub> to the unique group of magnetically flustrated systems. Discussion will be drawn in the context of recent theoretical approaches to geometrically flustrated magnetic systems.

 [1] J. Pospisil, M. Kratochvilova, J. Prokleska, M. Divis, V. Sechovsky; Phys. Rev. B 81, 024413 (2010)

### QJ03

# Interplay of rare-earth and iron sublattices in NdFeAsO

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We report zero-field (ZF) muon spin relaxation ( $\mu$ SR) measurements on NdFeAsO polycrystalline samples. We investigate the interaction of the Fe magnetic sublattice with the RE (rare-earth) paramagnetic moment in the nonsuperconducting magnetically ordered Nd oxypnictide by a detailed analysis of the hyperfine fields at the muon site. At high temperatures, the fluctuations of the Nd moments are evident from the observed exponential depolarization of the µSR signal. Microscopic evidence of static commensurate magnetic order of Fe moments is revealed by  $\mu$ SR at temperatures below TN(Fe) = 136K. Drastic slowing down of RE moment fluctuations is observed only below TN(RE) = 2K. In the intermediate temperature range, a Schottky-like polarization of the Nd magnetic sublattice by the ordered Fe system is observed. In the Fe system, a transition between a high- and a low-temperature indicateuros is proposed on the basis of a detailed symmetry analysis and magnetic dipole-field calculations at the muon site. The observed behavior resembles the one of the system in the structurally-related Nd<sub>2</sub>CuO<sub>4</sub>.

[1] H. Maeter, H. Luetkens, Y. Pashkevich, et.al. Phys. Rev. B. 80 094524 (2009). [2] W. Tian, et.al. Phys. Rev. B. 82, 060514 (2010). [3] Y. Qiu, et.al. Phys. Rev. Lett. 101, 257002 (2008).

# QJ04

## Magnetic properties of C15 laves phase compound SmRu<sub>2</sub>

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We report the magnetic properties of polycrystalline C15 laves phase rare-earth compound SmRu<sub>2</sub>. The sample of SmRu<sub>2</sub> was made by an arc-melt method and was annealed at 500 °C in vacuumed quartz tube for 50 hours. The susceptibility  $\chi$  of SmRu<sub>2</sub> follows a typical Curie-Weiss law above 80 K. The effective magnetic moment peff is 1.39  $\mu$ B/Sm-atom, and the Weiss temperature  $\theta$  is a positive value of 42 K. The low-temperature  $\chi$  for T < 70 K increases rapidly with decreasing temperature. In the resistivity  $\rho$  and the specific heat Cp, we found a small anomaly and a jump at about 55 K. Furthermore, we observed a hysteresis in the magnetic field dependence of magnetization at 2 K. These experimental results suggest that C15 laves phase SmRu<sub>2</sub> is ferromagnetic material on Tc  $\sim$  55 K. Detailed discussion will be presented at the conference.

# QJ05

# Pressure effect on the electrical resistivity of a ferromagnetic clathrate $Eu_8Ga_{16}Ge_{30}$

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Eu<sub>8</sub>Ga<sub>16</sub>Ge<sub>30</sub> is a unique intermetallic clathrate in which the guest sites are fully occupied by a rare-earth element. The Eu<sup>2+</sup> ions in the tetrakaidecahedral cages in the type-I structure are rattling among off-center positions [1]. This compound exhibits a ferromagnetic transition at Tc=36K. Both the electrical resistivity  $\rho$  and the temperature deviative of magnetization exhibit peaks at Tc and T\*. Upon applying pressure of 6GPa at 290K, the off-center rattling state of Eu ions changes to the on-centered state. Furthermore, at Ps=14GPa, a structural transition occurs with the volume reduction of about 4% [2]. In order to examine the pressure dependence of Tc and T\*, we have measured the temperature dependence of  $\rho$  under various pressures up to 15GPa, by using a Bridgman anvil cell made of sintered diamond. With increasing pressure, Tc shifts to higher temperatures. Broading of the peak at T\* for P≥4GPa is attributed to the change from the off-center to on-center state. At 290K,  $\rho(P)$  jumps at 13GPa, which must be the manifestation of the structural transition.

[1] B. C. Sales, et al., Phys. Rev. B 63, 245113 (2001). [2] T. Kume et al., unpublished.

# QJ06

## Magnetic transitions under pressure in $GdCo_2B_2$

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The pressure effect on magnetic transitions of GdCo<sub>2</sub>B<sub>2</sub> compound was studied by dc susceptibility and ac heat capacity in the temperature range 2 K ~ 300 K. The temperature dependence of dc susceptibility for GdCo<sub>2</sub>B<sub>2</sub> under various pressure in 0.1 T magnetic field show two information. First, the Antiferromagnetic transition temperature TN decrease with pressure increasing from 12.5 K in OGPa decrease to 6.3 K in 1.07 GPa and disappeared near to 1.30 GPa with a slope of dTN/dP about 5.8 K/Gpa. Second, ferromagnetic transition TC increases alightly with pressure increasing from 20.4 K in 0 GPa to 22.1 in 1.27 GPa with a slope of dTC/dP about 1.5 K/Gpa. They were also found in ac heat capacity measurements. It can be understand that there is a strong interaction in the ferromagnetic state and a very weak interaction in antiferromagnetic tate which can easily be destroyed by pressure and magnetic field. The transitions are proved to be the second-order transition. In addition, a transition temperature TM which comes from domains of ferromagnetics m in heat capacity measurements was found increase from 17.5 K under 0 GPa to 22.0 K under 2.36 GPa with a ratio dTM/dP about 2.2 K/GPa.

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# QJ07

# New ferromagnetic compounds with noncentrosymmetric crystal structures

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Recently, physical properties in noncentrosymmetric crystal structures have attracted considerable attentions. The broken inversion symmetry in the crystal causes finite antisymmetric spin orbit coupling (ASOC). A spin-polarized electrons that originates from the ASOC effect is formed even for nonmagnetic materials. Generally, the electronic states of opposite spin orientation are considered to be degenerate in nonmagnetic materials due to the presence of both time-reversal and space-inversion symmetries. However, this degeneracy will be solved by the ASOC in broken space-inversion symmetries with noncentrosymmetric crystal structures. A pair of split bands is produced in the k-space. Ferromagnetic compounds with noncentrosymmetric crystal structures are interesting materials, because the time-reversal symmetries also broken. We discovered new ferromagnetic compouds,  $Pr_4B_{13}$  and  $Nd_4B_{13}$ , with noncentrosymmetric crystal structure (space group 1-43d), which has no inversion center. Curie temperature T\_C are 48K and 94K, respectively. Large ASOC is expected caused by heavy material of Bi. We will report further results in detail.

# QJ08

# Singlet-triplet crossover in the two-dimensional dimer spin system YbAl<sub>1</sub>C<sub>3</sub>

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YbAl<sub>3</sub>C<sub>3</sub> has the hexagonal ScAl<sub>3</sub>C<sub>3</sub> type structure in which magnetic Yb<sup>3+</sup> ions form a two-dimensional triangular lattice. At 80 K, it exhibits a structural phase transition to an orthorhombic structure with a slight displacement of the Yb atoms [1]. Interestingly, a dimer ground state with the singlet-triplet energy gap of about 15 K has been suggested from the magnetic susceptibility and the specific heat at low temperatures [2]. Indeed, a direct observation of the singlet-triplet excitations has been reported from the inelastic neutron experiments [3]. To investigate the variation of the ground states in magnetiz fields, we measured magnetization of a single crystal of YbAl<sub>3</sub>C<sub>3</sub> down to 0.07 K and in fields up to 14.5 T. At low temperatures below 0.8 K, several magnetization steps were clearly observed for H||c in 6 T <  $\mu$ 0H < 9 T, whereas only two transitions were observed at 4.7 and 6.6 T for H||ab. At fields above 9 T, the magnetization becomes almost saturated for both H||c and H||ab. Our results indicate that the singlet-triplet crossover occurs in a relatively narrow field range, suggesting a rather weak interdimer interaction in spite of the nearly triangular lattice of Yb ions.

[1] T. Matsumura et al., J. Phys. Soc. Jpn. 77, 103601 (2008). [2] A. Ochiai et al., J. Phys. Soc. Jpn. 76, 123703 (2007). [3] Y. Kato et al., J. Phys. Soc. Jpn. 77, 053701 (2008).

## QJ09

### Magnetic structure transition in PrPd<sub>3</sub>

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A complex magnetic phase diagram of  $PrPd_3$  with  $AuCu_3$ -type structure was reported [1]. The specific heat shows double peaks at TN = 0.88 K and 0.77 K, which are relatively low compared with the Curie-Weiss temperature  $\Theta p \sim 17$  K. The temperature of the peak at higher temperature is increased but another decreased with the magnetic field. A large tail above TN in the specific heat is also observed. Considering no geometric frustrations in the simple cubic structure of the Pr site, these behavior suggest a competition of inter-site interactions. The magnetic field dependence of the peak at higher temperature, which is typical in the quadrupole ordering system, also suggests an importance of the multipole effect arising from the crystalline electric field ground state. We have investigated the multipole effect using the single crystal by specific heat, magnetic susceptibility, ultra-sonic and also neutron scattering measurements. In this report, we focus on the magnetic phase at 0.5 K. The propagation vector observed in the neutron scattering is changed from the incommensurate to commensurate one. We will discuss the complex magnetic phase based of the multipole effect.

[1] S. Zhang, et al., Journal of Physics : Conference Series, 15 (2009) 042074

## QJ10

# Competition between magnetic ordering and random spin freezing in Dy<sub>2</sub>PtSi<sub>3</sub>

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We have been paying special attention to a new family of the temary intermetallic compounds R<sub>2</sub>PtSi<sub>3</sub>. The spin-glass or cluster-glass-like effect can be observed in some R<sub>2</sub>PtSi<sub>3</sub> compounds with heavy rare earth element, which have been confirmed to crystallize in the ordered hexagonal structure [1]. The mechanisms of the anomalous magnetic behaviors are very interesting in physics. Recently, we prepared a polycrystalline Dy<sub>2</sub>PtSi<sub>3</sub> sample and systematically measured its magnetic properties including ac susceptibility (Xac), FC and ZFC magnetization, magnetic relaxation, electrical resistivity and specific heat. The most important observation is the evident peak in Xac(T) curve near Tc~6 K, which shifts to high temperature with increasing frequency suggesting the presence of random spin freezing effect. The observed long-time magnetic relaxation and the low temperature irreversible magnetism give further evidence for this conclusion. However, specific heat measurement shows an evident peak at Tc and the electrical resistivity drops rapidly as decreasing the temperature just below Tc. In this sense, the long-range magnetic exchange interactions have also important influence on the magnetic behavior. Thus it seems to be appropriate to classify the Dy<sub>2</sub>PtSi<sub>3</sub> sample as a large magnetic cluster system with competition between magnetic ordering and random spin freezing.

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### QJ11

## Dilatometric investigations on the semimetallic ferromagnet EuB<sub>6</sub>

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EuB<sub>6</sub> is a semimetallic correlated electron system, which exhibits a complex sequence of electronic and magnetic phase transitions at ~ 15.5 K (Tc1) and 12.5 K (Tc2). The material also shows a colossal magnetoresistance effect which is largest at Tc1. A percolative transition of magnetic polarons, i.e., clusters of ordered Eu<sup>2+</sup> moments polarized by the spins of localized charge carriers (holes), is involved in the paramagnetic to ferromagnetic transition. Thermal expansion and magnetostriction measurements have been performed on this material to explore to which extent lattice degrees of freedom are involved in these phase transitions. We find two corresponding anomalies in the thermal expansion, the one occurring at Tc2 being much larger than that at Tc1. By applying a small magnetic field of less than 50 mT, the anomaly at Tc1 is fully suppressed, while the lower-temperature anomaly at Tc2 shifts to higher temperature as the field is increased and finally fades out at a field B > 5 T. These measurements for a set of temperatures from below Tc2 to above Tc1, highlighting the extraordinarily large magnetoelastic effects in this material.

## QJ12

# The magnetic anisotropy and magnetostriction of $RAl_2 \, (R{=}\, rare \, earth)$ compounds

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The RAI<sub>2</sub> (R : rare earth element) intermetallic compounds crystallize in the MgCu<sub>2</sub> cubic Laves phase structure. In most of these compounds the R ions are trivalent 3+ and ferromagnetic ordering has been observed with Curie temperature TC up to 170 K. It is expected that some of structural transition occurs due to magnetic symmetry braking below TC. In the present work, spontaneous magnetostriction of ferromagnetic RAI<sub>2</sub> compounds are studied from the X-ray diffraction measurement at low temperature and the thermal expansion of the single crystal. The X-ray diffraction pattern shows the d-spilliting at several peaks of TbAI<sub>2</sub> and DyAI<sub>2</sub> due to the large magnetic anisotropy below TC. From these results, the crystal structures are estimated at ferromagnetic phase of TbAI<sub>2</sub> and DyAI<sub>2</sub>. The anisotropic magnetoelastic constants have been calculated by using the results of the magnetization and the magnetostriction at the ferromagnetic phase and at the magnetic field up to 2 T.

# QJ13

### Influence of weak-magnetic, non-magnetic and disoriented grains on remagnetization processes of Nd-Fe-B alloys

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The uniaxial high-anisotropic materials, in which magnetostatic interaction between volumes have a significant influence were investigated. The analysis of domain structure formation during magnetization reversal these materials in the presence of weak-magnetic, non-magnetic and disoriented grains was carried out. The modeling ensemble simulating magnetic properties of sintered alloy neodymium- iron - boron  $(Nd_{12}Fe_{22}B_{2})$  was selected Research of influence of inclusions was conducted on the basis of model experiments with the received ensemble. The program, devised for uniaxial high-anisotropic ferromagnetics, in which the sample is presented by ensemble from 1000 particles having transitive domain structure and following individual characteristics was used. After selection of all parameters of the ensemble simulating this material, for researching the influence of inclusions on magnetic reversal processes. key parameters of some microvolumes (Ms. Homax, Hs. Angle) were changed. It was concluded that magnetic reversal is realized by the channels formed at the expense of directed magnetostatic interaction. Weak-magnetic, non-magnetic and disoriented grains are the magnetic reversal centers, i.e. they are the beginning of formation of the channel of magnetic reversal.

## QJ14

### Cobalt magnetism in HoCo2 under pressure

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HoCo<sub>2</sub> belongs to a group of RCo<sub>2</sub> (R = rare earth metal) compounds forming in the cubic Laves phase structure. The Ho localized 4f-electron magnetic moments coexist together with the itinerant Co 3d moments, which appear on the verge of magnetism, similar to the ErCo<sub>2</sub> system [1]. Below TC, the large exchange field due to the ferromagnetically ordered Ho moments polarizes the Co 3d-electron states and the emerged moments in the Co sublattice orient antiparallel to the Ho sublattice. Weak Co moments survive rather far above TC (= 79,5 K) in Co magnetic clusters remaining coupled antiparallel to paramagnetic Ho moments forming a parimagnetic phase. This causes a susceptibility anomaly at the parimagnetic transition temperature Tf (= 125 K). At low temperatures the easy magnetization direction changes at the spin-reorientation temperature TR = (16,4 ± 0,5) K. We will present variations of the three transition temperatures with applying hydrostatic pressure up to 3 GPa. TC and TR coincidently decrease with increasing pressure whereas TR increases. The results will be discussed in term of suppressing the Co magnetic moments and varying the exchange and crystal-field interactions with applying the hydrostatic pressure.

[1] J. Herrero-Albillos et al., Phys.Rev.B 76 (2007) 094409

## QJ15

### High-field magnetization study of ErCo<sub>2</sub>

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In this work are presented new results on the magnetic properties of the  $ErCo_2$  compound in which three magnetic phases (ferrimagnetic, parimagnetic and paramagnetic) were proposed very recently. This study is based on the investigation of the magnetization under continuous high magnetic fields (up to 330 kOe at 4.2K and up to 230 kOe in the 4.2-300K respectively). At low temperatures the stability of the ferrimagnetic structure is shown although the saturation was never found. Special attention was paid to the 35-100K temperature range where the transitions between the different magnetic phases occur confirming of the parimagnetic" orderings. The transition temperature Tc from the ferrimagnetic to the parimagnetic phase increases from ? 34 to 62 K with the applied field. For the values of H smaller than 100 kOe. At H values greater than this Tc remains constant at (62  $\pm$ 5) K. The atomic magnetic moment of Co is independent of the magnetic structure is and in the order of 1µBatom-1.

# QJ16

### Resonance, magnetic and neutron investigations of magnetic structures in $Pr_{1x}Y_{x}Fe_{3}(BO_{3})_{4}$ system

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The rare earth borates RFe<sub>3</sub>(BO<sub>3</sub>)<sub>4</sub> with the structure of the huntite mineral are interesting as materials with two magnetic subsystems with competing anisotropy. In this work we present data on antiFerromagnetic resonance (AFMR), magnetic properties and neutron investigation of Pr<sub>1,x</sub>Y<sub>1</sub>Fe<sub>3</sub>(BO<sub>3</sub>)<sub>4</sub>, system. Our AFMR and neutron investigations [1-3] show that the last points of the system, YFe<sub>3</sub>(BO<sub>3</sub>)<sub>4</sub>, and PFe<sub>4</sub>(BO<sub>3</sub>)<sub>4</sub>, are easy-plane (EP) and easy-axis (EA) antiFerromagnets, respectively, confirming the magnetic data for R=Pr [4]. A dilution of the Pt<sup>3+</sup> subsystem by diamagnetic yttrium reduces the magnetic anisotropy of the Pr subsystem and can lead to a transition from EA to EP state at some ytfrium content. The single crystals Pr<sub>1,x</sub>Y<sub>x</sub>Fe<sub>3</sub>(BO<sub>3</sub>)<sub>4</sub>, with x=0.25, 0.325, 0.5 and x=0.75 were grown. Magnetic and resonance investigations show that the crystals with x = 0.25 and 0.325 are EA antiferromagnets with the energy gap at T=4.2 K close to 75 and 60 GHz, respectively, but the crystals with x=0.5 and 0.75 are EP ones. Neutron investigations showed that the concentration from EP to EA state occurs via an inclined magnetic structure which was found for x=0.325.The concentration dependency of AFMR energy gap is studied at T=4.2 K. This work was supported by RFBR grant #10-02-00765.

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## QJ17

Anisotropic transport and magnetic properties of PrGe single crystal Pranab Kumar Das, Neeraj Goyal, Ruta Kulkarni, Arumugam Thamizhavel and Sudesh Kumar Dhar

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The compounds RGe (R = La, Ce, Pr and Nd) crystallize in the two closely related CrB and FeB-type orthorhombic crystal structures, depending upon R and heat treatment. We have recently reported the anisotropic magnetic properties of CeGe in detail which crystallizes in the FeB-type crystal structure [1]. In continuation of our studies on RGe compounds, we report here the anisotropic properties of a single crystal of PrGe grown by Czochralski method. PrGe exhibits dimorphism in which the high temperature FeB-type structure transforms into CrB-type structure at lower temperatures [2]. From the powder X-ray diffraction we have confirmed that the grown single crystal of PrGe possesses CrB-type crystal structure with the space group Cmcm(#63). The anisotropic magnetic properties have been studied along the three principal crystallographic directions. From the magnetic and transport properties we find that PrGe exhibits two magnetic transitions at 44 K and 41.6 K along [100] and [010] directions, whereas along [001] direction only a ferromagnetic ordering is antiferromagnetic and the low temperature one is ferromagnetic. The specific heat capacity data also confirm the two magnetic transitions in PrGe.

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## QJ18

# Theoretical investigation of the phase transition and the spin-gap behavior of the triangular antiferromagnet YbAl<sub>3</sub>C<sub>3</sub>

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Ytterbium aluminum carbide YbAl<sub>3</sub>C<sub>3</sub>, first reported in 1992,[1] has recently attracted much attention concerning the nature of its phase transition at T\* = 77 K.[2-4] YbAl<sub>3</sub>C<sub>3</sub> containing triangular sheets of magnetic ions Yb<sup>3+</sup> (13) undergoes a phase transition at T\* = 77 K and exhibits a spin-gap behavior. The two model structures of YbAl<sub>3</sub>C<sub>3</sub>, proposed to explain these observations, give unreasonably short C-Al distances. These model structures were optimized by density functional calculations to find that the Model 2 structure is more stable than the Model 1 structure, with chemically reasonable C-Al distances, and that the phase transition at T\* is a cooperative second-order Jahn-Teller distortion involving the layers of corner-sharing CAl5 trigonal bipyramids. The spin exchanges between nearest-neighbor Yb<sup>3+</sup> ions, evaluated by density functional calculations for the optimized Model 2 structure; show that YbAl<sub>3</sub>C<sub>3</sub> has a two-dimensional spin lattice described by three antiferromagnetic exchanges, but exhibits a spin-gap behavior because one exchange is substantially stronger than the other two.

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# QJ19

# First Principle Analysis for Pressure Effect on Charge Density Wave Phases of ${\rm SmNiC}_2$

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Using a first principles calculation, we have investigated pressure effect on the electronic structure anisotropies, Fermi surface properties, and transport properties of SmNiC<sub>2</sub> to check their effects on the charge density wave instabilities. Both band structure and electric conductivity results show that SmNiC<sub>2</sub> has quasi 1D electronic structure by Ni 1D chain along a axis. However hybridization between nearest Sm 1D chain makes its charge density wave (CDW) properties 2D like with CDW vector  $q_i$ =(0.5, 0.5, 0) and another weak CDW vector  $q_R$ =(0.5, 0.5, 0.5) at -6 GPa. With increasing pressure, highly anisotropic change of lattice induces these CDW vector to become incommensurate with flattening Fermi surface (FS) in the [110] plane and suppressing FS in [111] plane. The FS nesting vector  $q_i$  is gradually shifted from  $q_i$ =(0.5, 0.5, 0, 0) at -6 GPa to  $q_i$ =(0.5, 0.56, 0) at 14 GPa, which is significant local maximum at -6GPa, becomes almost saddle point like with its nesting vector shifting another points. We conclude that CDW strength along  $q_i$  vector become stronger with flattening (more 1D like) FS, but CDW strength along  $q_k$  vector become weak with highly anisotropic change of lattice.

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# QK01

### The magnetoresistance of sandwich-structure organic spin-valve Feng Li\* Fapei Zhang and Yu Xiao

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We studied the magnetoresistance (MR) effect in an organic spin-valve structure of ferromagnetic/organic semiconductor/ferromagnetic system theoretically using the spin-diffusion theory and Ohm's law. Spin polarons and spinless bipolarons are assumed to be the main carriers in organic semiconductors. From the calculation, it is obtained that MR value increases with the increase of the polaron proportion and rapidly decreases with the increase of the thickness of organic layer. MR ratio can be enhanced remarkably when the interfacial resistances are spin related. We also investigated MR value dependence of the effects of the matching conductivity and the spin-polarization of ferromagnetic layer.

## QK02

### First-principles calculations investigation of interfacial roles in spindependent transport properties in OMTJs

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Organic spin valves and organic magnetic tunnel junctions (OMTJs) are new and promising research fields and show advantages for spin-dependent transport due to relatively weak spin-orbital coupling and hyperfine interaction. Long-chain alkanoic acids usually form close-packed monolayer films with alkyl chains highly oriented on substrates. The characterization of stearic acid adsorbed on Ni(111) surface has been studied. Our results suggest that the stearic acid is chemically adsorbed on the Ni surface via a bidentate interaction with a distance of about 1.8 A. We also investigated spin-dependent transport characterization through two kinds of organic magnetic tunneling junctions (OMTJs), as Ni/1-stearic acid radical (1-SAR)/Ni and Ni/1 18-stearic diacid radical (1.18-SDR)/Ni. We found magnetoresistance (MR) of Ni/1-SAR/Ni is -19 6% and it is 13 7% for Ni/1 18-SDR/Ni OMTJ. There is a magnetic proximity effect in the interface, and the magnetic decay rate is dissimilar in different interfaces. Due to interaction at hybrid organic/magnetic interface the electronic properties of interfacial atoms are modified. These effects influence the spin-dependent transport properties, and there appears positive or negative tunneling magnetoresistance (TMR) due to different interface. Our results suggest, the interface between organic barrier and FM electrode play an important role in the transport behavior.

# QK03

### GMR properties on flexible polymer film with bending stress

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Recently, demands and researches about variable flexible devices have been constant. In other to satisfy such demands, researches about magnetic properties in flexible environment also have been continued. In these researches, the simple giant magnetoresistance (GMR) spin valve structure was deposited on flexible polymer films then was induced by external tensile stress from repetitive bending fatigue. As bending samples, bending directions were decided whether parallel or transverse directions to easy axis directions of GMR structure. As a result, when increasing parallel bending times switching field interval range of free layer also increased inversely as increasing transverse bending times, coercivity of free layer also increased. Such tendency had opposite properties to inverse magnetostriction effect theory from the view of bending direction, because flexible films had larger elastic recovery forces than tensile forces after bending. Therefore we anticipate that such elastic recovery forces was the reason of opposite tendency from the aspect of direction. Based on these results, we clarified possibilities that the coercivity could be controlled without curvature, moreover axis behavior, e.g., easy axis or hard axis, also could be controlled without curvature and extra annealing processes.

### QK04

### Giant magnetoresistance in graphene nanoribbons: Geometry, interface and dephasing effects Stefan Krompiewski\*

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Motivated by the ongoing debate on the role of external contacts on electronic transport through carbon nanostructures [1-3], we theoretically estimate the combined effects - due to (i) geometry, (ii) the way the external contacts are attached and (iii) possible dephasing processes - in graphene nanoribbon (GNR). The method is based on the tight binding model, the Green function technique and the Landauer-type formalism. It is shown that giant magnetoresistance in GNRs depends strongly not only on the aspect ratio of the sample and the current direction, but also on whether ferromagnetic electrodes are attached in the end- or side-contacted manner. Typically the former attachment is favorable for conductivity. However the observed tendency is that the higher conductive systems reveal the lower giant magnetoresistive effect, and vice versa. Accordingly, the dephasing usually decreases conductivity, thereby increasing the GMR.

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## QK05

### Negative magnetoresistance in ferromagnet/semiconductor/ ferromagnet structures with cubic dresselhaus spin-orbit-interaction Kenii Kondo\*

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We have investigated the spin transport in Ferromagnet (FM)/Semiconductor (SC)/ Ferromagnet (FM) structures with a central SC barrier exhibiting cubic Dresselhaus spin-orbit-interaction (SOI). The energy profile of the barrier is assumed to be a square with height V and thickness d along z-direction. The magnetoresistance (MR) ratio has been calculated for three different barriers, GaAs, GaSb, and GaAs without SOI as a function of barrier thickness. Both ferromagnets are assumed to be Fe. Intriguingly, MR ratios exhibit negative values for two different barriers in some range of thickness. For GaAs, the MR ratio is negative when the thickness is larger than 0.03 nm. For GaSb, the MR ratio is negative in the range of 0.02 to 0.9 nm, changing the sign gradually from negative to positive at a thickness of about 0.9 nm. Afterwards, its value increases monotonically and reaches a constant value at a thickness of 2.5 nm. We can attribute these phenomena to spin filtering effects caused by both the spin precession and the energy-splitting in the central SC region. Without SOI, the MR ratio shows positive value in all the range since the spin does not rotate in the barrier.

# QK06

### Spin valve effect of NiFe/graphene/NiFe junctions Muhammad Zahir Iqbal, Muhammad Waqas Iqbal and Jonghwa Eom\*

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Novel linear dispersion in graphene makes unique electronic properties such as ambipolar transport and Dirac particle quantum Hall effect. Although graphene was previously considered to be a physically unstable form until it was shown to exist on the Si substrate, various methods to synthesize graphene have been developed. Exploiting such superior electronic properties combined with recently developed large scale growth of graphene by chemical vapor deposition (CVD), has led to field effect devices integrated on a wafer scale showing promise for future electronic applications. In this paper, CVD-grown graphene is used as intervened spacer between two ferromagnetic electrodes. The device is composed of 55 nm thick NiFe film, graphene and 75 nm thick NiFe film. The current-to-perpendicular-to-plane magnetoresistance (MR) shows the spin valve effect. The increased resistance in the antiparallel configurations of magnetizations is due to the spin dependent scattering inside the magnetic films or at interfaces. A positive MR ratio of the order of 0.045% has been observed at 300 K and this signal grows up to 0.14 % at 10 K. The junction resistance is found to be monotonically increasing with temperature. Graphene functions as not a tunnel barrier but a resistive thin film between two NiFe electrodes.

### QK07

# Giant magnetoimpedance and photoinduced magnetoresistance effects in ferromagnet/SiO<sub>2</sub>/p-Si hybrid structures

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We report the electron transport properties in the Fe/SiO2/p-Si hybrid structure. The results of dc studies show that the features of the transport properties of the structure are caused by a metal/insulator/semiconductor transition with a Schottky barrier formed at the SiO2/p-Si in-terface [1]. The magnetoresistance effect is determined by contributions of the processes occurring in the volume of a semiconductor and in a inversion layer near the SiO<sub>2</sub>/p-Si inter-face. Study of the impedance Z = R + iX and magnetoimpedance allows separating the dynamic contributions determined by charge and spin transport in different regions of the hybrid structure. The strongest effect of a magnetic field on R and X is caused by the presence of magnetic impurity centers in the semiconductor near the SiO2/p-Si interface. Magnetoimpedance originates from spin-dependent tunneling between a ferromagnetic electrode and magnetic impurity centers via a potential barrier (SiO2). The magnetic impurity centers determine also the change in photoconductivity in a magnetic field. The structure, as a matter of fact, exhibits giant magnetoresistance induced by optical radiation. As a mechanism of this effect, we consider the spin-dependent tunnel photocurrent, i. e., the current of photoexcited electrons, from p-Si to Fe via the SiO2 barrier.

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## QK08

# Charge imbalance with the same decaying length as spin accumulation ${\it Yao-hui}\ Zhu^*$

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Starting from the Valet-Fert theory of the current-perpendicular-to-plane giant magnetoresistance, we study the charge imbalance in ferromagnetic metals by solving Poisson's equation. Our results show that, in ferromagnetic layers, the charge imbalance has two exponential terms with different decaying lengths: the Thomas-Fermi screening length (on the order of angstrom) and the spin diffusion length (tens of nm in 3d ferromagnetic metals). The charge imbalance on the scale of the screening length has been well studied and it is also present in spin-unpolarized transport in nonmagnetic multilayers. However, the charge imbalance on the scale of the spin diffusion length has never been studied before and it shows up only in ferromagnetic layers with spin accumulation, which also decays on the scale of the spin diffusion length. This charge imbalance is essential to the giant magnetoresistance effect and thus one should be cautious when using quasi-neutrality condition, which neglects the charge imbalance.

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## QK09

### Simulation of spin-dependent transport in GaAs: Effect of electronelectron interactions

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We investigate the effect of electron-electron interaction on electronic spin transport in bulk GaAs. We use Ensemble Monte Carlo (EMC) techniques to simulate the electron and spin dynamics. EMC has been successfully used in the study of charge transport and more recently in spin transport. However electron-electron interactions are often neglected, also due to computational costs. Here, in addition to electron-electron scattering. Spin decoherence occurs due to the spin-orbit interaction (Dresselhaus term), which arises from the lack of inversion symmetry in GaAs. The electron-electron interaction is described by the screened Coulomb potential. At variance with previous simulations, we implement both spin dependent (Mott) and spinless (Rutherford) scattering. We also propose a new, more realistic, method for choosing partners undergoing electron-electrons. We show that, in EMC simulations, electron-electron interactions should not be discarded, as doing so may result in an anomalous drop in the spin relaxation time at temperatures as high as 100K.

## QK10

# Electrical detection of spin-polarized current in InAs quantum point contacts

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The controlled generation, manipulation and detection of spin-polarized currents in semiconductor is an important issue in spintronic devices [1]. The Rashba spin-orbit interaction has been focused to manipulate spin currents. Furthermore, the quantum point contact (QPC) having the Rashba interaction was recently proposed as a spin-current generator. [2], [3] In this work QPCs were fabricated from InAs quantum well. Gates were installed on the sides of the quantum well, and narrow conducting channels, i.e. QPC, were formed by application of voltages on theses gates. Two QPCs were connected in series and the distance between them was less than 1.6 micro-meter. An external magnetic field was applied in parallel with the line connecting the two QPCs. Oscillations of resistance according to the magnetic field have been observed in this device at 1.8K. We believe that these data can be an evidence for our QPCs to work as a spin polarizer with a following scenario; the first QPC works as a spin-polarizer, and the second does as a detector. The spin-current produced by the first QPC precesses by the external magnetic field, which makes the oscillatory resistance as the magnetic field changes.

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## QK11

# Gate dependence of spin-orbit interaction in a two-dimensional hole gas structure

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In order to implement spin-Field Effect Transistor (spin-FET), gate control of spinorbit interaction parameter ( $\alpha$ ) is a key factor in a quantum well structure because spin precession is influenced by strength of the  $\alpha$  [1]. Many papers utilized n-type channel for the spin transport, however, for the complementary logic device p-type channel should be also necessary. We have investigated the spin-orbit interaction parameter and the effective mass using the Shubnikov-de Haas (SdH) oscillation measurement in a GaSb two-dimensional hole gas (2DHG) structure. The  $\alpha$  of  $1.71 \times 10^{11}$ eVm and effective mass of 0.98m0 are obtained at T = 1.8 K, respectively. We also found the gate dependence of the  $\alpha$  and the hole concentration are reduced, which indicates the  $\alpha$ increases with the carrier concentration in p-type channel. On the order hand, n-type channel shows opposite gate dependence [1, 2]. Therefore, the combined device of p and n-type channel spin transistor would be a good candidate for the complementary logic device.

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# **OK12**

Perpendicular spin transport in ferromagnet/MgO/GaAs structures Joohyung Bae<sup>1</sup>, Kyung-ho Kim<sup>1</sup>, Hyun Cheol Koo<sup>1</sup>\*, Hyung-jun Kim<sup>1</sup>, Joonyeon Chang<sup>1</sup>, Suk Hee Han<sup>1</sup> and Sang Ho Lim<sup>2</sup>

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The electrical spin transport from TbFeCo/CoFeB into GaAs using MgO tunnel barrier is investigated. The 20 nm thick TbFeCo layer produces perpendicular magnetization and the 0.8 nm thick CoFeB layer, which contacts the MgO/GaAs layer, enhances the spin polarization of the injector. The MgO layer has a role of spin filter for efficient spin injection. In order to reduce the contact resistance, the heavily doped GaAs layer is inserted between TbFeCo/CoFeB/MgO and GaAs substrate. We measured spin accumulation and spin lifetime using three-terminal Hanle measurements with temperature range from 1.8 K to 300 K. In this geometry, the bias current is applied from the top TbFeCo layer to the GaAs channel and the voltage is measured at the interface. The perpendicular spin is injected so the in-plane magnetic field is swept to observe the spin relaxation at the interface. For a bias current of 1.5 mA, the spin signal of 0.12 mV and 0.03 mV are obtained at T = 1.8 K and 300 K, respectively. From the analysis of the Hanle curve, a spin life time of 0.4 ns is extracted at 300 K.

# **OK13**

### Spin hall effect in 2DEG in the presence of Rashba spin-orbit interation

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Spin Hall effect using an InAs-based 2DEG system with a strong Rasba effect is researched. To make all electric devices, permalloy is used as spin injector on Hall bar geometry. Two kinds of equivalent measurement, inverse and direct spin Hall, are performed. Inverse and direct spin Hall are the conversion of spin to charge and charge to spin, respectively. Also to confirm the effect of precession, dependences of channel length and magnetization direction are systematically investigated. The bias current between ferromagnet and neutral region of channel induces the injection of in-plane spin. Inside the channel, spin precession occurs and then spin orientation is changed to the perpendicular direction at the Hall cross. The spin signal as a function of channel length shows oscillation behavior which is also evidence of spin precession. At 1.8K, a spin Hall voltage of 4 µV is detected for a bias current 0.1 mA and a channel length of 0.64 µm. From this experiment, we confirm the spin Hall effect and spin precession at the same time.

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## **OK14**

### Spin transport and spin injection into turbostratic graphene

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The remarkable physical properties of graphene, a monolayer of carbon atoms packed into a twodimensional hexagonal lattice, make it a promising material for applications. In particular, the high electron mobility [1] and the relatively long spin lifetime (weak interaction of the electrons spin and orbital degree of freedom [2]) resulting in long spin diffusion lengths  $\lambda$  up to 2 µm [3]. make it an interesting candidate for spintronics. However,  $\lambda$  is currently limited by the strong interaction between graphene and the underlying substrate and by the intrinsic corrugation of mechanically exfoliated graphene sheets. This can be overcome by using turbostratic graphene (TG), a multilayer graphene stack consisting of tens of electronically decoupled lavers. TG combines the exciting properties of graphene with the higher robustness to environmental influences and the absence of inner corrugations of a microstructured material [4] In this study we have successfully performed spin injection from ferromagnetic electrodes into TG discs and detected the spin transport using a non-local spin valve configuration [3]. Moreover we can also investigate the correlation between the charge carrier transport properties and the spin injection efficiency.

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## **OK15**

### Spin signal in metallic lateral spin valves made by a shadow evanoration technique

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Amidst the large variety of proposed spin structures, Lateral Spin Valves (LSV) devices attracted recently an increasing attention. These lateral devices consist in two ferromagnetic electrodes connected by a non-magnetic metallic wire allowing to tune the transport properties according to the respective magnetic states of the two ferromagnetic electrodes. They allow realizing non-local transport measurements, by separating the spin current from the charge current using a four probe connections to an interface. We are developing non-local spin injection and detection technique in lateral spin valve structure made by electron beam lithography and angle evaporation techniques with a special emphasis on the study of the spin Hall effect. The spin signal depends mainly on the geometry of the LSV, on the quality of the interfaces, and on the choice of the used materials. We present a method used to optimize the spin signal of lateral spin valves using a technique to form and connect nanowires in a single step in vacuum, avoiding interface contamination and oxidation. By using different evaporation angles of materials through a resist mask and by reducing the devices geometry, we succeeded in obtaining spin signals close to 25 mQ at 77K in devices with transparent interfaces.

P. Łaczkowski et al., Applied Physics Express, Volume 4, Issue 6, pp. 063007-063007-3 (2011)

### **OK16**

Hanle effect with in-plane magnetic fields in metallic lateral spin valves Juan-carlos Rojas Sanchez, Laurent Vila, Matthieu Jamet, Piotr Laczkowski, Murat Cubukcu, Williams Savero-torres, Van Dai Nguyen, Alain Marty, Cyrille Beigne and Jean-philippe Attane

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When a charge current crosses the interface between a ferromagnetic (FM) and a nonmagnetic (NM) material, a spin current appears that occurs on the length scale of the spin diffusion. It is possible to detect this spin accumulation by measuring the voltage between the NM wire and the other FM wire. We show that the spin signal amplitude can be enhanced by inserting a natural Al oxide between the interface of Pv and Al. achieving spin signals above 150 mOhms at 77 K (instead of 24mOhms with transparent contacts) and 38 mOhms at RT. When the magnetic field is applied perpendicularly to the LSV plane, the injected spins in the normal channel precess around an axe parallel to H The decoherence and/or the precession induced by the magnetic field can lead to oscillations and disappearance of the spin signal (Hanle effect). Hanle measurements are usually performed by applying H perpendicular to the sample plane [1], the spin precessing in plane. Here we show that it is possible to carry out Hanle measurements by applying H in the sample plane, generating out-of-plane spin precession, and thus gaining a degree of freedom for the control of the spin current.

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# **OK17**

### Detecting the magnetization switching of a ferromagnetic dot using non local spin injection by means of lateral spin valve structures

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In most researches on Spin Torque, the charge and spin currents cannot be separated. To solve this problem, spin injection using lateral spin-valves, made of magnetic and non magnetic materials, have been proposed [1]. In this context, we study the spin torque generated by pure spin current on a ferromagnetic dot, using lateral spin-valve structures. The nanostructure is composed of two parallel ferromagnetic electrodes, one transversal metallic channel, and one ferromagnetic dot at the middle of the structure. The width of wires and dots is 50 nm. To detect the dot magnetization switching, non-local measurements are performed. The spin polarised current is injected on one side using a ferromagnetic electrode. The spin accumulation produces the splitting of the electrochemical potential, enabling the diffusion of the pure spin current along the channel. The pure-spin current is found to be adsorbed by the ferromagnetic dot, decreasing the spin signal amplitude, which is detected using the second ferromagnetic electrode. The non local geometry enables the spin current and the charge current separation. giving the possibility to test the magnetization switching induced by a pure spin current.

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# **OK18**

Perfect spin filter and highly spin-polarized current induced by fano antiresonance effect in the multiple-quantum-dot nanodevices

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To realize perfect spin filter or to generate highly spin-polarized current is one of the fundamental requirements of any applications on spintronics, and in what material or model structure to achieve this goal effectively is still one of the central issues in spintronics. To realize these aims, multiple-quantumdot system which can be man-made in a controllable way have been attracted much attention. In this report, we propose a practical design to realize perfect spin filter in a periodic diamondlike network composed of multiple quantum dots (QDs). In our design, by producing an energy difference between the site energy between the upper and down QDs in the network, the conductance spectra show a valley structure with a zero point due to the Fano antiresonance effect. As the energy difference increases or the network size increases to a not very large value, the antiresonance valley changes to a well-defined insulating band. Moreover, the conditions for the occurrence of the Fano antiresonance and its relation with the formation of the well-defined insulating band are determined. By a Zeeman effect, the spin-splitting conductance spectra show several highly to 100% spin-polarized windows, strongly proposing that this network can be applied as a perfect spin filter.

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## **OK19**

### Transmission of spin polarized photoelectrons across ferromagnet/ semiconductor interfaces using oblique Hanle effect

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We demonstrate spin-dependent transmission of photoelectrons from a GaAs quantum well (QW) into a ferromagnetic FeGa layer under optical spin orientation in a small oblique magnetic field. The application of the oblique magnetic field enables us to measure the spin-dependent electron transmission across the interface due to the inplane spin component of photoelectrons caused by the Larmor precession of electron spins about the magnetic field. A clear field dependence of the spin-dependent photocurrent is observed for a 20-nm-thick FeGa/GaAs QW sample, showing that the filtering of in-plane polarized electron spins which are generated by the Larmor precession occurs at the interface. To get concrete evidence for the spin filtering effect, we measure the magnetic field dependence of the phase shift between the circular polarization of light irradiated and the spin-dependent photocurrent. The phase shift exhibits a hysteresis behavior as a function of magnetic field, in compatible with the inplane magnetization process of the FeGa layer. These results clearly demonstrate that a very efficient spin filtering effect occurs at the FeGa/GaAs interface even at a low magnetic field, providing a promising basis for designing spin detection technology based on spin filtering effect.

## **OK20**

### Electric field tunning and spin coulomb drag in spin field effect transistors (spin-FETs)

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The quasi-bound states which appear as a consequence of the Rashba spin-orbit (SO) coupling, introduce a strongly irregular behavior of the spin-FET conductance at large Rashba parameter [1]. However, by introducing stray electric fields in addition to the SO couplings, the effect of the SO induced quasi-bound states can be tuned. The oscillations of the spin-resolved conductance become smoother, resulting the possibility to control the spin-FET characteristics [2]. Besides these effects, the Coulomb interaction gives rise to new and interesting manybody phenomena. D'Amico et al. discussed recently the physical origin of spin Coulomb drag effect (SCD) in rather general three-, two- and one dimensional systems. Here, a step closer to a realistic device is made: the SCD effect will be investigated in a nanowire spin-FET structure within the framework of self-consistent calculations for the open quantum system [3]. The spin drag trans-resistivity as a measure for SCD effect will be obtained, in the context of adjusting the spin injection efficiency from the contacts. In addition, the SCD effect in the presence of Rashba spin-orbit coupling will be analyzed. The Coulomb interaction will also be introduced in an exact manner for a few electron system, using the exact diagonalization method [4].

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# **OK21**

## The Rashba-type Spin splitting in Pb monolayer on Si and Ge surfaces: a First-principles study

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Rashba spin-orbit splitting of the surface states due to structural inversion asymmetry is a rapidly growing field of research. Specifically, owing to a technological importance in spin-dependent electronics, spin-polarized metallic surface states on semiconductor surfaces would be highly interesting. Along this line, we have performed first-principles calculations of electronic structures of monolayer of heavy element deposited on semiconductor surfaces, Pb/Si(111) and Ge(111)- $\sqrt{3}\times\sqrt{3}$  surfaces, by including the spin-orbit coupling in the form of additional nonlocal pseudopotential projectors. We investigate and compare energy band dispersions and spin polarizations of these systems and simulate angle-resolved photoemission spectra for direct comparison with experimental results. We also discuss our results focusing on the difference from the Rashba splitting in the Shockley surface states on Au(111) surface. This work was supported by the NRF of Korea (Grant No. 2011-0018306), and computational resources have been provided by KISTI Supercomputing Center (Project No. KSC-2011-C3-06)

## **OK22**

### Geometry effect on quasi-twodimensional electron system

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We calculate the local charge density and local spin density in a deformed square lattice system with Rashba spin orbital interaction. Use nonequilibrium Keldysh Green's function formalism, within the tight-binding framework. The system consists of normal metal lead and deformed central part. We treat the deformed system as adding effective magnetic field that introduces Zeeman effect. We also present the charge and spin current. The deformation we adopt is hemispherical shape and Gaussian shape. We will compare the result with the case that deformation vanished.

# **OK23**

### Electrical measurement of spin accumulation and transport in Fe/ AlGaAs heterostructures

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Utilizing the persistent photoconductivity effect, we report on spin accumulation and transport in AlGaAs at different carrier densities under the same experimental conditions without the necessity for making replicas to access different doping levels. We conducted non-local 3 and 4 terminal Hanle effect measurements in AlGaAs to measure spin accumulation and spin transport respectively at different carrier densities. We will present our recent results. especially the carrier density dependence of the spin lifetime derived from our non-local spin measurements. We also report the bias current dependence of the spin lifetime. From optical studies, the spin life-time at zero bias and at low temperature was reported to be larger than 100 ns in n-GaAs on the insulating side and ~80 ns on the metallic side [1]. Based on our non-local 3 terminal measurement in Si:Al<sub>03</sub>Ga<sub>03</sub>As, the extrapolated spin life-time at zero bias and at 5 K was found to be as large as ~2 ns on the insulating side and decreased with increasing bias current. This work was supported by NSF DMR-0908625 and NSFC 10920101071.

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# **OK24**

Pure spin current injection into a Gd wire Seiji Nonoguchi, Tatsuya Nomura and Takashi Kimura\* Kyushu University, Japan

In magnetocaloric effects, the spin entropy is transformed to the thermal entropy because of an interaction between spin and heat. Since the spin entropy can be controlled by the external magnetic field, the environment temperature also can be tuned by the external magnetic field. Although such methods can be operated efficiently, it is difficult to reduce the size of the entire system. An effective magnetic field can be produced also by injecting spin currents via the interaction between the conduction electrons and localized magnetic moments. Therefore, when a pure spin current is injected into the magnetocaloric material, we may control the spin entropy because of the spin transfer torque and/or spin accumulation. Such a cooling effect has a possibility for a highly efficient nanosized refrigeration device. To demonstrate this effect, we prepared a lateral spin valve with the Gd middle wire, which is a magnetocaloric material. We observed the reduction of nonlocal spin signal by inserting the middle Gd wire, meaning the pure spin current was injected into the Gd wire. We also found that the background resistance decreases with increasing DC current. This implies that the device was cooled by the pure spin current injection into the Gd wire.

**OL01** 

Synthesis and characterization of Ba<sub>3</sub>Co<sub>2</sub>Fe<sub>24</sub>O<sub>41</sub> by the proteic sol-gel process

Bruna Andrade, Petrucio Silva and Marcelo Macedo\* Physics, Federal University of Sergipe, Brazil

In this work, we show the production of hexagonal Z-type barium ferrite ( Ba<sub>3</sub>Co<sub>2</sub>Fe<sub>24</sub>O<sub>41</sub>) prepared via the proteic sol-gel process as well as their structural characterization and magnetic properties. The precursor materials were iron nitrate, cobalt nitrate and barium carbonate which formed a sol after being stoichiometrically dissolved in coconut water. We carried out a heat treatment at 100°C/24 h for the formation of an amorphous xerogel and a calcination at 1000°C/1 h for phase formation of Z-type barium ferrite. The material exhibited a magnetization curve as a function of temperature characteristic of ferrimagnetic material with Tc > 300 K and a hysteresis at room temperature with low coercive field (142 Oe) and saturation magnetization around 48 emu/g for a field of 7 T. Our results indicate that the proteic sol-gel process may be a route of great potential for obtaining Z-type ferrites.

# OL03

The magnetic and dielectric properties of multiferroic Bi1-xGdxFeO3 Takuya Yanoh\*, Naoki Sakai, Liming Zhu, Akinobu Kurokawa, Hiromasa Takeuchi,

Shinya Yano, Kazuki Onuma, Takaya Kondo, Kazunari Miike, Toshiki Miyasaka and Yuko Ichiyanagi Physics, Yokohama National University, Japan

BiFeO3 is known as multiferroic material with antiferromagnetic and ferroelectric properties. Especially, ferroelectricity at room temperature is remarkable and BiFeO3 is expected for new device materials. However, BiFeO, has antiferromagnetic ordering as the result, magnetic values such as remanent magnetization Mr and coercive force HC are rather small for further application. In this study, in order to improve magnetic property and to evaluate dielectric property, we tried to synthesize Bi1.xGdxFeO3 particles where x = 0 to 1.0 by our original wet chemical method. The structural, magnetic and dielectric properties of prepared samples were investigated. Crystal structure change from rhombohedral to orthorhombic was found around x = 0.2 from the X-ray diffraction (XRD) measurement. Lattice constants a and c estimated by XRD decreased as Gd ion increased, supporting that Gd ions were doped exactly in BiFeO3. The magnetic properties of Bi09Gd01FeO3 which has rhombohedral (BiFeO3 type) structure exhibited canted antiferromagnetism at 5 K. Both remanent magnetization Mr and coercive force HC significantly increased as Gd ion increased compared with BiFeO3. The dielectric properties of Bi09Gd01FeO3 showed that dielectric loss (tanb) has been drastically improved compared with that of BiFeO3 for about 90%, while real dielectric constant ɛ' decreased about 15%

### **OL04**

Double-exchange interaction in heavily Mn-doped CuO thin films Li Li, Bin Lv, Shen Wang, Jinzhu Cai, Wenqin Zou, Fengming Zhang\* and Xiaoshan

Physics, Nanjing University, China

RKKY exchange double-exchange interactions are two main mechanisms which are proposed to explain the origin of ferromagnetism of Mn-doped CuO films. Consider that double-exchange interactions cannot produce long-range magnetic order at concentrations of magnetic cations of low percent, Heavily Mn-doped CuO films have been prepared by PEVCD in air and in O2. The M-T and M-H curves show that the average magnetic moment of the films prepared in O<sub>2</sub> is larger than those prepared in air, which may attribute to the variation of Mn chemical valence state. Doubleexchange between Mn<sup>2+</sup> and Mn<sup>3+</sup> can be used to understand the ferromagnetism of CuO under heavily Mn-doping.. The Curie temperature varies slightly with the growth atmosphere. The structure and magnetic properties vary with the annealing parameters obviously, which may indicate that the strain and carriers may play roles on the double exchange

## **OL02**

### Ferromagnetism in vanadium doped ZnO thin films grown by pulsed laser deposition

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This paper investigates the structural, compositional and magnetic properties of vanadium doped ZnO films  $((ZnO)_{1x}(V_2O_5)_x$  where x = 0.02, 0.03 and 0.05) deposited by pulsed laser technique (PLD). Thin films were deposited for a constant duration of 1 hr at 2×10-2 mbar O2 ambient gas pressure by using frequency doubled Nd:YAG laser at 750 °C substrate temperature. The preferred orientation in thin films is along the (002) plane of ZnO. The lattice parameter c derived from the (002) diffraction peak increases as vanadium content increases, suggesting that vanadium substitutes for Zn in ZnO lattice. The PL measurement at low temperature shows the emission peak at 3.30 eV which is usually caused by acceptor-related transitions such as free electron to neutral acceptor (FA) and donor-acceptor-pair (DAP). The XPS results show that vanadium exist in V2+ and V4+ valence state which is in agreement with the XRD and PL results and points to the fact that the ZnVO phase, desirable in DMS, is formed in thin film samples. Furthermore, the hysteresis curves exhibiting the ferromagnetism for all thin films samples further support the formation of ZnVO phase in thin films samples.

## **OL05**

### Magnetic properties of Cu-doped GaN films grown by MBE

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The magnetic properties of Cu-doped GaN films have been investigated in dependence of Cu doping level and film thickness. The films were epitaxially grown at different Cuto-Ga beam-equivalent-pressure ratios (BEPR) by molecular beam epitaxy on sapphire substrates with an AIN buffer and were characterized by scanning electron microscopy energy dispersive x-ray spectroscopy, and x-ray diffraction. Within a narrow range of BEPR around 1 % a ferromagnetic behavior with a Curie temperature higher than 400 K is observed. For BEPR > 1 %, nonferromagnetic Cu-Ga islands are predominantly formed at the surface. The saturation magnetization of ferromagnetic films with BEPR = 1 % decreases with increasing film thickness. This suggests that the ferromagnetic behavior is due to defects created during the growth process and by the incorporation of Cu into GaN for BEPR ≤ 1 % leading to a long-range magnetic polarization.

# **OL06**

Magnetism and electronic transport of MnAs single nanostructures Federico Fernandez Baldis<sup>1</sup>, Marina Tortarolo<sup>2</sup>, Martin Sirena<sup>1</sup>, Laura Steren<sup>3</sup>, Victor Etgens<sup>2</sup> and Mahmoud Eddrief<sup>2</sup>

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The magneto-structural properties of MnAs thin films grown onto GaAs have recently drawn much attention due to its potential application in spintronics devices. The MnAs thin films exhibit a coexistence of a low-temperature ferromagnetic hexagonal phase (a) and a paramagnetic orthorhombic  $(\beta)$  one in a finite temperature range below the Curie temperature. When MnAs is grown onto GaAs[100], both phases are self-organized in a regular stripes array. This pattern is affected when the MnAs films are confined laterally to micrometric sizes and depends on the direction of the confinement. In this work, we present an investigation of the temperature dependence of the magnetic domain structure and magneto-transport properties of MnAs nanoribbons in order to determine the anisotropic response of the system, regarding both the magnetism and electronic transport. The confinement effects were examined at the low-temperature pure ferromagnetic phase and at the a-b phase-coexistence regime. Nanoribbons confined along the (I)  $\alpha$ -MnAs[1120] and (II)  $\alpha$ -MnAs[0001] directions were fabricated by e-beam lithography combined with optical lithography. The magnetic domain imaging was performed with a MFM, home-adapted for variable temperature measurements between 284K and 320K. The magneto-resistance and Hall effect measurements were made between 50K and 320K using a 9 Tesla superconducting magnet.

## **OL07**

The generated antiferromagnetic and ferromagnetic states in nanocrystalline Cu-Cu<sub>2</sub>O system and the consequences for spin chemistry

Anatoly Yermakov1\*, Michael Uimin1, Alexandre Korolyov1, Ilya Byzov1, Alexey Mysik<sup>1</sup>, Vladislav Maikov<sup>1</sup> and Mamoru Senna<sup>2</sup> <sup>1</sup> Institute of Metal Physics, Ural Branch of RAS, Russia <sup>2</sup> Faculty of Science and Technology, Keio University, Japan

Under investigations of nanocrystalline Cu-Cu<sub>2</sub>O oxides it has been established that they actually possess heterogeneous magnetic states in nanoscale range. In nano-structured Cu-Cu<sub>2</sub>O oxides, ferromagnetic (FE) and anti-ferromagnetic (AF) states are coexisting at above room temperature. The new idea based on the magnetic phase transformation from AF to FE state at an external field was suggested as a hypothesis to understand the variation of chemical rate transformations. The point is that a metamagnetic phase transition is accompanied by a variation of an electronic state of carrier from localized to delocalized one which can result in a variation of nanoreagents reactivity. In fact, the influence of magnetic field on the rate of chemical transformations, e.g. Cu ammonium complex formation using Cu-Cu-O nanomaterials, as magneto-controlled reagents, has been reliably observed. Thus, the nanocrystalline copper oxides having a liquid-like potential are a remarkable example of energy-wise degenerated different magnetic states that can be operated by fairly weak external magnetic fields to control a rate of chemical transformations and other properties. We think the principle we here propose could be extended to other transition metal oxide nanoparticles with a similar electronic states Thanks to RFBR Grants # 10-02-00323 and # 10-02-546 for support.

### **OL08**

### Unidirectional anisotropy observed in Fe film grown on GaAs at low temperature

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The magnetic anisotropy properties of Fe films have been investigated by planar Hall effect (PHE) measurements. A 4 nm thick Fe film was grown on a (001) GaAs substrate by MBE technique at room temperature. The field scans of the planar Hall resistance (PHR) at room temperature show an abrupt transition with hysteresis centered at zero field during the magnetization reversal, indicating the presence of symmetric magnetic easy axes in the film plane. However, the hysteresis in the PHR appeared to be shifted toward one field direction when the temperature decreased to 3 K. This asymmetry in the magnetization reversal process at 3 K can be understood by introducing an unidirectional anisotropy field  $H_{ud}$ . Interestingly, the direction of the  $H_{ud}$ directly follows the direction of the cooling field. This indicates that the direction of the H<sub>ud</sub> is not fixed but can be changed by the cooling field. For the case of the cooling field along the [-1-10] direction, we obtained the values of cubic, uniaxial and unidirectional anisotropy fields by analyzing angular dependence of PHR. The magnitude of H<sub>ud</sub> was one order of magnitude smaller than other anisotropy fields although it caused a significant asymmetry in the PHR hysteresis

# **OL09**

### Magnetic properties of hydrothermally synthesized ZnO nanostructures

H Hadiyawarman, Kadek J. Parwanta, Bowha Lee and Chunli Liu\* Department of Physics, Hankuk University of Foreign Studies, Korea

We report the study on the ferromagnetism in ZnO nanostructures fabricated by hydrothermal method at low temperature and with three different pH values. No posttreatment was applied to all samples. X-ray diffraction shows that the samples possess typical wurtzite structure and have no other impurity phases. Magnetization properties of ZnO samples were measured. With the defect analysis based on low temperature photoluminescence spectroscopy the effective defects contributed to the magnetic properties were investigated. The result suggests that oxygen defects and synthesis conditions play an important role in mediating the magnetic property in the undopped ZnO nanostructures

### Electrical, magnetic and magnetoimpedance studies of LSMO thin film prepared by sol gel method

Pawan Kumar and Ramanathan Mahendiran Physics, NUS, Singapore

In this work, we have prepared La<sub>0.7</sub>Sr<sub>0.3</sub>MnO<sub>3</sub> by sol-gel method using Si (100) substrate and characterized the sample by XRD, dc magnetization, dc magnetoresistance, and atomic force microscopy. The sample is found to be ferromagnetic below 367 K from the magnetic measurement. Impedance spectroscopy study has been reported on the thin film for the first time. We have studied the ac electrical impedance (Z) in LSMO thin film as function of Temperature (T) and magnetic field in response to radio frequency (f=0.1-5MHz) ac current flowing directly through the sample. DC resistivity and Z (T) shows transition around 365K which is correspond to Tc of LSMO thin film. In Z (T) peaks around Tc decreases in magnitude with applied field (1KOe), which is not affected much in DC resistivity. Magnetoimpedance  $\Delta Z/Z=-30\%$  is observed in MHz frequency range. It is suggested that observed magnetoimpedance results from suppression of the high frequency transverse permeability by an external field. Our study suggests that magnetoimpedance in manganite thin films might find new application in radio frequency regime.

References: (1). Gang Zou, Xian You, Pingsheng He, Materials Letters 62 (2008) 1785-1788.

# 0L11

### Differential conductance measurements in Ni nanoscale contact fabricated by electromigration

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A zero-bias anomaly in differential conductance has been reported in atomic-size contacts fabricated from the ferromagnets, Fe, Co and Ni, which is considered to come from the Kondo effect [1]. This is interesting because the Kondo effect usually requires two species of atom and is thought to be incompatible with ferromagnetic interaction. In order to clarify the origin of Kondo effect in atomic ferromagnetic contact, we have studied the size dependence of the zero-bais anomaly in Ni nanoscale contacts by mechanical controllable break junction technique and observe the anomaly in a large size where the ferromagnetism appears[2]. In this presentation, we study differential conductance in Ni nanoscale contacts with decreasing the size by electromigration. In large size contacts, zero-bias anomaly is not observed. When the conductance is smaller than 200G0 (G0=2e2/h), the anomaly is clearly observed.

[1]M. Reves-Calvo et al., Nature 458, 1150 (2009) [2]K. Ienaga et al., to be published.

# QL12

Magnetoresistance measurements in Pd atomic-scale contact at 4.2K Koichiro Ienaga<sup>1</sup>\*, Naoya Nakashima<sup>1</sup>, Yuji Inagaki<sup>1</sup>, Hiroyuki Tsujii<sup>2</sup> and Tatsuya Kawae<sup>1</sup>

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# QL13 Differential conductance measurements in Cu-Mn atomic-scale

contacts Koichiro Ienaga<sup>1</sup>, Naoya Nakashima<sup>1</sup>, Yuji Inagaki<sup>1</sup>, Hiroyuki Tsujii<sup>2</sup> and Tatsuya Kawae<sup>1</sup>\*

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## **QL14**

# Tuning the magnetic interaction in carbon nanotube/NiO nanocomposite system

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We have prepared a series of nanocomposite materials with multi-walled CNTs and NiO nanoparticles by varying the concentration of nickel nitrate solution. Transmission electron microscopy shows that the average size of the nanoparticles is about 10 nm and they are predominantly embedded either on the inner or the outer wall of the multi-walled CNTs forming a quasi one dimensional arrangement. It is also observed that the particles are well separated and density of the product NiO nanoparticles increases systematically with increasing concentration of the nickel nitrate solution. DC magnetization measurement shows the presence of thermomagnetic irreversibility in the temperature variation of the magnetization data. In contrary to the antiferromagnetic bulk NiO, field dependence of magnetization reveals the presence of strong ferromagnetic interaction along with prominent exchange bias effect. These materials also show magnetization relaxation which is an indication of non-ergodicity in the magnetic ground state. This system provides a rare opportunity to externally fabricate the degree of isolation among the particles and to tune the interparticle magnetic interactions as well for performing a systematic investigation on the evolution of the magnetic ground state with the variation of the interaction strength

## QL15

### Bias-voltage dependence of magnetotransport properties in codeposited Co-C granular thin films

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Granular systems consisting of ferromagnetic nanoparticles dispersed in a nonmagnetic matrix exhibit giant magnetoresistance (MR) due to spin dependent transport. Carbon-based nanocomposite thin films have large application potential as spin transport materials because long spin relaxation time is expected due to its weak spin-orbit interaction [1]. We have found that co-deposited Co-C films exhibit a large negative magnetoresistance (MR: 27.6% at 2 K) [2]. It is expected that the bias voltage dependence of MR includes significant information the mechanism of the MR. In this study, the bias-voltage dependence of magneto-transport properties for Co-C granular thin films have been investigated. Co-C granular thin films were fabricated by a co-sputtering technique. The current (I)- bias-voltage (V) characteristics and the MR behavior of the Co-C films were measured at temperatures of 2-30 K in a two-terminal geometry with the electrical current and magnetic fields (up to 70 kOe) in the in-plane direction. It was found that the magnitude of the MR ratio changed remarkably with bias voltage. A clear increase in the MR ratio with bias voltage was observed, and it reached 66.6% for V=25V

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### QL16

# Switchable voltage control of the magnetic anisotropy in heterostructured nanocomposites of CoFe/NiFe/PZT

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Heterostructures of nanostructured ferromagnetic (FM) and ferroelectric (FE) layers are of increasing interest due to the coupling between the magnetic moment and electric polarizations. In particular electric voltage, rather than the conventional magnetic field, can be directly used to control the magnetic properties of the magnetic layers. This magneticelectric coupling may open promising applications towards novel spin electronic devices with low power consumption and high-speed data access. In this article, we study on the crystallographic structure and magnetic properties of CoFe/NiFe/PZT heterostructures. In these nanocomposites, the magnetic behaviors, such as magnetization significantly changes in the presence of an applied voltage. The change also depends on the FM layer thickness, the magnetic field direction and it can reaches up to more than 100%. Furthermore we observed the switching in magnetization at suitable applied voltage. Besides, we present a phenomenological approach to investigate stress-induced magnetic anisotropy, which originated from magnetoelectric coupling between two FM and FE phases.

# QL17

## Origin of the ferromagnetism in scandium-doped ZnO thin films

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In recent years, there have been more startling discoveries in diluted magnetic semiconductors [1] achieving possible room-temperature ferromagnetism in undoped (ZnO [2]) and non-transition metaldoped (ZnO:C [3]) systems. ZnO-based materials serve as model systems for the dilute magnetic oxides and are among the most intensely studied systems so far [4]. Interestingly, Venkatesan et al [5] have shown that room temperature ferromagnetism is observed in Sc-doped ZnO films. We note that in the bulk phase Sc has different crystal structures and nonmagnetic properties. How different are their magnetic properties, electronic structure, and magnetic properties theoretical investigations of energetics, electronic structure, and magnetic properties of Sc-doped ZnO (11-20) thin films. In this work, the ZnO thin film was modeled by a (1×2) seven-layer slab supercell having a wurtzite structure along the [11-20] surface orientation. Each slab was separated from the other by a vacuum region of 10 A along the [11-20] direction. Even though Sc is not magnetic in its natural phase, we found that it is likely to order ferromagnetically in ZnO thin film.

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# QL18

Magnetic control of the hydrogen storage of hydrogen-injected ZnCoO Bum-su Kim<sup>1</sup>, Seunghun Lee<sup>1</sup> Jong Moon Shin<sup>2</sup>, Yong-chan Cho<sup>1</sup>, Yong Nam Cho<sup>1</sup>, Hee-ju Lee<sup>3</sup>, Chae Ryong Cho<sup>2</sup>, Hideomi Koinuma<sup>1</sup> and Se-young Jeong<sup>1</sup>\* <sup>1</sup> Cogno-Mechatronics Engineering, Pusan National University, Korea <sup>2</sup> Nano fusion technology, Pusan National University, Korea <sup>3</sup> Korea Atomic Energy Research Institute, Korea

Many researchers have tried to apply ZnO for hydrogen storage through additional doping and nanostructure which increase hydrogen capability[1]. It has been regarded for Zn-O bond center and oxygen vacancy as practicable position for hydrogen storage. Co-Co dimer in ZnCoO has been suggested as appropriate position for its reversibility, and hydrogen placed in Co-Co dimer mediates ferromagnetic spin ordering by Co-H-Co complex[2,3]. In our recent research, we have reported the correlation between hydrogen content and ferromagnetism in Co doped ZnO (ZnCoO) and suggested that ZnCoO can be exploited for hydrogen storage and the tools for the measurement of hydrogen[4]. In this study, ZnCoO powder was fabricated by sol-gel method and annealing process was performed at 300 °C and subsequently 800 °C for removing organic compound. Pt was coated on ZnCoO powder as hydrogen catalyst. The hydrogen absorption and charge in structure are characterized by the press composition and temperature (PCT) system and neutron diffraction, in-situ. We investigate the correlation between hydrogen contents and magnetism in ZnCoO and discuss is applicability for hydrogens. We also show the hydrogen contents in ZnCoO depending on the Pt layer and the spillover effect between Pt layer and ZnCoO nancerystal.

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## QL19

# Ferromagnetism in Co-doped TiO<sub>2</sub> films probed by low-energy muon spin rotation

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Ferromagnetic semiconductors have been widely studied recently for potential applications in spintronics. In this experiment, we probe the magnetism in Co-doped TiO<sub>2</sub> ferromagnetic semiconducting films using Low-energy muon spin rotation (LE- $\mu$ SR). This technique is very powerful to probe the magnetic properties of thin films and interfaces on a nanometer length scale. We investigated Co-doped (5%) TiO<sub>2-8</sub> films grown in different oxygen pressures. We find the internal magnetism in these films to be strongly dependent on the oxygen pressure during growth. The oxygen vacancies, modified by pressure, act as electron donors and alter the electron density. As the pressure increases, the magnetic field and depth dependence of the internal magnetism, and discuss its origin.

## QL20

# Structural and magnetic changes induced by high energy ball milling of $CdFe_2O_4$ oxide

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Bulk CdFe<sub>2</sub>O<sub>4</sub> ferrite nanoparticles have been synthesized by solid state reaction. The fine powders of CdFe<sub>2</sub>O<sub>4</sub> oxide were produced by high energy ball milling. The structural and magnetic changes induced by high energy ball milling are investigated. The oxides have been characterised by XRD, Mossbauer spectroscopy and VSM measurements. A general decrease in lattice constant with increasing milling time is observed. The Mossbauer spectra indicate paramagnetic spin state in all the samples. Magnetization increases with reduction in particle size. This can be explained by the redistribution of Cd ions in both tetrahedral (A) and octahedral (B) sites. The coercive field of the milled CdFe<sub>2</sub>O<sub>4</sub> oxide are highly sensitive to temperature compared to that of the bulk compound. An increase in coercive field for 50 hours milled oxide from about 9 Oe at 300 K to 520 Oe at 10 K has been observed. A smaller increase from 9 Oe at 300 K to 200 Oe at 10 K occurs in the coercive field of the bulk oxide. The variation of field cooled magnetization as a function of temperature is also presented. An anomalous variation of field cooled magnetization as a function of temperature has been observed.

# QL21

Room temperature ferromagnetism in  $Zn_{1-x}Ni_xS$  diluted magnetic semiconducting nanocrystalline thin films

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We have investigated the magnetic properties of Ni doped Zn<sub>Ly</sub>Ni<sub>y</sub>S diluted magnetic semiconducting nanocrystalline thin films for different doping concentrations (0.005  $\leq x \leq 0.2$ ) synthesized by electron beam evaporation technique. X-ray diffraction patterns confirm the existence of single phase nature in all the Ni doped Zn. Ni S samples with cubic zinc blend type structure. Evidence of nanocrystalline nature of the films was observed from the investigation of surface morphology using scanning electron microscopy and atomic force microscopy. Magnetic domains were observed by using magnetic force microscopy at room temperature indicating the existence of ferromagnetism over the film surface. The temperature and field dependent magnetization measurements by using superconducting quantum interface device showed ferromagnetic behavior between room and low temperatures (5 K) with a Tc at or above room temperature for the nanostructure samples with  $0.1 \le x \le 0.2$ . The saturation magnetization for Zn. Ni S system is found to increase with the donant concentration (x). The exchange interaction between local spin polarized electrons (Ni ions) and conductive electrons according to RKKY mechanism, rather than from the Ni oxide impurities, is proposed to be the possible mechanism for ferromagnetism.

## QM01

### Extraordinary hall measurements of Co/Ni multilayers

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Interest in Co/Ni multilayers is driven by their unusual magnetic properties. They exhibit perpendicular magnetic anisotropy, relatively low magnetic damping and unusual magneto-transport characteristics making them well suited for spintronic devices. To understand their magnetic and magneto-transport properties we measured the temperature dependence of the Hall effect for a series of sputtered Ni/Co films: SiO2/Ta(40A)/Pd(40A)/[Co(2A)/Ni(Y)]x10/Ta(40A)), with Ni thicknesses from Y=2-14 A. Because the signs, and the temperature dependence, of the extraordinary Hall effects (EHE) differ for Co and Ni, the behavior for the multilayers is complex. At room temperature the sign of the EHE changes as the Ni thickness increases from 4 to 5A. For Y=4A, the sign of the EHE is temperature dependent, going to zero near 220K and changing sign for lower temperatures. The magnitude of the EHE increases with increasing temperature for Y<5A, decreases for Y=5 and 6A and then increases for Y>6A. We used the EHE to measure the magnetization properties finding a decrease in the coercive field for increasing Y. For Y<4A, a complicated EHE response does not reflect the average magnetization. We discuss this research in term of competing EHE properties of Co and Ni. Work at UCSD supported by DOE-BES Award # DE-SC0003678

## QM02

### Magnetic Spin Structure of Fe<sub>50</sub>Pt<sub>50-x</sub>Rh<sub>x</sub> Films

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Ordered FePt alloys with L10 structure are known as materials with FM order and a high magnetic moment of Fe providing a large magnetization. The large atomic number of Pt on the other hand results in a high magnetic anisotropy. If grown in thin films, the high anisotropy often results in perpendicular magnetization which is the preferred orientation for current magnetic recording media. One way to control the magnetic properties in these materials is through the introduction of a third element into the crystal matrix e.g. Rh. When Rh is added to replace Pt in the equiatomic allov. new magnetic phases emerge. Bulk samples of Fe<sub>50</sub>Pt<sub>50</sub>, Rh. studied by magnetization measurements refer to an AF-FM transition with decreasing Rh content. An additional temperature driven AF-FM transition was observed in the Fe<sub>50</sub>Pt<sub>40</sub>Rh<sub>40</sub> film. It is also well known that the magnetic properties of a thin film can differ drastically from the bulk behavior due to dimensionality effects. Here we present neutron diffraction studies on the magnetic properties of different 200nm thick Fe<sub>50</sub>Pt<sub>50-x</sub>Rh<sub>x</sub> films in dependence of temperature and external magnetic fields. The polarized and unpolarized neutron diffraction data allowed us to develop a detailed model of the magnetic structure of the films

## QM03

**Cross over from anisotropic magnetoresistance to magnon magnetoresistance in PLD grown permalloy nanowires** Vineeth Mohanan Parakkat and Anil P. S Kumar *Physics, Indian Institute of Science, India* 

We are able to observe and distinguish Anisotropic-Magnetoresistance(AMR) and Magnon Magnetoresistance(MMR) [1] contributions clearly in Permalloy nanowires by varying their width thereby tuning their shape anisotropy. Nanowires of thickness 20nm and varying widths down to 160nm were prepared by e-beam lithography and Pulsed Laser Deposition. A linear non-saturating longitudinal MR observed in high field regime for NiFe nanowires could never be explained using AMR but only MMR can account for it. MMR follows as  $p=\alpha(MMs)B$ , where M is magnetization along the easy axis, Ms saturation magnetization, B magnetic field applied along the easy axis,  $\alpha=(dp/dB)$  is the slope of p(B) at magnetic saturation. A cross over in magnetoresistance at low field from AMR-dominant to MMR-dominant one was observed as wire width reduces since the AMR value is getting reduced from 1.2% for 650nm wide wire to 0.26% for 160nm, however the contributions from the electron-magnon interactions still persist. At 280nm itself the AMR contribution to the longitudinal MR has substantially reduced so that only magnon effect is visible. The MMR proves to be an excellent way of determining

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# QM04

### Investigation of magnetic anisotropy of ferromagnetic GaMnAs film by planar Hall effect

M/Ms, understanding the magnetization reversal mechanisms in nanostructures, domain

wall dynamics and also fundamental problems like electron magnon interactions.

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Magnetic anisotropy properties of the GaMnAs ferromagnetic film have been investigated by using planar Hall effect measurements. The field scan of the planar Hall resistance (PHR) showed a two-step switching behavior indicating a strong cubic anisotropy along the <100> directions. The difference in the behavior of the PHR for the two applied field directions between the [140] and the [110] was understood via well-known uniaxial anisotropy along the [110] direction. In addition to such known effects, we also found the presence of asymmetry for the field directions between the [010] and the [400] directions. This new asymmetry phenomenon was explained by introducing an additional uniaxial anisotropy field  $H_{U2}$  along the [100] and the [400] directions, which coincides with the two directions of cubic anisotropy. The values of the anisotropy fields, cubic  $(H_i)$ , first uniaxial  $(H_{U1})$ , and second uniaxial  $(H_{U2})$ , were obtained by analyzing the angle dependence of the PHR. Although the value of is small, its effect is clearly observed at high temperature above 25 K, where the transition of magnetization occurred before the field direction is reversed during the magnetization reversal process.

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Keywords: Ferromagnetic semiconductor, Magnetic anisotropy, Planar Hall effect

### **OM05**

# A crossover between magnetic vortex state and strip domains in electrodeposited nanogranular nickel films

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We have investigated the change of magnetization mode of nanogranular nickel films electrodeposited on silicon. The electrodeposition of nickel at a distinct potential starts with the formation of isolated nanogranular tilm. Assuming nickel nanogranules as oblate spheroids and calculating the critical size for single domain state, we find that the nickel nanogranules are magnetized under vortex ground state. Micromagnetic simulations calculated for this isolated nanogranules and MFM images confirm the magnetic vortex states for electrodeposited nickel nanogranules. Upon the nanogranules joining each other, coercivity increases due to the reduction of intergranular space and strengthening of magnetostatic interaction. Further growth leads to the formation of continuous nanogranular film which is associated with decrease in coercivity and rotation of the easy axis of magnetization to in- plane mode which is characterized by the domain wall displacement and strip-like multidomain pattern. The crossover between vortex ground state and multidomain magnetic pattern leads to the significant decrease in coercive force. Continuous nickel films electrodeposited on silicon have anisotropic magnetoresistance the value of which increases as the thickness and, appropriately, uniformity of film is increased.

# QM06

### Non-linear susceptibility of nanogranular FeAg films at the verge of superferromagnetism

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A detailed magnetization study of Fe<sub>4</sub>Ag<sub>unx</sub>, films, comprising nanoparticles (3 nm) of Fe embedded in a Ag matrix, has recently revealed [1] a crossover from superspin glass state to superferromagnetic state at x?30. This changeover is presumably caused by an interplay between the direct exchange, dipolar and RKKY interactions operating between the particle magnetic moments. In order to gain further insight into the nature of inter-particle interactions and their role, first five harmonics  $\chi_1 - \chi_5$  of the ac magnetic response were measured at several frequencies and magnetic fields over the temperature range 5-300 K on DC-sputtered Fe<sub>4</sub>Ag<sub>100x</sub> (x = 25, 35, below and above the percolation threshold, respectively) 100 nm thick films. Susceptibilities  $\chi_1 - \chi_5$  measured in the absence of the superposed dc magnetic field, Hdc, exhibit features that are characteristic of a spin glass and a ferromagnet in x = 25, 35. The ferromagnetic fingeprints in the temperature variations of the linnear and nonlinear susceptibilities become all the more apparent when Hdc is applied. The data are analyzed on the same lines as those in the previous works on magnetically disordered CMR manganites [2, 3] and LiNiO insulators [4].

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### **QM07**

# Dynamics of Ni-Fe elliptical dot arrays based on CPW-FMR measurements

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The dynamics of nanomagnets have been studied intensively for high-frequency applications in spintronic devices. The dynamics closely correlate with the damping constant  $\alpha$  which determines the strength of damping torque in nanomagnets.  $\alpha$  has been evaluated by exciting ferromagnetic resonance (FMR) in nanomagnets using microwave cavities and coplanar waveguides (CPWs). Advantages of CPW-FMR measurements are a wide range of radio frequencies and external field intensity and micron-scale samples. Herein to clarify the dynamics of nanomagnets, we employed FMR measurements with CPW to evaluate a in 50-nm-thick Ni-Fe elliptical dot arrays. Effective in-plane anisotropy field Hk,eff increases markedly as the dot size decreases, which may be attributed to enhanced magnetostatic energy. Eeffective saturation magnetization  $4\pi$ Ms,eff decreases slightly as the dot size decreases because the demagnetization field coefficient changes as the dot size decreases. Additionally, a increases monotonically from 0.01 to 0.02 as the dot size decreases in a 500 Oe external magnetic field applied in the longitudinal axial direction of the dots. Their values are higher than that of the Ni-Fe film ( $\alpha$ = 0.009) likely because the inhomogeneity of the demagnetization field distribution near dot edges. Therefore, these results propose that CPW-FMR measurements can evaluate the dynamics of nanomagnets in detail.

### **OM08**

### Structure and magnetic properties of SiO2(Co) granular film on GaAs substrate.

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The effect of giant injection magnetoresistance (GIMR) was recently observed in a granular SiO<sub>2</sub>/ (S4 - 75 at % Co) film on a semiconductor GaAs substrate in a temperature range near T = 300 K. The magnetoresistance reaches the value of 10000% in the magnetic field of 1.9 T and at the voltage of 90 V [1]. The grazing-incidence small-angle X-ray scattering (GISAXS) data indicated the layer of lower concentration of cobalt particles at the granular film/semiconducter interface if compared to the SiO<sub>4</sub>(Co) granular film itself. The thickness of this interface layer is about 70 A [2]. The polarized neutron reflectometry confirms the existence of the additional layer on the granular film/substrate interface. The temperature dependence of the magnetization M(T) shows the presence of two types of magnetic nanoparticles with different blocking temperatures and magnetization dynamics. The hysteresis curve of the magnetization demonstrates two-loop structure in region of the fields 0 - 80 mT and 80 - 200 mT. Data obtained by polarized neutron reflectometry, GISAXS and SQUID measurements are well correlated with each other and give the possibility to develop the GIMR effect theory paying attention to the interface fatures.

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# QM09

# Magnetic properties and structure of electrodeposited nickel on thin niobium film

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Cyclic voltammetry, chronopotentiometry, atomic force microscope, magneto-optical Kerr effect, and X-ray diffraction were used to investigate the magnetic property and structure of thin Ni films electrodepositing onto a 100 nm Nb. Constant-current method was used to grow the Ni films in this study. During the reduction process of Ni, a small amount of  $\beta$ -Ni(OH)<sub>2</sub> is formed; it becomes intense when the applied potential is more cathodic. The average size of the Ni grains grows and the coercivity of the Ni film decreases during the steady-state depositing process. The preferred orientations of the electrodeposited Ni film are fcc-(111) and fcc-(200), the easy axis of the magnetization of it is in the in-plane direction, however. This may due to the formation of the  $\beta$ -Ni(OH)<sub>2</sub>. The effects of current density one make any significant change in the magnetic property if the current density is higher than 5.2 mA/cm<sup>2</sup>. This may due to the intense hydrogen evolution reaction during the high current density deposition. These results could be the foundation of further research in the hydrogen energy area.

## QM10

# Characterization of epitaxial EuS(111) thin films on BaF2(111) and SrF2(111) substrates grown by molecular beam epitaxy

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A ferromagnetic insulating EuS thin film has received considerable attentions in fundamental studies of spin injection devices because of its very high spin filter efficiency[1]. It has been reported that EuS thin films on Si(100) and GaAs(100) substrates become insulating at a growth temperature(Ts) of 300°C and above, and the resistivity of that decreases with decreasing Ts due to the sulfur deficiency[2]. In this study, we have successfully grown EuS(111) epitaxial films on BaF2(111) and SrF2(111) substrates using molecular beam epitaxy at Ts=100~500°C. Pole figures of x-ray diffraction indicate a high degree of in-plane orientation, and all of the samples show very high resistivity. The arithmetic mean deviation of the surface for 10 nm-thick EuS films on BaF2(111) substrates measured by AFM are 0.122 and 0.092 nm, respectively. The Curie temperature of EuS films increases up to ~16 K with increasing Ts. In order to achieve an anti-parallel state between two ferromagnetic layers in spin injection devices, we try to manipulate the coercive force by Te doping. The obtained coercive force of Te-doped film (110 Qe) is large compared with that of undoped one (20 Oe). Optical properties such as a magnetic circular dichroism are also presented.

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### QM11

### The cation distribution and electrical hopping in $Fe_{3-x}Co_xO_4$ (0<x<1.65) ferrite films on MgO substrate grown by molecular beam epitaxy Der-sheng Lee<sup>1</sup> and Gung Chem<sup>2</sup>\*

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# QM12

# Magnetic circular dichroism in near-threshold two-photon photoemission

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Magnetic circular dichroism in near-threshold photoemission (TPMCD) results in large asymmetries similar to values observed in the regime of x-ray magnetic circular dichroism (XMCD). Hence, TPMCD facilitates new routes for microscopy and time-resolved experiments. Here we investigate the microscopic origin of the TPMCD for ultrathin Co films [1]. TPMCD is measured for perpendicularly magnetize Cs/Co/Pt(111) films with work function adjusted by Cs adsorption. For one-photon photoemission (1PPE) TPMCD at a fixed photon energy of hv = 3.06 eV and varying work function results in an almost constant value of 6.2%. TPMCD in two-photon photoemission (2PPE) with the same photon energy results in 8.4% demonstrating that for 2PPE the first excitation step is the dominant asymmetry-generating process. For 2PPE with reduced work function TPMCD yields an asymmetry of 17% in the photon energy range  $hv = 1.53 \cdot 1.66$  eV, thus revealing the importance of a real intermediate state. The experimental results are in reasonable agreement with ab initio calculations considering energy conservation and all directions of excitation instead of only those conserving parallel momentum. Funded by DFG (EL175/15, EL175/16), Carl-Zeiss-Stiftung and the MAINZ Graduate School of Excellence.

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## QM13

## Structural and magnetic properties of pseudocubic $BaFeO_{\rm 3}\text{-}d$ thin films

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Bulk BaFeO<sub>3</sub>-d is known to assume different crystal structures (hexagonal, cubic, tetragonal, orthorhombic, monoclinic) depending on the oxygen stoichiometry. For d < 0.35[1] the compound has a hexagonal structure and is antiferromagnetic below 160 K[2], while for higher values of d it assumes a distorted cubic structure and is a magnetically ordered insulator at room temperature that has been considered a potential multiferroic[3]. In order to explore their structural and magnetic properties, thin films of BaFeO<sub>3</sub>-d were grown by Pulsed-Injection Metal Organic Chemical Vapour Deposition (PI-MOCVD) on cubic SrIiO<sub>3</sub> single crystal substrates using different deposition conditions. The films were characterized using X-Ray diffraction, atomic force microscopy (AFM), SQUID magnetometry and Rutherford Backscatering Spectroscopy (RBS). By tuning the deposition conditions, namely the substrate temperature and the relative concentrations of Ba and Fe precursors, it was possible to obtain films that have a pseudocubic structure, are insulating and show antiferromagnetic properties at room temperature. The results obtained from the experimental characterization of the pseudocubic BaFeO<sub>3</sub>-d thin films are discussed and compared with the results obtained using density functional theory (DFT) using local spin density approximation (LSDA) and LSDA+U.

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# QM14

# Exchange anisotropy and antiferromagnetic coupling in NiFe/FeMn/ Co trilayers

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Ferromagnetic Resonance (FMR) study of the exchange anisotropy of magnetron sputtered  $Si_{tut}Cu_{2tum}N$ iFe<sub>sum</sub>FeMn<sub>tutum</sub>Co<sub>2</sub>Cu<sub>sum</sub> trilayers reveals the existence of antiferromagnetic coupling between NiFe and Co layers. The samples studied were produced in the presence of a magnetic field of 400 Oe to set the undirectional anisotropy and the thickness of the Co layer varied from 3 to 5 nm. The FMR experiments were done at room temperature using a commercial spectrometer operating with microwave frequencies of 9.79 and 34.1 GHz, swept static magnetic field and the usual detection techniques. The FMR spectra display two resonance modes due to Co and NiFe layers with very distinct effective magnetizations. The in-plane angular dependence of the resonance field of the Co layer displays the characteristic bell shape curve due to an exchange-biased bilayer with unidirectional anisotropy larger than the magnetocrystalline anisotropy. At the FeMn/Co interface of the sample with 5 nm thick Co layer, the exchange bias field is 160 Oe. At the NiFe/FeMn interface it is about only 10 Oe and the angular dependence of the resonance field exhibits an asymmetry due to the presence also of uniaxial anisotropy. This additional anisotropy results from an antiferromagnetic coupling between NiFe and Co layers.

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# QM15

# Parallel spin wave resonance in exchange-biased NiFe/FeMn/NiFe trilayers

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The Ferromagnetic Resonance (FMR) technique was used to study the magnetic properties of asymmetrical NiFe<sub>stum</sub>FeMn<sub>stum</sub>NiFe<sub>titum</sub> trilayers magnetron sputtered under working pressures of 2, 5 and 10 mTorr, in the presence of a 460 Oe magnetic field to set the unidirectional anisotropy. The FMR experiments were done at room temperature using a commercial spectrometer operating with microwave frequencies of 9.79 and 34.0 GHz, swept static magnetic field and the usual detection techniques. At a microwave frequency of 34.0 GHz, the parallel FMR spectra reveal that spin wave and uniform resonance modes are excited by the microwave field. The study of the in-plane angular dependences of the assonption fields reveals for both modes the effect of the unidirectional anisotropy axis, give for the spin wave and uniform modes of the sample produced under the pressure of 2 mTorn the anisotropy fields of 66 and 37 Oe, respectively. The parallel FMR results and spin wave resonance theory imply for this sample the exchange constant of 0.92 x 10-6 erg/cm. This result agrees with perpendicular FMR measurements at a microwave frequency of 9.79 GHz.

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### QM16

### Studies on local structures and magnetism at buried Fe/ Fe<sub>3</sub>O<sub>4</sub> interfaces using synchrotron-radiation Mössbauer spectroscopy Ko Mibu<sup>1\*</sup>, Hideto Yanagihara<sup>2</sup>, Takaya Mitsui<sup>3</sup>, Ryo Masuda<sup>3</sup>, Shiori Hori<sup>1</sup>, Atsushi Murata<sup>1</sup>, Masaaki Tanaka<sup>1</sup>, Kazuya Suzuki<sup>2</sup>, Eiji Kita<sup>2</sup> and Makoto Seto<sup>4</sup> <sup>1</sup> Graduate School of Engineering, Nagoya Institute of Technology, Japan

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Strong antiparallel magnetic coupling was found recently in Fe/ Fe<sub>3</sub>O<sub>4</sub> bilayers [1], and attention has been paid as a new and noble-metal-free "synthetic antiferromagnet" for magnetic recording devices. In order to clarify the mechanism of this interfacial magnetic coupling, it is important to elucidate the local crystallographic structure and magnetism at the buried interface experimentally. Mössbauer spectroscopy is a useful method to obtain such information through absorption spectra of gamma-rays (or X-rays) by the constituent Fe nuclei. In this presentation-based Mössbauer spectroscopy with a nuclear Bragg monochromator at BL11XU, SPring-8 [2]. The structures and magnetism at the buried Fe/ Fe<sub>3</sub>O<sub>4</sub> interface are discussed on the basis of the obtained spectra.

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### QM17

# Dependence of in-plane magnetic anisotropy of Au/Co/Au heterostructures on thickness of Co-component layer: An FMR study

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Recently, there was a lot of interest in magnetic tunnel junctions built from electrodes characterized by in-plane magnetic anisotropy. In the present study; ultrathin Mo/Au/ Co/Au heterostructures where the cobalt layer is of the wedge shape (i.e. with thickness gradually changing along the sample) were grown on a saphire substrate by molecular beam epitaxy. Room temperature ferromagnetic resonance (FMR) measurements were used to evaluate the resonance field as a function of (i) the in-plane and the out-of-plane orientation of external magnetic field and of (ii) the cobalt layer thickness. The experiments demonstrates an interesting effect of enhancement of the in-plane anisotropy for heterostructures with wedge-shaped Co layer. This finding provides a method to modify the magnetic anisotropy, which can be of a great importance for the application in magnetic tunneling junctions. A comparison of the obtained results with earlier data on Co anisotropy will be presented.

## QM18

Depth dependent chemical and magnetic information of CoFeB/ MgO multilayered thin films studied by x-ray and polarized neutron reflectometry

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Magnetic Tunnel Junctions (MTJs), consisting of a crystalline MgO barrier sandwiched by amorphous CoFeB layers, have received much attention owing to novel magnetic and transport properties such as perpendicular magnetic anisotropy and large tunneling magneto-resistance at room temperature [1-2]. With a combined study of x-ray (XRR) and polarized neutron reflectometry (PNR), depth-dependent chemical and magnetic profiles in Co<sub>40</sub>Fe<sub>40</sub>B<sub>3</sub>/MgO multilayered thin films were investigated. The annealing was done at 400°C for half an hour in an applied field of 4 kOe after sample growth. In addition to XRR and PNR, structure characterization was also conducted using hish-angle theta/2theta x-ray diffraction, and high resolution tunneling electron microscopy (HR-XTEM); the magnetic property was studied using vibrating sample magnetometer. The born diffusion and magnetic moment around both CoFeBMgO interfaces will be discussed in detail.

[1] S. Ikeda, K. Miura, H. Yamamoto, K. Mizumuma, H. D. Gan, M. Endo, S. Kanai, H. Hasegawa, J. Hayakawa, F. Matsukura, and H. Ohno, Nature Mater. 9, 721 (2010). [2] D. D. Djayaprawira, K. Tsunekawa, M. Nagai, H. Maehara, S. Yamagata, N. Watanabe, S. Yuasa, Y. Suzuki, and K. Ando, Appl. Phys. Lett. 86, 092502 (2005).

### QM19

# Depth-resolved rotational hysteresis of exchange-coupled Fe/Cr multilavers

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## QM20

### Magnetization reversal behavior in exchange-coupled NiFe/FeMn/ CoFe trilayers: vectorial MOKE & PNR study

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With a combination of vector magneto-optic Kerr effect (MOKE) magnetometry and polarized neutron reflectometry(PNR), the detailed magnetization reversal mechanism of the exchange-biased Py(30-nm)/ FeMn[t\_AF=0-30nm]/CoFe(30-nm) trilayers was studied.[1] We found that Py/FeMn(15-nm) and FeMn/ CoFe bilayers show completely different magnetization reversal modes, whereas they become very similar to each other in the corresponding Py/FeMn/CoFe trilayers. This is convincing evidence that the 15nm FeMn layer mediates the magnetization reversal behaviors of Py and CoFe layers through interlayer exchange bias coupling. Furthermore, magnetization reversal of Py and CoFe layers are decoupled for t\_ AF=30 nm, indicating that both Py and CoFe layers separated by an intermediate FeMn layer are exchangecoupled within 30 nm distance. This is in reasonable agreement with the theoretically predicted domain-wall width such as 28 nm for the polycrystalline FeMn/Co bilayers and 50 nm for the perfect Fe<sub>30</sub>Mn<sub>30</sub> crystal.

[1]Ki-Yeon Kim, Ji-Wan Kim, Hyeok-Cheol Choi, A. Teichert, Chun-Yeol You, Sungkyun Park, Sung Chul Shin, and Jeong-Soo Lee, Phys. Rev. B. 84,144410 (2011).

# QM21

### Uniaxial magnetic anisotropy of La<sub>0.7</sub>Sr<sub>0.3</sub>MnO<sub>3</sub> film grown on stepterrace surface SrTiO<sub>3</sub> (100) substrate

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In this work, we used the step-terrace surface STO (100) substrates to control the spin direction of the La<sub>0.7</sub>Sr<sub>0.3</sub>MnO<sub>3</sub> (LSMO) film in the atomic level. The shape magnetic anisotropy in the nano scale was induced by using step-terrace surface STO (100) substrates with small vicinal angle of 0.1 degree. Atomically flat STO substrate was prepared by annealing after etching by HCL (hydrochloric acid), and the surface morphology of the STO substrate with molecular step-and-terrace was observed by atomic force microscopy. Here, we simultaneously deposited two La<sub>0.7</sub>Sr<sub>0.3</sub>MnO<sub>3</sub> films on the processed and unprocessed STO substrates by using pulsed laser deposition. The direction dependence of magnetic hysteresis within the film plane was investigated by using SMOKE (Surface Magneto-Optical Kerr Effect) system. The magnetic hysteresis loops of the LSMO film on the suprocessed STO substrate showed weak direction dependence within the plane, while those on the step-terrace surface STO did a significant uniaxial magnetic anisotropy along the longitudinal direction of the step-terrace of the STO substrate.

## QM22

# The effect of compositional ratio on the magnetocaloric effect Ho-sup Kim<sup>1</sup>, Sang-soo Oh<sup>1</sup>, Kiran Shinde<sup>1</sup>, Seung-kyu Baik<sup>1</sup>, Kook-chae Chung<sup>2</sup> and

Bhavesh Bharat Sinha<sup>2</sup> <sup>1</sup> Korea Electrotechnology Research Institute, Korea <sup>2</sup> Korea Institute of Materials Science, Korea

Film type Ni-Mn-Ga alloys were prepared using co-evaporation method in order to investigate the effect of compositional ratio on magnetocaloric effect (MCE). In the deposition process on the metal substrate, Ni was evaporated by an induction heating, and Mn and Ga were evaporated by a resistive heating. The deposition rate of each element was measured by Quartz Crystal Microbalance and the rate was controlled up to 0.1 mm / hour to get a thick Ni-Mn-Ga film. Three crucibles of Ni, Mn, and Ga were arranged triangularly to form the composition gradient of Ni-Mn-Ga film on the surface of whole substrate. Long tape samples with different composition of Ni, Mn, and Ga along the length were prepared. The magnetic property of prepared sample a round transition temperature was continuously evaluated by a reel to reel Hall probe measurement system with a temperature controller. The Hall probe signal is related to the relative value of MCE so that the optimal composition ratio could be obtained by analyzing it. The section which exhibited the highest value in Hall probe measurement was cut and analyzed by XRD and followed by MCE measurement using PPMS equipment. Detailed experimental results will be reported in the presentation.

1. V. Recarte et al., Appl. Phys. Lett., 95, 141908, 2009. 2. Yuepeng Zhang et al., Smart Mater. Struct., 18, 025019, 2009. 3. V.A Chernenko et al., J. Phys.: Condens. Matter, 17,5215-5224, 2005.

### QM23

# Magnetocaloric effect in Ni doped Zn ferrite nanoparticles grown by the combustion method

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We studied the magnetocaloric effect (MCE) in two different compositions of Ni-Zn ferrite nanoparticles (Ni,Zn<sub>1-x</sub>Fe<sub>2</sub>O<sub>4</sub>, x=0.2 and 0.3) which were prepared by the combustion method. From the temperature and the magnetic field dependent magnetization curves, the magnetic entropy changes ( $\Delta S$ ) and the adiabatic temperature changes ( $\Delta T$ ) were evaluated. Magnetocaloric effect in these nanoparticles is compared to that in bulk system with the same compositions to study the effects of grain size. Although the values of  $\Delta S$  and  $\Delta T$  at 2 T magnetic fields in these nanoparticles are lower than the values in the bulk one, wider effect transition windows were observed in nanoparticle systems. As the domain size of Ni-Zn ferrite nanoparticles affects the temperature and magnetic field dependence of magnetization, the maximum  $\Delta S$  and  $\Delta T$  are also influenced by the grain size around the magnetic transition temperature (TC). Therefore, the field dependence of  $\Delta T$  at TC, which can be analyzed by the mean field approximation denoted by the exponential p, was found to be different from the typical one.

# QM24

### Interlayer interaction In ReCoPO (La, Nd and Sm): 31P NMR Study

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31P NMR result is presented in Co based non-superconducting members of pnictides family revealing the importance of the interlayer interaction, a key factor that determines the ground state. In ReCoPO (Re = rare earth) series, La and Ce show only PM  $\rightarrow$  FM transition, whereas Nd and Sm show successive PM  $\rightarrow$  FM  $\rightarrow$  AFM transition with varying degree of TC and TN. 31P NMR spectrum in polycrystalline La, Nd and Sm CoPO respectively corresponds to a powder pattern for a spin 1/2 nucleus with axially symmetric local magnetic field showing varying degree of shift parameters. In particular, Nd and Sm compound show large temperature dependent shift anisotropy with similar magnitude of Kiso and Kax. However, relaxation rate (1/T1T) in Sm is atleast two orders of magnitude higher than that observed in other two compounds. Our earlier 31P and 139La NMR [1] measurements in LaCoPO (quasi 2D Fermi surface) reveal dominant 2D spin fluctuation of 3d electrons in PM state with non negligible 3D part and in the FM state spin fluctuation is 3D in nature. Whereas, in SmCoPO the present result is able to distinguish the FM in plane spin-correlations from the AFM type out of plane spin-correlations.

1. M. Majumder, K. Ghoshray, A. Ghoshray, B. Bandyopadhyay, M. Ghosh , Phys. Rev. B 82, 054422 (2010).

## QM25

Jahn-teller distortion and enhancement of curie temperature of Mn<sub>3-x</sub>Ni<sub>x</sub>O<sub>4</sub> films on MgO (001) by molecular beam epitaxy

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We report on x-ray diffraction, scanning electron micrograph, and magnetization measurements vs. external field and temperature of  $Mn_{3x}Ni_4O_4$  (x=0, 0.25, 0.5, 0.75, 1) films grown on MgO (001) and (011) substrates by molecular beam epitaxy. Magnetization vs. temperature measurements reveal that Curie temperature increases from 43 K to 103 K as the Ni composition increasing. This increase of Tc is qualitatively understood as a result of the decrease of the exchange coupling of the cations between B site and B site (JBB). The Jahn-Teller distortion also reduces as the Ni composition increasing and the cubic structure is estimated to be recovered as x approaches ~ 1.2. The cation distribution in the present films is different from the bulk and can be described as  $[Mn^2+1-XMn^2+2x/3NI^2+2x/3Mn^2+2-2x/3]BO_n$ .

## QM26

Anomalous switching of in-plane magnetic anisotropy of Fe and Co thin films grown on curved Pt(001) surface

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We investigated the step-induced magnetic anisotropy of Fe and Co thin films grown on the curved Pt(001) surface using Surface Magneto-Optical Kerr Effect(SMOKE) and Scanning Tunneling Microscopy(STM). The continuous variation of step density of the curved Pt(001 surface are confirmed by STM experiment. From SMOKE measurements on both of Fe and Co thin films, we observed the switching of the magnetic easy axis from the parallel-to-step direction to the step-normal direction with increase of the step-density. We attributed The origin of this anomalous transition to the abrupt change in the step structure at high-step density position, which was revealed by STM measurements.

## QM27

Magnetic-field and temperature-dependent relaxation in ferrofluids characterized with a high-Tc SQUID-based nuclear magnetic resonance spectrometer

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In this work the magnetic-field and temperature dependent spin-lattice relaxation rate, 1/T1, and spin-spin relaxation rate, 1/T2, of ferrofluids are characterized using a high-Tc SQUID-based magnetometer as the detector. The ferrofluids are consisted of detrain-coated Fe<sub>3</sub>O<sub>4</sub> dispersed in phosphate buffer saline solution. It was found that 1/ T1 measured in a magnetic field of 56.8 mT is significantly higher than 1/T2 measured in a magnetic field of 102  $\mu$ T. We attribute this to the magnetic-field gradients from magnetic nanoparticles that accelerate the T1-relaxation more in a high strength of magnetic fields than they are in a low strength of magnetic fields. Furthermore, 1/T1 and 1/T2 decrease when temperature increases. We attribute this to the improved fieldhomogeneity seen by protons' spins under the influence of enhanced Brownian motion of magnetic nanoparticles at high temperatures than they are at low temperatures. Characterizing the relaxation rates of ferrofluids in low concentration regimes will be helpful for their future applications such as the low field MR imaging.

# QM28

# Influence of metal precursor on synthesis and magnetic properties of nanocrystalline strontium hexaferrite thin films

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The effects of starting metal salts on the structure and magnetic properties of strontium hexaferrite thin films synthesized by sol gel process have been investigated. Fourier transform infrared and thermal analyses were conducted to determine the chelated species and thermal decomposition of the gels, respectively. X-ray diffraction, scanning electron microscopy and vibrating sample magnetometery techniques were also applied to evaluate the microstructure, composition, crystallite size and magnetic properties of the thin films. The results showed the film prepared from nitrate precursor offered a single phase strontium hexaferrite with nanocrystalline structure and higher saturation magnetization after calcination. However, the films obtained from hydroxide and chloride precursors had higher coercivity due to smaller crystallite size.

## QM29

### Planar Hall effects measurements of sensitive magnetization response in epitaxial Fe thin films

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The Fe thin film is a promising candidate for future magnetic microwave devices, owing to the high ferromagnetic resonance frequency with an order of 10 GHz [1-2]. A high sensitive magnetization response is crucial for the low power device operation. In the present study, we prepared Fe(100) thin films deposited on MgO(100) substrates using magneton sputtering system and investigate the magnetization process using the planar Hall effect (PHE). The PHE measurements were carried out with several Fe thicknesses and temperatures. In the PHE measurements, alternating external magnetic fields along vertical easy axis were applied with various amplitudes and the sense current was applied along the orizontal easy axis. With increase of alternating magnetic field, the PHE voltage increased due to the disposition of Fe magnetization from easy axis. When the alternating magnetic field exceed a threshold value, the PHE voltage decreased rapidly. The threshold value depend on the thickness of Fe film and the PHE voltage decreased with decreasing temperature. These results suggest that the magnetization notates exceeding the hard axis direction and reversible susceptibility can be realized. The stochastic aspect of the magnetization process also appears at 10 K.

[1] E. Schloeman et al., J. Appl. Phys., 63, 3140 (1988). [2] N. Cramer et al., J. Appl. Phys., 87, 6911 (2000).

# QN01

### Antiferromagnetic order and domains in Sr<sub>3</sub>Ir<sub>2</sub>O<sub>7</sub>

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Resonant scattering at the Ir L2, 3 edge has been used to determine that  $Sr_3Ir_2O_7$  is a long-range ordered antiferromagnet below TN~230 K with an ordering wavevector,  $q = (\frac{1}{2}, \frac{1}{2}, 0)$ . The energy resonance at the L3 edge was found to be a factor of -30 times larger than that at the L2. This remarkable effect has been seen in the single layer compound  $Sr_2IrO_4$  and has been linked to the observation of a Jeff= 1/2 spin-orbit insulator. Our result shows that despite the modified electronic structure of the bilayer compound, caused by the larger bandwidth, the effect of strong spin-orbit coupling on the resonant magnetic scattering persists. Using the programme SARAh, we have determined that the magnetic order consists of two domains with propagation vectors  $k1 = (\frac{1}{2}, \frac{1}{2}, 0)$  and  $k2 = (\frac{1}{2}, \frac{1}{2}, 0)$ , respectively. A raster measurement of a focussed

the effinited that the magnetic order consists of two domains with propagation vectors  $k1 = (\frac{1}{2}, \frac{1}{2}, 0)$  and  $k2 = (\frac{1}{2}, \frac{1}{2}, 0)$ , respectively. A raster measurement of a focussed x-ray beam across the surface of the sample yielded images of domains of the order of 100 µm size, with odd and even 1 components, respectively. Fully relativistic, monoelectronic calculations (FDMNES) have been employed to calculate the relative intensities of the L2, 3 edge resonances, comparing the effects of including spin-orbit coupling and the Hubbard, U, term.

### QN02

# High frequency permeability of Fe-Co and Co granular composite materials

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Relative complex permeability and permittivity spectra of  $Fe_{50}Co_{50}$  (Permendur) and Co granular composite materials have been studied in the microwave frequency range up to 18 GHz. Particles were heat-treated in air using an electric furnace in order to suppress the eddy current effect in the composite structure containing percolated particles. Composite materials were prepared by mixing these particles with Polyphenylene sulfide resin. The 78 vol.%  $Fe_{50}Co_{50}$  granular composite that the larger relative permeability value than that of the 79 vol.% Co composite at 100 MHz. The negative frequency dispersion of permeability was observed above 7 and 8 GHz for  $Fe_{50}Co_{50}$  and Co composites, respectively. The dc magnetic field effect on the permeability spectra revealed that the permeability dispersion of Co composites in the microwave range is mainly attributed to the gyromagnetic resonance. The relative permittivity spectra of 78 vol.%  $Fe_{50}Co_{50}$  and 79 vol.% Co composites show the insulating property up to several GHz; the permittivity value increases with increasing particle content in both composites.

### QN03

# Magnetic carbonyl iron suspension with nanoclay additive and its magnetorheological properties

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Soft magnetic carbonyl iron (CI) based magnetorheological (MR) fluids with different loadings of nanoclay were prepared, in which the MR fluid is a complex colloidal suspension consisting of magnetic particles dispersed in a liquid, showing its rapid, reversible and tunable change between a liquid-like and solid-like state with an applied external magnetic field. The MR characteristics were measured via rotational tests, in which the flow curves exhibited a non-Newtonian behavior for all investigated samples under applied magnetic fields. Flow curves showed not only the dynamic yield stress change measured as a function of magnetic field strength using a power law fit but also the existence of a solid-like character. Sedimentation of the MR fluid with and without an additive was also examined.

# QN04

### Hysteresis analysis: A study on the demanetization by using M-B preisach model for improved stability of numerical analysis Hvuk Won and Gwan Soo Park\*

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It is necessary to describe the hysteresis characteristic of magnetic material precisely for the analysis of design of system with ferromagnetic materials. Although Preisach model is regarded as the most accurate method to describe the hysteresis characteristics, it is not applied widely to the real systems because of some difficulties. The conventional Preisach model shows the numerical instabilities during the iterative computations because the density distribution obtained from the sets of M-H curves are strongly localized. To remove such numerical instabilities, if we use a M-B variable instead of M-H variable would be better the results. In this paper, we suggest method for a M-H variable to change a M-B variable, and also, from the computed results of used normal Preisach, M-H cure method and M-B curve method, we shows better than numerical stabilities by using two dimensional finite element method.

### QN05

A study on the design of transmitting coils and receiving coils on active magnetic snesor using finite element method Hye Sun Ju and Gwan Soo Park\* *Pusan National University, Korea* 

Proximity magnetic sensor is able to detect the object target accurately in close range and it has been widely used in the underwater guided weapon system because there is no countermeasures from the target. In order to increase the damage of target by shock wave due to explosion of the underwater guided weapon system, the maximum detection range of the proximity magnetic sensor needs to be increased. In this paper, we describe the techniques of the optimum transmitting and receiving coils design using the Finite Element Method for the output power enhancement of the transmitter and the sensitivity improvement of the receiver. Finally, the proposed design techniques of the transmitter and the receiver were verified using a experimental setup and a prototype.

## QN06

# What can we learn from isothermal remanent magnetization curves on diluted nanoparticle assemblies?

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Isothermal remanent magnetization (IRM) and direct-current demagnetization (DcD) curves are mainly used to characterize interactions in a granular magnetic media and to determine the switching field distribution. These curves are only sensitive to irreversible changes, i.e. to the particle magnetization switching for samples made of clusters embedded in a matrix. Moreover, at remanence there is no unsuitable contribution (superparamagnetic, diamagnetic or paramagnetic) to the signal. This makes IRM/DcD curves of great interest in the investigation of the intrinsic properties of nanoparticles. The underlying process, i.e. macrospin switching by applying a field, is complementary to the thermal switching observed with ZFC/FC curves where the particle size distribution has a critical influence. Interestingly, IRM curves have a high sensitivity to other characteristics such as an anisotropy constant distribution, non-uniaxial terms of the anisotropy, and of course inter-particle interactions. We show how IRM/DcD curves can be modelled for an assembly of randomly oriented and non-interacting macrospins, thus allowing efficient and realistic simulations of experimental curves (taking into account the size distribution, temperature, etc). This modelling is then confronted to a series of measurements on diluted Co nanoparticle samples prepared by low-energy cluster-beam deposition. Experiment/ theory comparison can provide new insights on their magnetic propertie

# QN07

# On real-time hysteresis compensation for magentostirictive sound transducers

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Magnetostrictive actuators show inherent hysteresis, and it hinders wide applications. To solve this problem, in the literature, 'Jiles-Atherton model' and 'Preisach model' have been used mostly. But the methods require huge CPU time, or need complex mathematics. Thus, for real-time, ability of hysteresis compensation using capacitor is proposed by park et al [1]. In this paper, class-D amplifier containing capacitor as hysteresis compensator is proposed. The Class-D amplifier consists of comparator, current supplier and R-C circuit. Hysteresis compensator is implemented by a R-C circuit, because the output voltage of capacitor has similarity to a hysteresis compensated input and the process of development is as follows: 1. Generation of Hysteresis compensated input by Preisach model: 2. Comparison between hysteresis compensated input and R-C circuit model by non-linear regression: 3. Application found capacitance to the class-D amplifier. Experimental setup is consisted of a capacitance displacement sensor, a digital oscilloscope and a function generator. Hysteresis loss observed about 19 percent before application of hysteresis compensator, and hysteresis loss observed less than 13 percent after application of hysteresis compensator at 10 Hz by using the 3 uF of capacitor.

 Ki-Hyun Ji, Hae-Jung Park, and Young-Woo Park, "Semiempirical Approach to Compensate Hysteresis in Magnetostrictive Actuator", IEEE International Magnetics Conference 2011, 2011

## **QN08**

# The investigation on structural and magnetic properties of Ni-Cu-Zn-Mn-Mg-Li ferrites

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In this work, addition of seven raw materials inclusive of nickel oxide, copper oxide, zinc oxide, manganese oxide, magnesium oxide, lithium oxide, and iron oxide to form a new multi-component ferrite system. The Ni-Cu-Zn-Mn-Mg-Li substituted ferrites were prepared by conventional solid-state-reaction process and sintered at 1100°C for 2 h in air atmosphere. The influences of Ni-Cu-Zn-Mn-Mg-Li substitution on structural and magnetic properties were investigated by scanning electron microscopy, RF impedance/material analyzer, B-H curve loop tracer, and vibrating sample magnetometer. The experimental results show that the additives with low melting point form precipitates at the grain boundaries.[1] The average grain size of sintered ferrites decreases with a decrease of Ni/Mn ratio. The sintered ferrites, Ni/Mn=0.79, showed largest initial permeability (µi=85), largest saturation magnetism (Bs=3930 gauss), and better saturation magnetization (Ms=74.91 emug-1).

[1] Guiqing Wang, Xiufang Bian, Weimin Wang, Junyan Zhang, Mater. Lett., 57 (2003) 4083-4087.

## **QN09**

Magnetorheology of xanthan gum coated soft magnetic carbonyl iron microspheres and their polishing characteristics

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Magnetorheological fluids, complex colloidal suspensions consisting of magnetic particles dispersed in a liquid, showing their rapid, reversible change between a liquid-like and solidlike state with an applied magnetic field. Among these applications, MR polishing arrests a lot of attention due to its smart controlling polishing characteristics for dedicated MEMS applications. In order to improve polishing characteristics of MR fluids, in this study, carbonyl iron (CI) microspheres coated with xanthan gum (XG) were fabricated via solvent casting method. The morphologies and densities of both pure CI and CI/XG particles were characterized via SEM and pycnometer, respectively. In addition, rheological characteristics of the MR fluids under applied magnetic field strengths were examined via a rotational rheometer, in which shear viscosity, storage modulus, normal force the variation in MR performance under applied magnetic fields. Especially, the normal force was measured for different gap distances to simulate the work condition of MR polishing. MR polishing characteristics were conducted using an MR polishing machine to investigate the surface roughness from the MR polishing with nano ceria slurry abrasives. A series of experiments were performed for fused silica glass using prepared slurry by changing the processing parameters such as a rotating wheel speed and polishing time.

# **QN10**

# Cluster-glass-like magnetic state in rare-earth intermetallic compound $\mathrm{Tb}_{s}\mathrm{Pd}_{2}$

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The rare-earth intermetallic Tb<sub>2</sub>Pd<sub>2</sub> compound crystallizes in a cubic crystal structure with site disorder and magnetic frustrations due to triangle-like arrangement of Tb atoms [1]. In this work the DC and AC magnetic susceptibility measurements together with powder neutron diffraction have been performed in order to study the Tb<sub>2</sub>Pd<sub>2</sub> magnetic state. The susceptibility measurements on a polycrystalline Tb<sub>3</sub>Pd<sub>2</sub>sample have revealed a complex highly inteversible behaviour below 58 K, which can be associated with the existence of a spin-glass-like magnetic state. According to our neutron diffraction data the Tb<sub>2</sub>Pd<sub>2</sub> compound doesn't exhibit a long-range magnetic down to 5 K. However below and above the freezing temperature TG ? 58 K, the appearance of the broad diffuse maxima of magnetic origin has been detected in the low angle region. The second anomaly of the susceptibility was observed at TB ? 25 K, which can be ascribed to the blocking process of uncompensated magnetic moments of antiferromagnetically shortrange correlated regions. The results obtained allow us to suggest that the Tb<sub>2</sub>Pd<sub>2</sub> compound exhibits a complex non-ergodic magnetic state of a cluster-glass type associated with firstrations of the KKY4-type exchange interaction. This work was supported by the program of RAS (Projects RAS UB 11-2-HIT-265).

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# QN11

### Microwave absorption properties of polymer composites with amorphous Fe-B and Ni-Zn-Co ferrite nanoparticles

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Recently, many communication devices use GHz-range microwave. It causes serious issue, such as electromagnetic interference and multipath interference. To prevent these problem, microwave absorbers (MA) consisting of polymer composites of magnetic powders are paid attention. With the current trend being to miniaturize devices, thinner MA are now required. We reported that polymer composites with amorphous Fe-B particles (AFBP) and Ni-Zn ferrite nanoparticles showed high permeability of  $\mu$ 'r = 9.5 at 1.0 GHz and good microwave absorption properties at 0.6-1.1 GHz and resulted in thinner MA than reported MA [1]. However, they cannot used as MA over 1.1 GHz, because of their low resonant frequency (fr = 1.2 GHz). In this study, Ni-Zn-Co ferrite nanoparticles (NZCFN) were mixed with AFBP for increasing fr in polymer composites. XRD analyses show that NZCFN exhibit typical spinel structure. The fr of NZCFN increases from 0.7 to 3.6 GHz with increasing Co contents. It indicates that the magnetocrystalline anisotropy of NZCFN is increased by Co addition. The fr of polymer composites with AFBP and NZCFN can be increased to 2.6 GHz. They exhibites good microwave absorption properties at 0.8-1.4 GHz with thicknesses of 2.4-3.9 mm and resulted in thin MA at 1-1.4 GHz.

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## QN12

# Magnetostrictive properties of Mn substituted sintered cobalt ferrite derived from nanocrystalline materials

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Magnetostrictive smart materials are of great importance for sensor applications. Cobalt ferrite is well known for its highest magnetostrictive strain at room temperature among the oxide based magnetostrictive materials. The substitution of manganese for iron in cobalt ferrite, CoFe<sub>2</sub>O<sub>4</sub>, increases the strain derivative even though the magnetostriction value is decreased for the material synthesized through the conventional ceramic route [1,2]. In the present study we have investigated the influence of the initial particle size of the powders of manganese substituted cobalt ferrite on the magnetostrictive coefficient of the sintered ferrite samples with compositions 0≤x≤0.3 in CoFe<sub>2</sub>. "Mn<sub>x</sub>O<sub>4</sub>. Nanoparticles of CoFe<sub>2-x</sub>Mn<sub>x</sub>O<sub>4</sub> of size less than 10 nm were synthesized by an autocombustion method and the resulting powders were sintered in air at 1300° C. Magnetic characteristics of the nanocrystalline powders and sintered products have been studied. Magnetostriction measurements were performed on the sintered samples to determine the influence of initial particle size of the starting powders. Higher magnetostriction at low fields is obtained along with enhanced strain derivative for the sintered Mn substituted cobalt ferrite derived from nanocrystalline materials compared to the bulk counterparts

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# QN13

### Magnetic Properties of Co<sub>2</sub>Z Ferrite Densified through High BET Powders

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Co<sub>2</sub>Z is a planar hexagonal ferrite that exhibits high permeability (µi) and high resonance frequency (fr) up to the GHz region. Z-type Ba<sub>3</sub>Co<sub>2</sub>Fe<sub>24</sub>O<sub>41</sub> hexagonal ferrite was prepared by two-step calcinations [1]. Single phase Co<sub>2</sub>Z ferrite powders were possible after post-calcination at 1300°C. Various sintering aids improved densification of ceramic bodies, but lowered magnetic permeability due to non-magnetic characteristics of sintering additives. In order to increase permeability, Co<sub>2</sub>Z specimen was sintered without sintering aids. The densities of sintered bodies increased with increasing specific surface area (BET) of calcined powders, and Co<sub>2</sub>Z with BET 7.80m2/g led to significantly high densification(5.23g/cm3, 97% theoretical density). Permeability of Co<sub>2</sub>Z gradually increased with increasing density, whereas their resonance frequencies decreased with increasing density, where the product of µi x fr was estimated to be ~10.3GHz. It was thus confirmed that the calcined powders with high BET could enhance sintering kinetics of Co<sub>2</sub>Z hexagonal ferrites without sintering additives and significantly increase magnetic permeability.

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### QN14

### Magnetic properties of Cr<sup>3+</sup> substituted Mg-Cd ferrites

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Magnetization and susceptibility of polycrystalline ferrites with general formula  $Cd_x$   $Mg_{1,x}$   $Fe_{2,y}$   $Cr_y$   $O_4$  (x = 0, 0.2, 0.4, 0.6, 0.8, 1.0; y = 0, 0.05 and 0.10) were studied. Study of saturation magnetization reveals that the Neels two sublattice model exists up to x = 0.4, for y = 0, 0.05 and 0.1 and a three sublattice model (YK-model) is predominant for x > 0.4, y = 0, 0.05 and 0.10. The saturation magnetization and magnetize moment were found to decrease with the increase of  $Cr^{3+}$  contents, which is attributed to the dilution of B-B site interaction. Temperature dependence normalized AC susceptibility study reveals that  $MgFe_2O_4$  exibits single domain(SD) structure. On substitution of  $Cd^{2+}$ , single to multi domain transition takes place. Curie temperature decreases with  $Cd^{2+}$ , was attributed to decrease in A-B intraction. On substitution of  $r^{2+}$ , peak obtained in  $MgFe_2O_4$  was supressed which is attributed to decrease in grain size and further decrease

### QN15

## Docking Speaker Based on Magnetostrictive Sound Transducer

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All docking speakers available in the market are equipped with iDevice cradle and with cone-type sound unit. This paper presents a novel docking speaker without them. It is based on the magnetostrictive sound transducer (MST), which transforms any flat surface into speaker. The docking speaker consists of one MST, two circuit boards, and aluminum housing. One circuit board contains a Bluetooth module, a power amplifier module, Aux, and a recharge adapter. The other board is called a control board, which ATmega128 is used to control a digital clock, a LCD module, four control switches, and a power amplifier module. The reluctance concept is utilized to the magnetic circuit for the MST. Allowable current and number of coil turns are determined through this procedure. The result is compared with the simulation result. Then, circuit board and housing are designed. The assembled docking speaker is subjected to the experiments. For the experiments, the MST is placed and bolted between the wood and glass panels. An accelerometer sensor, ADXL001, is firstly located beneath the center of the MST. Then, the location is placed awy by 50 mm from the center of the MST. It is proved that the developed docking speaker works well.

# QN16

### Characteristics of composite materials Ba<sub>0.5</sub>Sr<sub>0.5</sub>Fe<sub>11.7</sub>Mn<sub>0.15</sub>Ti<sub>0.15</sub>O<sub>19</sub>/ La<sub>0.7</sub>Ba<sub>0.3</sub>MnO<sub>3</sub> as a microwave absorber V. Vekky R. Repi\* and Azwar Manaf Graduate Program of Materials Science Study, Universitas Indonesia, Indonesia



## QN17

### The effect of shape anisotropy to acoustical performance of magnetostrictive sound transducer Hae Jung Park, Young Woo Park\* and Ok Kyun Oh

Department of Mechatronics Engineering, Chungnam National Univers, Korea

This paper presents the experimental approach to investigate the effect of the shape anisotropy of magnetostrictive material (MM) to acoustical performance of magnetostrictive sound transducer. The assumption is that the volume of MM is kept constant. In the experiments, four MM samples are used with a volume of 295 mm^3 and input currents are swept from 20 Hz to 20 kHz. The diameters of MM samples are 4, 5, 6 and 10 mm and the lengths are 23.5, 15, 10.5 and 3.5 mm, respectively. The displacements are measured from 2.1 to 0.8 um at 20 Hz. Forces are measured from 81 to 120 N at 20 Hz. Among the samples, the MM sample of 23.5mm length shows the largest displacement and force. SPL is obtained from forces of MM and a area of vibration panel. The largest SPL is observed at 20 Hz and the SPL is measured to 115dB. Also, the SPL of the samples show the level from 103 dB at 1 kHz to 81 dB at 20 kHz. The displacement and force are measured from 0.36 um and 29N at 1 kHz to 15.3 nm and 2.5 N at 20 kHz.

### QN18

### Size effects on magnetic properties of nanocrystalline Sr<sub>2</sub>CuCo<sub>2</sub>Fe<sub>24</sub>O<sub>41</sub> prepared by Co-precipitation method K Praveena and K Sadhana

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The nanocrystalline z-type hexaferrite powders were synthesized by Co-precipitation method using aqueous solutions of strontium, copper, cobalt iron and nitrates. The as synthesized powders were characterized using XRD, TEM and TG/DTA. The average particle size of the powders was found to be 48nm. The effect of calcinations temperature on crystallinity, morphology and magnetic properties have been investigated. The sintered samples were characterized using XRD, FE-SEM and EDS. The hexagonal ferrite formation was observed around 900°C which is confirmed with the TG/DTA curves. The plate shape grains were observed in FE-SEM at 950°C. The grain sizes of sintered hexaferrites are in the range of 220nm to 850nm. The value of Ms= 70 emu/g and He=540 Oe was obtained at 900°C/sh.

# QN19

## Structural and magnetic properties of Mn<sub>2</sub>Rh<sub>1-x</sub>Co<sub>x</sub>Sn Heusler alloys

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 $Mn_3$ -based Heusler compounds are promising candidates for spintronic applications as they are known to crystallize with cubic and tetragonal structures or exhibit cubic-tetragonal phase transitions. Sn-containing compounds have only a small lattice mismatch with MgO thus providing higher symmetry correlation between magnetic film and tunneling barrier. Quaternary  $Mn_3Rh_{1,2}Co_3Sn$ and  $Mn_{2,2}Rh_{1,3}Sn$  Heusler alloys have been synthesized with a step of x = 0,1 and their magnetic properties were experimentally investigated. In the present work the structural and magnetic properties are discussed depending on the Rh and Co content. The first series undergoes a cubic-tetragonal transition at x = 0,5, while the latter one experiences a hexagonal-tetragonal-tetragonal change of structure. The presence of Co increases the Curie temperature keeping the magnetic moment unchanged. This work is supported by the Graduate School of Excellence Materials Science in Mainz (MAINZ).

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## QN20

# Magnetic properties of $Mn_2Sb_{1,x}Ge_x$ (0.05 $\leq x \leq$ 0.2) in high magnetic fields

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The Mn<sub>2</sub>Sb compound with a Cu<sub>2</sub>Sb-type tetragonal structure is ferromagnetic (FRI) below TC ~ 550 K. The substitution of various elements (V, Cr, Co, Cu and Zn) for Mn as well as (As, Ge and Sn) for Sb results in a first-order magnetic (FOM) transition from the ferromagnetic (FRI) to an antiferromagnetic (AFM) states with decreasing temperature [1]. Mn<sub>2</sub>,CoSb, it has been reported that the FOM transition occurs accompanied with large change in unit cell volume, resistivity and magnetization etc. However, detailed results on Mn2Sb<sub>1</sub>,Ge, have not been reported yet. In this study, magnetization and electrical resistivity measurements were carried out for polycrystalline Mn<sub>2</sub>Sb<sub>14</sub>,Ge, (0.05  $\le$  x  $\le$  0.2) in magnetic fields up to 16 T in 4.2-600 K temperature range, in order to investigate the magnetic and electrical properties affected by magnetic fields. Mn<sub>2</sub>Sb<sub>0.92</sub>Ge<sub>0.08</sub> showed the Curie temperature Tc of 532 K and the FOM transition at 1 = 230 K in a zero magnetic field. With increasing x, Tc decreased and Tt increased. For Mn<sub>2</sub>Sb<sub>0.02</sub>Ge<sub>0.08</sub>, we observed a matamagnetic transition from the AFM to FRI with a hysteresis of 0.5 T at 200 K.

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# QN21

# Microwave magnetic properties of FeNi films prepared by electrodeposition

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Lanznou University, Chi

With the development of computers and communication devices using the electromagnetic wave of GHz range, the serious electromagnetic interference (EMI) have been emerged. The electromagnetic wave absorbers have been widely investigated to eliminate the EMI problems. Electrodeposition is one of the simplest, most flexible, and cheapest processes available for the fabrication of single component and multilayered thin films. In this study, FeNi alloys thin films have been successfully electrodeposited on Ag films. The morphology, structure, composition, and magnetic property of the FeCo films were characterized by scanning electron microscopy, x-ray diffraction, energy dispersive spectrometer, vibrating sample magnetometer and network analyzer. With the increase of depositon time, the coercivity of deposited FeNi films reduce until getting a minimum value and the the real part of permeability increase. At the same time, the increase of the thickness leads to the decrease of resonance frequency. We also study the effect of the composition on the coercivity and microwave magnetic properties. By changing the composition of FeNi alloy films, the effective permeability and the resonance frequency can be adjusted.

# **ON22**

### Preparation and characterization of ferromagnetic fluid and magnetorheological fluid

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Ferromagnetic fluid and magneto-rheological fluid are liquids that nanometer-sized particles of iron dispersed in the medium. PEG(polyethylene glycol) and silicon oil were used as the fluidic media. Fluids are flexible dispersion state under a nonmagnetic field. But, under a magnetic field, particles are oriented in the direction of the magnetic field and the fluid shows the solid-like behavior. In this study, we studied to know how the magnetic field impacts the yield stress of ferromagnetic and magnetorheological fluids. Yield stress of the fluid is evaluated by the viscosity measurement. Viscosity behaviors of ferromagnetic and magneto-rheological fluids are measured with a magneto-rheological viscometer (MCR 301, Anton paar, Austria) with applying magnetic field from the plate. Two types of experiment were conducted. One is a static test in which the viscosity of fluids is measured with increasing shear rate while the applied magnetic field is fixed. The other one is a dynamic test in which the viscosity

of fluids is measured with increasing magnetic field while shear rate is fixed. As the results, viscosity of ferromagnetic fluid and magneto-rheological fluid increased largely by applying magnetic field in all conditions.

**ON23** 

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## Microwave magnetic and absorbing properties of the planaranisotropy Ce<sub>2</sub>Fe<sub>17</sub>N<sub>3-6</sub> powder composite

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Ce2Fe17N3-8 powders were prepared by arc-melting, milling and nitrogenation. The microstructure and static magnetic properties of the powders were analyzed by x-ray diffraction (XRD) and vibrating sample magnetometer (VSM), respectively. The c-plane alignment of Ce2Fe17N3-6 was evaluated with the XRD data of an aligned composite sample. The effective complex permeability of Ce2Fe17N38 powders/epoxy resin composite was measured in the frequency range of 0.1-15 GHz. The intrinsic static permeability of Ce2Fe17N3-8 powder was calculated by Bruggerman (BG) effective medium theory. The natural resonance frequency of Ce2Fe17N3.8 was obtained using Landua-Lifshitz-Gilbert (LLG) equation. The Snoek's product of the effective static susceptibility and the resonance frequency is much larger than the Snoek's limit of the composite. The reflection loss of the composite sample was determined by simulation and by measuring the S11 parameters after rear face of sample was terminated by metal. The minimum reflection loss was found moving towards the lower frequency region with increasing the sample thickness. Quantitatively analysis reveals that the peak frequencies (fp) of the dips (minimum reflection loss), the  $\mu$ (fp), the  $\epsilon$ (fp) and the sample thickness obey the quarter-wavelength ( $\lambda/4$ ) formula. Accordingly, the quarter-wavelength formula have been successfully applied to explain and predict the microwave absorbing properties

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# **ON24**

### Co<sub>2</sub>Y-NiCuZn ferrite composites for high frequency applications Ruei-lin Lin and Hsing-i Hsiang\*

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The most commonly used materials in inductors for high-frequency applications are NiCuZn ferrites, or non-magnetic materials such as low-temperature cofired ceramics. The non-magnetic materials are used for high-frequency applications because NiCuZn ferrites (ferroxcube) typically exhibit severe property changes above 200MHz due to the Snoek limit. The maximum quality factor frequency in chip inductors made of nonmagnetic material is over 500 MHz. However, the quality factors at frequencies around 200-300MHz are much lower than those at higher frequencies. Because chip inductors are prepared by winding a wire around a non-magnetic material core it is necessary to have a larger number of coil winding turns to obtain the desired inductance; thus limiting miniaturization. Therefore, it is desirable to develop a magnetic material that has a higher initial permeability and quality factor than non-magnetic materials at 200-300 MHz, for use in making high frequency inductors. In this study, a  $Co_2$ Y-NiCuZn ferrite composite with high permeability was prepared by coating NiCuZn ferrites onto the surface of Co2Y ferrites using heterogeneous nucleation process. The effects of NiCuZn ferrites addition on the densification behavior, micro-structure and magnetic property of Co<sub>2</sub>Y-NiCuZn ferrite composites were investigated.

# **ON25**

### A novel low-temperature-fired multifunctional varistor-magnetic NiCuZn ferrites

Wei-hung Hsu, Hsing-i Hsiang\* and Li-then Mei Resources Engineering, National Cheng Kung University, Taiwan

A novel low temperature-fired(950°C) multifunctional varistor-magnetic ferrite materials can be obtained by adding Bi2O3 into Zn05Ni03Cu019Nb001Fe202O4 ferrites. The relationship between the microstructure and varistor properties were investigated using scanning electron microscopy (SEM), X-ray Diffraction (XRD), and impedance analysis. For Zn<sub>0.5</sub>Ni<sub>0.3</sub>Cu<sub>0.19</sub>Nb<sub>0.01</sub>Fe<sub>2.02</sub>O<sub>4</sub> ferrites with 2.5wt% Bi<sub>2</sub>O<sub>3</sub> content sintered at 950°C for 2h, the nonlinearity exponent ( $\alpha$ ) can reach 5.1 and the initial permeability( $\mu$ ) is about 200 to 250. Moreover, the avalue can be promoted to 67.8 for the Zn<sub>0.5</sub>Ni<sub>0.3</sub>Cu<sub>0.19</sub>Nb<sub>0.01</sub>Fe<sub>2.02</sub>O<sub>4</sub> ferrites added with 9Bi<sub>2</sub>O<sub>3</sub>-CuO sintered at 950°C for 2h

## **ON26**

### Synthesis and magnetic properties of transition metal ferrite nanoparticles

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The use of ferromagnetic metal ferrite particles for magnetic hyperthermia application has been extensively studied by many scientists. The use of ferromagnetic particles with low curie temperatures result in the formation of self regulated magnetic hyperthermia. In this context, multiple experiments on synthesizing ferrite nanoparticles with varying cobalt and nickel oxide concentrations have been done. The products were analyzed using scanning electron microscopy, energy dispersive spectroscopy, X-ray diffraction and vibrating sample magnetometer. The microscope images indicated the formation of nanoparticles while the spectroscopy and diffraction patterns indicated the formation of nickel and cobalt ferrite nanoparticles along with additional nickel oxide and cobalt oxide. Magnetometer analysis conducted on the samples indicated a reduced magnetic saturation and hysteresis. Hysteresis loss energy was determined by calculating the area under the hysteresis loop, frequency of the magnetic field and the curie temperature of the material. The results indicated a two-fold decrease in the hysteresis energy giving a hint for their use in controlled magnetic hyperthermia treatment.

## 0001

## High room temperature power factor (Z) in K<sub>x</sub>RhO<sub>2</sub>

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We discuss the thermoelectric and optical properties of layered KxRhO2 by means of the electronic structure determined by first principles calculations as well as Boltzmann transport theory. The main contribution to the optical spectra are due to intra and interband transitions between the Rh 4d and O 2p states. Our results of transport perproties of K<sub>x</sub>RhO<sub>2</sub> at room temperature show highest value among hole-type materials. Specially at 100 K, the Z value of K<sub>x</sub>RhO<sub>2</sub> is 3x10<sup>-3</sup> K<sup>-1</sup>. Our results show that the electronic and optical properties of KxRhO2 are similar to NaxCoO2 with an enhanced transport properties.

# 0002

### Optimized Halbach array based magnet systems for Lorentz force velocimetry purposes

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Lorentz Force Velocimetry (LFV) is in research since 2003. The Lorentz force acting on a source magnet system is providing a velocity dependent signal. The efficiency of the magnet system (ratio of generated Lorentz forces to the mass of the system) is an essential parameter for sufficient resolutions of the whole LEV sensor device. The full article will propose a very efficient combination of permanent magnet arrays, where the particular permanent magnets have different magnetization directions. Here an iron voke is not necessary, respectively ineffective, and so the arrays can be assembled with a very light carbon fiber composite fixture in order to reduce the overall mass further. The underlying principle of this array combination will be presented and the geometric optimization with commercial software, as well as the assembling method, will be explained. On an experimental test channel for LFV on electrolyte flows the magnet system is tested and compared to another simple standard magnet system with single permanent magnet blocks instead of the arrays. From previous finite element simulations an increase of a factor of three in efficiency is estimated and will be proven in the experiment. Further improvement using cladding technique will be introduced too

## 0003

### Growth and characterization of SmFe/Ge/single crystals

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Hexagonal single crystal plates of SmFe<sub>6</sub>Ge<sub>6</sub> have been grown by the flux method using a 4th flux element (Sn), extending the range of members in the RFe6Ge6 family of magnetic compounds, whose physical properties have so far been described only for rare earths R = Gd-Lu [1]. This family exhibits antiferromagnetic (AF II) ordering at 450 K of the layered Fe honeycomb sublattice, and independent ordering (or no ordering) of the R sublattice moments below 30 K. Here we present a basic characterization of the structural, magnetic and transport properties of the SmFe6Ge6 crystals and compare the general anisotropic behavior with other members of the family. We thank the support of Brazilian agencies FAPESP and CNPq.

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## 0004

### Anomalous magnetic and related properties of Nd<sub>5</sub>Ge<sub>3</sub>

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We have observed that the stoichiometric intermetallic compound Nd<sub>5</sub>Ge<sub>3</sub> shows many anomalies in the magnetization, resistivity and heat capacity data. This compound undergoes an antiferromagnetic (AFM) transition at 46 K and another one at 30 K on cooling. It is found that the strengths of these two antiferromagnetic states are different. Ultra-sharp jumps are seen in the low temperature isotherms of magnetization, heat capacity and electrical resistivity. Another important anomaly is the field-induced irreversibility in all these three properties, which suggests that the spin, electronic and lattices states are strongly coupled in this compound. The memory effect seen here is quite remarkable and is consistently seen in various measurements. Another striking observation is the spontaneous evolution of the ferromagnetic (FM) phase as a function of time, at a fixed temperature and at the critical field. Sweep rate and cooling field dependencies point towards the role of martensitic strains in the anomalous behavior. The compound appears to show the coexistence of FM and AFM components at low temperatures. The features seen in this compound indicate the presence of disorder in influencing the magnetic and related properties considerably. The correlation of various properties as seen from the martensitic scenario will be presented

# 0005

### Magnonic meta- and meta-meta-materials

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We report on magnonic metamaterials represented by a system of magnetic layers separated by those of a dielectric non-magnetic material. In such a metamaterial, the constituent elements are made from a magnetic material that provides an easy way to control its electromagnetic response by tuning the bias magnetic field. Firstly, we report a theory of the effective permeability of planar multilavered metamaterials containing thin ferromagnetic layers with magnetization pinned on either one or both surfaces. The pinning of the magnonic modes at film surfaces facilitates their coupling to the uniform magnetic field of the incident electromagnetic wave. We demonstrate that such structures, for magnetic parameters characteristic for conventional transition metal alloys, exhibit a negative refractive index at sub-THz frequencies. Furthermore, we show how the constituent magnetic layers of such systems could themselves be nanostructured into planar magnonic crystals, with the whole structure therefore representing a meta-metamaterial. Due to the additional nanostructuring, the negative permeability at sub-THz frequencies can be achieved as a result of a small lattice constant and lateral quantization of spin waves. Finally, we discuss the prospects for fabrication of such metamaterials and their experimental studies

## 0006

### Understanding the nature of magnetic phase transition in relevance to magnetic refrigeration

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It is important to understand the nature of magnetic phase transition for a material suitable for magnetic refrigeration. With this aim, we have studied a wide class of intermetallic compounds with diversity in their nature of magnetic ordering and phase transition, such as ferromagnetic ordering, ferrimagnetic ordering, coexistence of ferromagnetic and antiferromagnetic ordering, metamagnetic transition, and magnetostructural transition. The dependences of the nature of magnetic phase transition on temperature, magnetic field and chemical substitutions have been investigated. The role of the nature of exchange interaction and universality class in determining magnetocatoric effect (MCE) has been brought out by investigating the intermetallic compounds TbCo2, Fe, and Mn<sub>4</sub>FeGe<sub>3-x</sub>Si<sub>x</sub>. A phase coexistence of ferromagnetic and antiferromagnetic phases has been observed in Cu<sub>1</sub>  $_xNi_xMnSb$  semi-Heusler alloys around x = 0.15 where a large MCE is found as a low applied magnetic field causes an AFM to FM phase transformation. A broad operating temperature range due to coexistence of AFM and FM phases, and quenched disorder has also been obtained. The role of other interesting phenomena, such as metamagnetic transitions and domain wall pinning in obtaining a large MCE and an inverse MCE, respectively, has been brought out through the investigation of the intermetallic compounds NdMn2xCoxSi2

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## 0007

### Specific heat study on successive magnetic transitions in α-Dy<sub>2</sub>S<sub>3</sub> single crystal under magnetic fields

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Specific heat measurements in magnetic fields have been performed on  $\alpha$ -Dy<sub>2</sub>S<sub>3</sub> single crystal which shows successive magnetic transitions at TN1 = 11.5 K and TN2 = 6.4 K<sup>1</sup> The specific heat in no magnetic field exhibits sharp peaks at both temperatures of TN1 and TN2. The change of magnetic entropy at each transition is estimated as N0kBln2 per mole of Dy2S3, which suggests Dy moments on only one site between two crystallographically inequivalent Dy sites order at each transition temperature. When the magnetic field is applied along the b-axis of the orthorhombic system, two peaks of the specific heat shift toward lower temperatures. On the other hand, the magnetic field perpendicular to the b-axis shifts the peaks toward higher temperatures. The transition temperature TN1 shifts to 9.6 K (H//b) and 12.5 K (H-b) with keeping sharpness of the peak under the magnetic field of µ0H=2 T. The peak of TN2 broadens gradually with increasing the magnetic field for each direction, and the peak is consequently obscure under the field of 2 T.

<sup>1</sup>S. Ebisu, M. Narumi, M. Gorai and S. Nagata, J. Magn. Magn. Mater. 310, 1741 (2007).

# QO08

# Magneto-inductive wave in periodic chains of ferrite cores and chip capacitors

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In this paper, we present a new magneto-inductive wave generated in a chain of resonators fabricated with ferrite cores and chip capacitors. RF signal propagates to neighbor resonator one by one as a consequence of the magnetical coupling between two resonators in our device. The magnetical coupling is due to the mutual inductances along chains of resonators. So, the signal amplitude (? coupling intensity) is inversely proportional to the distance between two elements. In order to design the device, some electromagnetic simulations and experiments have been carried out systemically. The transmission characteristics of a magneto-inductive wave could be controlled by applied external magnetic field. For, the device composed of 5 resonance elements, the center frequencies were estimated to be 41.1 MHz and 34.8 MHz with the external magnetic flux density of 750.72 G and 2680 G, respectability. The reported result opens a promising way to a high variety of applications in one- and two-dimensional functional devices, such as transducers, delay lines, power dividers and couplers.

## Q011

### Reconstruction of cubic rs-ZnO on MgO (200) substrate through 100 plane of w-Zn for transparent electronic application

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ZnO crystallizes in various crystallographic forms. However, wurtzite (w) ZnO is generally observed structure while the rock salt (rs) structure of ZnO appear at high pressure. We present rs-ZnO structure grown on MgO (200) substrate at certain growth environments, not high pressure. The lattice constant of MgO. The possible configurational path for transformation is proposed, so that the compression of w-ZnO along c-axis and movement of Zn or O sublattice along (110) direction and (110)plane. These results are confirmed by X-ray diffraction and X-ray photoemission spectrum experiments. It has also been observed that rs-ZnO on MgO substrate is stable for certain minimum thickness (in this case, it is 150nm) and possible be stabilize at lower temperature for few nm thickness

# Q009

### Large refrigerant capacity in Ni<sub>2.9</sub>Co<sub>0.1</sub>MnIn type-Heusler alloy

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Magnetic refrigeration (MR) is a cooling method that is based on the magnetocaloric effect (MCE). Our research effort focuses on identifying magnetic structures that may have improved efficiency with suitable characteristics for MR i.e. good magnetocaloric properties with small or zero thermal and magnetic hysteresis. Some Heusler alloys may meet these criteria. The Ni<sub>2</sub>Mnln full-Heusler alloy exhibits a second-order paramagnetic to ferromagnetic transition at 300 K, but its MCE is small [1,2]. However, we observe a large improvement in the refrigerant capacity (RC) value when the Ni<sub>2</sub>Mnln is Co and Ni-doped. In this paper, we present the study of the magnetic, structural and magnetocaloric properties of the Ni<sub>2</sub><sub>2</sub>Go<sub>1</sub>,Mnln Heusler-type alloy. The samples crystallizes in a cubic structure. Magnetization measurements show a ferromagnetic transition at 207 K. The transition is almost 40 K wide, and thermal hysteresis is negligible. The magnetic entropy change,  $\Delta$ SM, for  $\Delta$ H=5T, was calculated from heat capacity data and the maximum value found is -2.75 J/kg K, with broad temperature magne which leads to RC=261.3 J/kg. The large RC value found and the absence of thermal hysteresis makes this material a good candidate for magnetic refrigeration applications. C.S.M. and A.M.G. would like to than CNPq and FAPERJ.

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### QO10

# Reduction of weak localization strength on controlling oxygen defects by ex-situ annealing

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We have studied Pb-based gapless semiconductors, which are very sensitive to external parameters such as temperature, pressure, magnetic field, electric field, etc [1]. Based on the theoretical paper, this material has a strong spin-orbit interaction. So, we have fabricated pure PbPdO<sub>2</sub>, Co- and Mn-doped PbPdO<sub>2</sub> thin films. Because a Pb is an easily vaporized element when we grow the films, the films have a lot of oxygen vacancies. It may induce other problems. For example, Pd is unstable state as Pd+1 and the oxygen vacancies act as additional hole carriers. To study intrinsic properties of doped PbPdO<sub>2</sub> thin films, we tried to various methods to enhance the films quality. Because Pb is hard to handle in growth process, we have annealed thin films with different annealing time, temperature, and environments after deposition. The carrier densitivities are changed in different way. More importantly, weak localization strength is most likely to be reduced by ex-situ annealing. Furthermore, magnetoresistance ratio is also reduced by increasing ex-situ annealing time.

### Q012

## Magnetic cooling machine prototype based on cold rolled Gd foil

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With the discovery of the "giant magnetocaloric effect" (Pecharsky and Gschneidner, 1997) the development of magnetic refrigeration gained increased momentum [1]. One of the ways of engineering MCE materials is tightly connected with preparing very thin (a few microns) ribbons of high value MCE alloys with good mechanical properties. At present rapid solidification is the main technique for producing this kind of materials [2]. In our work we present magnetic cooling device which is used Gd cold rolled samples of different thickness as a working body. Magnetocaloric, mechanical and other properties of this material are discussed Morphology of the surface was investigated by AFM methods. Magnetocaloric effect was measured by direct method on AMT&C setup. Mechanical properties was investigated on Netzsch DMA 242C. Magnetic properties was measured on SQUID magnetometer. X-ray analysis was performed on Brucker D8 Advance diffractometer. Several axial magnetic systems are proposed as the source of high gradient magnetic field. Depending on the air gap of magnetic system the residual magnetic induction is in 600 mT. 1900 mT range. It is shown, that cold rolling technique is the alternative way for producing thin ribbons of MCE materials with good mechanical properties for magnetocaloric application.

I.Gschneidner K.A. Jr. et al. Thirty years of near room temperature magnetic cooling: Where we are today and future prospects, Int. J. Refrig, 31, 945 (2008). 2.P. M. Shan et al. Magnetic behavior of melt-spun gadolinium, Phys. Rev. B. 77, 184415 (2008).

### QO13

# Research on the orbital state of IrO<sub>2</sub> with resonant x-ray scattering method

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In the 5d electron systems of iridium oxides, the crystal fields, the Coulomb repulsion, and the spin-orbit interaction compete with each other, which results in various exotic physical properties. For example, magnetic iridium oxides  $Sr_2IrO_4$  and  $CaIrO_3$  are known to be spin-orbit-induced Mott insulators: the reported resonant x-ray scattering experiments showed that 5d t2g orbitals of Ir are reconstructed by the spin-orbit interaction and realize the complex Jeff=1/2 orbital states, which form a narrow band and enhance the electron correlation.[1] On the other hand, the orbital states of metallic iridium oxides like  $IrO_3$ , which is attracting attention for a gigantic spin-Hall effect, are yet to be investigated. We performed a resonant X-ray scattering measurement for metallic iridium oxide  $IrO_3$  which has a rutile-type structure, and observed the 0 0 3 reflection coming from the anisotropic tensor of susceptibility (ATS). In the energy spectrum the peak corresponding to the excitation to 5d t2g orbitals appeared around the L3 absorption edge, but did not appear around the L2 absorption edge. These results indicate that the complex Jeff=1/2 states of 5d t2g orbitals are realized even in metallic iridium oxides  $IrO_2$ .

[1] B. J. Kim et al., Science 323, 1329 (2009).

# QO14

# Investigation of magnetic-field tunable properties of magnetorheological elastomers

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In this work we present the results of investigation of magnetic field influence on the value of dielectric permittivity and magnetic susceptibility in magnetorheological elastomers (MRE) based on the powders with various magnetic properties - Fe, FeNdB and Fe<sub>3</sub>O<sub>4</sub>. Both magnetic susceptibility and dielectric permittivity have been found to be considerably changed under magnetic field, the changing of permittivity value up to 150% has been observed. The field dependence of the dynamic dielectric permittivity was significantly non-linear. The anisotropy of field influence on the permittivity and susceptibility was observed as well. The saturation magnetization and coercivity of the magnetic fillers have been shown to be the most important factors affecting the field dependence of permittivity. In addition, the size and concentration of particles as well as conductivity of magnetic material have significant influence on the type of field effect. The influence of named parameters on the permittivity and susceptibility changing is discussing; we consider the mechanisms determining the features of magnetic field influence on studied physical properties in detail. It is shown that MRE based on various magnetic particles can be considered as promising materials with magnetic-field tunable properties. Supported by Russian Foundation for Basic Research.

### Q015

# Disagreement between modification methods for magnetism-absorbing agent

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Two methods are widely used to improve absorbing properties of the microwave magnetism-absorbing agent: mixing with the other material such as resistance carbon fiber absorbent, or wrapping the microwave magnetism-absorbing agent in that material. However, do these two methods coincide with each other? If not, which method should be used as better way to represent the absorbing properties of magnetism-absorbing agents? To answer these questions, both methods are carried out on the same magnetism-absorbing agents, of which the results are compared. Distinct difference between these measured results is found, and the possible reasons for the difference are also discussed.

### Q016

# Ordered magnetic arrays of cobalt SMM: properties and the relationship with crystal symmetry and SMM environment

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The modulation of the magnetic properties of SMMs and their organization in extended magnetic arrays are, at the present, some of the top areas in the field of molecular magnetism for the implication in the improvement of the physical properties of the SMM and the potential use in devices such as for data storage. Few polymers of SMM have been structurally characterized in spite of the number of known SMM. [Co4citrate4]8- cubanes can behave as SMMJ(1) We have reported recently(2) the union of these cubanes through "Co4(H<sub>0</sub>)<sub>2</sub>ETG" (ETG=ethylene glycol) bridges and the formation of anionic 2D layers. These layers crystallize in the tetragonal system, P-421c space group, with counterions (Cs' or Rb'), water and ETG molecules connecting the layers though electrostatic interactions and hydrogen bridges. The magnetic studies show the co-existence of two magnetic nets: the SMM net and the one formed by the bridged cobalt atoms and also shows the influence of the counterions in the layer magnetic meretics. Two intriguing relaxation processes with different characteristics timescales are observed. Moreover, [Co4citrate4]8- cubanes can also organize in 3D, diamond-like structural arrays showing SMM magnetic behaviour below 6K with an energy barrier of 54K and magnetic ordering below 2.7K.

I-M. Murrie, Chem. Soc. Rev. 2010, 39, 1986-1995. 2- E. Burzuri, J. Campo, L. Falvello, E. Forcen-Vazquez, F. Luis, I. Mayoral, F. Palacio, Cristina Saenz de Pipaon, M. Tomas. Chem. Eur. J. 2011, 17, 2818-2822.

# Q017

# Enhancement of the refrigerant capacity in $Gd_{65}Fe_{20}Al_{10}B_5$ alloys by partial crystallization of melt-spun amorphous ribbons

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The magnetocaloric effect (MCE) is attracting a great deal of scientific interest because it is a basis of the environmentally friendly magnetic refrigeration technology. Among the recently developed magnetic refrigerant materials, the Gd[Fe\_Mn]Al-based glassy alloys combine good magnetic entropy characteristics with sufficiently high effective magnetic moment per volume, which makes them promising candidates for magnetic refrigeration in the operating temperatures around 150 K [1,2]. In this work, we report on the beneficial effect of partial crystallization on magnetocaloric properties of melt-spun Gd\_eFe\_avAl\_uB\_3 alloys. The magnetic entropy charges,  $\Delta Sm$ , were calculated from the magnetization versus applied field dependences measured in the temperature range from 5 to 390 K. The value of the maximum magnetic entropy changes under 50 kOe for the Gd\_eFe\_avAl\_uB\_3 ribbon reached 5.01 J/kg K at 185 K. The corresponding value of the refrigeration capacity, RC, determined as the area below the  $\Delta Sm$  peak with the integration limits at its half maximum is 647 J/kg. We have found that this value can be further enhanced by suitable heat treatment of melt-spun rights and a fiscussed in relation with two-phase character of partially crystallized samples.

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## QO18

# Chemical pressure effects on half-metallic properties in single crystals of A<sub>2</sub>FeMoO<sub>6</sub> (A = Ca, Sr, and Ba)

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An ordered double perovskite oxide Sr<sub>2</sub>FeMoO<sub>6</sub> has been predicted and widely studied as a half-metal, in which electrons with only one spin direction present at the Fermi level (100% spin polarization), above room temperature. In this study, in order to improve the half-metallic properties, the chemical pressure effects, through A-site substitution, on the magnetic and electronic properties have been examined using single crystalline samples of A2FeMoO<sub>6</sub> (A = Ca, Sr, and Ba) grown by the floating zone method. In the A = Ca crystal, every property related to the half-metallic performance was found to be degraded as compared with the A = Sr and Ba crystals: the saturation magnetization and the Curie temperature were significantly decreased, metallic conductivity became more than twice larger, and the spin polarization (evaluated by Andreev reflection spectroscopy) was reduced below 50%. From the results in the A = Sr and Ba crystals, it is suggested that the optimized half-metallic performance will be obtained in between these two compositions. Optical measurements were also performed to gain insight into the change of electronic structures with respect to the A-site substitution.

# QO19

### Magnetic label field collector of biochip sensor

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We have proposed a new method called stray field collector (SFC) of magnetic labels for signal enhancement in the magnetoresistive biosensors. The proposed design related to qualitative as well as quantitative detection of magnetic beads. The integration of SFC in magnetic sensor will increase the amount of stray field via the increase of active surface area. Therefore, this SFC may have great potential to improve the performance of magnetic biosensor.

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# QO20

Cogging torque cancellation technique for dual rotor type motor using adjustment between outer and inner rotor magnet angle

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In this paper, we design the IPM motor and discuss the design method of double rotor permanent magnet motor for cancelling the cogging torque. According to downsizing the electric motor, high torque and power density are the most serious requirements for electric machines. Double rotor permanent magnet motor is proposed to solve the problem. But it has a large cogging torque that generates the large acoustic noise. A variety of techniques exist for reducing the cogging torque of conventional radial flux IPM machines. This paper introduces a new cogging torque minimization technique for radial flux multiple IPM motors. Optimization of the adjustment magnet angle between outer rotor and inner rotor which results in offset cogging torque and assessment of the effect on the maximum available torque using Finite Element Analysis (FEA) is investigated. In our experiments, Prototype dual rotor type motor has achieved the outer rotor and inner rotor cancellation of cogging torque by 70 percent.

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0021

## A broadband circuit analog absorber based on printed dipole antenna

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Based on the equivalent transmission line theory, a composite absorber with broadband planar array antenna is designed. The first layer is an array antenna, which uses printed dipole with balun feed as basic unit. The second layer a PIN diode array, which corresponded with antenna feeding units. The impedance match characteristic of the absorber can be adjusted by setting the microstrip balun. And the reflection characteristic of the absorber can be adjusted by adjusting the bias voltage of the PIN diode array. Simulated by electromagnetic simulation software HFSS, the results show that the absorber reflection coefficient is below -10dB at 3.5-15GHz. The experimental result, which replace PIN diodes with surface mount device resistors, is in good agreement with simulation data.

# QO22

Effect of oxygen pressure to magnetism and transport properties of epitaxial  $Fe_3O_4$  grown by molecular beam epitaxy

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 $Fe_5O_4$  is predicted to possess as half-metallic nature, ~ 100% spin polarization, and has a high Curie temperature and an ultrahigh room temperature magnetoresistance (MR) for tunneling magnetoresistance (TMR) junctions with  $Fe_5O_4$  electrodes [1]. However, for the TMR junction, a highly conductive under layer and the sharp between interfaces are required because the poor conductivity and roughness of films may lead to a non-uniform current distribution [2]. In addition, the Verwey transition (TV, a first order metal-insulator transition) of 120 K in bulk  $Fe_5O_4$  is still under controversy because many parameters such as orientation of substrate, buffer layer, thickness, pressure, and thermo-chemical treatment affect the TV [3, 4]. Epitaxial  $Fe_5O_4$  films were grown on MgO(001) substrates via molecular beam epitaxy. The growth modes, magnetism, and transport properties of the thin films were strongly dependent on the oxygen pressure during film growth. The average roughness decreased from 1.021 to 0.263 nm when the oxygen pressure dus increased from 2.3x10-7 to 8.2x10-6 Torr, respectively. In addition, the saturation magnetic moment reduced as increasing the oxygen pressure during grown. Moreover, the Verwey transition in the Fe<sub>2</sub>O<sub>4</sub> disappeared for a sample grown at high oxygen pressure.

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# QO23

# Magnetism and transport properties of epitaxial Mn<sub>5</sub>Ge<sub>3</sub> thin films on GaAs(001) and GaSb(001)

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The epitaxially ferromagnetic  $Mn_sGe_3$  thin films were grown on GaSb(001) and GaAs(001) substrates using molecular beam epitaxy. Our result revealed that the substrate facilitates to modify magnetic and electrical properties due to tensile strain effect. The Curie temperature was remained above 350 K for samples grown on GaAs(001), while was about 320 K for samples grown on GaAs(001). The anomalous Hall effect were observed up to 370 K and 320 K for films grown on GaAs(001) and GaSb(001), respectively, which are strong evidence for spin polarization. Our electronic calculation indicated that the difference in Mn-Mn bond length in  $Mn_sGe_3$  modified the magnetization and spin polarization. The strains are promised to control both magnetic properties and spin polarization.

## QO24

## Current-driven domain wall motion in artificial magnetic domain structures

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We report progress towards optimisation of artificial magnetic domain structures for efficient spin transfer torque domain wall (DW) motion [1]. Co/Pt multilayer samples have been sputtered on (100) Si/SiO<sub>2</sub> substrates and perpendicular magnetic anisotropy confirmed using polar magneto-optical Kerr effect (MOKE) measurements. The influence of the thickness of Co and Pt layers on the coercivity and switching behaviour was systematically investigated and the conditions established for realising well-suited structures with medium coercivity (~100 Oe) and sharp switching fields [2]. Optimised Co/Pt multilayer films were lithographically patterned into nanowire devices for time-resolved extraordinary Hall effect (EHE) measurements. Our devices are based on 50 W coplanar waveguides incorporating single and double Hall cross structures. The coercivity of the region surrounding the Co/Pt Hall crosses was reduced by local focussed ion beam (FIB) irradiation [3] allowing the controlled nucleation of domain walls at the edges of these regions by application of an appropriate field sequence. We describe polar MOKE experiments that show how DC currents lead to asymmetric switching of these artificial domains due to current-assisted DW motion across them. We will also present preliminary results of time-resolved EHE and MOKE measurements of current-driven DW motion in these structures.

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## QO25

# The cellular uptake mechanism of SPIONs: an in-vitro study

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The Superparamagnetic Iron Oxide Nanoparticles of sizes ranging from 10-50 nm are being used in a large number of biological studies because of their peculiar characteristics of inducing local hyperthermia, MR imaging, specific targeting and drug delivery. An in-vitro study of the cytotoxicity and an understanding of the specific pathway of cellular uptake will enable manipulation of conditions for optimal cellular uptake of SPIONs for targeted therapy. The objective of the present study was to identify the endocytotic pathway through which the SPIONs are taken up by C6 glioma cells. The cells were pre-incubated with different concentrations of pharmacological inhibitors and then exposed to SPIONs for a few hours. The endocytosed particles were localized and quantitatively estimated using Perl's or Prussian Blue reaction. There was significant reduction in the uptake of SPIONs when incubated with the inhibitor indicating the uptake of nanoparticles is being inhibited. This reduction in SPION uptake was found to be dependent on the concentration of the inhibitors, the data can be narrowed down to the final pathway that may be involved in the SPION uptake.

# QO26

# Theoretical and experimental studies of valence states in Fe-Mo compounds

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The half-metallic ferromagnetic double perovskite compound Sr<sub>2</sub>FeMoO<sub>6</sub> is considered as an important material in view of its potential spintronic applications. It appears to be fundamental to understand the role of electronic parameters controlling the half-metallic ground state. In particular, half-metallic ferromagnetic compounds have been searched as a source of spin polarized currents. In this work, we study the valence states of cations in Fe-Mo based double perovskites from theoretical and experimental point of view. In the experimental case, we apply the electron energy-loss spectroscopy (EELS) technique. For transition metals, L2,3 edges of EELS are characterized by two sharp peaks, known as white lines. EELS experiments have shown that the change in valence states of cations introduces significant modification in the energy and shape of the white lines, leading to the possibility of identifying cation valence states. These results are compared with those coming from theoretical calculation using the Green functions and the renormalization perturbation expansion method. The theoretical model is based on a correlated electron picture with localized Fe-spins and conduction electrons interacting with the local spins via a double-exchange-type mechanism.

## QO27

# Structural modifications in magnetic MWCNT by 100 MeV SHI irradiations

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Carbon nanotubes (CNTs) possess unique mechanical and electronic properties suitable for fabricating the nano-scale building blocks of nanodevices [1]. One of the requirements for applications is to cut the CNT in small dimensions in the order of few nanometers. To this respect, we investigate the effect of heavy ion beams, namely, 100 MeV of O and Ni ion beam on CNT synthesized by chemical vapor deposition (CVD) technique[2]. The ion fluence selected for this study was 1e13 ions per cm<sup>2</sup>. High energy heavy ion irradiation was carried out at Inter-University Accelerator Centre, New Delhi, India. This paper reports the experimental results obtained from X-ray diffraction pattern and images of scanning electron microscopy (SEM) and transmission electron microscope (TEM) measured at Nano Materials Analysis Center, KIST, Korea. These results show that heavy ions can induce damage or cut in an ordered periodicity of few nanometers. Apart from above studies, we also carried out near-edge x-ray absorption spectroscopy investigations for electronic structure modifications.

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## QO28

# Magnetic and magneto-transport properties of $(Mn_{0.8},Zn_{0.2})_{\rm 1-x}Ga_xFe_2O_4$ ferrites

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A significant enhancement of low-field magnetoresistance (LFMR) could be achieved from polycrystalline samples of Ga<sub>2</sub>O<sub>3</sub>-doped Mn-Zn ferrites. For this study, the magnetic and magnetotransport properties of samples with two different nominal compositions of ( $Mn_{0.8}$ ,  $Zn_{0.2}$ )<sub>1.4</sub>Ga, Fe<sub>2</sub>O<sub>4</sub> (x = 0.05, 0.1) were carefully investigated. The X-ray diffraction analyses revealed that all Ga<sub>2</sub>O<sub>3</sub>-doped Mn-Zn ferrites of the spinel structure were a pure phase without the second phase. By doping Ga<sub>2</sub>O<sub>3</sub> to the pure Mn-Zn ferrite, the resistivity level was abruptly decreased over four orders of magnitude. The 5 mol% Ga<sub>2</sub>O<sub>3</sub>-doped Mn-Zn ferrite sample exhibited the LFMR ratio over 1.7 % at 300 K in 0.5 k Oe without appreciable increase in its resistivity. Detailed doping effects of Ga<sub>2</sub>O<sub>3</sub> on the magnetic and maneto-transport properties of the Mn-Zn ferrite will be presented and discussed in this paper.

# QP01

### Magnetic coupling at the CoO/Ni interface

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We study the CoO/Ni interface. Since the lattice mismatch between Ni and CoO is 21%, we use a large supercell for our first principles calculations which reduces the lattice mismatch to 0.8%. We investigate the structural, electronic, and magnetic properties of two interface configurations: (1) An O layer mediates the coupling between Ni and Co and (2) a direct Ni-Co contact. Our results indicate that the magnetization is reduced by 19% in the first case, while in the second case it increases by 100% as compared to bulk Ni. The magnetic moments obtained at the interfaces can be explained by the local environment of the interface atoms. In addition, we find effects of charge transfer between the interface atoms.

## **QP02**

# Electronic, structural, and magnetic properties of O and Py deficient CoO/Py interfaces.

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The development of magnetic devices, which are based on coupling between ferromagnet and antiferromagnet (such as spin-valves), depends strongly on the interfacial magnetic structure. Since vacancies at the interface are expected to be important in determining the macroscopic properties of such systems, a quantitative determination of their influence on the interfacial magnetization is essential. We therefore investigate the CoO/Py interface ( $Py = N_{100} Fe_{20}$ , Permalloy). Py has a high permeability, low magnetostriction, and high anisotropic magnetoresistance. It is a good conductor for spin-majority electrons and poor conductor for spin-minority electrons. At the same time the Neel temperature of CoO is approximately room temperature and its magnetocrystalline anisotropy is high. Since the lattice mismatch between Py and CoO is about 20%, we use a large supercell for our first principles calculations which reduces the lattice mismatch to less than 0.1%. We describe the dependence of the interface magnetization on the amount of O and Py vacancies and address the exchange bias.

# QP03

### Effect of the external fields on SpinRAM switching time

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Spin current driven magneto resistive random access memory (SpinRAM) is one of the candidates for the next generation non volatile memory. There are still many problems to realize it. One of the problems is the switching speed. The method for fast switching with lower current is required. Field assist switching is one way. In this paper, the effect of the DC and AC external field on the switching time of the SpinRAM with perpendicular anisotropy is investigated by simulation. The effect of the thermal noise is also investigated. From the results without thermal noise, the switching time was decreased by both of the DC and AC field. However, AC field is effective for fast switching compared with DC field. From the results with thermal noise, the switching time was not decreased by weak DC field and more strong DC field was required to reduce the switching time. On the other hand, the switching time was decreased by weak AC field. From these results, it was found that AC field is more effective than DC field to reduce the switching time. This study was supported by New Energy and Industrial Technology Development organization (NEDO) partly.

## **QP04**

# A study on the perpendicular toggle-MRAM system by using new combined hysteresis method for high Gb/Chip Hvuk Won and Gwan Soo Park\*

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Conventional (toggle) MRAM has unlimited read/write endurance but has a low capacity than flash technology. A lot of work has been done for the commercialization of submicrometer MRAM in recent years. A low-power 1Mb MRAM with copper interconnects and cladding layer, and about 35 ns access and cycle times tested in the previous work show the potential of this technology. A new multibit MRAM cell of toggle switching type has been reported and also, a heat interaction investigation in thermally assisted MRAM has been reported. In this paper we presents a new design that has advantages conventional MRAM on injected current and cell size. It is pole type perpendicular MRAM (PTP MRAM). PTP MRAM uses perpendicular magnetic field in order to change the state of the free layer in a perpendicular magnetic tunnel junction (pMTJ). In this paper, new efficient algorithm to combine two models are presented and tested. This approach reduces computing time and gives more precise results because the result of 2 dimensional Preisach model becomes an initial values of 3 dimensional micromagnetics.

# QP05

### Room temperature exchange bias in the multiferroic BiFe<sub>0.8</sub>Mn<sub>0.2</sub>O<sub>3</sub> nanoparticles with a core-shell structure S M Yusuf\* and P K Manna

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The compound BiFeO3 shows a multiferroic character. We have shown that the BiFeO3-based compound viz. BiFe0.8Mn02O3 nanoparticle system exhibits other important multifunctional behaviour as well. We have observed the exchange bias effect, in the form of a shift in the field-cooled hysteresis loop, and a training effect in a dc magnetization study. From transmission electron microscopy, neutron diffraction, thermoremanent magnetization, and isothermoremanent magnetization measurements, the nanoparticles are found to be core-shell in nature, consisting of an antiferromagnetic (AFM) core, and a two-dimensional diluted AFM (DAFF) shell with a net magnetization under a field. Our calculation using the Binek's model on the training effect yields a theoretical understanding of the observed loop shift which arises entirely due to an interface exchange coupling between core and shell. The intrinsic contribution of the DAFF shell to the total loop shift is zero. The present study is useful to understand the origin of exchange bias in other DAFF-based systems as well. The significantly high value of the exchange bias field (~ 0.16 kOe), observed at room temperature, in the present nanoparticle system has made it a blue-eyed candidate for the next generation advanced nano-technological applications in the area of magnetic storage devices.

### **QP06**

# Magnetic and Transport Properties of Mn<sub>3-x</sub>Ga/MgO/Mn<sub>3-x</sub>Ga Magnetic Tunnel Junctions: A First-Principles Study

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Magnetic and transport properties of  $Mn_{3-x}Ga/MgO/Mn_{3-x}Ga$  ( $0\le x\le 1$ ) magnetic tunnel junctions are studied using first-principles approach based on density functional theory and non-equilibrium Green's function. Perpendicular magnetization, of which the magnetic anisotropy energy reaches more than 1 meV/unit-cell, is confirmed to be energetically favoured by both  $Mn_2Ga$  and  $Mn_3Ga$  thin films. Furthermore, despite high spin-polarization at Fermi energy for both these compounds as reported, our transport calculation shows considerable disparity in the transmission behaviour between  $Mn_2Ga/MgO/Mn_2Ga(001)$  and  $Mn_3Ga/MgO/Mn_3Ga(001)$  magnetic tunnel juctions: huge optimistic tunneling magnetoresistance artio ~ 1000% for the former, and nevertheless, no tunneling magnetoresistance effect absolutely for the latter. This phenomenn is attributed to the spin symmetry filtering effect of the MgO spacer. On this premise,  $Mn_{3-x}Ga$  compounds with low Mn concentration are predicted to be promising candidate materials to serve as the electrodes of spin-transfer toque devices in

### **OP07**

### Concomitant memory effect in CrO<sub>2</sub>/Cr<sub>2</sub>O<sub>3</sub> core-shell nanorods Ashish Chhaganlal Gandhi and Sheng Yun Wu\*

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The concomitant memory effects are intensively studied in last few decades and are mostly observed in spin-glass and superparamagnetic systems [1-3]. The memory effects were observed in the interacting permalloy nanoparticles [4] and isolated nanoparticles [5], providing us interesting information about the investigation of nature of the low temperature phase. In the present work, we succeeded in preparing CrO<sub>2</sub>/Cr<sub>2</sub>O<sub>3</sub> core-shell nanoparticles can be treated as a ferromagnetic-CrO<sub>2</sub>core and antiferromagnetic-Cr<sub>2</sub>O<sub>3</sub>O-shell compound coupled by an exchange interface coupling and raised magnetic anisotropy energy between them. The evolution of core shell magnetization triggered by an applied magnetic field at 1.8 K was observed, and a core-shell anisotropic energy model was applied to simulate the competition between the surface magnetic anisotropy and exchange interface coupling. Memory and aging effects were observed and overcame the superparamagnetic limit [3] up to sufficiently high temperature which makes it an important candidate in magnetic data storage, as an application. Acknowledgments: This research was supported by a grant from the National Science Council of Taiwan, the Republic of China, under Grant Number NSC-100-2112-M-259-003-MY3.

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### **OP08**

### High resolution probes for magnetic force microscope

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Magnetic force microscopy is a convenient and efficient method to detect magnetic domain structures in magnetic recording media or nano-magnet. In this work, we present a simple process to fabricate probes for magnetic force microscopy with sub-ten nanometer resolution. A thin layer of FeCo is deposited onto commercially available silicon probes by a facing targets sputtering system. We choose facing targets sputtering because this kind of system can minimize the damage of the films from high energy particle bombardment. The radius of the curvature of the probes is depended on the FeCo film thickness. High resolution scanning microscope image show the radius of the curvature is less than 20 nm for a 20 nm thick FeCo coating. Magnetic force microscopy images show our probes have resolution better than 10 nm. Due to the high magnetic moment of FeCo, we can get magnetic contrast for FeCo coating as thin as 10 nm. Micromagnetic simulations indicate than the formation of a magnetic vortex at the apex of the probe is crucial for the high resolution.

### OP09

# Selective magnetization switching by microwave assistance for layered magnetic pillar

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Microwave assisted magnetization reversal (MAMR)[1] is a promising candidate for future magnetic recording technology. In MAMR, magnetization switches only when DC and microwave fields are applied to magnetic material at particular microwave frequency. Utilizing such a feature, possibility of selective switching of magnetizations was estimated by micromagnetic calculation for realizing 3-dimensional magnetic recording. In the calculation, a magnetic pillar composed of 3 layered materials was considered. Pillar diameter is 11 nm, and thickness of the 1st, 2nd, and 3rd layers are 5, 13, and 14 nm, respectively. Hk<sub>1</sub> (= 24kOe), Ms<sub>1</sub> (= 600 emu/cc), Hk<sub>2</sub> (= 13kOe), Ms<sub>2</sub> (= 400 emu/cc), and Hk<sub>3</sub> (= 6 kOe), Ms<sub>5</sub> (= 800 emu/cc) represent anisotropy fields and saturation magnetizations of 1st, 2nd, and 3rd layers. These values give sufficient thermal stability indices (greater than 60 at 300 K) for every layer. Magnetization switching probabilities equal to 1 were obtained at particular microwave frequencies for every layer, selective switching occurs at 1st, 2nd, and 3rd layer when microwave frequency is 8, 18, and 28 GHz, respectively. In this calculation, the single pillar has shown to store 3 bits information. These results imply MAMR has a potential in 3-dimensional magnetic recording.

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# QP10

### Thermal effect in microwave assisted magnetization reversal Yoshitoki Furomoto, Yuto Otsuka, Terumitsu Tanaka\* and Kimihide Matsuyama ISEE, Kyushu University, Japan

Microwave assisted magnetization reversal (MAMR)[1] has potential application in future magnetic recording. The critical curves for MAMR in a single-domain particle is theoretically estimated by Bertotti et al.[2, 3], which well explains steady state magnetization switching at 0 K and the equations allow three solutions "stable", "unstable" and "saddle". In this study, MAMR properties in a single-domain particle were simulated solving the Landau-Lifshitz-Gilbert equation taking account of temperature elevation due to ferro-magnetic resonance. Magnetization switching probability at 400 K widely distributes in HMW in unstable region. The critical values of HMW are greater than that at 0 K nevertheless thermal activation decreases energy barrier height. This implies stochastic thermal fields reduce efficiency in MAMR in the region. On the other hand, thermal activation gives 25-35 % of reduction in critical value of HMW in stable region, which is achievable for recording head of recent hard disk drive.

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### QP11

# Stable magnetization switching with microwave assistance for exchange coupled composite grain

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Microwave assisted magnetization reversal (MAMR)[1] is a promising candidate for future magnetic recording. Li et al. proposed ECC medium for solving the potential problems in microwave power and frequency for MAMR[2]. Bertotti et al. theoretically analyzed MAMR process taking a single spin for example[3], which showed the LLG equation can be written in the other forms[4]. The equation allows three solutions "stable", "unstable" and "saddle". However, the theory cannot be applied to exchange coupled composite (ECC) structured grain. In this study, distributions of the regions defined by the three solutions were estimated for an ECC grain. In the ECC grain, stable switching region, in which magnetization reversal modes are hardly affected by disturbance such as thermal fluctuation, distributes at lower fields region comparing to that for the single spin particle. This result means that stable magnetization reversal occurs with relatively weak magnetic fields even for ultra high Hk materials assigned to future high density recording.

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### QP12

### Low-temperature epitaxial growth of FePt on glass substrates for ultrahigh magnetic recording densities

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Magnetic FePt alloy thin film with ordered L10 structure is a promising material for perpendicular magnetic recording (PMR) application, duo to its huge uniaxial magnetocrystalline anisotropy (Ku), high saturation magnetization (Ms), high coercivity (Hc) and outstanding corrosive resistance. Ku value of L10 FePt is about 7 MJ/m3 which sets the superparamagnetic limit of L10 FePt as small as 2.8 nm. In PMR application the preferred orientation has to be transformed into (001) that makes the [001]-axis perpendicular to the film. Accordingly, developing an effective way to prepare L10 FePt (001) thin film with perpendicular magnetic anisotropy at low order-disorder transformation temperature is the concern of first priority.Although excellent properties can be obtained in sputtered FePt on single crystal MgO (001) much effort is needed to obtain PMR properties of sputtered FePt films on glass substrates. A number of underlayers have been used to induce texture such as CrRu, MgO and very recently TiN. Here we report a systematic study of the system { Hoya glass/CrRu/MgO/FePt/Ta} and through optimization of layer thickness, deposition temperature around 400 0C.

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# QP13

# Controlling nanostructure in FePt films: Co-sputtering of FePt and C or SiO,

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Achieving magnetic recording densities in excess of 1Tbit/in2 requires not only perpendicular media with anisotropy larger than 7 MJ/m3, which makes the FePt alloy as the best choice, but also narrow distribution and small sizes below 10 nm for a reduced S/N ratio along with high thermal stability. Previous work has shown that the L10 FePt grain size can be controlled by alloying FePt with materials such as C, Ag, and insulator like AlOx, MgO act to segregate and magnetically decouple the FePt grains. Better results were obtained with C with respect to the uniformity of grains and SiO<sub>2</sub> with respect to the shape, but far from optimization. We present our results by co-sputtering FePt with C or SiO<sub>2</sub> (up to 30 vol %) on MgO (001) single crystal substrates at 350 and 500 0C. With C or SiO<sub>2</sub> addition we achieved the grain size reduction, the shape control and isolated structure formation, by producing continuous films with high uniformity and narrow main size distribution down to 10 nm. This also reflected in the variation of coercivity as a function of the additional phase, thus giving us a possibility to simultaneously control the corecivity and the S/N ratio.

## QP14

# Retention time under currents and magnetic fields in a CoFeB/MgO perpendicular magnetic tunnel junction

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CoFeB/MgO based magnetic tunnel junctions with perpendicular anisotropy (p-MTJs) [1,2] have been investigated extensively because of their potential for realization of nonvolatile logic circuits and random access memories [3-6]. For ensuring their nonvolatility, effects of current (I) and magnetic field (H) on cumulative distribution of retention time (tr) of magnetization direction in the recording layer should be clarified. Theoretical expression for the distribution under I and H has been proposed [7]. However it has not been compared with experimental results in CoFeB/MgO p-MTJs. In this work, we investigated tr under I and perpendicular H in the same CoFeB/MgO p-MTJ with 40 nm in diameter and compared it with the theoretical expression [7]. It was found that the cumulative distribution of tr under I (H) can be fitted well with the theoretical expression in the range of tr = 30  $\mu$ s - 100 s (5 - 960 s). From tr under I (H), thermal stability factor is evaluated to be 44 (39), which agrees well with previous results in similar MTJs [1,4]. This work was supported by the FIRST program from JSPS.

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# QP15

# Thermal diffusion in magneto-optic collinear volumetric hologram memory

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Collinear holography memory is a good candidate of the next generation of optical memory. Recently we have investigated the magneto-optical recording with thick polycrystalline magnetic garnet films by collinear holography. Rewritable properties and shift multiplexing of the magnetic holograms were confirmed experimentally. However, it has still high noise level in the retrieved image. We investigated the suppression method of thermal diffusion in magnetic hologram recording numerically and experimentally. The recording principles of the magnetic holography are basically similar with thermomagnetic recording. The interference patterns with signal information are recorded as magnetization patterns. Due to thermal diffusion in recording process, the interference patterns were in disorder. It could be shown in results of the low diffraction efficiency and high noise level. Numerical simulation result shows the influence of thermal diffusion was suppressed with below 10-9 second of laser pulse width. Experimental result with two-beam interferometer also shows over 200 % higher diffraction efficiency recorded by 50 ps pulse width laser than 25 ns one in high spatial frequency range.

# QP16

# Resolution of magnetic garnet films for magneto-optic collinear volumetric hologram memory

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Collinear holography memory is promising optical memory of the next generation. The photopolymer materials have received much attention because of high sensitivity, high dynamic range and low costs as recording materials. On the other hands, magnetic garnet films have the advantage of rewritable, no necessity of shading, and no shrinkage, as compared to photopolymer which has write-once property. In this study, we fabricated the different grain sizes of magnetic garnet films by annealing time. The fundamental properties of samples and the SNR of the retrieved images with single hologram were investigated. Bi-substituted iron garnet films (3.3  $\mu$ m thick) were fabricated by RF magnetron sputtering on SGGG substrates. Samples were annealed at 750 °C in air for several annealing time. The retrieved images were evaluated by collinear holographic system. The resolution of magnetic recording media was related with the grain size of magnetic recording media. The results show SNR of the retrieved image was improved by reduction of the grain size.

**OP17** 

### Room temperature ordering perpendicular magnetic anisotropy L11 CoPtCu thin film on glass substrate

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Rhombohedral unit cell of L11 CoPt films have attracted a lot of attention due to their large magnetocrystalline anisotropy (Ku ~ 107 erg/cm3) as well as relatively low ordering temperature (~ 200 oC) [1,2]. The addition of Cu significantly improves the alignment of c-axis, chemical ordering and enhances perpendicular corcrivity (Hc<sup>L</sup>) of L11 CoPtCu phase deposited at 350 oC [3-5]. In this study,  $Co_{27}Pt_{50}Cu_{23}$  films with a Pt underlayer were deposited on glass substrates at room temperature (R1) to 500 oC by sputtering. The thickness of the film was 5 m. Interestingly, the high perpendicular magnetic anisotropy (PMA) was observed when the film was deposited at RT. Furthermore, PMA was maintained until the substrate temperature (Ts) exceeded 450 oC. Hc was largely enhanced from 0.25 to 2.6 kOe when 15 was increased from RT to 350 oC. Beyond 350 oC Hc started to decline. The appearance of high PMA suggests that the ordered L11 structure is formed in CoPtCu film even prepared at room temperature. This finding promotes the applications of L11 CoPtCu films in future perpendicular magnetic recording and spintronic device. The variation of phase structure and microstructure with Ts will be presented.

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## QP18

# Magnetic anisotropy and chemical ordering in Fe-Pt films prepared by low-temperature growth technique

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We investigated the chemical ordering and perpendicular magnetic anisotropy (PMA) in L1<sub>0</sub> FePt and L1<sub>2</sub> FePt<sub>3</sub> films. FePt films with different thickness of 35 and 70 A were grown, at the substrate temperature  $T_s$  = 300 °C by MBE co-evaporation and layer-by-layer growth, onto a MgO layer deposited on Si wafer. In order to optimize the growth of highly ordered FePt(001) films in conditions of low-temperature fabrication, the influence of a thin (001)-oriented *fcc* Pt buffer layer on the c-axis oriented texture was explored by high resolution TEM and X-ray diffraction (XRD). Our results showed a strong enhancement of PMA without high-temperature annealing, for 35 A thick FePt films, while a decrease in PMA was observed for L1<sub>2</sub> FePt<sub>3</sub>. By introducing a small percentage of Boron into the L1<sub>2</sub> ordered FePt<sub>3</sub>, enhanced structural ordering was seen in XRD. However, no significant change of magnetic properties was observed, in contradiction with the expectation to induce PMA by altering the chemical ordering in FePt<sub>3</sub> L1<sub>2</sub> phase. In this work, we discuss the origin of spin reorientation transition phenomena related to our results. \*corresponding authors : tachee@ewha,ac,kr; Anny.S.Michel@univ-poitiers.fr

# QP19

### Nanoscale ion beam etching process for reducing damage and leakage path of magnetic tunnel junction

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To integrate ultrahigh density spin-transfer torque magnetic random access memory (STT-MRAM), a stable process to pattern metallic magnetic tunnel junctions (MTJs) is required. [1] We investigated the ion milling process to fabricate the MTJ cells for perpendicular STT-MRAM application with minimized redeposition layer by a gradational stage tilting process and low damage of MTJ by reduced beam supply voltage. The hard mask used a 30nm diameter negative electron-beam resist patterned by electron beam lithography. Thereafter, ion milling process was conducted through regulation of beam/accelerate supply voltage and stage tilting angle. We found that the highest MTJ slope angle was attained at the stage tilt angle of 80° which was the angle of the incident ion flux to the stage surface. However, high angle etching resulted in a thick redeposition layer. For reducing redeposited layers, we have utilized variable-angle multi-step ion milling process. Also for reduced MTJ sidewall damage by high beam supply voltage, reduced the beam supply voltage. Consequently, sidewall damage of MTJ layer was minimized. The TMR ratio and the resistance of the fabricated MTJ cells were 13% and 1kΩ respectively. Our results showed the potential of using ootimized ion milling for fabricatine nanoscale MTJs for ultrahieh density STT-MRAMs.

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### OP20

### A new planer patterned media with anti-ferro / ferro transformation Hiroaki Ono and Hiroyuki Awano

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For next-generation HDD, the planer patterned media which need neither pattern etching nor smoothing processing is attractive.1),2) As a proposal of the new media, ferro-magnetized patterned area in the flat anti-ferro-magnetized layer has been prepared by using near field laser irradiation or masked ion implantation. On the ferro-magnetized area, required data can be recorded. We prepared several kinds of TbFeCo/ Rh/TbFeCo tri-layers. When the Rh thickness is some period, strong anti-ferro-magnetic coupling (AFC) is realized. In the combination, Rh thickness of the 1st peak of the anti-ferro-coupling was 0.45nm. Since the thickness of the TbFeCo was thick and Rh thickness was very thin, the transformation from AFC to FC can be control easily. The coercive force was 1.5kOe. Also the Hc can be designed by controlling Tb concentration. For example, in this case, the transformation temperature was 100 °C.It is also easy to control by changing TbFeCo thickness. Moreover, the patterned media can be applied as a write once HDD, since it is impossible to change recorded area.

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## QP21

# Topology optimization of perpendicular magnetic recording head considering magnetic saturation effect

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The topology optimization method is an attractive scheme for magnetic device design because an initial conceptual design can be obtained without any prototype model. However, studies concerning the three-dimensional magnetic device design by topology optimization are rare. Even more, topology optimization of magnetic systems has been generally accomplished by using the two-dimensional finite element method while the nonlinear saturation effect of the material property between magnetic field strength and magnetic flux density has not been fully considered. This study proposes a topology optimization technique of the three-dimensional design as well as the two-dimensional design based on the density method considering the nonlinear saturation effects. The proposed method makes it possible to provide useful and practical process for magnetic system design any lead to improved design with light system weight. The proposed method is applied to the optimization problem of a magnetic recording system. During the design variables are accomplished by the adjoint variable method and the sequential linear programming, respectively.

Topology optimization; finite element analysis; magnetic saturation effect; three dimensional design; density method; adjoint variable method.

# QP22

# The effect of capped layer thickness on switching behavior in coupled granular/continuous media

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Coupled granular/continuous (CGC) media (a CoPtCrB continuous layer exchange coupled to a granular oxide CoPtCr: SiO<sub>2</sub>TiO<sub>2</sub>Ta<sub>2</sub>O<sub>3</sub> layer) exhibited improvement of SNR, thermal stability and switching field distribution (SFD)[1, 2]. TEM images together with intrinsic SFD for capped structure are studied to investigate the mechanism of the SFD reduction in CGC media. We find out that in the bottom part, CoPtCr magnetic grains are separated by nonmagnetic oxide grain boundaries. The amorphous grain boundary phase in the granular layer propagates to the top surface of the capped layer. The grain size increases with capped layers thickness. A critical capped thickness exists. At thinner capped thickness, capped layer inherits structural irregularity from the granular layer, while at thicker layer, the upper portion of capped layer is less influenced by the granular layer. ReportedAH (M, AM) method[3] is used to calculate intrinsic SFD. Both intrinsic and macroscopic SFD decrease linearly with capped layer thickness. As we know, the intrinsic SFD is due to the local variations of the grains properties. The reduction of macroscopic SFD is caused by homogenous grains and uniform exchange coupling at thicker capped layer.

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### QP23

# The effect of magnetic field on FePt nanoparticles during annealing process

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FePt, L10 nanocrystals are expected as one of promising candidates for the magnetic recording media with ultrahigh densities. The 3nm FePt NPs have been synthesized by the polyol method and analysed by TEM analysis. The composition of the films was estimated to be approximately  $Fe_{cx}$   $OPt_{37:00}$  by EDXS. Surface of Si and SiO<sub>2</sub> wafers have been covered by FePt nanoparticles and heated in 600° for 1h to form L10 phase nanoparticles [1], FE-SEM and XRD analyses show uniform surface distributions of FePt, L10 nanoparticles. The SiO<sub>2</sub> [2] substrate prevents nanoparticles from coalescence during annealing and prevents them from growth [3], hence, causing their size not to become more than 40nm. SiO<sub>2</sub> wafer is more suitable than Si wafer in separating the nanoparticles, which may be due to oxygen of surface bonding with nanoparticles. Presence of 30mT magnetic field controls the size of nanoparticles around 25nm, which can be reduction of random mobility in magnetic field and causes the nanoparticles satis of ranoparticles is aligned in this direction [4]. The magnetic field also controls the surface distribution during heating process.

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## QP24

# Novel soft lithography technique for fabrication of Ni nanodots for use as bit patterned media

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Magnetic disk drive technology has successfully reduced the size of multi-grain bits

to ~30 nm, and there are intensive industrial efforts to shrink the bit size further. This work reports the development of a cheap, scalable polymer template lithography technique of fabricating large arrays of magnetic nanostructures towards use as bit patterned media, one of the prominent next generation data storage technologies. The technique has been standardized and physical properties such as bit sizes, pitch sizes and the storage densities have been calculated. Critical magnetic properties such as magnetization reversal have also been studied and discussed to ascertain the plausibility of its use in futuristic high density magnetic data storage devices

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# RA01

### Magnetic resonance and ferroelectricity of BaTi<sub>1-x</sub>Fe<sub>x</sub>O<sub>3</sub>

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BaTiO<sub>3</sub> is known as an excellent candidate for multiferroics. In this work, we present magneticresonant and ferroelectric properties of polycrystalline BaTi<sub>1x</sub>Fe<sub>x</sub>O<sub>3</sub> ceramics ( $0 \le x \le 0.12$ ) prepared by conventional solid-state reaction. Measurements have been based on an X-ray diffractometer, electron spin resonance (ESR) spectroscopy, and Precision Premier-Radiant. X-ray diffraction patterns reveal the tetragonal-hexagonal phase transformation as increasing Fe-doping content. ESR spectra prove that many samples to be ferromagnetic, excepting for the  $x = 0.02 \cdot 0.06$  samples, which are paramagnetic. Such the results are in good agreement with the previously report on the magnetic behavior. Room-temperature P-E hysteresis loops of BaTi<sub>1x</sub>Fe<sub>x</sub>O<sub>3</sub> reveal that with increasing Fe concentration, the feature of ferroelectric P-E loops changes in the shape to the paraelectric regime, indicating the decrease in ferroelectricity. Consequently, the values of the remnant polarization Pr and the coercive field Ec gradually are also decreased. Keywords: Fe-doped BaTiO<sub>3</sub>; Magnetic phase separation, Ferroelectricity \*Electronic mail: ptlong2512@yahoo.com Phone:

Daniel Khomskii, Physics, 2.20 (2009) W. Eerenstein, N.D. Mathur & J.F. Scott, Nature 442(17) (2006) N.V Dang, Thanh, L.V.Hong, V.D Lam, The-Long Phan, Applied Physics 110, 043914, (2011). QIU Shen-Yu, LI Wang, LIU VU, LIU GUi Hua, WU Y-Qiang, Chen Nan, Trans. Nonferrous McSos. China 20,1911-1915(2010), Manfred Fiebig, J. Phys. D: Apply:Phys.38, R123-R152. S.Angappane, G Rangarajan, and K. Sethupathi, J. Apply. Phys.93,8334 (2003).F. Lin, D. Jiang, X. Ma, and W. Shi, Physica B 403, 2525 (2008). S. Ray, P. Mahadevan, S. Mardal, S. R. Krishnakamar, C. S. Kuroda, T. Sasaki, T. Taniyama, and M. Ioh, Phys. Rev. B 77, 104416 (2008).

## **RA02**

### Growth of high pure BiFeO<sub>3</sub> single crystals

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BiFeO<sub>3</sub>(BFO) is a heavily beaten compound in recent publications. BFO exhibits simultaneous polar and magnetic order at room temperature, promising important applications, but growth of pure BFO single crystals is really a great challenge. To grow extremely pure BFO single crystals, we have systematically investigated  $Bi_2O_3$ -Fe<sub>2</sub>O<sub>3</sub> pseudo-binary system from the long-term crystal growth experiments, high performance powder X-ray diffraction analysis and high temperature differential thermal analysis instrument. The peritectic transition temperatures in  $Bi_2O_3$ -Fe<sub>2</sub>O<sub>3</sub> pseudo-binary phase diagram were reevaluated. It has been found that purity of grown BFO single crystals depends greatly on crucible materials. When suitable kind of crucibles were used, extremely high pure BFO single crystals could be easily obtained with maximum size above one centimeter.

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# **RA03**

# Structural, electrical and magnetic properties of La-doped BiFeO<sub>3</sub> ceramics.

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Polycrystalline ceramic samples of Bi<sub>1-x</sub>La<sub>x</sub>FeO<sub>3</sub> (x=0.0, 0.1 and 0.2) have been prepared by standard solid state reaction method using high purity oxides and carbonates. The structure and formation of the compounds have been checked by X-ray diffraction followed by energy dispersive X-ray microanalysis (EDAX) techniques. There is a good agreement between observed and calculated X-ray diffraction patterns obtained by performing the Rietveld refinement with a structural model using the noncentrosymmetric space group R3c. The lattice parameters have been refined but the over all structure remained the same. The microstructural studies have been carried out using scanning electron microscopy. The dc conductivity of all the samples has been measured and their activation energies were calculated from  $\log\sigma$  vs 103/T curves. Vibrating sample magnetometer (VSM) has been used to study the magnetic behaviour of the compounds. It has been observed that with the increase of substitution concentration of La, the insulating behaviour of the materials have been improved and showing the antiferromagnetic to weak ferromagnetic behaviour. The results are discussed in detail.

### **RA04**

## In situ X-ray absorption spectroscopy study on $BaTiO_{\rm 3}$

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Ferroelectric systems are under intensive study for both the point of view of fundamental materials physics and the potential applications such as memories, sensors, and transducers. The Ti<sup>4+</sup> (d0-ness) systems which are the generally well known ferroelectric as form ATiO<sub>2</sub> structure (A=Ba and Pb), inter alia, have been extensively investigated. To understand this system, both of structural and electronic investigations must be completed together. Structurally, Ti ions in the TiO<sub>4</sub> octahedra are moved along the applied electric (E)-field, resulting in a symmetry lowing. However, we still do not understand much about a change in the electronic structure between Ti 3d and O 2p states. Although the soft x-ray spectroscopy study at the oxygen absorption edge that is regarded as a good approach for exploring the electronic structure, such study accompanies an experimental difficulty: electron motion distortion under in situ E-field condition. In this work, we introduce an approach for overcoming such difficulty in conventional x-ray absorption spectroscopy (XAS) measurement. By looking at the O K-edge XAS of BaTiO, via a fluorescence vield, we studied electronic structure and reveal that the observed spectra features are related with the structural changes under the applied E-field.

### RA05

**Chiral skyrmions and magnetic bubbles in multiferroic materials** Xiuzhen Yu<sup>1</sup>\*, Y. Tokunaga<sup>1</sup>, S. Seki<sup>2</sup>, Y. Kaneko<sup>3</sup>, S. Ishiwata<sup>2</sup>, M. Mostovoy<sup>4</sup>, N.

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Nanometric chiral skyrmions and magnetic bubbles have been successfully realized in multiferroic materials  $Cu_2OSeO_3$  and Sc-doped barium ferrite by means of Lorentz transmission electron microscopy. Under the normal magnetic fields, chiral skyrmions with a topological number of -1 and a unique chirality, 50 nm in diameter, emerge in a non-centrosyemmetric helimagnet  $Cu_2OSeO_3$ . In analogy with the transition from a helical spin structure to chiral skyrmions, a normal magnetic field turns the magnetic stripes into a bubble lattice with a positional order and random chiralities in a hexaferrite BaFe<sub>0.79</sub>Sc<sub>0.16</sub>Mg<sub>0.06</sub>O<sub>3</sub> or mini diameter, emerge in a non-centrosyemmetric helimagnet chiral turns the magnetic stripes into a bubble lattice with a positional order and random chiralities in a hexaferrite BaFe<sub>0.79</sub>Sc<sub>0.16</sub>Mg<sub>0.06</sub>O<sub>3</sub> or mini alisotropy. By systematically investigating the dependences of magnetic configurations on magnetic field, temperature and sample thickness, we reveal the differences between the spin textures of chiral skyrmions induced by Dzyaloshinskii-Moriya interaction, however, show less size dependence on the external factors. Second is a difference in chirality. The chiral skyrmions have a unique chirality determined by the crystal chirality, while magnetic bubbles have two opposite chiralities, resulting in a variety of zoological magnetic structure with varying external magnetic field.

### RA06

### Studies on Bi<sub>1-x</sub>La<sub>x</sub>FeO<sub>3</sub> crystals in pulsed high magnetic fields Masashi Tokunaga<sup>1</sup>, Mitsuru Akaki<sup>2</sup>, Hideki Kuwahara<sup>2</sup>, Kengo Oka<sup>1</sup> and Takumi Kihara<sup>1</sup>

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Department of 1 hysics, sophia Oniversity, supan

BiFeO<sub>3</sub> is known as the unique material that shows multiferroic nature at room temperature [1,2]. Application of high magnetic fields to this material induces a magnetic transition accompanied by a significant change in the electric polarization [3-5]. Although the occurrence of this transition is explained in the finamework of the Ginzburg-Landau theory, in detail study is highly desirable to understand the microscopic origin of the magnetoelectric effects in this material. In this work, we synthesized crystals of Bi<sub>1</sub>, La, FeO<sub>3</sub> with using the flux method, and studied their magnetic and dielectric properties in pulsed magnetic fields up to 55 T. Early experiments on the polycrystalline samples showed that partial substitution of La for Bi reduces the magnetic field needed to cause the magnetic transition [6]. Our isothermal magnetization measurements showed that the transition field at 4.2 K changes from 18 T to 15 T by La substitution. Our dielectric study revealed that the electric polarization shows significant change accompanied with this magnetic transition. Therefore, the phenomena observed in BiFeO<sub>3</sub> are reproduced in the reduced field scale in the La substituted system, which enables us more systematic studies to clarify the microscopic origin of the magnetolectric effects associated with the magnetic transition.

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# RA07

### Enhanced magnetic properties of Ni-doped BiFeO<sub>3</sub> compounds

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The structure and the magnetic properties of polycrystalline  $BiFeO_3$  and  $BiFe_{0.95}Ni_{0.05}O_3$ bulk ceramic compounds which were prepared by solid-state reaction and rapid sintering were investigated High-purity Bi<sub>2</sub>O<sub>2</sub> Fe<sub>2</sub>O<sub>4</sub> and NiO powders were mixed with the stoichiometric proportions, and calcined at 450°C for 24 h to produce BiFe, Ni.O., The crystalline structure of samples was investigated at room temperature by using a Rigaku Miniflex powder diffractometer. The field-dependent and the temperature-dependent at high temperatures magnetization measurements were performed with a vibratingsample magnetometer. The temperature-dependent magnetization at low temperatures was analyzed with Ouantum Design superconducting quantum-interference-device (SOUID) magnetometer. The x-ray diffraction study demonstrates the compressive stress due to the Ni substitution at the Fe site. The lattice constant of BiFeagNiagO<sub>2</sub> is smaller than of BiFeO<sub>3</sub> because of the smaller ionic radius of Ni<sup>3+</sup> than that of Fe<sup>3+</sup>. The fielddependent magnetization of BiFeoseNioosO2 exhibits a clear hysteresis loop at 300 K. The magnetic properties of BiFe<sub>0.95</sub>Ni<sub>0.05</sub>O<sub>3</sub> were improved at room temperature because of the existence of structurally compressive stress. The zero-field-cooled temperature-dependent magnetization reveals that the magnetic transition temperature and the magnetic moment of BiFe<sub>0.95</sub>Ni<sub>0.05</sub>O<sub>3</sub> are higher than those of BiFeO<sub>3</sub>.

### RA08

# Magnetic modulation of electrical impedance in Bi-doped $La_{0,7}Sr_{0,3}MnO_3$

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In recent years, magnetic control of dielectric polarization in multiferroics and magnetoelectrica has become a topic of considerble interest. Dielectric studies in these materials are often done using an impedance analyzer in the frequency range f = 10 to 10 kHz. As the frequency increase, dieletric permitivity decreses. Here, we show that an interesting phenomenon occurs in MHz range in Bi doped La<sub>0.7</sub>Sr<sub>0.3</sub>MnO<sub>2</sub>. With increasing frequency of the ac current or ac voltage, both in- and out-of phase compnents of electrical impedance show an abrupt increase that is accompanied by a peak around the Curie temperature. The peak decreases in magnitude and is suppressed in a field of 1 kOe. This results in a huge ac magnetoresistance (30-35 % for H = 1 kOe) and magnetoreactance (= 30-40 %). The origin of the magnetoreactance is not necessarity due to change in capacitance but but it is suggested to magnetoindutance effect, which has not been seriously considered so far. The combination of huge ac magnetoresistance and magnetonductance in a single material provides multifunctional capabilities. Results from our studies are compared to magnetocapacitance behavior found in BiMnO.[1]

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### RA09

# Theoretical study of PCAR measurement with d-wave superconductors.

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High spin polarization ferromagnetic metals are essential for creating high performance spintronics devices, therefore it is important to estimate the value of spin polarization quickly and efficiently. Point contact Andreev reflection (PCAR) spectroscopy is a powerful method for measuring spin polarization of ferromagnetic metals. However, most PCAR measurement were performed at very low temperatures, around 1.5 K, since the conventional s-type superconductor is used. If we use high-Tc superconductors, PCAR measurement can be performed at higher temperature. In this study, we theoretically analyze the Andreev reflection in the ferromagnetic metal (FM) / insulator (I) / d-wave superconductor (dS) junctions. The superconducting gap in momentum space is not uniform in the d-wave superconductor, Especially, in the case of  $d_{x2y2}$  model, we found that the bias voltage at which the normalized conductance of the FM/I/dS junction drops is expressed as a function of the exchange field of the ferromagnetic metal.

# RA10

# $T_{\rm C}$ Evolution of Bulk and Optical Spectra of Nanocolloidal Fe-doped Manganate CaMn<sub>1-x</sub>Fe<sub>x</sub>O<sub>3</sub> (x = 0, 0.01, 0.03, 0.05)

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Recently, the effect of magnetic reversal at low magnetic field was observed with the multiferroic Fe-doped CaMnO<sub>3</sub> [Sol. Stat. Commun. 142 (2007) 525]. We report here the optical behaviors of this class of compounds (CaMn<sub>1-x</sub>Fe<sub>x</sub>O<sub>3</sub>, x = 0, 0.01, 0.03, 0.05) which were prepared by using traditional ceramic method with parent oxides as precursors. The analysis of structure showed the predominant orthorhombic phase with slight increased cell constants according to doping content (from 3.743 to 3.746 Å). The investigation of the Raman spectra also agreed with the increased Mn-O bond length according to doping, therefore suggested the weakening of ferromagnetic exchange between Mn<sup>3+</sup> and Mn<sup>4+</sup> cations (exactly, from 0.61 to 0.52 eV). This weakening also developed together with the reduction of Curie temperature (T<sub>c</sub>) (from 155 to 135 K), the shifts of infra-red (IR) absorption maxima towards the longer wavelengths and the narrowing of band-gaps (from 0.45 to 0.14 eV). The absorption spectra also showed a clear absorption line near 1.2 eV which is characteristic for Fe, and is composed mainly from 3d electron transition with a change in spin.

PACS number: 75.47.Lx; 75.50.Ee; 74.25.Fy; 75.30.Kz Keywords: perovskite, manganate, structure, optical

## RA11

### Evidence of magnetic phase separation in LuFe<sub>2</sub>O<sub>4</sub>

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The magnetic properties of a polycrystalline sample of multiferroic LuFe<sub>2</sub>O<sub>4</sub> [1] have been investigated by means of magnetization measurements combined with neutron scattering, between 1.5 and 300K [2]. Magnetic Bragg peaks appear below  $T_{\rm N} = 240$ K, and can be indexed using two propagation vectors, corresponding to either a ferromagnetic or an antiferromagnetic stacking of the iron bilayers. The magnetic arrangement in the ab plane follows a 1:2 pattern, which obeys the charge ordering one observed in transmission electron microscopy [3]. Neutron inelastic scattering data confirm the strong easy-axis single-ion anisotropy of the iron spin in its trigonal environment, of about 8 meV at 5 K. Modeling of the diffraction data, based on a CO-type modulation of the spin ordering, shows coexisting magnetic phases with opposite signs of the intra-bilayer interaction, corresponding to two distinct magnetic transition temperatures. Coexisting antiferromagnetic and ferrimagnetic phases, along with the enlarged axial magnetic anisotropy, are necessary to describe the macroscopic magnetic properties of this compound.

 N. Ikeda et al., Nature (London) 436, 1136 (2005) [2] J. Bourgeois et al., submitted to Phys. Rev. B (2012) [3] J. Bourgeois et al., to appear in Phys. Rev. B (2012)

## RA12

# Magnetoelectric effect in mechanically mediated structure of TbFe<sub>2</sub>, Pb(Zr,Ti)O<sub>3</sub>, and nonmagnetic flakes

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For most investigations on layered ME composites with excellent ME effect compared to single phase materials and multiphase bulk composites, the classical structure is constructed layer by layer through interface bonding. In this work, the proposed ME structure is made up of two magnetostrictive TbFe<sub>2</sub> alloys, one piezoelectric PZT ceramic and two nonmagnetic glass flakes. When an external ac magnetic field Hac is applied along the length direction of magnetostrictive TbFe<sub>2</sub> flakes, the TbFe<sub>2</sub> flakes will produce an extensional strain. The strain then transfer to the PZT flake because of the end parts of PZT and TbFe<sub>2</sub> flakes bunded with glass slides in a vertical plane. Due to the piezoelectric effect, the PZT flake will induce a charge output. The interface between magnetostrictive and piezoelectric phases is not required to achieve ME coupling. Increasing H<sub>dc</sub> from zero,  $\alpha_{E,31}$  increases linearly until a maximum value is reached at 350 Oe, and then decreases subsequently. With further increase in Hdc,  $\alpha_{E,31}$  reaches a maximum value, and then decreases subsequently. There are several resonance peaks at 200 Oe and 1<f<150 kHz and a giant ME voltage coefficient as high as 2.7 Vcm<sup>-1</sup> Oe<sup>-1</sup> can be obtained at the resonance frequency of 44.5 kHz.

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# RA13

# Enhancement of magnetization in sulfur doped BiFeO<sub>3</sub>

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Sulfur doped BiFeO<sub>3-x</sub>S<sub>x</sub> are synthesized by high pressure solid state reaction method with the high purity original materials of Fe<sub>2</sub>O<sub>3</sub>, Bi<sub>2</sub>O<sub>3</sub>, Bi<sub>2</sub>O<sub>3</sub>, S, and FeS. The pressure remains 5 GPa during the sample preparations. We have compared the effects of original materials on the final products, and got the optimum products with Fe<sub>2</sub>O<sub>3</sub>, Bi<sub>2</sub>O<sub>3</sub>, and FeS as the original materials. Slight impurity of Bi<sub>2</sub>Fe<sub>4</sub>O<sub>9</sub> is detected in the products. A magnetization enhancement is confirmed to roughly more than 2 orders with the above materials reaction. The saturation magnetization more than 0.75 emu/ g at 1.5 Tesla, and the Hc is roughly 50 Oe. The good ferroelectric property is also predicted. We argue that the enhancement of magnetization may due to the variation of the chemical valence of iron.

# RA14

# Effects of chlorine and fluorine on the structure and magnetism in BiFeO<sub>1</sub>

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Clorine and fluorine doped BiFeO<sub>3.x</sub>M<sub>x</sub> are synthesized by solid state reaction method under high pressure and high temperature with the high purity original materials of Fe<sub>2</sub>O<sub>3</sub>, Bi<sub>2</sub>O<sub>3</sub>, and FeCl<sub>2</sub>, or FeF<sub>2</sub>. The pressure remains 6 GPa and the temperature is set to be 750 oC during the sample preparations. They are almost the single phase. We have compared the chlorine and fluorine effects on the structure and magnetization on the final products. Both doping samples show the weak magnetism at room temperature, much larger than the pure BiFeO<sub>3</sub>, although the paramagnetic background is strong. Chlorine doping show much larger magnetization than that of fluorine. We argue that the magnetization results from the variation of the chemical valence of iron ion and the distortion of FeO<sub>6</sub> due to the element of oxygen replaced by fluorine or chlorine. The dielectric properties show us the small leakage may indicate the increase of carriers. We expect a good dielectric properties may obtained by adjust the dopants in BiFeO<sub>3</sub>.

## RA15

# Investigation of electricity coercive behavior of LSFMTO system using ultrasonic mixing method

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We have investigated electricity coercive behavior of  $La_{0.5}Sr_{0.3}Fe_{0.9}Mn_{0.05}Ti_{0.05}O_3$ (LSFMTO) system using ultrasonic mixing method. LSFMTO system is driven by 40 kHz frequency and calcinate process up to 750°C for 30 minutes. Then sintering process also carried out by systematic various of temperature 900 °C, 1000 °C, and 1100 °C respectively, with time preparation is same previously. Very interestingly, the LSFMTO system exhibits ferroelectric (La,Sr)(Fe,Mn,Ti)O<sub>3</sub> and ferromagnetic SrFe<sub>12</sub>O<sub>19</sub> properties and minor phase La<sub>2</sub>O<sub>3</sub> by XRD analyzation. We have also carried out measurement of LSFMTO system using a systematic applied the external magnetic field and electric polarization. The hysteresis curve of LSFMTO system shows the electricity coercive (Ec) increases while saturation polarization (Pc) and remanent polarization are tendency constant at the magnetic strength of 0.12 T. I believed that our finding of LSFMTO system is potensial to multiferroic material.

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## **RA16**

# Enhancement of multiferroic properties of solid state prepared La doped $BiFeO_{\rm 3}$

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BiFeO<sub>3</sub> (BFO) has generated widespread interest since it is a room-temperature multiferroic. La<sub>x</sub>Bi<sub>1x</sub>FeO<sub>3</sub> (x= 0.05-0.4) are prepared through solid state reaction method. X-ray diffraction for x  $\leq$  0.15 shows R3c and for x> 0.15 the structure changes to Pbnm. The percentage of these phases present is estimated through Rietveld analysis. Raman spectra for x  $\leq$  0.15 exhibit 8 peaks (4A+4E modes). There is no shift in Raman modes indicating that La substitutes Bi in BFO. The dielectric constant increases and leakage current decreases with x. The enhanced polarization may be due to the distortion of the crystal lattice with La doping. With increase in La doping the  $T_N$  decreases. The magnetization (M) decreases up to 250 K and increases on further cooling indicating the spin reorientation. A cusp in M-T is observed close to 50 K indicating the spin glass behavior. M-H loops do not saturate even for field up to 6 T. With La doping,  $H_C$  and the magnetization increases dramatically in comparison to x=0. Both Hc (~ 1.6 T) and Mr (0.27 emu/g) have maximum values for x =0.2 at 5 K. The increase in the magnetic parameters is attributed to the break in the spin cycloid structure.

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# RA17

## Magnetostructural coupling at the metal-insulator transition in YBaCo<sub>2</sub>O<sub>5.5</sub> as seen by synchrotron x-ray diffraction and absorption Jessica Padilla-pantoja, Javier Herrero-martin, Carlos Frontera and Jose Luis Garcia-

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In LnBaCo<sub>2</sub>O<sub>3.5</sub> (Ln: lanthanide) cobaltites Ln and Ba cations order in alternating planes along the c-axis yielding an ordered arrangement of CoO<sub>5</sub> pyramids and CoO6 octahedra. The presence of two local environments for Co may be at the origin of proposed Co spin-state changes whose interplay with the structural and electronic changes at metal-insulator transitions (MIT) are currently attracting much attention [1-5]. We present a study of the magnetic, structural and thermoelectric properties of YBaCo<sub>2</sub>O<sub>3.5</sub> as a function of temperature. Our results evidence the presence of three magnetic transitions with a strong magnetostructural coupling. We have used synchrotron x-ray powder diffraction to show that the Pmmm orthorhombic structure transforms (on cooling) into a monoclinic (P112/a) phase below  $T_{MI} \sim 295$  K by doubling the a-axis [6]. Concomitantly a ferromagnetic moment appears and the Seebeck coefficient is enhanced. A further cooling leads to the successive appearance of Co AFM structures ( $T_{NI} \sim 267$  K,  $T_{NZ} \sim 231$  K), the crystal structure becoming again orthorhombic. Finally, we present temperature dependent soft x-ray absorption data at the Co L<sub>3</sub> edge. They discard the possibility of a high to low spin state change in Co ions at octahedral sites across the MIT [7].

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### **RA18**

### Magnetic field dependence of dielectric properties in LuFe<sub>2</sub>O<sub>4</sub> Takashi Kambe\*, Yukimasa Fukada, Tomoko Nagata and Naoshi Ikeda *Physics, Okayama university, Japan*

The multiferroic properties in LuFe<sub>2</sub>O<sub>4</sub> are of interest, in which contains equal amount of Fe<sup>2+</sup> and Fe<sup>3+</sup> in triangular lattice. This material shows a ferri-magnetic transition at 240K and three-dimensional charge ordering at around 320K. The dielectric property of this material is suggested with a valence fluctuation process of iron ions in the polar charge ordered domain. Here we report a magnetic field variation of the dielectric constant of single crystal LuFe<sub>2</sub>O<sub>4</sub>, and discuss with the valence fluctuation process of iron charges in the manetic domain wall.

# RA19

**Pressure studies of LaAgSb<sub>2</sub> utilizing new integrated pressure cell** Sven Friedemann<sup>1</sup>\*, Zhou Feng<sup>1</sup>, Takao Ebihara<sup>2</sup>, Christophe Thessieu<sup>3</sup> and F Malte Grosche<sup>1</sup>

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### RA20

# Crystal & Magnetic structure studies of doped BiFeO<sub>3</sub> multiferroic compound

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Multiferroics, materials combining multiple order parameters, offer an exciting way of coupling phenomena such as electronic and magnetic order. The coexistence of different order parameter permits potential applications in information storage, spintronics, and magnetic or electric field sensors. The perovskite BiFeO<sub>3</sub>(BFO) is known to be antiferromagnetic below the Neel temperature of 640 K and ferroelectric with a high Curie temperature of 1100 K. According to the previous doping studies of BFO<sub>3</sub> it is likely that non-stoichiometry and second-phase formation are the factors responsible for leakage current in BFO. It has been suggested that oxygen nonstoichiometry leads to valence fluctuations of Fe ions in BFO, resulting in high conductivity. To reduce the large leakage current of BFO<sub>3</sub> one attempt is to make donor-doped BFO compounds. In this study, we have tried to generate the new multiferroic material functioning on at room temperature. The candidate systems are BiFeO<sup>3+</sup>BaInO<sub>3</sub> and BiFeO<sup>3+</sup>BaITO<sub>3</sub>. These system have been fabricated by a solid-state reaction method and flux method, respectively. The crystal and magnetic structure of these systems are studied using XRD and neutron powder diffraction as functions of temperature. We will discuss the magnetic and electric property of new multiferroic system working on at room temperature.

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### RA21

# The analysis of structure, magnetic and ferroelectric properties of $Ba_{1\_x}Bi_xTi_{0.95}Fe_{0.05}O_3$

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The multiferroic material is a material which shows more than one primary ferroic order parameter simultaneously and can be made by introducing the magnetic impurities into ferroelectric material [1]. BaTiO<sub>3</sub> is known as a classical ferroelectric material, whereas the Fe-doped BaTiO<sub>3</sub> is reported to have a magnetic ordering. Although the Fe doping can induce the ferromagnetism in BaTiO<sub>3</sub>, the ferroelectric properties of Ba<sub>1-x</sub>Bi<sub>x</sub>T<sub>0.95</sub>Fe<sub>0.05</sub>O<sub>3</sub> ( $0 \le x \le 0.05$ ). XRD results indicate that the tetragonal phase is merged with the hexagonal phase for x < 0.05. At x = 0.05, the pure tetragonal phase is observed. The magnetic field dependence of magnetization at room temperature for x = 0 shows a kind of antiferromagnetism. Introducing 2 % Bi doping, it shows the ferrimagnetic properties. However, with the 5 % doping of Bi, the ferrimagnetic loop is disappeared and becomes paramagnetic. It suggests that the superexchange between Fe<sup>3+</sup> ions in otahedral and pentrahedral sites associated with oxygen vacancy produces the ferromagnetism [2]. The ferroelectric hop with low leakage current is only observed for x = 0.05. Reservents indicate that the multiferroic properties are observed in Ba<sub>0.98</sub>Bi<sub>0.02</sub>Ti<sub>0.95</sub>Fe<sub>0.05</sub>O<sub>3</sub> sample.

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# RA22

### Spin dynamics of multiferroic BiFeO3 single crystal

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Multiferroic compounds are one of most challenging topics in the condensed matter physics [1] not only for its fundamental importance but also for huge future applications. Although several systems are known to have multiferroic behavior, there are few multiferroic materials that have both magnetic and ferroelectric transition temperatures above room temperature. BiFeO<sub>3</sub> is probably the only exception with Neel temperature at  $T_c$ =f50 K and Curie temperature at  $T_c$ =1100 K [2]. In order to understand the spin dynamics of BiFeO<sub>3</sub> we have carried out inelastic neutron scattering experiments on ten co-aligned single crystals of BiFeO<sub>4</sub> with Avee taried out inelastic neutron scattering experiments on ten co-aligned single crystals of BiFeO<sub>3</sub> we have carried out metastic neutron scattering experiments on ten co-aligned single crystals of BiFeO<sub>4</sub> with Avee taried out metastic neutron scattering experiments on ten co-aligned single crystals of BiFeO<sub>4</sub> with avee and the full three-dimensional spin wave [3]. We have also calculated the spin wave dispersion using Heisenberg Hamiltonian with two exchange parameters between the nearest and next nearest neighbors. A Dzyaloshinskin-Moriya term that arises from the spiral magnetic structure of BiFeO<sub>3</sub> was also examined. By carefully examining the AMATERAS and MERLIN data, we could find, for the first time, the full spin wave dispersion of BiFeO<sub>3</sub> and detarmine the schange parameters twith the experimental results.

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## RA23

### Magnetic and dielectric properties of the single crystals Sm<sub>1</sub>. xHo<sub>x</sub>Fe<sub>3</sub>(BO<sub>3</sub>)<sub>4</sub> and Sm<sub>1-x</sub>La<sub>x</sub>Fe<sub>3</sub>(BO<sub>3</sub>)<sub>4</sub>

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The trigonal rare-earth ferroborates with hantit ctructures is interested due to very strong of the electric polarization induced by magnetic field. The research of single crystals with double rare-earth subsystems such as Sm<sub>1</sub>, R, Fe<sub>3</sub>(BO<sub>3</sub>)<sub>4</sub> (R = Ho, La) is important for analysis microscopic mechanism of magnetoelectric interaction. The single crystals of Sm<sub>1</sub>, Ho<sub>2</sub>Fe<sub>3</sub>(BO<sub>3</sub>)<sub>4</sub> with x = 0; 0.3; 0.5 and Sm<sub>1</sub>, La, Fe<sub>3</sub>(BO<sub>3</sub>)<sub>4</sub>, with x = 0.25; 0.5 were grown by a flux method. The magnetic and magnetodielectric properties have been investigated. It has been revealed what magnetic behaviour and magnetodielectric effect essential depended on x. The magnetic field influence on the change in degree of dielectric penetrability is decreased with increasing x for both under study families. The decreasing of magnetodielectric effect is probably related with influence of the racemic twinning.

### RA24

### Far infrared spectroscopy of EuFe<sub>3</sub>(BO<sub>3</sub>)<sub>4</sub>

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We present the results on the polarized far infrared absorption spectra of EuFe<sub>3</sub>(BO<sub>3</sub>)<sub>4</sub> crystal in a wide range of temperature (3.5 - 300K). Rare-earth iron borates RFe<sub>3</sub>(BO<sub>3</sub>)<sub>4</sub> undergo the antiferromagnetic phase transition at temperatures below 40 K, and some of them exhibit a significant magnetoelectric and magnetoelastic effect. EuFe<sub>3</sub>(BO<sub>3</sub>)<sub>4</sub> is known to exhibit two phase transitions, the structural one at  $T_c = 58$  K [1], and the antiferromagnetic ordering at  $T_N = 34$  K [2]. The phonon frequencies demonstrate pronounced changes at  $T_C$  consistent with the observed peculiarities in the dielectric constant and thermal expansion [3]. Smaller but clearly visible phonon shift were registered below  $T_N$  thus evidencing the spin-lattice interaction in the title compound.

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# RA25

# Electric polarization, toroidal moments, spin chirality, spin canting at avoided level crossing induced by Dzyaloshinsky-Moriya interaction in $V_3$ nanomultiferroics in transverse magnetic field

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The spin-frustrated V<sub>3</sub> and Cu<sub>3</sub> nanomagnets with the Działoshinsky-Moriya (DM) interaction have received great attention for potential application in quantum computation [1-8]. V<sub>3</sub> with the out-of-plane Dz and in-plane Dn DM interactions demonstrates the multiferroic behavior in longitudinal magnetic field Bz|Z; field-induced polarization Pz(Bz) and toroidal moments Tz(Bz), which depend on the vector Kz(Bz) [scalar z(Bz)] chirality, and complicated individual spin dynamics at avoided level crossing (ALC) [7]. In order to investigate the multiferroic behavior in transverse field Bx<sup>⊥</sup>Z and B-reversal/rotation effects, P(Bx), Tn(Bx), M(Bx)magnetization, Kz(Bx), g(Bx) and spin canting were calculated for B<sup>⊥</sup>Z. The spin current P-Is mechanism [9] is the driving force of Pz(Bx) with a maximum at ALC that corresponds to Ig(max). The in-plane DM leads to the in-plane Tn(Bx) which exhibits T-flop under B-rotation/reversal. Tn(B), z(B), orbital moment M'(B), circular orbital current Io-z [Pz(B), Kz(B), Is] show the switch on/off [reduction] under B-rotation within vertical XZ plane. Field behavior Pz(Bx), Tn(Bx), M(Bx), Kz(Bx), Is, Io, and individual spin dynamics for Bx<sup>⊥</sup>Z strongly differ from that for Bz[Z v<sub>3</sub> exhibits no-linear reduction of Kz(Bx)/vector chirality, z(Bx)=0, and the entanglement effect. The V<sub>3</sub> nanomultiferroics

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## **RA26**

### Mössbauer studies of bismuth ferrite

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Perovskite bismuth ferrite (BiFeO<sub>3</sub>) is a multiferroic material, having ferroelectric order  $(T_c = 1,103 \text{ K})$  and antiferromagnetic order  $(T_v = 643 \text{ K})$  [1]. It is a possible candidate for information storage, spintronics, and sensors applications [2]. Here, we have measured the magnetic properties of BiFeO<sub>3</sub> using x-ray diffractometer, field emission scanning electron microscope (FE-SEM), vibrating sample magnetometer, and Mössbauer spectroscopy. Single phased BiFeO3 was synthesized by a low-temperature hydrothermal method. FE-SEM results showed that the grain sizes of microsphere and microcube were about 80 µm and 50 µm, respectively. BiFeO3 powder showed weak-ferromagnetic behavior with a small magnetization value (Ms ~ 100 memu/g) at 295K. Rietveld refinement analysis showed that Fe-site and O-site splits into two pairs, respectively. The Mossbauer spectra of BiFeO<sub>3</sub> were analyzed with 2-set sextet. Isomer shift ( $\delta$ ) values showed that Fe<sup>3+</sup> ions existed in octahedral site. Electric quadrupole shift (EQ) values suggested different lattice distribution at octahedral site. Based on EQ values (0.02 ~ 0.1 mm/s), splitting into two partially occupied sites results in a distorted FeO6 octahedral. The weak ferromagnetism observed in the hydrothermallyprepared BiFeO<sub>3</sub> particles comes from distorted FeO<sub>6</sub> octahedral, leading to incomplete counterbalance between the antiferromagnetic sublattices of the Fe ions.

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# RA27

### Synthesis and electrical characterization of Bi<sub>1-x</sub>Y<sub>x</sub>FeO<sub>3</sub> ceramics

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Compositions within range  $0.05 \le x \le 0.20$  in the system  $Bi_{1x}Y_xFeO_3$  (BYFO) have been synthesized by solid state reaction method. Structural and electrical properties of yttrium doped bismuth ferrite ceramics have been investigated. Single-phase formation has been confirmed by x-ray diffraction, which has shown rhombohedral symmetry with R3c space group at room temperature. Temperature dependent dielectric behavior of BYFO ceramics with in composition range  $0.05 \le x \le 0.20$  at few frequencies have shown an anomaly near the temperature which corresponds to magnetic phase transition ' $T_A$ ' for pure BiFeO<sub>3</sub> material (i.e.370°C). Substitution of yttrium in BFO has shown a shift in  $T_A$ : The temperature dependence of ac conductivity indicates that the conduction process may be influence by defects which occurs due to oxygen vacancies.

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## **RA28**

### Lattice engineering on transition metal oxide thin film Chang Uk Jung\*

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Transition metal oxide having a perovskite structure showed a large spectrum of physical properties. YBa2Cu3O7 showed high temperature superconductivity, (La,Ca,Sr)MnO3 showed ferromagnetism with half metallicity, Pb(Zr,Ti)O3 showed Ferroelectricity, and BiFeO3 showed coexistence of antiferromagnetism and ferroelectricity. Recent advance in interface control revealed a 2d electron gas, interface superconductivity, interface magnetism, oribital reordering, interface thermoelectric effect at the interface of transition metal oxide. The effect of lattice distortion has been studied for a long time and recently showed that lattice distortion during thin film growth can create new ground state. Especially ATiO3, (A=Sr, Y, Eu) was interesting. [In bulk state, YTiO3 and EuTiO3 are magnetic insulator, and SrTiO3 is not ferroelectric.] I will introduce the some current research on the progress on these materials In this case, only lattice mismatch between substrate and film was used. On the contrary the available substrate is rather narrow. I will also introduce efforts to overcome this problem.

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## **RB01**

Structural and magnetic properties of doped iron oxo-selenides

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In the Cu oxide superconductors the Cu resides in a square planar geometry[1], while in the Fe pnictide superconductors the Fe ion resides in a tetrahedrally coordinated position[2]. Both compounds may show a separating LaO layer between the superconducting sheets. La<sub>2</sub>O<sub>3</sub>Fe<sub>2</sub>Se<sub>2</sub> shows square planar Fe O sheets and offers a direct super exchange pathway to the next nearest Fe neighbour via the Se, comparable to the iron arsenides[3]. These super exchange pathways are similar to the superconductors despite the difference and provide a similar local environment for the Fe as is known for the Fe-superconductors[4]. We investigated the changes in the physical properties in the iron oxo-selenides by varying the electron density and structural components. We doped the parent compound with F, Cd and rare earths with smaller radii (La<sub>2</sub>O<sub>3</sub>, FxFe<sub>2</sub>Se<sub>2</sub>, and La<sub>2</sub>, Cd/RE), O<sub>3</sub>Fe<sub>2</sub>Se<sub>2</sub>, respectively). Temperature dependent neutron diffraction of the samples mostly show an antiferromagnetic order along the a-direction with a propagation of (0.5 0 0.5). We find indications for an evolving transition in the physical properties, as decreasing TN for fluorine and rare earth electron deficient Cd doped samples that show an in creasing TN.

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### **RB02**

# Phase separation of antiferromagnetism and superconductivity in $Rb_xFe_{2x}Se_2$ observed by Rb NMR

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The observation of the superconducting (SC) transition in A<sub>x</sub>Fe<sub>2-y</sub>Se<sub>2</sub> (A: alkali metal elements; Tc of ~30K) has been one of surprising things, because the transition was found in the antiferromagnetic (AF) system with a large AF moment of ~3.3  $\mu$ B/Fe. Moreover, their electrical resistivity just above Tc is too large to conceive the occurrence of the superconductivity. Here, to clarify whether the AF and SC states are microscopically coexistent or just phase-separated, we have carried out NMR measurements on both crystals with and without SC transition. The Rb-NMR spectra below 300 K observed for these crystals are clearly separable into two parts with different widths. The broader spectrum can be explained by the presence of the superconducting spin proposed by the neutron magnetic scattering studies [1]. The narrower spectrum for the crystal with SC transition can be understood from its temperature dependence and spectral shape to stem from the superconductivity and the antiferromagnetism in the present system are phase-separated.

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### RB03

### Superconductivity of EuFe<sub>2</sub>As<sub>2</sub> under high pressure Shugo Ikeda and Hisao Kobayashi

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EuFe<sub>2</sub>As<sub>2</sub> with the tetragonal ThCr<sub>2</sub>Si<sub>2</sub>-type structure exhibits a spin density wave (SDW) transition at 190 K, with a simultaneous structural change [1]. The SDW transition is suppressed by applied pressure and bulk superconductivity appears around 2.5 GPa [1-3]. Furthermore, magnetic moments of Eu sublattice order magnetically at To = 19 K. The value of To reveals no changes at pressures of up to 3 GPa [1-3]. We have studied the superconductivity and the magnetism of EuFe<sub>2</sub>As<sub>2</sub> by measuring dc magnetization and 57Fe nuclear forward scattering (NFS) under high pressure Magnetization at 2.6 GPa rapidly increases below 24 K and then gradually decreases by the appearance of superconductivity. The size of the drop at low temperatures is consistent with nearly 100 % shielding. Namely, bulk superconductivity appears under the magnetic ordered state of Eu sublattice. The pressure dependence of 57Fe NFS spectra shows the decrease of the hyperfine field Bhf of Fe sublattice with increasing pressure Furthermore it is found that the direction of Bhf varies around 2.6 GPa implying the changes of magnetic structure. In this presentation, we will discuss the relation between the superconductivity and the magnetism of EuFe<sub>3</sub>As<sub>3</sub>, considering these results

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### **RB04**

### Superconductivity in 4d, 5d pnictides

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After the discovery of the iron arsenide superconductors, a tremendous effort has gone into the search for new high-Tc superconductors. We explored 4d, 5d pnictides, hoping to find superconductors which share important characteristics with Febased superconductors. We report two new superconducting systems with such characteristics, but sufficiently different from Fe pnictides. (i) APt<sub>3</sub>P (A = Ca, Sr and La) The new compounds APt<sub>3</sub>P crystallize in an antiperovskite structure with distorted PPt6 octahedra. All of them display superconductivity below 6.6, 8.4 and 1.5 K for Ca-, Sr- and LaPt<sub>3</sub>P, respectively. The highest-Tc compound SrPt<sub>3</sub>P exhibits strongcoupling s-wave superconductivity. The presence of multiple pocket Fermi surface was inferred in SrPt<sub>1</sub>P, which might enhance the coupling with low-energy phonons. (ii) RuPn (Pn = P, As) Binary RuPn crystallize in an orthorhombic MnP-type structure. RuP and RuAs were found to show a metal to non-magnetic insulator transition at 270 and 200 K, respectively. The structural analysis evidenced another phase transition at higher temperatures, indicative of pseudogap formation. The two transitions can be suppressed by Rh-doping. We discovered superconductivity with a maximum Tc = 3.7 K and 1.8 K at the critical doping level for suppressing the pseudogap-phase in Ru1xRhxP and Ru1-xRhxAs, respectively.

D. Hirai et al., arXiv:1112.0604

### **RB05**

# Elastic softening and electric quadrupole in iron pnictide superconductor Ba(Fe<sub>1-x</sub>Co<sub>x</sub>)<sub>2</sub>As<sub>2</sub>

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We have performed ultrasonic measurements on iron based superconductor Ba(Fe<sub>1-x</sub>Co<sub>3</sub>)<sub>2</sub>As<sub>2</sub> in order to investigate the mechanism of the superconductivity. The end-material BaFe<sub>2-AS<sub>2</sub></sub> shows structural and antiferromagnetic simultaneous transitions at-140K. With increasing Co concentration x, the temperatures of these transition are lowered and separated from each other. The superconductivity appears at x>0.03. In our recent work [1], the softening of the elastic constant C66 as large as 21% was observed for optimized samples Ba(Fe<sub>0-x</sub>Co<sub>4</sub>)<sub>2</sub>As<sub>2</sub> with decreasing temperature from 300K down to superconducting transition temperature of TSC=23 K, while other elastic constants do not show softening. Here, we have measured elastic constants of Ba(Fe<sub>1-x</sub>Co<sub>4</sub>)<sub>2</sub>As<sub>2</sub> samples with x=0,0.03 and 0.07, where only the C66 shows huge softening larger than 70% with decreasing temperature down to the structural transition temperature. The large softening of C66 in Ba(Fe<sub>1-x</sub>Co<sub>4</sub>)<sub>2</sub>As<sub>5</sub> with several C0 concentration x≤0.10 has been reported by other groups [2,3]. This huge elastic softening of C66 is caused by the quadrupole fluctuation associated with the degenerate dy'z and dz' band aslant  $\Gamma$  and X points of the Brillouin zone. Quadrupole effects on the superconductivity [45,6].

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# **RB06**

### Concentration dependence of magnetic and transport characteristics in EuFe<sub>2</sub>As<sub>2</sub>, P, single crystals

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EuFe<sub>2</sub>As<sub>2</sub> exhibits the spin-density wave of Fe spins at TSDW=190 K and subsequently magnetic order of Eu<sup>2+</sup> moment at TN=19 K. It is also known that this compound shows the superconductivity by isovalent substitution of As to P and the Tc value increases up to 28 K. We have studied magnetic and transport properties in single crystal of the EuFe<sub>2</sub>As<sub>2-x</sub>P<sub>x</sub> system grown by the self-flux method. From the magnetization measurements at 5 K, we observed a saturation of magnetization of Eu<sup>+</sup> at HM ~ 1 T and a spin-flop like behavior at Hsf ~ 0.6 T for H||ab-plane. These fields seem to decrease as x is increased, indicating reduction of the exchange interaction and the uniaxial anisotropy. We also performed electron spin resonance measurements at X-band on single crystals of EuFe<sub>2</sub>As<sub>2-x</sub>P<sub>x</sub>. The result shows that the resonance field increases with decreasing temperature toward to TN=19 K below which a new line was observed. We will discuss these results including electron spin resonance experiments with emphasis on the phase diagram of magnetism and superconductivity in EuFe<sub>2</sub>As<sub>2</sub>. xP<sub>x</sub>.

### **RB07**

### Resonance-like response in antiferromagnetically ordered Fe<sub>1.02</sub>Te<sub>0.95</sub>Se<sub>0.05</sub>, studied by polarized inelastic neutron scattering

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Unconventional superconductivity in iron pnictides and chalcogenides [1] shows intriguing similarities with both copper-oxide- and rare-earth-based superconductors. Common is a proximity to a magnetic order and an existence of magnetic fluctuations [2]. Magnetic fluctuations change below the superconducting transition and are readily detectable by neutron inelastic scattering. The character of these has been shown to be related to the pairing mechanism and gap symmetry [3]. In all superconducting pnictides and chalcogenides one finds welldefined magnetic excitations centered close to the Q = (0.5, 0.5, 0) reciprocal position that strongly disperse with increasing energy and peaking in the 6.5-9.8 meV range [2]. Polarization neutron experiment allows for a determination of the in-plane and out-of-plane inelastic responses that are proportional to the relevant imaginary parts of the dynamical susceptibility. While the fluctuations in FeTe<sub>1.2</sub>Se<sub>0</sub>, x = 0.40 are nearly isotropic with a slightly enhanced component along the c axis [4,5], in BaFe<sub>1.9</sub>Ni<sub>01</sub>As<sub>2</sub> they are have an in-plane 2D character [6]. In this contribution we report on a polarized neutron study showing that similar fluctuations also exist also in antiferromagnetically ordered Fe<sub>1.02</sub>Te<sub>1.09</sub>Se<sub>0.09</sub> above the magnetic phase transition.

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### **RB08**

# Influence of filament diameter on superconducting properties of MgB<sub>2</sub> multi-core wires

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Multi-core MgB<sub>2</sub> wires with 19 filaments of different diameter were produced [1]. The filaments are covered by a Ti shell and embedded in a Cu matrix. From outside the wires are stabilized by stainless steel tubes. Critical current densities were determined from hysteresis loops using Bean's model [2]. Magnetic relxation measurements were performed at 1, 3, and 5 T both in increasing and decreasing field for up to 30 min. From these measurements mean effective activation energies were determined within the flux creep theory of Anderson [3]. All samples were prepared in identical way. They only differ in diameter of the filaments, which ranges between 30 and 15 µm. Influence of the filament diameter on the superconducting properties, like critical current density, irreversibility line and activation energies are discussed.

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## **RB09**

# Kinetic energy density of cooper pairs in Sr doped $YBa_2Cu_3O_7\text{-}\delta$ single crystals

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The kinetic energy, K(T) is a relevant tool to sort the distinct proposal for the normal state and their consequent paring mechanisms in HTSC[1,2] Motivated by this potential we report on reversible DC magnetizations as a function of temperature of a series of YBa<sub>2x</sub>Sr,Cu<sub>3</sub>O<sub>7</sub> $\delta$  (x = 0, 0.02, 0.10, 0.25, 0.37 and 0.50) single crystals with the purpose to study the role of ion size chemical introduced disorder on the average kinetic energy density (AKED) of Cooper pairs in the YBa<sub>2</sub>Cu<sub>3</sub>O<sub>7</sub> $\delta$  superconducting state[3] The Sr ion size chemical introduced disorder is introduced in YBa<sub>4</sub>Cu<sub>3</sub>O<sub>7</sub> $\delta$ ; superconducting state[3] The Sr ion size chemical introduced disorder is mitoduced in YBa<sub>4</sub>Cu<sub>3</sub>O<sub>7</sub> $\delta$ ; superconducting state[3] The Sr ion size chemical disorder is introduced in YBa<sub>4</sub>Cu<sub>3</sub>O<sub>7</sub> $\delta$ ; superconducting state[3] The Sr ion size chemical introduced the magnetization data is supported by viral theorem of superconductivity. [1,2,4] The AKED preliminary results of the YBa<sub>12</sub>Su<sub>3</sub>(Cu<sub>4</sub>O<sub>7</sub> $\delta$  single crystal show that the series of the isofield K(T) curves, obtained to the H = 10 kOe, scales in a common behavior that smoothly develops as the temperature is increasing toward to the Te, disappearing inside a common behavior flat smoothly develops as the temperature is increasing toward to the Te, disappearing inside a common behavior flat smoothly develops as the temperature is increasing toward to the Te, disappearing inside a common behavior flat smoothly develops as the temperature is increasing toward to the Te, disappearing inside a common behavior flat smoothly develops as the temperature is increasing toward to the Te, disappearing inside a common behavior flat smoothly develops as the temperature is increasing toward to the Te, disappearing inside a normon behavior flat smoothly develops as the temperature is increasing toward to the Te, disappearing inside a normon behavior flat smoothly develops as the temperature is increasing toward to the Te, disappearing inside a normon behavior flat smoo

I. S. Salem-Sugui, Jr., M. M. Doria, A. D. Alvaraga, V. N. Vieira, P. F. Farinas and J. P. Sinnecker, Phys. Rev. B 76, 132502 (2007). 2. M. M. Doria, S. Salem-Sugui, I. G. de Oliveira, L. Ghivelder and E. H. Brandt, Phys. Rev. B 65, 144509 (2002). 3. V. N. Vieira and J. Schaf, Phys. Rev. B 65, 144531 (2002). 4. M. M. Doria, J. E. Gubernatis and D. Rainer Phys. Rev. B 39, 9573 (1989).

## **RB10**

### Interplay between superconductivity and antiferromagnetism in BaFe<sub>2</sub>. Ni,As, single crystals studied by 57Fe Mossbauer spectroscopy

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We have performed magnetic susceptibility and 57Fe Mossbauer spectroscopy measurements on single crystal mosaics of BaFe<sub>2x</sub>Ni<sub>x</sub>As<sub>2</sub> (x=0.065, 0.075, 0.085). Our findings show antiferromagnetic order below TN obtained from magnetic hyperfine field phase diagram, with superconductivity below TC Obtained by susceptibility measurements. Careful analysis of our Mossbauer data show that a decrease in the magnetic hyperfine field is observed below TC and the effect is proportional to the difference between TN and TC, being the largest when both temperatures become close. We argue this is caused by coexistence of superconductivity and magnetism, since both phenomena come from the same Fe 3d electrons, and two possible scenarios are presented. Magnetic volume fraction does not show any variation at or below TC, but a strong dependence with doping level, most likely due to short range magnetic order that cannot be sensed by our 57Fe local probe, probably caused by local inhomogeneties. A phase diagram is shown to illustrate how the variation of the magnetic volume fraction and the crossover between magnetism and superconductivity are related.

M. Wang, H. Luo, J. Zhao, C. Zhang, M. Wang, K. Marty, S. Chi, J. W. Lynn, A. Schneidewind, S. Li and P. Dai, Physical Review B 81, (2010) 174524. M. Alzamora, J. Munevar, E. Baggio-Saitovitch, S. L. Bud'ko, Ni Ni, P. C. Canfield and D. R. Sanchez, Journal of Physics: Condensed Matter 23, (2011) 145701. arXiv:1111.5853

## **RB11**

### Magnetism in superconducting Ba<sub>0.78</sub>K.<sub>22</sub>Fe<sub>2</sub>As<sub>2</sub> and EuFe<sub>2</sub>As<sub>1.4</sub>P<sub>0.6</sub> single crystals studied by Mossbauer spectroscopy

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We studied magnetism in superconducting Fe-pnictide single crystals using 151Eu and 57Fe Mossbauer Spectroscopy. Neutron studies on Ba, K, Fe,As, single crystals revealed a decrease in Bragg peak intensity due to iron moment or magnetic volume fraction reduction. 57Fe Mossbauer measurements on Ba<sub>0.7</sub>K<sub>0.2</sub>Fe,As<sub>2</sub> single crystal mosaics were performed down to 1.5 K, revealing a decrease in the magnetic hyperfine field below Te, without change in magnetic volume fraction. Our data confirm a reduction of Fe magnetic moment below Te explaining also the neutron results. Such a decrease is caused by a spectral weight transfer when entering the superconducting state. For EuFe<sub>2</sub>As<sub>1</sub>, P<sub>06</sub> compound, magnetism originates from Eu<sup>+</sup> moments. Mossbauer spectra reveal magnetic hyperfine fields below TM = 18 K of Eu<sup>+</sup> moments. Data analysis shows coexistence of ferromagnetism, from Eu<sup>+</sup> moments ordered along crystallographic e-axis, and superconductivity below TSC ~ 15 K. We find indications for a change of small Fe magnetic moment dynamics (~ 0.07 µB) at superconductivity onset: below TSC the Fe magnetic moment sate when entering the SC state, and discussion about coexistence of magnetism and superconductivity will be provided.

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# **RB12**

# Doping evolution of the in-plane London penetration depth in $Fe_{1+s}(Te_{1+s}Se_s)$ single-crystals

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Our work focuses on the doping evolution of the superconducting properties as derived from penetration depth measurements as a function of temperature in single crystals of Fe<sub>1+x</sub>(Te<sub>1-x</sub>Se<sub>x</sub>) with y<0.02 and Se concentration spanning from 25% to 45%. The London penetration depth is one of the most important characteristic parameter in type II superconductors, as it can give information about the pairing mechanism, with its zero temperature value being directly related to the density of superconductive electrons in the material. Precise measurements of the in-plane penetration depth  $\lambda ab$  as a function of temperature using a tunnel diode oscillator technique were performed in a dilution refrigerator down to a temperature of 30mK. By using a set of two mutually coupled planar inductors the probing ac field is uniform across the sample along the c axis with the variation in susceptibility solely due to supercurrents flowing in the ab crystallographic plane. A complete study of the temperature and Se concentration dependence of the physical properties derived from these measurements is presented.

## RB13

Potassium doping effect in double-chain BaFe<sub>2</sub>Se<sub>3</sub>

Jinke Bao, Yunlei Sun, Wenhe Jiao, Yongkang Luo, Zhu'an Xu and Guanghan Cao\* Department of Physics, Zhejiang University, China

Compared to iron-based superconductors with two-dimensional FeAs(Se) layers, BaFe<sub>2</sub>Se<sub>3</sub> is a one-dimensional material consisting of Fe<sub>2</sub>Se<sub>3</sub> double chains with Ba<sup>2+</sup> cations intercalated. It exhibits antiferromagnetic order with ferromagnetic Fe<sub>4</sub> block of 2.8µB / Fe<sup>2+</sup> below 255 K.[1][2][3] Considering the close relationship between antiferromagneitsm and superconductivity, we performed a chemical doping with potassium in this compound. Ba<sub>1-x</sub>K<sub>x</sub>Fe<sub>2</sub>Se<sub>3</sub> single crystals with different K content were successfully grown. Resistivity measurements showed that all the crystals were semiconducting with no superconducting transition above 3K. The magnetization of Ba<sub>1-x</sub>K<sub>x</sub>Fe<sub>2</sub>Se<sub>3</sub> is anisotropic, and the sample of x=0.4 exhibits spin glass behavior.

[1] A. Krzton-Maziopa, E. Pomjakushina, V. Pomjakushin, D. Sheptyakov, D. Chernyshov, V. Svitlyk, and K. Conder, J. Phys. Condens Matter 23, 402201 (2011), [2] J. M. Caron, J. R. Neilson, D. C. Miller, A. Llobet, and T. M. McQueen, Phys. Rev. B 84, 180409(R) (2011). [3] Hechang Lei, Hyejin Ryu, Anatoly I. Frenkel, and C. Petrovic, Phys. Rev. B 84, 214511 (2011).

### **RB14**

# $\mathrm{Sr}_2\mathrm{VO}_3\mathrm{FeAs}$ : Hybrid of a magnetic $\mathrm{Sr}\mathrm{VO}_3$ and a FeAs superconducting layers

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We investigate the upper critical fields up to 14T of Sr<sub>2</sub>VO<sub>3</sub>FeAs single crystals. For Sr<sub>2</sub>VO<sub>3</sub>FeAs, where the perovskite SrVO<sub>3</sub> layers are inserted between SrFe<sub>2</sub>As<sub>2</sub> layers, we observed a strong anisotropy of the upper critical field, compared with Hc2(T) of other pnictides. Strong anisotropy in Hc2 of Sr<sub>2</sub>VO<sub>2</sub>FeAs suggests that SrVO<sub>3</sub> layer is insulating. Furthermore, we found the anomalous magnetoresistance in the normal states up to T~150K indicating possible magnetic ordering in the SrVO<sub>3</sub> layer. Sr<sub>2</sub>VO<sub>3</sub>FeAs have a structural peculiarity that means each magnetic ordered layers and charge reservoir layers. Thus, unlikely usual layered pnictide compounds, this compound provides the magnetic interaction.

### **RB15**

# Static and dynamic properties of nearly optimaly doped superconductor $SmFeAsO_{0.86}F_{0.14}$

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We have reported resistivity, heat capacity, magnetization and 19F and 75As NMR in polycrystalline sample of nearly optimaly doped superconductor SmFeAsO<sub>1.4</sub>F<sub>x</sub> (x = 14%). The resistivity shows linear dependence above 200 K and 50 - 160 K. The linear dependence reveals non Fermi liquid behavior. The heat capacity data for SmFeAsO<sub>0.86</sub>F<sub>0.14</sub> exhibits a small jump at T\_C-48K, indicating the bulk nature of superconductivity. The field cooled magnetisation data deviates from Curie Weiss law below 160 K. A small peak around 163 K was observed in 19<sup>6</sup> spin-spin and spin-lattice relaxation rate showing evidence of structural phase transition driven by antiferromagnetic spin fluctuation. 75As NMR signal is detectable only above 180 K All measurements point towards a possible existence of nematic ordering where structural phase transition have been driven by anti-ferromagnetic spin-fluctuations at 163 K. Their is a considerable FeAs and SmO layer interaction present in the system. The results also revealed the coexistence of local and timerant electronic states below 90 K where two fluid model of heavy fermion holds. Coexistence of static anti-ferromagnetic ordering due to Sm 4f electrons and superconductivity has also been observed.

### **RB16**

### Effect of d-orbital characters of the Fe magnetic moment on the electronic and magnetic properties of BaFe<sub>2</sub>As<sub>2</sub> Hyungju Oh and Hyoung Joon Choi\*

Department of Physics and IPAP, Yonsei University, Korea

There have been many studies detailing the orbital character of band structures in the iron-based superconductors. However, the orbital characters of the Fe magnetic moment still remain unrevealed up to now. By performing first-principles calculations of the electronic and magnetic properties with constraint on the real space shape of Fe magnetic moments, we study the d-orbital characters of the Fe magnetic moment in BaFe<sub>2</sub>As<sub>2</sub>. We find that, depending on the d-orbital characters of the Fe magnetic moment, the electronic properties of BaFe<sub>2</sub>As<sub>2</sub> such as band gap, band dispersion, and dzx/dyz occupation change distinctively. Furthermore, we compare obtained band structures with published angle-resolved photoemission spectroscopy (ARPES) result, and propose that the Fe magnetic moment in BaFe<sub>2</sub>As<sub>2</sub> has in-plane dxy character. This work was supported by the NRF of Korea (Grant Nos. 2011-0018306). Computational resources have been provided by KISTI Supercomputing Center (Project No. KSC-2011-C3-05).

## RB17

# $Ca(Fe_{1:x}Co_x)_2As_2$ under pressure: studies using single crystal neutron diffraction

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CaFe<sub>2</sub>As<sub>2</sub> and related compounds exhibit extremely strong coupling of the magnetic and structural phase transitions and sensitivity to an applied pressure. Undoped system undergoes a structural transformation from a high temperature tetragonal (T) phase to a structure with orthorhombic (O) symmetry that is associated with the appearance of an antiferromagnetic order [1]. Upon application of a ranther modest pressure, both transitions are suppressed and superconductivity appears below 12 K [2]. For pressures larger than 0.3 GPa another, so-called collapsed tetragonal (CT) phase with greatly reduced c parameter appears [3]. Three different Ca(Fe<sub>1</sub>, Co,)<sub>2</sub>As<sub>2</sub> single crystals with concentration x = 0.032, 0.051 and 0.063 were investigated by means of neutron diffraction under uniaxial and hydrostatic pressure. Region of existence of the high-temperature tetragonal, low-temperature orthorhombic and pressure-induced collapsed tetragonal phases have been determined. Structural details of the tetragonal phases at various conditions have been determined and the corresponding p-T phase diagrams have been constructed. The critical pressure of the T (or O) to CT transition decreases with increasing Co doping level. The behavior is highly hysteretic and structural studies show no variation of the As positional parameter with pressure.

[1] D. C. Johnston, Advances in Physics 59, 803 (2010). [2] M. S. Torikachvili, et al., Phys. Rev. Lett. 101, 057006 (2008). [3] A. Kreyssig, et al., Phys. Rev. B 78, 184517 (2008).

# **RB18**

## Monte Carlo study of the magneto-elastic effects in Fe pnictides

Cesar Jose Calderon Filho $^{1} \mbox{{\sc {*}}}$  Alex Antonelli $^1$ , David Vaknin $^2$  and Gaston Eduardo Barberis $^1$ 

<sup>1</sup> IFGW, State University of Campinas - UNICAMP, Brazil

<sup>2</sup> Ames National Laboratory and Department of Physics, Iowa State University, USA A (J1, J2) model of the infinite square lattice of classical spins coupled by Heisenberg

A (17, 22) index of the infinite square faited of classical spins coupled by fields (Eq. 27) exchange interactions that includes a spin-lattice interaction of the form A (Exx - Eyy) (Sx2- Sy2) is used in Monte-Carlo simulations to determine a phase diagram in the J1, J2, and A parameter-space. To facilitate phase transitions in the 2D system, we also include a weak out-of-plane coupling to spins in adjacent planes (Jz). This phase diagram is relevant to the parent materials of the newly discovered iron-pnictides superconductors. Using first principles band calculations, both for the low and high temperatures, we obtain realistic parameters to the model that allow the prediction of pressure induced magnetic transitions.

### RB19

**Geometrical Vortex Transition in the iron pnictide SmFeAs(O,F)** Philip J.w. Moll<sup>1</sup>\*, Luis Balicas<sup>2</sup>, Janusz Karpinski<sup>3</sup>, Nikolai Zhigadlo<sup>4</sup> and Batlogg Bertram<sup>5</sup>

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We have observed a distinct vortex matter transition in SmFeAs(O,F) (Te~50K) upon cooling below T\*-41K(=0.81Tc) into a highly mobile state for fields well aligned with the FeAs layers. Below T\*, the vortex core is confined between two adjacent FeAs planes by periodic modulations of the superconducting order parameter within the unit cell. Vortex motion parallel to the planes is even more pronounced at lowest temperatures, well below Tc and Hc2, as the vortex cores avoid the highly effective pinning sites located in the FeAs layers. For fields slightly out-of-plane (<0.3°) the vortices are again completely immobile as they cross the planes and are hence strongly pinned by defects within the FeAs layers resulting in large critical current densities at high fields and low temperatures (>10<sup>6</sup> A/cm<sup>2</sup> @5K, 14T[1]). These results indicate strong pinning localized in the FeAs layer to be the dominant factor for the highly effective pinning and thus show the pathway to improve the technological prospect of the pincides.

[1] P.J.W. Moll et al., 'High magnetic field scales and critical currents in SmFeAs(O,F) crystals. Nature Materials 9, 628 (2010)

### **RB20**

# SmFeAsO<sub>1-x</sub>F<sub>x</sub>: Raman scattering and x-ray diffraction study under low temperature and high pressure conditions

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We present a comparative low-temperature Raman study on single crystals of the parent non-superconducting SmFeAsO and the fluorine-doped superconducting SmFeAsO<sub>00</sub> $F_{0.35}$ [1]. As in the compounds from the 122 family [2], clear renormalization of the c-axis polarized A1g(Sm) and the B1g(Fe) Raman-active modes are observed through the magneto-structural transition in the parent compound. Splitting of the in-plane Eg modes is also observed (like the reported in [3,4]). In the fluorine-doped compound, x-ray diffraction measurements reveal the retention of the structural transition (similar to the one reported in [5]). Despite the absence of a coupled magnetic transition, a renormalization of the c-axis polarized modes similar to the one seen in the parent compound is observed, followed at lower temperature by an anomalous increase of the linewidth of these phonons below the superconducting transition temperature. The effects of the structural transition on the phonon modes are further investigating by applying high pressure on the parent compound, allowing a direct comparison of the effect of pressure and doping on the lattice dynamics of the system.

[1] Z.-A. Ren et al., Chinese Physics Letters 25 (6), 2215 (2008) [2] M. Rahlenbeck et al., Phys. Rev. B 80, 064509 (2009) [3] T. Yildirim, Phys. Rev. Let. 101, 057010 (2008) [4] L. Chauviere et al., Phys. Rev. B 80, 094504 (2009) [5] A. Martinelli et al., Phys. Rev. Let. 106, 277001 (2011)

# **RB21**

### Mass renormalization in isostructural Ru- and Fe-pnictides

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We have studied the angular-dependent de Haas-van Alphen oscillations of LaRu<sub>2</sub>P<sub>2</sub> using magnetic torque in pulsed magnetic fields up to 60 T. The observed oscillation frequencies are in excellent agreement with the geometry of the calculated Fermi surface. The temperature dependence of the oscillation amplitudes reveals effective masses m\*( $\alpha$ ) = 0.71 and m\*( $\beta$ ) = 0.99 me, which are enhanced over the calculated band mass by  $\lambda$ cyc of 0.8. We find a similar enhancement of  $\lambda\gamma \sim 1$  in comparing the measured electronic specific heat ( $\gamma$  =11.5ml/mol K<sup>2</sup>) with the total density of states from band-structure calculations. Remarkably, very similar mass enhancements have been reported in other pnictides, LaFe<sub>2</sub>P<sub>2</sub>, LaFePO (Tc ~ 4 K), and LaRuPO, independent of whether they are superconducting or not. This is contrary to the common perceptions that the normal-state quasiparticle renormalizations reflect the strength of the superconducting pairing mechanism. To separate mass enhancement due to electron correlations find e.g. electron-phonon interactions, high energy bandwidth measurements such as ARPES would be highly desirable, thus leading to new questions about competing pairing in isostructural and isoelectronic Ru- and Fe-pnictide superconductors.

PHYSICAL REVIEW B 84, 224507 (2011)

## **RB22**

The role of spin fluctuation for the high Tc superconductivity in LiFeAs, LiFeP, NaFeAs

Geunsik Lee, Hyo Seok Ji and Ji Hoon Shim\* Chemistry, POSTECH, Korea

Based on fully self-consistent density functional theory and dynamical mean field theory (DMFT) methods, we have studied the electronic spectrum of three isovalent 111 compounds, LiFeAs, LiFeP, NaFeAs. It has been known that while both LiFeAs and LiFeP are intrinsic superconductors with the nodeless and nodal gap symmetries respectively, NaFeAs exhibits spin density wave phase where superconducting transition occurs with chemical doping. With our simulated electronic spectrums, we examine the spin fluctuation mechanism for the high Tc superconductivity.

## RB23

DFT+DMFT study on the electronic structure and anisotropy of Sr<sub>2</sub>VO<sub>3</sub>FeAs superconductor. Hyo Seok Ji and Ji Hoon Shim\* Chemistry. POSTECH. Korea

Electron correlation effect on Fe 3d orbital has very important role in describing electronic structure and anisotropy of iron-based superconductors.[1] By using combination of the density functional theory (DFT) and dynamical mean field theory (DMFT), the electron correlation effect can be described well. In this study, we have investigated the electronic structure and anisotropy of Sr<sub>2</sub>VO<sub>3</sub>FeAs superconductor. Simple DFT result shows mixed band structures of Fe 3d and V 3d orbital, but the experimental result shows only Fe 3d band near the Fermi level [2] The DFT+U method explains insulating V layer structure correctly[2, 3] but it cannot describe whole electronic structure such as renormalized Fe 3d bands. So we apply the electron correlation effect on the V 3d orbital as well as the Fe 3d orbital using DFT+DMFT. With antiferromagnetic(AFM) ordering, V 3d bands split into -1 eV and 1 eV region, which is consistent with experimental ARPES result.[2] Also we calculate electrical anisotropy  $\gamma = \sigma xx/\sigma zz$  which is important for determining dimensionality effect in iron-based superconductors. Calculated anisotropy value is much smaller than simple DFT result and the result is well consistent with recent experimental result. This trend is also consistent with the trend of previous work[1].

[1] H. S. Ji, G. Lee, and J. H. Shim, Phys. Rev. B 84, 054542 (2011). [2] T. Qian et al., Phys. Rev. B 83, 140513(R) (2011). [3] H. Nakamura and M. Machida, Phys. Rev. B 82, 094503 (2010).

# **RB24**

# Behaviour of Magnetic and structural transitions upon Sr doping in CaFe<sub>2</sub>As, and EuFe<sub>2</sub>As,

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We have grown single crystals of Ca<sub>1-x</sub>Sr<sub>x</sub>Fe<sub>2</sub>As<sub>2</sub> and Eu<sub>1-x</sub>Sr<sub>x</sub>Fe<sub>2</sub>As<sub>2</sub>. Our aim was to study the progression of SDW transition and the effect of dilution on Eu magnetism in EuFe<sub>2</sub>As<sub>1-x</sub>Sr<sub>x</sub>Fe<sub>2</sub>As<sub>2</sub>. CaFe<sub>2</sub>As<sub>2</sub> and SrFe<sub>2</sub>As<sub>2</sub> undergo magneto-structural transition at 170 K and 205 K respectively[1,2]. With introduction of Sr in the CaFe<sub>2</sub>As<sub>2</sub>, magneto-structural transition temperature increases, and it reaches 200 K in nominal Sr<sub>0-3</sub>Ca<sub>0-3</sub>Fe<sub>2</sub>As<sub>2</sub> indicating that Sr dominates Ca in determining the SDW and structural transition. Doping Fe with Co induces superconductivity in Ca<sub>0-5</sub>Sr<sub>0-5</sub>Fe<sub>2</sub>As<sub>2</sub> below 12 K, which surprisingly shows a signature of SDW transition temperature of Eu sublattice decreases with increasing x while the structural/SDW transition temperature gradually varies from 190 K to 200 K. Field induced ferromagnetic behaviour of the parent EuFe<sub>2</sub>As<sub>2</sub> is also observed in the Sr doped compounds, the magnetic field needed to induce ferromagnetic meters from decreases with increasing Xr.

[1] Neeraj et. al. Phys. Rev. B, 79, 012504 (2009). [2] Krellner et al, Phys. Rev. B, 78, 100504(R) (2008)

### **RB25**

### **Specific heat of the vortex lattice in iron-pnictide superconductors** Miguel Araujo<sup>1</sup> and Pedro Sacramento<sup>2</sup>

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<sup>2</sup> CFIF, Instituto Superior Tecnico, TU Lisbon, Portugal

We study a vortex lattice in a three-orbital model for iron-pnictide superconductors. A gauge transformation, first introduced by Vafek-Tesanovic [1], is employed in order to obtain an effective spatially periodic Hamiltonian that produces a spectrum of Bloch states and automatically includes quantum effects, such as the Berry phase and Doppler shift of quasi-particles, without resorting to semi-classical approximations. We calculate the energy spectrum and specific heat as function of applied magnetic field. The effect of disorder due to impurities and positional disorder of the vortices is also addressed

[1] O. Vafek, A. Melikyan, M. Franz and Z. Tesanovic, Physical Review B 63, 134509 (2001)

### **RB26**

### Hexagonal superconducting pnictide SrPtAs: an ab initio study Sonny S. H. Rhim<sup>1</sup>, S. J. Youn<sup>2</sup>, Daniel Agterberg<sup>3</sup>, Michael Weinert<sup>3</sup> and Arthur J

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<sup>3</sup> Physics, U. Wisconsin-Milwaukee, USA

We present an ab initio study on SrPtAs, a recently discovered hexagonal pnictide superconductor. SrPtAs is distinct from other pnictides: (i) it has a hexagonal structure and (ii) a large spin-orbit interaction, whose consequence is not trivial. Based on electronic structure calculations, a tight-binding Hamiltonian is constructed to analyze the character of the spin-orbit interaction associated with the crystal structure. Further, the pairing mechanism by phonon and spin fluctuation are analyzed from phonon dispersion and susceptibility with matrix element fully taken into account. Finally, based on the band structure, we propose that electron doping could increase TC.

## RB27

Korea

### Josephson effect in BaKFeAs inter-grain boundary junctions Sung-hak Hong<sup>1</sup>, Sung Hoon Lee<sup>1</sup>, Soon-gul Lee<sup>1</sup>\*, Soon-gil Jung<sup>2</sup>, Nam Hoon Lee<sup>2</sup>

and Won Nam Kang<sup>2</sup> <sup>1</sup> Department of Display and Semiconductor Physics, Korea University, Korea <sup>2</sup> BK21 Physics Division and Department of Physics, Sungkyunkwan University,

We have studied superconducting transition properties of  $Ba_{0.6}K_{0.4}Fe_2As_2$  intergrain junctions. The junctions were fabricated by focused ion beam (FIB). Prior to FIB patterning, films were prepatterned into microbridges by a standard photolithography and argon ion milling technique. The lowest-possible beam current of 1.5 pA was used for the FIB pattern to minimize the etching damage to the bridge. The nominal dimensions were 200 nm in width and 100 nm in length. Resistive transition, currentvoltage (I-V) characteristics, the temperature-dependent critical current (Ic), and the normal state resistance (Rn) were measured. Measured transition properties indicated that the junctions were dominated mainly by Josephson coupling, showing resistivelyshunted-junction (RSJ) behaviors. At higher currents, multiple transitions were observed, which might be due to dynamics of Josephson vortices driven by transport currents. Details of the measurement results will be discussed.

### **RB28**

Coexistence of different electronic phases in the  $K_{0.8}Fe_{1.6}Se_2$ superconductor: a bulk-sensitive hard x-rays spectroscopy study

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In A<sub>x</sub>Fe<sub>2x</sub>Se<sub>2</sub> (11)-type chalcogenides superconductivity appears to occur only in Fe-deficient samples, where the alkali metals are intercalated between the FeSe layers [1-3]. It has been demonstrated how the ordering of the Fe vacancies can be correlated to superconductivity can be tone on a single sample from an insulating state by post-annealing and fast quenching, and it is tempting to conclude that the superconductivity is achieved when the Fe-vacancies are in a disordered state [6]. Here we report on the electronic and magnetic structure of  $K_{0x}Fe_{1x}Se_{2x}$  superconductor by x-ray emission and high resolution absorption spectroscopy. We report a study where the electronic and magnetic properties are investigated at the same time as a function of temperature in several consistent thermal cycles. We discuss the effects of ordered and disordered Fe vacancies on the electronic and magnetic structure, the existence of memory effects on thermal cycles, and the relation between electronic and magnetic properties and superconductivity. The results obtained are finally compared with the electronic and magnetic properties of (11)-type chalcogenides [7].

 Jing Guo, et al, arXiv: 1101.0092; Y. Kawasaki, et al, arXiv: 1101.0896; J J Ying et al, New J. Phys. 13, 033008 (2011), [2] D. MWang, et al, arXiv: 1101.0789, [3] A. M. Zhang, et al, arXiv: 1101.2168v1, [4] Ricci et al. Supercond. Sci. Technol. 24, 082002 (2011), [5] Bao Wei et al, Chin. Phys. Lett. 28, 086104 (2011). [6] F. Han et al., arXiv: 11031347. [7] L.Simonelli et al, submitted to PRB.

### **RB29**

# Step towards a ferromagnetic Josephson junction in YBCO/LCMO heterostructures

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High quality high-Tc superconductor/half metal ferromagnet/high-Tc superconductor YBCO/LCMO/YBCO trilayers (SFS) with different orientation of  $CuO_2$  planes in superconducting electrodes and different thicknesses of LCMO layers have been successfully fabricated by employing pulsed laser deposition techniques. It has been found that a LCMO barrier with thickness of 2-3 unit cells is ferromagnetic with a Curie temperature close to the bulk value. The strong difference in magnetization loops for different SFS-structures measured at T = 5K and the magnetic field parallel to the F-layers is caused by the presence of a superconducting current across the ferromagnet in the (110) oriented SFS-structure with a thickness of the ferromagnetic layer =1nm. This has been confirmed by measuring current-voltage characteristics in patterned structures. This sample structure fulfils all prerequisites for the realization of high-Tc SFS-Josephson junctions.

# RC01

### **DMFT study of the correlation effects on a topological insulator** Tsuneya Yoshida\*, Satoshi Fujimoto and Norio Kawakami

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Recently, topological insulators are theoretically proposed in correlated electron systems (e.g. Iridium oxides, Heusler compounds, and filled skutterudites)[1]. The correlation effects on the topological insulators are expected to produce exotic states [2] and are extensively studied. The strong Coulomb interaction could drive the non-trivial band insulator into a Mott insulator. In this paper, we study an extended Bernevig-Hughes-Zhang model[3] including the on-site Coulomb interaction with DMFT. We calculate the double occupancy at each orbital and the spin-Hall conductivity which clearly distinguishes the trivial and non-trivial phases in our model. From the analysis of these quantities, it is concluded that the topological insulator changes into a trivial nonmagnetic Mott insulator via a first-order transition [4]. Therefore, the gap closing does not occur at the transition. In the presentation, we also address the antiferromagnetic instability.

A. Shitade et al., Phys. Rev. Lett. 102, 256403 (2009), S. Chadov et al., Nature Materials 9 541-545 (2010), B. Yan et al., arXiv 1104.0641.
 D. A. Pesin et al., Nature Physics 6, 376 - 381 (2010).
 B. A. Bernevig et al., Science 314 1757 (2006). [4] T. Yoshida et al., arXiv 1111.6250.

### **RC02**

# Self-consistent treatment of fully relativistic effect in the small bismuth clusters Bin $(2 \le n \le 7)$

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The energy scale of spin-orbit (SO) interaction reaches to the order of 0.1-1eV at the elements in the fifth low of periodic table. Such interaction could cause effects on structural properties as well as on electronic ones. Bismuth has been investigated by many researchers due to its largeness of SO coupling. Rashba effect, which appears as the spin state splitting of surface bands, has been discussed intensively in the system which consists of bismuth [1]. The phenomena which are related with new physics of topological insulator have been investigated in a lot of bismuth systems [2]. In this context, bismuth clusters are interesting to investigate including relativistic effects. In this work, we have studied structural and electronic properties in small bismuth clusters; Bi n ( $2 \le n \le 7$ ). The Car-Parrinello molecular dynamics [3] and self-consistent fully relativistic calculations [4,5] have been employed to investigate effects of SO coupling effects. It has been found that the clusters of odd number atoms were found to carry magnetic moment with non-collinear magnetic configuration.

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### RC03

### Hidden topological order in URu<sub>2</sub>Si<sub>2</sub>

Tanmoy Das Theoretical Division, Los Alamos National Laboratory, USA



# RC04

# Ground-state properties of a two-dimensional correlated topological insulator

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Recently, the topological band insulator (TBI) has attracted much interest in condensed matter physics. The TBI is characterized by a nontrivial band structure, which has a full insulating gap in the bulk and gapless edge states on the boundaries. While the essential features of the TBI can be described as a one-body problem, there is much interest in the correlation effects on the TBI. In this study, we investigate a generalized Bernevig-Hughes-Zhang model having electron correlations with the variational Monte Carlo (VMC) method. In order to study how interactions affect the edge states and bulk states, it is important to consider inhomogeneous nature due to the lack of translational symmetry perpendicular to the boundaries. For this purpose, we introduce spatially-dependent variational parameters. We calculate the ground-state energy and the momentum distribution to discuss how the correlation effects affect electronic properties in the TBI.

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## RC05

Magnetic impurity doping effect on bulk Rashiba spin splitting system BiTeI

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# **RC06**

### Topological charge pumping effect by the magnetization dynamics on the Surface of 3D Topological Insulators

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We discuss a current dynamics on the surface of a 3-dimensional topological insulator induced by magnetization precession of a ferromagnet attached. It is found that the magnetization dynamics generates a direct charge current when the precession axis is within the surface plane. This rectification effect is due to a quantum anomaly and is topologically protected. The rectified current is useful in the sense that a direct current is easy to detect experimentally. To see the effect of the fast-varying exchange field, we study the current dynamics by the topological field theory. We adopt the dimensional regularization to regularize the ultraviolet divergence, which is inevitable to study an electromagnetic response of a topological insulator. As a result, the robustness of the rectification to the fast-varying exchange field is confirmed.

# **RC07**

Tuning of carrier type in Mn-doped Bi<sub>2</sub>Se<sub>3</sub>

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Bi<sub>2</sub>Se<sub>3</sub> is a good candidate for studying thermoelectric material as well as topological insulator. [1] In order to promote the thermoelectric performance, it is good to connect two materials in series with different charge carrier type. In Bi<sub>2</sub>Se<sub>3</sub>, the Se vacancies are dominant, which leads to n-type charge carriers.[2] In order to tune the carrier type and carrier density, we have doped Mn into Bi<sub>2</sub>Se<sub>3</sub> and have studied their physical properties, mainly focusing on the thermoelectric properties. The single crystals of MnxBi<sub>2</sub>xSe<sub>3</sub> (x=0.03, 0.05, 0.09 and 0.15) were grown by melting the stoichiometric elements in an evacuated quartz ampoule. The Hall and thermoelectric measurements illustrate the change of carrier type from n to p type between x=0.03 and 0.05. The transport measurements have shown the metallic behavior for x $\leq$ 0.03 and nonmetallic behavior for x $\geq$ 0.05 at low temperatures. This supports that the Fermi level of p-type samples lies in the surface state between the bulk valence band maximum and the bulk conduction band minimum.[3] The Seebeck coefficient increases linearly up to ~100µV/K at room temperature both n- and p-type samples. This implies that Mn-doped Bi<sub>2</sub>Se<sub>3</sub> is possible materials for thermoelectric application.

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# RC08

Magnetic edge profile in the Kane-Mele-Hubbard model

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We investigate the magnetic state on the edge of the Kane-Mele-Hubbard model. The ferromagnetic edge states are induced in a graphene nanoribbon by electron-electron interactions while certain spin-orbit interactions are known to give rise to a gapless current-flowing state on the edge without electron-electron interactions. We use the Hartree-Fock approximation to examine the nature of edge states in the presence of the Hubbard interaction in the Kane-Mele model. Computing the local magnetization of a nanoribbon with various widths, we obtain the edge phase diagram with four different phases for the ribbons of finite width and study how the magnetic properties are affected by electron-electron interactions. In particular, we show that the edge magnetic order survives for weak spin-orbit interactions in the thermodynamic limit. The edge magnetic order penetrates towards the ribbon center. The resulting magnetic profile gives its characteristic length which is found to be strongly affected by both electron-electron and spin-orbit interactions of its characteristic length turn out to be closely related to the phase transition between the topological insulator and the antiferromagnetic insulator. Electron dispersions are also discussed.

### RC09

### Theoretical study of spin texture in the Bi thin film

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The Spin-Orbit interaction (SOI) is very important effect for heavy atom like Bismuth. Recently, first-principles study revealed that significant SOI effects on the electronic structure and lattice bistability of Bi nanofilms.[1] In this study, we perform firstprinciples electronic structure calculations under the electric field for Bismuth thin film and Bi (001) multi-layer surface. We revealed that the SOI combined with electric field (so called Rashba effect) is very important in Bi surface state and Bi edge state.[2] We calculate spin polarization and spin direction for each band and k-points. Recently, giant out-of-plane spin component on Bi (001) surface is discovered SR-ARPES experiment.[3] We also calculate the spin structure of multi-layer Bi (001) surface, and our calculated result shows the detail of spin features at each k-point.

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## RC10

### A recipe for new Topological Insulators based on bonds, bands, symmetry and heavy atoms

- Lukas Muechler<sup>1</sup>, Binghai Yan<sup>2</sup>, Stanislav Chadov<sup>3</sup>, Haijun Zhang<sup>4</sup>, Frederick Casper<sup>1</sup>, Shoucheng Zhang<sup>4</sup> and Claudia Felser<sup>3</sup>\*
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In this work we will present a recipe to find new Topological Insulators (TIs) based on bonds, bands, symmetry and heavy atoms. A big issue concerning the compounds known up to now is the control of the bulk carrier density to produce truly insulating samples in the bulk. Using concepts from chemistry and supported by densityfunctional calcu- lations, we want to motivate an extended search for new compounds with tunable bulk properties.

### RC11

# Optimizing the Bi<sub>2-x</sub>Sb<sub>x</sub>Te<sub>3-y</sub>Se<sub>y</sub> solid solutions to approach the intrinsic topological insulator regime

Zhi Ren\*, Alexey Taskin, Satoshi Sasaki, Kouji Segawa and Yoichi Ando ISIR, Osaka University, Japan

Three dimensional (3D) topological insulators represent a new quantum state of mater realized in band insulators, which are predicted to present novel surface transport phenomena. In reality, while a number of materials have been identified to be 3D TIs, most of them are poor insulators in the bulk. Given this situation, search for new TI materials better suited for achieving the bulk insulating state is of obvious importance. In this presentation, we show that there exist a series of "intrinsic" compositions for Bi<sub>2.5</sub>Sb<sub>x</sub>Te<sub>3.5</sub>Se<sub>y</sub> solid solutions where the acceptors and donors compensate each other and present a maximally bulk-insulating behavior. At such compositions, the resistivity can become as large as several  $\Omega cm$  at low temperature, and one can infer th role of the surface-transport channel in the nonlinear Hall effect[1]. In particular, in the composition of Bi<sub>1.5</sub>Sb<sub>0.5</sub>Te<sub>1.5</sub>Se<sub>1.3</sub>, one can achieve a surface-dominated transport where the surface channel contributes up to 70% of the total conductance[2]. In addition, angle-resolved photoemission spectroscopy reveal that the Dirac cone dispersion changes systematically so that the Dirac remers at x-0.9[3].

[1] Z Ren, A.A. Taskin, S. Sasaki, K. Segawa, and Y. Ando, Phys. Rev. B 84, 165311 (2011). (Editors' Suggestion) [2] A.A. Taskin, Z. Ren, S. Sasaki, K. Segawa, and Y. Ando, Phys. Rev. Lett. 107, 016801 (2011). [3] T. Arakane, T. Sato, S. Souma, K. Kosaka, K. Nakayama, M. Komatsu, T. Takahashi, Z. Ren, K. Segawa, and Y. Ando, Nat. Commun., in press.

## **RC12**

### **Exploration of Three-dimensional Rashba Materials**

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Strong spin-orbit interactions under the broken space inversion symmetry can induce spontaneous spin polarization without external magnetic fields, which is known as the (surface/interface) Rashba effect. On the other hand, due to the lack of a model compound, the Rashba effect in three-dimensions (bulk) has been poorly explored. In this study, it is demonstrated that ternary compounds of bismuth tellurohalides, BiTeX (X = Cl, Br, and I), are promising candidate compounds as ideal threedimensional Rashba materials. As the essential backbone for both fundamental and applied research, this study established the growth of their sizable single crystals (>  $1 \times 1 \times 0.2$  mm<sup>3</sup>). Furthermore, it was found that, for BiTeI, that electronic state (the carrier concentration) can be adjusted from metallic to insulating ones by the choice of the growth techniques. Realization of the 3D Rashba states, including the Fermi surface topology for the corresponding carrier concentrations of the obtained crystals, was confirmed by relativistic first-principles calculations. We will discuss that our 3D Rashba materials can be a leading player in spin-involved novel phenomena, ranging from the metallic extreme (unconventional superconductivity) to the transport intermediate (spin Hall effects) to the novel insulating variant (3D topological insulating states)

# RC13

# Quantum phase transition from normal to topological insulator phase in Na, $IrO_3$

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 $Na_{2}IrO_{3}$  has been suggested to have a possible quantum spin Hall effect based on the novel jeff=1/2 state arising from the 5d Ir states. From the electronic structure calculation analysis of the band topology of  $Na_{2}IrO_{3}$ , however, it turns out that  $Na_{2}IrO_{3}$ as it is should be a trivial insulator. We investigated the electronic structure of by employing a tight-binding model for the layered iridium oxides with honeycomb lattice. Our analysis for the layered iridates reveals that the topological nature of the spin-orbit coupled ground state depends on the trigonal crystal field and long-range hoppings. From the first-principles-derived tight-binding Hamiltonian we determine the phase boundary through the parity analysis and predict a quantum phase transition from normal to topological insulator. It is suggested that  $Na_{2}IrO_{3}$  can be a candidate material which can have both non-trivial topology of bands and strong electron correlations.

## RC14

## Self-assembled nanowire with giant Rashba spin splitting

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We found gaint Rashba-type split bands for Pt-induced nanowires on the Si(110) surface by angle resolved photoemission spectroscopy (APRES) measurements. The observed Rashba parameter is about 2.1eVA, which is among the largest ever reported for Rashba systems and is unique in one-dimensional systems. Pt-induced nanowires were also investigated by scanning tunneling microscopy (STM) and scanning tunneling spectroscopy (STS) at low temperature (78K). High resolution STM topography images show that the wire has a width of 1.6mm and a X3 superstructure along the wire. On the other hands, the STS measurements reveal well resolved local density of states (LDOS) with a peak at -200meV and a characteristic dip at -300meV on the nanowire. These peak and dip energies corresponds well to the ARPES data; the top edge and the Dirac-like crossing point energy of the Rashba-type spin split bands, respectively. Furthermore, based on the dI/dV spectroscopic mapping, we confirmed noticeable one-dimensional electron channels along the nanowire edges, which are related to the gaint Rashba band. This implies a future application of the Pt-induced nanowire structure in silicon-based spintronics devices.

## RC15

# Ab initio study of topological surface state on Sb (111) surface with iron impurities

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We study iron impurities on Sb (111) surface and their effects on topological surface state by using an ab-initio pseudopotential density-functional method. We implemented the spin-orbit interaction into the SIESTA in a form of additional fully non-local projectors. To calculate electronic structure of topological surface states, we consider a slab of Sb using a supercell containing 20 atomic layers with experimental bulk Sb lattice parameters. We determine atomic positions of Fe impurities on Sb (111) surface by minimizing the total energy, and calculate surface band structures near the Fermi level. To find the impurity effects on the electronic structure of the surface states, we simulate ARPES spectra as a function of impurity density on the surface and calculate local density of states (LDOS) near the Fermi level. From the results, we find that Fe impurity states are present near Fermi level and they strongly interact with the surface states and produce a characteristic feature in LDOS on the top surface Sb layer. This work was supported by the NRF of Korea (Grant No. 2011-0018306) and KISTI Supercomputing Center (Project No. KSC-2011-C3-06).

# **RC16**

# Topological insulating phase in cubic system: tight-binding approach

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Recently, much attention has focused on three-dimensional strong topological insulators as a new quantum state of matter, such as  $Bi_2Se_3$  and  $Bi_2Te_3$ . Here we suggest topologically non-trivial electronic structures in cubic systems based on the tight-binding analysis. Assuming the well separated s and p states and hopping parameters, we can set up an 8x8 tight-binding Hamiltonian, and the trivial to non-trivial topological phase transition occurs by tuning the hopping parameters or spin-orbit coupling strength. Our tight-binding Hamiltonian is reduced to a 4x4 continuum model which is the minimal model to describe the topological phase transition. Topologically protected surface states in the slab geometry are also shown within our model.

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### **RC17**

# Topological insulator phase and Kitaev-like anisotropic exchange interactions in Li,IrO<sub>3</sub>

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From first-principles electronic structure calculations results, we set up an effective tight-binding Hamiltonian for Li<sub>2</sub>IrO<sub>3</sub> by constructing Wannier orbitals of lowenergy states near the Fermi level. The Wannier orbitals are found to be quite close to the  $S_j_{eff}^{=1/2}$  state, which are recently suggested to exist in many other iridate compounds. The model Hamiltonian is characterized bythe nontrivial second and third nearest neighbor hopping inside Ir honeycomb lattice as well as relatively small but significant inter-layer hopping terms. It is shown that the ratio between the nearest and third nearest neighbor hopping determines the  $S_2_2$ -sinvariant of this material, and that it can be controlled by the lattice strain in our first-principles calculations. We also suggest realistic exchange interactions from effective Hamiltonian, which has various nontrivial terms including the most strong Kitaev-type anisotropic exchange terms within the same sublattice, rather small but significant Dzyaloshinskii-Moriya terms, and Heisenberg exchange interactions.

## **RC18**

### Structural and electrical properties of (111) oriented half-Heusler La-Pt-Bi thin films

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Topological insulators are candidate materials for low-power and high-speed spintronics devices, due to the dissipation less and spin-polarized electrons on the surface state. Recently, we have reported structural properties of LaPtBi(001) thin films, one of candidate materials of half-Heusler topological insulator[1-2], grown by three-target co-sputtering[3]. In this paper, we report preparation and properties of LaPtBi with different crystalline orientation.By optimizing deposition rate for each element, (111) oriented taPtBi thin films were obtained below 500°C on Sapphire-C substrates. On the other hand, for films on YAlO<sub>3</sub> substrates, both (001) and (111) orientations were observed below 500°C. X-ray photoelectron spectroscopy study is revealed that (111) oriented films exhibited chemical shifts and appearance of satellite peaks of La 3d, which are attributed to charge transfer from Bi p and Pt d states to La d and f states due to the mixture of these states in the valence band. Investigation and comparison of the behaviors of LaPtBi thin films with different orientations will enable us to uncover the electron and spin transport properties in the surface state due to the topological nature.

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# RC19

# First-principles study of spin texture in the multilayer graphene on Ni(111)

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Graphene, a two-dimensional honeycomb lattice consisting of C atoms, shows useful properties for spintronics applications. Recently, spin-split electronic states in the graphene on Ni(111) is discussed extensively [1-3]. By using noncollinear spindensity functional theory implemented in openmx code[4], we predict spin-polarized electronic states of the graphene on Ni(111) in real- and momentum-space. The Diraccone shaped band structure was broken by strong hybridization between monolayer graphene and substrate Ni(111). On the other hand in the bilayer graphene on Ni(111), the Dirac-cone shaped band recover only in minority spin. The relativistic spin-orbit effect make non-collinear spin vortex (Rashba Effect) in the momentum space. This spin noncollinearity in momentum space may make significant effects on the electrical transport properties of the graphene/Ni(111).

[1] Yu. S. Dedkov, M. Fonin, U. Rudiger, and C. Laubschat, Phys. Rev. Lett. 100, 107602 (2008). [2] O. Rader, A. Varykhalov, and J. Sanchez-Barriga, D. Marchenko, A. Rybkin, and A. M. Shikin, Phys. Rev. Lett. 102, 057602 (2009), [3] K. Sawada, F. Ishii, M. Saito, Phys. Rev. B 82, 245426 (2010). [4] T. Ozaki, H. Kino, J. Yu, M. J. Han, N. Kobayashi, M. Ohfuti, F. Ishii, T. Ohwaki, H. Weng, M. Toyoda, and K. Terakura, http:// www.openmx-square.org/

## RD01

### Single Crystal Growth, Electrical and Magnetic Properties in ROs<sub>2</sub>Al<sub>10</sub> (R=rare earth)

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The ternary rare-earth compounds  $ROs_2AI_{10}$  (R = rare earth) crystallize with orthorhombic structure [1]. Among them,  $CeOs_2AI_{10}$  have attracted much attention because of the heavy-fermion behavior accompanied by a novel long-range order (LRO) at 28.7 K [2]. In order to understand the mechanisms of LRO, the information of physical properties of the reference compounds  $LaOs_2AI_{10}$  and  $PrOs_2AI_{10}$  is required. In this work, we have succeeded in growing single crystals of  $LaOs_2AI_{10}$  and  $PrOs_2AI_{10}$  by Al self-flux method, and investigated the physical properties by the X-ray diffraction, electrical resistivity, Hall effect and magnetic susceptibility. Both of  $LaOs_2AI_{10}$  and  $PrOs_2AI_{10}$  show the typical metallic temperature dependence of electrical resistivity without any indication of phase transitions.  $LaOs_3AI_{10}$  exhibits the Pauli paramagnetism with small temperature dependence, while  $PrOs_2AI_{10}$  exhibits the Van Vleck paramagnetism with the curie-Wiess behavior at high temperature with a large magnetic anisotropy, which is most probably due to the crystalline electric field effect. All of these results indicate that the electronic structure of  $CeOs_2AI_{10}$  is different from those of  $LaOs_2AI_{10}$  and  $PrOs_2AI_{10}$  a

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### **RD02**

# Anisotropic transport and magnetic properties of CeZn<sub>11</sub> single crystals1

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Anisotropic measurements of the magnetic and transport properties have been performed on single crystals of CeZn<sub>11</sub> in three different crystallographic directions. M(H) for H\$\perp c\$ shows a broad metamagnetic transition at about 1.9 T with Ms~ 1.6 \$\mu\_B\$ per Ce. Analysis of the temperature dependence of M/H indicates strong magnetic anisotropy. Zero-field resistivity of CeZn<sub>11</sub> shows a broad maximum, characteristic of that of Kondo compounds, at around 10 K followed by a sharp change of the slope and a kink corresponding to antiferromagnetic (AF) transition at 2.0 K. Anisotropic temperature and field dependent magnetoresistance and H-T phase diagrams will be discussed. \$^1\$ Work at the Ames Laboratory was supported by the Department of Energy, Basic Energy Sciences under Contract No. DE-AC02-07CH11358.

# RD03

### Shubnikov-de Haas oscillation in PuIn<sub>3</sub>

Yoshinori Haga<sup>1</sup>, Oscar Ayala-valenzuela<sup>2</sup>, Ross Mcdonald<sup>2</sup>, Chuck Mielke<sup>2</sup>, Eric D. Bauer<sup>2</sup>, J N Mitchell<sup>2</sup>, P. H. Tobash<sup>2</sup>, Joe D. Thompson<sup>2</sup> and Zachary Fisk<sup>1</sup> <sup>1</sup> Advanced Science Research Center, Japan Atomic Energy Agency, Japan <sup>2</sup> Los Alamos National Laboratory, USA



### RD04

## Quadrupolar ordering in a caged compound PrOs<sub>2</sub>Zn<sub>20</sub>

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Pr-based intermetallic compounds with cubic crystal structures are expected to show various phenomena arising from quadrupolar degrees of freedom, when the crystal electric field (CEF) ground state of  $Pr^{3+}$  ion is the nonmagnetic F3 doublet. Recently, it has been found that a caged compound PrIr<sub>2</sub>Zn<sub>20</sub> with the F3 doublet ground state undergoes a superconducting transition at Tc=0.05 K in the presence of antiferroquadrupolar order below TQ=0.11 K [1, 2]. We have studied the isostructural compound PrOs<sub>2</sub>Zn<sub>20</sub>, with 5d electron one less than in PrIr<sub>2</sub>Zn<sub>20</sub> by the measurement of the electrical resistivity  $\rho$ , magnetic susceptibility  $\chi$ , and specific heat C in the temperature range from 0.4 K to 300 K. On cooling below 5 K,  $\chi(T)$  approaches a constant value, indicating a nonmagnetic CEF ground state. C(T) shows a peak at To=0.6 K, where no anomaly appears in the  $\chi(T)$ . These findings strongly suggest a nonmagnetic phase transition at To arising from the quadrupolar degrees of freedom. Furthermore, the To is five times higher than TQ in PrIr<sub>2</sub>Zn<sub>20</sub>, which is possibly attributed to enhancement of quadrupolar exchance interaction.

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### **RD05**

# Spin-triplet superconductivity induced by ferromagnetic fluctuations in UCoGe

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After the discovery of superconductivity in the ferromagnetic (FM) UGe<sub>2</sub> under pressure, FM superconductors have attracted much interest, since in an itnerant FM superconductor with the presence of a large energy splitting between the majority and minority spin Fermi surfaces, exotic spin-triplet superconductivity has been anticipated. Among such FM superconductors, UCoGe [1] is one of the most readily explored experimentally, because of its high superconductors, UCoGe [1] is one of the most readily explored experimentally, because of its high superconducting transition temperature ~ 0.7 K and low Curie temperature ~ 2.5 K at ambient pressure. Large anisotropic behaviors of superconductivity [2, 3] and longitudinal FM fluctuations along the c axis (direction of the ordered moment) [4] are thought to link with the mechanism of the superconductivity. From precise angle-resolved NMR and Meissner measurements we found that the magnetic field along the c axis strongly suppresses the FM fluctuations and that the superconductivity is observed in the limited magnetic field region where the FM spin fluctuations tuned by H [] c induce the unique spin-triplet superconductivity in UCoGe. As far as we know, this is the first clear experimental example that FM fluctuations are intimately related with superconductivity.

N. T. Huy et al., Phys. Rev. Lett. 99 (2007) 067006. [2] N. T. Huy et al., Phys. Rev. Lett. 100 (2008) 077002. [3]
 D. Aoki et al., J. Phys. Soc. Jpn. 78 (2009) 113709. [4] Y. Ihara et al., Phys. Rev. Lett. 105 (2010) 206403.

# **RD06**

# Competitive magnetic properties between the different anisotropic

SDW phases in heavy-fermion system Ce<sub>0.87</sub>La<sub>0.13</sub>(Ru<sub>1-x</sub>Rh<sub>x</sub>)<sub>2</sub>Si<sub>2</sub> Hiroaki Okamoto<sup>1</sup>, Eiichiro Harada<sup>1</sup>, Yusuke Amakai<sup>1</sup>, Shigeyuki Murayama<sup>1</sup>\* Hideaki Takano<sup>1</sup>, Naoki Momono<sup>1</sup>, Kazuyuki Matsubayashi<sup>2</sup> and Yoshiya Uwatoko<sup>2</sup> <sup>1</sup> Graduate School of Engineering, Muroran Institute of Technology, Japan <sup>2</sup> Institute for Solid State Physics, University of Tokyo, Japan

Typical heavy-fermion compound CeRu<sub>2</sub>Si<sub>2</sub> shows an SDW transition by the substitution of La at the Ce site with the gap perpendicular to the c-axis and the moment along the c-axis(1). It shows another SDW transition with the gap and moment both parallel to the c-axis by the substitution of Rh at the Ru site(2). In order to investigate the two different SDW phases, we have measured the susceptibility  $\chi$  and resistivity p on single-crystalline Ce<sub>0.87</sub>La<sub>0.13</sub>(Ru<sub>1-x</sub>Rh<sub>2</sub>)Si<sub>2</sub> for  $0 \le x \le 0.4$  in directions perpendicular (2) and parallel (//) to the c-axis(3). The  $\rho \angle$  for  $0 \le x < 0.05$  shows a gap-type anomaly at TN, whereas the  $\rho //$  for  $0 \le x < 0.05$  indicates a downward inflection or a shoulder at around TN, suggesting the transverse SDW transition. For  $0.05 \le x \le 0.075$  these anomalies become very small and a hysteresis is further seen between temperature decreasing and increaseing process below TN. For  $0.07 \le x \le 0.3$  the gap-type anomaly in  $\rho \angle$  disappears and a clear gap-type jump has been detected in  $\rho //$  at TN, suggesting the longitudinal SDW transition. These behaviors would suggest that a competitive magnetic phase occurs between the two anisotropic SDW phases.

(1) S. Quezel et. al., J. Magn. & Magn. Mater., 76 & 77, 403 (1988). (2) S. Murayama et. al., Phys. Rev. B 56, 11092 (1997). (3) Y. Amakai et. al., J. Phys. Soc. Jpn. 80, SA062 (2011).

### **RD07**

### Elastic constant of SmOs<sub>4</sub>Sb<sub>12</sub> under high magnetic field

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The filled skutterudite compound SmOs<sub>8</sub>Sb<sub>12</sub> exhibits a large electronic specific heat coefficient Y ~ 820 mJ K<sup>2</sup> mol<sup>-1</sup>, which is insensitive to an applied magnetic field H || <001>, and also shows a phase transition at 2.5 K with tiny ferromagnetic moment. The origin of the magnetically robust heavy fermion state and an exotic phase transition are unsolved issues. The crystalline electric field (CEF) ground state of this compound has also not been settled, thus far. Owing to clarified whether the CEF ground state is the  $\Gamma$  8 quartet or  $\Gamma$  7 doublet for Sm<sup>3+</sup> ionic configurations, we have measured the elastic constant C\_11 on single crystalline SmOs<sub>8</sub>Sb<sub>12</sub> by using pulsed magnetic fields up to 62 T. As a result, C\_11 shows a minimum at around 10 T and gradually increase with increasing magnetic field at 4.2 K. In the cubic symmetry, C\_11 can be described as the sum of the bulk modulus C\_B and the  $\Gamma_3$ -type quadrupole susceptibility. Assuming that C\_B is only dominated by magneto-elastic effects of the weak ferromagnetic moment in the low magnetic field region, the increasing tendency of C\_11 at high magnetic field region could be explained by a quadrupole susceptibility with  $\Gamma_8$  quartet CEF ground state.

[1] S. Sanada et al., J. Phys. Soc. Jpn. 74 (2005) 246, [2] W. M. Yuhasz et al., Phys. Rev. B 71 (2005) 104402, [3]
 Y. Aoki et al., Physica B 378-380 (2006) 54, [4] T. Yanagisawa et al., J. Phys. Soc. Jpn. 80 (2011) 043601.

### **RD08**

# 27Al-NMR study for critical phenomena of metamagnetic transition in UCoAl

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UCoAl shows a first-order paramagnetism (PM)-ferromagnetism (FM) transition (metamagnetic transition) in small external magnetic field about 0.6 T applied only along its easy magnetization axis, e-axis [1]. This first order transition line terminates at a "critical-end-point" at Tcr ~ 13 K. Above Tcr the first order transition from PM to FM changes to crossover. This feature is apparently similar to gas-liquid transition, where the order parameter and the tuning parameter are density of molecules and pressure, respectively. In fact, a gas-liquid transition is regarded to be mathematically equivalent to 3D-Ising FM transition [2], therefore we expect the metamagnetic transition in UCoAl has the same physical framework as a gas-liquid transition. In order to investigate static and dynamic properties in the critical phenomena on UCoAl, we performed 27AI-NMR measurement for a single-crystal UCoAl and measured Knight-shift K and nuclear spin-lattice relaxation rate 1/ T1 by varying temperature T and c-axis magnetic field Hc. In the poster presentation, I will show the result of c-axis magnetization Mc (order parameter) and c-axis dynamical susceptibility Sc (fluctuations of order parameter) calculated from K and 1/T1T, as a function of (T, Hc) and discuss similarities and differences from comparison with gas-liquid transition.

[1] N. V. Mushinikov et al., Phys. Rev. B 59, 6877 (1999). [2] T. D. Lee et al., Phys. Rev. 87, 410 (1952).

## RD09

# Hybridization gap and the hidden order in the heavy fermion Kondo lattice $URu_2Si_2$

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What drives the second order phase transition that occurs at 17.5 K in the Kondo lattice heavy fermion URu<sub>2</sub>Si<sub>2</sub> remains unknown after two and half decades of research despite ubiquitous observation of gap-like behaviors. Determining the nature of the extracted gaps, whether they are the order parameter or not, is a crucial task to resolving this hidden order enigma. Quasiparticle scattering spectroscopy (QPS), better known as point-contact spectroscopy, is a powerful technique to probe the bulk electronic properties by exploiting a ballistic junction between two electrodes. Employing QPS, we have carried out conductance measurements on URu<sub>2</sub>Si<sub>2</sub> [1]. The differential conductance data exhibit a distinct double-peak structure with pronounced asymmetry, a signature for a Fano resonance in a Kondo lattice [2]. The extracted gap, a hybridization gap between the renormalized bands, opens well above the hidden order transition temperature, indicating this hybridization gap is not the hidden order transition in URu<sub>2</sub>Si<sub>2</sub>. The work at UIUC is supported by the U.S. DOE under Award No. DEF-FO02-07ER46453. The work at LANL is carried out under the auspices of the U.S. DOE, Office of Science.

[1] W. K. Park et al., PRL 108, 246403 (2012) [2] M. Maltseva, M. Dzero, P. Coleman, PRL 103, 206402 (2009).

### **RD10**

### Resonant Raman effect on LaRu<sub>2</sub>Al<sub>10</sub> and CeRu<sub>2</sub>Al<sub>10</sub>

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LaRu<sub>2</sub>Al<sub>10</sub> and CeRu<sub>2</sub>Al<sub>10</sub> with space group Cmcm have the orthorhombic YbFe<sub>2</sub>Al<sub>10</sub> type structure [1], where the number of Raman active phonons is estimated as 10Ag + 9B1g + 6B2g + 8B3g by group theory. CeRu<sub>2</sub>Al<sub>10</sub> shows the transition at T0 = 27 K [2, 3, 4], while LaRu<sub>2</sub>Al<sub>10</sub> shows no anomaly. We reported the temperature dependence of the phonon Raman spectra at the SCES 2011, where we could not observe all phonon modes by only 514.5 nm excitation and no anomaly at T0 in CeRu<sub>2</sub>Al<sub>10</sub>. We concluded that the transition at T0 does not accompany with the structural change. Recently we have employed several excitations light at 488.0 nm, 514.5 nm, 568.2 nm, and 647.1 nm. We have observed the B1g mode missed by only 514.5 nm, that is, we observe all phonon modes by changing the excitation energy. In the guest atoms modes, the vibration along b axis shows the intensity enhancement at 647.1 nm, but that along a or c axis does not show the case. Remarkable differences between LaRu<sub>4</sub>Al<sub>10</sub> and CeRu<sub>4</sub>Al<sub>10</sub> is observed for the phonons at around 100 cm-1 in (y, z) polarization. These differences are onisinated by the difference of the electronic states between them.

[1] Thieda V M T, et al, 1998 J. Mater. Chem. 8 125. [2] Strydom A M 2009 Physica B 404 2981 [3] Nishioka T, et al 2009 J. Phys. Soc. Jpn. 78 123705 [4] Muro Y, et al 2010 J. Phys.: Conf. Series 200 012136

# RD11

### Magnetic properties of β-US2 single crystals

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Uranium chalcogenides compounds UXS (X:S, Se and Te) have the same orthorhombic (Pnma) structure and those electronic state can easily be modified by changing the constituting elements. UTeS is an Ising-type ferromagnet with a Curie temperature 87 K and shows metallic conductivity at all temperatures. USeS shows a semimetallic resistivity. Electrical resistivity of USeS decreases with decreasing temperature but increases steeply below 50 K, with an anomaly at the ferromagnetic ordering temperature 24 K.US2 is similar to USeS at high temperature and the electrical resistivity of US2 decreases with decreasing temperature and increases steeply below about 80 K, but US2 does not order magnetically. The electrical resistivity of US2 at low temperature suppresses under a high pressure or a high magnetic field [1]. We study the magnetic properties of US2 in detail by measuring magnetic susceptibility and electric resistivity under a magnetic field.

[1] S. Ikeda et al., J. Phys. Soc. Jpn. 76 (2009) 114704.

# RD12

# Evidence of nodal gap structure in the skutteride superconductor $\mbox{PrPt}_4\mbox{Ge}_{12}$

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In order to study the gap function of the newly discovered skutterudite superconductor PrPt\_4Ge\_{12} (T\_{c}=7.9K), we have measured the London penetration depth  $\Delta\lambda(T)$  of the single crystalline PrPt\_4Ge\_{12} down to 0.05Tc by using a tunnel diode oscillation (TDO) based technique. Observation of a quadratic temperature dependence of  $\Delta\lambda$  below 0.3Tc indicates the existence of point nodes in the superconducting gap. The derived superfluid density  $\rho_{s}(T)$  can be nicely fitted by the A-phase gap function of 3He  $(\Delta(\theta, \phi) = \Delta_{s}(0) | K_{s}(x) + iK_{s}(y)|)$  with a fitting parameter of  $\Delta_{s}(0) = 2.2k_{s}(B)T_{c}(c)$ , providing new evidence for a possible p-wave superconductivity in PrPt\_4Ge\_{12}.

# RD15

## Effect of pressure on the YbNi<sub>3</sub>Ga<sub>9</sub> single crystal

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We have investigated the effect of pressure on the physical properties of singlecrystalline YbNi<sub>3</sub>Ga<sub>9</sub>. This compounds crystallize in a trigonal ErNi<sub>3</sub>Al<sub>9</sub>-type structure (space group R32). From the thermal, magnetic and transport study, YbNi<sub>3</sub>Al<sub>9</sub> turns out to be an intermediate valence compound with r = 30 mJ/molK [1]. The electrical resistivity and ac susceptibility of nonmagnetic heavy fermion compound YbNi<sub>3</sub>Al<sub>9</sub> have been measured under hydrostatic pressure up to 11 GPa. With increasing pressure, the ferromagnetic ordering state was appeared over 9 GPa. These results suggested that nonmagnetic Yb state change to the localized magnetic state, via quantum critical point, by tuning the pressure. In this paper, we will discuss the electrical state of YbNi<sub>3</sub>Al<sub>9</sub> under pressure.

[1]T. Yamashita and S. Ohara; submitted to J. Phys. Soc. Jpn.

## RD13

Mossbauer spectroscopy of Fe doping valence fluctuating α-YbAlB<sub>4</sub> Yui Sakaguchi<sup>1\*</sup>, Shugo Ikeda<sup>1</sup>, Kentaro Kuga<sup>2</sup>, Keita Sone<sup>2</sup>, Satoru Nakatsuji<sup>2</sup> and Hisao Kobayashi<sup>1</sup>

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Valence fluctuating YbAlB<sub>4</sub> compounds have two different crystal structures; α-YbAlB<sub>4</sub> and β-YbAlB<sub>4</sub> [1]. β-YbAlB<sub>4</sub> is the first compound which shows a Yb-based heavy fermion superconductor with Tc = 80mK [2] and exhibits a pronounced non-Fermi-liquid behavior above Tc. Meanwhile, α-YbAlB<sub>4</sub> shows normal metallic behaviorsat low temperature. Recently, the small substitution of Al by Fe in α-YbAlB<sub>4</sub> revealed a breakdown of the Fermi-liquid behavior with an antiferromagnetic order. We have studied magnetic properties of  $\alpha$ -Yb(Al<sub>0.75</sub>Fe<sub>0.25</sub>)B<sub>4</sub> at low temperatures using 57Fe Mossbauer spectroscopy. The observed Mossbauer spectrum at room temperature using single crystalline samples is typical for a paramagnetic state. We analyzed the spectrum using a doublet. From the intensity ration of the doublet, the principle z-axis of the diagonalized electric-field gradient tensor is perpendicular to the c-axis. We observed a small magnetic hyperfine field at 4.5K whitch comes from the antiferromagnetic ordering of Yb magnetic moments. This small magnetic hyperfine field disappears around TN (~9K). These results indicate that the magnetic structure is not simple in α-Yb(Al<sub>075</sub>Fe<sub>025</sub>)B<sub>4</sub>. Above TN, the observed spectra change drastically with increasing temperature. Some of changes in the spectra were related to anomalies in the temperature dependence of magnetization.

[1] R. T. Macaluso et al., Chem. Mater. 19. 1918 (2007) [2] S. Nakatsuji et al., Nature Phys. 4, 603 (2008)

### **RD14**

### Neutron scattering study on f-electron states of PrCu<sub>4</sub>Au

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Pr-based compounds have recently been studied for heavy electron properties originating from 4f<sup>2</sup> electron configuration, as in PrFe<sub>4</sub>P<sub>12</sub> [1]. We extracted a large electronic-specific-heat-coefficient of 0.77 J/mol K<sup>2</sup> from specific heat measurements on PrCu<sub>4</sub>Au [2]. This material is expected to undergo an antiferromagnetic transition below 24 K. We investigate the role that f-electron states play in these properties using neutron scattering techniques. We used the spectrometers SEQUOIA at ORNL, 4F1 at LLB, and TOPAN at JAEA to measure the temperature dependence, between 1.6 K and 100 K, of the magnetic excitations in PrCu<sub>4</sub>Au. The spectra are composed of several inelastic peaks. The crystal-fieldsplitting scheme with the triplet ground state, as suggested from the magnetization result [2], is a fair match to the observed spectra. A quasielastic component was also detected. The temperature dependences of the spectral intensity and width are similar to those in other heavy electron systems. Diffraction experiments for the powdered sample on the diffractometer G4-1 at LLB reveal antiferromagnetic peaks characterized by the wave vector (1/2, 1/2, 1/2). This result can be interpreted as the type-II magnetic ordering on the fact sub-lattice of Pr ions. To summarize, PrCu<sub>4</sub>Au exhibits localized and strongly-hybridized Felectron states.

[1] H. Sugawara et al.: Phys. Rev. B 66 (2002) 134411. [2] S. Zhang et al.: J. Phys.: Condens. Matter 21 (2009) 205601

### Fermi surfaces in the mixed valent Yb system Hisatomo Harima\*

**RD16** 

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Many Yb-compounds have been crystalized recently in good quality, so that the de Haas-van Alphen effect is detected in those including the Yb-based heavy fermion compounds, i.e., YbCu<sub>2</sub>Si<sub>2</sub>, YbAlB<sub>4</sub>. Non-magnetic Yb compounds are regarded mostly as divalent Yb system, in which Fermi surfaces are well reproduced by using conventional band structure calculations. In some of non-magnetic Yb compounds, such as YbAl<sub>3</sub> the valence of Yb-ion is diverted from Yb<sup>2+</sup>. Actually, the Fermi surfaces of YbAl<sub>3</sub> are very different in topology from those of divalent Ybln<sub>3</sub>. This is the clear evidence of the 4f electrons really affect the formation of the Fermi surfaces in YbAl<sub>3</sub>[1]. In Yb compounds, the dispersion of the 4f electrons. Recently the LDA+U approach has shown that only the doublet ground states participate the Fermi surfaces in YbAl<sub>8</sub>[2]. The role of 4f electrons in mixed valent Yb system is discussed, for YbAl<sub>4</sub>, YbCoGa<sub>5</sub> and the related compounds.

[1] T. Ebihara, et. al., J. Phys. Soc. Jpn. 69 (2000) 895. [2] Nguyen Duc Dung, et. al., J. Phys. Soc. Jpn. 78 (2009) 084711.

### RD17

# Unusual heavy fermion behavior in $PrTr_2Al_{20}$ (T = Nb, Ta) associated with $\Gamma_3$ quadrupolar degrees of freedom

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Recently,  $PrTr_{2}X_{20}$  (Tr = transition metal, X = Zn, Al) compounds have attracted much attention as a new candidate for investigating the quadrupolar Kondo effect because PrTr<sub>2</sub>Al<sub>20</sub> (Tr = Ti, V) show non Fermi liquid behaviors and multipolar orderings attributed to nonmagnetic  $\Gamma_3$  doublet ground state[1]. Recently, we found that  $PrNb_2Al_{20}$  also has a nonmagnetic  $\Gamma$  3 doublet ground state and does not show any phase transitions down to 0.15 K[2]. In addition, this compound has a large Sommerfeld coefficient  $\gamma$ (~ 1.2 J/molK<sup>2</sup>). This result suggests a possibility of the formation of unusual heavy fermion state attributed to the quadrupole moments of  $\Gamma_3$  state. In order to investigate the quadrupole-induced strongly correlated electronic properties in PrTr2Al20 systematically we have grown single crystals of other Pr based compound, PrTa<sub>2</sub>Al<sub>20</sub>, and performed specific heat, magnetization and electronic transport measurements. In C/T, PrTa2Al20 shows a peak around 0.6 K attributed to the ordring of quadrupole moment and has a large gamma value of 1.6 J/molK<sup>2</sup> at low temperatures with little field dependence. In this presentation, we show recent results for  $PrT_{2}Al_{20}$  (T = Nb. Ta) and discuss the possibility of the formation of heavy fermion formation induced by quadrupolar degrees of freedom.

[1] A. Sakai and S. Nakatsuji: J. Phys. Soc. Jpn. 71 063701 (2011). [2] R. Higashinaka et al.: JPS Autumn meeting 23pPSA-37 (2010), JPS Spring meeting 28aEC-11 (2011).

# RD18

# Tuning of the heavy-fermion ground state in YbNi<sub>3</sub>X<sub>9</sub> (X=Al, Ga) by substitution

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We have succeeded in synthesizing new Yb-based Kondo lattice systems YbNi<sub>3</sub>X<sub>9</sub> (X =Al, Ga) which crystallizes in the trigonal ErNi<sub>3</sub>Al<sub>9</sub>-type structure [1]. YbNi<sub>3</sub>Al<sub>9</sub> shows typical features of a heavy-fermion antiferromagnet with Neel temperature of TN=3.4K and Sommerfeld coefficient  $\gamma$ -100mJ/(molK2). In contrast, YbNi<sub>3</sub>Ga, is an intermediate valence compound with  $\gamma$ =30mJ/(molK2) and a high characteristic temperature of TK~600K. YbNi<sub>3</sub>X<sub>9</sub> (X =Al, Ga) is a great opportunity to investigate the tunability of heavy-fermion ground states in Yb-based compound for using external control parameters such as pressure or doping. In this regard, we are investigating physical properties of mixed crystals of YbNi<sub>3</sub>Al<sub>9</sub>. For YbNi<sub>3</sub>Al<sub>9</sub>, there we report Ni-site substitution effects on the antiferromagnetism of YbNi<sub>3</sub>Al<sub>9</sub>. For YbNi<sub>3</sub>Al<sub>9</sub>, the resistivity shows a -logT dependence below 100K and a broad peak at about 40 K arising due to the Kondo effect combined with the CEF excited state. In the lower-temperature region, a strong reduction in the extent of magnetic scattering associated with the antiferromagnetic ordering was observed below TN=3.4 K. By Fe-substitution TN is decreased, while is increased by Cu-substitution. The tunability of TN in Yb(Ni<sub>13</sub>Al<sub>9</sub>, J<sub>13</sub>A<sub>9</sub>, will be discussed.

[1] S. Ohara et al., J. Phys.: Conf. Series 273, 012048 (2011); T. Yamashita et al., to be published in J. Phys. Soc. Jpn.

### **RD19**

# Pressure effect studies in Ce<sub>2</sub>T<sub>3</sub>Ge<sub>5</sub> (T=Rh, Pd, Ir) by electrical resistivity

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In the Ce<sub>2</sub>T<sub>3</sub>Ge<sub>5</sub> system (T: transition metal) with the orthorhombic crystal structure, Ce<sub>2</sub>Ni<sub>3</sub>Ge<sub>5</sub> is an antiferromagnet with a Neel temperature  $T_N = 5.2$  K.  $T_N$  decreases with increasing pressure and becomes zero at a critical pressure  $P_c = 3.9$  GPa. The A and  $\rho_0$  values in the Fermi liquid relation (the low-temprature electrical resistivity  $\rho = \rho_0 + AT^2$ ) increase steeply above 3 GPa. It seems to be the heavy fermion state around Pc, in which pressure region superconductivity was found below 0.26 GPa [1]. We measured temperature dependence of electrical resistivity of Ce<sub>2</sub>Rh<sub>3</sub>Ge<sub>5</sub>, and Ce<sub>2</sub>Pd<sub>3</sub>Ge<sub>5</sub> under pressure in this study. These theree compounds are all antiferromagnets at ambient pressure, showing kinks of resistivity at  $T_N$  (=4.7, 9.5 and 3.8 K respectively)[2, 3, 4]. In Ce<sub>2</sub>Rh<sub>3</sub>Ge<sub>5</sub> and Ce<sub>2</sub>Ir<sub>3</sub>Ge<sub>5</sub>  $T_N$  decreases with increasing pressure and seems to be zero at a critical pressure  $P_c=0.4$  and 2.0 GPa, while  $T_N$  increases in Ce<sub>2</sub>Pd<sub>3</sub>Ge<sub>5</sub> by pressure. We also found another anomaly at lower temperature, indicating  $T_N$  splits into two ordering  $T_{N1}$  and  $T_{N2}$  under pressure more than 0.2 GPa.  $T_{N1}$  and  $T_{N2}$  are both increasing pressure up to 2.3 GPa.

[1] M. Nakashima, et al., J. Phys.: Condens. Matter 17 (2005) 4539. [2] C. Godart, et al., Solid State Commun. 67 (1988) 677. [3] B. Becker, et al., Phys. B 230-232 (1997) 253. [4] Z. Hossain, et al., Phys. Rev. B 60 (1999) 10383.

### **RD20**

### Photoemission study on new Kondo lattice compounds YbNi<sub>3</sub>(Ga<sub>1-x</sub>Al<sub>x</sub>)<sub>9</sub>

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Recently, new Kondo lattice compounds YbNi<sub>3</sub>X<sub>9</sub> (X=Al, Ga) were discovered [1]. YbNi<sub>3</sub>Al<sub>9</sub> is a heavyfermion antiferromagnetic compound (TN=3.4 K), while YbNi<sub>3</sub>Ga<sub>9</sub> is a valence fluctuating compound from the magnetic susceptibility measurements. These compounds are considered to be very suitable for research of Yb-based Kondo lattice systems because they both possess the ErNi<sub>3</sub>Al<sub>9</sub>-type crystal structure, and the Al and Ga ions are isovalent. In this study, we have carried out hard x-ray (hv-6 keV) and low energy (hv-7 eV) excited photoemission spectroscopy (HAXPES and LEPES) on YbNi<sub>3</sub>X<sub>4</sub>, and mixed crystals YbNi<sub>3</sub>(Ga, a,Al<sub>3</sub>), to investigate the Yb valence and the Kondo resonance behavior. The Yb 3d HAXPES experiments revealed that the Yb ion in YbNi<sub>3</sub>Al<sub>9</sub> exists with almost trivalent state while that in YbNi<sub>3</sub>Ga, strongly fluctuates with the averaged Yb valence of ~2.5. The Al substitution for YbNi<sub>3</sub>Ga, (x-0.1) changes the Yb ion toward trivalent state. The LEPES spectra of YbNi<sub>3</sub>Ga, clearly exhibit the Kondo resonance peak near the Fermi level (EF), and the peak is shifted toward EF and the intensity increases with lowering temperature. The extrapolated peak energy at zero temperature provides TK~600 K for YbNi<sub>3</sub>Ga, and the Al substitution lowers TK.

[1] S. Ohara et al., J. Phys.: Conf. Series 273, 012048 (2011); T. Yamashita et al., to be published in J. Phys. Soc. Jpn.

# RD21

### Heavy-fermion properties of YbCu<sub>5-x</sub>Au<sub>x</sub> (x = 0.5, 0.6, 0.7)

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Yb compounds are very attractive for research because of their low temperature anomalous behaviour. Recently, emphasis was given to the investigation of the solid solution  $VbCu_{x,M_x}(M = Ag, Au, In)$  crystallizing in the cubic AuBe<sub>c</sub>-type structure. The substitution of Cu by M offers the possibility of studying the different evolutions of the ground state depending on the M element (e.g. a valence transition in M = In, Kondo-lattice behavior in M = Ag, magnetic ordering in M = Au) [1,2,3]. Here we present the measurements of heat capacity and resistivity in the temperature scale 0.4 -300 K for different applied magnetic fields 0 - 9 T of the  $VbCu_{x,A}u_x$  (x = 0.5, 0.6, 0.7) system. The polycrystalline samples were prepared by induction-melting method, heat-treated and then quenched in cold water. Phase purity was checked by electron probe microanalysis (EPMA) and X-ray diffraction (XRD)

K. Yoshimura et al., J. Alloys Comp 262-263 (1997) 118 [2] M. Galli, E.Bauer, St. Berger, Ch. Dusek et al. Physica B 312-313 (2002) 489 [3] J. He, N. Tsujii et al., J. Phys. Soc. Jpn 66 (1997) 2481

## **RD22**

## $Magetic \ properties \ of \ heavy-fermion \ compounds \ Ce_{1.x}Lu_xRu_2Si_2$

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Typical heavy-fermion compound CeRu<sub>2</sub>Si<sub>2</sub> with a tetragonal structure shows an antiferromagnetic transition by the substitution of La at the Ce site due to the negative chemical pressure by the larger La ion(1). The substitution of Y introduces a decrease of the specific heat coefficient  $\gamma$  and an increase of the Kondo temperature TK due to the positive chemical pressure by the smaller Y ion(2). In order to study substitution effect of Lu with the closed f-electron shell and the ion smaller than Y, we have measured the susceptibility  $\chi$  and resistivity  $\rho$  on single-crystalline Ce<sub>1x</sub>Lu<sub>x</sub>Ru<sub>5</sub>Si<sub>2</sub> for  $0 \leq x \leq 0.2$  perpendicular ( $\angle$ ) and parallel (//) to the c-axis.The  $\chi$ // shows a maximum and the Tmax increases with x. Since the Tmax is supposed to be the heavy-fermion formation temperature proportional to  $1/\gamma$ , the effective mass meff is considered to increase with x. This indicates an increase of TK due to the positive chemical pressure by Lu. The low temperature resistivity follows the T<sup>2</sup> law and the coefficient A decreases with x. Since A is proportional to  $\gamma^2$ , meff is again considered to increase with x.

(1) S. Quezel et. al., J. Magn. & Magn. Mater., 76 & 77, 403 (1988). (2) M.J. Besnus et. al., J. Magn. & Magn. Mater., 63 & 64, 323 (1987).

# RD23

### Moment-bearing Tb substitution in CePt<sub>2</sub>Si<sub>2</sub>

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The effects of substituting moment-bearing Tb in CePt<sub>2</sub>Si<sub>2</sub> are reported through X-ray diffraction (XRD), electrical resistivity ( $\rho$ (T)), magnetic susceptibility ( $\chi$ (T)) and magnetization ( $\sigma$ ( $\mu$ 0H)) measurements. XRD results for all compositions of the (Ce<sub>1</sub>,  $_x$ Tb<sub>x</sub>)Pt<sub>2</sub>Si<sub>2</sub> system indicate a tetragonal CaBe<sub>2</sub>Ge<sub>2</sub>-type structure with space group P4/ mmm.  $\rho$ (T) results indicate the evolution from a coherent Kondo lattice to incoherent single-ion Kondo behavior up to 90% Ce substitution  $\chi$ (T) data at higher temperatures follow the Curie-Weiss relation for all alloy compositions and given effective moment values  $\mu$ eff which increase gradually from the expected value of 2.54  $\mu$ B for the Ce<sup>3+</sup>-ion to the expected value of 9.72  $\mu$ B for the Tb<sup>3+</sup>-ion. At low temperature  $\chi$ (T) data exhibit antiferromagnetism in the concentration range 0.75 x  $\leq$  1. No evidence of metamagnetic behavior was observed from  $\sigma$ ( $\mu$ 0H) data for all compositions.

# **RD24**

**RD25** 

Heavy Fermion Behavior of Yb<sub>2</sub>Ni<sub>12</sub>P<sub>7</sub> Studied by 31P NMR

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We present results of the nuclear magnetic resonance measurement (NMR) of Ybbased intermetallic compound Yb\_2Ni<sub>12</sub>P<sub>7</sub>, which crystallizes in the hexagonal Zr\_2Fe<sub>12</sub>P<sub>7</sub>type structure [1]. Yb<sub>2</sub>Ni<sub>12</sub>P<sub>7</sub> does not order magnetically down to 50 mK and forms the heavy fermion state. The electronic specific heat of Yb<sub>2</sub>Ni<sub>12</sub>P<sub>7</sub> is markedly enhanced at low temperatures (200 mJ/K2 mol-Yb). The magnetic susceptibility shows the Curie-Weiss like behavior at high temperatures and has a noticeable hump at around 50 K. We have performed NMR measurement to investigate the electronic properties of Yb<sub>2</sub>Ni<sub>12</sub>P<sub>7</sub>. The observed NMR line shows a complicated shape, reflecting the presence of three non-equivalent P sites and local symmetries around them. The nuclear-spin lattice relaxation rate (1/T1) shows nearly T1T=const. behavior, which is consistent with the delocalization of the 4f electrons.

K. Zeppenfeld and W. Jeitschko, J. Phys. Chem. Solids, 54 (1993) 1527.

### **R** 1 1

Engineered p-d exchange interaction of coupled double diluted magnetic quantum dots

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**RE01** 

This paper presents the effects of inserting a thick InGaN layer which was grown between n-GaN and active layer by MOCVD. By adding the thick InGaN layer, we expected the compensation of lattice mismatch between n-GaN and active layer. From AFM results, QDs (Quantum Dots) were formed at around 200 A InGaN thickness by TD (Threading Dislocation) termination effect. According to the result of TD termination, the electrical and optical properties were improved due to the decrease of non-radiative recombination area in active layer. The result of PL data shows PL intensity increased by inserting the thick InGaN layer.

Quasi-particle localization by disorder in an incompressible fractional

The strongly correlated electrons, together with disorder, organize themselves in a non-

trivial way to form incompressible collective states for the filling fraction v = 5/2 in

a dirty bi-layer semi-conductor quantum well system. The theoretical analysis of the

quantum phase transition(QPT) exhibited by this system has been carried out including

the single-particle quantum lifetime sensitive to the long range scattering and the usual

zero field short-range Drude scattering (transport lifetime). Since the tunneling terms

can be controlled by changing gate voltages, we find that, in principle, a completely

reversible sequence of events, viz. a crossover from the bi-layer fractional quantum

Hall state to the single-layer one followed by an entry to quantum critical state (and

the QPT to the bi-layer state beyond) and eventually a crossover to the single-layer

state, could be realized with the increase in inter-laver tunneling. The sequence, which

is manifestation of quasi-particle localization by disorder, could prove to be useful in

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### RE02

quantum Hall state

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### Electronic structure of RCu<sub>2</sub>Si<sub>2</sub> (R=Yb, Y) studied by soft x-ray angleresolved photoemission spectroscopy

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In recent years, electronic structure of YbCu\_Si<sub>2</sub>, which is a classical example of Yb-based valence fluctuation compounds, has been studied by various experimental methods such as a quantum oscillation measurement [1] and an angle-integrated photoemission spectroscopy (AIPES) [2]. However, these results are inconsistent each other. In fact, the band structure calculation which well describes the results of the quantum oscillation measurement [1] cannot explain the AIPES spectrum. In the present study, we have performed the soft x-ray angle-resolved photoemission spectroscopy (SX-ARPES) for YbCu\_Si<sub>1</sub> to understand the valence band structure and the Fermi surface (FS) in a unified way. The SX-ARPES measurements were recorded in BL23SU of SPring-8. We have observed that hybridized bands of Yb 4f electrons and conduction states cross the Fermi level at 20 K, which is well below the Kondo temperature of TK = 50 K, in line with the valencefluctuating state of YbCu\_Si<sub>2</sub>. We will discuss effects of Yb 4f electrons on its band structure by comparing the ARPES spectra with those of a non-4f reference compound YCu\_Si<sub>3</sub>.

[1] N. D. Dung et al., J. Phys. Soc. Jpn. 78, 084711 (2009). [2] M. Matsunami et al., Phys. Rev. B 78, 195118 (2008).

### **RD26**

### Anomalous increase of TC in UGa2 under pressure

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UGa<sub>2</sub> belongs to compounds quite thoroughly studied, although there is still no consensus about the character of the 5f states. In particular,  $5f_2$  or  $5f_3$  localized and band-like 5f states were considered, none of those being able to explain all its features [1]. Response of magnetic properties to external pressure is an important indicator of the situation of the 5f states. While band systems tend to a suppression of magnetic moments and their ordering, magnetism based on localized states has insensitive size of moments and their ordering temperatures can weakly increase. Previous experiments [2] indicated a quite rapid increase (approx. 3 K/GPa) of the Curie temperature from TC = 126 K up to p = 0.8 GPa. New resistivity experiments on UGa<sub>2</sub> single crystal, extending the pressure range up to 8 GPa, reveal a continuous linear increase, reaching 160 K without any sign of saturation. Such anomalously large increase can be only understood in the framework of a two-band model, in which the compression leads to a stronger hybridization of the 5f at other conduction-electron states.

[1] T. Honma et al., J. Phys. Soc. Japan 69 (2000) 2647. [2] A.V. Kolomiets et al., Physica B 259-261 (1999) 238.

# RE03

### NpCoGe, near quantum criticality?

performing memory function and gating operation.

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The orthorhombic isostructural series AnTGe (An = actinide ; T = transition metal) offers a unique opportunity to study the properties and trends of a system where ferromagnetism and superconductivity coexist at ambient pressure in URhGe (TC = 9.5K,  $\overline{TSC} = 0.26K$ ) [1] and UCoGe (TC = 3K, TSC = 0.7K) [2]. However, due to the difficulty to handle them, transuranium systems are much less documented. In an effort to bridge this gap, we have undertaken the study of the neptunium analogues and have previously reported that NpRhGe exhibits an antiferromagnetic ground state (TN = 21K,  $\mu$ Np = 1.14 $\mu$ B) [3]. In the present work, we have investigated NpCoGe by dc magnetization, ac susceptibility, specific heat, electrical resistivity, 237Np Mossbauer spectroscopy and LSDA calculations. We find that NpCoGe orders antiferromagnetically at TN=13K with an average ordered moment <mNp> = 0.8 µB. The weak antiferromagnetic interactions ( $\theta p = -5.5K$ ) are overcome by the application of a moderate magnetic field (B ~ 3T) that induces a metamagnetic phase. NpCoGe appears as a more delocalized antiferromagnet than NpRhGe, which is consistent with the trend observed in UTGe analogues. The proximity of NpCoGe to a quantum critical point and its implications will be discussed.

# **RE04**

# Pressure and magnetic-field induced non-Fermi-liquid behavior in $YbCo_2Zn_{20}$

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In heavy fermion systems, applying pressure and magnetic field controls the electronic configuration, and thus anomalous behaviors appear at around magnetic quantum critical point (QCP), such as an unconventional superconductivity and Non-Fermi-liquid behavior. Recently, YbCo<sub>2</sub>In<sub>30</sub> have attracted much attention because of their heavy fermion state and proximity to the pressure-induced QCP [1]. It is worth noting that the critical pressure of YbCo<sub>2</sub>Zn<sub>30</sub> is lower than that of other Yb based compounds. Therefore YbCo<sub>2</sub>Zn<sub>30</sub> has an advantage of investigating the behavior of Yb- based compounds near QCP. More interestingly, it is reported that a metamagnetic crossover appears at H ~ 0.6 T [2], however, the origin of this metamagnetic behavior remains unclear. In order to shed light on the nature of the pressure- and field-induced quantum critical behavior, we have measured the specific heat of YbCo<sub>2</sub>Zn<sub>30</sub> under pressure and magnetic field. We found that the low temperature-spressure-magnetic field phase diagram will be presented to sum up the physical properties.

[1] Y. Saiga et al., J. Phys. Soc. Jpn 77 (2008) 053710 [2] M. Ohya et al., J. Phys. Soc. Jpn 79 (2010) 083601

### RE05

# Relationship between single-particle excitation and spin excitation at the Mott transition

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The Mott transition in the one- and two-dimensional Hubbard models is investigated using the Bethe ansatz, dynamical density-matrix renormalization group method, cluster perturbation theory, and random-phase approximation. I will show that the dispersion relation of single-particle excitation in a metallic phase leads continuously to that of spin excitation of the Mott insulator.

### **RE06**

# Quantum criticality in Kondo quantum dot coupled to 2D topological insulator

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We investigate theoretically the quantum phase transition between the one-channel Kondo (1CK) and two-channel Kondo (2CK) fixed points in a quantum dot coupled to edge states of interacting 2D topological insulators (2DTI) with Luttinger parameter 0 < K < 1. For K < 1, the strong coupling 2CK fixed point was argued to be stable for infinitesimally weak tunnelings between dot and the 2DTI based on the simple scaling dimensional analysis. We re-examine the model beyond the scaling dimension analysis via a controlled 1-loop renormalization group (RG) approach on the effective Kondo model via bosonization and re-fermionization near the 2CK fixed point. We find for K<1 the 2CK fixed point can be unstable towards the 1CK fixed point and the system may undergo a quantum phase transition (QPT) between 1CK and 2CK fixed points. The QPT in our model comes as a result of the combined Kondo and the helical Luttinger physics in 2DTI, and it serves as the first example of the 1CK-2CK QPT accessible by a controlled theoretical approach. We extract quantum critical and crossover behaviors from various thermodynamical quantities near the quantum critical point.

# **RE07**

# Nature of insulator-metal-insulator transitions in the ionic Hubbard model

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It is known that the ionic Hubbard model possesses a metallic phase between band and Mott insulators in high dimensions. We study the phase diagram in the Bethe lattice for various strengths of the staggered lattice potential within the dynamical mean-field theory (DMFT) combined with the continuous-time quantum Monte Carlo method. We also investigate its zero-temperature limit by means of exact diagonalization as a DMFT solver. Observed at finite temperatures are the crossover between the metallic and band insulating phases and the first-order transition between the metallic and Mott insulating phases, ending up with the critical point. It is discussed how such transition behaviors evolve as the lattice potential is increased at low temperatures and whether the finite- and zero-temperature results are consistent with each other.

### RE08

# Variational cluster approach to the Hubbard model on the honeycomb lattice

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Physics of electrons on the honeycomb lattice, or graphene, is one of the topical issues in condensed matter physics. Especially, the influence of electron correlations on the magnetism and electronic states at the Fermi level, which is characterized by the linear dispersion relation and vanishing density of states in the non-interacting case, is of particular interest. In this study, we apply the variational cluster approach to the half-filled Hubbard model on the honeycomb lattice. In order to study the semimetalinsulator transition and paramagnetic-antiferromagnetic insulator transition, we introduce the bath hybridization strength and staggered magnetic field as variational parameters. We calculate the single-particle spectra and magnetic order parameter as a function of Coulomb interaction U, and discuss the implications of the results.

## RE09

## Onset of magnetic order in $U_2(Ni_{1-x}Fe_x)_2Sn$

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Single crystal of U\_2Ni\_SN and polycrystals U\_4(Ni\_k\_Fe\_b\_SN and their hydrides were prepared so as to investigate the character and stability of magnetism in U\_3Ni\_SN (antiferromagnet with T\_N = 26 K). We show that its magnetism is uniaxial, with 3 metamagnetic transitions detected in magnetic fields up to 60 T along the tetragonal c-axis. Such anisotropy manifests also in the paramagnetic state. Estimated difference of paramagnetic Curie temperatures, 175 K, quantifies the anisotropy. The AF order disappears rapidly with the Fe substitution. Magnetic susceptibility and specific heat of the sample with 20 % Fe reveals a great similarity to previously studied NFL heavy-fermion U\_2Co\_SN [1]. The C/T uptum at low temperatures, very resistant to magnetic fields, is suppressed with further increase of Fe concentration. Exposure to 100 bar H\_2 leads to the H-absorption forming U\_2(Ni,Fe)\_SnH\_2<sub>sc</sub>, U\_2Ni\_2SnH\_1<sub>s</sub> is AF (T<sub>N</sub> = 87 K), which is at least partly due to the lattice expansion (dV = 7-10%). In the middle of the Fe concentration range, hydrogenation induces a ferromagnetic order. A detailed study of the Fe-poor end, for which the approach to the NFL behaviour can be followed, is under way.

## **RE10**

### Magnetic phase diagram of UCoAl

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Heavy fermion UCoAl is known to be a unique system with metamagnetic transition from paramagnetic to ferromagnetic ground states in uranium-based compounds. This kind of metamagnetism and quantum criticality in strongly correlated electron systems have attracted much attention because of expectation of new quantum phases. Recently, we succeeded to grow a high quality single crystal of UCoAl by the Czochralski pulling method and performed the magnetic measurements. In this presentation, we will show the magnetic phase diagram of UCoAl in detail.

# **RE11**

**Renormalization-group exponents for competitions between interelectronic and phonon-mediated interactions in ladder systems** *Wen-min Huang<sup>1,\*</sup>, Yiwei Cai<sup>2</sup> and Hsiu-hau Lin<sup>2</sup>* 

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After developing a scaling ansatz for electron-electron interactions under renormalization group transformation[1], we show that, with the inclusion of phononmediated interactions the ansatz, characterized by the divergent logarithmic length and a set of renormalization-group exponents, also works rather well. The superconducting phases in a doped two-leg ladder are studied and classified by these renormalization-group exponents as demonstration. Finally, nontrivial constraints among the exponents are derived and explained.

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## **RE12**

### Crystal Growth and Magnetic Order of Ni-doped CePdAl

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CePdAl is a stoichiometric, antiferromagnetic compound (Neel temperature = 2.7 K) which can be tuned to a quantum critical point (QCP) by hydrostatic [T. Goto et al., J. Phys. Chem. of Solids 63, 1159 (2002)] or chemical pressure [Y. Isikawa et al., J. Phys. Sco. Jpn. 65, 117 (1996)]. The latter can be achieved by substituting Ni for Pd. Neutron-scattering experiments [A. Doenni et al., J. Phys.: Condens. Matter 8, 11213 (1996)] pointed towards a partial fustration of the Ce moments in CePdAl, making this system a promising candidate for investigating the influence of frustration on quantum criticality. We have grown large single crystals of Ni-doped CePdAl by the Czochralski method. The samples were characterized by x-ray powder diffraction, atomic absorption spectroscopy and x-ray Laue diffraction. Specific heat and inelastic neutron scattering data of CePdAl yield a Kondo temperature of approximately 5 K. Magnetization measurements display a strong magnetic anisotropy which is preserved in the Ni-doped compounds, where the magnetic order is suppressed. Neutron-diffraction experiments indicate that short-range correlations are present well above the Neel temperature. Below the Neel temperature, one third of the Ce moments display short-range order only, confirming the frustration in this system.

# **RE13**

## Transport properties of Ho<sub>1-X</sub>Lu<sub>X</sub>B<sub>12</sub> solid solutions

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Recent studies of Ho<sub>1-x</sub>Lu<sub>x</sub>B<sub>12</sub> solid solutions have shown [1] that the antiferromagnetic (AF) order in geometrically frustrated system of HoB<sub>12</sub> (x = 0, TN = 7.4 K) is linearly suppressed to zero temperature, i.e. TN  $\rightarrow 0$ , as lutetium concentration increases to x  $\rightarrow$  xC 0.9. In this contribution, we present original results of electrical resistivity measurements on Ho<sub>1-x</sub>Lu<sub>x</sub>B<sub>12</sub> single crystalline samples with x = 0, 0.2, 0.5, 0.7, 0.9, 1 in the temperature range 0.05 - 300 K and in magnetic fields (B) up to 8 T. Complex B vs TN phase diagrams were received from precise temperature and field dependences of resistivity with more AF phases for x  $\geq$  0.5 pointing to a possible quantum critical point at xC 0.9. The scattering of conduction electrons in the AF phase and in the paramagnetic phase as well as Hall effect results are analyzed and discussed for various concentrations x, when magnetic dilution increases with the increasing content of nonmagnetic Lu ions in the Ho<sub>1-x</sub>Lu<sub>x</sub>B<sub>12</sub> system.

[1] S. Gabani, I. Batko, M. Batkova, K. Flachbart, E. Gazo, M. Reiffers, N. Shitsevalova, K. Siemensmeyer and N. Sluchanko, 17th International symposium on boron, borides and related materials, Sept 11-17, 2011, Istanbul, Turkey.

### RE14

### Low temperature thermal and electrical transport properties of ZrZn<sub>2</sub> in high magnetic field

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# RE15

# Non fermi liquid behaviour in YFe<sub>2</sub>Ge<sub>2</sub>

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# RF01

# Magnetic orderings at the interface between Mott insulator and band insulator

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Recent experimental progress in crystal growth techniques enables us to investigate the complex oxide hetrostructures as a new frontier for studies of strongly correlated electron systems. In such systems, the main interest is put toward the metallic state around the interface, which is realized between two different insulating compounds, such as SrTiO<sub>2</sub>LaTiO<sub>3</sub> and SrTiO<sub>2</sub>LaAlO<sub>3</sub> [1]. Intrguingly, magnetism and superconductivity are also reported in such metallic interfaces [2]. Stimulated by these findings, we theoretically study the magnetic properties of the strongly correlated heterostructure composed a Mott insulator and band insulators. For this purpose, the single-band Hubbard model with long-range Coulomb interaction is investigated using the Hartree-Fock approximation [3]. We find the electronic and magnetic phases specific to the interface, including a canted antifierromagnetic state and a charge-ordered state with checkerboard charge distribution. We elucidate the origin of these phases is closely related to the spatial inhomogeneity which induces unconventional couplings between spin and charge degrees of freedom. Additionally, the applied magnetic fields also reveal other aspects of correlated heterostructures; a first order metamagnetic transition and field-induced checkerboard charge-ordered state, which are also driven by the non-uniform charge distribution.

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### **RF02**

## Strong Electron Correlation in Cu-doped CaO Nanocolloid

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Theoretical prediction of magnetism induced by defects or doping in non-metallic colloids has gained a renewed interests recently. In this work, we investigated the possible appearance of magnetism in Cu doped CaO nanocolloids activated by SPAN-80 in the framework of density functional theory (DFT). Despite of strong antiferromagnetic super-exchange interaction between Cu2+ ions, the local magnetic moment of Cu may appear due to colloidal agent attachment onto the surface of CaO nanocluster (surface modification). The ferromagnetism attributes to the degeneration of Cu 3d orbitals in CaO crystal fields, the aspects of electron correlation and quantum spin fluctuation.

PACS number: 42.30.R, 42.40.Ht, 42.30.Kq Keywords: DFT, band-gap, CaO, calcium-oxide, mBJ, PBE, functional

### RF03

# The book-keeping fermion analysis of the double exchange model with antiferromagnetic background

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We analyse the double exchange model with antiferromagnetic background by the book-keeping fermion method. First, we consider the motion of a single conduction electron. We assume that the localized spins form the Neel state and the added conduction electron forms only three types of states under the strong Hund-coupling assumption and we regard these states as the book-keeping fermion. Under these assumptions, we calculate the spectral functions which show good agreement with the results given by the exact diagonalizations for finite size clusters. We also study the relation between the strength of the Hund coupling and the bandwidth in the lowerenergy region. The results show that the bandwidth approaches a finite value as the Hund-coupling constant becomes close to infinity. The bandwidth even increases with increasing the Hund coupling, which is in strong contrast to the expectation from the semi-classical theory. Next, we consider the situation where there are finite numbers of spin flips occurring in the down-spin sublattice in the Neel ordered state. Considering the effects of the spin defects, we add a new book-keeping fermion and estimate the energy of the spin defects by a perturbation theory. This theory gives the phase diagram comparable with the numerical results.

# **RF04**

## Phonon Induced Thermodynamic Properties of $La_{1\mbox{-}x}Ca_xCoO_3$

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Perovskite type cobaltate,  $La_1$ -xCaxCoO<sub>3</sub> ( $0 \le x \le 0.3$ ) have been studied intensively because of their wide range of unique physical properties. The fact that rare-earth perovskite-type cobaltate are most suitable as a cathode material in solid oxide fuel cells (SOFC), makes the thermal behavior of these compounds highly important. As these compounds are predominantly ionic in nature hence the lattice contributions to the specific heat at constant volume (C<sub>v(lattice)</sub>) of pure and Ca doped LaCoO<sub>3</sub> has been studied and thereby thermal expansion is computed as function of temperature by means of Rigid Ion Model (RIM). We have systematically investigated the effect of phonons on thermal properties, Debye temperature  $(\Theta_{D})$ , molecular force constant (f), Reststrahlen frequency (v), cohesive energy ( $\varphi$ ), and gruneisen parameter ( $\gamma$ ) for La<sub>1</sub> $xCaxCoO_3$  ( $0 \le x \le 0.3$ ). Also the effect of phonons on the bulk modulus is studied using the atoms in molecules (AIM) theory for pure and ca doped LaCoO3. We have found that the computed properties reproduce well with the available experimental data, implying that RIM represents properly the perovskite cobaltate La<sub>1x</sub>Ca<sub>x</sub>CoO<sub>3</sub> (0  $\leq x \leq 0.3$ ). To our knowledge some of the properties for these complicated compounds are reported for the first time.

## **RF05**

# Dynamical antiferromagnetic phase transition after the quantum quench in the fermionic Hubbard model

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Spontaneous symmetry breaking is a universal concept in physics that stretches widely from condensed matter to cosmology. Recently it is getting feasible to study in detail how symmetry is dynamically broken in correlated electron systems using an ultrafast optical spectroscopy technique [1], which works as a key probe for electron correlation and fluctuations. Theoretically the phenomenon has been discussed with the Ginzburg-Landau (GL) theory, a macroscopic phenomenology. To capture the evolution of the electronic structure, however, one needs a microscopic description, which is not well understood. Here we take the fermionic Hubbard model, a typical microscopic model for correlated electrons, and study how the antiferromagnetic initial state after the quantum quench using the nonequilibrium dynamical mean-field theory. The system shows a multistep time evolution: (i) a fast thermalization within the paramagnetic phase, (ii) an exponential growth of the order parameter, and (iii) decay of collective excitations. We show how these features are reflected in the single-particle and optical-conductivity spectrum, and discuss the nature of antiferromagnetic fluctuations.

[1] R. Yusupov et al., Nat. Phys. 6, 681 (2010).

## **RF06**

# Itinerant magnetism in the hubbard model within the dynamical cluster approximation

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There have been a lot of efforts to understand the itinerant magnetism because ferromagnetism in many metallic magnets like iron, cobalt, and nickel is due to the itinerant magnetism. However, the mechanism and the stability of the itinerant magnetism are not clear yet because the strong quantum mechanical correlation in time and space is crucial in the itinerant magnetism. The Stoner model, which is the most popular model for the itinerant magnetism, ignores those correlation effects and so is inadequate [1]. In this paper, we studied the itinerant magnetism within the dynamical mean-field theory (DMFT) and its cluster extension, dynamical cluster approximation (DCA). These two methods are non-perturbative and fully include strong correlation effects [2]. Particularly, the DCA incorporates spatial correlation (nonlocal effects) systematically and gives numerically exact solution for the many-body problem [3]. We applied these two methods to study the ferromagnetic and antiferromagnetic instability of the Hubbard model in the 2-dimensional square lattice, the 3-dimensional cubic lattice, and the face-centered-cubic (fcc) lattice. We show that the correlation in time and space reduces the critical temperature remarkably and is essential to understand the itinerant magnetism.

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# **RF07**

# First Principles DFT+U Method for Strongly Correlating Electronic Structure Systems

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First principles density functional+U pseudopotential (DFT+U PP) method free from any empirical parameters has been developed and applied to the electronic structure calculations of strongly correlating electronic structure systems such as FeO and LaVO. The method can calculate the electronic structure of the systems from first principles, being contrary to the conventional DFT+U PP method, which uses empirical Ueff paramaters. The method primary calculates the Hubbard on-site interaction parameter Ueff for localized electrons, using an approximate constrained DFT technique and then uses the calculated parameter in the DFT+U PP electronic structure calculation. Examples of calculations resulted in Ueff of Fe 3d electrons of anti-ferromagnetic FeO to be 5.11 eV (the empirical value is 5eV) and Ueffs of La 4f and V 3d electrons of paramagnetic LaVO<sub>3</sub> to be 15.6 and 0.47 eV, respectively (the empirical values are 20.0 and 0 eV, respectively). The band gap structure of these materials was correctly described. Our method seems to be a promising tool for investigating strongly correlating electronic structure systems for spintronics applications, correctly describing the localized and itinerant nature of electrons of the systems from first principles.

## **RF08**

**Correlation effect in ferromagnetic 3d transition metals** Muneyuki Nishishita<sup>1</sup>, Sudhakar Pandey<sup>2</sup> and Dai Hirashima<sup>1</sup>\* <sup>1</sup> Nagoya University, Japan <sup>2</sup> APTPC, Korea



### **RF09**

### Mott transition in frustrated Hubbard model with spatial anisotropy: Cellular dynamical mean field study

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Geometrical frustration has attracted much interest in the field of strongly correlated electron systems. In the frustrated systems, a lot of intriguing phenomena such as the heavy fermion behavior in LiV<sub>2</sub>O<sub>4</sub>, spin liquid ground state in organic materials, etc. have been observed and they have stimulated theoretical investigations of frustrated electron systems. In particular, the reentrant behavior in the Mott transition in moderately frustrated systems, the heavy Fermi-liquid behavior around the Mott transition, etc. are new aspects of the Mott transition and very interesting. In our recent study on the kagome lattice system with spatial anisotropy [1] and also Ref. [2], it has been found that the anisotropy is strongly enhanced around the Mott transition, which stabilizes antiferromagnetic spin configurations. In this study, we apply the cellular dynamical mean field theory combined with a continuous-time quantum Monte Carlo solver [3] to the Hubbard model on the anisotropic kagome and triangular lattices. We systematically investigate the effects of the enhanced anisotropy around the Mott transition on the quasiparticle spectra and spin correlation functions.

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### RF10

# Spin-nematic and -singlet states in the Mott insulator phase of the S=1 two-dimensional Bose-Hubbard model

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Because of the recent development in experimental cold-atom physics, a quantum gas microscope technique has opened the door for detecting and manipulating single bosonic atoms at single site level in an optical lattice, just like scanning tunneling microscopy in solid-state physics. Quite recently, Endres et al [1], used this technique to track the superfluid-Mott insulator (SF-MI) transition of spinless bosonic-atoms on an optical lattice in more detail, and found that correlated pairs of a doubly populated site and an unpopulated site which representing the excitations in the MI, fundamentally determine the properties of the SF-MI transition, which is consistent with recent numerical studies [2,3]. Theoretically, spin-1 bosonic atoms on a two-dimensional lattice is also of prime importance because it is a bosonic analogue of the two-dimensional fermionic (S=1/2) Hubbard model, which has been exhaustively studied in the context of magnetism and high-temperature superconductivity. Thus in this work, we study the ground state properties and quantum phase transitions of S=1 two-dimensional Bose-Hubbard model, focusing on correlated pairs excitations based on a variational Monte Carlo approach. We discuss the spin-structures in the MI, and present a consistent description of the transition from the spin-mematic state to the spin-singlet state beyond mean-field theory.

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### RF11

# Dynamical instability in two-component bosonic systems in an optical lattice

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Because of the recent development in experimental technique, there has been an emerging interest in dynamical properties of Bose-Einstein condensates (BEC) on an optical lattice. In particular, quantum gas microscope technique has opened the door for real-time observation of single bosonic atoms at single site level in an optical lattice, which can act as scanning tunneling microscopy in solid state physics. Quite recently, Greiner and his co-workers used this technique to track the superfluid-Mott insulator (SF-MI) transition of bosonic-atoms in real-space and time [1,2]. Further studies for dynamical properties of BEC may be only a matter in time. Dynamical analysis of BEC in strongly correlated bosonic systems is also of interest in theoretical point of view. In particular, the dynamical instability of SF in multi-component bosonic systems is of prime importance because there would be entangled interplay among superfluidity, Mott physics, and possibly also Anderson localization. Thus, in this work, we study the dynamical instability of the SF phase of the two-component Bose-Hubbard model based on the dynamical Gutzwiller approximation. We mainly focus on the effects of the dispersive current of one-component on the stability of SF current of another component.

[1] W. S. Bakr et al., Science 329, 547 (2010). [2] M. Endres et al., Science 334, 200 (2011).

## **RF12**

### Mechanism for the high Neel temperature in SrTcO3

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The microscopic origin [1] of the high Neel temperature observed experimentally in SrTcO<sub>3</sub> [2] and CaTcO<sub>3</sub> [3] has been examined using a combination of ab-initio electronic structure calculations and mean-field solutions of a multi-band Hubbard model. The G-type anti-ferromagnetic state is found to be robust for a large region of parameter space, with large stabilization energies found, surprisingly, for small values of intra-atomic exchange interaction strength as well as large bandwidths. The microscopic origin of this is traced to specific aspects associated with the d3 configuration at the transition-metal site. Considering values of interaction strengths appropriate for SrTcO<sub>3</sub> and the corresponding 3d oxide SrMnO<sub>3</sub>, we find a ratio of 4:1 for the Neel temperatures as well as magnitudes consistent with experiment.

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# RF13

# Spin-spectral-weight distribution and energy range of the parent compound La<sub>2</sub>CuO<sub>4</sub>

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The spectral-weight distribution in recent neutron scattering experiments on the parent compound La<sub>2</sub>CuO<sub>4</sub> (LCO) [1], which are limited in energy range to about 450 meV, is studied in the framework of the Hubbard model on the square lattice [2]. We find that the higher-energy weight extends to about 566 meV and is located at and near the momentum [pi,pi]. Our results confirm that the U/t value suitable to LCO is in the range between U/t=6 and U/t=8. The continuum weight energy-integrated intensity vanishes or is extremely small at momentum [pi,0]. This behavior of the intensity is consistent with that of spin waves, which are damped at [pi,0]. Our study combines a number of theoretical and numerical approaches, including, in addition to standard treatments [3], a new spinon approach for the spin excitations [4,5] and density matrix renormalization group (DMRG) calculations for Hubbard cylinders [6].

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### RF14

### J<sub>1</sub>J<sub>2</sub> anti-ferromagnetic heisenberg model on bilayer honeycomb lattice Mojtaba Shoja Shabankah and Farhad Shahbazi *Physics, Isfahan University Of Technology, Iran*

Recent experiment on spin-3/2 bilayer honeycomb lattice antiferromagnet  $Bi_3 Mn_4 O_{12}$  (NO<sub>3</sub>) shows a spin liquid behavior down to very low temperatures. This behavior can be ascribed to the frustration effect due to competitions between first and second nearest neighbor antiferromagnet interaction. Motivated by the experiment, we study J<sub>1</sub>2-Antiferromagnet Heisenberg model, using Mean field theory and Monte Carlo simulation. This calculation shows highly degenerate ground state. We also calculate the effect of second nearest neighbor through z direction and show these neighbors also increase frustration in these systems. In addition calculation of the effect of third nearest neighbor shows that this interaction can lift the degenerate ground state to four discrete states. The next step of these calculations is to consider quantum effect of these interactions by Word line quantum Monte Carlo method.

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## RF15

### A hybrid exchange density functional study of La<sub>1-x</sub>Ca<sub>x</sub>MnO<sub>3</sub> Romi Kaur Korotana<sup>1</sup>, Leandro Liborio<sup>1</sup>, Giuseppe Mallia<sup>1</sup>, Zsolt Gercsi<sup>2</sup> and Nicholas Harrison<sup>1</sup>

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Manganites exhibit a rich variety of crytallographic, electronic and magnetic properties [1] manifested in ferromagnetic-, antiferromagnetic (collinear and canted)-, metallic-, insulating-, charge- and orbital-ordered states. An insight into the competition and coupling between various degrees of freedom in doped manganites is of interest for the optimisation of the magnetocaloric effect [2, 3] in such materials. In this work, hybrid-exchange density functional theory (DFT) calculations have been carried out to determine the effects of A-site doping on the structural, electronic and magnetic properties of perovskite-type manganites La<sub>1-x</sub> A<sub>x</sub> MnO<sub>3</sub>, where A=Ca and x=0, 0.25, 0.50, 0.75 and 1. A magnetic stability plot for La<sub>1-x</sub> Ca<sub>x</sub> MnO<sub>3</sub> is complied, in which the ferromagnetic (F), antiferromagnetic A-type (A-AF), G-type (G-AF) and C-type (C-AF) configurations are considered. At this level of theory, the structural, magnetic ground states and electronic states at each of the compositions studied are compared to available experimental data. This provides a basis for a first principles description of the magnetocaloric effect in La<sub>1-x</sub> Ca<sub>x</sub> MnO<sub>3</sub> systems.

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# **RF16**

# Magnetic properties of $GdFe_{\rm 11}Ti$ via first principal calculations.

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# RF17

Mott transition of ultracold Fermi-Fermi mixtures in optical lattices Takahiro Ooi\* and Seiichiro Suga

University of Hyogo, Japan



# **RG01**

### Anisotropy in a high Landau level due to effective electron-electron interactions Orion Ciftia

Prairie View A&M University, USA

Quantization of Hall resistivity in strongly correlated two-dimensional electronic systems at high magnetic fields generally indicates the stabilization of novel electronic quantum liquid phases of matter. This is the nature of the integer and fractional quantum Hall states that stabilize at integer and fractional odd-denominator (not always, though) filling factors of the Landau level. Away from certain filling factors that represent quantum Hall liquid states, different phases, some of them with unusually high magneto-transport anisotropy have been known to stabilize specially in high Landau levels. In this work, we try to understand the anisotropy of such quantum phases in terms of effective electron-electron interaction potential. To this effect, we implement a full projection of the original Coulomb interaction potential in the suitable Landau level. We find out that, in high Landau levels, thus for relatively weak magnetic fields, a semi-classical description of the interaction potential between electrons appear to be an adequate choice. The features of this semi-classical interaction potential in this limit suggest ways how the energetic balance between density waves and/or liquid crystalline phases might be sensitively affected.

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# **RG02**

# Bipolaron-Bipolaron interaction in many electron Holstein-Hubbard model

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Polarons and bipolarons have been long-term subjects of interest in many areas of condensed matter. The Holstein model describing the strong electron-phonon interaction and the Hubbard model describing the strong electron-electron interaction have played pivotal roles in various important physical systems, such as high Tc superconductors and the CMR manganites. In real systems, the polaron-polaron and polaron-bipolaron interactions might be of crucial importance. In order to investigate the physics of polaron and bipolaron systematically, we have studied the combined Holstein-Hubbard model treating both the electron-phonon and the electron-electron interactions as a function of the electron-phonon coupling strength and the electron-interaction U. We have calculated the ground state energy and wave functions of the system by applying the conjugate-gradient technique on a variationally constructed basis. The obtained correlation functions, the bipolaron binding energy, and the effective mass enables us to capture the intricate mechanisms at play and their consequences.

## RG03

# Partial disorder in an Ising-spin Kondo lattice model on a triangular lattice

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Partial disorder, that is coexistence of paramagnetic and magnetically ordered sites, is an intriguing example of exotic orders induced by geometrical frustration [1]. Experimentally, several metallic magnets with a triangular lattice structure, such as UNi<sub>4</sub>B and Ag<sub>2</sub>CrO<sub>2</sub>, were reported to present such unique order. On the other hand, theoretical attempts to explore the partial disorder in two-dimensional localized spin systems turned out to be unsuccessful so far; it is fragile against thermal fluctuations [2] and only exists as a quasi-long-range order at most [3]. To explore the possible realization of the true long-range partially disordered phase in two dimensions, here we theoretically investigate the effect of the coupling of localized moments to itinerant electrons in an Ising-spin Kondo lattice model on a triangular lattice. We study the low-temperature magnetic phases by a real-space Monte Carlo technique. As a result, we obtained a long-range order of the partially disordered state at finite temperatures in the vicinity of electron density n=1/3, where a two-sublattice stripe-type order and three-sublattice ferrimagnetic order compete with each other and an electronic phase separation takes place at a lower temperature. Further details on the phase diagram and the partially disordered phase will be given in the presentation.

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### RG04

# Basis reduction in the exact diagonalization method for the dynamical mean-field theory

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The exact diagonalization (ED) method for the dynamical mean-field theory is nearly the only method to calculate electron spectra involving core levels, such as core-level photoemission spectra, x-ray absorption/emission spectra, resonant inelastic x-ray scattering, etc., in a strongly correlated electron systems. In order to overcome the problem of the exponentially increasing Hilbert-space dimension by the number of sites, the number of configurational bases for the ED method for the single-impurity Anderson model is systematically reduced by their perturbative expansion starting from the lowest-energy configuration and consecutively allowing the hopping between an impurity level and a conduction band. Comparison of resulting Green's functions with exact ones in an imaginary frequency axis shows that bases up to the fourth order is enough to render reliable results. Since the reduced Hilbert-space dimension increases according to the power law, the available size of involved sites could be greatly enhanced to simulate core-level-related spectra with a large spin-orbital degree of freedom.

## RG05

# Electronic Structure of ternary stannides RRu<sub>4</sub>Sn<sub>6</sub> (R=Y, La, Pr, Ce, and Gd) compounds

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We report density functional calculations of the band structure and density of states of the RRu4Sn6 (M=Y, La, Ce, and Gd) compounds. Our investigation is carried out within the framework of the local density approximation, using a relativistic, full-potential band-structure method. For the Ce-146 compound we find that LSDA calculations predict a semiconducting ordered ground state whereas the other compounds ordered in metallic states, in accordance with available experimental measurements .

## RG06

Local correlation effects in Mn doped GaAs Igor Di Marco\*, Olle Eriksson and Patrik Thunstrom Physics and Astronomy - Materials Theory, Uppsala University, Sweden



### RG07

# First-principles study on noncollinear magnetism and effects of spin-orbit coupling in 5d pyrochlore oxide $Cd_2Os_2O_7$

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Recently, increasing attention has been focused on 5d pyrochlore oxides  $A_2B_2O_1$  (B=Ir, Os, etc.) in search for unconventional phenomena induced by strong spin-orbit coupling. Cd<sub>2</sub>Os<sub>2</sub>O<sub>7</sub> is a typical member of these compounds. Experimentally, it exhibits a metal-insulator transition at 227 K, below which a Neel order appears [1]. Despite many experimental attempts, the nature of the low-temperature phase remains to be clarified; e.g., the magnetic structure is unknown. Furthermore, several peculiar properties have been reported experimentally. (1) The transition temperature is the highest among pyrochlore oxides. (2) No clear charge gap is seen in the temperature dependence of the resistivity despite the high transition temperature. Motivated by these experiments, we investigate electronic structure of Cd<sub>2</sub>Os<sub>2</sub>O<sub>7</sub> [2] using the local spin density approximation + U method. We use a fully relativistic two-component first-principles computational code [3] based on the projector augmented-wave method. We show that the so-called all-infall-out magnetic order is the most stable in a wide range of U. By analyzing an effective spin model, we show that he local <111> easy-axis anisotropy removes geometrical flustration to stabilize this non-collinear magnetic order. We also compute the density of states and the electronic band structure and compare the results with experiments.

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# **RG08**

### Nonequilibrium states and I-V characteristics in one-dimensional band and Mott insulators attached to electrodes

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Nonequilibrium states induced by an applied bias voltage (V) and the corresponding current-voltage characteristics of one-dimensional models describing band and Mott insulators are investigated theoretically by using nonequilibrium Green's functions[1]. We attach the models to metallic electrodes whose effects are incorporated into the additional long-range interaction are calculated self-consistently within the Hartree approximation. For both models of band and Mott insulators with length L\_C, the bias voltage induces a breakdown of the insulating state, whose threshold shows a crossover depending on L\_C[2]. It is determined basically by the bias V\_{th}- $\Delta$  for L\_C smaller than the correlation length  $\xi=W/\Delta$  where W denotes the bandwidth and  $\Delta$  the energy gap. For systems with L\_C $\Box \xi$ , the threshold is governed by the electric field, V\_{th}/L\_C, which is consistent with a Landau-Zener-type breakdown, V\_{th}/L\_C $\Box \Delta^{-2}/V$ . We demonstrate that the spatial dependence of the scalar potential, where the breakdown occurs at V {th}

K. Yonemitsu, J. Phys. Soc. Jpn. 78 (2009) 054705.
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### **RG09**

# Enhancement of charge ordering by Zeeman effect in one-dimensional molecular conductors

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A recent experiment observes increase of resistance with increasing the external magnetic field in one-dimensional molecular conductor TPP[Co(Pc)(CN)\_2]\_2 [1]. The positive magnetoresistance is independent of the direction of the magnetic field, and then the origin of it is considered to be the Zeeman effect. Therefore, the experiment indicates that charge ordering (CO) is enhanced by the Zeeman effect. Stimulated by the experiment, we theoretically investigate effect of uniform magnetic field on CO in one-dimensional extended Hubbard model. In a previous study, the ground state with coexistence of CO and  $2k_F$ -SDW was investigated, but the amplitude of CO does not show noticeable enhancement under the uniform magnetic field [2]. In the present study, the parameter region where only CO appears is studied. We find the following results; (1) The amplitude of CO shows continuous enhancement under the external magnetic field. (2) Enhancement of CO is accompanied by both uniform and staggered (4 k\_F) component of magnetic moment. Our conclusion is consistent with experimental finding in TPP[Co(Pc)(CN) 2] 2.

[1] N. Hanasaki et al., J. Phys. Soc. Jpn. 75 (2006) 104743. [2] K. Kishigi and Y. Hasegawa, J. Phys. Soc. Jpn. 69 (2000) 2145.

### RG10

# Contribution of electron-lattice and spin-orbit coupling to the insulator-metal transition in VO<sub>2</sub> and $Sr_2IrO_4$

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Mott insulating state in the strongly correlated electron system is often realized in competition and/or cooperation with interactions other than the on-site Coulomb repulsion energy, such as the electron-lattice coupling and the spin-orbit coupling  $\Box$  One of representative examples is VO<sub>2</sub> where the electron correlation and the Peierls instability have been debated as a driving force to realize its metal-insulator transition near room temperature. Another example is Sr<sub>2</sub>IrO<sub>4</sub> which is now believed to have a spin-orbit-induced Mott insulator ground state. In this presentation, we discuss based on optical spectroscopy of thin films how the insulating ground states of each system evolve into the metallic states with a decrease of Peierls instability and the spin-orbit coupling strength which together can give an insight in to the complex phases in strongly correlated electron system [1,2].

[1] J. S. Lee, Y. Krockenberg, K. Takahashi, M. Kawasaki, and Y. Tokura, Phys. Rev. B 85, 035101 (2012), [2] J. S. Lee, K. Shibuya, M. Kawasaki, and Y. Tokura, submitted.

# RG11

# Multipole moments at Pr-site and electric field gradients at Sb-site in $PrOs_4Sb_{12}$

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 $PrOs_4Sb_{12}$  is well known as an unusual superconductor with Tc = 1.85 K [1]. The superconductivity might be related with the multipole moments of 4f<sup>2</sup> electrons of Pr, because the antiferro-quadrupole ordering is realized under applied magnetic fields [2]. So, it is important to study a role of the multipole moments of Pr to make clear the system of unusual superconductor PrOs<sub>4</sub>Sb<sub>12</sub>. The NQR frequency at Sb site shows temperature dependence corresponding to the crystalline electric field splitting below 10 K, associated with the Pr<sup>1</sup>+(4f<sup>2</sup>)-derived ground state [3]. This experimental result suggests that the change of multipole components affect the electronic field gradients (EFG) at Sb-site. The LDA + U calculations have described the Fermi surfaces of PrOs<sub>4</sub>Sb<sub>12</sub> well [4]. It has been shown that EFG at Sb-site can be calculated by using an FLAPW method and the calculated EFG reproduce well the experimental NQR frequency [5]. Moreover, the multipole moments of Pr can be classified from the density matrix, which is used in the LDA + U method. we have analyzed how each multipole component affects the EFG at Sb-site. Then we discuss the relationship between multipole moments at Pr-site and EFG at Sb-site.

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## RG12

# The effect of pairing fluctuations and disorder on the BCS-BEC crossover

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A fermi liquid with weak attractive interaction undergoes a BCS transition to a superconductor with reducing temperature. With increasing interaction strength the transition temperature displays non-monotonic behaviour as the system heads towards Bose condensation of preformed pairs. We use a new Monte Carlo tool, incorporating amplitude and phase fluctuations of the pairing field, to map out this BCS-BEC crossover in the attractive Hubbard model on large (40 X 40) two dimensional lattices. Havingestablished this benchmark we discuss the role of disorder in this system. This includes (i) the role of non-magnetic impurities where we track the superconductor to insulator transition with increasing disorder, and (ii) the emergence of 'gapless' superconductivity, the suppression of T\_c, and the induced correlation between scatterers themselves, in the case of magnetic impurities. Our method reproduces the standard results in all these cases but allows unprecedented spatial resolution, and access to the dynamical response. We compare the results to imaging experiments on disordered films.

## RG13

### Transport properties of ferromagnetic material with Anderson-Hubbard centers

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A competition between itinerant behavior and localization effects in strongly correlated electron system of periodically spaced Anderson-Hubbard centers introduced into narrow-band metal is studied. The configurational representation for localized subsystem is used and the effective Hamiltonian is obtained by canonical transformation. Besides the spin-spin interactions and strong on-site Coulomb interaction, the model takes into account the hybridization with conduction band which results in the indirect hopping and indirect exchange interactions. Lattice deformation under the external pressure is taken into account and equilibrium value of lattice strain is found, having effect on both electrical and magnetic properties of the system. Green functions for band and localized electrons are calculated. On this base, the energy spectrum and magnetization are investigated as function of model parameters, temperature and external pressure. For temperatures close to Curie temperature and partial filling of conduction band the peculiarities of electronic conductivity are studied. Concentration dependence of static conductivity exhibits different transport regimes. Our results show that in the considered model the effects of localization are enhanced even if external pressure promotes electrical conductivity

# **RG14**

First principles studies of organic charge transfer salts

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We study a number of BEDT-TTF based charge transfer salts in the framework of ab initio density functional theory and determine the parameters of the effective Hubbard Hamiltonian describing the low energy excitations[1]. For the spin liquid candidate kappa-(BEDT-TTF)2Cu<sub>2</sub>(CN)<sub>3</sub>, we investigate the question if the structures determined at various temperatures lead to different effective Hamiltonian representations[2]. For the Fabre salts (TMTTF)2X, we compare the properties of the underlying Hamiltonian for a series of different anions X.

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# RG17

Magneto-polaronic effects in molecular transistors as the consequence of quantum uncertainty of the displacement of vibrating quantum dot. Glib A(alexandrovich) Skorobagatko<sup>1\*</sup>, Sergev I, Kulinich<sup>2</sup>, Ilva V, Krive<sup>3</sup> and Robert

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The influence of magneto-polaronic effects induced by the quantum uncertainty of the displacement of quantum dot on electron transport through a single-level vibrating quantum dot subjected to a transverse (to the current flow) magnetic field is considered . It is shown that the effects are most pronounced in the regime of sequential electron tunneling, where a polaronic blockade of the current at low temperatures and an anomalous temperature dependence of the magnetoconductance are predicted.

I. G. A. Skorobagatko, S. I. Kulinich, I. V. Krive, R. I. Shekhter, and M. Jonson, Low Temp. Phys. 37, 1032 (2011).

# RG15

Iuly 12 (Thu)

### Dynamics of strongly correlated Fermi systems: The efffects of pairexcitations and exchange

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Understanding the emergent properties of many body systems from the underlying microscopic interactions is the major goal of condensed matter physics. Recent progress, in the understanding of the dynamics of strongly correlated systems, has been made by approximating the excited wave function by a superposition of correlated particle-hole and 2-particle-2-hole excitations. The amplitudes, describing these excitations, are obtained by minimizing the action integral. A surprising result of this approach is the confirmation/prediction of the reemergence of collective mode below the PHB in two dimensional 3He for density fluctuations. By including exchange the agreement of the calculated spectrum with the experiment was further improved for 3He. The next step is to apply this tool to spin-density fluctuations, which are of inherent exchange character. Due to the proximity of the ferromagnetic phase transition the fluctuations are at very low energies for long wavelenth. This experimental finding is also present in a first implementation of our theory. The presented approach is also applicable to other correlated fermi systems e.g. electrons and dipols in two dimensions. We predict this transition of the collective mode seen in 3He also for two dimensional electrons. The capability in describing these systems is investigated.

Bohm, H. M., Holler, R., Krotscheck, E. & Panholzer, M. Dynamic many-body theory: Dynamics of Strongly Correlated Fermi Fluids. Phys. Rev. B 22, 224505 (2010).

### **RG16**

# The induced effects of the Dzyaloshinskii-Moriya interaction on the thermal entanglement

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The pairwise thermal entanglement in one-dimensional (1D) spin-1/2 XX model with added Dzyaloshinskii-Moriya interaction and transverse magnetic field is studied through fermionization technique. We have showed that, for fields less than quantum critical point, the thermal entanglement decreases by increasing the temperature and vanishes at a field-independent critical temperature. On the other hand for fields more than quantum critical value of the magnetic field (Fig. (1)(b)), NN spins are not entangled up to the first critical temperature  $T_{c_{1}}$ , the thermal entanglement regain and takes a maximum value and then decrease and reaches to zero at the second critical temperature  $T_{c_{2}}$ ).

### RG18

### **Defect states and electron correlations in multi-orbital Mott insulators** Adolfo Avella<sup>1\*</sup>, Peter Horsch<sup>2</sup> and Andrzei Oles<sup>3</sup>

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We address the role played by defects in doped Mott insulators with active orbital degrees of freedom. We observe that defects are characterized by rather complex and rich physics which is well captured by a degenerate Hubbard model extended by several terms that describe crystal-field splittings, the orbital-lattice coupling, as well as local terms generated by defects such as the Coulomb potential terms that act both on a doped hole and orbitals of undoped sites (orbital polarization). We show that the multiplet structure of excited states generated in such systems by strong electron interactions is well described in an optimized unrestricted Hartree-Fock approximation, taking into account the usual symmetry breaking by the onset of magnetic and orbital order. More importantly, we show that defect states are responsible for new features that arise within the Mott-Hubbard gap and in the multiplet spectrum at high energy. These states involve active orbital flavors at atoms b eing nearest neighbors of the defect state, which are modified by the local defect-orbital Coulomb interactions. The present study suggests a new mechanism for the Coulomb gap realized in the presence of defect states and investigates the dependence of the orbitals on the orbital polarization

### **RH01**

# Magnetism of the noncentrosymmetric compound ${\rm CeNiC_2}\mbox{-} pressure effects$

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The system RNiC<sub>2</sub> (R=Rare-earth element) has a simple but a characteristical noncentrosymmetric crystal structure, and shows various novel physical properties such as charge density waves (CDW) and complex magnetic structures. In the system the compound of R=La shows unconventional superconductivity, and has been intensively studied very recently. For CeNiC2 a few reports indicated that the incommensurate antiferromagnetic state (ICAF) appears at around 20 K and this changes to the commensurate antiferromagnetic structure (CAF) at about 10 K, then a ferromagnetic state (F) takes place at around 2 K [1, 2]. However, since there are some clear discrepancies in these studies the magnetism of this compound has been reinvestigated in detail using a carefully prepared polycrystalline sample. It is found that there are definite magnetic moments in the paramagnetic state. With these moments the complex antiferromagnetic states undergo, which replaces to a ferromagnetic state at low temperatures. To investigate these magnetic states further, their pressure effects have been studied up to 2.5 GPa. The results show that the IC and CAF states at the higher temperatures is fairly stable, whose transition temperatures increase at a rate of about 2 K/GPa. On the other hand the F state is changed little

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# **RH02**

Atomic scale disorder driven bicritical region in Sm<sub>0.5</sub>(Ca<sub>1-x</sub>Sr<sub>x</sub>)<sub>0.5</sub>MnO<sub>3</sub> Saurav Giri<sup>1\*</sup>, Sk. Sabyasach<sup>1</sup>, S. Majumdar<sup>1</sup>, S. Das<sup>2</sup> and V. S. Amaral<sup>3</sup> <sup>1</sup> Solid State Physics, Indian Association for the Cultivation of Science, India <sup>2</sup> Department of Physics and CICECO, University of Aveiro, Portugal <sup>3</sup> Department of Physics and CICECO, University of Aveiro, India

On the verge of collapse of charge and orbital ordering a bicritical region occurs in hole doped manganites which displays rich electronic phase diagrams in R0.5(Ca,Sr)0.5MnO<sub>3</sub> [1]. Here, we demonstrate a fascinating consequence of A-site disorder in intermediate members of a new series Sm0.5(Ca1-xSrx)0.5MnO3. The Rietveld refinement of X-ray diffraction studies reveal maximum structural distortions at  $x = \frac{1}{2}$ . This is indicated by the highest value of  $\varepsilon$  [= v2b/(a+c), a, b, c being lattice constants] and most deviation of Mn-O-Mn bond-angle from 180° compared to end compounds. An unusual behaviour in magnetic field-induced ultrasharp transition to ferromagnetic metallic (FMM) state and collapse of charge ordering is observed both in magnetization and magnetoresistance curves for intermediate compositions at much lower field compared to end compounds. When magnetic field is applied above a critical field, FMM state appears and it retains high-field state, although magnetic field decreased to zero. With decreasing field sweep rate, the transition shifts to higher field indicating a meta-magnetic transition. The atomic-scale local inhomogeneity/distortion arising from the difference in ionic radii and/or the random Coulomb potential from ion mixture is supposed to be critical behind such striking phenomena. [1] Y. Takura, Rep. Prog. Phys. 69, 797 (2006)

### RH03

### 11B-NMR study on Shastry Sutherland system TbB<sub>4</sub>

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In TbB<sub>4</sub>, the arrangement of Tb ions within the c-plane is topologically equivalent to the Shastry Sutherland (S-S) lattice with geometrical frustration [1]. A well-known example of S-S lattice is SrCu<sub>2</sub>(BO<sub>3</sub>)<sub>2</sub> which shows the characteristic magnetization plateaus caused by the frustration. However, in TbB<sub>4</sub>, due to the quadrupole interaction and RKKY interaction between classical spins, the appearance of the novel magnetism quite different from SrCu<sub>2</sub>(BO<sub>3</sub>)<sub>2</sub> is expected. The magnetization of TbB<sub>4</sub> for H // c-axis shows feld-induced multi-step jump while the magnetization in the c-plane on which we focus attention shows only one large jump at Hc = 15.9 and 12 T for H//[100] and H//[110] respectively [2]. In order to investigate the magnetic structure in high magnetic fields for H//(100], we performed 11B-NMR experiments up to 17.5T for the single crystalline TbB<sub>4</sub>. 1B-NMR spectra observed at low magnetic field changes drastically at Hc = 15.9 T, where the magnetization jump occurs. This result indicates the change of the magnetic structure at the field-induced magnetic plase transition. In order to determine the magnetic structure, we have compared the observed NMR spectra with the calculation based on the classical dipole-dipole interaction.

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### RH04

### Magnetotransport Property of the Hole-doped Delafossite CuCr<sub>0.97</sub>Mg<sub>0.03</sub>O<sub>2</sub> with a Spin-3/2 Antiferromagnetic Triangular Sublattice

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Magnetotransport property of a hole-doped delafossite  $CuCr_{0.97}Mg_{0.05}O_2$  with a spin-3/2 antiferromagnetic (AF) triangular sublattice was investigated by a pulsed high magnetic field up to about 56 Tesla. [1] A dramatic change of magnetoresistance (MR) was observed with a variation of temperature (T), which is due to a development of an AF correlation with a decrease in T. A negative MR is observed and the ratio is well proportional to a square of magnetization above a Curie-Weiss temperature ( $\theta$ ), which seems to suggest that the magnetic state above  $\theta$  is well explained by a Kondolattice model [2] with a weak ferromagnetic Hund's coupling between itinerant holes and localized spins. However, below  $\theta$ , the negative MR gradually deviates from the scaling relation with a decrease in T and is much enhanced around an AF transition temperature (TN). Furthermore, a component of positive MR appears around TN. Such a MR behavior is due to a critical enhancement of spin fluctuation around TN, which may cause a nontrivial promotion of a 120° AF state by the Mg substitution [3,4].

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 T. Okuda, Y. Beppu, Y. Fujii, T. Onoe, N. Terada, S. Miyasaka, Phys. Rev. B 77, 134423 (2008).

# RH05

# Lifshitz transition with interactions in high magnetic fields: application to CeIn<sub>3</sub>

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### Novel colossal magnetoresistance in NaCr<sub>2</sub>O<sub>4</sub>

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We recently discovered new chromium oxide, NaCr<sup>2</sup>O<sup>4</sup>, to have calcium ferrite type structure, in which spin frustration is expected. The compound shows canted-antiferromagentic transition at TN = 125 K and Arrhenius-type temperature dependence of electrical resistivity. The resistivity below TN is greatly suppressed by magnetic field and it remains to be a finite value at 0 K under H = 9 T. Thus, the magnetoresistance ratio, {rho(H)-rho(0)}/rho(0), becomes closer to -100% at low temperatures. The similar behavior has been observed in some manganites, such as  $P_{0.7}Ca_{0.3}MnO_3$  and  $Nd_{0.5}Sr_{0.5}MnO_3$ , for which the competition between charge-ordered insulating phase and ferromagnetic metallic phase causes colossal magnetoresistance (CMR) through the first-order transition. Unlike them, however, NaCr2O4 shows no thermal or field hysteresis in its physical properties, and there is no ferromagnetic metallic phase compositionally close to NaCr<sub>2</sub>O<sub>4</sub>. Therefore, CMR of the compound is considered to be caused by a new mechanism. The magnetism of Ca1-xNaxCr<sub>2</sub>O<sub>4</sub> strongly suggests unusual electronic state of Cr<sup>4+</sup> ions. We will present a possible mechanism of CMR in NaCr<sub>2</sub>O<sub>4</sub>, in which spin frustration and the electronic state of Cr<sup>4+</sup> ions play a key role.

# **RH07**

### Low field study of Hall effect in GdB<sub>6</sub>

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GdB<sub>6</sub> plays the important role for understanding of the mechanisms responsible for the antiferromagnetic (AF) state formation in the rare earth hexaborides. However the magnetic structure of GdB<sub>6</sub> which is characterized by two successive AF transitions (AF(I) ordering below T\_N-15.5K, and AF(II) phase at T^\*-4.7K) is still the subject of discussion [1]. In current work we present the comprehensive study of Hall effect on the high quality single crystals of GdB<sub>6</sub> in the wide temperature interval (2-120K) at low magnetic fields (H≤1T). In paramagnetic state Hall coefficient R\_H is practically temperature independent and value of R\_H--4-10^{4}-4)cm^3/C detected with the high accuracy is in accord with the results of light hexaborides RB<sub>6</sub> (R-La, Ce, Pr, Nd) published previously [2],[3]. The temperature decrease below T\_N induces the drastic enhancement of negative Hall coefficient (AR\_H/R\_H~ 64%). The magnetic ordering at T^\* (AF(I)-AF(II) transition) is accompanied by the appearance of anomalous Hall effect which is detected in GdB<sub>6</sub> for the first time. The complex analysis of charge transport parameters suggests the possible effect of 5d-states spin density polarization in AF and PM states of GdB<sub>6</sub>.

[1] M.Amara, S.E.Luca et al., Phys. Rev. B v. 72, p. 064447 (2005), [2] Y. Onuki, A.Umezawa et al., Phys. Rev. B v. 40, p. 11195 (1989), [3] M.A.Anisimov, A.V.Bogach, V.V.Glushkov et al., arxiv cond-mat:1006.0124v1

## **RH08**

# Theoretical study of structures and mechanical properties of M-type hexagonal ferrites $BaFe_{12}O_{19}$

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M-type hexagonal ferrite,  $BaFe_{12}O_{19}$ , has extensively studied in a wide range of applications such as microwave devices and permanent magnets due to high Curie temperature, large saturation magnetization and good chemical stability. In this paper, we have investigated the structure and elastic modulus of  $BaFe_{12}O_{19}$  using atomistic simulations. An interatomic potential based on the shell model is used for the structure and properties of  $BaFe_{12}O_{19}$ . The shell model has been extensively used to a wide range of oxide materials since it is a simple and useful potential with the Coulomb interaction incorporated with the short-ranged term for ionic characteristics, though it is incapable of describing the elemental components. The calculated results are in agreement with experiments and previous theoretical studies. To our knowledge, although many theoretical investigations have so far played an important role in understanding the properties of magnetic materials, there have been few theoretical studies using the interatomic potential for the structure and properties of  $BaFe_{12}O_{19}$ . The interatomic potential employed has been shown to be capable of describing the Ba-based hexagonal ferrites.

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### RH09

### Nonlinear susceptibility of gadolinium near curie temperature Takashi Shirane\* and Shohei Sakurai

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The ac linear and nonlinear susceptibilities of an elemental ferromagnetic gadolinium have been studied near the Curie temperature, to clarify the feature of the magnetic phase transition with special emphasis on the behavior below the Curie temperature. The temperature dependence of linear susceptibility is in good agreement with that of an experimental result reported in ref. [1]. The first observation for the critical behavior of the nonlinear susceptibility is gadolinium has been made at various field amplitudes. The behavior of the nonlinear susceptibilities are explained on the basis of a mean field theory taken into account the domain structure[2]. The detailed discussion will be shown the full paper in comparison with magnetization curves at various temperature below the Curie temperature.

[1] J. M. D. Coey, K. Gallangher, and V. Skumryev, J. Appl. Phys. 87 (2000) 7028. [2] T. Shirane, Statistical and Condensed Matter Physics: Over the Horizon, ed. S. Fujita, T. Obata, and A. Suzuki (New York: Nova Science), p. 189-197, (2007).

### **RH10**

# Electron correlation and dynamical Jahn-Teller effect in orbitally degenerate system

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Quantum spin liquid state is one of the attractive subjects in correlated electron systems. Recently, a possibility of the quantum spin liquid state is experimentally reported in a layered copper oxide  $Ba_3CuSb_2O_9$  where the orbital degree of freedom in a Cu ion and the dynamical Jahn-Teller effect (DJTE) are suggested to play some roles on an origin of the spin liquid state. Motivated from these recent experimental results, we study the DJTE in a spin-orbital coupled system. In particular, we focus on competitive or cooperative phenomena between the superexchange interaction and the DJTE. A superexchange interaction part is derived from the d-p model and the DJTE part for the low lying vibronic states is described by the orbital pseudo-spin and the lattice vibration. We analyze the model which combines the two interactions on a honeycomb lattice by using the Bethe approximation and the exact diagonalization method. We find that the magnetic order is unstable in a wide parameter region and a spin-singlet dimer state associated with an orbital order is realized. With increasing the DJTE, furthermore, the orbital order is strongly suppressed and a resonance state of the spin-orbital dimers appears.

H. D. Zhou, et.al., Phys. Rev. Lett., 106 (2011) 147204.

## RH11

# High-field NMR study on the charge in stability in quantum spin system $Cu_3Mo_5O_9$

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The quasi-one-dimensional quantum spin system  $Cu_3Mo_2O_9$  with the two spin-degrees of freedom, antiferromagnetic chain and dimer-like sites clinging around the chain undergoes the long-range antiferromagnetic order at 7.9 K, associated with a weak ferromagnetic moment. The system shows a small magnetization jump of 0.0033 µB at H\* = 8.2 T, where the anomalous charge instability appears [1], suggesting an existence of tight coupling between spin and charge. In order to inverstigate the microscopic spin and charge states in the high magnetic field, we have utilized Cu-NMR technique, which separately probes the spin and charge in the chain site  $Cu_1$ , the dimer like sites  $Cu_2$  and  $Cu_3$  through hyperfine coupling and electric quadruple interaction respectively. In this presentation, we show by detailed analyses of spectra taken under a wide range of filed up to 17.5 T that the ordered antiferromagnetic moment is dominated by  $Cu_1$ , while the dimer-like sites only feel a uniform hyperfine field but staggered one, and that the change inhomogenization takes place in  $Cu_2$ , and  $Cu_3$  in the vicinity of the field of magnetization jump H\*.

1 T. Hamasaki et al., Phys. Rev. B77 134419 (2007).

### **RH12**

# Anisotropic spin excitations in spin-Peierls CuGeO<sub>3</sub>

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Even though CuGeO<sub>3</sub> is a well-known 1-dimensional quantum spin system that shows magnetic transition into a dimerized state with a spontaneous lattice distortion, it is not clear whether it is an archetypal material of a spin-Peierls system. In fact, neutron scattering experiment has still not succeeded in assignment of the soft phonon mode accompanying the magnetic transition. Actually, due to an antiferromagnetic coupling between the neighboring spin chains, the dimerization in the spin chain is close to the critical state [1][2][3], and it should not have a strong effect on a phonon. In this time, we performed inelastic neutron scattering experiment on CuGeO<sub>3</sub> with the Fermi chopper spectrometer 4SEASONS at MLF/J-PARC in Japan. In order to further understand the inter-chain coupling, we measured intensity-maps in a 4-dimensional ( $q_{c0}$ ) space including the inter-chain direction due to rotation of the aligned 17.2 g CuGeO<sub>3</sub> single crystals in horizontal plane. In this experiment, we have clarified an anisotropy of the magnetic excitations below the spin-Peierls stransition temperature. This result suggests the inter-chain coupling correlates with the dimerization in the spin-Peierls state on CuGeO<sub>3</sub>. We will present the detail of the magnetic excitations together with new other features recently discussed in several theoretical models.

[1] M. Nishi et al., PRB, 50,6508 (1994). [2] J. Riera et al., PRB, 51,16098 (1995). [3] M. C. Martin et al., PRB, 56, 3173 (1997).

### RH13

## Neutron inelastic scattering on spin-peierls system TiOBr

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Newly proposed spin-Peierls system TiOX (X: Cl, Br) has been revealed showing exotic structural and magnetic properties such as a successive phase transition, onedimensional (1D) nature associated with orbital ordering of Ti ions and super-lattice structure being related to the Peierls instability. It is pointed out that resulting only from an arrangement of Ti dxy orbital, the formation of 1D spin chains and the spin-Peierls transition will be realized. Recently, it has been demonstrated that TiOBr also exhibits two successive phase transitions similar to TiOC<sub>1</sub> at Tc1=27K and Tc<sub>2</sub>=47K. Here we carried out inelastic neutron experiments in order to see the evidence of spin-Peierls transition. The inelastic spectrum with a large amount of poly crystalline sample of TiOBr shows the localized signal in the vicinity of the magnetic zone center Q=0.9A-1. Observed spin gap like signal lies at energy of  $\Delta$ E=10meV. The gap energy in TiOBr is expected much higher from measured thermodynamic properties and by analogy with TiOCI. Constant Q cuts of the observed S(Q,E) map show some Q-dependent structure in its intensity indicating the signal is sample or intend. The Q structure quite reveals the intensity is well explained by the powder averaged dynamical structure.

# RH14

# Physical properties of the novel triangular-lattice silver oxides Ag<sub>2</sub>MO<sub>2</sub> (M = Co, Ga, Rh)

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Triangular antiferromagnets have been extensively studied by many scientists in terms of geometrical frustration. According to current theories, triangular antiferromagnets show various exotic magnetic properties, e.g. chirality order, Kosterlitz-Thouless transition, a partially disordered state, etc. A series of layered Ag<sub>2</sub>MO<sub>2</sub> compounds, where M is a transition metal, are possible model compounds of triangular antiferromagnets. The crystal structure of Ag<sub>2</sub>MO<sub>2</sub> consists of alternate stacks of an  $M^{5}O_2$  layer and a (Ag<sub>2</sub>)+ layer. The former includes a triangular lattice of the  $M^{5+}$  ions, while the latter provides itinerant electrons from the quarter-filled Ag 5s band. The triangular lattices are well separated from each other by the non-magnetic double Ag layer. A path does not exist for superexchange interaction between the  $M^{5+}$  spins on the adjacent triangular lattices. In addition, no cation disorder is expected between the M and Ag ions. Ag<sub>2</sub>MO<sub>2</sub> is, therefore, considered a suitable system to study the magnetic properties of triangular antiferromagnets. We recently succeeded in synthesizing Ag<sub>2</sub>MO<sub>2</sub>, M = Co, Ga, Rh, for the first time using a high-pressure technique. The details of physical properties of Ag<sub>2</sub>MO<sub>2</sub> will be reported.

### RH15

# Electrical and thermal transport properties of the polycrystalline $(Cr_{86}Ru_{14})_{1x}V_x$ alloy system

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Thermopower (S) and electrical resistivity ( $\rho$ ) measurements on the (Cr<sub>86</sub>Ru<sub>14</sub>)1-xVx alloy system with  $0 \le x \le 0.144$  reveal unique behaviour[1,2,3]. The onset of a commensurate (C) itinerant-electron antiferromagnetic spin-density-wave (SDW) structure is marked by anomalies in  $\rho$ (T) and S(T), associated with the Neel temperature. Focus is on the thermopower as it is more sensitive to the changes in the electronic structure and scattering mechanisms that are of importance at the SDW transition[2]. For this series S(T) shows a broad valley for  $x \le 0.032$ . The valley gradually becomes more shallow with increasing x and is eventually replaced by a surplus of positive thermopower for x = 0.074, for which parallelism is seen between  $\rho$ (T) and S(T) angly correlates and decreases with increasing x; attributed to the addition of V to the system that destabilizes the CSDW structure. The effect of the addition of V will be explained in terms of a model that relates the behaviour of S(T) to the energy dependence of the decrease in the scattering rate of electrons by phonons and an increase in the resistivity on decreasing temperature due to a decrease in the Fermi surface area[3].

 [1] L. Reddy et al., J. Appl. Phys. 103 (2008) 07C903-1. [2] E. Fawcett et al., Rev. Mod. Phys. 66 (1994) 25. [3] A. L. Trego et al., Phys. Rev. 166 (1968) 495.

### RH16

Magnetic properties of the layered triangular-lattice antiferromagnets CsM(MoO<sub>4</sub>)<sub>2</sub> (M=V, Fe) Masahiko Isobe and Yutaka Ueda *ISSP, Univ. of Tokyo, Japan* 

Two-dimensional triangular lattice antiferromagnets are of great interest concerning the ordering in frustrated spin system. We have prepared powder samples of the layered triangular-lattice antiferromagnets  $CsFe(MoO_4)2$  and  $CsV(MoO_4)2$  and have investigated the magnetic properties. In  $CsFe(MoO_4)2$ , a structural phase transition at around 365K from P3m1 to P3 was detected. Also we observed the magnetic transition at 4.4 K for  $CsFe(MoO_4)2$  and 26 K for  $CsV(MoO_4)2$ . The magnetic properties of  $CsFe(MoO_4)2$  are similar to that of the multiferroic RbFe(MoO\_4)2.

# RH17

### Kitaev-Heisenberg magnetism in honeycomb iridates A2IrO3(A=Li,Na)

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5d transition-metal oxides can have unique properties because of the strong spin-orbit coupling and relatively weak electron correlations. Layered perovskite iridate Sr<sub>2</sub>IrO<sub>4</sub>, for example, was revealed to be Jeff=1/2 Mott-insulator where spin and orbital degrees of freedom are entangled with including complex phase i[1]. The magnetism of such Jeff=1/2 Mott-insulator might be different from 3d transition-metal Mott-insulator.Indeed in iridates, the superexchange coupling between Ir-O-Ir was proposed to be anisotropic ferromagnetic when IrO<sub>6</sub> octahedra form edge-sharing network, due to the complex phase[2]. As a consequence, the spin-liquid state based on the Kitaev-model was theoretically predicted when such edge-shared IrO6 octahedra compose a honeycom network. Such unique lattice topology can be found in layered iridatesA,IrO<sub>2</sub>(A=Li,Na). Na,IrO<sub>2</sub> is a frustrated antiferromagnet with Tn=17K, and the  $\theta$ w=-125K. The strong antiferromagnetic interaction is likely due to direct interaction between Ir 5d electrons. Meanwhile, Li, IrO, exhibited less frustrated magnetism with Tn=12K and 0w=-12K, even with the similar structure. We argue that this seemingly non-frustrated magnetism of Li,IrO, likely originates from a coexistence of direct antiferromagnetic and ferromagnetic superexchange interaction, exemplifying the Kitaev-Heisenberg model[3]. The difference might come from the local structure, the honeycomb of Na<sub>2</sub>IrO<sub>3</sub> is distorted but that of Li,IrO, is closer to ideal honeycomb

[1] B.J.Kim et al., Science 33,1329(2009) [2]G.Jackeli and G.Khaliullin, Phys.Rev.Lett. 102,017205(2009) [3] J.Chaloupka et al., Phys.Rev.Lett. 105,027204 (2010)

## **RH18**

# $Effects \ of \ the \ annealing \ conditions \ on \ the \ magneto-transport properties \ of \ La_{0,7}Sr_{0,3}Mn_{1+4}O_3-manganese \ oxide \ composites$

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The effects of the annealing conditions on the magneto-transport properties of the pure La<sub>0.7</sub>Sr<sub>0.3</sub>MnO<sub>3</sub>(LSMO) and LSMO-manganese oxide composites were carefully investigated in this study. While the ferromagnetic-paramagnetic transition temperatures of all polycrystalline samples were insignificantly varied in the range of 359–368 K, low field magnetoresistance (LFMR) properties and dMR/dH values were remarkably improved for the composite samples annealed at relatively high temperature. For instance, with increasing the annealing temperature from 1300 to 1450°C, the LFMR ratios of 95 mol% LSMO - 5 mol% Mn2O3 composite sample were increased from 0.8 % to 1.1 % at 300 K in 0.5 kOe, suggesting that a magnetic disorder at the LSMO grain boundary, known as a key factor suppressing effective spin-dependent scattering, could be greatly improved by high temperature annealing process. Detailed annealing effects on the magneto-transport properties of samples will be presented for a discussion. This work was supported by the Korea Research Foundation Grant funded by the Korean Government (MOEHRD, Basic Research Promotion Fund) (KRF-0417-20100021)

## RH19

# Temperature dependence of spin lattice relaxation time of proton NMR in mixed antiferromagnets $A_{1x}B_xCl2-2H_2O$

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In the mixed antiferromagnet systems N<sub>i1-x</sub>Co<sub>x</sub>C<sub>12</sub>\_2H<sub>2</sub>O and Ni1-xMnxC<sub>12</sub>\_2H<sub>2</sub>O, the Co substitution increases a little the transition temperature and the Mn substitution decreases the transition temperature rapidly. In Ni-Co system, the concentration dependence of the phase transition temperature is well explained by simple molecular field theory. But, in Ni-Mn system, molecular field theory cannot explain it sufficiently. Thus Mn spins in NiC12\_2H2O crystal show the peculiar behavior. Then, we have measured the temperature dependence of the spin-lattice relaxation time of proton NMR, in order to understand the specificity of the Mn spins. We have prepared the mixed antiferromagnets A1-xBxC12-2H2O (A, B=Co, Mn, Ni). In Co-M (M=Ni, Mn) system, in comparison with the Ni substitution, the Mn substitution induces a significant impact on the relaxation time. In Mn-M (M=Ni, Co) system, the relaxation time is affected by a small amount of substitution. In Mn-Co system, the reentrant spin-glass phase appears at low temperature. The strange behavior of the relaxation time at low temperature is believed to reflect that. A similar trend is seen at the low temperatures in the case of Ni substitution, too. We suppose that this may be attributed to a kind of the instability of Mn spins

## **RH20**

Electronic Structure and magneto-optical properties of Co<sub>2</sub>MnX alloys where X = Ge, Sn and Pb: a first-principles investigation in LDA+U approach

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We report our ab initio calculational results on the electronic structures and magnetooptical (MO) properties of the ferromagnetic Co2MnX alloys where X= Ge, Sn and Pb. Employing the +U corrections for the 3d valence bands of transition metal atoms in addition to the conventional local density approximation (LDA), we investigate the correlation effect on the MO spectra in polar geometry as well as the detailed electronic structures using the all-electron FLAPW [1] method. Results show that the correlation effect results in a blue-shift of the peak positions and the large enhancement of the MO spectra compared to the LDA results. We find that our results can be attributed to the increased t2g-eg splitting of spin minority d-bands of both Co and Mn atoms, which indicates the suppression of diagonal elements of optical conductivity at energy region of 1~2 eV where the interband transitions are forbidden.[2]

[1] E. Wimmer, K. Krakauer, M. Weinert, and A. J. Freeman, Phys. Rev. B 24, 864 (1981). [2] Support by Basic Science Research Program (NRF-2010-0005387) and Priority Research Centers (2010-0029617) through the National Research Foundation of Korea.

# **RH21**

### 2D Heisenberg antiferromagnetism in spin-orbit Mott insulator Sr<sub>2</sub>IrO<sub>4</sub> Akiyo Matsumoto1 Tomohiro Takayama1 and Hide Takagi2\*

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In 5d iridates, the interplay between spin-orbit coupling and modest electron correlation produces novel electronic phases.  $Sr_2IrO_4$  was recently found to be a Jeff =1/2 Mott insulator where spin and orbital degrees of freedom are entangled[1]. The magnetic coupling between such Jeff = 1/2 state, which could be critically different from that in 3d-based Mott insulators, is of particular interest. For example, the magnetic coupling between Ir was theoretically proposed to be isotropic Heisenberg-type when IrO6 octahedra form corner-sharing network[2]. We focused on Sr<sub>2</sub>IrO<sub>4</sub> to reveal the magnetism of such spin-orbit Mott insulator. Sr<sub>2</sub>IrO<sub>4</sub> crystallizes in a K2NiF4-type perovskite structure, and magnetically orders at 240K. We investigated the detailed magnetic susceptibilities on the single crystals. The ground state was found to be antiferromagnetism, and the inplane susceptibility steeply increases just above the transition temperature likely due to Dzialloshinski-Moriya interaction. Meanwhile, at higher temperatures, the in- and out-ofplane susceptibilities were almost isotropic. By analyzing the susceptibilities, we found that they can be well described by the 2D Heisenberg antiferromagnetism. This indicates that even in the limit of strong spin-orbit coupling, Sr<sub>2</sub>IrO<sub>4</sub> displays the isotropic Heisenberg magnetism. This situation contrast with the 3d oxides, where spin-orbit coupling is responsible for magnetic anisotropy.

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### RH22

### Magnetoresistance and magnetic properties of oxygen deficient (Sr,Y) (Fe,Co)O<sub>3</sub> perovskites

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The magnetic and magnetoresistive properties of oxygen deficient (Sr,Y)(Co,Fe) O3 perovskites have been investigated in relation with the composition. The valence state of Fe has been probed by 57Fe Mossbauer spectrometry. Sr1-xYxFe0.8Co0.2O3-d and Sr<sub>0.8</sub>Y<sub>0.2</sub>Co<sub>1-x</sub>FexO<sub>3-d</sub> series were more particularly investigated. It appears that the progressive replacement of Sr2+ by Y3+ and the Co/Fe substitution lead to a change of the valence state of Fe and to a drastic change of the magnetic and magnetoresistive properties, in relation with the variation of the Fe4+/Fe3+ ratio and the change in the oxygen content. It is shown that the  $Sr_{0.8}Y_{0.2}Co_{0.5}Fe_{0.5}O_{1.4}$  reveals the highest magnetoresistance effect, in relation with the fact that the magnetic field induced ferromagnetic alignement of the Fe/Co magnetic moments allows more metallic pathaways to be created. This work was supported by the French Agence Nationale de la Recherche (project ANR-08-BLAN-0005-01).

### **RH23**

### High field element selective magnetometry in erbium iron garnet Cornelius Strohm\*, Thomas Roth, Peter J. E. M. Van Der Linden, Olivier Mathon and Sakura Pascarelli

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Ferrimagnets play an important role for applications as well as from a fundamental point of view. The rare earth iron garnets for example were at the heart of the demonstration of the local molecular field theory [1] and later used for bubble memories, microwave devices and optival isolators. The ferrimagnetic superexchange between the rare earth and iron cations being mainly determined by bond angles and distances in the oxygen polyhedra, results in molecular fields of the order of tens of Tesla [2, 3]. We have performed element selective magnetometry on samples of Erbium Iron Garnet, using energy dispersive x-ray magnetic circular dichroism at the Fe K-edge and the Erbium L-edges. Pulsed fields of 30 T were generated by a high duty cycle miniature coil [4] and the signal was acquired using a multiframe detection scheme exploiting the entire pulse duration of 1 ms [5]. Here we discuss mainly the results obtained at the Fe K-edge. At 65 K, below the compensation point, two successive field induced phase transitions and the reversal of the net magnetization of the Fe sublattices in the intermediate canted phase were observed, while the Fe signal remains unchanged at 100 K, above the compensation point.

[1] L. Neel Science 174, 985 (1971), [2] A. E. Clark and E. Callen J. Appl. Phys. 39, 5972, (1968), [3] K. Nakao, T. Goto and N. Miura J. of Magn. Magn. Mat. 54, 1363, (1986). [4] P. J. E. M. van der Linden et al. Rev. Sci. Instrum. 79, 075104, (2008). [5] C. Strohm et al. J. of Synchrotron Rad. 18, 224, (2011).

### **RH24**

### THz and infrared excitation spectrum below the Jahn-Teller transition in Sr<sub>2</sub>Cr<sub>2</sub>O<sub>6</sub>

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We report on optical excitations observed recently in Sr<sub>3</sub>Cr<sub>2</sub>O<sub>8</sub> by THz and infrared spectroscopy. Low-energy excitations below 3 THz are detected by THz time domain spectroscopy including the spin singlet to triplet excitations.[1] These excitations can be divided into two different classes according to the temperaturedependent properties. One is emergent right below the Jahn-Teller transition temperature, which is determined by specific heat measurement to occur at 285 K. [1,2] The other appears only below 100 K, where the fluctuations are sufficiently suppressed, consistent with the temperature dependence of low-energy Raman modes. [3] Infrared transmission measurements reveal a broad crystal-field excitation, which can be associated with an electronic transition from E to  $T_2$  orbital states.

[1] Zhe Wang et al., Phys. Rev. B 83, 201102 (2011) [2] D. L. Quintero-Castro et al., Phys. Rev. B 81, 014415 (2010) [3] D. Wulferding et al., Phys. Rev. B 84, 064419 (2011)

### **RH25**

### High-field multi-frequency ESR in the S=2 Heisenberg antiferromagnetic chain compound MnCl<sub>3</sub>(bpy)

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We report the results of high-field multi-frequency ESR experiments at 1.3 K on a powder sample of MnCl<sub>3</sub> (bpy) (bpy=2, 2'-bipyridine). This compound is one of the rare examples of the spin 2 quasi-one-dimensional Heisenberg antiferromagnets, and the magnetic properties on tiny single crystal samples were reported previously [1]. In our previous paper on this compound [2], we attained good agreement between experiment and calculation on the temperature dependence of magnetic susceptibility and the high-field magnetization process at 1.3 K. In this work, we have observed some high-frequency resonance modes with a zero-field gap of about 800 GHz that may correspond to the excitation modes at q=0 in the energy dispersion higher than the triplet excitation mode of about 50 GHz at q=pi expected from the evaluated exchange constant and the g-value.

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# **RI01**

### Distribution of magnetization in the random ising models.

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We investigate the two- and three-dimensional random Ising models by the Monte Carlo method. The integrated distribution functions of site-magnetization are calculated above and below the lower Griffith temperature. It is found that these functions have distinct character above or below this temperature for the two-dimensinal random Ising model. However, for the three-dimensional random Ising model, this distinction cannot be observed on contrast to the two-dimensional random Ising model. This behavior may be explained by the percolation theory.

### **RI02**

## Berezinskii-Kosterlitz-Thouless trasition in two-dimensional p-state clock model

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In two dimensions, the Berezinskii-Kosterlitz-Thouless (BKT) transition occurs in models with the U(1) symmetry. The helicity modulus, which measures the rigidity against a global twist across the system, is useful for a detection of the transition because of its distinctive feature, the universal jump at the transition point. Meanwhile, in the p-state clock model which has the Z p symmetry, it is believed that the BKT transition occurs when p is greater than 4 since the study by Jose et al. in 1977[1]. However, a recent numerical study by Baek et al. in 2010[2] detects an absence of the universal jump in the helicity modulus for p=5. This contradicts the Jose's result. In this study, we show that the helicity modulus as defined by Baek (HMB) remains finite for any p with the high-temperature expansion. Therefore, the universal jump in the HMB cannot be observed whether the transition is the BKT or not. This casts a question on the physical meaning of the HMB. We will present a systematic study of the helicity modulus to clarify the universality class of the transition for p=5.

[1] J. V. Jose, L. P. Kadanoff, S. Kirkpatrick and D. R. Nelson, Phys. Rev. B 16, 1217 (1977) [2] S. K. Baek and P. Minnhagen, Phys. Rev. E 82, 031102 (2010)

### **RI03**

### Influence of interplanar coupling on the entropy and specific heat of the bilayer ferromagnet

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The Ising-Ising and Ising-Heisenberg bilayer with a simple quadratic lattice in each plane is studied using the pair approximation (PA) method. The foundation of formalism used has been developed in [1], where the phase diagrams for the anisotropic bi- and multilayer system have been obtained. In this paper the method (leading to a complete self-consistent thermodynamics) is adopted for the entropy and magnetic specific heat calculations. An interesting case is demonstrated when the exchange integrals in both planes are of unequal strength, and the interplanar coupling (Ising or Heisenberg) is relatively weak. In such a case, two peaks in the temperature dependence of the specific heat can be observed. The isothermal change of the entropy caused by application of the external magnetic field has been additionally calculated for a wide temperature range. The entropy change in the field also shows a double peak structure, similarly to the specific heat. Such a phenomenon may be of potential interest for investigations of the magnetocaloric effect. The results of the PA method have been compared with those obtained for two Ising-type planes by means of exact diagonalization for finite clusters

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## **RI04**

### Nontrivial ferrimagnetism of the heisenberg model on the union jack strin lattice

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In low-dimensional quantum spin systems with frustration nontrivial magnetic phenomena occur due to strong quantum fluctuation. One of the phenomena is a non-Lieb-Mattis (NLM) ferrimagnetism[1][2], which is different from the well-known Lieb-Mattis (LM) ferrimagnetism[3][4]. In contrast to the LM ferrimagnetism. the spontaneous magnetization of the NLM ferrimagnetism changes gradually with respect to the strength of frustration. Incommensurate modulation with longdistance periodicity in local magnetizations is also a characteristic behavior of the NLM ferrimagnetism. However, there are only a few cases revealing the NLM ferrimagnetism, which prevents us with understanding the mechanism of the occurrence of the NLM ferrimagnetism. In this study, we examine the ground-state properties of the S=1/2 Heisenberg model on the Union Jack strip lattice. We find the existence of the NLM ferromagnetism by the numerical diagonalization and density matrix renormalization group methods. We also discuss the relationship between the NLM ferromagnetic state of this strip lattice and the intermediate canted state of the S=1/2 Heisenberg model on the two-dimensional Union Jack lattice[5].

[1] S. Yoshikawa and S. Miyashita: J. Phys. Soc. Jpn. 74 (2005) Suppl. 71. [2] K. Hida: J. Phys.: Condens. Matter 19 (2007) 145225. [3] W. Marshall: Proc. Roy. Soc. A 232 (1955) 48. [4] E. Lieb and D. Mattis: J. Math. Phys. 3 (1962) 749. [5] R. F. Bishop, P. H. Y. Li, D. J. J. Farnell, and C. E. Campbell: Phys. Rev. B 82 (2010) 024416.

### R105

### Control of quantum critical points in bond disordered spin ladder materials

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Spin ladders are model systems of low dimensional quantum spins. There is now considerable interest in materials which realise these ideal systems with solvable Hamiltonians and rich phase diagrams. An archetypal spin ladder material in the lowenergy regime is  $(C_5H_{12}N)2CuBr_4$  whose phase diagram is completely accessible in laboratory magnetic fields. The rung (Jr) and leg (Jl) exchange are 12.8 K and 3.4 K, respectively, and two field-controlled quantum critical points (OCPs) exist at 6.9 T and 14.5 T. (C<sub>e</sub>H<sub>12</sub>N)2CuC<sub>14</sub> is a related spin ladder material. The superexchange between quantum Cu<sup>2+</sup> spins is controlled by the halide atoms and is lowered to Jr=3.4 K and Jl=1.3 K. The QCPs appear at the lower fields of 1.7 T and 4.3 T. A common structure and rich chemical flexibility in the superexchange pathways allows for partial halide substitution to create the bond disordered spin ladder materials (C5H12N)2CuBr4(1-x)Cl4x, while retaining ideal spin ladder geometry. We present recent neutron scattering work and measurements of the magnetic and thermodynamic properties on the substitution series, and demonstrate explicitly substitution-controlled modification of exchange interactions as well as the effects of bond disorder on the spin ladder physics and OCPs.

# **RI06**

### Triangular spin tubes with bond randomness

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We studied X-ray diffraction and magnetic susceptibility experiments on nonequilateral triangular spin tubes composed of  $\alpha$ -KCr<sub>1</sub>-xFexF<sub>4</sub> (x=0-0.13), which consists of one-dimensional Heisenberg antiferromagnets coexisting with geometrically frustrated spin systems with bond randomness. For x=0, an antiferromagnetic longrange order occurred at TN1 = 2.5(1) K and another phase transition occurred at TN<sub>2</sub> = 4.0(1) K, because superexchange interactions through the three Cr-F-Cr paths in each non-equilateral triangle lost their equilibrium at low temperatures [1]. As a result, the values of spin-flop transition field drastically decreases with increasing x. This is probably due to the close correlation between the spin structure in the antiferromagnetic ordered state and the crystal structure as theoretically predicted by Nenert and Palstra, i.e., a magnetoelectric linear effect in which a magnetic field in an antiferromagnetic ordered state induces electrical polarization. Thus we carefully verified the crystal structure for x = 0.0.13.

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## **RI07**

### A new approach to the characterization of aging, rejuvenation, and memory effects in magnetic systems

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We introduce a new approach to study non-equilibrium dynamics in magnetic systems. We found that simple ferromagnets such as pure terbium and nickel aluminide foils actually resembles typical spin glasses such as Cu<sub>27</sub>Mn<sub>2</sub> and Cd<sub>42</sub>Mn<sub>4</sub>. Te, in behavior of out-of-phase component of ac susceptibility and magnetic viscosity, when they are thermally perturbed after isothermal aging. Briefly, it is not enough to conventionally measure such dissipative properties for characterizing the non-equilibrium dynamics of these systems. In contrast, in the spin glasses, the relaxations of the uniform magnetization were accelerated at a constant magnetic field, when the spin glasses were cooled/heated after the isothermal aging. Surprisingly, such relaxations were reversed when they were subsequently reheated/recooled to the original temperature, despite the persistence of the magnetic field. Because the uniform magnetization mirrors the global evolution of the spin configuration, these reversions indicate that the spin configurations are spontaneously restored to the original [1,2]. On the other hand, such reversions were never observed in the ferromagnets. This fact indicates that ferromagnetic correlations are unaffected by the thermal perturbations, unlike rejuvenation in spin glasses. Thus, the observations of the uniform magnetization would be an additional key to clarifying the features of non-equilibrium dynamics in magnetic systems.

[1] H. Mamiya and S. Nimori, New J. Phys. 12 (2010) 083007. [2] H. Mamiya and S. Nimori, J. Appl. Phys, in press.

### **RI08**

### The linear soliton generated by $Z_2$ vortex in quantum antiferromagnet Pawel Rusek

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The order parameter of noncollinear Heisenberg antiferromagnet is a triad of orthonormal vectors, i.e. the space of available states is isomorphic to SO3 group. The topology of SO<sub>2</sub> manifold allows the existence Z<sub>2</sub> vortices in a system. We show that the anisotropy essentially changes the structure of Z<sub>2</sub> vortex in 2D spin system. On the scale larger than the scale la defined by weak anisotropy energy the Z2 vortex is a termination of a linear soliton with the topological charge  $\pi 1$  (R) = Z<sub>2</sub>. In the soft core of this soliton the rotation angle of triad  $\theta$  rapidly changes from the anisotropy locked value  $\theta a$  to  $\theta = \pi$ , jumps on the line to  $\theta = -\pi$ , and then decreases to the value  $\theta a$  on the la distance from the line. The linear soliton can terminate on the other Z<sub>2</sub> vortex or it should go out to surface sample. The rapid change of  $\theta$  in the soft core of linear soliton causes the change of plaquette chirality there- the chirality change the sign on the line initiating the creation of the chirality domain.

### **RI09**

# Exact results of a mixed spin-1/2 and spin-1 Ising model with bilinear and three-site four-spin interactions on decorated planar lattices

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A mixed spin-1/2 and spin-1 Ising model on decorated planar lattices is investigated using a generalized decoration-iteration transformation technique [1]. In addition to the standard bilinear nearest-neighbor interactions, we also consider three-site four-spin interactions and the crystal field effects. In this contribution, exact expressions for the partition function, critical temperatures and other thermodynamic quantities are obtained and analyzed. Performing numerical calculations, we have found that in addition to the standard ferromagnetic and paramagnetic phases there exists in the system also a peculiar partly ordered phase, which possesses non-zero entropy at T = 0 and appears only for zero bilinear exchange interaction. In this phase, the decorating atoms occupy spin states ±1 and 0 with equal probability at all temperatures, while the nodal spins-1/2 are perfectly aligned along easy axis and generate non-zero magnetization. Our calculations primarily clarify the influence of many-body interactions in localized spin systems, which is apparently very different from the standard pair interactions due to the different symmetries of relevant terms in Hamiltonian. However, the presented results may be of wider interest since the origin of the three-site four-spin interaction is related to the magnetoelastic effect [2].

This work has been supported under grant VEGA No. 1/0234/12. [1] M. Jascur, Physica A 252 (1998) 217. [2] T. Iwashita, N. Uryu, J. Phys. C: Solid State Phys. 17 (1984) 855.

# **RI10**

### Anomalous spin diffusion on the two-dimensional percolating network in Rb2Mn06Mg04F4

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Inelastic neutron scattering experiments were performed on a two-dimensional dilute antiferromagnet, Rb2Mn06Mg04F4, with the magnetic concentration being close to the percolation concentration,  $c_p = 0.593$ , well above the Neel temperature, and with a high energy-resolution of  $\Delta E = 17.5 \ \mu eV$ . The energy spectrum obtained from the observed dynamical structure factor S(q,E) integrated over wave number, q, throughout almost the entire Brillouin zone, showed a power law dependence,  $S(E) \sim E^{\{-x\}}$ . The diffusion on a percolating network is anomalous and it has been predicted that the mean square displacement of a random walker is described by  $\langle R^2(t) \rangle \sim t^{2/(2+\theta)}$ with an exponent  $\theta$  as a function of the time t. The exponent in S(E) is described as  $x = 1 - Df/(2+\theta)$  with Df being the fractal dimension of the medium. The observed exponent, x, was in good agreement with a theoretical value.

## **RI11**

Quantum phase transitions in 1/3 plateau of the quantum spin tube Kouichi Okunishi1\*, Masahiro Sato2, Toru Sakai3, Kivomi Okamoto4 and Chigaku Itoi<sup>5</sup>

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We study quantum phase transitions in the 1/3 plateau state of the three-leg spin-1/2 tube, where the low-energy effective chirality degree of freedom plays an essential role. Using density matrix renormalization group and the effective chirality model, we find that the chirality liquid, a novel spin imbalance phase and the vector-spinchirality ordered phase emerge without closing the plateau spin gap, as the leg coupling increases. We also report the role of S\_3-symmetry of the spin tube behind these quantum phase transitions in detail.

K. Okunishi, M. Sato, T. Sakai, K. Okamoto and C. Itoi, arXiv:1109.0063

### **RI12**

### Cs<sub>2</sub>CoC<sub>14</sub> - an effective XY-spin-½ compound in transverse magnetic fields

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Cs<sub>2</sub>CoCl<sub>4</sub> is a model system for studying the magnetism of one-dimensional spin chains with XY anisotropy. It contains CoCl4 tetrahedra, which form one-dimensional chains along the crystallographic b-axis. The orbital groundstate of Co2+ is split by the crystal field into doublets and an easy-plane anisotropy of the magnetization is established. The ground-state doublet is separated from the first excited doublet state by approximately 15 K, such that at temperatures between 0.3 and 4 K the compound is well described by the one-dimensional XXZ model. We compare experimental data of specific heat and thermal expansion to numerical calculations of the XXZ model for transverse magnetic fields; i.e., fields applied within the XY planes. Decreasing temperature below 0.3 K, magnetic order arises at field-dependent temperatures TC(H). Measuring thermodynamic properties with a magnetic field applied along the crystallographic b-axis, we observe a series of magnetic transitions within the ordered state. This work was supported by the DFG through SFB 608.

# **RI13**

### Critical phenomena at the antiferromagnetic phase transition of Azurite Cu<sub>2</sub>(CO<sub>2</sub>)2(OH)2

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The natural mineral Azurite, Cu3(CO3)2(OH)2, has been considered as a model substance for the 1D distorted diamond chain, where unconventional magnetic properties result from the interplay of strong quantum fluctuations, low dimensionality and frustrating interactions. The system exhibits 3D long-range antiferromagnetic order at TN = 1.88 K. Our measurements of the sound velocity, ultrasonic attenuation and thermal expansion show pronounced signatures at TN indicating a strong coupling to the lattice degrees of freedom. We present a detailed investigation of the critical behavior of the sound velocity v22 and ultrasonic attenuation  $\alpha$  in the vicinity of TN. Our analyses reveal a power-law behavior of the critical contributions v22cr  $\sim$  t- $\zeta$  and  $\alpha$ cr  $\sim$  t-n over a decade of reduced temperature t = (T - TN)/TN both above and below TN. While the value of  $\eta = 1.21$ , obtained above TN, is close to the theoretical prediction for an anisotropy 3D Heisenberg antiferromagnet, the critical exponent derived from the sound velocity of  $\zeta = 0.056$  being identical for temperatures above and below TN, is at odds with this universality class. However, this value is consistent with the critical exponent observed in recent neutron scattering measurements [1].

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### **RI14**

### Dynamical properties of supersolid states in spin systems Yuta Murakami\*, Takahi Oka and Hideo Aoki

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It has been recognized that spin systems can be effectively thought of as boson models as in magnon BEC[1]. Specifically, a possibility to realize supersolid(SS) states (in the boson language) in spin systems is attracting recent attention, where a candidate is a frustrated spin-1/2 dimer model[2]. Here we study how the SS state accommodates collective excitations with the spin-wave theory. [3] We find that the excitation spectrum changes discontinuously as we go from a superfluid phase to an SS phase (that accompanies an SDW; CDW in boson language). Quantitatively, an analytic relation between the velocity of the Goldstone mode and the superfluid density is given. We also calculate the spin-spin correlation function,  $Szz(q, \omega)$ , which translates to the dynamical structure factor in cold-atom systems. We find the intensity of the mode as estimated from the peaks in  $Szz(q, \omega)$  sharply blows up toward the boundary of SS, which is accompanied by a drastic change in the character of the roton mode. We also mention the possibility to enlarge the SS phase region. These results obtained in the effective boson model should also be directly applicable to cold bosonic atoms.

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### **RI15**

### Magnetic phase transition of antiferromagnetic Cs<sub>3</sub>V<sub>2</sub>Cl<sub>9</sub>

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Cs<sub>3</sub>V<sub>2</sub>Cl<sub>9</sub> belongs to the family of hexagonal compounds A<sub>3</sub>M<sub>2</sub>X<sub>9</sub> (A=Cs, Rb, M=transition metal ions, X=Cl, Br) in which  $M_2X_9^{3}$  dimers are oriented along crystallographic c-axis. Former work11 on polycrystalline sample of Cs3V2Cl9 reported that intradimer interaction is ferromagnetic although ground state is nonmagnetic because of relatively large zero-field-splitting. Dimers of Cs<sub>3</sub>V<sub>2</sub>Cl<sub>9</sub> form triangular lattice in their c-plane so that the geometrical spin frustration effect is expected to occur if interdimer interactions are antiferromagnetic. We measured magnetic susceptibility, specific heat, high field magnetization and 133Cs-NMR using single crystal to study magnetic properties of Cs<sub>3</sub>V<sub>2</sub>Cl<sub>0</sub>. The specific heat shows two anomalies at around 15 and 4 K suggesting successive phase transitions. Results will be discussed in terms of the spin frustration.

1) B. Leuenberger et al. Inorg. Chem. 25, (1986) 2930.

# **RI16**

### Novel Field-Induced Quantum Phase Transition of the Kagome-Lattice Antiferromagnet

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The kagome-lattice antiferromagnet is investigated using the numerical exact diagonalization up to 42-spin clusters. From our investigation of the magnetization process, we found a novel field-induced quantum phase transition, called the "magnetization ramp", at 1/3 of the saturation magnetization [1]. The finite-size scaling analysis of the critical exponents[2] indicated a significant difference of the magnetization ramp from the magnetization plateau observed in the triangular-lattice antiferromagnet. The critical magnetization behavior at lower-field side of the magnetization ramp is same as that of the one-dimensional spin liquid. In order to clarify a difference between the magnetization ramp and plateau, we study a generalized anisotropic triangular-lattice model including the regular-triangular- and the kagome-lattice antiferromagnets in the parameter space. It revealed a quantum phase transition between the triangular- and kagome-lattices at 1/3 of the saturation magnetization. It means that the magnetization ramp of the kagome-lattice antiferromagnet is an essentially different phenomenon from the magnetization plateau of the triangular-lattice one. The relation to the recent magnetization measurements on the volborthite and vesignieite[3] is also considered. In addition we discuss about the spin gap issue of the kagome-lattice antiferromagnet [4].

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## **RI17**

### Antiferromagnetic phase transition of K-Rb alloy nanoclusters incorporated in sodalite

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Aluminosilicate sodalite possesses a bcc arrangement of nanospaces (beta-cages). Alkali-metal nanocluster (A4)3+ can be formed in each beta-cage, where an s-electron is shared by four alkali cations. (Na4)3+ clusters in sodalite are known to show antiferromagnetism below the Neel temperature of 48 K [1]. By substituting heavier alkali atoms for Na ones, the Neel temperature systematically increases, namely, 72 K in (K4)3+ and 80 K in (K3Rb)3+ clusters [2]. In the present work, we have succeeded in preparing K-Rb alloy clusters with a higher Rb concentration, (K1.5Rb2.5)3+. Magnetic susceptibility obeys Curie-Weiss law and shows a peak at ~ 100 K. ESR spectrum is suddenly broaden below ~ 90 K to be an asymmetric shape which can be explained by powder pattern of antiferromagnetic resonance. These results clearly indicate that the present sample exhibits an antiferromagnetic phase transition at approximately 90 K which is higher than that in (K3Rb)3+. Substituting Rb for K makes the size of s-electron wave function of the cluster larger because of the larger ion size and the lower ionization energy of Rb. This reduces the U/t in the Mott-Hubbard insulating state resulting in an increase in the kinetic exchange coupling between the adjacent clusters.

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## **RI18**

### A novel scaling method for critical phenomena studies: finite size effects

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We present a novel approach for critical phenomena studies. By quantifying scaling plot data overlap, we introduce a 'goodness of scaling' parameter, allowing an iterative search for the 'best' critical exponent values for a given set of magnetization data. The main advantage of this method is going beyond the usual power-law fits of the critical isotherm or low field susceptibility curve, by considering the complete equation of state and all data within the critical region. We consider Monte-Carlo simulated data of the 3D Ising model, and evaluate finitesize effects using our scaling method. Our results are in good agreement with the bibliography, including the dependence of pseudo-critical exponents on system size [1]. Interestingly, we have found that even with a strong dependence of these pseudo-critical exponents on system size, the equation of state is not affected by finite-size effects. The resulting scaling function is in good agreement with recent theoretical predictions for both exponent values and equation of state [2,3]. Our results suggest that while some effects may change the values of observed critical exponents of a given magnetic system, the state equation may retain its identity. In a sense, the state equation can be more universal than exponent values.

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# RI19

The transverse-field quantum Ising model on infinite-dimensional structures using quantum Monte Carlo method and finite-size scaling

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In a number of classical statistical-physical models, there exists a characteristic dimensionality called the upper critical dimension above which one observes the mean-field critical behavior. Instead of constructing high-dimensional lattices, however, one can also consider infinite-dimensional structures, and the question is whether this mean-field character extends to quantum-mechanical cases as well. We therefore investigate the transverse-field quantum Ising model on the globally coupled network and on the Watts-Strogatz small-world network by means of quantum Monte Carlo simulations and the finite-size scaling analysis. We confirm that both of the structures exhibit critical behavior consistent with the mean-field description. In particular, we show that the existing cumulant method has difficulty in estimating the correct dynamic critical expectation value can be a practically useful numerical observable to determine critical behavior when there is no well-defined dimensionality.

### **RI20**

# Study of cluster heterogeneity scaling in the two-dimentional Ising model

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The cluster heterogeneity, defined as the number of distinct cluster sizes, was suggested in [1] as an indicator of the phase transition in the percolation model. Noh et al. [2] found that the conventional finite-size scaling(FSS) form using the correlation-length exponent v does not describe the scaling of the cluster heterogeneity in the percolation model and that the other exponent v\_H should be used instead. We apply the FSS form suggested in [2] for cluster heterogeneity in the Ising model. Since there are two types of clusters, (+) and (-) spin clusters, we calculate the heterogeneity of each type (H+ and H-, respectively), and carry out FSS near critical temperature. We find that total heterogeneity (H = H+ +H-) scales with the exponent v\_H, as expected from [2] for the percolation model.

[1] H.K. Lee, B.J. Kim, and H. Park, Phys. Rev. E 84, 020101(R) (2011). [2] J.D. Noh, H.K. Lee, and H. Park, Phys. Rev. E 84, 010101(R) (2011).

### RI21

### Zero-temperature phase transition in a one-dimensional Ising ferromagnet using Glauber dynamics with a synchronous update Il Gu Yi and Beom Jun Kim\*

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Resently Glauber dynmaics for the one-dimensional Ising spin system at lowtemperature is studied variously not only experimently but also theoritically in the nano molecular system. Sznajd-Weron [Phys. Rev. E 82, 031120 (2010)] study the ferro anti-ferro phase transition of the density of active bonds, and the behavior of the relaxtion time which means the time needed to reach to steady state in the generalized zero temperature Glauber dynamcis. They suggested in that paper that the one-dimensional Ising model subject to the zero-temperature synchronous Glauber dynamics exhibits a discontinuous phase transition. We show here that the phase transition instead of a continuous nature, and we identify critical exponents, \beta \approx 0, \nu \approx 1, and z \approx 2 via a systematic finite-size scaling analysis.

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# **RI22**

### Stochastic treatment of magnetic moment relaxation in spin echo models

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The method of nuclear spin-echo amplitude calculation based on the density matrix technique is improved. Spin echo model [1] was supplemented with stochastic consideration of the magnetic moment relaxation provided by the interaction nuclear spin with fluctuating magnetic field. The theory of reducible systems [2] was applied for solving stochastic differential equations. This mathematical approach was adapted to find the solution of the Schrödinger equation for the time-evolution operator which describes behavior of a nuclear spin in the presence of a radiofrequency pulsed magnetic field and fluctuating magnetic field. NMR spin echo for I = 1/2 is considered as the simplest illustration of the method. Unlike classical approaches [3,4], It allows to calculate spin-echo amplitude correctly over whole period of relaxation without any estimates. The possibility of application of well-known Ito calculus [5] is also discussed. It was obtained more accurate formulae for spin-lattice relaxation time T1 and spin-spin relaxation time T2.

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### **RJ01**

### Vortex core switching on notched circular disks

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The vortex core manipulation by field has been frequently reported. However, the minimization of switching field is still being debated. In general, the vortex core switching is performed by annihilation of vortex and anti-vortex, however the nucleation of the pair requires large magnetic field. In this report, we proposed the new switching mechanism to decrease switching field by using circular disks with V-shaped notch. When the in-plane pulse field is applied, the vortex core moves to the notch and annihilates automatically. After the field is switched off, the vortex core gets back by demagnetizing field. The vortex core polarity is controlled by the external field perpendicular to the disk. In this switching mechanism, the vortex core switching field is decreased compared with the circular disks without notch. We investigated the vortex core switching by micromagnetic simulation. The field perpendicular to the disk was fixed at 50 Oe. The result reveals that the minimum switching field decreases 60% compared with the circular disk without notch.

### RJ02

## Effect of oersted field on magnetic vortex core gyration

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Recently, the vortex core manipulations by spin current are studied intensively. Because spin current creates the Oersted field, there are some reports, which study the effect of the field on the vortex core motion. When the vortex core is manipulated by the current, the electrodes are usually attached on top of the disk. Therefore, the current goes in and out from the top surface of the disk mainly. It makes the current non-uniform, and it makes the Oersted field complex. In this work, we report the effect of the Oersted field and the current distribution on the vortex core gyration by AC current by micromagnetic simulation. A Permalloy circular disk with 1.5µm diameter and 40nm thickness is used in the simulation. The AC current density is 2.66x1011A/ m2 mainly. The current frequency is varied around the resonant frequency of the disk. The result reveals that the effect of the Oersted field on the driving force is 8% in maximum. Furthermore, the vortex core is affected by not only the in-plane components but also the out-of-plane component.

# RJ03

### Mutual spin-transfer torque in vortex nano-oscillators

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We studied the emission response of spin transfer vortex oscillators with dynamic polarizer. In these systems, mutual spin transfer torque between the two ferromagnetic layers is at the origin of a very rich dynamics. By changing the injected current, we could gain access to up to three dynamic modes. At high currents, mutual spin transfer torque is at play and high-frequency dynamics is simultaneously excited in both layers. The detected mode corresponds to a Slonckzewski's windmill-like mode, which is stable in a large range of experimental conditions. The mode observed at low currents corresponds to the commonly observed dynamics in nano-contact vortex oscillators, in which the vortex core is far away from the contact and current-in-plane spin-transfer torque is at play. Surprisingly, we could gain access to a third mode at intermediate currents, which we ascribe to the precession of a vortex in the thin ferromagnetic layer under the contact area. This mode is of particular scientific interest because it allows one to experimentally study the vortex dynamics in the small amplitude limit.

### **RJ04**

### Key role of temperature in ferromagnetic bloch point simulations Kristof M. Lebecki<sup>1\*</sup>, Denise Hinzke<sup>1</sup>, Oksana Chubykalo-fesenko<sup>2</sup> and Ulrich Nowak<sup>1</sup>

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A Bloch point (BP) is known to play a crucial role during vortex core switching process [1]. Vortices themselves attract recently much attention since it was suggested to use them as a magnetic storage medium. For such an application BP-related vortex core properties like the switching speed or stability of the final state are essential. This makes the BP an important object of study. We present here results of numerical simulations, where the BP at elevated temperature is investigated. Our approach makes it possible to model its behavior without problems related to its singular character [1,2] - as opposed to all BP simulations presented in the literature so far. We have included temperature effects via Landau-Lifshitz-Bloch equation [3]. Our implementation, basing on the OOMMF code, allow us to compare results of analytical theories with modern numerical approach. We focus our attention on permalloy and conduct studies in the full temperature range. Firstly, we will present results of modeling an artificially pinned BP. This reveals details of its geometry, including its temperature dependent radius. Secondly, we will present results of dynamical studies, where the vortex core switches via nucleation, propagation and annihilation of a BP.

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### **RJ05**

### Polarization-selective signal propagation in a chain of vortices

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Dynamics on the subnanosecond time scale and potential technological applications give rise to a broad scientific interest in the dynamic properties of ferromagnetic microstructures with vortex magnetization configuration. An excitation of the low frequency mode of gynation can be transferred via dipolar interaction between neighboring ferromagnetic elements with the transfer efficiency strongly depending on their separation and the relative configuration of the vortex-core polarizations [1]. The dependence on the core polarizations can be understood considering the time-dependent shape of the external stay field. We employ time-resolved scanning transmission x-ray microscopy to study the vortex-core dynamics in a chain of three stray-field coupled permalloy squares. After exciting the first element via a short in-plane magnetic field pulse, the excitation can be transferred through the chain. No response is observed in the third square for equal core polarizations of all elements but the excitation can be transferred through the chain just by changing the polarization in the central element. For alternating polarizations, a transfer efficiency of about 56% to the third square is achieved. The chain can be switched back and forth between the transmitting and a locking state [2].

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# **RJ06**

# Theoretical study on frequency of vortex-antivortex pairs rotation in a magnetic thin-film with multi-contacts

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Spin transfer torque (SST) induced magnetization dynamics is one of the key concepts in spintronics and intensive studies have been conducted in the last two decades to realize STT based memories (Spin-RAM) and microwave oscillators. Finocchio et al. studied the STT induced magnetization dynamics in magnetic thin-film nanopillar which has one aperture and found that the rotational dynamics of a vortex-antivortex (VA) nair produces the resistance oscillation [1] Therefore, it is interesting to study how the VA-pair dynamics in a multi-contacts system is influenced by the distance between contacts. We performed micromagnetic simulations of a magnetic thin-film where the spatially localized spin-polarized current was injected through two nanocontacts and found the rotation of VA pairs. If the distance between two contacts is so large that the VA pair cannot interfere with each other, they independently rotate around each contact. In contrast, when the distance becomes small, the single VA pair rotates around a large area in which the two contacts are included and the rotation frequency drops down to about 60% of that of a single contact. Our results show that we can control the frequency of the VA pair rotation by changing the distance between the contacts

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## **RJ07**

# Parametric excitation and subcritical phase-locking in spin-transfer vortex oscillators

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In the last decade several systems have been proposed to design low power consuming nano-oscillators. One promising example is a MgO tunnel junction with a vortex free layer[1,2] (STVO). Classically, oscillations at its natural frequency f\_0 are obtained for currents larger than a critical value Ic. However, by applying a rf-excitation current at 2f\_0, i.e., parametric excitation[3,4], it is possible to force the system into oscillation even in the subcritical current regime (1-slc). The characteristics of such oscillation (emitted power and linewidth) are not comparable to the ones obtained in the super-critical regime. If I<4c, no oscillation is observed. If we inject a rf-current (-25 dBm) at 2f\_0, a small signal (parametric excitation) appears. When the rf-power is further increased (up to -15 dBm), the parametric signal gets synchronized. The resulting oscillation has the characteristic of the synchronized signal even though is obtained at sub-critical currents. Both mechanisms can be explained in the frame of non-linear vortex dynamics. Comparisons between the analytical approach and experimental results are presented. Support from ANR VOICE PNANO-09\_P231-36, EU MASTER NMP-P7-212257 and CANON-ANELVA for MTJ film are acknowledged.

A. Dussaux et al., Nat. Comm. 1 1 (2010) [2] A.V. Khvalkovskiy et al, Phys. Rev. B 80 140401(R) (2009)
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## **RJ08**

### Micromagnetic simulation for controlling the magnetic vortex chirality by current-induced Oersted field

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The polarity and the chirality are the magnetization direction and the magnetization curling of the magnetic vortex core, respectively. It is an important issue for developments of spintronics to control these quantities. The magnetic polarity is easily controlled by an external magnetic field. However, the magnetic chirality (MC) is not easily controlled. Recently, the MC in permalloy disks has been controlled by the shape of the disk [1] and a current-induced Oersted field[2]. We study magnetization dynamics for magnetization switching under the current-induced Oersted field at the permalloy thin film. By performing micromagnetic simulations based on the Landau-Lifshitz-Gilbert equation, we investigate an optimum condition for controlling the MC. We have calculated the magnetization reversal process in the permalloy thin film under an external uniform magnetic field. Three types of magnetization structure have been observed during the magnetization switching without an electrical current: a clockwise structure, a counterclockwise structure, and non-vortex core. Using the Oersted field induced by an electrical current perpendicular to the thin film, we have successfully obtained the endy clockwise structure (or counterclockwise structure). Details of the results and analysis will be reported.

S. Yakata, M. Miyata, S. Nonoguchi, H. Wada and T. Kimura, Appl. Phys. Lett., 97, 222503 (2010) [2]
 S. Yakata, M. Miyata, S. Honda, H. Itoh, H. Wada, and T. Kimura, Appl. Phys. Lett., 99, 242507 (2011)

## **RJ09**

**Stability of the vortex structure on the core switching by AC current** Tomonori Sato\* and Yoshinobu Nakatani

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Recently, magnetic disks with a vortex core have been considered as one of the candidates for the storage devices, and many reports relating to the core switching by spin current have been reported. Because the storage capacity depends on the magnetic disk size, it is important to investigate the stability of the vortex structure on its size. Until now, the disk size that the core can switch stable has not been reported. In this work, we investigate the stability on the core switching by micromagnetic simulation. In the simulation, two methods are used, the method by the core switching simulation by AC spin current and the method by the evaluation of the magnetic energy. The switching simulation reveals that the core does not switch around the transition dimension of magnetic disks between single-domain structure and the vortex structure on a remanent state, and the disk diameter needs to be increased about 20 to 80 nm for the stable core switching compared with the stability limit on the remanent state. Furthermore, the result of the energy evaluation agreed with the result obtained by the switching simulation. Therefore, the stability on the core switching can be estimated by the magnetic energy.

# RJ10

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## Time-averaged observation of magnetic vortex resonated in squareshaped NiFe films

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A dynamics of magnetic vortex confined in a patterned ferromagnetic thin film attracts much attention for both the fundamental and the technological points of view. The observation of the trajectory of magnetic vortex core excited with an application of ac magnetic field or ac current is significant for the quantitative understanding of spin dynamics. In this study, a time-averaged observation of magnetic vortex resonated in square-shaped NiFe thin film was demonstrated by means of magnetic force microscopy (MFM) to obtain the trajectory. A 30 nm-thick NiFe film with the lateral size of  $1 \times 1 \text{ um}^2$  was fabricated on a coplanar waveguide that generates an ac magnetic field of 1mT in amplitude. The resonant frequency of the magnetic vortex formed in the NiFe film was calculated to 260 MHz that is much higher than the sampling frequency of MFM. When the ac magnetic contrast appeared in the vicinity of center of the NiFe film. The diameter of the circle is comparable with that of the trajectory of magnetic vortex core calculated with micromagnetic simulation.

### **RJ11**

# Diverging-converging spin vortex pairs in biquadratically interlayer exchange coupled elements

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Spin structures have been an interesting topic of magnetism research for many years. Within this field, magnetic vortices have attracted much attention, due to their non-trivial topology and the various dynamic modes they exhibit [1]. A magnetic vortex consists of a planar, flux-closing magnetization curl that turns out of the plane in the central nanoscopic core region. In a single layer structure, the curl's radial components typically cancel each other out. Recent investigations show that this also holds true for multilayer vortex systems with bilinear interlayer exchange coupling [2]. Here we report on pairs of diverging-converging spin vortices occurring in biquadratically coupled systems. Using magnetic scanning transmission x-ray microscopy (STXM) we directly observe that the individual vortices of such pairs possess a residual radial magnetization component, i.e.  $\bigtriangledown Mxy \neq 0$ . This implies an additional perpendicular magnetization divergence  $\bigtriangledown Mz$ , for which we compare a continuous model with discrete micromagnetic simulations.

[1] S.-B. Choe et al., Science 304, 420 (2004). [2] S. Wintz et al., Appl. Phys. Lett. 98, 232511 (2011).

# RJ12

### **Magnetic vortex dynamics in exchange-biased micron-sized structures** Sofia De Oliveira Parreiras<sup>1</sup>\*, Flavio Garcia<sup>2</sup> and Maximiliano Delany Martins<sup>1</sup> <sup>1</sup>Applied Physics Laboratory, CDTN/CNEN, Brazil

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The study of magnetic dots with magnetic vortex spin configuration has recently attracted great scientific interest [1]. The great potential of applications of magnetic vortices (as for example magnetic memories and nanoparticles for cancer treatment) draws attention for the investigation of vortex proprieties. In this work, we studied the dynamics proprieties of exchange-biased vortex using the code OOMMF (NIST) [2] that applies the Landau-Lifshitz-Gilbert equation to simulate the spin configuration and compute the energy and magnetization of microstructures. A series of micromagnetic simulations for Permalloy/Fe50Mn50 disks with 0.5 ?m of diameter was done varying the magnetic coupling constant between the layers. We had observed that the vortex gyrotropic movement has a variable frequency that increases with the time, which is not observed when exchange bias is absent. Under a rotating magnetic field acting in the disks, the critical velocity for vortex polarity reversion increases with the coupling constant and frequency. Our results show that the critical velocity can be adjusted in a wide range by selecting the magnetic coupling constant and the oscillating frequency, i.e., it is possible to control the critical velocity for vortex polarity inversion through the exchange bias coupling.

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## RJ13

### **Magnetic Vortex Echo**

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The interest of magnetic vortices has greatly grown in the last years since these systems present a myriad of promising applications, which are mainly related on their dynamics aspects. When the vortex core is excited from its equilibrium position, it relaxes in a spiral motion, with a very well defined frequency, which is inversely proportional to the disk aspect ratio. The sense of the gyrotropic motion is determined by the core polarity. So, controlling the polarity, it is easy to control the sense of gyrotropic motion. In this work, we analytically shown that, properly manipulating the dynamic properties of the vortex, namely its polarity, in an analogous way as it is done in Nuclear Magnetic Resonance, it is possible to generate a magnetic vortex echo (MVE). This echo is similar to the spin echo, and it may provide fundamental information about the dynamic properties of real arrays of vortices, e.g., magnetic coupling, inhomogeneities, magnetic stability etc. To illustrate the MVE, we performed micromagnetic simulation of arrays of nanodisks, where it has been varying their most significant parameters to prove that the MVE can give significant information about the dynamic properties of the real array of disks.

ArXiv:1201.3553

## **RJ14**

### Origin of the dipolar coupling between vortex-state disks Ki-suk Lee and Sang-koog Kim\*

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Coupled vortex gyrations in spatially separated nanodisks have been studied theoretically[1,2], experimentally[3-8], and numerically[9] owing to their possible applications as an alternative to current signal processing devices with the advantage of negligible energy loss[7]. Fundamental phenomena such as eigenfrequency splitting and its dependence on dot-to-dot interdistance[3,4,7] have been verified experimentally and explained well by the normal mode representation[7,9]. In this work, we studied dynamic dipole interaction of coupled vortex gyrations observed in two disks by micromagnetic simulation and analytical approaches. Shibata et al.[1] reported that, based on rigid vortex model [10], the dynamic dipolar interaction of the coupled vortex gyrations originates from the magnetostatic interaction energies of the two dipolar-coupled disks and compared those that are obtained directly from the micromagnetic simulations and obtained from spin configurations based on two different side-charge-free-[10] and rigid-vortex models. The simulation result is in better agreement with the side-charge-free-model-based spin configurations of neighboring vortex-studed low. These results imply that the dipolar energy of the coupled gyrations of neighboring vortex-studed lows dipolar displar displares divertating the gyrations of the two dipolar energy of the coupled systems of neighboring vortex-studed lows from not the side-charge-free-model-based spin configurations than the rigid-vortex-model one. These results imply that the dipolar energy of the coupled gyrations of neighboring vortex-stude disks originates dominantly from not the side surface charges of each disk, but the volume charges.

This work was supported by Basic Science Research Program through the National Research Foundation of Korea(NRF) funded by the Ministry of Education, Science and Technology(No. 20110000441), [1] J. Stibtata et al.,Phys. Rev. B 67,224404(2003). [2] O.Y. Sukhostavets, et al.,Appl. Phys. Express 4,065003(2011), [3] A. Vogel et al.,Phys. Rev. Lett.105,037201(2011), [4] S. Sugimoto et al.,Phys. Rev. Lett.106,197203(2011). [5] A. Barman et al.,J. Phys. D43,422001(2010), [6] H. Jung et al.,Appl. Phys. Lett.972:2520(2010), [7] H. Jung et al.,Sci. Rep. 1.59;DOI:10.1038/srep00059(2011), [6] A. logel et al.,Phys. Rev. Lett.106,137201(2011), [9] K.-S. Lee et al,J. Appl. Phys.110,113903(2011), [10] K. Y. Guslienko et al.,J. Appl. Phys.91,8037(2022).

# RJ15

### Switching dynamics of vortex cores in nanodots by azimuthal-spinwave-mode excitation

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Owing to a possibility of implementations of magnetic vortices in future information-storage devices, vortex-core reversals have attracted much attention [1]. Very recently, ultrafast, low-power-driven core switching by azimuthal spin-wave mode was experimentally demonstrated [2]. Since their quantitative interpretation and deeper understanding remain clusive, we start to elucidate the underlying physics of the switching dynamics. In this study, we performed micromagnetic simulations on a Permalloy disk of 302 nm diameter and 50 nm thickness. We correlate the critical velocities (vc) and gyrofields (hz,cri) to the core switching driven resonantly by azimuthal mode excitations. For given disk's dimensions and geometry used here, the vc increases to ~400 and ~700 m/s for m = -1 (clockwise) and +1 (counter-clockwise) modes, respectively, although vc is known to be ~330 m/s for the gyration-mode-assisted-core-switching mechanism [3]. This difference originates from the difference of hz,cri between the gyration mode and azimuthal mode driven core switching mechanisms because the magnetization di of mz = -1 is necessary for the switching. From our calculations, hz,cri are estimated to be ~7 kOe and ~11 kOe for CW and CCW azimuthal mode excitations commared to the value of hz cri = 3 3 kOe for the syntation for CW and CCW

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### RJ16

### Logic operations based on magnetic-vortex-state networks

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Logic operations based on coupled magnetic vortices were experimentally demonstrated. We utilized a simple chain structure consisting of three physically separated but dipolar-coupled vortex-state Permalloy disks as well as two electrodes for application of the logical inputs. We directly monitored the vortex gyrations in the middle disk, as the logical output, by time-resolved full-field soft X-ray microscopy measurements. By manipulating the relative polarization configurations of both end disks, two different logic operations are programmable: the XOR operation for the parallel polarization, and the OR operation for the antiparallel polarization. This work paves the way for new-type programmable logic gates based on the coupled vortex-gyration dynamics achievable in vortex-state networks. The advantages are as follows: a low-power input signal by means of resonant vortex excitation, low-energy dissipation during signal transportation by selection of low-damping materials, and a simple patterned-array structure.

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### RJ17

### Vortex-gyration-mediated magnonic crystals

### Dong-soo Han and Sang-koog Kim\*

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Recently, mutual energy transfer between spatially separated magnetic disks has been studied and interpreted based on a coupled-vortex-oscillator model [1-5]. This novel energy-transfer mechanism provides the advantages of tunable energy transfer rate, low-power input signal, and extremely low energy dissipation for the case of using negligible damping materials. Here, we extended this study into magnonic crystals of one-dimensional (1D) and 2D arrays of dipolar-coupled magnetic disks. We analytically and numerically explored the fundamental modes of collective vortex-gyration excitations in such 1D and 2D disk-array magnonic crystals of four different ordering of vortex polarization (p) and chirality (C) states. Both the analytical and numerical calculations show excellent quantitative agreement on contrasting dispersion branches for different orderings of p and C states. Such collective vortex-gyration modes can be understood in terms of the energy variation of dynamic dipolar interaction between the neighboring vortex-state disks. These results provide for the possibility of vortex-gyration-mediated magnonic crystals as applications for future signal processing devices, offering the advantages of low energy dissipation and low power input. This work was supported by Basic Science Research Program through the National Research Foundation of Korea (NRF) funded by the Ministry of Education, Science and Technology (No. 20110000441).

# **RJ18**

# Vortex-gyration transfer rate and energy attenuation in coupled nanodisks

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One of the most important research goals in current information-signal processing technologies is to enhance signal-processing speed and to reduce energy loss[1]. In recent electronic devices, information-signal transports are based on the motion of electron charges. Interestingly, coupled vortex oscillations in nanodots can be used as an alternative mechanism of information-signal processing[2-4]. Here, we report a study of, by analytical and micromagnetic numerical calculations, information-signal transfer in vortex-state nanodisks that is available with assistance of vortex-gyration transfer through the dynamic dipolar interaction. We analyticall parameters for controlling signal processing speed and energy attenuation that are both technologically essential parameters for controlling signal processing speed and energy loss in signal processing devices. It is found that the transfer rate can be controlled by the relative polarization configuration p1p2, the saturation magnetization Ms, and the radius (R)-to-thickness (L) ratio R/L of given magnetic disks as well as interdistance. The energy attenuation is governed by not only the intrinsic damping constant of a given material, but also the values of Ms, L, and R. The analytical results are in reasonable agreements with micromagnetic simulation results. This work provides a foundation of manipulating information-signal processing between vortex-state networks.

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## RJ19

# Effect of spin-motive force and spin-diffusion on a vortex dynamics Jung-hwan Moon<sup>1</sup>, Aurelien Manchon<sup>2</sup> and Kyung-jin Lee<sup>1</sup>\*

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Theoretically, it is known that the dynamics of magnetic texture highly depends on damping ( $\alpha$ ) and non-adiabatic parameter ( $\beta$ ). However, the value of  $\beta$  is controversial both theoretically and experimentally. Moreover, it was recently reported that  $\alpha$  and  $\beta$  are also affected by the spin configuration. The former comes from spin-motive force(SMF) [1, 2] and the latter comes from torque which is exerted from spin-diffusion [3]. In order to understand the spin dynamics under complex spin texture, it is needed to study the effect of SMF and spin-diffusion on  $\alpha$  and  $\beta$ . In this work, we performed micromagnetic simulation using LLG equation with spin torque term, SMF, and spin-diffusion. The current-induced dynamics of a vortex core is micromagnetically modeled using a computational framework based on the fourth-order Runge-Kutta method. The model system is a circular Permalloy disk with the thickness of 20 nm and the diameter of 270 nm which is vortex favored dimension. Initial trajectory with spin-diffusion must affect and using a computation is governed by non-adiabaticity [4]. The further studies will be discussed in detail.

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# **RJ20**

# Soft X-ray microscopy of non-linear magnetic vortex core motion

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Magnetism on the nanometer length scale and its fast spin dynamics is a scientfically highly attractive and technologically relevant topic. It addresses fundamental magnetic length scales e.g. magnetic exchange lengths in the sub-10mm range and fast time scales in the sub-ns regime where precessional and relaxation phenomena, domain wall motion and vortex dynamics occur [1,2]. New technological concepts such as spintronics require precise control the electron spin on a nanoscale with psec timing. We use time magnetic soft X-ray microscopy providing a spatial resolution down to 10nm and a temporal resolution below 100ps [3]. As one example, we report experimental results showing non-linear behavior in the resonant excitation of magnetic vortex cores with increasing amplitude of the exciting magnetic AC fields [4]. Supported by the Director, Office of Science, Office of Basic Energy Sciences, Materials Sciences and Engineering Division, of the U.S. Department of Energy.

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### **RK01**

#### Dynamics of successive minor hysteresis loops

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Cumulative growth of successive minor hysteresis loops in Co/Pd multilayers with perpendicular anisotropy was studied in the context of time dependent magnetization reversal dynamics. We show that in disordered ferromagnets, where magnetization reversal involves nucleation, domains' expansion and annihilation, differences between the time dependencies of these processes are responsible for accumulation of nuclei for rapid domain expansion, for the asymmetry of forward and backward magnetization reversals and for the respective cumulative growth of hysteresis loops. Loops stop changing and become macroscopically reproducible when populations of upward and downward necleation domains balance each other and the respective upward and downward reversal times stabilize.

### RK02

Gilbert damping constants of exchange biased NiFe/FeMn bilayers Jungbum Yoon, Hyeok-cheol Choi and Chun-yeol You\* Department of Physics. Inha University: Korea

The physics of exchange bias effect has been studied intensively for the past decades, however, a few experiments have been reported about the spin dynamics and damping mechanism for the exchange biased systems. Recently, the vector network analyzer ferromagnetic resonance (VNA-FMR) is employed to research the spin dynamics of ferromagnetic thin films with wide range microwave frequencies. In this study, the dynamics magnetic properties of the NiFe/FeMn bilayers is investigated by VNA-FMR with various external static field. The exchange bias is verified by vibrating sample magnetometer and VNA-FMR with varying the thickness of FeMn. Spin dynamics and the Gilbert damping constants of exchange biased NiFe/FeMn bilayers are investigated by the analysis of FMR spectra. The exchange bias field induced asymmetry in the magnetization hysteresis loops. It implies the spin dynamics must be asymmetry for positive and negative field region. We perform VNA-FMR measurement for positive and negative fields. In results, we find that the apparent damping parameters are different in the both field directions in the exchange biased NiFe/FeMn bilayers. Therefore, we conjecture that the exchange bias layer acts differently, depends on the relative direction of the ferromagnetic layer magnetization to the exchange bias field.

#### **RK03**

# Non-linear susceptibility and influence of the applied magnetic field on ZFC/FC curves

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ZFC/FC curves are widely used to characterize assemblies of magnetic nanoparticles. They reflect the crossover between the blocked and superparamagnetic (SP) regime with increasing temperature. With a low applied field a linear response can be assumed, the shape of the curves is then independent of the applied field H, and a simple theoretical modeling is possible: this allows efficient theoretical fits of experimental curves. We have studied the influence of the applied field magnitude on the ZFC/ FC curves shape, both theoretically and experimentally. While the effect of H on the energy barriers has already been discussed, its effect on the response of SP particles has not been considered. However, this non-linearity manifests itself much before the modification of the switching energy distribution. In addition to experimental measurements on a diluted Co nanoparticle assembly, we have simulated ZFC/FC curves for different applied fields, including the third-order susceptibility in the SP response. The later depends on the anisotropy and does not at all correspond to a Langevin function around the blocking temperature. We find that the curves can be significantly affected (in particular the low temperature limit of the FC) for quite low applied fields, which are usually used in experiments.

### **RK04**

#### On the relation between the magnetoelastic effect and the damping constants of $(Ni-Fe)_{1x}M_x$ (M = Ag, Cr, Ga, Au, Pd, and Pt) films Yasushi Endo\*, Yoshio Mitsuzuka, Yutaka Shimada and Masahiro Yamaguchi

*ECEI, Graduate School of Engineering, Tohoku University, Japan* The dynamics of magnetic thin films have received attention in magnetic device applications such as recording heads, media, and MRAM. Although the dynamics significantly depend on the damping constant *a*, which determines the strength of

damping torque in magnetic thin films, details about their damping mechanism, especially the magnetoelastic effect on damping remain unclear. Herein to clarify the effect of  $\lambda s$  on  $\alpha$  in (Ni-Fe),  $\mathcal{M}_{x}$  films in detail, we evaluated  $\alpha$  and  $\lambda s$  in these films. For M= Au, Pd, and Pt,  $\alpha$  increases linearly and negative  $\lambda s$  increases as x increases. In contrasts,  $\alpha$  and positive  $\lambda s$  tend to increase almost linearly with x when M=Ag, Cr, and Ga. These results provide clear evidence that  $\alpha$  is correlated with  $\lambda s$  in (Ni-Fe)-M films. Furthermore, increments of  $\alpha$  and the absolute value of  $\lambda s$  to x are markedly enhanced in the order of M = Pt, Au, Pd, Cr, Ag, and Ga, suggesting that 5d transition metal dopants are more influential on both  $\alpha$  and  $\lambda s$  than 3d and 4d transition metal dopants due to the strong spin-orbit interaction of 5d dopants. Consequently, these results demonstrate that a change in magnetostriction energy via transition metal dopants can effectively control  $\alpha$ .

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#### RK05

#### Ultrafast magnetization dynamics of ferromagnetic systems induced by mid infrared laser pulses

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The aim of the present study is to show how the ultrafast demagnetization and the subsequent rapid re-magnetization occur when exciting a ferromagnetic material with low energy infrared pulses. We have used mid-infrared femtosecond laser pulses ( $\lambda$ = [3-10 µm]) to excite CoPt3, Ni and Co ferromagnetic thin films. The magneto-optical response is then probed in the visible ( $\lambda$  = 798nm). Our results show that even though only intraband transitions occur, the demagnetization process and its subsequent relaxation to the lattice and to the environment are still the dominant processes involved in the magnetization dynamics. We also show that the material band structure is important to interpret the thermalization dynamics of the spins that occur before the heating of the lattice. For specific experimental configurations, we show that it is possible to induce a motion of precession of the magnetization dynamics induced at 6.5 µm in nickel shows an oscillatory behaviour with a period of 2 ps. We attribute this result to the excitation of a two-magnons mode on the NiO by an acoustic mode generated in nickel.

#### RK06

#### Ferromagnetic resonance of bilayer CoFeB/NiFeSiB thin film

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Spin transfer torque magnetic random access memory (STT-MRAM) is one of the candidates for next generation random access memory. For practical application, the device should have high thermal stability and low critical current density. They have strong correlation with saturation magnetization (Ms) and damping constant ( $\alpha$ ) so that we need to determine Ms and  $\alpha$ . To study the damping mechanism of the thin films, we have performed ferromagnetic resonance (FMR) experiments through vector network analyzer with bilayer thin films of CoFeB(10-x nm)/NiFeSiB(x nm) (x=0,1,2,...,10). We have fitted the FMR data with Lorentz function to get the information of resonance frequency and line width. From these two fitting parameters, we could get Ms and  $\alpha$  by Kittel formula and Full Width at Half Maximum. As increasing x values, the Ms tends to increase, whereas the  $\alpha$  suddenly decrease and then saturate. This implies that the added NiFeSiB layer in CoFeB plays a role to prevent the spins align along the field direction.

# RK07

# Observation of non-kittel ferromagnetic resonance in Co/Cu multilayer system

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In this research, we study the ferromagnetic resonance behavior (FMR) of Co/Cu multilayer system. The samples stacking Ta (5nm) / [Co (4nm) / Cu (t)]\*10 / Ta (5nm) with t = 0.5 ? 6 nm were prepared by dc magnetron sputtering in a multitarget sputtering chamber at room temperature and 1 mTorr of pure Ar-pressure. The magnetoresistance (MR) was measured by the standard four point measurement with current in plane (CIP) configuration using Quantum Design Physical Properties Measurement System (PPMS). The magnetic hystersis loops were recorded using a Vibrating Sample Magnetometer (VSM). FMR responses were measured at microwave frequencies by means of a Vector Network Analyzer (VNA) equipped with a low pass filter circuit. Clear FMR responses were recorded for all samples, however, the noise increased for samples with thicker spacer. Then resonance frequencies (f) acquired from FMR responses. Theoretically, this FMR should follow Kittel frequency given by: f=( $\eta/2\pi$ )/H(H+4 $\pi$ M). However, our results shows that the resonace linewidth and the deviation of resonance frequency from kittel behavior will increases as the spacer thickness increased. This behavior is the result of RKKY interaction evident by the increased GMR and will be dicussed within the manuscript.

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#### RK08

#### Neighboring layer dependence of ultrafast thermo-magnetic property in GdFeCo films

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Femtosecond pulsed laser light allow excitation of magnetic systems much shorter than the time scale of thermal diffusion represented by conventional Fourier's law. In this time scale, heating and demagnetization phenomena arise via strongly non-equilibrium non-adiabatic way, and cannot be explained by conventional equilibrium thermo dynamics description[1]. In this study, we investigated the neighboring layer dependence of ultrashort laser-induced thermal/magnetic response in layered GdFeCo films by all-optical pump-probe method. Simultaneously, change of normalized reflectivity was measured for monitoring the time evolution of electron temperature Te. We designed layered structures as same 20 nm thick GdFeCo with different neighboring layers (conductive AITi and insulating SiN). We found two characteristic time region from magnetic behavior: (A) rapid step-like demagnetization and (B) following recovering process with precessional motion. The time scale of (A) is conformed as within picoseconds range (time constant ~100 fs) independently with film structure, which is much shorter than ferromagnetic resonance (period ~0.1) ns at 10 GHz ) and hundreds fs delayed with respect to the increase of Te. Following regime (B), film structural dependency of precessional motion of lattice temperature in magnetic layer.

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#### RK09

#### Magnetization dynamics in perpendicular magnetic anisotropy CoFeB/ MgO system

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We study the ultrafast magnetization dynamics of perpendicularly magnetized MgO/CoFeB/Ta and Ta/CoFeB/MgO stack structures which gains wide attention from a viewpoint of perpendicular magnetic tunnel junctions. An all-optical time-resolved magneto-optical Kerr effect measurement reveals that effective Gilbert damping  $\alpha$  stays unchanged at ~0.02 in high external field regardless of pump fluence, but it declines drastically with the increase of pump fluence in weak external field. This can be explained by the enlarged apparent relaxation time due to slow remagnetization. Genuine damping tends to attract the precessing magnetization vector towards the effective equilibrium axis, while the tip of the reduced magnetization vector after pump pulse heating recovers growing away from the effective axis. These two competing contributions determine the apparent relaxation time. In a weak field regime, slow recovery of the magnetization wector results in the increased relaxation time and low effective Gilbert damping. We believe that low Gilbert damping found in CoFeB/MgO structure will be expected to reduce the critical current for current-induced magnetization switching.

# RK10

# Composition dependence of the gilbert damping constant for co-based heusler alloy

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Spin polarized current attracts much attention to operate the magnetic memory devices. To reduce the spin current density, the materials with small Gilbert damping are needed. Recently, it is reported that the Gilbert damping constant of Co-based Heusler alloy known as half-metallic ferromagnet is less than 0.01[1], although the Gilbert damping factor of Co-based Heusler alloy are closely related to crystalline structure and concentration ratio. In this study, to understand the damping mechanism in high-quality L<sub>2</sub>1-type Co-based Heusler alloy we measured ferromagnetic resonance (FMR) and estimated the Gilbert damping factor 25 nm thick L<sub>2</sub>-type Co<sub>2</sub>-Fe<sub>3</sub>Si (x=0, 1, 1.5) films for which degree of crystalline order were found to be about 70%/[2] were grown on Si(111) substrates. The saturation magnetization of Co<sub>2</sub>-Fe<sub>3</sub>Si (x=0, 1, 1.5) films for which degree of crystalline order were found to be about 70%/[2] were grown on Si(111) substrates. The saturation magnetization of Co<sub>2</sub>-Fe<sub>3</sub>Si (x=0, 1, 1.5) films for which degree of crystalline order were found to be about 70%/[2] were grown on Si(111) substrates. The saturation magnetization of Co<sub>2</sub>-Fe<sub>3</sub>Si (x=0, 1, 1.5) films for which degree of crystalline payer and coplanar waveguides (CPW) which consists of Ti (5 nm)/Au (60 nm) were prepared. Following that, FMR was measured using a vector network analyzer. The results show the Gilbert damping constant inversely proportional to saturation magnetization. Considering the electron density of state from the ab initio calculation, the broadening of FMR spectra is considered to be related to the electron density of state.

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### **RK11**

# Detection of picosecond magnetization dynamics of 50 nm magnetic dots down to the single nanodot regime

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We present the detection of the picosecond dynamics in arrays of 50 nm permalloy dots down to the single nanodot regime by an all-optical time-resolved magneto-optical Kerr effect microscope. The inter-dot separation (S) varies from 200 to 50 nm and simulated magnetostatic fields shows a transition from magetostatically isolated to strongly coupled regime as S decreases. Consequently, we observe a single precessional mode for S down to 75 nm, whose frequency increases with the decrease in S. At the smallest separation S = 50 nm, we observe a mode splitting. The simulated mode profile reveals that the dynamics of a single 50 nm dot is dominated by the edge mode. In sparsely packed arrays we primarily observe the isolated dynamics of the constituent dots in phase. For S = 50 nm, we observe an additional backward volume magnetostatic mode of the array. The damping is minimum for S = 200 nm but increases linearly with the decrease in S as a result of the dynamic dephasing of the proposition of two resonant modes results in a sudden increase in the apparent damping.

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# RK12

# Femtosecond demagnetization in Ni: Electron-phonon spin flip scattering from first principles

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The femtosecond demagnetization discovered in 1996 [1] represents a critical test of magnetization dynamics theories and may lead to many interesting applications. It is still far from being understood on a microscopic level. Electron-phonon spin-flip scattering in Ni was suggested to be the microscopic explanation of its femtosecond laser-induced demagnetization [2]. We have calculated the spin-flip Eliashberg function [3] based on ab initio electron-phonon coupling matrix elements, which allows us to obtain the spin-flip probability with much higher accuracy. We extend this method also to the regime of non-equilibrated electron distributions relevant for ultrafast processes. We have found that the spinflip probability depends strongly on electron energy. We consider two cases for system excited by a laser: thermalized very hot electron distributions, as well as highly non-equilibrium electron distributions that are expected to be present immediately after the fs laser excitation. Employing this approach we compute the electron-phonon SF rates. We find that the demagnetization rate is very low for any thermalized electron distribution as compared to non-equilibrium distributions present within first femtoseconds following the pump laser [3]. This is due to the density of states and the specific energy-dependence of SF probability.

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# **RK13**

#### Ultrafast magneto-acoustic pulses in a nickel film Jiwan Kim, Mircea Vomir and Jean-yves Bigot\*

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We report about the ultrafast magnetization dynamics induced by magneto-acoustic pulses in a 200-nm-thick Ni film generated with femtosecond laser pulses. The magneto-acoustic pulses result from the coupling between the magnetization and the acoustic waves generated by the laser pulses. In order to distinguish between dynamical magnetic effects induced either by the thermal excitation or by the magnetoacoustic pulse, the spin dynamics is measured from both the front and rear sides of the film via the magneto-optical Kerr technique. It is found that the acoustic pulses excite the magnetization on both sides of the film and the perturbation of the magnetization is very efficient at the rear side (10% of the static one). Using a detailed modeling of magneto-acoustic pulses combining the concepts of acoustic pulse propagation and ultrafast magnetization dynamics, we reproduce the magnetization dynamics on both sides of the film. In addition, our results imply that the magnitude of magnetoacoustic pulses can be controlled and maximized by selecting proper substrates with same ferromagnetic materials. We forecast that our results will have a strong impact for making ultrafast magneto-acoustic devices, with the capability of sensing the magnetization at relatively long distances from acoustic pulses generated by the laser pulses.

### RK14

# Minimal precessional and switching currents for relaxing-precessional magnetization reversal within a spin valve

Jui-hang Chang\*, Hao-hsuan Chen and Ching-ray Chang Physics, National Taiwan University, Taiwan

The relaxing-precessional magnetization reversal [1] is studied from the point of view of nonlinear dynamics. The solution of the Landau-Lifshitz-Gilbert equation with spin-transfer toque [2] shows that there are two critical values,  $\alpha_{-}$ es and  $\alpha_{-}$ cp, for the damping constant, and  $\alpha_{-}$ cs  $< \alpha_{-}$ cp. Above  $\alpha_{-}$ cs the minimal switching current  $a_{-}$ s is the same as the modified Stoner-Wolfarth (SW) limit [3]  $a_{-}$ cs for the switching, and above  $\alpha_{-}$ cp the minimal precessional current  $a_{-}$ p is the same as the modified SW limit  $a_{-}$ cp for the onset of precession. For a given magnetic anisotropy and an arbitrary in-plane bias field, condition  $a_{-} < a_{-}$ cs  $< a_{-}$ p  $< a_{-}$ p always happen, where  $a_{-}$ s,  $a_{-}$ p are functions of  $\alpha_{-}$ s,  $\alpha_{-}$ p, respectively. These investigations will be of importance for the design of spin-torque-transfer magnetic random access memories [4] and nano oscillators [5].

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### RK15

# Time dependent dichroism induced near the surface plasmon of Au nanoparticles

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We have studied the Surface Plasmon (SP) dynamics of Au Nanoparticles, with a diameter of 50 nm, excited by circularly polarized femtosecond laser pulses. We report about a new effect observed in a time resolved pump-probe experiment, analyzing the polarization state of the probe pulses. It manifests as a circular dichroism and an optical rotation induced in the vicinity of the SP when pumping with left ( $\sigma$ +) or right ( $\sigma$ -) circularly polarized pulses. We attribute this effect to a time dependent change of the orbital momentum of conduction electrons as the effect is more pronounced when the nanoparticles are excited with a pump wavelength -pump = 800 nm as compared to 400 nm. Indeed, in that case the interband transitions from the d-band to the conduction band are minimized with respect to the Drude electrons. The induced dichroism is resonant on the SP (560 nm). Its lifetime is comparable to the energy relaxation time of the quasiparticles to the lattice. It suggests that the electron-phonon interaction is the main mechanism for the dissipation of this pump induced orbital momentum. The detailed behavior of the SP dynamics as a function of probe wavelength, pump polarization and pump-probe delay will be discussed.

# RK16

#### The effect of surface anisotropy on the switching of a particle magnetic moment

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The effect of surface anisotropy on the switching of a particle magnetic moment Shuang Guo, An Du\* College of Sciences, Northeastern University, Shenyang 110004, China \* Corresponding author(du\_an\_neu@126.com) The dynamic precession of the moment for a spherical particle in a microwave field was studied by using Landau-Lifshitz-Gilbert (LLG) equation. The spins inside the particle have singleion anisotropy and the ones on the surface of the particle have the surface anisotropy normal to the surface. The switching field threshold was calculated for different surface anisotropy with definite microwave frequency. It is found that the surface anisotropy the threshold increases the switching speed of the magnetic moment increases obviously.

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#### RK17

# Magnetization dynamics of GdFeCo nanostructures revealed with PEEM

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Radboud University Nijmegen, Institute for Molecules and Materials, Netherlands

The manipulation of spins is a very exciting topic from fundamental point of view as well as for practical applications. Combining experiments and simulation, we have been able to demonstrate that a fs-optical excitation is sufficient to trigger magnetization reversal in GdFeCo nanostructures on very short timescales. Employing a photoemission electron microscope (PEEM) at the SIM beamline, we have proved that we can manipulate the magnetization of nanostructures by using a heat pulse only. Performing time resolved X-ray magnetic circular dichroism (TR-XMCD) measurement we have observed that the magnetization reversal within the structures occurs on a timescale faster than 100 ps and evidenced that the reversal occurs against an external applied magnetic field. In our experiment the reversal happens only by heating the system on the time scale of the exchange interaction of the two sublatices and does not require any other external stimulus.

'Ultrafast heating as a sufficient stimulus for magnetization reversal' T. A. Ostler, J. Barker, R. F. L. Evans, R. Chantrell, U. Ataita, O. Chutykalo-Fesenko, S. El Moussaoui, L. Le Guyader, E. Mengotti, L. J. Heyderman, F. Nolting, A. Tsukamoto, A. Itoh, D. Afanssiev, B. A. Ivanov, A. M. Kalashnikova, K. Vahaplar, J. Mentink, A. Kirihyak, Th. Rasing and A. V. Kimel, Nature Comm. 3, 666 (2012).

#### RK18

# A study of magnetic domain and magnetization reversal in L-shaped Py

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The evaluation of characteristic of nanometer-sized magnetic domain is one of the most important issues in the fields of spintronics. We made the Py stripe with and without the L-shaped edge structure. First, magnetic domain and its reversal reversal behavior were simulated with the use of OOMMF program. Depending on the thickness, width of this L-shaped structure, magnetic domain and its reversal behavior could be classified in several groups. We also used scanning electron microscope with polarization analysis(SEMPA, or spin-SEM) to probe the magnetic domain pattern predicted in the simulation.

# RK19

# Ultrafast dynamics of ferromagnetic copd thin film by various polarized probe beam

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We study ultrafast dynamics of the ferromagnetic CoPd thin film by using timeresolved magneto-optical Kerr effect. In the case of using linearly polarized probe beam, the spin precession is observed in the frequency of about 14 GHz in the range of large time scale. On the other hand, when the probe beam changes to circular polarization, the large time scale precession disappears, and ultrafast precession (~100GHz) is observed in short time scale. We suppose that these result from the momentum change of probe beam from 0 to 1, and the ultrafast precession is related to interaction between spin and circularly polarized photon.

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#### RK20

# Low temperature time domain THz spectroscopy of terbium gallium garnet crystals

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We have used the terahertz (THz) time domain spectroscopy to study high frequency magnetic excitations in terbium gallium garnet (TGG) crystals cut along <111> and <001> crystallographic planes. We demonstrate that a THz bandwidth transient electromagnetic pulse can efficiently couple to magnetic moments in TGG. By comparing the spectrum of the pulse before and after transmission through the crystal, we are able to isolate the absorption corresponding to magnetic resonance modes of TGG. We measure and discuss the dependence of the observed modes upon the temperature and the strength and orientation of the bias magnetic field with respect to the crystallographic axes of the crystals. The magnetic modes are present at temperatures above the Neel point, which is interpreted in terms of the field-induced magnetic close to the TGG absorption band destroys the magnetic ordering. Thus, the light induced demagnetization of TGG is observed. Our findings demonstrate that the time domain THz spectroscopy can be a powerful tool by which to study high frequency properties of dielectric magnetic materials.

#### RK21

# Ferromagnetic resonance of a single micron dot using vector network analyzer

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Magnetization dynamics of small ferro-magnets in the gigahertz region have been investigated intensively for the applications in microwave devices. It is based on the precessional motion of magnetization and it is necessary to control one of each small magnets in devices. In this paper, we investigated ferromagnetic resonance (FMR) of a single miconscale dot of permalloy (Py) using a vector network analyzer (VNA). A micron-scale Py dot and a coplanar wave guide were fabricated using electron beam lithography, electron beam evaporation and lift off technique. The thickness of the Py dots was 30 nm, and the shapes were square and rectangle with a width and length of 1~10 um and 1~40 um, respectively. The FMR measurement was performed using the VNA and probe station. In the square shape dot of 10 um, the resonant frequency depends on magnetic fields, which was good agreement with the Kittel's equation of a thin film. In the rectangle shape dots, the resonant frequency shifted to higher frequency with at demagnetization effect becomes large when the width is less than 1 um for a thin film.

# RK22

Relation between gilbert damping constants and perpendicular magnetic anisotropy in Ti buffered Co/Ni multilayers

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Recently, there has been a growing interest in spin-transfer-torque magnetic random access memory utilizing perpendicularly magnetic anisotropy (PMA) materials in order to overcome thermal stability problems. Since the critical current density depends on the Gilbert damping parameter, it has been an important issue to understand and manipulate the Gilbert damping parameter in the PMA materials. In this work, we have investigated PMA [Co/Ni] multilayers with Ti buffer layers by an all-optical pump-probe time resolved magnetic optical Kerr effect (TR-MOKE). In particular, we have studied the variation of Gilbert damping constant ( $\alpha$ ) and PMA as a function of the thickness of Ti buffer layer thickness (). Since the PMA and damping constant are strongly related with the spin-orbit coupling, both physical quantities must be correlated. Clear damped oscillations of the magnetization are observed in TR-MOKE measurements. After background subtraction, the signal is fitted with a damped harmonic function, from which the precession frequency (f) and the decay time ( $\tau$ ) are deduced. We obtained f and  $\tau$  by fitting with Landau-Lifshitz-Gilbert equation, and we could be estimated  $\alpha$ . We find that The  $\alpha$  and PMA values increased monotonically with increasing oft. This result clearly shows close relationship between PMA and  $\alpha$ .

### **RK23**

#### Chaotic motion of magnetic domain structure under alternate field

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Magnetic domain motion in a garnet thin film under alternate magnetic fields up to 5000 Hz has been investigated at room temperature. Domain structures and motions are observed by using a high-speed video camera with the help of magneto-optical Faraday effect. When a field frequency is low, the magnetization changes periodically and a domain pattern has a labyrinth structure. By increasing the field amplitude and driving frequency, irregular oscillations of a magnetization appear. Under a rapidly oscillating field, chaotic motions of domain are observed. In this region, domain structures have a disk-like shape. These disks grow from some crystal defects. And a growing point shows a branch-like form.

# **RK24**

# Demagnetization dynamics observed by spin-resolved ultrafast x-ray photoemission

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We report on the experiment and the results of time- and spin-resolved photoemission. Ultrafast demagnetization was first observed by Beaurepaire et al. in 1996 and many approaches where done to support these findings. Our method allows for measuring the whole valence band mediated by the cascade electrons emerging from the sample and therefore gives a direct measure of the sample's magnetization. Ultrafast demagnetization is observed on thin films of Fe on W(110) by optical pumping at 800 mm and x-ray probing at 7 nm. The measured demagnetization time of 450 fs is limited by the experiment. Although space charge effects limits the photoelectron gain the measured spin asymmetry stays almost constant with increasing x-ray flux and only drops at very high x-ray fluxes (> 4 nJ/pulse). We also show the feasibility of single shot magnetic measurement. The experimental setup consists of a completely mobile two chambers ultra-high vacuum system (preparation and measurement chamber) with Mott-polarimeter which can be brought to the free electron laser FLASH, Hamburg.

# **RL01**

#### Synchronized modes of in-plane/out-of-plane spin-torque oscillators in MTJ with synthetic ferrimagnetic free layer

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Javan

There has been considerable interest in the phenomena of spin-torque oscillators (STO) in magnetic tunnel junctions (MTJs) for device applications as well as the current induced magnetization switching (CIMS) resulting from the spin transfer torque (STT). In particular the MTJs with synthetic ferrimagnetic (SvF) free layer is important since we can expect the cooperation and/or competition of the magnetization dynamics due to the interlayer coupling between magnetizations in the ferromagnetic bilayer. We analyze the magnetization dynamics self-consistently in the MTJs with SyF layer by iterative calculation following two steps. The STT in the ballistic regime is estimated, and the magnetization reversal is simulated by the LLG method at the finite temperature. We discuss the effect of the interlayer coupling on the STO. As the interlayer coupling increases, the change of the STO mode is observed, where the STO mode moved from the out-of-plane precession (OPP) to in-plane precession (IPP). From the mapping of the STO behavior on the plane with applied current and interlayer coupling, synchronized mode of STO appears in the region where the effective filed corresponding to the interlayer coupling is larger than the anisotropy field. These results suggest that the STO appears even without external magnetic field.

### **RL02**

#### Interface material effects on magnetic anisotropy and its electric field induced variation in thin films

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The spintronics has grown up intensively to realistic applications in the technology of magnetic random access memory (MRAM) development. Such development has been remarkable in memory density, reading-writing speed, and non-volatile property in cooperation with the technologies of spin-injection and physics of spin transfer torque. The basic physics about magnetism has been developing in the response to electric field (EF). This has emerged as the connection with low power consumption device and small energy scale of magnetoelectric effects. For the thin films [1,2] as a memory, sensitivity and large response may be required in device applications. Due to the limitation of EF penetration into the thin film, interface with a few metallic layers is critical to determine the response to the EF. We have investigated MgO/Fe/M(001) (M = Au, Pt) [3,4] with using the density functional calculation, in which the substrate of Pt was found to enhance the EF effect on magnetic anisotropy. In this work, in order to investigate influences of stacking structure for the interface, we investigate MgO/Fe/Pt/Au(001) and MgO/Fe/Pd/Au(001). In addition, results in MgO/Au/Fe/Au(001) are also discussed in the connection with segregation effects of film preparation, which may be expected in experiments.

[1] Y. Shiota et al., Appl. Phys. Express 2, 063001 (2009). [2] T. Nozaki et al., Appl. Phys. Lett. 96, 022506 (2010). [3] M. Tsujikawa et al., J. Appl. Phys., 109, 07C107 (2011). [4] M. Tsujikawa et al., submitted for publication.

### **RL03**

#### Effect of spin relaxation rate on the interfacial spin depolarization in ferromagnet/oxide/semiconductor contacts

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The electrical injection and detection of spin-polarized carriers in semiconductors (SCs) has been successfully achieved by employing spin tunnel contacts and the Hanle effects. However, many aspects of the spin phenomena in these systems, e.g., (i) the location, magnitude, and sign of the induced spin accumulation, (ii) the unusual bias and temperature-dependence of the spin signal, and (iii) the unexpected short spin lifetime and its weak variation with temperature, require additional investigation. Here, we report the effect of spin relaxation rate on the interfacial spin depolarization (ISD) from the local fields in ferromagnet (FM)/oxide/SC contacts [1]. The combined measurements of normal and inverted Hanle effects reveal the effect of spin relaxation rate on the ISD [2] in CoFe/MgO/ Si and CoFe/MgO/Ge contacts. We have observed, despite the similar amplitudes of the interfacial roughness and local magnetic fields, significant differences of the ISD [2] depending on the host SC; the spin accumulation exposed to similar local fields in different SCs give rise to a clearly different ratio of the inverted Hanle signal to the normal one. This can be understood in terms of two competing mechanisms in the host SCs, namely the spin relaxation and spin precession due to the local fields.

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# **RL04**

#### Spin-pumping and revelation of inverse spin-Hall effect in n-type Si at room temperature

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Since spin-orbit interaction in n-type Si is very week due to its hand structure and lattice inversion symmetry, observation of the spin-Hall effect in n-type Si has been believed to be difficult. Here, we report successful observation of the inverse spin-Hall effect (ISHE) in n-type Si at room temperature. The sample composes of a NisoFeon/ n-type Si film with a doping concentration of 1.0×1019 1/cm3. The n-type Si surface was etched by hydrofluoric acid to remove the naturally-oxidized Si A ferromagnetic Ni<sub>80</sub>Fe<sub>20</sub> was formed by using electron beam lithography and deposition methods on the n-type Si. Two electrode-pads were attached to the n-type Si laver to detect the electromotive force in the Si. In a ferromagnetic resonance condition of the Ni<sub>10</sub>Fe<sub>30</sub> a pure spin current was injected into the n-type Si layer by the spin pumping. If ISHF is induced in the n-type Si, the spin current is converted to a charge current. Here, the ISHE in the n-Si was observed. The output voltage due to the ISHE was estimated to be 0.4µV. The voltage sign was inverted by the magnetization reversal of the Ni<sub>00</sub>Fe<sub>20</sub>. In this presentation, we discuss details of the inverse spin-Hall effect in the n-type Si

#### **RL05**

#### Critical current density and domain wall mobility in nanowires with exchange coupled hard-soft magnetic layers Xiaoxi Liu1\*, Liangqiu Gao2 and Akimitsu Morisako

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In this study, critical current density and domain wall mobility of current driven domain wall motion in nanowires composed of TbFeCo and NiFe layers have been studied. TbFeCo layer has perpendicular magnetic anisotropy. Wires and electrode are fabricated by photolithographic method. Current pulses with width less than 100 ns are applied to drive the domain wall in the wires. The domain structure is observed by a high resolution Kerr microscope. We found critical current density decreases in the soft-hard bi-layered structure compared to a single TbFeCo layer. Nanowires of single TbFeCo layer have critical current density of 5×106 A/cm2; however, the critical current density in bi-layered structure has critical current density as small as 2 ×106 A/cm2. The domain wall mobility is deduced from the dependence of domain wall velocity on the applied current density. We found that the domain wall mobility improved by almost two times with the introducing of the NiFe layer. Of particular interesting is that, micromagnetic simulations show that the Bloch walls in perpendicular TbFeCo wires change to Neel walls with the introducing of NiFe layer. Our experimental results suggest nano-wires composing of soft-hard layers suitable for race track memory.

#### **RL06**

#### Spin Seebeck Effect in SiO<sub>2</sub>/Py structures

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Recently spin version of Seebeck effect, the spin-Seebeck effects (SSE) was observed experimentally at room temperature[1]. For the SSE, it is expected that the voltage sign at the hot end is opposite to that at the cold end and the absolute magnitudes of the voltages at the hot and cold ends are the same. However, it is often observed that an asymmetric magnitude of the voltage at the hot and cold ends [2]. The origin of this asymmetric behavior is attributed to an additional temperature gradient along the thickness direction, which generates an additional voltage signal due to the anomalous Nernst-Ettingshausen effect(ANE) [2]. In this study, we investigated the dependence of the SSE and the ANE on the thickness of the ferromagnetic layer in SiO<sub>2</sub>/Py samples. In our results, the sign change in the voltage signal between the hot and cold ends was not observed. Instead, we obtained an offset voltage due to the mixture of SSE and ANE signals. We attribute this offset to the longitudinal SSE caused by an additional temperature gradient along the thickness direction of Py layer. A possible origin of the dependence of offset voltage on Py thickness will be discussed in detail

[1] K. Uchida et al., Nature, 455, 778-781 (2008). [2] S. Bosu et al., Phy. Rev. B 83, 224401 (2011).

# **RL07**

#### Domain wall pinning by stray field from NiFe on Co/Ni nano-wire Ryo Hiramatsu<sup>1</sup>, T. Koyama<sup>1</sup>, D. Chiba<sup>1</sup>, S. Fukami<sup>2</sup>, N. Ishiwata<sup>2</sup>, Y. Nakatani<sup>3</sup> and T.

#### Ono<sup>1</sup>\*

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Control of a magnetic domain wall (DW) displacement in a ferromagnetic nano-wire is essential in its potential application for nonvolatile magnetic memories. Parkin et al. demonstrated a race track memory device by controlling DW displacement in a NiFe nano-wire with notch structure [1]. However, the structure with notch has two problems such as i) transformation of DW structure and ii) increase of current density due to the change of the wire width between notch and wire part. Here, to make new DW pinning method without the change of wire width, we focus on the DW pinning in Co/Ni nano-wire by the stray field from ferromagnetic stack.We prepared two types of devices; one was a simple wire of Ta/Pt/[Co(0.2 nm)/Ni(0.6 nm)]4/Co(0.2 nm)/Pt/ Ta, and the other was a wire of Ta/Pt/[Co(0.2 nm)/Ni(0.6 nm)]<sub>4</sub>/Co(0.2 nm)/Pt/Ta on which NiFe/SiO2 was stacked partly. To investigate an influence of the stray field from the NiFe stack, DW depinning field ( $H_{dep}$ ) was measured for both types of wires.  $H_{dep}$ of the wire without the NiFe stack was 200 Oe, while the wire with that had larger  $H_{dep}$ of 700 Oe, indicating the effectiveness of the use of the stray field to control the DW position.

[1] S. S. P. Parkin, et al., Science 320, 190 (2008).

#### RL08

#### Compositional dependence of critical current density for spin transfer torque switching of amorphous GdFeCo for thermally assisted MRAM Bing Dai1\*, Takeshi Kato1, Satoshi Iwata1 and Shigeru Tsunashir

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Spin transfer torque (STT) switching has been demonstrated in magnetic tunneling junctions (MTJ) with perpendicular magnetic anisotropy (PMA), which exhibits a low critical current density compared to their in-plane counterparts [1]. In this paper, we fabricated giant magnetoresistance (GMR) devices consisting of bottom electrode/ [Pd (1.6)/Co (0.4)]<sub>6</sub>/Co<sub>40</sub>Fe<sub>40</sub>B<sub>20</sub> (0.5)/Cu (3)/ Gd, (Fe<sub>st</sub>Co<sub>10</sub>)100-x (10)/Cu (5)/top electrode (thickness in nanometer). The R-I loops were measured for 120×180 nm<sup>2</sup> cell with Gd<sub>x</sub>(Fe<sub>90</sub>Co<sub>10</sub>)<sub>100-x</sub> memory layers, and the STT switching property of these devices was estimated from the resistance measurement after applying current pulses with a duration of 100 ms. In TM-rich compositions, the average value of the critical current density Jc increased from  $2.0 \times 10^7$  A/cm<sup>2</sup> for x = 22.3 at% to  $4.5 \times 10^7$  A/ cm<sup>2</sup> for x = 24.0 at.% with increasing x, while the Jc reduced to  $1.4 \times 10^7$  A/cm<sup>2</sup> for RE-rich Gd<sub>289</sub>(Fe<sub>90</sub>Co<sub>10</sub>)71.1 It is noted that the values of Jc are comparable to the results in conventional MTJs[2]. The difference of Jc values may be due to the difference of the effective perpendicular anisotropy and/or the Curie temperature of the GdFeCo layer. The dependence of Jc on the current pulse width and temperature will be discussed in the presentation.

[1]S. Ikeda, et al. Nature Materials 9, 721-724 (2010). [2]Y. Huai, et al. Appl. Phys. Lett. 82, 222510 (2005).

#### **RL09**

#### Microscopic theory of magnon-drag thermodynamic transport in ferromagnetic metals

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We study the thermodynamic transport carried by magnons in ferromagnetic metals. In this study, the Peltier effect caused by a magnon heat current (magnon-drag Peltier effect) is described in a microscopic model. Magnons can indirectly interact with electric fields via electron-magnon interaction; we evaluate the magnon heat current perturbatively with respect to the electron-magnon interaction in the 2nd order. As a consequence, we found out that the magnon heat current is proportional to the spin current carried by the spin polarized electric current drifted by the electric filed, that is, the magnon-drag Peltier coefficient is governed by the spin polarization of the electric current. This is in contrast to the phenomenological result derived by Grannemann and Berger [1] in which the coefficient is proportional to the electric current directly. In addition, we show the temperature dependence of the Peltier coefficient is T5/2 in low temperature. These results indicate that, in the inverse effect (Seebeck effect) of the Peltier effect, the magnon heat current caused by a temperature difference induces the spin current, and the temperature dependence of the Seebeck coefficient is T<sup>3/</sup>

[1] G. N. Grannemann and L. Berger, Phys. Rev. B 13, 2072 (1976).

# **RL10**

Electric-field control of magnetic properties in cobalt by means of electric double layer

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Electric-field effect on magnetism using a capacitance structure (gate electrode/insulator/ferromagnetic) has actively investigated to open up a new technology for an electrically manipulation of magnetization [1-10]. Here, we show electric field control of magnetic properties in Co using an electrical double layer (EDL) formed at the interface between the Co layer and a polymer film containing an ionic liquid. The polymer film containing the ionic liquid was put on a 0.4-nm-thin-Co-film and Pt-thin-film was placed on the top of the polymer film as a gate electrode. For the direct detection of the magnetization under the gate-voltage (Vg), superconducting quantum interference device [6] was used. Magnetization curves were measured at 300 K under the application of Vg from -2.0 V to +2.0 V. Dramatic change of coercivity was observed depending on Vg. To investigate the mechanism of this, we measured temperature dependences of the remanent magnetization under the different Vg. The significant increase of Curie temperature from ~320 K at Vg = 0 V to ~370 K at 2.0 V was observed. Thus, the dramatic change in corecivity at 300 K presented here is attributed to this large modulation of Curie temperature [10].

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#### **RL11**

#### Spin-torque magnetic resonance of superparamagnetic Fe nanoparticles in Fe/MgO/Fe magnetic tunnel junctions

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During the past decade, electric detection and manipulation of submicron-scale magnet have been conducted by using magnetoresistance (MR) effect and spin-torque. Now, a great deal of interests in manipulating nano-scale magnets (or single-spin) and physics of spin-torque in such a small system are taken. In this study, we employ superparamagnetic Fe nano-particles and the single-crystal MgO based magnetic tunnel junctions (MTJs) [1]. The magnetic resonance of super-paramagnetic Fe nano-particles is electrically excited by spin-torque and detected by the spin-torque diode effect [2,3]. The structure of the MTJs is as follows: MgO(100) substrate/Fe film (50 nm)/MgO (1.35 nm)/Fe nano-particles (2.1 nm)/capping layer. All of the layers were grown by molecular beam epitaxy. The RF detection mechanism is explained as the homodyne detection of the applied RF current because of the oscillation of resistance at the same frequency. In the measurements, the resonant peak was observed around 5 GHz under a magnetic field of 5 kOe at room temperature, for example. This represents that the magnetic resonance of superparamagnetic Fe nano-particles was electrically excited and detected. We thank Y. Shiota, H. Tomita, G. Shiomi for valuable discussions. This work is supported by Grant-in-Aid for Scientific Research (S), MEXT, Japan. (No. 23226001)

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# **RL12**

# Spin Coulomb drag and optical excitations in low dimensional systems

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Within the remit of new quantum technologies, an intense effort is devoted to improving our understanding of spin dynamics, with the aim of building novel spintronics devices. In this context the theory of spin Coulomb drag (SCD) was recently developed. It shows that Coulomb interactions are an intrinsic decay mechanism for spin currents. As confirmed by experiments, SCD can be substantial in semiconductors and it is bound to become one of the most serious issues in spin polarized transport, since, due to its intrinsic nature, it cannot be avoided even in the purest material. More recently the influence of SCD on optical spin-injection and spinresolved optical experiments has been considered. Here we report on SCD effects on intersubband optical spin excitations in III-V quantum wells, where SCD may contribute substantially to the linewidth of spin plasmons. By going beyond the usual local density functional approximation and properly including the effects due to the inhomogeneity of the system in the growth direction, we show that the quantization of states in the growth direction may strongly reduce the intrinsic plasmon linewidth.

# RL13

# Perpendicular magnetic property and magnetic damping of very thin CoFeB films

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A very thin CoFeB film with a perpendicular magnetic anisotropy is a promising material for applications of spintronics devices, such as Spin-RAM and spin-logic circuit [1]. Development of materials with high magnetic anisotropy and low magnetic damping is an important issue to achieve high thermal stability and low power consumption in the spintronics devices. However, the relationship between magnetic damping and perpendicular anisotropy has not been systematically investigated. In this work, we have investigated capping layer dependence of perpendicular magnetic property and magnetic damping in very thin CoFeB films. The films were prepared by sputtering system. The structure of the films were SiO<sub>2</sub>-sub/Ta(5)/MgO(0.9)/CoFeB(d)/X(5 nm), (X = Ta, Pd, Cu, Ru, V, MgO). After the deposition of the films, annealing process was carried out at 250 - 350°C. Magnetic property and magnetic damping were respectively characterized by SOUID and FMR [2]. Perpendicular magnetic anisotropy and damping strongly depends on the capping layer material. For the Ta capping layer, CoFeB film below 1.4 nm showed a perpendicular anisotropy. However, magnetic damping was drastically enhanced for the films with perpendicular anisotropy. Similar behavior was observed in the films with capping layers of Pd and Ru. This work was supported by the FIRST program.

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# RL14

# Perpendicular magnetic tunnel junctions with TbFeCo-based pinned layer and CoFeB-MgO free layer

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We have fabricated perpendicular magnetic tunnel junctions (pMTJs) with a CoFeB/ TbFeCo pinned layer and Ta/ CoFeB/ MgO free layer using in-situ annealing process. The TbFeCo-based hybrid pinned layer is a useful approach to obtain a high perpendicular magnetic anisotropy (PMA). However, a high temperature annealing, often required for MgO-based MTJs to obtain a high tunneling magnetoresistance (TMR), may deteriorate the PMA of TbFeCo films. Our experimental results show a possibility to obtain a reasonable TMR and high PMA simultaneously by conducting an annealing process to achieve a grain-to-grain epitaxy in CoFeB/MgO/CoFeB layers before depositing the TbFeCo layer. The pMTJs prepared by this process show a tunneling magnetoresistance (TMR) ratio of 15% to 19% at room temperature.

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#### **RL15**

# Thickness and magnetic anisotropy dependence of anomalous Nernst effect in L1a-FePt films

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The anomalous Nernst effect (ANE) is a thermomagnetic effect which gives rise to an electric field perpendicular to both temperature gradient and saturation magnetization. Although ANE has been known phenomenologically for a long time, only a few studies have been reported, and the mechanism has not fully been understood. Ll<sub>0</sub>-FePt with high magnetic anisotropy (K<sub>u</sub>) attracted much attention for recent years as a promising material for "spin caloritronics" [1]. In this paper, we have made systematic investigation on the thickness and Ku dependence of ANE in perpendicularly magnetized FePt(001) films. FePt(001) films were deposited on MgO(001) single crystal substrates at 500°C. The thickness was varied in the range from 7 to 40 nm. Ku was changed in the range from  $3.0 \times 10^6$  to  $2.7 \times 10^7$  erg/cc, depending on the degree of L10 order which was controlled by the deposition temperature from 300 to 500°C. It is found that the anomalous Nernst coefficient (Q<sub>s</sub>) does not depend on the film thickness. On the other hand, Qs increases with decreasing K<sub>w</sub>. These results suggest a contribution of spin wave in ANE.

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### RL16

#### Spin-torque efficiency and magnetization reversal in three-dimensional Rashba materials

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In addition to the magnetization reversal by the spin-transfer torque, the magnetization reversal driven by a strong spin-orbit coupling has been intensely investigated experimentally and theoretically [1,2]. In this presentation, we focus on 3D Rashba models coupled with localized spins, and study the spin torque theoretically. This is motivated by the recent discovery of BiTel as 3D Rashba systems [3]. As a result, we find that the spin torque in 3D Rashba models is largely enhanced in the high-carrier-density regime compared with 2D Rashba models. We also find that the spin-torque efficiency defined as the ratio between the spin torque and the electric current is enhanced when the Fermi energy lies on only lower band, both in 3D and 2D. As we change the Rashba spin-orbit coupling, the spin-torque efficiency becomes maximum when the Rashba spin-orbit coupling is comparable to the exchange coupling to the localized spins. The optimum spin-torque efficiency becomes large when the Rashba spin-orbit coupling is large, and it is preferable for the magnetization reversal with smaller amount of current injection.

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#### RL17

#### A new circuit model for spin-torque oscillator including the perpendicular torque of magnetic tunnel junction

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Spin-torque oscillator (STO) is a promising new technology for the future RF oscillator, which is based on the spin transfer torque (STT) effect in magnetic multilayered nanostructure. It is expected to overcome the limitations of the semiconductor-based device technology. In our previous work, we have proposed a circuit-level model of the Giant Magneto-Resistance (GMR) STO. In this paper, we present a physics-based circuit-level model of the Magnetic Tunnel Junction (MTI)-based STO. The new model includes the effect of perpendicular torque that has been ignored in the GMR STO model. The variations of three major characteristics, generation frequency, mean oscillation power, and generation linewidth of an MTJ STO are investigated by changing the amount of perpendicular torque. The model is fully compatible with circuit-level simulators such as SPICE. The accuracy of new model was verified by HSPICE simulation. The simulation results show an excellent agreement with the experimental data. With our circuit model, we could compose a current mirror circuit that contains an STO element and a multi-stage amplifier for signal amplification, such that a full circuitlevel simulation of MIT STO was accomplished for the first time.

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#### **RL18**

# Current induced localized domain wall oscillators in NiFe /Cu /NiFe nano-stripe

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We study domain wall (DW) oscillations induced by the injection of a dc current through a nanowires containing artificial symmetric protrusions with Permalloy based pseudo spin valve (NiFe 12 nm/Cu 8 nm/NiFe 24 nm). An experimental analysis was carried out to describe the variation of differential resistance dV/dI by applying in-plane dc current with pinned DWs. When the DWs are pinned at the traps, we have observed a reversible motion of the DWs by changing the applied in-plane dc current. The peak in dV/dI associated with the reversible change of magnetoresistance has been taken as evidence of current induced coherent DW oscillation between the pinned regimes which gives rise to the peak. The excited critical current shifts symmetrically with respect to *H*. It is expected from the DW trap with symmetrical potential landscape. The dynamics of localized DW oscillations under application of dc current was examined by micromagnetic simulations (OOMMF). The frequency f of DW oscillation decreases from 1.5 GHz to 0.95 GHz as *H* increasing. This simple concept of DW based oscillators could be used as radio frequency assisted writing in magnetic recording technology.

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# RL19

# Current induced transverse field versus Joule heating in Co/Pd nanowires

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There has been a tremendous effort in understanding of the electric current effect in the ultra thin film of ferromagnetic materials[1, 2]. We have recently found that in the nanowire made of a multilayer structure of Co/Pd, there is a current induced transverse magnetic field[3]. Electric current could also cause Joule heating especially in the high current densities. The Hall resistance has been measured at different current densities for the magnetic fields in the easy axis direction. By increasing the current density, there is a reduction in the switching field of the nanowire. We further performed the current induced transverse field. Furthermore, by using four-probe measurements, the nanowire resistance has been measured at different temperatures and the temperature coefficient of the Co/Pd nanowire has been extracted. By measuring the ranswire te nanowire at different current densities, we can approximate the nanowire at different current densities, we can approximate the nanowire the sufficient of the correl at different temperatures and the temperature different at different current densities, we can approximate the nanowire temperature different current densities, we can approximate the nanowire temperature during the measurement at different current densities.

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#### RL20

#### Improvement of generation efficiency of pure spin current using multiterminal spin injection

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Lateral spin valve structures have several advantages for the generation, manipulation and injection of the spin current because of their flexible probe configurations. Particularly, nonlocal spin injection techniques enable us to create a pure spin current which carries the spin information with extremely low propagation loss. However, in general, the generation efficiency of the pure spin current is quite low. To solve this serious obstacle, we propose a nonlocal spin injection with multi spin injectors. In the present study, we fabricated the lateral valve structure consisting of a quadruple Permalloy spin injectors and nonmagnetic Cu strip and show that the developed device has two advantages. One is that the maximum generated spin current is strongly enhanced. The other one is that the efficiency for generating a certain quantity of the spin current is significantly improved. We also show that the optimization of the device dimension yields the further enhancement of the maximum generating spin current.

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### RL21

#### Electrical transport properties of Co2MnSi Schottky diode

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Recently, the spintronic device has attracted considerable attentions due to its possibility to lead a new generation of electronic devices by utilizing spin degree-of-freedom. Especially, silicon is a very attractive material for the spin channel of spintronic devices due to a weak spin-orbit coupling [1] and a long spin lifetime of ~10 ns [2, 3], while Co<sub>2</sub>MnSi being a half-metallic Heusler compound is regarded as a promising candidate for the spin injection/detection electrodes due to a possible complete spin polarization. [4, 5] In this work, the electonic transport chracteristics of Co<sub>2</sub>MnSi/Si Schottky diodes are studied for the investigation of spin electrode/channel interface quality, in which Co<sub>2</sub>MnSi decreased from 0.56 to 0.52 eV and 0.44 to 0.36 eV for electron and hole, respectively, with increasing RTA temperature from 450 to 550°C. Also, the ideality factor close to unity for RTA temperature of 450°C significantly degraded as it changed to 550°C. Also, the ideality is expected as a promising ferromagnetic material for the applications in silicon spintronic devices.

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# RL22

#### Negative electron-beam resist hard mask ion beam etching process for the fabrication of nanoscale spin transfer torque magnetic random access memory device

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To conduct MTJs etching process, Ta[1] hard masks are generally used for reactive ion etching and ion beam etching (IBE) process due to relatively good selectivity over magnetic metals. However, tapering of the Ta mask is unavoidable which results in the underlying MTJ to become trapezoidal. Here we utilized negative electron-beam resist (NER) hard mask process to overcome the hard mask induced nano-patterning difficulties. Arrays of 30 nm dot patterns 200 nm apart were defined by electron beam lithography with a NER thickness of 90 nm. The IBE was performed using the NER as the hard mask layer. During the IBE process, The IBE etch rate of the NER was 3 nm/min and since the total MTJ etching thickness was 36 nm with 2.1 nm/min etch rate, a 30 nm diameter NER pillar with 90 nm thickness was enough to produce the 30 nm diameter MTJ nanopillars. And redeposition of the etched material on side walls of the NER protects the resist from tapering, resulting in vertical side profiles in the etched MTJs.

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### **RL23**

# Magnetization reversal process of a Py nanodot under pure spin current injection

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A new type lateral spin valve structure based on a pair of closely located Permalloy nanopillars has been developed. This structure enables us to flow a large amount of current, resulting in the injection of a large pure spin current into a ferromagentic nanodot entirely. We investigated the magnetization reversal process of the Permalloy nanodit under dc pure spin current injection from the field dependence of the nonlocal spin valve signal. The dc pure spin current was generated by flowing the dc current with an ac excitation current in the nonlocal current terminal. The nonlocal spin valve signal without the dc current exhibits a clear bipolar signal typically observed in the conventional lateral spin valve.[1,2] The nonlocal signal curve shows the systematic changes with increasing the dc pure spin current and its variation strongly depends on the current intensity and polarity. This means that the reversal process of the Py dot is strongly modified by the pure spin current injection.

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### **RL24**

# Spin transfer effect in FePt nanowires: controlling the stochasticity of domain wall depinning using constrictions

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In this contribution, we present dynamic measurements of domain wall (DW) depinning from structural defects, in nanowires processed from single FePt layers and FePt/Pd(Pt)FePt spin-valves. Statistical analysis reveals that in such systems, the DW depinning is stochastic, and that this stochasticity has two origins. The first one is thermal activation, which helps the wall to cross the energy barrier associated to the defect. The second one originates from the ability of the DW to get pinned randomly along different micromagnetic configurations, each one corresponding to a given depinning fields. We study the influence of different constriction types on the dynamic of DW depinning, using spin-valve-based nanowires. The depinning probability function is found to exhibit very different behaviors, which depend on the types of constriction: it is partially possible to control the stochastic behavior of DW depinning can be varied by applying a DC current, the effect of the spin torque on thermal and configurational stochasticines being similar to those of a field. The measured spin-transfer torque efficiency is very high (~10-13T.m\_A<sup>-1</sup>), and is similar for single FePt layers and FePt-based spin-valves.

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# RL25

# Electric field effect on magnetic coercivity of $Fe_3O_4/BaTiO_3$ heterostructures

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Electric field control of the magnetic coercivity of half metallic Fe<sub>3</sub>O<sub>4</sub> epitaxial thin films on BaTiO<sub>3</sub>(001) is demonstrated at room temperature, with a view to developing a new type of highly spin polarized electron source, where the magnetization orientation of the spin source can be controlled by electric field. Fe<sub>3</sub>O<sub>4</sub> films were grown on BaTiO<sub>3</sub>(001) substrates by reactive molecular beam epitaxy in an oxygen atmosphere at a pressure of 1×10<sup>6</sup> Torr. The Fe<sub>3</sub>O<sub>4</sub> films show a clear Verwey transition at ~120 K, ensuing the high quality of the films grown. The electric field variation of the in-plane magnetic coercivity of Fe<sub>3</sub>O<sub>4</sub> films was measured using Kerr effect, exhibiting a hysteretic behavior, where the magnetic coercivities for positive and negative electric fields at ±10 kV/cm show distinct values. Since a possible strain effect due to ferroelectric domain switching of BaTiO<sub>3</sub> should be equivalent for both electric field polarities over the switching field, another possible mechanism such as polarization charge accumulation at the interface likely has its origin in the electric field effect, associated with carrier induced modulation of the electronic states of Fe<sub>3</sub>O<sub>4</sub>. Work supported in part by Industrial Technology Research Grant Program in 2009 from NEDO, Japan.

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### RL26

# Dependence of current-induced effective rashba field and perpendicular magnetic anisotropy on annealing temperature

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Manipulation of local magnetization using electrical current has attracted considerable interest due to its rich physics and potential applications for a new class of spintronic devices. It was theoretically proposed that Rashba-type spin-orbit coupling (SOC) caused by inversion symmetry breaking yields a new type of current-induced effective magnetic field [1], i.e. Rashba field of which direction is perpendicular to both the direction of current-flow and that of the inversion symmetry breaking. The existence of this Rashba field in the structure consisting of non-magnetic metal | oxide was recently confirmed experimentally [2]. We previously investigated the effect of the ferromagnetic film thickness on the Rashba spin transfer torque (STT), as well as its effect on the perpendicular magnetic anisotropy. The magnetic anisotropy clearly showed surface properties, while the surface nature of the Rashba field remained elusive. In this work, we further investigate the dependence of current-induced effective Rashba field and perpendicular magnetic anisotropy on annealing temperature. Experiments were carried out on Pt(3mn)/C0(0.6nm)/MgO(2nm) samples with the annealing temperature from 300°C up to 450 °C. In this presentation, the effect of annealing transment will be discussed.

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#### **RL27**

#### Spin motive force driven by magnetization dynamics Junichiro Ohe

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The current-induced magnetization dynamics realized in spintronics devices involve both of charge and spin degrees of freedom. Recently, it has been pointed out that the magnetization dynamics induces an effective electric field acting on the conduction electrons through the spin Berry phase. The effective electric field, or a 'spin motive electric field', was investigated for a simple one-dimensional domain wall. It is difficult to estimate analytically this effective electric field in actual systems, because the magnetization dynamics obeys the non-linear Landau-Lifshitz equation. In this report, we investigate the the spin motive force and the measurable voltage by solving Landau-Lifshitz-Gilbert equation numerically. We propose several structures that show measurable spin motive force. We also compare the our numerical results and experimental results.

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#### **RL28**

# Electrical detection of the spin Hall effects in the InAs quantum well structure

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Spin Hall effect (SHE) becomes intriguing because it generates and manipulates carrier spins in non-magnets without external magnetic field. Use of perpendicularly polarized spin currents for spin field effect transistor is a quite promising because precession of perpendicularly polarized spins can be manipulated by Rashba spin orbit coupling in the absence of the applied field along momentum direction [1]. The SHEs are electrically measured for an InAs quantum well (QW) structure where Rashba field of 8.5 Tesla is induced. The device is composed of an InAs channel and Pd/CoFe multilayers with perpendicular magnetization, which allows to detection of spin Hall voltages. The spin-polarized electrons with perpendicular orientation cause a charge accumulation induces a transverse voltage. We observe large spin Hall resistance up to 9.3 mΩ and spin Hall angle is found to be about 0.01. According to the theoretical reports [2, 3], one of its origins is spin-orbit scattering of conduction electrons through side-jump or skew scattering. The SHEs are emerging from the side-jump scattering of flowing spins in the InAs QW and the scattering direction is determined by spin orientation with spontaneous magnetization.

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#### RL29

#### **Bias-voltage controlled tunneling resistance in a ferromagnet-metalinsulator-ferromagnet tunneling junction** Sui-pin Chen\*

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We adopt the spin-polarized free-electron model [1-3] and extend our previous work [4] to study the tunneling resistance change in a ferromagnet/metal/insulator/ferromagnet magnetic tunneling junction [5], and find a method for changing the tunneling resistance. Unlike the traditional method of varying the magnetizations configuration between the two ferromagnetic layers, the proposed method uses the polarity of a bias voltage with a small strength to change the tunneling resistance. Under suitable conditions, we show that both tunneling resistance changes resulting from the polarity of the bias voltage and from the magnetizations configuration are equal in magnitude, and are larger than that in a conventional ferromagnet/insulator/ferromagnet magnetic tunneling junction.

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### RL30

# A highly (001) textured Ge/MgO/bcc-ferromagnet system prepared by ultra-high vacuum sputtering

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Observation of spin accumulation in semiconductor is very fascinating for fundamental physics and applications as well. Many important progresses have recently been made in this field. For example, the insertion of a tunnel barrier at the interface [1]. Ge/MgO/bco-FM has been intensively studied because of an epitaxial growth of MgO on a Ge substrated ue to small lattice mismatch (~4%) between Ge(001)[100] and MgO(001)[110]. As a result, the symmetry-related spin-dependent tunneling is expected to occur at the Ge/MgO/bco-FM junction. The successful growth of the single-crystal Ge(001)MgO(001)bco-FM(001) stack has been reported [2,3]. However, these results have been obtained using a molecular beam epitaxy technique which may not be suitable for applications in the industry. Here, we show that a highly (001) textured Ge/MgO/CoFeB junction can be achieved by an ultra-high vacuum sputtering technique. We will discuss the effect of various deposition conditions on the (001) texture of MgO/CoFeB layers will be presented based on the results of x-ray diffractometry and transmission electron microscopy. Finally, we will examine the tunneling is characteristics of the junction.

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# RL31

# Electrical spin injection and detection in GaAs with ferromagnetic metal/MgO junctions

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The intrinsic conductance mismatch between ferromagnetic metals (FM) and semiconductors is a critical obstacle in spin injection, and inserting a thin tunnel barrier, especially MgO, is expected to hurdle this problem. [1] In this study, we report electrical measurements of spin polarization in GaAs coming from efficient spin injection through MgO tunneling barrier. The whole epitaxial layers of FM/MgO/GaAs structures were prepared in-situ process using cluster MBE. To reduce the depletion layer of Schottky barrier, doping modulated GaAs layer has been inserted in the vicinity of the interface between GaAs and MgO layers. Large voltage drop ( $\Delta V = 6.3$  mV at 10 K when 1 = -0.9 mA in Fe/MgO/GaAs structure) between the baseline and the dip of Hanle curve is observed at various temperatures up to 400 K in a 3-terminal Hanle measurement geometry. The Lorentzian fitting of the Hanle curve gives 234 ps spin lifetime of n-doped GaAs and 952 nm spin diffusion length at 300 K.

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#### RL32

#### Transistorless 3D STT-MRAM Architecture Weizhong Wang\*

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The MTJ based devices are fabricated through thin film deposition and patterning which are intrinsically capable for 3D vertical multiple layers integration. However, current MRAM architecture requires at least one isolation transistor in each MRAM cell(1). Since there is no mature 3D CMOS technology available, isolation transistor prevent the progress on developing 3D MRAM. Therefore the races on MRAM area density improvements are all based on aggressive horizontal cell size scaling. This paper introduces a new transistor-less MRAM architecture to enable vertical 3-D integration. Since the MTJ devices are essentially two terminal resistors, an array of MTJ devices form a connected resistor network. We introduce a 3D architecture with new read/write operation to address individual MTJ cells. Comparing with conventional STT-MTJs, our preliminary simulation results show (1) little penalties on both speed and power consumption during read operation;(2) little penalty on writing speed; and (3) 2x to 3x penalty on writing power consumption depending on TMR ratio of MTJs. The proposed STT-MRAM architecture can greatly enhance the area density of STT-MRAM and beat CMOS or any single layer based memory technologies.

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#### RM01

# Collision of cobalt atom with Alq3 molecule thin film: a molecule dynamics study

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Yun-Peng Wang, Xiu-Feng Han and Ling-Ling Tao Beijing National Laboratory of Condensed Matter Physics, Institute of Physics, Chinese Academy of Sciences, Beijing 100190, China Hai-Ping Cheng Department of Physics and Quantum Theory Project, University of Florida, Gainesville, Florida 32611, USA Molecular dynamics simulations on deposition of cobalt atom onto tris(8-hydroxyquinoline) aluminum (Alq3) molecule thin film are presented. A Lennard-Jones potential calibrated by firstprinciple calculations was used to describe Alq3-Co interactions. Microscopic collision process was analyzed by looking into one trajectory. Statistics of final positions and penetration depths of cobalt atoms show that cobalt atoms with low initial energy penetrate little into Alq3 thin film and stay near the surface of Alq3 thin film while cobalt atom with high initial energy penetrates significantly into Alq3 thin film. Consequently, thickness of Alq3/cobalt interface is equal to roughness of Alq3 surface if initial energies of cobalt atoms are small, but 1 nanometer larger than roughness of Alq3 thin film if initial energies of cobalt atoms are large. Injected cobalt atoms lose their kinetic energies via collision with 8-hydroxyquinoline ligands of Alq3 molecules. DFT calculations show hybridization of carbon and oxygen due to cobalt and sizable charge transferred from cobalt to nearest Alq3 molecule

# RM02

# Magnetism and Electronic Structures of SiC nanoribbons: Role of defects

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Magnetism in diamagnetic materials such as defective SnO2, ZnO, or SiC got particular attention to probe and understand the source and nature of magnetism in these materials. Particularly, the role of intrinsic defects, e.g. a vacancy which is the main source of magnetism, motivated many scientists around the world. The defect formation energies of vacancies in these diamagnetic materials are very large which limit its experimental synthesis. Here, we focused on SiC nanoribbons (NRs) to probe the native defects and their interacations with the alien light element impurities, e.g., B, N, in terms of magnetism. Using ab-initio calculations based on density-functional theory (DFT), our extensive results indicate that silicon or carbon vacancy per cell induces magnetic moment in SiCNRs, whereas SiCNR with Si+Si divacancy has almost zero magnetic moment. The half-metallic behavior, which is very important for spintronics devices, of SiC NRs is maintained in the presence of one C or Si vacancy per cell. The Si or C vacancies have high defect formation energies, howevere, we found that foreing impurities B or N significantly reduced the defect formation energy of vacancies in SiCNRs.

### **RM03**

#### Bulk and surface half-metallic ferromagnetism in transition-metal chalcogenides with rocksalt phase from first-principles calculations Guoying Gao\* and Kailun Yao

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In the past decade, most of transition-metal pnictides and chalcogenides with metastable zincblende structure were predicted to be half-metallic ferromagnets. So far, however, only four of them, zinc-blende CrAs, MnAs, CrSb, and CrTe, have been realized experimentally in the form of thin films or multilayers because of the instability of the zinc-blende phase. Here we predict with first-principles calculations that all three rocksalt CrTe, VPo, and CrPo are half-metallic ferromagnets with the half-metallic gaps of 0.03, 0.18, and 0.31 eV, respectively. Importantly, the rocksalt phases are approximately 0.19-0.33 eV per formula unit lower in energy than the corresponding zinc-blende phases for these three half-metallic compounds, which indicate it is more possible to fabricate these binary compounds with rocksalt phase than the zinc-blende phase in the form of thin films or multilayers. We also study the (111) surface properties of rocksalt CrTe and CrPo, and find the Te- (Po-)terminated (111) surfaces preserve the bulk half-metallicity. The present work will not only add to the understanding of these binary half-metals, but will also stimulate experimental efforts toward fabrication and utilization of these systems in spintronics.

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### **RM04**

# Magneto-transport properties of Fe thin films in an external electric field

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- <sup>3</sup> Physics and Astronomy, Northwestern University, USA



### **RM05**

# Role of interfacial B impurity in magnetocrystalline anisotropy at MgO/Fe interface

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An interface induced-perpendicular magnetocrystalline anisotropy (MA), as observed in the MgO/CoFeB,[1] has attracted much attention in magnetic application. From first-principles, previously, we demonstrated a perpendicular MA at the MgO/Fe interface due to a weak Fe d\_z-O p\_z hybridization[2], but it still need to discuss a role of interfacial B impurity. Here, we extend investigation to include the interfacial B impurity. Calculations were carried out by using FLAPW method[2]. First, the phase stability at the MgO/Fe interface was determined based on calculated total energy with a model of an MgO monolayer on a five-layer Fe, where the B were incorporated in interstitial positions at top, second, and third Fe-interface-layers. Results indicate that the B favors to segregate to the interface so as to form an MgO/FeB/Fe interface. Then, the MA energy (E\_(MA)) was calculated by using the second variation SOC method and the force theorem. The calculated E\_(MA) has a large positive value of 0.69 eV for the ideal MgO/Fe, indicating a perpendicular MA. For the MgO/FeB/Fe, the E\_(MA) turns out to be almost zero (-0.02 eV), which indicates no interface MA, due to an enlargement of the Fe-MgO interfaver distance by the interface IB atoms.

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### RM06

# First-principles study on magnetic anisotropy of Co/Pt(111) film in electric field

Sho Yasuda\* and Shugo Suzuki

Division of Materials Science, Faculty of Pure and Applied Sciences, University of Tsukuba, Japan Electric field (EF) control of magnetism is attractive for spintronics devices. In recent research, it has been shown that EF affects not only magnetic moment but also

recent research, it has been shown that EF affects not only magnetic moment but also spin conduction and magnetic anisotropy[1-4]. In particular, it has been found that combination of the 3d and 5d metals in EF causes large effects on magnetic anisotropy. In this study, we investigate the dependence of the magnetic anisotropy energy (MAE) and the magnetic moment of Co/Pt(111) on EF by first-principles calculations. We employ the fully relativistic full-potential linear-combination-of-atomic-orbitals method based on the density-functional theory within the local-spin-density approximation. The exchange-correlation energy functional is the Perdew-Wang parameterization of the Ceperley-Alder results. We have found that the magnetic moment of the Co atom is in proportion to EF. The results of the calculations of the MAE and the other details will be shown at the conference.

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#### **RM07**

#### Magnetic ground state of TM/Graphene/TM films

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Using the full potential linearized augmented plane wave (FLAPW) method, we have investigated the magnetic properties of transition metal (TM)/Graphene/transition metal (TM) structure. We have explored magnetic properties of metal/grapheme/ metla systems, which is Co(1 ML)/Graphene/Ni(1 ML). Co(1ML)/Graphen/Ni(1ML) film have the AFM ground state. The calculated magnetic moment of Ni is -0.22 µB, while the Co layer has 1.50 uB. In calculation of optimized structure, we have that the distance between carbon and Co is 2 07 A and it is 2 04 A for carbon and Ni. The relative height difference between two carbon atoms in grapheme layer is 0.29 A. To understand the effect of buckling geometry of graphene, we have performed the calculation whithout any buckling factor of graphene layer. The magnetic ground state is not changed and calculated magnetic moment of Ni and Co are -0.60 and 1.70. respectively. This implies that the AFM coupling is an intrinsic property of Co(1ML)/ Gr/Ni(1ML) system. Also, we will present the magnetic properties of Co(2ML)/Gr/ Ni(1ML) and Co(1ML)/Gr/Ni(2ML) systems. This work was supported by KCAP located in Sogang University funded by the Ministry of Education. Science and Technology (MEST) through the National Research Foundation of Korea (NRF-2011-C1AAA001-2011-0030278) and by KOSEF (No. R01-2008-000-20014-0).

#### **RM08**

# A model of chain of ellipsoid-rings for magnetic nanotubes

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Numerical calculation of magnetic properties is a very effective way to understand the whole magnetic behavior of nanotubes. Currently, the most studies of calculation of magnetic properties of nanotubes are mainly grounded on the Stoner-Wohlfarth (S-W) model, starting from an elongated prolate ellipsoid with single-domain. But, it is hard to imagine how such an ellipsoid is arranged in the hollow tubular structure and hence the realization of predicted magnetic properties has been hindered by the experimental difficulties. In the present article, an alternative model of chain of ellipsoid-rings is proposed to calculate the magnetic properties of nanotube. Based on this new model, we calculate the magnetic properties of nanotube and further discuss the influence of tubular geometric parameters on the magnetic properties. All the results are well consistent with the experimental dato of Ni nanotube and moreover are available for the Ni nanowire. Consequently, our model provides an easy and general approach to both magnetic nanotubes and magnetic nanotube and moreover are available for the Ni nanotube and magnetic nanotubes and magnetic nanotube and superior and the superimental data of Ni nanotube and moreover are available for the Ni nanotubes and magnetic nanotubes

#### RM09

#### First-principles GGA+U calulations of half-metallicity in wurtzite NiO/ZnO(0001) superlattices X. H. Zhou

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Wurtzite NiO has been predicted to display half-metallicity and have only small lattice mismatch with wutzite ZnO [1]. It is interested to explore the possibility of the so-called all oxides spintronics obtained by the hybrid structrues consisted of the wurtize NiO and ZnO [2]. In this contribution, the structural, electronic, and magnetic properties of wurtzite NiO/ZnO(0001) superlattice were investigated by means of density functional calculations using spin-polarized generalized gradient approximation GGA and GGA+U schemes. We demonstrate that these superlattices retain their half-metallic behavior leading to a complete, i.e., 100%, spin polarization of the conduction electrons. This property makes the wurtzite NiO/ZnO(0001) heterostructures to be excellent candidates for high-efficiency magnetoelectronic devices.

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#### **RN01**

#### Magnetic properties of La doped nanocrystalline Z-type ferrite nanopowders synthesized via co-precipitation method

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Co2Z type hexagonal ferrites powders doped with La<sup>3+</sup> ions (Ba<sub>3</sub>-XLaXCo<sub>2</sub>Fe<sub>24</sub>O<sub>41</sub> where X=0.0, 0.05, 0.10 and 0.15) were prepared using co-precipitation method. The results revealed that single phase CO2Z type ferrite was formed from the precipitated precursor in presence of SDS as anionic surfactant and then preheated at 6000C for 4h and then post annealed at 1300oC for 6h. Moreover, the crystallite size and the porosity,% of the formed powders were increased whereas the unit cell volume were decreased by increasing the La<sup>3+</sup> substitution. The microstructure of the formed powders appeared as hexagonal platelet like structure. The DC resistivity of the obtained CO2Z was decreased as La<sup>3+</sup> content increased. The saturation magnetization (Ms= 53.7-55.5 emu/g) was slightly increased with substitution of La<sup>3+</sup>. Moreover, two resonance frequencies peaks are observed through the imaginary part of complex permeability meanwhile a decrease in the real part of magnetic permeability through X-band frequencies. The reasons are discussed using electromagnetic theory.

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# **RN02**

#### Magnetic properties of Fe-doped NiO nanoparticles

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NiO is known as a P-type semiconductor with wide band gap between 3.2 eV and 3.5 eV, and expected as a new material of diluted magnetic semiconductor. Magnetic properties of NiO nanoparticles with several nanometers were observed and blocking temperature(TB) was determined to be 17 K, while Neel temperature of NiO bulk crystal exhibits TN=523 K. In this study, Fe ions were doped into NiO nanoparticles in order to obtain higher blocking temperature. Ni1-xFexO (x=0, 0.05, 0.1) nanoparticles with several nanometers encapsulated with amorphous SiO2 were prepared by our novel preparation method. NiO single phase structure was confirmed from the x-ray diffraction measurements. It is considered that Ni ions are replaced to Fe ions exactly. This fact was also supported by X-ray absorption fine structure measurements. Magnetization measurements were performed by SQUID magnetometer for obtained samples. Temperature dependence of magnetization showed that TB shifted from 17 K to 57 K as the amount of Fe ions increased, and below TB, ferromagnetic behaviors were observed. Coercive force (HC) increased from 0.8 kOe to 1.5 kOe as the amount of Fe ions increased. This phenomenon can be explained by characteristics of extremely small particles in this system.

#### RN03

# Uniaxial strain effects on spinel ferrite nanoparticles containing Nd and B elements

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We studied the uniaxial strain effects on spinel ferrite nanoparticles that containined Nd and B elements (Nd : Fe : B = 10.4 : 83.5 : 6.1) and had the particle size of 5.1 nm. In diffused magnetic nanoparticles, the magnetic easy-axis is distributed randomly, and it is impossible to add the strain along a specific structural axis of each particle. However, in order to control the magnetic property of the magnetic nanoparticles, the application of the uniaxial strain for the particles with the uniform easy-axis vector is preferable to that of hydrostatic pressure for randaomly diffused particles. In this study, magnetic easy-axis of each particle was arranged by solidifying an epoxy resin containing the particle at the dc magnetic field of 1 T, resulting in creating an assembly of uniformly oriented nanoparticles. In both situations, the magnetic blocking temperature had the maximium at the strain of about 3-5 kbar. The difference between two strain experiments appeared at pressure of above 10 kbar.

#### RN04

# The magnetic proximity effect in $Fe_3O_4$ core/ $\gamma$ -Mn<sub>2</sub>O<sub>3</sub> shell nanoparticles

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The phenomenon of 'proximity effect' is very popular in the context of superconductor/ non-superconductor heterostructures. In the present study, we report the magnetic proximity effect in a core-shell nanoparticle system by employing dc magnetization and neutron diffraction techniques. In a heterostructure, the substrate, used for its deposition, often modifies the magnetic properties of the deposited layers. However, a core-shell nanoparticle system, being substrate-free, rules out any such possibility. We have observed magnetic proximity effect in a Fe<sub>3</sub>O<sub>4</sub> core/γ-Mn2O3 shell nanoparticle system, in terms of an enhancement of the Curie temperature (Tc) of the  $\gamma$ -Mn<sub>2</sub>O<sub>3</sub> shell (~66K) compared to its bulk value (~40K), and the presence of magnetic ordering in its so-called paramagnetic region (i.e. above 66K). The origin of these two features has been ascribed to the proximity of the  $\gamma$ -Mn<sub>2</sub>O<sub>3</sub> shell with a high-Tc Fe<sub>3</sub>O<sub>4</sub> core (~858K in bulk form) and an interface exchange coupling between core and shell. Interestingly, no exchange bias effect is found due to a weak interface exchange coupling between core and shell. The present study brings out the importance of the relative strength of interface coupling in governing the simultaneous occurrence of the magnetic proximity effect and the exchange bias phenomenon in a single system

# **RN05**

#### **Covalent immobilization of biotin on superparamagnetic nanoparticles** Long Giang Bach<sup>1</sup>, Md. Rafiqul Islam<sup>2</sup> and Kwon Taek Lim<sup>2</sup>\*

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Among the magnetic materials, the nano-scale structured magnetite, Fe3O4 is widely applicable magnetic nanparticles (MNPs). MNPs offer attractive possibilities in biomedicine, biosensor, and separation processes because of their unique properties of superparamagnetism, biocompatibility and low toxicity. The ease of synthesis and subsequent surface functionalization of MNPs with a desired material improve biocompatibility which creates an avenue of delivering multimodal and multifunctional MNPs. In this study, a simple protocol of covalent immobilization of biotin on the surface of MNPs for improving biocompatibility and cell viability of the original MNPs has been demonstrated. MNPs were prepared using a solution of ferrous and ferric ions. The anchoring of 3-aminopropyltrimethoxysilane with MNPs resulted due to the formation of Fe-O-Si covalent linkage which is quite strong and having free amino group. The hydroxyl groups of the biotin were activated by N,N'-disuccinimidyl carbonate followed by conjugation with MNPs-NH2 to afford Biotin. MNPs. The conjugation of biotin to MNPs was confirmed by FT-IR, XPS and EDX. The Biotin-MNPs showed superparamagnetic character as investigated by SQUID.

### RN06

# Fabrication of (Mn-Al)/Pd or Ni with Core-Satellite Structured Magnetic Particles

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### **RN07**

# The effect of hydrostatic pressure on the Morin transition in hematite nanoparticles

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Recently, hematite particles have been the subject of intensive studies in areas as diverse as magnetic storage media, photo electrodes and catalysis[1,2]. Hematite is an antiferromagnetic system that displays a spin reorientation from uniaxial to canted antiferromagnetism at the Morin temperature[3]. Particle shape, size and microstructural defects have strong influence on the Morin temperature and thus can be used to tune the system's magnetic response. It is known that the Morin transition temperature and thus can be used to tune the system's observed in the bulk material[5]. In a novel approach, we report on the effect of hydrostatic pressure son hematite rhombohedral nanoparticles of narrow size distribution<DTEM>=93 ± 2 nm[6]. Pressure was found to increase the Morin temperature and there an inder effects of pressure on both the Morin temperature and thermal hysteresis. These results strongly contrast with the effects of pressure on the bulk material where a milder influence of pressure on both the Morin temperature and thermal hysteresis is observed.

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### **RN08**

#### Ultra-thin MgO coating of superparamagnetic magnetite nanoparticles by combined co-precipitation and sol-gel synthesis

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Superparamagnetic magnetite nanoparticles coated with an ultra-thin (~1 nm) magnesium oxide (MgO) layer have been synthesized by combining co-precipitation and sol-gel methods. A thorough chemical and structural characterization has been carried out by means of HRTEM, XRD, EDS, DLS and TGA. Aberration corrected HRTEM experiments with sub-angstrom spatial resolution have allowed us to distinguish the ultrathin MgO shell that grows epitaxially on the magnetic cores. The capability of the MgO shell to protect the magnetic nuclei from oxidation up to 600°C has been demonstrated. The magnetic properties of the material have been studied before and after the coating procedure. The superparamagnetism of the magnetic nuclei at room temperature is preserved even after calcination. The possibility of obtaining particles of controlled size coated with an isolating layer of nanometric thickness and high thermal stability makes the combination of the two synthesis methods used in this work a starting procedure to obtain nanometric powders suitable for technological applications, such as high-frequency electronics. These particles, after the appropriated functionalization, are also potential candidates to be used in biomedical application

### **RN09**

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#### Synthesis and characterization of ultra-small magnetic FeNi/G and NiCo/G nanoparticles

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Ultra-small magnetic nanoparticles consisting of NiCo and FeNi alloys coated by graphitic shells (NiCo/G and FeNi/G) have been synthesized by a procedure similar to that described by Seo et al. [1]. The cores, which retained the face-centered cubic (fcc) symmetry of the original bulk metals, together with the graphitic coating have been characterized by means of aberration corrected scanning transmission electron microscopy (STEM) The particles have mean sizes of 2.6 nm and 6.2 nm for NiCo/G and FeNi/G, respectively. The samples contain a large fraction of superparamagnetic nanoparticles at room temperature. The thermal stability of the particles is enhanced by the graphite shell, and the graphite coated FeNi and NiCo are stable under oxygen atmosphere up to 170 C. Non-coated bimetallic FeNi and NiCo have also been prepared. Their chemical characterization revealed the massive oxidation of the bimetallic nanoparticles in absence of the graphite coating (oxygen content higher that 60 at. %). This fact confirms the effectiveness of the graphite shell in preventing the magnetic cores oxidation and therefore preserving their magnetic properties even at high temepratures

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#### **RN10**

#### Magnetization measurements and blocking temperature distribution in magnetic nanoparticle systems

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We present magnetization measurements on a variety of magnetic nanoparticle ensembles. The samples are; (1) Artificailly synthesized Helicobacter Pylori ferritin reconstituted with three different Fe ions/ferritin ratio, 160, 2000, and 4000. (2) Commercial horse spleen ferritin (3) CoPt nanoparticles with average particle size of 2 nm. Temperature dependence of magnetization has been measured using a SQUID magnetometer. Measurement sequence has been modified from a typical field-cool sequence and the measurements have been performed with field "stop-resume" steps at the regular temperature interval. Using a phenomenological model, we could extract the moment weighted blocking temperature distributions in the samples from the experimental data. The model is based on the assumptions that; 1) the particles in the ensembles can be uniquely categorized in terms of the blocking temperature. 2) the particles with the same blocking temperature will have the same field and temperature dependence of the magnetization. 3) the particles are non-interacting with each other. With the obtained distributions, the numerical simulations are performed to reproduce the field-cool magnetization data and zero-field-cool magnetization data.

# 298 | The 19<sup>th</sup> International Conference on Magnetism

### **RN11**

### Structural and magnetic anomalies in CoAl<sub>2</sub>O<sub>4</sub> nanocrystals

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In this study, we investigated magnetic properties of spinel CoAl2O4 nanocrystals (NCs) synthesized by a supercritical hydrothermal synthesis. We confirmed that all the samples are composed of the single phase CoAl<sub>2</sub>O<sub>4</sub> by the X-ray diffraction measurements. Crystalline sizes are controlled to be 4 ? 8 nm. We also found that the inversion parameters of the samples, describing the cation distribution are more than 0.3 [1,2], being much higher than bulk specimens previously reported. We conducted magnetization and, dc- and acsusceptibility measurements. All the samples show paramagnetic behavior above 70 K. The effective Bohr magnetons and the Weiss temperatures are 2.6 - 2.8 µB and -40 - -50 K. The effective Bohr magnetons are smaller than 4.4 - 4.6 µB reported on specimens with the low inversion parameters [3,4]. The susceptibilities deviate from Curie-Weiss law below 40 - 50 K. From the results of the magnetization measurements, it was found that spontaneous magnetization, 0.1 µB/Co, emerges below 40 K in the samples with the size of 4 nm. The reduced Bohr magnetons and the emergence of the spontaneous magnetization indicate that in our NC samples, a certain magnetic phase, being different from those of bulk specimens, appears below 40 K.

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#### **RN12**

#### Preparation of chains of single-domain Ni nanoparticles with collinear direction of magnetization

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Networks of ferromagnetic nanoparticles have been proposed to form quantum cellular automata [1]. We present a fabrication method based on self-organized growth on stepped surfaces for fairly regular nanoparticles arranged in a one-dimensional chain, that are suitable for this application. We studied Ni nanoparticles with sizes in the range 80-130 nm on a highly-ordered pyrolytic graphite (HOPG) substrate and succeeded to prepare one-dimensional chains of these single-domain particles characterized by the same direction of their magnetization. The nanoparticles were found to decorate atomic cleavage steps on the basal surface and their distances are small enough to give rise to strong dipolar coupling. The necessary and sufficient conditions are: (i) 3D growth mechanism along with non-wetting behaviour, (ii) the dielectric substrate has atomic cleavage steps, (iii) the size of the ferromagnetic particles is smaller than the characteristic domain size but still above the transition into the superparamagnetic state and, finally, (iv) the distances between the particles are small enough for dipolar coupling. The study was carried out using atomic-force (AFM) and magnetic-force (MFM) microscopy. The approach of ordering ferromagnetic nanoparticles into one-dimensional chains is of particular interest because chains are a basic block for storage and transport of information.

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#### **RN13**

#### Influence of the morphology on the magnetic properties of dodecanethiol-capped Au anoparticles

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It has been proved that Au nanoparticles (NP) coated with certain molecules can present different magnetic characters, ranging from superparamagnetism to permanent magnetism at room temperature (RT) [1-3]. We have prepared ferromagnetic dodecanethiol-capped Au NPs with a high magnetic signal of 1.7 emu/gr and coercive field of 85 Oe at RT. An analysis by high resolution electron microscopy reveals that the NPs are composed by ~ 200 atoms forming multi-twinned structured nanocrystals of ~2 nm with icosahedral or decahedral morphologies where well ordered fcc twinned entities are joined forming an only nanocrystal with a spatial 5-fold symmetry. The external morphology is conditioned by the final stabilization of the NP and, with a surface to bulk atoms ratio up to ~75% and a magnetism located exclusively in the surface atoms, the morphology plays a crucial role in the definition of the magnetic character of the NP. In a comparative view with the multidomain structures presented in bulk materials, we argue that borders among the twinned nanosized entities would act as pinning centers, hindering the alignment of the magnetic moments between entities and promoting the coercive field.

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# **RN14**

#### Structure effects on the magnetic behavior in Fe oxide based nanonarticle

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A key question in magnetic nanoparticles (NP) is how the nanostructure modifies their magnetic and electronic properties. As the particle size is reduced, deviations from bulk behaviour have been widely reported and attributed to both the appearance of a magnetically disordered surface layer and finite-size effects, thus giving rise to a decrease in the particle magnetization, an exchange field acting on the core and an increase in the anisotropy. In order to elucidate those phenomena we have carried out STEM, XAS, XMCD and magnetization measurements in Fe oxide based NP models, aiming at the real-space characterization at the sub-nanometer scale of a single NP to characterize the particle structure. surface magnetization, and the dependence of the latter on the strength of the surface bond to a bio-molecule coating. We also aim at the quantification of the orbital and spin contributions to the NP magnetic moment. The key role of the crystal quality is thus suggested because particlelike behaviour above about 5 nm in size is observed only when NP are structurally defective. These conclusions are supported by Monte Carlo simulations. It is also shown that thermal decomposition is a chemical route capable of producing high-crystal quality NP with bulklike properties.

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#### **RN15**

#### Magnetic properties of surface-functionalized nano-particles

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#### **RN16**

#### Study of aqueous dispersions of magnetic nanoparticles by magnetic and rheological measurements

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The observed magnetic tunability of light transmission through a ferrofluid [1] can be effectively understood in terms of the inter-particle interaction that can be estimated from the magnetic and rheological properties of these fluids. The present study reports complementary magnetic and rheological measurements of synthesized aqueous dispersions of ferrite nanoparticles, using a commercial ferrofluid as a standard. The aqueous dispersions of the ferrite nanoparticles were synthesized by co-precipitation using suitable precursors and surfactants. The dispersions were stable for reasonably long durations and the particle size was estimated to be 10-28 nm from XRD measurements. The room temperature magnetization measured in a SQUID magnetometer up to fields of 1 Tesla showed superparamagnetic behaviour of the powder and the dispersion with the background signal of the liquid varying from a diamagnetic to a paramagnetic behaviour. The room temperature rheological behaviour in zero magnetic field of both the standard and synthesized dispersion was investigated by measuring the viscosity as a function of shear rate from 1-1000s-1. The viscosity of the dispersion increased with increasing shear rate reaching a constant value at higher shear rates; while the standard sample showed almost no change in viscosity with shear rate conforming to a Newtonian behaviour.

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# **RN17**

#### Magnetodielectric effect of Fe<sub>2</sub>O<sub>3</sub> nanoparticles embedded in SiO<sub>2</sub> glass matrix

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Magnetic iron-oxide (doping 0.5 mole %) nanoparticles (NPs) have been synthesized in silica glass matrix, prepared by the sol-gel method with different annealing temperatures. XRD, TEM and XANES show the presence of NPs of  $Fe_2O_3$  and  $Fe_3O_4$ crystalline phase (< 5 nm). An interesting colossal enhancement of dielectric constant with diffuse phase transition is observed around room temperature due to the thermally activated oxygen vacancies. The magnetodielectric effect observed in the glass composites is considered to be associated with magnetoresistance changes, depending on NPs size

# **RN18**

#### Tuning the concentration of magnetic Co nanoparticles in In<sub>2</sub>O<sub>3</sub> with oxygen pressure and concentration of tin

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This study shows the power of combining magnetic measurements with magnetic circular dichroism (MCD) to detect the contribution to the magnetism from metallic Co nanoparticles, and the result that the formation of nanoparticles is suppressed by the inclusion of tin. Thin films of  $(In_{1,0.95}Co_{0.05})_2O_3$  were deposited using pulsed laser deposition (PLD) on sapphire substrates. The magnetism and the amount of cobalt in nanoparticles increased with the oxygen deficiency in the PLD chamber and also after annealing in vacuum as expected and as confirmed using Extended X-ray Absorption Fine Structure. The MCD spectra show two features that are well fitted by a combination of the spectrum from the nanoparticles in the In2O3 matrix, as calculated using Maxwell-Garnett theory, and a contribution from polarized carriers [1]. The changes in the magnetization and the MCD are used to get a quantitative estimate of the contribution of the nanoparticles magnetization to the total magnetization in each sample. Adding 5% of Sn made little difference to the saturation magnetization but the signal due to Co nanoparticles was completely suppressed and that from polarized carriers enhanced strongly. This indicates that, the magnetization of the carriers in Sn and Co co doped In<sub>2</sub>O<sub>3</sub> is particularly high.

Hakimi, et al. Physical Review B 84, 085201 (2011)

### **RN19**

#### Synthesis, structural, and magnetic properties of strontrium hexaferrite nanoparticles with La, Sm doping and core/shell structure by the sol-gel hydrothermal process

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SrFe12O19 nanoparticles have been prepared by the sol-gel hydrothermal method. These SrFe<sub>12</sub>O<sub>19</sub> nanoparticles have been doped with La, Sm and coated with NiFe<sub>2</sub>O<sub>4</sub>, CoFe<sub>2</sub>O<sub>4</sub> nanoparticles by the sol-gel and the self-assembly method, respectively. Properties of the nanoparticle samples were charaterized by powder X-ray diffraction with the Rietveld refinement, transmission electron microscope, scanning electron microscopy with energy dispersive spectrometer, and vibrating sample magnetometry. SEM micrograhs of the samples showed the spherical shape of NiFe<sub>2</sub>O<sub>4</sub>, CoFe<sub>2</sub>O<sub>4</sub> nanoparticles located above on the surface of SrFe<sub>12</sub>O<sub>19</sub> hexagonal platelet shape. TEM micrographs showed the particle size of core and shell. With the core/shell nanoparticles, the saturation magnetization was decreased but coercivity was increased with SrFe., O., /CoFe,O. sample.

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### **RN20**

# A facial fabrication of the superparamagnetic $Fe_3O_4$ (2) $TiO_2$ microspheres and its photocatalytic application

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Over the past few decades, semiconductor photocatalysts have attracted much attention because of their applications in new energy fields where they represent a promising strategy for splitting water to provide clean hydrogen energy, and their use in the photocatalytic degradation of polluted atmosphere and wastewater in environmental remediation[1,2]. Among the reported photocatalysts, ittanium dioxide (TiO<sup>2</sup>) is one of the most studied semiconductors for photocatalytic neactions because of its high activity, chemical stability, robustness against photocorrosion, low toxicity and availability at low cost, especially for the detoxification of water and air [3,4]. In this study, we have successfully synthesized of multifunctional core/shell (Fe<sub>2</sub>O<sub>4</sub>/TiO<sub>2</sub>) submicron particles with controlled shell thickness. In the first step, superparamagnetic Fe<sub>2</sub>O<sub>4</sub> submicron particles was coated with TiO<sub>2</sub> by atomic layer deposition (ALD) method. Through this method, the thickness of shell was accurately tunable. The recyclable Fe<sub>2</sub>O<sub>4</sub>@TiO<sub>2</sub> particles exhibit excellent catalytic properties for the oxidation reaction of 2.4.6-trichlorophenol (2.4.6-TCP) in aqueous solution.

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# RN21

# Synthesis of high moment magnetite ( $Fe_3O_4$ ) nanoparticles by simple modified polyol method

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(Synthesis of high moment magnetite (Fe<sub>3</sub>O<sub>4</sub>) nanoparticles by simple modified polyol method) Mohamed Abbas and CheolGi Kim Center for NanoBioEngineering and Spintronics, Department of Materials Science and Engineering, Chungnam National University, Daejeon 305-764, Korea Abstract Single phase magnetite with high saturation magnetization and 50 nm as a medium range particles size has been synthesized by in expensive and surfactant-free simple modified polyol method using FeCl<sub>2</sub>.4H<sub>2</sub>O as a precursor of iron and Poly Ethylene Glycol (PEG) as a solvent and reducing agent simultaneously. X-ray diffraction (XRD) and EDS analysis indicated that, formation of purely magnetite (Fe<sub>3</sub>O<sub>4</sub>) without presence of any other phase. Transmission electron microscopy (TEM) also show that the synthesized nanoparticles are almost sphere in shape. The magnetic properties of the synthesized magnetite nanoparticles were measured by means Physical Property Measurement System (PPMS-VSM) at different temperature and shows that the as-synthesized NPs were ferromagnetic with high moment 91.7 emu/g at 5 K. Corresponding author. Tel:+ 82428216236 E-mail address: cgkim@cnu.ac.kr (Prof/ CheolGi Kim)

### **RN22**

#### Study of Structural, morphological and optical properties of $SrFe_{12-x}$ Co<sub>x</sub> O<sub>19</sub> (x= 0, 0.1, 0.2) hexaferrite nanoparticles

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A series of M-type strontium hexaferrite samples having nominal composition SrFe<sub>12x</sub> Co<sub>x</sub> O<sub>19</sub> (x = 0.0, 0.1 and 0.2) with different morphologies have been synthesized via sol gel method. For synthesis of samples, first a precursor gel was prepared using the sol-gel method. The mixture of KCl and dry-gel was calcined at 1000 °C for 2 h to obtain the needle-like nano- SrFe<sub>12x</sub> Co<sub>x</sub> O<sub>19</sub>. The rod-like and bubble nano- SrFe<sub>12x</sub> Co<sub>x</sub> O<sub>19</sub> was prepared by calcining the mixture of KBr and dry-gel or KI and dry-gel, respectively, and spherical nano- SrFe<sub>12x</sub> Co<sub>x</sub> O<sub>10</sub> was prepared only by calcined dry-gel. The structure and morphology of samples were systematically characterized by X-ray diffraction (XRD), scanning electron microscopy (SEM). The optical properties of samples were investigated by UV-Vis. The result of XRD of samples showed that the M-type hexagonal structure can be maintained for all the samples without the segregation of secondary phases. The FT- IR frequency bands in the range 560 - 580 cm<sup>-1</sup> and 430 - 470 cm<sup>-1</sup> in the all of samples corresponds to the formation of tetrahedral and octahedral clusters of metal oxides in ferrites, respectively.

### **RN23**

#### Light induced ferromagnetism of nanocrystalline CuCr<sub>2</sub>Se<sub>4</sub> particles Dongsoo Kim\*, Kookchae Chung and Choljin Choi Korea Institute of Materials Science, Korea

A novel method to synthesize ternary chalcogenide nanoparticles was developed with microwave-assisted polyol reaction in this study. Copper chloride, chromium acetate and metallic selenium were used as precursors with stoichiometric ratios in poly ethylene glycol. The system was purged with argon and then the mixed reactants were irradiated for 60 minutes with 60% of the instrument power. Products were collected, centrifuged, washed with ethyl alcohol in order to remove poly ethylene glycol and dried overnight under inert atmosphere. Magnetic field of 500 Oe in argon atmosphere and Curie temperature was calculated as 445 K. From the hysteresis loop, magnetic property was ferromagnetic and its saturation magnetization and coercive force were measured to 15 emu/g and 80 Gauss respectively. Nanocrystalline CuCr\_Se<sub>4</sub> particles were dispersed in the matrix of light transparent resin so as to investigate light induced effect on the magnetic properties. Er-laser with the wavelength of 1.5  $\mu$ m was employed as light source. Light illumination in an external magnetic field increased the magnetization and this increment of magnetization ( $\delta$ I) to initial magnetic icol (J) showed a maximum.

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#### **RN24**

#### Magnetite nanoparticles in hybrid aerogel and PEG encapsulated of magnetite nanoparticles for the hyperthermia application Eunhee Lee and Chang-yeoul Kim *KICET. Korea*

Magnetite nanoparticles have received intensive attractions because of their applications as magnetic carriers in drug targeting and hyperthermia in tumor treatment. However, magnetite nanoparticles with 10 nm in diameter easily aggregate themselves to form large secondary particles. To disperse magnetite nanoparticles, this research paper suggests the infiltration of magnetite nanoparticles into hybrid silica aerogels. The PEG encapsulate of magnetite is to disperse magnetite up in PBS solution. The well-dispersion of magnetite is necessary for target the tumor cells and hyperthermia treatment. Methyltriethoxysilicate (MTEOS) based hybrid silica aerogels were synthesized by a supercritical drying method. To incorporate magnetite nanoparticles into the hybrid silica aerogels, we immersed it into magnetite precursor solution. The surfaces of magnetite nanoparticles were modified by PEG(polyethylene glycol). The infiltration and PEG coating of magnetite nanoparticles were identified by XRD, FT-IR and TEM. Vibrating Sample Magnetometer analysis showed that the composite of the magnetite and hybrid silica aerogel is superparamagnetic. The temperature behaviors of the magnetite composite and the surface coated magnetite nanoparticles were discussed in the context of the change of magnetite particles with PBS for Hyperthermia application.

#### RO01

# Effect of the Cu content on the microstructural and magnetic properties of Nd-Fe-B sintered magnets

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The Cu addition in the Nd-Fe-B sintered magnet improves the coercivity after the post-sintering annealing (PSA) because the continuous Nd-rich grain boundary phase (GBP) and the metastable C-Nd<sub>2</sub>O<sub>3</sub> phase were formed [1-3]. However, the magnetic properties of the commercial magnet were deteriorated when the Cu content of the magnet is higher than 0.15 at.% (optimized Cu content). In this study, we investigated the Cu content effects on the microstructural and magnetic properties of sintered 12Nd-2.7Dy-(76.45-x)Fe-xCu-GB-2.65M (at.%, x=0.2, 0.3, 0.4, 0.5) magnets. The coercivity was decreased (28.7–26.1 kOe) with increasing the Cu content, but the remanence was unchanged. Two kinds of Cu enriched Nd-rich triple junction phases (Cu-rich TJP), such as 5~10 at.% Cu containing Nd-rich TJP (Culow-rich) and 35-45 at.% Cu containing Nd-rich TJP (Culigh-rich), were observed. The Culow-rich TJP was the stable h-Nd<sub>2</sub>O<sub>3</sub>, but the Culigh-rich TJP was the metastable C-Nd<sub>2</sub>O<sub>3</sub> phase. The thickness and Fe content of Nd-rich GBP was increased with increasing the GBP/Nd<sub>2</sub>Fe<sub>44</sub>B interface. As a result, the magnetic decoupling through the Nd-rich GBP was deteriorated and the coercivity was decreased when the Cu content of the magnet is high.

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# RO02

# Effects of the D<sub>y</sub>H<sub>3</sub> and D<sub>y2</sub>O<sub>3</sub> powder addition on the magnetic and microstructural properties of Nd-Fe-B sintered magnet

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Effective way for enhancing the magnetic properties as well as saving the Dy content in the Nd-Fe-B based magnet is to develop the core-shell microstructure using the Dy<sub>2</sub>O<sub>3</sub>, DyF<sub>3</sub> or DyH3 powder doping [1, 2]. The DyH<sub>3</sub>-doped magnet shows the most improved results [2]. However, the role of DyH<sub>3</sub> powder has not been identified yet. In this study, we investigated the magnetic and microstructural changes of the DyH<sub>2</sub>-doped (2.4 wt%) Nd<sub>2760</sub>Dy<sub>489</sub>FebalB<sub>10</sub>M<sub>24</sub>(wt%) sintered magnet comparing with those of the Dy<sub>20</sub>-doped (2.0 wt%) magnet. The coercivity of Dy<sub>20</sub>-doped magnet slightly increased (30.13→30.71 kOe) with a little remanence reduction (11.59→11.56 kG) comparing with the undoped magnet. However, the coercivity of the DyH<sub>3</sub>-doped magnet. The hydride powder doping enhanced the diffusion of Dy because the dissolution of the hydrogen in the Nd<sub>2</sub>Fe<sub>14</sub>B increased the lattice parameter and the point defects. The hydride was also known as a solid lubricant role at the grain boundaries [3]. As a result, the core-shell microstructure magnet.

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#### RO03

# Magnetic properties of nano-composite Nd-Fe-B thick-film magnets prepared by vacuum arc deposition

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In order to upgrade the torque of the small motor comprising a multi-polarly magnetized rotor with isotropic thick-film magnets [1], enhancement in remanence of the isotropic thick-film is indispensable. We, therefore, prepared nano-composite Nd-Fe-B/Alpha-Fe isotropic thick-films by using a vacuum arc deposition method [2]. This contribution reports the effect of the composition and the power of arc deposition on the magnetic properties. Samples were prepared by using several targets with the compositions of NdXFe<sub>14</sub>B + Nb<sub>0.5</sub> at.% (X =1.4-2.0). The voltage and capacitance varied in order to control the power range between 11 and 44 J. As the power was fixed at 44 J, use of a Nd<sub>1.6</sub>Fe<sub>14</sub>B + Nb<sub>0.5</sub> at.% target enabled us to obtain the highest (BH)max value. Although the deposition rate decreased from approximately 20 to 5 microns per hour due to the power reduction from 44 to 11 J, the value of (BH)max could be improved by approximately 15 kJ/m3 at the target composition of Nd<sub>1.6</sub>Fe<sub>14</sub>B + Nb<sub>0.5</sub> at.% Resultantly, we succeeded in obtaining a 29 micron-thick Nd-Fe-B thick-films with the remanence, coercivity and (BH)max values of 0.92 T, 390 kA/m and 79 kJ/m3.

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### RO04

# Corrosion resistance and corrosion behaviors of sintered rare-earth magnets in different corrosive environments

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Sintered rare-earth magnets possessing outstanding magnetic properties are widely applied in many fields. However, corrosion resistance has been a problem with the magnets owing to the high chemical activity of the phases rich in rare earth elements [1]. The Nd-rich phases of the Nd-Fe-B magnets can react with water vapor and oxygen in a damp and humid atmosphere. The electro-chemical corrosion occurs for Nd-Fe-B magnets in a high humidity atmosphere because of the difference of corrosion potential between the matrix Nd\_Fe<sub>4</sub>B phase and the boundary Nd-rich phase [2]. In this article, Corrosion resistance of sintered Nd-Fe-B and Sm-Co magnets was investigated in steady-state damp heat (HH), neutral salt spray (NSS) and pressure cooker (PCT) climates. The effects of alloy composition and different corrosive environments on magnetic flux loss were discussed. It shows that the corrosion rates of sintered Nd-Fe-B magnets nave a approximate linear trend. The corrosion resistance of sintered Nd-Fe-B magnets can be improved by Dy partially substitution for Nd and minor Co addition. The sintered Sm-Co magnets show very good corrosion resistance in the above three kinds of climates.

Sintered rare-earth magnets; corrosion resistance; corrosion behavior; different corrosive environments

# **RO05**

# Effects of grain size and interface state on the coercivity in Nd-Fe-B/ Nd thin films

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High coercivity in sintered Nd-Fe-B magnets is the crucial subject, especially in view of a high demand for developing Dy-saving magnets [1]. In order to approach this matter, the grain refinement and the interfacial microstructure control are known long to be the promising methods [2, 3]. In this work, the effect of Nd coating on the coercivity has been investigated in the Nd<sub>2</sub>Fe<sub>14</sub>B thin films with different grain sizes. As the Nd<sub>2</sub>Fe<sub>14</sub>B layer thickness tNFB was decreased from 70 nm to 5 nm, the coercivity exhibited a gradual increase behavior from 6.5 kOe to 16 kOe, and from 17.5 kOe to 26.2 kOe for the Nd-Fe-B and Nd-Fe-B/Nd films, respectively. It should be noted that the amount of Hc increase by the Nd coating was about 10 kOe irrespective of Nd-Fe-B layer thickness tNFB. AFM observations showed that the average Nd2Fe14B particle size of the tNFB = 5 nm film is about 60 nm, which is much smaller than the critical size for the single domain in the Nd2Fe14B phase. These results therefore suggest that a proper interface state is extremely important to achieve high coercivity not only in larger grain system but also in smaller particle one.

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### **RO06**

#### Interface state and coercivity in Nd-Fe-B/Dy films

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It is believed that the Nd-rich phase plays an important role for the coercivity emhancement in the Nd-Fe-B magnets [1-2]. In the case of Dy-diffusion-processed Nd-Fe-B magnets [3], however, it is still not clear whether only the Dy-rich shell or both Dy-shell and the Nd-rich phase is indispensible for the high coercivity. Therefore, we fabricated a model interface system, which consists of the Nd\_Fe<sub>14</sub>B layer with the Dy and Nd overlayers in order to study the relationship between microstructure near the interface and the coercivity. Nd\_Fe<sub>14</sub>B particles with the c-axis normal to the film plane were constructed by using HV or UHV sputtering systems. We introduced high-melting-point bcc metals, Ta or Mo, as an underlayer on the heat-resistant glass or sapphire c-plane single-crystal substrates. When the temperature of the substrate during a deposition of the 1 µm thick Nd-Fe-B layer was 620°C, the demagnetization curve of this film without rare-earth overlayers exhibited a nice squareness with coercivity Hc = 3.3 kOe. On the other hand, preliminary experiments on the post-annealed Nd-Fe-B lims with the Dy overlayer showed a significant increase in Hc of up to above 15 kOe.

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### **RO07**

# Investigation on the magnetic and crystalline structures of die-upset Nd-Fe-B magnets

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Die-upsetting is used to prepare anisotropic Nd-Fe-B magnets due to high magnetic texture in which the c-axis of Nd2Fe14B grains is parallel to the loading direction. [1] Though many papers have reported on the structures of high-textured Nd-Fe-B magnets [2-5], there are still some issues needed to be clarified, such as how about the uniformity of the magnetic structures and how to improve the uniformity. In this paper, x-ray diffraction (XRD) and magnetic force microscopy (MFM) are used to reveal the crystalline and magnetic structures of die-upset Nd<sub>116</sub>Fe<sub>5.73</sub>Cd<sub>4.6</sub>B<sub>3.5</sub>Cd<sub>0.4</sub>m magnets. It is found that the ratio of the XRD peak intensity 1(006)/ 1(105) from center to edge of the sample decreases from 1.86 to 1.35. The magnetic configurations of the central of the sample are typical interaction domains. The interaction domains are formed due to the strong inter-granular exchange interaction and magnetostatic interaction between grains. However, on the edge of the sample the invertex during the die-upsetting process. This indicates we need to improve the uniformity of the strain distribution for the development of high-repromance magnets.

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### **RO08**

#### A novel approach -microwave assisted sintering - for preparation of high performance permanent magnets

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A large portion of the cost for the fabrication of permanent magnets is due to the energy needed for sintering. Any improvement for the sintering process with respect to time and energy will reduce the cost and will increase the market of permanent magnets considerably. Based on the most recent development in microwave applications reported by R. Roy (1), where it was demonstrated that sintering of metal powdersa surprising application- is possible with microwaves (in view of the known fact that bulk metals reflect microwaves), we have employed microwaves to sinter Nd-Fe-B at a fraction of time and energy, compared to the conventional resistive/radiant heating in excess of 1000 0C for many hours, and in many cases homogenizing also at high temperatures. Microwave heating is a volumetric heating involving conversion of electromagnetic energy into thermal energy, which is instantaneous, rapid and highly efficient. In this work, we will present the emerging possibilities for sintering Nd-Fe-B and Sm-Co permanent magnets as well as the opening of new possibilities for metal bonded magnets by employing microwave radiation, that is already used for ferrites by various groups. Supported by the project NANOPERMAG of the EU

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#### **RO09**

### Preparation of Nd-Fe-B thin films with columnar structure and their structure and magnetic properties

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The addition of heavy rare earth elements is essential for Nd-Fe-B sintered magnet to obtain high-performance characteristic at high temperature. However, a reduction of these additive elements is strongly required since the price of these elements has increased in recent years. It was reported that the additive elements such as Cu and Ga were diffused and concentrated into the grain boundary phase (GBP) [1, 2], and the magnetic coupling between the main phases is thought to reduce by the GBP. In order to improve the hard magnetic properties, therefore, the control of the GBP is indispensable. In this study, the Nd-Fe-B thin films with well-defined columnar structure have been prepared and structure and magnetic properties have been also investigated. The samples were prepared by using an ultra-high vacuum sputtering system. The thickness of the Nd-Fe-B layer was varied in the range of 50 ~ 300 nm. From the X-ray diffraction patterns, the peaks from Nd2Fe14B (004), (006), and (008) were observed for all the samples. This indicates that the films grew with a strong (001) texture. A good magnetic squareness and relatively high coercivity of 8.3 kOe were obtained for the Nd-Fe-B films with t = 100 and 200 nm

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#### **RO10**

#### Microstructure of (Nd,Dy)-Fe-B permanent magnet by spark plasma sintering

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Sintered (Nd,Dy)-Fe-B magnets are widely used in various parts such as motor, generators, actuators, and so on because of their outstanding magnetic properties. (Nd,Dy)-Fe-B sintered magnets are generally composed of the Nd2Fe14B hard magnet phase as a matrix and a Ndrich phase at the triple junction [1,2]. However, typical sintered (Nd,Dy)-Fe-B magnets are limited in their applications due to low curie temperature[3]. Many researchers have shown that the magnetic properties of (Nd,Dy)-Fe-B magnets are structure sensitive, and dependent on the distribution of the rare-earth-rich phases, average grain size and size distribution. Particularly the grain size of Nd<sub>2</sub>Fe<sub>14</sub>B was essentially important for high magnet properties. In this study, we focused on controlling of the grain size to enhance magnetic properties by the spark plasma sintering process. The starting powder with the composition of Nd<sub>13</sub>Dy<sub>2</sub>Fe<sub>76.2</sub>TM<sub>2.8</sub>B<sub>6</sub> was prepared by strip casting and jet milling process. The raw powder was pre-sintered for preventing collapsibility of anisotropy. After the pre-sintering, (Nd,Dy)-Fe-B magnet was fabricated by spark plasma sintering process under 30 MPa with various temperatures. Consequently we had made the (Nd,Dy)-Fe-B sintered magnet with grain size of 5.9 µm. The grain size was effectively controlled by spark plasma sintering.

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### **RO11**

#### Effect of small Dy-alloy powder additions on the coercivity of NdFeB sintered magnets

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Recently, demand of high performance NdFeB sintered magnets is rapidly increasing mostly due to expansion of wind generators and hybrid electric vehicles. For traction motors of hybrid electric vehicles, the coercivity of NdFeB sintered magnet is required to exceed 25 kOe. For such high coercivity, partial substitution of Dy for Nd in Nd-Fe., B of NdFeB sintered magnet is practically inevitable, with the sacrifice of magnetization value and the increase of material cost Because the price of Dy is very high and keeps increasing due to its scarcity, it is important to reduce the amount of Dy in NdFeB magnets as much as possible without losing coercivity. Recent researches on NdFeB magnets therefore are mainly focused on Dy-saving methods. In this study as an effort to develop a Dy-saving method, small Dy-TM alloy powder was mixed with NdFeB magnet powder and then a sintered magnet was fabricated by a conventional process. When 1.9 wt% of DvCo powder was mixed, the coercivity of the sintered magnet increased about 4 kOe, indicating 20% of Dy saving effect. More experimental results will be discussed in this presentation

#### **RO12**

Coercivity of near single domain size Nd-Fe-B-type alloy particles Hae-woong Kwon1 and J H Yu2 <sup>1</sup> Pukvong National University, Korea <sup>2</sup> KIMS Korea

Fine particles of Nd-Fe-B alloy in near single domain size have increasingly found new applications for high performance micro-magnet. However, because of high oxygenaffinity and high specific surface area, the ultra-fine Nd-Fe-B-type particles are readily oxidized in air, hence causing a radical coercivity loss. In the present study, an ultrafine Nd-Fe-B-type particles in near single domain size was prepared by ball milling of HDDR-treated Nd12.5Fe80.6B6.4Ga0.3Nb0.2 alloy. The prepared near single domain size Nd-Fe-B-type powder (
0.3 um) had high coercivity over 9 kOe. However, the coercivity was radically reduced as the temperature increased in air (< 2 kOe at 200 °C). Room temperature long-term stability of coercivity of the fine powder in air was investigated for 1 month, and the coercivity was decreased by rate of 0.5 kOe/week Feasibilty of surface nitrogenation of the fine powder for improving the long-term stability of coercivity was studied. Long-term stability of coercivity of the nitrogenated fine powder was improved markedly. The nitrogenated powder showed no coercvity reduction even after 1 month in air at room temperature. In this article, the surface passivation of near single domain size Nd-Fe-B-type particles by nitrogenation for improving long-term stability of coercivity is to be discussed

### **RO13**

#### Effect of annealing temperature on microstructure, magnetic properties and corrosion resistance of NdFeB/α-Fe nanocomposite magnets

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REFeB based nanocomposite alloys have attracted much attention in both the scientific community and bonded magnet industry. However, lower coercivity and poor corrosion resistance limit their applications. Both the practical coercivity and corrosion resistance are sensitive to the character of phase constitute in the microstructure. Sintered NdFeB magnet is the poor corrosion resistance due to the presence of the highly corrosive Ndrich phase. As no Nd-rich phase is formed in the NdFeB/a-Fe nanocomposite magnets. these alloys should be more corrosion resistant than the NdFeB sintered magnet. In this work, the effects of annealing temperature on the microstructure, magnetic properties and corrosion resistance in NdFeB/a-Fe alloys were investigated. It is shown that after heat treatment, the alloys mainly consist of two phases: a hard magnetic phase of  $Nd_2Fe_{14}B$  (space group  $P_{42}/mnm)$  and a soft magnetic phase of  $\alpha\mbox{-}Fe$  (Im3m), and that no other phases are identified. Moreover, the optimum corrosion resistance of NdFeB/ a-Fe magnets was observed with the annealing temperature at 610°C. Keywords: Nanocomposite, Magnetic property, corrosion resistance

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# **RO14**

Effect of magnetic heat-treatment on magnetic properties and corrosion behavior of Nd<sub>6</sub>Fe<sub>72,x</sub>Co<sub>x</sub>B<sub>22</sub>(x=10, 20, 30) nanocomposite ribbons

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Nanocomposite rare-earth magnets consisting of a hard and a soft magnetic phase have attracted considerable attention because the remanence enhancement in exchange coupled nanocomposite alloys has potential importance as a way to achieve high values of the maximum energy product (BH)max. It is well established that the magnetic properties of nanocrystalline two-phase magnets are strongly dependent on its microstructure. A uniform microstructure and a well-coupled grain interface should be two essential conditions for a well-performing nanocomposite. Magnetic annealing has been tried to perfect the microstructure and enhance the magnetic properties. In this paper, Nd<sub>6</sub>Fe<sub>72-x</sub>B<sub>22</sub>(x=10,20,30) ribbons were heat-treated with or without an external magnetic field of 3 kOe in argon atmosphere in the temperature range of 600-750°C for 10 min. It is found that magnetic heattreatment enhances both coercivity and remanence ratio of the Nd<sub>6</sub>Fe<sub>72-x</sub>CoxB<sub>22</sub>(x=10,20,30) ribbon, and decrease the annealed temperature obtained the optimum magnetic properties due to the uniform distribution and preferential orientation of  $\alpha$ -Fe and Nd<sub>2</sub>Fe<sub>14</sub>B grains. The corrosion behavior of Nd<sub>6</sub>Fe72, CoxB22 (x=10, 20, 30)nanocomposite ribbons with or without an external magnetic field is discussed in detail.

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#### **RO15**

### Study on magnetization reversal behavior for the annealed $Nd_2Fe_{14}B/$ α-Fe nancomposite alloys

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Effect of thermal annealing on the magnetization reversal behavior of α-Fe/Nd<sub>2</sub>Fe<sub>14</sub>B alloys has been investigated. A drastic increase of the remanence Mr from 0.67T up to 0.87T and remanence ratio Mr/Ms from 0.66 up to 0.76, respectively, is observed in the  $\alpha$ -Fe/Nd<sub>2</sub>Fe<sub>14</sub>B alloys annealed at 610 oC as compared with the as-quenched sample. Whereas the further annealing at 680 oC results in a strongly increase of the corecivity Hc as high as 491 kA/m but a slight decrease in Mr. The analysis result of the magnetization reversal behavior shows that the maximum value of the integrated recoil loop area about 1.58 kJ/m<sub>3</sub> is obtained in the α-Fe/Nd<sub>2</sub>Fe<sub>14</sub>B alloys at the annealing temperature of 610 oC, significantly lower than other annealed samples. This indicates a significant advantage for the application of this material as permanent magnets in electrical machines and generators due to a low energy loss. Keywords: Nanocomposite; Magnetization reversal behavior; Recoil loop

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### **RP01**

#### Construction of a versatile neutron-scattering spectrometer HERMES-E using renovated Ge monochromator crystals

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We propose an efficient way, which increases a number of neutron spectrometers in a research-reactor facility. The first key is to making use of a characteristic Gecrystal monochromator, which discretely emits Bragg reflections to different scattering angles at the same time. The second is to save the instrumental space by using curved and deformed Ge-wafer crystals as a monochromator or an analyzer. Indeed, we are planning to double a neutron-powder diffractometer of Tohoku University (HERMES [1]) that is working at a research reactor JRR-3, Tokai, Japan; a versatile neutronscattering spectrometer HERMES-E is now under construction next to HERMES.

[1] K. Ohoyama et al., Jpn. J. Appl. Phys. vol. 37 (1998) 3319.

# **RP02**

### NMR study of the phase transition behavior in Ce<sub>3</sub>Co<sub>4</sub>Sn<sub>13</sub>

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We present a study of the phase transition behavior in Ce<sub>3</sub>Co<sub>4</sub>Sn<sub>13</sub> using 59Co nuclear magnetic resonance (NMR) spectroscopy. The quadrupole splitting, Knight shift, and spin-lattice relaxation rate (1/T1) below and above the phase transition temperature To have been identified. All measured NMR quantities exhibit pronounced features at around To=155 K except for the Knight shift. It thus excludes the magnetic origin for the observed transition. In addition, the x-ray diffraction results below and above To confirm the absence of a crystal structural change, suggesting that the peculiar phase transition is possibly driven by the charge-density-wave (CDW) formation. As a matter of fact, the double-peak feature in the 59Co NMR central line smears out below To which can be associated with the spatial modulation of the electric field gradient (EFG) due to incommensurate CDW superlattices. Furthermore, a distinct peak found in the spin-lattice relaxation rate near To can be accounted for by the thermally-driven normal modes of the CDW

### RP03

### Analysis of 1/f noise characteristics of magneto-optical Kerr effect measured from Co/Pt and NiFe/Pt multilayers thin film

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We have investigated 1/f noise characteristic of the polar and longitudinal magnetooptical Kerr effect (MOKE). The hand-made MOKE measurement system was set using diode laser as a light source. Si photodetector was used as a detector where the bias has been systematically controlled. To increase a signal-to-noise ratio, the laser light was modulated with the lock-in amplification frequency referenced from the mechanical chopper in the frequency ranges up to 4 kHz. Noise characteristics of the optical/electrical/mechanical error were systematically examined. To separate the noise contribution from the ferromagnetic layer and from the intrinsic measurement error, various configurational changes were made for the lock-in frequency as well as the bias voltage of the Si photodetector. Magnetic noise properties under a different magnetic field were systematically investigated as well. For a perpendicular anisotropy sample, Co/Pt multilayer film was investigated while for an in-plane anisotropy sample, NiFe/ Pt bilayer film was investigated. The noise characteristics with respect to the multilayer thickness variations were examined as well. The 1/f noise characteristics observed for a wide range regardless of the lock-in frequency change or the film thickness change implies that there might be a universal characteristics in the MOKE signal noise, arising from the sample itself.

### **RP04**

#### Photoemission electron microscopy of three-dimensional magnetization configurations in core-shell nanostructures

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We present a photoemission electron microscopy method that combines magnetic imaging of the surface and of the inner magnetization in three-dimensional core-shell nanostructures [1]. The structure investigated consists of a cylindrical nickel core that is completely surrounded by a shell of iron oxide and silicon oxide layers. The method enables one to image the magnetization configuration of the nickel core even though the shell is thicker than the mean-free path of the photoelectrons. Characteristic L<sub>3</sub> and L<sub>2</sub> edges can be observed not only in the yield of the photoelectrons emitted from the surface of the nanostructure but also in its shadow. X-ray magnetic circular dichroism in the electron yield of the x rays absorbed and transmitted by the multilayered nanowire allows for the individual imaging of the magnetization configurations of the iron oxide tube and the nickel core. The method suggests novel approaches for the characterization of the magnetic and material properties of complex three-dimensional nanostructures. Financial support by the DAAD via the Project 50725506 and by DFG via the SFB 668 and the GrK 1286 as well as the Forschungs- und Wissenschaftsstiftung Hamburg via the Exzellenzcluster "Nano-Spintronik" is gratefully acknowledged.

[1] J. Kimling, F. Kronast, S. Martens, T. Boehnert, M. Martens, J. Herrero-Albillos, L. Tati-Bismaths, U. Merkt, K. Nielsch, and G. Meier, Phys. Rev. B 84, 174406 (2011).

### **RP05**

Development of high-field ESR system using SQUID magnetometer and its application to measurement under high pressure

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We have developed high-field ESR system using commercial SQUID magnetometer, in which ESR is observed as a change of magnetization. [J. Phys.: Conf. Ser. 334 (2011) 012058.] The advantage of this method as compared with conventional high-field ESR is that the measurement can be done very easily. Moreover, the macroscopic magnetization measurement can be made simultaneously. Although the sensitivity of this system is lower than that of ESR system using cavity, it is comparable with that of transmission type ESR system. Recently, we have succeeded in performing ESR measurement under high pressure by this ESR system. The pressure is generated by a widely used clamped type piston-cylinder pressure cell whose cylinder is made of non-magnetic CuBe alloy. Since all inner parts are made of zirconium oxide, the electromagnetic wave can be introduced into the sample space. Therefore, we can measure the magnetic moment of sample and obtain ESR signal under pressure. The pressure is available up to 1.5 GPa. This is the easiest ESR system which gives us pressure dependence ESR spectra as far as we know. The setup of this system and several results will be shown in detail.

### RP06

Alternating magnetic force microscopy: direction detectable imaging of static and alternating magnetic field with high spatial resolution Hitoshi Saito<sup>1</sup>, Ito Ryoichi<sup>1</sup>, Kodai Hatakeyama<sup>1</sup>, Zhenghua Li<sup>2</sup>, Genta Egawa<sup>1</sup> and

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We have developed a new functional MFM named as "Alternating Magnetic Force Microscopy (A-MFM)" which enables us to measure static and alternating magnetic fields with the polarity of field direction [1-5]. A-MFM utilized the frequency modulation of a cantilever oscillation generated by an off-resonance alternating magnetic force between a magnetic sample and a magnetic tip. The alternating magnetic force was extracted by a lock-in amplifier for the frequency demodulated signal of the cantilever oscillation. For the static magnetic field from a sample, a soft magnetic tip which ware driven by an AC magnetic field was used to detect the static magnetic field with high sensitivity. On the other hand, for the alternating magnetic field from a sample, such as a magnetic writing head, a hard magnetic tip was used A-MFM can measure a magnetic force without surface short-range forces and directly detect the direction of perpendicular magnetic fields. When the tip-sample distance became short, spatial resolution was improved. In the conference, we will demonstrate high performance of A-MFM and show high-resolution magnetic field images with a resolution of less than 10 nm for highdensity recording media and magnetic writing heads. This study was supported by SENTAN, JST.

[1] H. Saito et al. J. Appl. Phys. 105 (2009) 07D524, [2] H. Saito et al. J. Appl. Phys. 107 (2010) 09D309, [3] W. Lu et al. Appl. Phys. Lett. 96 (2010) 143104, [4] W. Lu et al. IEEE Trans. Magn. 46 (2010) 1479, [5] H. Saito et al. J. Appl. Phys. 109 (2011) 07E330

#### **RP07**

# 30 T pulsed-high-magnetic-field and element-selective magnetization studies using soft x-ray magnetic circular dichroism

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We demonstrate the applicability of 30 T pulsed-high-magnetic-field to observe magnetic circular dichroism (MCD) of antiferromagnetic compounds in soft x-ray energy regions, in which there are magnetically important transition energies, L-edge of 3d transition metal and M-edge of 4f rare earth elements. We have developed a specially designed pulse magnet equipped with an ultra-high-vacuum environment and a portable capacitor bank for synchrotron soft x-ray experiments [1]. A time-resolved total electron yield method is used to record bank for synchrotron sin synchronization with pulsed-high-magnetic-fields and periodic switching of circularly polarized soft x-rays. We have measured elements-selective magnetizations of Ni, Co and Mn in the NiCoMnIn alloy which shows a large field-induced strain accompanied by a drastic change of the magnetization [2]. All 3d transition metal elements show steep non-linear increases of the MCD intensities around 15 T corresponding to the metamagnetic transition. This high-magnetic-field soft x-ray MCD result directly evidences that all magnetic elements to the field-induced shape memory effect in the NiCoMnIn alloy.

[1] T. Nakamura et al., Appl. Phys. Express 4, 066602 (2011). [2] R. Kainuma et al., Nature 439, 957 (2006).

#### **RP08**

#### Imaging magnetic responses of nanomagnets by X-ray PhotoEmission Electron Microscopy

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The Spin-resolved PhotoEmission Electron Microscope (SPEEM) is a permanently installed setup at the X-ray synchrotron radiation source BESSY, Helmholtz-Zentrum Berlin. The capabilities of excellent spatial resolution and an energy- and polarization-tunable light source make this end-station ideal for magnetic imaging and micro-spectroscopy with quantitative analysis. The application of magnetic fields while measuring significantly expands the technique of magnetic imaging. A dedicated magnetic sample holder incorporating both DC and AC magnetic field with additional temperature control has recently been developed. The SPEEM end-station, in combination with this novel sample holder enables the insitu control of the magnetic state and direct observation of the corresponding magnetic response. Two selected examples demonstrate these capabilities: magnetization curves of individual Fe nanocubes, and the local magnetic as susceptibility in and microstructures as a function of temperature. In the latter example, a surprising onset of finite size effects at the micro-scale has been observed, with a significant reduction of the ordering temperature in microstructures compared to the extended film. These results were observed specifically as a result of these improvements in the PEEM technique. In combination with other possibilities, such as depth resolved imaging, new insights to magnetic nanostructures can be realized.

#### RP09

### Development of high-sensitivity cantilever-detected ESR measurement using a fiber-optic interferometer.

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Cantilever-detected high-frequency electron spin resonance (ESR) is a powerful method for terahertz ESR spectroscopy of a tiny magnetic sample at low temperature [1,2]. In this technique, a small magnetization change associated with ESR transition is detected as deflection of a sample-mounted cantilever. So far, we have succeeded in ESR detection at 370 GHz using a commercial piezoresistive microcantilever. The spin sensitivity was estimated to ~1012 spins/ gauss. In order to further increase the sensitivity, we adopt a fiber-optic-based detection system using a Fabry-Perot interferometer in place of piezoresistive system. Fabry-Perot cavity is formed between an optical-fiber end and microcantilever surface, and a change in interference signal, corresponding to the cantilever deflection, is sensitivity detected. This system is suitable for low-temperature and high-magnetic field experiments because of its compact setup and less heat dissipation. We will combine our fiber-optic detection system with a 15 T superconducting magnet and BWO light sources covering a wide frequency range 200-1200 GHz. Test measurement with Co Tutton salt is now in process to estimate the system performance. In the presentation, details of our experimental setup and test results will be presented.

[1] E.Ohmichi, N. Mizuno, M. Kimata, and H.Ohta, Rev. Sci. Instrum. 79, 103903 (2008) [2] E.Ohmichi, and T.Osada, Rev. Sci. Instrum 73, 3022 (2002)

#### **RP10**

#### Study of the epitaxial growth and perpendicular magnetic domain structure of ordered FePt thin film on MgO substrate using HRTEM and electron holography

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The epitaxial film of ordered FePt, showing the c-axis normal to the film plane, was successfully grown on top of MgO substrate using DC-magnetron co-sputtering technique. The electron holography study showed that a columnar epitaxial film displayed only a weak stray-field signal due to the presence of columnar boundaries. The growth of a continuous (full) epitaxial film however revealed characteristic Lorentz TEM images for the perpendicular magnetic domains. The electron hologram indicated the presence of the stray-fields of induction originating from the perpendicular magnetizations. It was shown that the sense of stray fields and thus the sense of perpendicular magnetizations could be determined by analyzing the contraction and expansion behavior of hologram finges which were parallel to the reference fringes. Furthermore, direct hologram information was detected in a thicker epitaxial film by applying high resolution electron holography. Perpendicular magnetic domain structure was directly revealed in the unwrapped phase image and was confirmed by the measurement of the line profile of phase-shift. The estimation of the local strength indicated that the perpendicular induction strength was varying from about 0.33T to about 0.61T depending on the domain size. The variation of induction strength was due to a weak ordering of the film.

# RP11

# Asteroid curve of GMR films on the practical substrate under the stress

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Magnetic thin films usually have some of magnetostriction effect. After forming electric devices, this effect sometimes deteriorates the sensitivity of sensors such as read heads of hard disk drives, HDD, or the bit stability of memories such as magnetic random access memories, MRAM, through the inverse magnetostriction phenomenon. We need, therefore, to understand the magnetostriction effect on a film deposited on an actual wafer and the influence from the stress to it. In the previous paper, we showed the novel method of measuring the magnetostriction constant on the practical substrate [1]. In this paper, by using this method we measured the asteroid curves of the magnetic thin film of the giant magnetoresistance, GMR, on the practical wafer with varying applied stress on it. This result showed that the coercive force parallel to the easy axis, Hc, of a magnetic thin film with a negative sign of magnetostriction constant was decreasing with increasing tensile stress and the coercive force parallel to the hard axis, Ha, was also decreasing. For the film with the positive sign one, Hc and Ha were increasing with increasing tensile stress. And we showed the Hc was roughly proportional to the applied stress.

[1] K. Okita, K. Ishiyama and H. Miura, 'Magnetostriction Measurement of GMR Films on the practical substrate', Journal of Physics, Conference Series 200 (2010) 112008

#### RP12

Micromagnetic study on the perturbative effect of magnetic force microscopy probes on 90° asymmetric Neel walls in a soft magnetic material Hironori Asada<sup>1\*</sup>, Hidenori Kubo<sup>1</sup>, Hazrina Abu Seman<sup>1</sup>, Takashi Manago<sup>2</sup> and Hiromi Kuramochi<sup>3</sup>

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It is important to clarify the dependence of perturbative effect of MFM probes upon probe parameters and experimental conditions[1]. In this study, 3D micromagnetic simulation has been performed to evaluate the perturbative effects on magnetization states and force gradient signals of 90° asymmetric Neel walls in a 100nm-thick patterned permalloy film. The signals are computed by varying the position of the probe assuming the stripe shape. The different signal distortions are observed depending upon the combination of the out-of-plane component of the magnetization direction of wall and the magnetization direction of the probe. In the parallel case, signal asymmetry originated from the asymmetric wall structure is enhanced and the normalized contrast defined as the difference of the maximum and minimum values of the signal divided by probe saturation magnetization(Ms) is increased. In the antiparallel case, signal asymmetry is suppressed and the normalized contrast is decreased. When probe Ms is increased firther (antiparallel case), the out-of-plane component direction of wall is reversed being accompanied by the jump of the Neel cap positions. The probe Ms dependences of signal distortions such as contrast variation and distance change between the signal peaks which correspond to the measured wall width are also presented.

[1] J. M. Garcia, et al., Appl. Phys. Lett. 79, 656 (2001).

### RP13

# Development of a non-conventional ESR spectrometer with a composite antenna system and an electronically controlled tuning and matching circuit

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We present a new development of non-conventional Electron Spin Resonance (ESR) spectrometer, in which the sweeping frequency is available in the wide range of frequency. This is a modification of the previous system without using a cavity and with two antennas one for the microwave transmission and the other for the ESR signal detection.[1] In contrast to previous work[1], utilization of microstrip antenna as a transmitter provided us the capability of wide-frequency operation between 0.8 and 10 GHz. On the other hand, the use of loog antenna as a detector with tuning and matching circuit enhanced ESR signal. In order to simplify tuning and matching circuit and make resonance condition easily reproducible, conventional capacitors were replaced with varactor diodes. This allowed us to avoid any mechanical action during tuning and matching procedure, since the capacitance of the diodes was changed by applying different DC voltages. Compared to Ref[1], such a schemen increased SN ratio several times in a wide frequency range. For testing the developed system, a homemade magnet with water-cooling system was used. The efficiency of the ESR spectrometer was checked by measuring a signal of 1,1-diphenil-2-picrylhydrazyl (DPPH) sample at different frequencies at room temperature.

[1] Z. H. Jang, B. J. Suh, M. Corti, L. Cattaneo, D. Hajny, F. Borsa, and M. Luban, Rev. Sci. Instrum. 79, 46101 (2008).

### RP14

A new type of spin-polarized scanning tunneling microscopy for observing an in-plane magnetization component with high resolution

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To obtain a high resolution image routinely even from the rough surface of a ferromagnetic sample with a spin-polarized scanning tunneling microscope (SP-STM) operating in modulating tip magnetization mode [1], we have developed an SP-STM using a micrometersized magnetic tip integrated onto the magnetic recording head of a hard disk drive. The tip apex was formed into a round shape with a diameter of less than one micrometer by using a focused ion beam instrument. Driving the recording head with a signal generator to switch the magnetization of the tip apex periodically generates an unnecessary artificial current flow into a current-to-voltage (IV) converter through stray capacitance, which prevents detection of the spin-polarized tunneling current precisely. To reduce the artificial current, we added a homemade circuit that generates a countercurrent against the artificial current into the input of the IV converter. We tried to image both the topography and the spin polarization of an iron thin film deposited on a magnetic recording medium whose local roughness was on the order of several nanometers due to the crystal grains of the medium. This is the first spin image of a magnetic thin film having such a rough surface obtained with an SP-STM.

[1] U. Schlickum, W. Wulfhekel, and J. Kirschner, Appl. Phys. Lett. 83, 2016 (2003).

#### RP15

**Violation of Hund's third rule in structurally disordered ferromagnets** Vassilios Kapaklis<sup>1</sup>\*, Panagiotis Korelis<sup>1</sup>, Bjorgvin Hjorvarsson<sup>1</sup>, Athanasios Vlachos<sup>2</sup>, Iossif Galanakis<sup>3</sup>, Panagiotis Poulopoulos<sup>3</sup>, K. Ozdogan<sup>4</sup>, Makis Angelakeris<sup>5</sup>, Fabrice Wilhelm<sup>6</sup> and Andrei Rogalev<sup>6</sup>

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Violation of Hund's third rule caused by structural disorder is observed for the induced magnetic moment of Zr, using x-ray magnetic circular dichroism. The induced spin and orbital magnetic moments are antiparallel in the crystalline state, but parallel in an amorphous state of the investigated Co- and Fe-based materials. First-principles calculations are used to provide physical insight into the dependency of the spin-orbit coupling on the interatomic distance and coordination number.

V. Kapaklis, P. T. Korelis, B. Hjorvarsson, A. Vlachos, I. Galanakis, P. Poulopoulos, K. Ozdogan, M. Angelakeris, F. Wilhelm, and A. Rogalev, Violation of Hund's third rule in structurally disordered ferromagnets, Phys. Rev. B 84, 024411 (2011).

### RP16

# Bulk Cr tips with full spatial magnetic sensitivity for spin-polarized scanning tunneling microscopy

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Spin-polarized scanning tunneling microscopy (SP-STM) is a powerful technique to access the magnetic properties of a surface on the atomic scale. Here, magnetic tips are usually prepared by extensive in-situ metal evaporation, demanding for a tip exchange mechanism inside the microscope. Using magnetic bulk tips is an elegant way to circumvent these efforts. For example, bulk chromium tips are reported to be suitable for SP-STM, exhibiting an in-plane magnetic sensitivity [1]. Since they are antiferromagnetic, their influence on the sample magnetization in terms of stray field interaction is negligible. We perform SP-STM experiments on the sample system of 1.8 atomic layers Fe/W(110) with the magnetization lying in the plane on the monolayer and pointing perpendicular to the plane on the double layer [2]. Our study demonstrates that Cr bulk tips can be used to image the complete magnetic structure and hence are sensitive to both the in-plane as well as the out-of-plane-component of sample magnetization [3]. Analyzing magnetic SP-STM maps enables the full spatial characterization of the tip.

[1] A. L. Bassi et al., Appl. Phys. Lett. 91, 173120 (2007). [2] M. Pratzer et al., Phys. Rev. Lett. 87, 127201 (2001). [3] A. Schlenhoff et al., Appl. Phys. Lett. 97, 083104 (2010).

### **RP17**

#### Polarization state of scattered light in apertureless reflection-mode magneto-optical scanning near-field optical microscopy

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Magneto-optical scanning near-field optical microscopy (MO-SNOM) has attracted attention because of its high spatial resolution in magnetic imaging. We are developing a reflection-mode apertureless scanning nearfield optical microscopy (a-SNOM). The sample is a patterned Chromium film with a thickness of 20nm deposited on a glass substrate. The a-SNOM images are measured with Si tip that has an extremity's radius of 7nm, and with an illumination of 408 nm with an incident angle of 45 degree. Scattered light caused by an interaction between the near-field wave and the sample was measured as a signal for a-SNOM image. Finitedifference time-domain (FDTD) method was used to calculate the polarization state of the scattered light. The a-SNOM images of Cr patterns were successfully obtained with a spatial resolution of 30-40nm. In addition, we found that the extinction ration of the scattered light back to the same direction with the incidence was better than 100 which is an enough polarization property for obtaining MO images FDTD simulation reproduced that the scattered light preserved its polarization state. These results showed that the a-SNOM is a promising technique to obtain high-resolution magnetic images with a resolution of several tens nm.

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### **RP18**

#### Application of image processing to determine size distribution of magnetic nanoparticles

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Digital image processing has increasingly been implemented in analyzing micro- and nanostructured materials and would be an ideal tool to characterize the morphology and position of self-assembled magnetic nanoparticles for high density recording. In this work, magnetic nanoparticles were synthesized by the modified polyol process using Fe(acac)<sub>3</sub> and Pt(acac)<sub>2</sub> as starting materials. Transmission electron microscopy (TEM) images of an as-synthesized product were taken with a resolution of 800 x 800 pixels, 72 dot per inch and then inspected using image processing algorithms on MatLab. The grayscale image was converted to a binary image by using Otsu's thresholding. An individual particle was then detected by using the closing algorithm with disk structuring elements and the canny edge detection. The areas of detected particles were filled and small objects were removed. A centroid, diameter and area of each particle were finally evaluated. The average diameter is 4.62 nm with a standard deviation of 0.52 nm. The degree of polydispersity of magnetic nanoparticles can then be compared using size distribution from the image processing.

#### **RP19**

#### Microfabrication of a MEMS cantilever for mechanically detected high-frequency ESR measurement

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now in progress and details will be reported in the presentation.

Mechanically detected high-frequency electron spin resonance (ESR) is a promising tool for high-resolution terahertz (THz) ESR spectroscopy of tiny magnetic samples. So far we succeeded in multi-frequency ESR detection for 80-370 GHz and achieved a spin sensitivity on the order of 10<sup>9</sup> spins/Gauss[1,2]. In this study, a prototype microcantilever is fabricated in order to further improve the sensitivity and flexibility of our ESR measurement system. Microcantilevers are fabricated from silicon-oninsulator (SOI) wafers using standard MEMS techniques such as lithography, wet etching, and RIE. By using commercial SOI wafers, fabrication cost and the number of process can be substantially reduced. Two types of cantilevers for capacitive and optical detection are fabricated in this study. The former has lateral dimensions of 4 mm x 8 mm, and is anodically bonded to a glass substrate to form capacitive electrodes. The latter has dimensions of 50  $\mu m$  x 200  $\mu m$  , and is equipped with a gold mirror for Fabry-Perot detection. Application to high-frequency ESR measurement is

[1]E. Ohmichi, N. Mizuno, M. Kimata, and H. Ohta, Rev. Sci. Instrum. 79 (2008) 103903. [2]E. Ohmichi, N. Mizuno, S. Hirano, and H. Ohta, J. Low Temp. Phys. 159 (2010) 276

### **RP20**

#### 2D reflection-type electron spin filter increasing the detection efficiency in spinresolved spectroscopy by 4 orders of magnitude

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Spin-resolved electron spectroscopy is characterized by a low figure of merit FoM= S<sup>2</sup>I/lo typically 10<sup>-4</sup>. S spin sensitivity. I/Io reflectivity. Electron diffraction from W(001) in the (00) LEED spot facilitates naralle detection of 3800 data points in our imaging spin filter behind a PEEM [1], and 960 behind a hemispherical energy analyzer [2]. We achieved a "2D FoM" of 1.7 being four orders of magnitude higher than the previous value. We further studied the dependence of the spin sensitivity and reflectivity as a function of scattering energy and angle of incidence. The experimental setup includes a spin-polarized GaAs electron source hemispherical analyzer and a Delaylinedetector. Intensities. spin-orbit-coupling induced asymmetries and FoM were calculated via a relativistic layer KKR SPLEED code [3]. The superior performance of multichannel spin detection facilitates experiments on highly reactive surfaces like in-situ prepared Heusler films [4] or radiationsensitive organic layers [5] and payes the way to single-shot experiments at ultra bright fs-sources like FELs. Funded by DFG (Scho341/9), Stiftung Innovation (project 886) and centre of complex materials COMATT.

[1] C. Tusche et al., APL 99 (2011) 032505; [2] M. Kolbe et al., PRL 107 (2011) 207601; [3] R. Feder, Polarized Electrons in Surface Physics (World Scientific, Singapore, 1985); [4] M. Hahn et al., APL 98 (2011) 232503: [5] K. Medianik et al., JACS 134 (2012) 4694

#### **RP21**

#### Element selective magnetization measurements under high magnetic field

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X-ray Magnetic Circular Dichroism (XMCD) spectroscopy is a well-established experimental tool to study the microscopic origin of magnetism allowing one to determine separately spin and orbital magnetic moments of each element in both amplitude and direction. So far, XMCD has been extensively used to investigate mainly ferro- or ferrimagnetic materials, and only very few studies have been performed on paramagnetic compounds. In this presentation we describe first a new experimental set-up dedicated to high field XMCD measurements that has been recently installed at the ESRF beamline ID12. Static magnetic field of up to 17 Tesla is generated by a superconducting solenoid. The sample is mounted on a cold finger of a He continuous flow cryostat allowing to set the temperature in the range from 2.2 K to 300 K with a stability of about 100 mK. Performances of this set-up are illustrated with results of element selective magnetization measurements in various systems: (i) Intrinsic magnetic moment in gold nanoparticles grown onto naturally thiol-containing proteinatious archaeal surface layer has been evidenced by field dependent XMCD at the Au L2,3-edges[1]; (ii) Extrinsic origin of room temperature ferromagnetism in diluted magnetic semiconductors has been proven with field and temperature dependent XMCD measurements[2].

[1] S. Selenska-Pobell et al., Nanomater. nanotechnol., 1 (2011), 8 [2] A. Ney et al., New J. Phys., 13 (2011) 103001

#### RP22

#### POLI: The new single crystal polarized neutron diffractometer for investigation complex magnetic structures at FRM-II Vladimir Hutanu\*, Martin Meven, Gernot Heger and Georg Roth

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POLI is a new single crystal polarized neutron diffractometer [1] at the hot source of the 'Forschungsneutronenquelle Heinz Maier-Leibnitz' (FRM-II) in Garching, Germany. It is designed to perform spherical neutron polarimetry (SNP) in zero field as well as classical polarized neutron diffraction (PND) measurements (flipping ratio) under applied magnetic fields. Among the intended applications of this instrument are: complex magnetic structures and magneto-electric coupling in multiferroics, magnetic structure and spatially resolved spin densities in molecular magnets, magnetically ordered superconductors, strongly correlated transition metal oxides and spin chain compounds. The instrument is unique in the sense that it combines a variable wavelength focussing monochromator with the use of 3He spin-filter cells (SFC) to create polarized, short wavelength neutrons. A new zero-field polarimeter Cryopad [2, 3] has been built in cooperation between RWTH and ILL. On the detector side, the polarisation-analysis and detection unit DECPOL (again with 3He-spin filter cells) is used. This setup called 'POLI-HEIDI' is operational and open for users at FRM-II over JCNS proposal system. Its properties will be demonstrated in this talk.

[1] V. Hutanu, M. Meven, E. Lelievre-Berna, G. Heger, Physica B 404, 2633 (2009). [2] F. Tasset, Physica B 156-157, 627 (1989). [3] E. Lelievre-Berna, "Novel polarized neutron tools", I. Anderson & B. Guerard, Editors, Proceedings of SPIE 4785, 112-125 (2002).

# RP23

#### Observation of the magnetic domain using scanning electron microscopy with polarization analysis

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We have developed a scanning electron microscope with polarization analysis(SEMPA or spin-SEM) for the nano-scale observation of magnetic domain structure. A typical micro Mott type spin detector with 25 kV operating voltage, shows the effective Sherman function in the range between -0.3~-0.35, which is relatively low valued. In order to get higher Sherman function, we optimized all geometrical factors and the accelerating voltage. First, we tried Monte Carlo simulation by increasing operating voltage from 25 kV to 50 kV. At 40 kV of acceleration voltage with optimized 120 degrees of scattering angle, a better effective Sherman function was obtained compared with the value operated at 25 kV(-0.5 : -0.3). Base on this simulation, we have developed the high-efficiency Mott spin detector with 40 kV operating voltage. Also a secondary electron deflector, and transfer lens system has been optimized. In order to test performance of the new spin detector, the thick Fe film sample was prepared on Si substrate using an evaporation method. We will show several different images of Fe films depending on the oxidation condition

#### RO01

#### Precision broadband ac measurement system for magnetotransport. application of magnetoimpedance effect for protein biomarker magnetopolarization and magnetoelectric properties detection Jun Lu\*, Baogen Shen and Xiaoping Shao

State Key Laboratory of Magnetism, Institute of Physics, Chinese Academy of Sciences, China

For complex or strongly correlated magnetic materials, broadband ac measurements provide much richer dynamic information than static or quasistatic measurements, which would help one to understand more clearly how magnetic materials of interest respond to various external excitations. After many-year researches, we have developed a suite of precision broadband ac magnetic measurement system based on novel broadband digital lock-in amplifier techniques, which can work from sub 1 Hz to above 1 GHz. In addition to multi-channel lock-in amplifier techniques, to realize broadband ac transporting and polarizability measurement from sub 1 Hz to above 100 MHz, transmission-line bridge compensation techniques have been used so as to measure complex ac magnetoresistance and Hall coefficient simultaneously in a precision way. Broadband ac magnetoelectric measurement system has also been set up for evaluating complex magnetoelectric coupling coefficient, which including not only magnitude but also phase lag. For various magnetotransport and magnetoelectric measurements, we have developed measuring software which supports fully automatic tests.

[1] J. Lu et al., Meas. Sci. Technol., 19, 045702-6(2008) [2] J. Lu et al., IEEE Trans. Magn., 44(9), 2127-9(2008) [3] J. Lu et al., A kind of precision broadband anti-noise lock-in frequ Chinese Patent, 201110380805.X

### **RO02**

#### Detection of magnetic beads using an extraordinary magnetoresistance sensor fabricated with unpatterned semiconductor substrate Jian Sun\* and Jurgen Kosel

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A strong geometry dependent magnetoresistance effect, the so-called extraordinary magnetoresistance (EMR), has been observed in semiconductor/metal hybrid structures. It is based on the redistribution of the current from the metal shunt into the semiconductor causes by the Lorentz force. A typical EMR device consists of a patterned semiconductor mesa and metal shunt. When the dimensions of the device are reduced, for example, to increase the sensitivity, the alignment of the metal contact with the semiconductor bar becomes more difficult. Furthermore, the patterning process inescapably causes some damage to the semiconductor thereby reducing the performance. In this work, we report a new EMR device fabrication process, which does not require patterning of the semiconductor substrate reducing the complexity of the fabrication and omitting any kind of damage of the semiconductor. The sensitivity of this device is 0.58  $\mu T/\sqrt{Hz},$  which is twice as high as the one of the patterned sensor. The device is employed to detect superparamagnetic beads of 2.8 µm in diameter, which is the first time this is demonstrated with an EMR sensor. Applying the beads to the sensor surface results in a voltage signal of 100 µV.

# RO03

### Magnetization of single ferromagnetic-grain obtained from observation of field-induced-translation in a chamber-type µg system.

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A new principle to measure saturated magnetization Ms of a single ferromagnetic-grain is proposed, which is based on free translation of the grain caused by field-gradient force. The motions were observed in a diffused (~10Pa) uG condition (0.01G). It is confirmed for the first time that a classical conservation rule between field-induced potential and kinetic energy is conserved for a translating magnetic-grain. Accordingly, Ms is obtained from the relationship between sample velocity and field intensity observed at different sample positions. The Ms value is obtained even for a nano-sized sample, provided that translation is observed by an ultra-violet fluorescence microscope [1]. This is because the method does not require mass measurement: it is also free of the interfering signal emitted from a sample holder [1]. The above two factors has been the major problem in detecting magnetization of single grain in conventional methods. A chamber-type drop shaft to produce µG, which can be introduced in an ordinary laboratory, was newly developed to realize routine Ms measurement. Here, size of drop cansule was reduced to 30cm in diameter by introducing a magnetic circuit composed of an NdFeB permanent magnet.

[1] C.Uyeda et al: Jpn. Phys. Soc. Jpn. 79, 064709 (2010).

### **RO04**

D.g. Park and Hoon Song

korea atomic energy research institute, Korea

Ultra-low concentration detecting of biomolecules, has gain increasing interest since various protein biomarkers for cancer or chronic diseases are present at very low levels at the early stages of the disease development [1]. In this work, we present the results of the development of a sensor prototype working in a principle of label free for monitoring the effects on the immobilization of the biomolecular using giant magnetoimpedance. Amorphous ribbons were tested in different biomoleculars. The MI response of the sensive elements made from ribbons was measured with a various biomolecular in order to demonstrate the capacity of monitoring the immobilization of biomolecule. To demonstrate the sensitivity and specificity of this GMI sensor platform for biomolecule detection and quantification, we have used this platform to successfully detect and quantify the biomolecules. To apply to the real biomarker and verify the sensitivity of GMI-based magnetic sensor, we used the biomarker scheme using the applied the real protein on the magnetic sensor surface. The capture antibodies (P21 mono antibody) are first immobilized on the positive magnetic sensor surfaces. The advantage of this approach consists in the possibility to design a biosensor with enhanced sensitivity and the application for detection of small amounts of biomolecules.

[1] M. Gomez and G. Silvestri, "Lung cancer screening," Am. J. Med. Sci., vol. 335, pp. 46-50,

### **RO05**

#### Ac Calorimetry under Pulsed High Magnetic Field

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Applied quantum physics, kvushu university, Japan

<sup>2</sup> Institute for solid state physics, the university of tokyo, Japan

Ac calorimetry measurement adapted for use of the pulsed high magnetic field will be presented. This technique is quite useful to investigate the thermal properties under high magnetic field beyond the steady magnetic field range. In addition, it is also possible to capture a transient phenomena under fast sweep magnetic field. Resultant examples will be presented in detail.

### **RQ06**

**Specific heat and thermal expansion of Sr<sub>1-x</sub>Ca<sub>x</sub>RuO<sub>3</sub>** Rasna Thakur<sup>1\*</sup>, Archana Srivastava<sup>2</sup>, Rajesh K. Thakur<sup>1</sup> and N.k. Gaur<sup>1</sup> <sup>1</sup>Department of Physics, Barkatullah University, Bhopal, India <sup>2</sup>Department of Physics, Sri Satya Sai College for Women, Bhopal, India

We have investigated the thermodynamic properties of perovskite ruthnate Sr1  $_x$ Ca $_x$ RuO $_3$  ( $0 \le x \le 1$ ) probably for the first time by means of Rigid Ion Model (RIM). The lattice contributions to the specific heat and thermal expansion of pure and Ca doped SrRuO<sub>3</sub> as a function of temperature (0 K $\leq$  T  $\leq$  1000 K) are reported. The systematic trend of variation of specific heat and the closer agreement with the experimental data reveal the suitability and appropriateness of RIM, for Sr<sub>1-x</sub>Ca<sub>x</sub>RuO<sub>3</sub>  $(0 \le x \le 1)$  perovskite ruthnates. The calculated lattice specific heat gives Debye temperature (478.4 K) which is consistent with experimental value (480 K). Our calculated value of 494.4 K for CaRuO3 is in very good agreement with the reported value of 495 K. The Atoms in Molecules (AIM) theory is used to determine the bulk modulus of these compounds. Further, the value of Gruneisen parameter is calculated which is within the range 2-3 as reported earlier for perovskite family. In addition, the results on cohesive energy ( $\phi$ ), molecular force constant (f), restsrahlen frequency ( $\upsilon$ ), Debye temperature( $\theta_D$ ) and gruneisen parameter ( $\gamma$ ) are also presented. Our results on cohesive and thermal properties revealed by using RIM reproduces well with the available experimental data

### RQ07

Voltage-current characteristics of superconductor-normal metal contact junctions measured by a picovoltmeter

Wan-seop Kim\*, Mun-seog Kim, Po Gyu Park, Kyu-tae Kim and Danbee Kim *KRISS, Korea* 

A picovoltmeter based on a DC SQUID has been developed for improvement of Josephson voltage measurement uncertainty to better than  $\Delta V/V = 50 \times 10$ -12. An output to input ratio of 108 is achieved using a proper current feed-back system. The performance of the picovoltmeter demonstrated by the Voltage(V)-Current(I) measurements on a Cu-wire with a diameter of 0.25 mm and a length of 10 mm shows that a voltage resolution is of about 10-12 V (pV) and a corresponding resistance resolution is order of 10  $\mu$ Ω. In addition, the current sensitivity of the picovoltmeter was found to be 2x10-7 A/Φ0, where Φ0 is the magnetic flux quantum. Here we present voltage steps induced by magnetic flux trap in the V-I curve measurements for a bulk cylinder-type Pb wire with a diameter of 0.25 mm. The step height was observed to change abruptly in an irregular manner, but in multiple of 6 pV. Moreover, dissipation behavior of the remnant voltage observed in the V-I curves for thick films of superconductor/normal metal/superconductor contact junctions will be discussed.

#### **RQ08**

# Development of high resolution cryogenic particle detectors using a magnetic calorimeter

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Micro-calorimeters operating at low temperatures have become important tools in many aspects of science because of their high energy resolution much beyond the limit of semiconductor based detectors. These sensitive detectors measure the temperature dependent properties of sensor materials such as their resistance or magnetization. In metallic magnetic calorimeters, a small concentration of erbium doped in gold host provides a paramagnetic system, in which the magnetization is inversely proportional to the temperature. A metallic film of the sensor material is fabricated on a superconducting meander-type pickup coil connected to an input coil of a SQUID current sensor. An absorber, another gold foil, is employed to detect incident alpha particles. The gold foil absorber and the magnetic temperature sensor are thermally connected by gold bonding wires. The kinetic energy of each particle is converted into a temperature increase of the absorber/sensor assembly. In the present report, the recent progress on high resolution alpha spectrometers using these detectors is discussed together with an application to radionuclide analysis.

# RQ09

#### **High pressure inductive measurements using microcoils in anvil cells** Swee K. Goh<sup>1</sup>\*, Thomas Meissner<sup>2</sup>, Patricia Alireza<sup>1</sup> and Juergen Haase<sup>2</sup> <sup>1</sup> University of Cambridge, United Kingdom

<sup>2</sup> University of Leipzig, Germany

Diamond or Moissanite anvil cells are routinely employed to reach pressures higher than 35 kbar. However, due to the presence of the gasket and the limited sample space, inductive type measurements have been challenging. To overcome these challenges, we place a microcoil inside the gasket hole of the anvil cell. With an increase in the filling factor and the close proximity between the sample and the coil, this arrangement has so far enabled us to perform AC susceptibility [1], the de Haas-van Alphen effect [2], tunnel-diode oscillator [3] and Nuclear Magnetic Resonance [4,5] experiments. Technical aspects associated with the preparation of these pressure cells will be presented, with particular emphasis placed on recent NMR experiments performed using this technique.

[1] P. L. Alireza et al., Rev. Sci. Instrum. 74, 4728 (2003) [2] S. K. Goh et al., Curr. Appl. Phys. 8, 304 (2008) [3] S. K. Goh et al., arXiv:1107.0689 (2011) [4] J. Haase, S. K. Goh et al., Rev. Sci. Instrum. 80, 073905 (2009) [5] T. Meissner, S. K. Goh et al., Phys. Rev. B 83, 220517(R) (2011)

#### RQ10

#### Microwave synthesis and characterization of the series of co<sub>4</sub>-xfexsb<sub>12</sub> high temperature thermoelectric materials

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An alternative route has been followed for a rapid synthesis of skutterudites as it was suggested by Biswas et al. (1) using microwave assisted synthesis. Microwave heating and sintering (involve microwaves only at 2.45 GHz and 915 MHz frequencies) is fundamentally different from the conventional sintering, which involves radiant/resistance heating followed by transfer of thermal energy via conduction to the inside of the body being processed. Microwave heating is a volumetric heating involving conversion of electromagnetic energy into thermal energy, which is instantaneous, rapid and highly efficient. The use of microwave energy for materials processing has major potential and real advantages over conventional heating. These include: - Time and energy savings. - Rapid heating rates (volumetric heating vs. conduction). - Considerably reduced processing time and temperature. - Fine microstructures - Lower environmental impact. Using the microwave-assisted synthesis, here we report our effort to synthesize the series of Co4-xFexSb12 using this novel approach, which gave high quality materials with little or no impurity in a fraction of time compared to the conventional synthesis. We will present data for structural approach, Supported by the project NEXTEC of the EU

1. K. Biswas, et al. MRS Bulletin 46, pp. 2288-2290 (2011). 2288-229 2. R. Roy, et al. Nature 399, pp. 668-670 (1999).

#### RO11

# Development of SI-STM optimized for 3D(x,T,B) phase-diagram-wide FTSTS mapping on high Tc superconductors

Seokhwan Choi<sup>1</sup>, Jimin Kim<sup>1</sup>, Chanhee Kim<sup>1</sup>, Jaewook Kim<sup>1</sup>, Hwansoo Suh<sup>2</sup> and Jhinhwan Lee<sup>1</sup>\*

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We are constructing an SI-STM optimized for 3D(x,T,B) phase-diagram-wide QPI mapping on high Tc superconductors, Kondo systems and SCESs. It is based on Panstyle STM design for high magnetic field compatibility and yet has enough symmetry required for covering a wide temperature range with minimal drift. Also it has multiple sample storage at 4K for unlimited reuse of samples with various dopings. Its primary goal is to extend the SI-STM measurement coverage over the most significant regime of the 3D phase diagram for unique determination of the mechanisms of high Tc superconductivity using the fully-phased Green-function-based cross-sectional FTSTS analysis method.

# RQ12

#### Homemade microcalorimetry equipment, with magnetic fields up to 9 Teslas, for magnetocaloric measurements

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The magnetocaloric effect (MCE) has gains relevance as a cheap and environmentally friendly alternative to the current household vapor cycle refrigerators. The most common method for calculating the MCE of a material, due to its technical and mathematical practicality, is the application of Maxwell's thermodynamic equation to a set of magnetic isotherms so as to calculate the magnetic entropy change [1]. The other current alternative for this calculation is through the adiabatic temperature change. Unfortunately, even though a working magnetocaloric refrigerator should rely much more on an adiabatic (Brayton) cycle [2], such calculation is usually hindered by the lack of a proper measurement system that would allow for reliable specific heat measurements with an applied magnetic field. Our group has recently constructed such a device, using a 9 Tesla cryostat, allowing for liquid helium temperatures, and microcalorimetry chips from the company Xensor Integration. This equipment, besides the usefulness in magnetocalorics research, is an invaluable tool in the study and investigation of any form of phase transitions where the application of a magnetic field may play a significant role. We now report the overall characteristics of this equipment and further illustrate it with a few early measurements performed on it.

[1]Debye P, 1926, Some observations on magnetisation at a low temperature, Ann. Phys. 81: 1154-1160; [2] Tishin A.M, Spichkin Y.I, 2003, The Magnetocaloric Effect and its Applications, MPG Books Ltd., Bodmin, 2003;

### RQ13

# Nucleation and development of clustered state in hole doped manganites and cobaltites

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The origination of the ferromagnetic (F) metallic (M) clusters in a paramagnetic (P) state of the hole doped manganites was well established [1-3]. Study of this state in the hole doped cobaltites is the topic of the current research activity [4,5]. It is found by measurements of the nonlinear magnetic responses (the second and third orders) that an universality exists in the nucleation and development of the clustered states (CS) in these systems which proceed into three stages. On cooling, during the first stage, the F clusters nucleate at the preferential sites that are likely produced by variation in oxygen and doping stoichiometry. Second stage is characterized by a sharp increasing concentration of the isolated F clusters. A coalescence of the clusters into large-scale complexes containing some amount of the F domains is associated with the third stage that can end with formation of a percolative FM network. This scenario is demonstrated by using the data obtained on the single-crystalline magnanites Pr1-x CaxMnO<sub>3</sub> (x = 0.2, 0.25; the F insulators) and single-crystalline cobalitie La<sub>0</sub>, Sft<sub>0.3</sub>COO<sub>3</sub> (a F metal). A new effect is found in the latter, in which the CS in the P metallic regime is observed for the first time.

[1] M. B. Salamon, M. Jaime, Rev. Mod. Phys. 73 (2001) 583. [2] V.A. Ryzhov, A.V. Lazuta, I.D. Luzyanin et al., Zh. Eksp. Teor. Fiz. 121 (2002) 678. [3] A.V. Lazuta, V.A. Ryzhov, V. P. Khavronin et al., Funct. Mater. 17 (2010) 11 and Ref. therein. [4] C. He, S. El-Khatib, S. Eisenberg et. al., Appl. Phys. Lett. 95 (2009) 22511. [5] A. V. Lazuta, V.A. Ryzhov, A. I. Kurbakov et al., Solid State Phenom. 168-169 (2011) 457.

### RQ14

Measurements and analysis of core loss including higher harmonic induction waveforms using superposition principle and steinmetz's law Duhyung Yeon and Derac Son\* *Physics, Hannam university, Korea* 

Non-oriented electrical steels have been used in rotating electric machines. Magnetic induction waveforms of stator cores of induction motors have ac major hysteresis loop of driving frequency and ac minor loops due to the slip. Prediction of core loss including higher harmonic is one of important task for high efficiency induction motor design[1]. In this work, we have constructed core loss measuring system which could control waveform of magnetic induction the same as generated from waveform synthesizer using analog negative feed-back system and test specimen was used air flux compensated ring core. Using the developed measuring system, we can measured core losses of ac minor loops which depend on the peak amplitude and position in major hysteresis loop and theses core losses obeyed Steinmetz's law. Core loss including minor loops could be calculated superposition principle of the core loss from major hysteresis loop and the core losses from minor loops which come from higher harmonic induction.

 [1] D. Son "AC hysteresis loop measurement of stator-tooth in induction motor", IEEE Trans. on Magn., Vol.35, No.5, p.3931-3933(1999)

# RQ15

# Highly sensitive cantilever magnetometryin static and dynamic modes for micro-scale samples

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Recently, dynamic modecantilever magnetometriesusing an ultra-soft cantilever with attonewton dynamic force sensitivity arebecoming a very powerfultool for measuring the magnetic properties of micrometer-sized samples. In spite of its high sensitivity, the dynamic mode measuring a cantilever's resonance frequency shifthas a limitation that it requires anisotropic magnetic samples, andhigh or moderate external magnetic fields. Therefore, we suggest a new type of a cantilever magnetometry with no limitation on magnetic anisotropy and external magnetic field, a cantilever forcemagnetometrywhich measuresastatic displacement of a cantilever in a well-defined magnetic field gradient. Its principle and experimental details including the magnetometry 'sstructure are presented, which are accompanied by the results of measurement at T = 3.5 K on a small magnetic moment of  $1.84 \times 10-14$  A m<sup>2</sup> using both static and dynamic modes revealing the high sensitivity of the static mode.

### RQ16

**Pulsed high magnetic fields for synchrotron and neutron applications** Fabienne Duc\*, Xavier Fabreges, Paul Frings, Marc Nardone, Julien Billette, Jerome

Beard, Abdelaziz Zitouni and Geert Rikken Laboratoire National des Champs Magnetiques Intenses - CNRS Grenoble-Toulouse, France

As is well known, a magnetic field, together with temperature and pressure, is a very important and efficient thermodynamic parameter for the investigation of condensed matter. It can be used as an external variable to tune the ground and excited states properties of a given magnetic system. The last ten years have seen a growing interest for the combined use of pulsed high magnetic fields (up to 40 T) at synchrotron and neutron sources. Because of the larger flux of synchrotron x-ray sources compared to neutron facilities, the development efforts have been first focused on synchrotron methods. Neutron diffraction has complementary features to x-ray diffraction, and unique capabilities for studying microscopically the magnetic properties of materials. Here, we will present the various devices developed by the LNCMI. Toulouse to combine pulsed high magnetic fields with synchrotron and neutron techniques. These developments are performed in close collaboration with the ESRF. Grenoble for synchrotron applications and with the ILL and CEA, Grenoble for neutron diffraction. We will discuss the main limitations for the use of pulsed magnetic fields related to their low duty cycle. Those instrumental developments will be illustrated by some recent results

# RQ17

#### Characteristics of an SQUID system with a superconductive shield for biomagnetic measurements

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We have fabricated a magnetoencephalography (MEG) system having a superconductive shield helmet. By using the perfect diamagnetic effect of the superconductor (Pb), superconducting quantum interference device (SQUID) magnetometers inside the Pb enclosure can be shielded from external magnetic disturbances, and this effect can be used to enhance the signal quality of biomagnetic signals. Along the surface of the helmet, the shielding factor (SF) was from 28 to 57 dB, depending on the position of the SQUID from the edge. SF reduced rapidly as the distance from Pb surface increased. The shielded helmet-shape MEG system has a noise level of about 7 fT/Hz1/2 at 1 Hz, and 3 fT/Hz1/2 at 100 Hz. Auditory evoked brain signals were measured and were compared in the case of superconductive shield and without shield, and confirmed big improvement of the signal quality by using superconductive shield.

[1] Y. H. Lee, "A low noise multichannel magnetocardiogram system for the diagnosis of heart electric activity", Korean society of medical and biological Eng., Vol. 27, pp. 154-163, 2006. [2] R. H. Kraus, "First Results for a Superconducting Imaging-Surface Sensor Array for Magnetocardiography", Recent Advances in Biomagentism Proceeding, pp. 33-46, 1999. [3] D. B. Hulsteyn, " Superconducting Imaging Surface Magnetometery", Rev. Sci. Instrum, Vol. 66, pp. 3777-3784. [4] H. Ohta, "A 64-channel whole head SQUID system in a superconducting magnetic shield", "Supercon. Sci. Technol., 161, 2p. 762-765, 1999.

# **RQ18**

# Radio-frequency atomic magnetometer for sensitive susceptibility detection

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Alkali-metal magnetometers use the coherence precession of polarized atomic spins to detect and measure magnetic fields. Previously, the interaction between light and atoms allows magnetometers to detect and measure magnetic fields, with sensitivity of 10 TT/Hz [1]. Atomic magnetometers such as spin-exchange relaxation free (SERF) regime magnetometer and coherence population trapping (CPT) magnetometer have been designed to detect of varying magnetic fields mainly at low frequency. Therefore, detection of weak fields at high frequencies up to several megahertz requires an alternative detection method. To observe the Nuclear magnetic resonance (NMR) signals at several tens of kilohertz, RF (Radio-Frequency) atomic magnetometer can be used. RF magnetometer has a fundamental sensitivity limit on the order of 0.01 TT/Hz [2, 3]. In RF magnetometer, the oscillating magnetic field causes the polarization of each spin to totate along the pumping axis. The resonant RF field eventually makes all of the processing spins coherent. In this study, the measured SNR corresponded to a sensitivity of 50 TT/Hz at 25 kHz and such results will be useful to atomic NMR detector.

J. C. Allred, R.N. Lyman, T.W. Kornack and M.V. Romalis, Phys. Rev. Lett. 89, 130801 (2002) [2]
 S.-K. Lee, K. L. Sauer, S. J. Seltzer, O. Alem, and M. V. Romalis, Appl. Phys. Lett. 89, 214106 (2006) [3]
 I.M. Savukov, S.J. Seltzer and M.V. Romalis, Journal of Magnetic Resonance 185 (2007) 214-220

### RQ19

# Cancellation coil allows precision magnetic measurements with strong magnetization field inside a shielded environment

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A great number of precision magnetic measurements can benefit significantly from, or in some cases, even require strong prepolarization fields ( $B_p$ ) and magnetically shielded environments. We devised a cancellation coil (CC) to neutralize the  $B_p$  on electrically conductive shield walls that may otherwise induce currents on the walls to produce a lingering transient residual field ( $B_p$ ) inside the shielded environment and disrupt the measurement operations. The CC was designed using the inverse problem method to effectively neutralize magnetic fields generated on the shield walls by the  $B_p$  coil. The implemented CC was evaluated by measuring the resulting  $B_p$  after strong pulsed  $B_p$  using a fluxgate magnetometer at different magnetometer positions and cancellation coil currents (ICC). We have conducted multi-mode component analysis on the  $B_p$  measurements to reveal two dominant components, where the component with shorter time constant from the current induced on the ceiling and floor of the MSR. The analysis also allows optimization of ICC for each of the top, side, and bottom sections of the CC to enable significantly easier fine-tuning of individual sections of the CC to enhance CC performance.

Seong-min Hwang, Kiwoong Kim, Chan Seok Kang, Seong-Joo Lee, and Yong-Ho Lee, Appl. Phys. Lett. 99, 132506 (2011) Seong-min Hwang, Kiwoong Kim, Chan Seok Kang, Seong-Joo Lee, and Yong-Ho Lee, submitted to J. Appl. Phys. 111, 083916 (2012)

#### **RQ20**

#### Development of a SQUID based ultra-low-field MRI system

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We obtained 2-dimensional Magnetic Resonance Imaging (MRI) image of a phantom in fields of a few microtesla by using a dc Superconducting QUantum Interference Device (SQUID) based MRI system. Microtesla MRI technique is a challenging application based on SQUID technology. The high sensitivity of the SQUID magnetometer enables measurement of very weak magnetic resonance signals even for the low Larmor frequency in microtesla fields. Measuring the nuclear magnetic resonance (NMR) signals at such a low field gives many benefits like development of an open-type low-cost surgery-monitoring MRI system, high-contrast cancer detector, etc. In this presentation, we introduce experimental details for ultra-low-field (ULF) MRI.

### RQ21

# **26 T+ steady magnetic field for neutron science at HZB Berlin** P. Smeibidl, Karel Prokes\*, H. Ehmler, O. Prokhnenko and A. Tennant

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The Helmholtz Zentrum Berlin is building in collaboration with the National High Magnetic Field Laboratory, Tallahassee, FL, USA a new series-connected hybrid magnet system for neutron scattering experiments [1]. Two conical endings of 30° opening angle enable neutrons to hit the sample and to be detected either in the forward or backward direction. To achieve a maximum field with the present technology, a 13 T superconducting Nb<sub>3</sub>Sn cable-in-conduit coil is combined with resistive insert coils of 13 T to 19 T, depending on electric power, to give a maximum of 26 to 32 T. The magnet that provides a 50 mm diameter room temperature bore in horizontal orientation will be permanently mounted at the dedicated time-of-flight instrument ExED 76 m away from the neutron source. The contribution describes the most important design features of the system (the 20 kA, 8 MW power supply, the helium refrigerator system for cooling of the resistive coil), the status of the fabrication of the components, the outline of the building for the technical infrastructure and the status of the installation of the most important components.

[1] P. Smeibidl et al., J. Low Temp. Phys. 159, 402 (2010).

#### RQ22

#### A simple method for measuring blocking temperatures Mansor Hashim<sup>1</sup>, Ghazaleh Bahmanrokh<sup>2</sup> and Ismavadi Ismail<sup>1</sup>

<sup>1</sup> Institute of Advanced Technology, Universiti Putra Malaysia, Malaysia
 <sup>2</sup> Physics Department, Faculty of Science, Universiti Putra Malaysia, Malaysia

We describe a new method for measuring a blocking temperature, much simpler than that involving a SQUID. Nanoparticles of Co, Co coated with Au,Co-Pt and Co-Pt coated with Au were prepared by the reverse micelle microemulsion method. The four magnetic powder specimens were each pressed to yield a disc. A Hall-effect probe blade surface was stuck intimately against one flat surface of the disc. Properly analysed, this arrangement allows the probe to detect the normal component, BN, projecting from the disc since, at the interface, BN (air) = BN (disc). Use of this relation led to the determination of a magnetic moment near the disc surface. By simply cooling this arrangement of disc + probe in a probe station to 15K, a graph of the magnetic moment was clearly peaked at a temperature which, unmistakably, was the blocking temperature, TB. For the four different samples TB lies between 40K and 45K.

#### RQ23

#### Optimization of operation condition of orthogonal fluxgate sensor fabricated with Co based amorphous wire

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Agency forDefense Development, Korea

Orthogonal fluxgate (OFG) sensors have an additional advantage comparing with conventional fluxgate sensors in their simple structure [1-2]. Because the sensitivity of an OFG is typically proportional to its winding number N, cross section area of its magnetic core A, and operation frequency f the f should be as high as possible to prevent reduction of the sensitivity in the case of small-sized OFG sensors in which the N and A are restricted. The f is limited due to the LC resonance occurred by inductance and stray-field capacitance of the pickup coil around the magnetic core and transmission line, coaxial line in this study, as well as performance of the signal conditioning circuit in common. And the sensitivity could be also reduced according to increase of operation frequency due to the skin effect if the sensor is operated in the frequency range of near MHz. We present how to optimize the operation condition including frequency of the OFG fabricated with a CoFeSiB amorphous wire with the diameter of 100 µm in this paper. The LC resonance frequency according to external capacitances was calculated with an equivalent circuit and FEM analysis to calculate the magnetization distribution in a high frequency.

[1] E. Paperno, E. Weiss and A. Plotkin, IEEE Trans. Magn., 44, 4018 (2008). [2] I. Sasada, J. Appl. Phys., 91, 7789 (2002).

# RR01

#### Preparation of γ-Fe4N and ε-Fe3N particles with high magnetization for electromagnetic wave absorption applications

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High-Ms ferromagnetic iron nitrides have vast applications on microwave absorbing and magnetic recording. In this study, we successfully used a simple nitrogen treatment to prepare cubic  $\gamma$ -Fe<sub>4</sub>N and hexagonal  $\epsilon$ -Fe<sub>3</sub>N powder. Both kinds of powders are synthesized by the nitriding process of the commercial micro-scale Fe and Fe<sub>2</sub>O<sub>3</sub> powders, at 550°C for 5 hours during a flowing mixture of NH<sub>3</sub> and H<sub>2</sub> gas. Our pure  $\gamma$ -Fe<sub>4</sub>N powder exhibits a high saturation magnetization of 15.03 kG at room temperature, which is slightly lower than that of bulk Fe (22.0 kG @ 4 K) but is considerably higher than that of  $\gamma$ -Fe<sub>2</sub>O<sub>3</sub> (4.0 kG @ 4 K). The  $\epsilon$ -Fe<sub>3</sub>N powder including a slight amount of  $\gamma$ -Fe<sub>4</sub>N was also synthesized by heating the commercial Fe<sub>2</sub>O<sub>3</sub> powder with the same procedure. The chemical stabilities of both  $\gamma$ -Fe<sub>4</sub>N and  $\epsilon$ -Fe3N powder are superior to that of the pure Fe and have high level of mechanical hardness. The mircostructures and electron spin resonances of our iron nitrides were also studied

This work was supported by the National Science Council, Taiwan under the NSC 98-2112-M-007-016-MY3 project.

#### **RR02**

#### Methods for determining the quality of magnetic fluids Viorica Chioran

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The conversion parameter of the magnetic properties of magnetic fluids is a physical quantity describing the quality of those fluids from a technological point of view. This study determines the value of this parameter for three samples of magnetic fluid. Many technical applications require magnetic fluids with specific magnetic properties such as very high magnetization at saturation point or rheological properties involving high fluidity (very low viscosity). The sample with bigger density (1070 kg/m3) also has a larger solid volume fraction (0, 059). Increasing the solid part/fluid volume increases the density and viscosity of that fluid. A magnetic fluid with a larger magnetic part/ volume (0,031) has a higher magnetization at saturation point (0,0138 A/m). The parameter always has sub-unitary values because inside the volume of the magnetic fluid, the percentage of the solid part is larger then that of the magnetic part. Responsible for this is the nonmagnetic layer at the surface of the particles. For the fluid samples studied, regardless of the magnetic and physic diameter, the nonmagnetic layer of particles has the same value: a = 0,83 nm. Good quality of a magnetic fluid is indicated by a high value of the conversion parameter.

#### RR03

#### A study on magnetic fluid viscosity

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This paper shows the results of a study of the parameters that the viscosity of magnetic fluids depends on, in the absence of magnetic fields. Viscosity's and shearing tension's dependence on the rheologic component and shearing speed was verified for several samples of kerosene based magnetic fluid with the same dimensional distribution but various concentrations, at different temperatures. Experimental results show one situation where the sample P1 has a non-Newtonian character over a small interval of shearing speeds. An increase in fluid viscosity indicates that bonds between particles are formed inside the fluid. As shearing speeds increase these bonds disappear. Samples P and P<sub>2</sub> show Newtonian character over the shearing speed interval where viscosity remains constant. By comparing viscosity values of samples at different temperatures and shearing speeds, decreasing viscosity while increasing temperature is observed due to the activation energy of viscous flow. A linear increase of shearing tension with the increase of shearing speed was observed for the studied samples at different temperatures. The slope of this line can be used to determine the dynamic viscosity's value. For the considered samples, viscosity strongly depends on several factors: density, concentration, temperature (by Arrhenius's Law) and shearing speed (by Newton's equation).

### **RR04**

Magneto-motive force and torque analysis of squirrel cage induction motor with rotor or stator faults

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<sup>2</sup> Electrical Engineering, Chungnam National University, Korea

Torque characteristics of high voltage induction motor are simulated under the condition of under voltage and broken rotor bars. The motor is high voltage squirrel cage induction type for driving force of pumps and valves in nuclear power plant. In the simulation, a rotor with health and broken bars, stator winding deterioration and voltage drop 25% alone and combination with bar faults were used. Then the results were analyzed. The effects of rotor faults on the torque-speed curve were compared according to the number of the rotor faults. The results revealed that voltage drop and the number of the rotor fault at high speed span is directly proportional to its effect on the curve, but the curve indicated the opposite at the low speed span due to the rotor impedance increase. In stator, winding faults made a distortion of the phase magneto-motive force and reduction of co-energy at the gap which cause vibration, harmonics and total power decrease.

### **RR05**

# Evaluation of materials degradation of ferromagnetic steels for various magnetized states using a hysteresis scaling law

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We have examined a hysteresis scaling behavior of magnetic minor hysteresis loops for ferromagnetic low carbon steels under DC-biased magnetic field and/or remanent magnetization state. For all minor loops with dominant contribution of irreversible Bloch wall motion, a scaling relation between the hysteresis loss and maximum magnetization shows the same curve from very low to high magnetization regime, suggesting the invariant pinning mechanism for the irreversible wall motion. In the intermediate magnetization regime, the relation follows the usual Steinmetz law with a power-law exponent of about 1.5, whose coefficient increases with defect density. The method using a hysteresis scaling can be a useful technique of evaluating materials degradation of ferromagnetic steels in an unknown magnetic state, which is commonly encountered on on-site measurements.

### **RR06**

#### Theoretical design of magnetic energy harvesting module

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Energy harvesting is a key technology to build a sustainable society. In previous paper, we focused on a magnetic power-line noise as a reusable energy. In this paper, we present a theoretical design method of magnetic energy harvesting module. This module consists of an air-core coil and resonant capacitor. With a simple RLC circuit model, we derive a equation of the harvesting energy as a function of coil size. In order to demonstrate the magnetic field, a uniform magnetic field is generated by our developed coil system. From the experimeted results, we successfully demonstrated the energy harvesting of 100 mW from a magnetic field of 0.09 mT at 60 Hz. This value is good in agreement with the estimated results. The harvested energy is proportional to the square of the magnetic flux density. However, ICNIRP2010 provide a guide line for human helth that an acceptable level in a public space is 0.2 mT at power-line frequency. We also reveal the required coil size to harvest a demanded energy.

### **RR07**

### A study on detecting and determining the shape of small axial cracks by using magnetic flux leakage in ndt system of pipe

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There is the source of energy that is used widely all over the world, which is oil and natural gas. Oil and gas is supplied for industries or home consumers through pipelines. Most pipelines are laid to underground and exposed the external environment. By these conditions pipelines have some bad effects like the metal loss, corrosions and any other damages. For preventing accidents, these defects have to be detected by the Non-destructive testing(NDT). The magnetic flux leakage(MFL) method among the NDT methods is suitable for testing pipelines which the magnetic characteristic of has high magnetic permeability. The system that MFL method applies to is called the pipeline inspection gauge, MFL PIG. The previous MFL PIG has high performance for detecting the metal loss and corrosions. But MFL PIG cannot detect the cracks which occurred by the difference of pressure between the inside and outside of pipelines and the shape of is long and very narrow. Cracks occures frequently in the pipelines and the risk of the accident from the cracks is higher than that from the metal loss and corrosions. The circumferential MFL(CMFL) PIG performs magnetic field circumferentially and can maximize the magnetic flux leakage at the cracks.

#### **RR08**

Control of working temperature of isothermal magnetic entropy change by hydrogen absorption into  $La_{0.8}Nd_{0.2}(Fe_{0.88}Si_{0.12})_{13}$  for magnetic refrigerant

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Hydrogen absorption into La<sub>0.8</sub>Nd<sub>0.2</sub>(Fe<sub>0.88</sub>Si<sub>0.12</sub>)<sub>13</sub> has been investigated to obtain large magnetocaloric effects (MCEs) near room temperature for applications to magnetic refrigeration. La(Fe<sub>0.88</sub>Si<sub>0.12</sub>)<sub>13</sub> exhibits large MCEs just above the Curie temperature T<sub>c</sub> = 195 K because of the itinerant-electron metamagnetic (IEM) transition [1]. The MCEs due to the IEM transition are enhanced by the partial substitution of Nd for La [2]. It was confirmed that T<sub>c</sub> of La<sub>0.8</sub>Nd<sub>0.2</sub>(Fe<sub>0.88</sub>Si<sub>0.12</sub>)<sub>13</sub> is increased from 193 K to 318 K by hydrogen absorption with keeping the IEM transition because the unit cell volume is expanded. Note that the isothermal magnetic entropy change  $\Delta$ Sm due to the IEM transition in a magnetic field change of 2 T for La<sub>0.8</sub>Nd<sub>0.2</sub>(Fe<sub>0.88</sub>Si<sub>0.12</sub>)<sub>13</sub>H<sub>1.1</sub> was evaluated to be about -23 J/kg K at T<sub>c</sub> = 288 K, which is larger than  $\Delta$ Sm = -19 J/kg G of the former also is larger than 160 J/kg of the latter. Accordingly, it is concluded that the partial substitution of Nd improves the MCEs near room temperature.

[1] A. Fujita et al., Phys. Rev. B 67 (2003) 104416. [2] S. Fujieda et al., Mater. Sci. Forum 561-565 (2007) 1093.

#### **RR09**

**Characteristic analysis of induction motor for electric vehicle according to electric loading and magnetic loading** Ki Young Sung<sup>1</sup> and Ki-chan Kim<sup>2</sup>\*

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In the paper, an inductive motor for electric vehicle that was designed for compact car modification was selected as the basic model, and its performance change was examined according to the change in the electric and magnetic loading related to the rotor slot. The characteristic changes according to the change in the loading were analyzed for two speed ranges: the low speed (near the base speed) and the high speed (maximum speed with field weakening control). To provide the characteristic analysis conditions according to the change in the slot shape of the reference model, the rotor bar depth was changed with the rotor teeth width fixed to vary the electric loading, and the rotor teeth width was changed with the rotor bar area fixed to vary the magnetic loading. The characteristics of the inductive motor were analyzed via the analysis of the characteristics according to the parameters of the equivalent circuit based on the magnetic circuit method and via the electromagnetic analysis based on the finite element method to consider the nonlinear characteristics of the motor parameters.

[1] ZZhang, Eprofino, a. tenconi, Improved design for electric vehicle induction motors using an optimisation procedure'. Electric Power Applications, IEE Proceedings, vol. 143, pp.410-416, 1996 [2] T. Wang, P. Zheng, and S. Cheng, Design Characteristics of the induction motor used for hybrid electric vehicle'. IEEE Tracnsaction on Magnetics, vol. 41, pp. 505-508, 2005 [3] I. Boldea and S. A. Nasar, The induction machine handbook'. Boca Raton, (RC Press, 2002

### **RR10**

#### Torque characteristics of interior permanent magnet synchronous motor for electrical hydraulic power steering system

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In the paper, a 12 V-DC interior permanent magnet synchronous motor for electric hydraulic power steering (EHPS) was designed for the purpose of improving its torque performance. The EHPS is a steering system that uses the battery power of the vehicle, which is weaker than the power of the hydraulic power steering system. Because of the spatial limitation, the IPMSM, which has high-power characteristic and excellent performance within the field weakening area at a high rotating speed, is advantageous. To improve the torque performance, the average torque was increased by increasing the rotor size. In addition, because the torque ripple is detected as vibration when the steering wheel is operated, the shape of the barrier was modified to change the inductance element of the axis and reduce the torque ripple. The optimally designed model characteristics were analyzed using the finite element method to comparatively analyze the parameters. It is expected that this study will be important reference data to understand the characteristics of the IPMSM as well as of the EHPS.

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#### **RR11**

# Edge auxiliary teeth design of stationary discontinuous armature PM-LSM with concentrated winding

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Recently, in order to resolve the problem of high cost, we have proposed the stationary discontinuous armature PM-LSM in which the armature is engaged only when accelerated and decelerated operations are necessary, when PM-LSM is used with long-distance transportation systems in factories [1]. As the armatures are arranged discontinuously in the PM-LSM, the number required is dramatically decreased. In this case, several stator blocks can be arranged at certain intervals within the overall running route. The space between the blocks without the stator is called the free-running section, and the mover in this section drives under its own inertia. However, the stationary discontinuous armature PM-LSM contains the edges which always exist as a result of the discontinuous arrangement of the armature [2]. For this reason, cogging force generated between the "entrance end" and the "exit end" has become a problem. Therefore, we have examined the edge cogging force by installing the auxiliary teeth at the armature's edge in order to minimize the cogging force generated when the armature is arranged discontinuously. We obtain the edge cogging force by unsequence analysis with a FEM and also by adjusting the width, height and the length of pitch of auxiliary teeth.

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#### **RR12**

#### Inductively coupled LC resonators as displacement sensor

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In this paper, we present novel displacement sensor composed of a pair of LC resonators. The LC resonator was formed by serial connection of 300-turn Cu wire wound on a rectangular NiZn ferrite ( $20 \times 3 \times 4$  mm) core and SMD capacitor ( $3.2 \times 1.6$  mm). The resonance frequencies of two resonators are adjusted to be identical with each other. By this, we could obtain the sensor output linearly dependent of distance between two LC resonators, because the mutual inductance is proportional to distance between two resonators in a few centimeters. The sensor was designed by a sequence of magnetic field simulation with a FEM tool (COMSOL) and circuit calculation with a SPICE. The sensitivity estimated by S = fr/D was about 480 kHz/m, where, fr and D are resonance frequency and distance between two resonators, respectively. The proposed sensor would be a promising displacement sensor when it takes an active part in an environment surrounded by serious electrical/magnetic noise, since it could be operated in extremely narrow frequency range.

# **RR13**

#### Magnetic NDE for sensitization of Inconel 600 alloy

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Inconel 600 alloy, Ni base alloy, is used as steam generators in nuclear power plants and sensitization occurs along grain boundary with heat treatment on this material. The sensitization causes a degradation of corrosion resistance and generates stress corrosion crack [1], [2]. Therefore nondestructive evaluation (NDE) technique for sensitization of Inconel alloy is quite important so as to keep the integrity of plants. Inconel 600 alloy has usually paramagnetic property, however, it shows ferromagnetic properties along grain boundary when sensitization occurs: this means NDE using magnetism is possible. Thus, in this study, Inconel 600 alloys were heat treated at 873 K from 0 to 400 hours and their magnetic properties were investigated in detail. Magnetization increases with increasing time of heat treatment and takes a maximum. On the other hand, coercivity decreases with the increase in time. We confirmed that only characteristics at the grain boundaries changes ferromagnetic by MFM observation. As trial for industrial application, heat treated Inconel 600 alloy has scanned by magnetic field sensor, and the changes of magnetization nondestructively were successfully obtained. The result indicates the feasibility of magnetic NDE for sensitization of Inconel 600 alloy

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#### RR14

#### Comparison of characteristics between the PM synchronous motor and the induction motor for electric vehicle

Mi-jung Kim, Ik-sang Jang, Ki-doek Lee, Jae-jun Lee, Jeong-ho Han, Tae-chul Jeong and Ju ${\rm Lee}^\ast$ 

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This paper presents the characteristics over a wide speed range of the interior permanent magnet synchronous motor(IPMSM) and the induction motor(IM) for an application of electric vehicle drive. In generally, IPMSM have higher performance than IM. The performance: torque, efficiency, power factor etc. are compared at rated speed of 2,843rpm and maximum speed of 10,000rpm. The induction motor operating constant speed is not need to consider the control property; however, the adjustable speed drive applications such as traction motor are should be deliberated for control. In this paper, the characteristics of IPMSM and IM controlled by space vector PWM are illustrated and verified by experiment. In case of IPMSM, the efficiency map was made up through driving simulation considering the actual control. It was carried out MTPA control until constant torque operating region and applied the field weakening control under constant power region. In case of induction motor, V/f control maintained a constant ratio of voltage versus current was carried out until constant torque operating region and applied the constant voltage variable frequency control under constant power region. As illustrated in Figure 1, overall efficiency of IPMSM has high level than IM about 5%.

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#### RR15

# Comparison of simultaneously measured pulse waveforms from both hands using permanent magnet-hall pulsimeter sensors

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Two radially arterial pulses of dual hands using the prototype of a clamping clip pulsimeter equipped with permanent magnet and Hall device are compared and analyzed. The phase difference of two pulse wave signals is dominantly presented from the simultaneous measuring clinical pulse wave signals for twenty two male participants at their 20's. It is possible to analyze that the fast and slow pulse wave for right hand and left hand depend on the muscle property of arms rather than the total length of blood vessel due to cardiovascular circulatory system.

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# **RR16**

# A study on new permanent magnet configuration for high thrust density in permanent magnet synchronous linear motor

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#### **RR17**

#### Design of stationary pole pieces in a coaxial magnetic gear

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Magnetic gears have been proposed by unique advantage compared with mechanical gears, such as contactless power transmission. In addition, they can limit a torque under overloaded condition and have excellent performances on noise, vibration and reliability [1]. However, their use in industries has been very limited due to poor torque density, low transmission ratio and transmitted toque performances. Of magnetic gears, a coaxial magnetic gear is a common configuration and extensively studied for its relative high performances [2]. It is consist of inner rotor, outer rotor and stationary pole pieces in between the rotors. The stationary pole pieces have an important role in torque transmission by modulating the air gap flux density distribution adjacent to inner and outer rotor. However, they also generate unwanted harmonics resulting in torque ripples and vibration. This paper proposes variants of stationary pole pieces to gain high performances in torque transmission. It also has mechanical and manufacturing advantages in transmitting torque by a structure combining each pole pieces. The finite element analysis shows the produced torque has reduced harmonic components in both sides of rotors.

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### **RR18**

# Design techniques for reducing torque ripple in permanent magnet flux switching motor

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Permanent magnet flux switching motor (PMFSM) is a novel double salient machine which employ the PM instead of the field winding for excitation and contains only one set of armature winding so has the advantage of low cost power converter and extreme high efficiency. It has a potential for taking place of BLDC in some applications, such as fans and air ventilation machines. Due to the different structure and operation principle from traditional permanent magnet machine, the generated torque ripple in PMFSM is very critical and rather unique compared to the common PM machines. In this paper, the design techniques of reducing the torque ripple in PMFSM has been presented. Different kinds of design approaches have been presented to reduce the torque ripple with comparing the mean output torque. Analytical method is first presented to investigate the influence and determine the design parameters thus reduce the computational effort. Then FEA simulation is employed to validate the analysis above. Finally, time stepped FEA simulation is being employed to investigate the performance with different approaches adopted.

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# **RR19**

# The analysis of line-start permanent magnet machine with saliency ratio

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# **RR20**

Analysis of non-linear characteristics of linear compressor Park Daegeun and Cho Yunhyun\* Dong-A University, Korea



### **RR21**

#### Torque harmonics and reduction design characteristics of induction motor for electric vehicle propulsion

Kyung-won Jeon<sup>1</sup>, Seungho Lee<sup>1</sup>, Yong-jae Kim<sup>2</sup> and Sang-yong Jung<sup>1</sup>\* <sup>1</sup> School of Electronic and Electrical Engineering, Sungkyunkwan Univ., Korea <sup>2</sup> Dept. of Electrical Engineering, Chosun University, Korea

This paper deals with torque harmonic characteristics of induction motor for electric vehicle (EV). For calculate phase stator harmonic leakage flux of a three phase induction motor, an analysis method is determine relationships with stator teeth design. In particular, torque harmonics including EMF harmonics at whole speed ranges are analyzed by the finite element method (FEM). Torque harmonic components are identified with Fourier fast transform (FFT). Torque harmonics are generated due to the different phase windings interaction and they can be reduced by applying a stator teeth design. The electromagnetic field and corresponding harmonics of an electrical machine are calculated using the FEM, the harmonic winding factors show how the torque harmonics can be reduced significantly by good stator teeth design. The novel design methods to compensate the specified torque harmonics are proposed. Then, its effectiveness is clarified according to the representative control strategies for induction motor such as maximum torque per ampere (MTPA) and flux-weakening control.

[1] Lee S, Kim Y.J., Jung S.-Y., "Numerical Investigation on Torque Harmonics Reduction of Interior PM Synchronous Motor With Concentrated Winding", IEEE Trans on. Magnetics, vol.48, no.2, pp.927-930, Feb. 2012, [2] Yamazaki K., Statuki A., Ohto M., Takakura T., Nakagawa S., "Equivalent Circuit Modeling of Induction Motors Considering Stray Load Loss and Harmonic Torques Using Finite Element Method", IEEE Trans on. Magnetics, vol.47, no.5, pp.986-989, May: 2011, [3] Bianchi N., Alberti L., "MMF Harmonics Effect on the Embedded FE Analytical Computation of PM Motors", IEEE Trans on. Industry Applications, vol.46, no.2, pp.812-820, Mar-Apr. 2010.

# RR22

#### Numerical analysis on iron loss and pm loss of permanent magnet synchronous motor considering the carrier harmonics

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In this paper, decisive influences of the inverter harmonic on iron loss characteristic for PMSM (Permanent Magnet Synchronous Motor) are numerically investigated with the FEM (Finite Element Method). Particularly, analysis is carried out based on a multi-layer buried PMSM, which is newly designed for EV drive system and its rotor contains ferrite magnets. First of all, not only the design criteria and the specification of the designed PMSM but also the features of the motor like torque ripple, Back-EMF are represented. In addition, the remarkable iron loss analysis result, that is obtained using 3-D FEM considering skewed stator structure, is covered in priority. Then its each harmonic component is identified using the FFT (Fourier Fast Transform) and that is compared with iron loss characteristic in case of sinusoidal input current which is not include harmonic component. Finally, the efficiency at each operating conditions using loss analysis result, reality real input current values implemented by testing performance of the manufactured model.

[1] Z.Q.Zhu, Y.S.Chen, D.Howe, 'Iron Loss in Permanent-Magnet Brushless AC Machines Under Maximum Torque Per Ampere and Flux Weakening Control', IEEE Trans. on Magnetics, vol. 38, No. 5,pp. 3285-3287, Set. 2002. [2] Katsumi Yamazaki, Shinjiro Watarim 'Loss Analysis of Permanent-Magnet Motor Considering Carrier Harmonics of PWM Inverter Using Combination of 2-D and 3-D Finite-Element Method', IEEE Trans. on Magnetics, vol. 41, No. 5,pp. 3285-3287, May: 2005. [3] G. Ombach, J. Junak, 'Torque ripple optimization of skewel IPM motor for field weakening operation', Electrical Machines and Systems (ICEMS) 2011, pp. 1-7, Aug. 2011.

### RR23

# Defect depth estimation based on the analysis of interference defects on the underground gas pipelines

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Underground gas pipelines have corrosions and deformations due to both the natural causes such as moisture and ground pressure and the artificial causes such as faulty construction, etc. So, periodic pipeline inspection is essential. In NACE (national association of corrosion engineers) standard, it is advised to replace the pipeline when a defect's depth is over 50% of the pipeline's thickness. Therefore, the study to analyze the defect's depth and its progress is meaningful. Especially when defects more than two are closely located, the interference defects on underground pipelines are analyzed and their depths are estimated effectively. Experimental results show that the proposed method can effectively estimate the depth of the interference defects.

### RR24

# Comparison on electromagnetic losses of super high speed PM motor/

generator with slot and slotless stators Jin Hak Jang<sup>1</sup>, Jian Li<sup>2</sup> and Yun Hyun Cho<sup>1</sup>\* <sup>1</sup> Dong-A University, Korea <sup>2</sup> Dong-A University, China



# **RR25**

# A study on Voltage-PWM control of switched reluctance generator at low speed

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This paper is a study on switched reluctance generator's control method----Voltage-PWM. It can be used to regulate the current to a desired value by varying the average applied voltage in the switched reluctance generator . The advantages of Voltage-PWM control at low speed are more than high speed which may make the application of PWM difficult. The switched reluctance machine can operate as a generator as well as a motor by simply changing the firing angles. Some main control methods we often used are Chopped Current Control (CCC), Angular Position Control (APC) and PWM Control . The controller could regulate control variables, such as current level, tum-on and turn-off angles, supply voltage and so on. Different control methods will be carried out by regulating some variables in different conditions. For example, at low speed, the d.c link voltage is controlled by varying the set-point current or the duty-cycle and at high speed by varying the control angles. Because of excitation current is closely linear relation to duty cycle, so many applications determine this control method.

[1]Miller, T.J.E. Electronic Control of Switched Reluctance Machines. MPG Books Ltd, Bodmin, Cornwall. 2001. [2]K.Kannan, Dr.S.Sutha. The 4th International Power Engineering and Optimization Conference (PEOCO2010), ShahAlam, Selangor, Malaysia.23-24 June 2010. [3] Jianzhong Sun, Fengxian Bai, Special type electroc machine and control China waterpub press, 2005.137-143.

### RR26

#### **The measurement procedure of the equivalent core loss in a PM motor** Guo Jhih Yan<sup>1\*</sup>, Ming-hung Jian<sup>1</sup> and Chia-sheng Huang<sup>2</sup>

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For the development of high efficiency motor, the core loss decrease is one of the most important tasks. However, the calculated core loss in the design stage is not so accurate because the specification provided by material manufacturer usually obtained by ideal test condition in simple geometry, i.e. single sheet tester and toroidal tester. Therefore, a method that can measure the core loss with the practical shape of motor stator/rotor core includes every aspect of the manufacturing effects is required. This paper proposed a refined procedure based on torque metric method. Among several proven measuring methods, it is a directly measuring method that can be applied to the actual motor. Typical torque method uses a motor driving a specimen at a specific speed and then measures the drag torque. However, it is time consuming and requires higher precision torque meter. Instead, the proposed method measures the speed-power characteristic curve of the driving motor with without the test motor. From the difference of the speed-power characteristic, the core loss at various speeds can be obtained. Several conventional PM motors are taking as examples to demonstrate this procedure and to investigate the core loss influenced by manufacturing.

[1]N. Stranges, and R. D. Findlay, "Measurement of Rotational Iron Losses in Electrical Sheet," IEEE TRANSACTIONS ON MAGNETICS, Vol. 36, No. 5, Sep. 2000, pp. 3457-3459, [2]W. Arshad, etc., "Incorporating lamination processing and component manufacturing in electrical machine design tools," in Industry Applications Conference, 2007. 42nd IAS Annual Meeting, Conference Records, 2007, pp. 94-102, [3] Nouguang Guo, etc., "Measurement and Modeling of Rotational Core Losses of Soft Magnetic Materials Used in Electrical Machines: A Review," IEEE TRANSACTIONS ON MAGNETICS, Vol. 44, No. 2, Feb, 2008, pp. 279-291.

### **RR27**

#### MFL signal enhancement based on exponential smoothing

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Nowadays, as industrialization is in progress worldwide and economy grows rapidly, demand for the energy supply of oil and gas has also increased. Currently, over millions of kilometers of underground gas pipeline has been established worldwide. The maintenance of pipeline is needed for the efficient supply of energy, but in reality, it is impossible for human experts to investigate all the underground pipelines in order to confirm whether a pipe has deterioration or damage because of the pipe's geographical position. So, automatic nondestructive testing methods by using MFL (magnetic flux leakage) signal are currently studied actively to detect and extract defects on the pipeline. The raw signals acquired from MFL sensors contain the signals of defects, as well as noise and distortion caused by various underground environments. In this paper, MFL signal is enhanced by using several smoothing techniques. Through experiments, it is confirmed that the distortion of raw signal by noise canceling can be reduced.

# **RR28**

# Electromagnetic separation of the brown coal ash of thermal power stations

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Tens millions of brown coal ash have accumulated near thermal power stations (TPS) in Uzbekistan, so their utilization has become quite actual. So far the ash is used for cement. However the ash is well known to contain a large amount of valuable elements, including ~3.13%Fe<sub>2</sub>O<sub>3</sub> in Russian coal. Element analyses of the brown coal ash from Angren TPS were done with the neutron activation techniques and 3.6%Fe, 0.084%Mn, 0.063%Zr and also 0.023% of rare earth elements have been determined. The ash samples were undergoad to electromagnetic separation, and the magnetic fraction was 3.7%, where the content of Fe increased to 58%. Then the magnetic fraction was separated with the standard screen grader into 7 classes from 0.4 to 0.05mm. Two times increase of the contents of Fe and rare earth elements have been found in the fine fractions has revealed the presence of Fe<sup>2+</sup> and Fe<sup>2+</sup> ions. Microscopic analysis of the fine fractions has shown that there are black iron microspheres of 0.02-0.05mm sizes, which can be used for cleaning of technological liquid waste by means of the high gradient magnetic field, and powder metallurgy.

### **RR29**

#### Studies on viscidity in ferrofluids of Fe<sub>3</sub>O<sub>4</sub>

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Fe<sub>3</sub>O<sub>4</sub> nanoparticles with superparamagnetic properties are synthesized by sol-gel approach. Fe<sub>3</sub>O<sub>4</sub> nanoparticles are coated using the amorphous silica. The viscidity of the coated Fe<sub>3</sub>O<sub>4</sub> nanoparticles increases as the particles is incorporated in the impregnant, which may due to the magnetic interactions. At room temperature, the ferrofluid viscidity increases momotonically with increasing the concentration of Fe<sub>5</sub>O<sub>4</sub> nanoparticles, the ratio of surface activity. With rising the temperature, the ferrofluid viscidity becomes larger. The applied magnetic field has also effects on the ferrofluid viscidity. The switching process of the applied magnetic field may increase the ferrofluid viscidity, although the ferrofluid viscidity increases with increasing the applied magnetic field. Dynamic properties of the magnetic particles in the impregnant are also discussed in the present.

### **RR30**

Non-contact magnetic evaluation of ferromagnetic plate and its compensation of unknown air gap

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This paper represents a magnetic hysteresis model-assisted measurement method for non-contact evaluation of ferromagnetic plates in the case of the air gap between the test sample and the sensor probe is unknown. Our method involves measuring hysteresis loops from a test sample using a sensor probe in close proximity to the sample and measuring inductance spectra according to the various air gap. The sensor outputs were simulated based on a magnetic hysteresis model (Jiles-Atherton model) to describe the magnetic property of the sample. In this study, the measurement was applied to characterize a series of Fe-C samples. However, it is hard to confirm the permeability of a sample when the air gap is unknown, since the magnetic resistance of magnetic circuit composed of the sample and the sensor probe is affected by air gap size. The main idea of our study is that the frequency spectrum of inductance of the magnetic circuit is unaffected when the air gap is adequately small, even if the amplitude of inductance is directly dependent on air gap. The magnetic hysteresis measured with unknown air gap was compensated by measured inductance spectra, and was found to agree with nonlinear FEM analysis.

# RR31

# The several analysis techniques of high speed induction motor for copper die casting

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# SA02

# Magnetoelectric $CoFe_2O_4$ -PZT thin film composites grown by pulsed laser deposition

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### **RR32**

Effect of magnetic reynolds number variation on MHD convection inside an enclosure

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Numerically investigation is carried out for magnetohydrodynamics natural convection in an enclosure. The vertical walls are maintained isothermal at different temperatures and the other walls are adiabatic. The magnetic field is applied horizontally. The non linear governing equations for the fluid flow and heat transfer are solved for three magnetic Reynolds numbers (10-5, 10-3 and 0.1). The Rayleigh number, the Hartmann number and the Prandtl number are considered 105, 80 and 0.01, respectively. A finite volume code based on Patankar's SIMPLER method is utilized. It is found that when the magnetic Reynolds number (Rem) is very low (10-5 and 10-3) the magnetic field distribution in the cavity is not affected by flow field and as Rem increases the magnetic field in x and y direction is not constant. Also it is observed that as Rem increases the rate of heat transfer changes.

#### SA01

# Magnetic annealing effects on properties of the multilayer $BaTiO_3 / CoFeO_4$ thin films

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Multilayer BaTiO<sub>3</sub>/CoFe<sub>2</sub>O<sub>4</sub> thin films are prepared via chemical solution deposition method with and without magnetic field annealing process, respectively. The microstructural, magnetic, dielectric and the ferroelectric properties of the films are investigated. It is observed that the dielectric constant, polarizations and magnetizations of the multilayer thin films are improved with magnetic field annealing process. Moreover, for the magnetic field annealing films, the saturation magnetization with measuring magnetic field vertical to the film surface is much higher than that parallel to the film surface.

# SA03

# Synthesis and magnetic property of multiferroic DyMnO<sub>3</sub> nanoparticles in mesoporous silica

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Multiferroic compound DyMnO<sub>3</sub> shows antiferromagnetic ordering with  $T_N \sim 40$  K and Dy spin ordering below about 6 K. The ferroelectricity is attributed to lattice modulation

Dy spin ordering below about 6 K. The ferroelectricity is attributed to lattice modulation accompanied by the antiferromagnetic order. We are interested in the size effect on DyMnO<sub>3</sub> nanoparticles with particle size of a few nano-meters. We synthesized DyMnO<sub>3</sub> nanoparticles and investigated their magnetic property. The DyMnO<sub>3</sub> nanoparticles were synthesized in the pores of a mesoporous silica SBA-15 with pore diameter of about 8 nm. SBA-15 was used as a template to equalize the particle size in the fabrication of the DyMnO<sub>3</sub> nanoparticles. The X-ray diffraction pattern for DyMnO<sub>3</sub> nanoparticles shows some broad Bragg peaks corresponding to crystal structure of DyMnO<sub>3</sub> nanoparticles. The size of nanoparticles was estimated to be about 10 nm based on the Bragg peaks and using Scherrer's equation. The temperature dependence of DC magnetic susceptibilities for the nanoparticles exhibited a magnetic irreversibility between field-cooled susceptibility and zero-field-cooled one below about 40 K and an increase in susceptibilities due to magnetic ordering among the Dy spins below about 5 K. The magnetization curves exhibited hysteresis loop below blocking temperature. DyMnO<sub>3</sub> nanoparticles were successfully synthesized with size of about 10 nm and showed superparamagnetic behaviors.

#### SA04

# Magnetic properties of $Co_{1-x}Mn_xFe_2O_4$ (x=0-0.5) - PZT thin films fabricated by sol-gel spin coating method

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 $Co_{(1:s)}Mn_xFe_2O_4$  (x=0-0.5) thin films were fabricated on PZT layer by sol-gel spin coating method in order to investigation of influence of Mn-substitution in  $C_sFe_2O4$  on the magnetic properties. The films were fabricated on the PZT layer which was previously fabricated on Pt(111)/Ti/Slo\_/Sl substrate using the same method. The phase identification of the samples was studied by X-ray diffractometry analysis. Magnetic properties of the samples were also measured using a vibrating sample magnetometer.

### SA05

# M doping element localization by the molecular field theory in the Ga,Fe<sub>2</sub>,O<sub>3</sub>:M thin films

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Magnetoelectric materials open the way towards new applications through controlling magnetic or electric polarization with electric or magnetic field, respectively.  $Ga_{0.6}Fe_{1.4}O_3$  is predicted magnetoelectric with non zero magnetization at room temperature. High quality GFO thin films having similar magnetic properties to those of the bulk materials have been produced only recently. However their electric properties failed to be properly characterized until now because of important leakage currents originating from a hopping mechanism between  $Fe^{2+}$  and  $Fe^{3+}$  within the structure. Recently, we succeeded to strongly lower the leakage currents of these films by chemical doping. The determination of the cationic distribution within the structure is a key point for the knowledge of these phases, because it is at the origin of the evidenced magnetic and electrical properties of these materials. We present here an original approach to determine the localization of the M doping element based on the molecular field theory. Different modeling associated with the different substitution possibilities have been performed within the frame of this theory and are compared with the experimental values.

#### SA06

# Interplay between magnetization and polarization in epitaxial (Ga,Fe)<sub>2</sub>O<sub>3</sub> thin films with additional ion-substitutions into Fe sites

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(Ga,Fe)<sub>2</sub>O<sub>3</sub> (GFO) is a promising material for room-temperature multiferroic because of high magnetic  $T_c$  as 370 K when Ga = 0.6 and Fe = 1.4. However, multiferroic properties of GFO in thin film type is controversial because most of polarization-electric field hysteresis curves in GFO showed leaky behaviors due to large leakage current. In this research, we focused on the improved magnetic and electric properties of GFO thin films as multiferroic by decreasing their leakage current. The main origin probably lies on electron hopping between Fe<sup>3+</sup> and Fe<sup>3+</sup> octahedron sites which can co-exist due to charge neutrality by oxygen vacancy. Therefore, we tried to decrease leakage current by substituting Fe<sup>3+</sup> site with divalent cation. The ion-substituted GFO thin films were epitaxially deposited along the b-axis by pulsed laser deposition at 750°C for 15 min in oxygen partial pressure of 200 mTor on SrRuO<sub>2</sub>/SrTiO<sub>3</sub> (111) substrates. The films showed the strongly reduced leakage current up to 5–6 order of magnitude. Magnetic properties of the doped GFO thin films were investigated by SQUID measurement. Ferroelectric polarization as a function of external field was measured in macroscopic and microscopic schemes. Magnetoelectric interactions of the epitaxial thin-films will be addressed in this talk.

#### SA07

# Fabrication and characterization of heusler-alloy/perovskite heterostructures

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Multiferroic heterostructures composed of ferromagnetic half-metal and ferroelectric materials have attracted significant interest due to the potential applications for nanoscale devices. Heusler-alloy / perovskite heterostructures are expected to emerge the interface magnetoelectric1) and turnel electrore-sistance effect. In this study, to explore the potential of Heusler-alloy / perovskite heterostructures are expected to emerge the interface magnetoelectric1) and turnel electrore-sistance effect. In this study, to explore the potential of Heusler-alloy / perovskite heterostructures, we have investigated multiferroic and transport properties of heterostructures composed of half-metallic Heusler-alloy Fe\_CrSi (FCS) and ferroelectric perovskite  $Ba_0,Sr_{0.3}TiO_3$  (BSTO). BSTO was firstly deposited on MgO and LaAIO, substrates by magnetron sputtering, followed by the deposition of FCS at 550-750°C. Structural characterization revealed that FCS grew epitaxially on BSTO (100 JIIO)][FCS(001)[110]. From M-H curve of the BSTO (100 nm) / FCS (150 nm) bilayer, saturated magnetization of 417 emu/cc was obtained. The ferroelectric properties of La $_0,Sr_0,MnO_3$  (LSMO) / BSTO (100 nm) / FCS (50 nm) were measured, and the maximum remaneant polarization was 73  $\mu$ Ccm<sup>2</sup>. From these results, ferromagnetic and ferroelectric properties were preserved in BSTO/FCS theterostructures. From the analysis of 1-V characteristics, Schottky-emission was found to be a dominant leakage mechanism in LSMO / BSTO / FCS. The BSTO / FCS

1) K. Yamauchi et al., Appl. Phys. Lett. 91, 062506 (2007)

# **SA08**

Optical properties of  $(SrMnO_3)_n/(LaMnO_3)_{2n}$  superlattices: An insulator-to-metal transition observed in the absence of disorder

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We measure the optical conductivity,  $\sigma_1(\omega)$ , of  $(SrMnO_3)_n/(LaMnO_3)_{2n}$  superlattices (SL) for n=1, 3, 5, and 8 and 10 < T < 400 K. Data show a T-dependent insulator to metal transition (IMT) for n=3, driven by the softening of a polaronic mid-infrared band. At n=5 that softening is incomplete, while at the largest-period n=8 compound the MIR band is independent of T and the SL remains insulating. One can thus first observe the IMT in a Manganite system in the absence of the disorder due to chemical doping. Unsuccessful reconstruction of the SL optical properties from those of the original bulk materials suggests that  $(SrMnO_3)_n/(LaMnO_3)_{2n}$  heterostructures give rise to a novel electronic state.

A. Perucchi et al., Nano Letters 10 (12), 4819-4823 (2010)

#### **SA09**

#### Preperation of hexagonal YFeO<sub>3</sub> powder and thin film

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In the past few years, multiferroic materials, characterized by coexistence of two or more ferroicity, have attracted much attention due to their interesting properties. Among the magnetic ferroelectric materials, there is an important class, hexagonal manganese oxides (RMnO<sub>3</sub>, R: rare-earth element), which have been extensively studied because of its unusual property. However, hexagonal iron oxides (RFeO<sub>3</sub>), which have a similar structure to RMnO<sub>3</sub>, have been rarely reported due to a preferred formation to the orthombic or garnet phase. Here, we report that the hexagonal phase can be easily prepared as a powder and thin film from chemical solution deposition method (CSD). The crystallization temperature, crystal structure, magnetic and electric properties of the thin film and powder were investigated. The hexagonal phase was observed at 750-800, and at higher temperature (>850) orthombic phase, rather than ferroelectric phase. Samples showed paramagnetic property at room temperature caused by geometrically fluctuated magnetism in tri-angular spin lattice.

#### SA10

#### Magnetic hysteretic effects in LaAlO<sub>3</sub> /SrTiO<sub>3</sub> Heterostructures

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A two-dimensional electron system (2DES) with remarkable electronic properties can be realized at the interface between two band insulators LaAlO<sub>3</sub> and SrTiO<sub>3</sub> (1). Depending on growth conditions, this 2DES has been found to show a superconducting ground state (2), magnetic ordering (3) and even co-existence of these two properties (4-5). The origin of the magnetic behaviour at the interface is still unclear, and to try to understand it we have performed magneto-transport studies on LaAlO<sub>3</sub>/SrTiO<sub>3</sub> heterostructures in high magnetic fields up to 33 T, and low temperatures down to 0.3 K. We observed negative magnetoresistance (nMR) which is nearly identical for all field orientations in the temperature range from 4 K to 350 mK. This nMR could be related to reduced spin scattering due to the alignment of spins at the interface (3). At low magnetic fields, the nMR is very sensitive to temperature and at temperatures below 1 K an unusual time dependent magnetic hysteresis appears for field below 3 T. We propose a magnetocaloric effect model to explain the hysteretic behaviour, and propose that the possible underlying origin of the hysteresis is the co-existence of ferro- and antiferromagnetic belows in system.

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### **SA11**

# Synthesis and magnetoelectric properties of multiferroic composites on the cobalt ferrite - pzt system

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In the present study we prepared and investigated the properties of bulk particulate composites made from cobalt ferrite (COF) and a soft piezoelectric material (PZT). We chose COF and soft PZT due to their high magnetic and piezoelectric parameters respectively. Composite samples, according to the formula xCOF(1-x)PZT with  $0 \le x \le 1$  were prepared by mixing the COF and PZT powders, pressing them as discs and then sintered at 1200 C for 3 hours. Structural, magnetic, piezoelectric and magnetoelectric characteristics of the sintered composites were then determined as a function of the weight fraction of COF into PZT and magnetic fields. The ME coefficient shows a maximum value about 70 mV/cmOe for composites containing 40 % COF and 60% PZT, in a DC magnetic field of 3 kOe at a frequency of 1KHz of the ac magnetic field. Some of the dielectric and piezoelectric properties of the composites are aligned to the remove of the recent theoretical cube model of different connectivities of particles of one phase embedded into the matrix of the second phase.

### SA12

# Detailed structural study of BiFeO<sub>3</sub>/SrRuO<sub>3</sub> heteroestructures grown on SrTiO<sub>3</sub>(001) substrates

Claribel Dominguez<sup>1</sup>, John Edward Ordonez<sup>1</sup>, Maria Elena Gomez<sup>1</sup>, Wilson Lopera<sup>1</sup> and Pedro Prieto<sup>2</sup> <sup>1</sup> Thin Film Group, Department of Physics, Universidad del Valle, A.A. 25360, Cali, Colombia <sup>2</sup> Center of Excellence for Novel Material CENM, Calle 13 # 100-00 320-1026, Cali, Colombia

# WITHDRAWN

#### SA13

### Lattice engineering on transition metal oxide thin film Chang Uk Jung\*

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### SA14

#### Magnetisation in different multiferroic YMO (Yttrium Manganese Oxide) thin films

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Recently Yttrium manganese oxides (YMO): hexagonal YMnO<sub>3</sub> and orthorhombic YMn<sub>2</sub>O<sub>5</sub> have attracted considerable attention as both of these compounds exhibit multiferrocity but the microscopic origins are different. Here we stabilize thin films of different phases of YMO by pulsed laser deposition from a single h-YMnO<sub>3</sub> target and characterize them with various techniques and perform magnetization measurements. It is observed that the phase stability and crystallinity of YMO thin films depend on the substrate used and oxygen partial pressure (OPP). Oriented and polycrystalline growth of h-YMnO<sub>3</sub> and O-YMn<sub>2</sub>O<sub>5</sub> phase film are observed on different substrates. ZFC,FC (Magnetisation verses Temperature ) and MH magnetization measurements of these multiferroic films were performed. All the multiferroic thin films show the antiferromagnetic transition in ZFC measurements which is forbidden in FC measurements. MH behavior at both sides of magnetic transitition is in good agreement in these multiferroic thin films.

 Bas B. Van Aken, Thomas T. M. Palstra, Alessio Filippetti and Nicola A. Spaldin,Nature Materials, Vol 3, 2004, pp.164-170 2 L. C. Chapon, P. G. Radaelli, G. R. Blake, S. Park and S.W. Cheong, Phys.Rev.Letter 96, 097601 (2006) 3. W. Prellier, M P Singh and P Murugavel, J. Phys.: Condens. Matter 17 (2005) R803-R832 4. A Munoz et al 2002 J. Phys.: Condens. Matter 14 3285

#### SA15

#### Dependence of magnetoelectric response on magnetostrictive content in composite multiferroics

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Magnetoelectric multiferroics exhibit a strong coupling between ferromagnetic and ferroelectric ordering. The coupling between these two types of orderings may occur naturally or be mediated via mechanically. This coupling is very weak in single phase multiferroics. The design and investigation of composites of magnetostrictive and piezoelectric phases have shown magnetoelectric coupling orders of magnitude higher than that of single phase multiferroics. The mechanical deformation in magnetostrictive and piezoelectric phases results in the coupling between the two constituents. The multiferroic CoFe<sub>2</sub>O<sub>4</sub>-BaTiO<sub>3</sub> (CFO-BTO) composites with nanoscale mixture of the two constituents were prepared by solgel process. Four different compositions, with varying percentage of constituents, were investigated. X-ray diffraction revealed the coexistence of the two phases. Magnetic properties of the CFO-BTO nanccomposites were examined which were consistent with the composition. The magnetoelectric coupling between CFO and BTO was demonstrated by an external magnetic field induced change in capacitance. The magnetoelectric coupling was found to be not only dependent on the applied magnetic field but also on the composition of the composites. It has a nonlinear dependence on magnetostrictive content.

#### SA16

### Raman analyses of oxygen defects in hexagonal HoMnO<sub>3</sub> thin films

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Oxygen defects are usually unavoidable in synthesizing oxide thin films. In this study, we investigated the origins of the oxygen defects in hexagonal manganites HoMnO<sub>3</sub> epitaxial thin films through Raman spectroscopy. We found that the oxygen defects in hexagonal HoMnO<sub>3</sub> thin films have different effects on different phonon modes and on magnon scattering. Our analyses indicated that the oxygen defects in hexagonal HoMnO<sub>3</sub> thin films would be mainly correlated with the basal O<sub>3</sub> and/or O<sub>4</sub> oxygen ions. Furthermore, our analyses of oxygen defects predicted that the Mn 3d orbitals would be more strongly hybridized with the apical O<sub>1</sub> and/or O<sub>2</sub> 2p orbitals than the basal O<sub>3</sub> and/or O<sub>4</sub> 2p orbitals. This prediction is consistent with our resonant Raman scattering study and earlier first-principle calculations of electronic structures of hexagonal manganites.

# SA17

Preparation and properties of inverse perovskite Mn<sub>3</sub>GaN thin films Hiroki Tashiro, Ryosuke Suzuki\*, Tetsuya Miyawaki, Kenji Ueda and Hidefumi Asano

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There has been considerable interest in antiperovskite Mn<sub>3</sub>GaN because of its intriguing physical properties such as negative thermal expansion, magneto-volume effect and gnetic effect(1). Mn<sub>3</sub>GaN is expected to be advantageous to fabricate thin-film hybrid devices with various perovskite oxides, since it has a cubic antiperovskite structure virtually identical with the perovskite oxides. However, there is no report on growth of Mn<sub>3</sub>GaN thin films and heterostructures. In this paper, we report on the preparation and properties of Mn<sub>3</sub>GaN thin films as well as ferroelectric Ba<sub>0.7</sub>Sr<sub>0.3</sub>TiO<sub>3</sub>/Mn<sub>3</sub>GaN heterostructutes. Mn3GaN films were deposited on (001) SrTiO3 and LaAlO3 substrates at 450~650°C by ion beam sputtering of a sintered stoichiometric target. Ba07Sr03TiO3films were prepared at 650~750°C by magnetron sputtering. In-plane and out-of-plane X-ray diffraction analysis revealed that on SrTiO3 substrate Mn3GaN films were grown coherently with the in-plane lattice parameter aMGN of 0.3903 nm. On LaAlO3 substrate a Ba07Sr03TiO3 layer and subsequently a Mn3GaN layer were grown epitaxially with the relationship of LaAlO<sub>3</sub>(001)[100]//Ba<sub>0.7</sub>Sr<sub>0.3</sub>TiO<sub>3</sub>(001)[100]//Mn<sub>3</sub>GaN (001)[100]. These Mn<sub>3</sub>GaN thin films showed a metallic temperature dependence of resistivity with 0.8-0.9 m $\Omega$ cm at room temperature. It was found that the Ba<sub>0.7</sub>Sr<sub>0.3</sub>TiO<sub>3</sub>/Mn<sub>3</sub>GaN heterostructutes exhibited a colossal magnetocapacitance effect.

(4) P. Lukashev, et al., Phys. Rev. B 78 184414 (2008)

#### SA18

#### Ferroelectric and magnetic properties of BiMnO<sub>3</sub> thin films Min-hwa Jung and Yoon-hee Jeong\*

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BiMnO<sub>3</sub> is heralded widely as a rare case of multiferroicity with simultaneous existence of both ferroelectricity and ferromagnetism. Ferroelectricity in BiMnO<sub>3</sub>, however, still remains controversial as various results do not converge on a definite conclusion [1-3]. In order to shed light on this issue, we have synthesized BiMnO<sub>3</sub> thin films on Nb:SrTiO<sub>3</sub> substrates by PLD method. Measurements of various physical properties such as dielectric properties, PE hysteresis loop, and piezoelectric d33 do not show any significant evidence of ferroelectricity for all films. Thus, we are led to conclude that BiMnO<sub>3</sub> is not ferroelectric in accordance with a theoretical calculation [4-5]. Also of interest is the ferromagnetism of BiMnO<sub>3</sub> and, in particular, its magnetic easy direction. We also investigate the magnetic properties of the films with in and out of the plane orientations. From these measurements, the magnetic easy direction is inferred for the monoclinic phase BiMnO<sub>3</sub>. In conclusion, the present study offers a full characterization of the ferroelectric and magnetic properties of multiferroic BiMnO<sub>3</sub>.

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#### SA19

# Ferroelectric polarization induced magnetic anisotropy in $Co_{40}Fe_{40}B_{20}/$ $YMnO_3$ multiferroic heterostructure

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Multiferroic (MF) materials have attracted much attention due to their interesting new physics and potential for exploring novel multifunctional devices. The research of magnetic anisotropy in MF heterostucture is an important aspect. One of the key issues in the MF heterostructures is the effect of ferroelectric (FE) polarization on magnetic anisotropy. So far, however, a lot of work has been carried out on MF heterostructures with crystalline ferromagnetic films. Actually, MF heterostructures involving amorphous ferromagnetic materials are also very interesting due to the absence of magnetocrystalline anisotropy and low correive force Hc. We have studied the coupling between amorphous magnetic Coa<sub>4</sub>Fea<sub>40</sub>Ba<sub>3</sub> (CFB) and YMnO<sub>3</sub> (YMO) in the heterostructure with CFB film grown on YMO (110) substrates. A large uniaxial magnetic anisotropy (UMA) was observed from room temperature to 10 K with magnetic easy axis (MEA) along the polarization direction of YMO. More interestingly, the angular dependence of coercive force Hc shows a non-monotonous behavior. The results were discussed by considering the coupling between the FE domain in YMO and magnetism in CFB. This work is helpful for understanding the effect of FE polarization on magnetic anisotropy in MF heterostuctures.

# SA20

Ferroelectric-domain-switching controlled magnetism in CoFeB/PMN-PT multiferroic heterostructure

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Multiferroic materials have attracted much attention due to their interesting new physics and potential applications in new multifunctional devices. In multiferroic materials, the magnetization can be tuned by electric fields, which open a new way to reach high integration density and low power consumption in information storage. In the shorter term, niche applications are more likely to emerge in the two-phase systems [1], so there have been lots of work on the electric-field control of magnetism via the pizzo-strain in ferromagnet-ferroelectric heterostructures[2, 3]. However, the reports related to magnetoelectric (ME) effect via ferroelectric domain switching are still limited. Thus, we have studied the electric-field control of magnetism in Co<sub>20</sub>Fe<sub>40</sub>B<sub>20</sub>/Pb(Mg<sub>12</sub>Nb<sub>22</sub>)<sub>0.7</sub>Ti<sub>03</sub>O<sub>3</sub>(PMN-PT)(001) two-phase-system. Nonvolatile and volatile changes of magnetization related to different types of ferroelectric field control of magnetism indicates that some magnetic domains rotate 90°, which is consistent with our in situ electric-field dopendent magnetic force microscopy data. The results were explained by considering the ferroelectric domain switching induced magnetic anisotropy. This work is helpful for understandine the coupling between magnetism and ferroelectric anisotropy. This work is helpful for

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#### SA21

#### Multiferroic properties of $BaTiO_3 - Zn_{1-x}Co_xO$ multilayer thin films

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Near room temperature ferroelectric and ferromagnetic hysteresis is observed in horizontal superlattice system of  $BaTiO_3 - Zn_{1,x}Co_xO$  (x > 0.2) [BTO - ZCO] multilayered thin films. The multilayered thin films were prepared on Si substrate by alternate deposition of BTO and ZCO layers to a thickness of 350 nm using pulsed laser deposition technique. Analysis of the prepared thin film by x-ray powder diffraction and electron microscopic techniques confirm the polycrystalline nature of the multilayered film. Dielectric and the magnetic properties studied on the multilayered thin film at different temperatures exemplified the coupling between their dual order parameters and it is mediated through an electric field effect at the interface.

 Multiferroic properties of Bi0.8La0.2FeO3/CoFe2O4 multilayer thin films. Shan-Tao Zhang, Yi Zhang, Zhen-Lin Luo, Ming-Hui Lu, Zheng-Bin Gu, Yan-Feng Chen., Applied Surface Science 255 (2009) 5092-5095. 2. Multiferroic BaTiO3-CoFe2O4 Nanostructures., H. Zheng, J. Wang, S. E. Lofland, Z. Ma, L. Mohaddes-Ardabili, T. Zhao, L. Salamanca-Riba, S. R. Shinde, S. B. Ogale, F. Bai, D. Yiehland, Y. Jia, D. G. Schlom, M. Wuttig, A. Roytburd, R. Ramesh. Science 303, (2004) 661.

### SA22

#### Study of structural phase transition and multiferroic properties of Samarium substituted BiFeO, thin films

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In this study we report the effect of samarium substitution on the structure and multiferroic properties of BiFeO<sub>3</sub> thin films. The Bi<sub>1-x</sub>Sm<sub>x</sub>FeO<sub>3</sub> thin films of various thicknesses were sputter deposited on to Si substrates. Magnetic and electric properties of the thin films were studied by vibrating sample magnetometer, impedance analyzer and I-V characterizer. Samarium substitution-driven structural phase transition was studied by x-ray diffraction at room temperature. Samarium doping was found to modify structural properties of Bi<sub>1-x</sub>Sm<sub>x</sub>FeO<sub>3</sub> thin film. The phase transition was found to destruct cycloid spin structure of BiFeO<sub>3</sub> that resulted in amplified magnetoelectric interactions. The substitution of Sm was found to be a promising way to control the conductivity and leakage current of film by suppressing the formation of oxygen vacancies.

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# SA23

Study of ferroelectricity in Eu substituted BiMnO<sub>3</sub> films Shivaraman Ramaswamy\*, Helen Annal Therese and C Gopalakrishnan

Nanotechnology Research Center, SRM University, India In this paper, we report evidence of ferroelectricity in europium substituted manganite (Bi<sub>1-x</sub>Eu<sub>x</sub>MnO<sub>3</sub>) thin films synthesized via sol gel synthesis and dip coating method on a single-crystal (1 0 0) substrates. The films were post treated by annealing at 500°C. X-ray diffraction measurements were used to analyze the crystal structure and phase purity of the thin films. AFM measurements were performed to investigate surface morphology; quantitative values of roughness and grain size which are in the range between 30 and 150 nm. Ferroelectric characterization was conducted at low

and phase purity of the thin lims. AFM measurements were performed to investigate surface morphology; quantitative values of roughness and grain size which are in the range between 30 and 150 nm. Ferroelectric characterization was conducted at low temperatures and at 300 K. Hysteresis loops (polarization vs. voltage) were obtained, showing enhanced ferroelectricity at room temperature in the 10% Eu substituted BiMnO<sub>3</sub> film. Resistance vs. temperature measurements were performed, which indicated this to be very robust insulating material. Moreover, all films show small dielectric dispersion and dielectric loss.

# **SB02**

**SB03** 

Charge disproportion, spin and orbital states in the tri-layered nickelate  $La_4Ni_3O_8$  from first principles Hua  $Wu^*$ 

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Carrier doping to the novel layered nickelate

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Several mixed-valent layered nickelates were synthesized very recently [1,2], in the effort to search an analog to the hole doped cuprates. Here we have studied the electronic structure and the magnetism of the tri-layered La<sub>4</sub>Ni<sub>5</sub>O<sub>8</sub>, using the configuration-state constrained LDA+U calculations. Our results show that La<sub>4</sub>Ni<sub>5</sub>O<sub>8</sub> would be a charge homogeneous low-spin half metal in the ferromagnetic state. In contrast, the C-type antiferromagnetic state is insulating and has a charge disproportion of the formal 2Ni<sup>+</sup>/Ni<sup>2+</sup> type, thus accounting for several experimental observations. The high-spin (S=1) Ni<sup>2+</sup> state and its orbital configuration are found to be against the crystal-field level picture. We note that it would be interesting to explore, for novel magnetism and superconductivity, the Ni<sup>2+</sup> low-spin S=0 state and the partially occupied x<sub>2</sub>-y<sub>2</sub> band structure of the Ni<sup>+</sup>/Ni<sup>2+</sup> mixed-valent nickelate.

[1] V. V. Poltavets et al., Phys. Rev. Lett. 104, 206403 (2010). [2] V. V. Poltavets, M. Greenblatt, G. H. Fecher, and C. Felser, Phys. Rev. Lett. 102, 046405 (2009).

Recently a series of new layered nickelate (Lnn+1NinO2n+2) (Ln=La, Nd) has been

synthesized with using CaH2 as a reducing agent or hydrogen gas. The crystal structure

is the analogue of so-called Ruddlesden-Popper phase, and basically the same crystal

structure with high-Tc superconducting cuprate, consisting of insulating block layer

and two-dimensional NiO2 layer. Interestingly Ni ions in these materials can have the

same electronic configuration with Cu<sup>2+</sup> 3d<sub>9</sub> in high-Tc cuprate, suggesting that this is a

promising candidate for another high-Tc material. Therefore, we have tried to carrier-

dope to Ln and Ni sites of (Ln<sub>n+1</sub>Ni<sub>n</sub>O<sub>2n+2</sub>) (Ln=La, Nd) with n=2, 3, and ∞ using

Ca. Cu. Co ions in order to tune the electronic state of Ni. If the Ni electronic state

is adjusted to 3d<sub>9.15</sub> or 3d<sub>8.85</sub> which are the same cases with those of optimum doping

level for the emergence of high-Tc in cuprate. We will present the detailed result in the

### SA24

# Experimental evidence for Exchange bias in polycristalline $BiFeO_3 / Ni_{01}Fe_{10}$

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Multiferroic (MF) materials have been the subject of extensive studies in the past decade since it allows a coupling of ferroelectric and magnetic orders in a single phase. Among all the multiferroic materials studied so far, BiFeO<sub>3</sub> (BFO) exhibits the coexistence of ferroelectric and antiferromagnetic orders at room temperature. Magnetic heterostructures with a ferromagnetic thin film exchange biased to an antiferromagnetic-BFO film may allow an electric field influence on the magnetic hysteresis properties of the bilayer. The exchange coupling and exchange bias involving MF/ systems has previously been obtained with either epitaxial BFO [1] or BFO bulk materials [2]. However, the exchange coupling between a ferromagnetic film and a polycrystalline multiferroic was not studied. Furthermore, previous studies have shown that polycrystalline BFO exhibits multiferroic properties. We report experimental evidence for exchange bias in polycrystalline BiFeO<sub>3</sub>/Ni<sub>81</sub>Fe<sub>19</sub> bilayers, grown by RF sputtering [3]. Our measured 17 Oe shift corresponds to exchange energy per unit area of interface coupling the two films of  $11 \times 10^{-3}$  eng/cm<sup>2</sup> comparable with similar epitaxial grown systems. The azimuthal behavior of longitudinal and transverse components of magnetization reveals the presence of induced undirectional and biquadratic anisotropies.

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#### **SB01**

# Parit variation in the rectangular array of periodic holes on superconducting thin film

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Superconductivity has been full of surprises and wonder. One of the most elegant properties of superconductors is the quantization of magnetic flux in units of  $\phi 0$ =h/2e, where h is the Planck constant of quantum mechanics, and 2e the electron-pair charge serving the fundamental building block of the superconducting condensate. Here we report our discovery of a new twist of this phenomenon in a film geometry perforated with a periodic lattice of sub-micron holes, a variation of the flux quantization based on the parity of the number of the flux quanta. Measured in terms of magneto-resistance at temperatures just below the transition temperature of the bulk superconductor, we find that superconductivity is most enhanced at odd multiples of the flux quantum. We explain the phenomenon by noting that thermal quasi-particles, which see the superconducting flux quantum only as half, are in fact suppressed at these flux values. In fields larger than the thermodynamic critical field, where quasi-particles dominate, magneto-resistance minima only occur at such odd values.

#### SB04

conference

#### Occurrence of superconductivity and structural variations in the $Sr_{1,x}T_{2}$ . <sub>25</sub>Ge<sub>2</sub> layer system (T = Ni, Pd, and Pt; $x \ge 0$ )

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Superconducting transition temperature Tc of the body-centered-tetragonal (bct) 122 Sr(Ni<sub>1,y</sub>Pdy)<sub>2</sub>Ge<sub>2</sub> ( $0 \le y \le 1$ ) layer system (space group 14/mmm) increases monotonically from 0.87 K for 3d-Ni SrNi<sub>3</sub>Ge<sub>2</sub> to 3.12 K for 4d-Pd SrPd<sub>3</sub>Ge<sub>2</sub>. However, due to larger atomic radius of 5d-Pt, the bet stoichiometric Sr(Pd<sub>1,y</sub>Pty)<sub>2</sub>Ge<sub>2</sub> phase (14/mmm) is unstable and a structural transformation to a vacancy stabilized superstructure Sr<sub>1,x</sub>(Pd<sub>1,y</sub>Pty)<sub>2,z</sub>Ge<sub>2</sub> with orthorhombic space group Pmna may be formed. Trace of superconductivity with Tc ranging from 5-10 K in the Sr<sub>1,x</sub>Pt<sub>2,2x</sub>Ge<sub>2</sub> system was observed but very difficult to be stablized. This structural transformation is very similar to the high temperature superconducting system K<sub>1,x</sub>Fe<sub>2,2x</sub>Se<sub>2</sub> with 14/mmm and two superstructure space groups Pmna and 14/m, where superconductivity Tc up to 32 K was reported only in the bct 14/m system. This work was supported by Grant No. NSC98-2112-M-007-013-MY3 of National Science Council of Republic of China.

### **SB05**

Fermi surface study on the rattling-induced superconductor KOs<sub>2</sub>O<sub>6</sub> Taichi Terashima<sup>1</sup>\*, Nobuyuki Kurita<sup>1</sup>, Andhika Kiswandhi<sup>2</sup>, Eun-sang Choi<sup>2</sup>, James S. Brooks<sup>2</sup>, Kota Sato<sup>3</sup>, Jun-ichi Yamaura<sup>3</sup>, Zenji Hiroi<sup>3</sup>, Hisatomo Harima<sup>4</sup> and Shinya Uji<sup>1</sup>

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#### SB06

#### Transient analysis of the current density and temperature distribution of the MgB<sub>2</sub> superconductor in the He atmosphere H. M. Iftekhar Jaim<sup>1</sup> and Klaus Barner<sup>2</sup>

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In this paper, the transient analysis is carried out for the type II superconductor material MgB<sub>2</sub> in liquid Helium atmosphere. The strong nonlinearity of the E-J relationship near critical current poses a challenge in such calculations for getting practical values in other models employed for Ohmic materials. For this reason, the curl-curl equation and the variable magnetic field quantity are employed as the unknown in solving the PDE by the finite element method. The temperature distribution by the generated heat in the superconductor is calculated by the coupled conduction and convection module. Using the nominal experimental values for the exponential current fundamental characteristics like the magnetic flux, current distribution density and temperature profile in the superconductor are calculated for 2D axis symmetric cylindrical cases. The study gives the limiting current values to avoid heat generation and shows the development of the mixed state and finally the normal state with the time. Effect of changing the magnitude of current and time constant is also studied. Ripple effect at the surface is observed due to difference in the inner and outer magnetic flux distribution.

#### SB07

# Meissner-like effect on normal- superfluid interface of imbalanced Fermi gas

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The observation of phase separation between the superfluid (SF) core and a normal (N) shell around the core in ultra-cold Fermi gases has been reported by the new experiments of MIT and Rice groups[1-2]. Much researches have been done about the N-SF interface[3-5]. In this article, we examine the N-SF interface of a imbalanced Fermi gas with two spin species, in the presence of a weak magnetic field. In our analysis, we consider a constant magnetic field in the N side and anticipate the possibility of the Meissner-like effect in the SF side. We take into account the various scattering regions of the BCS regime and analytically obtain the transmission coefficients by using the perturbed Bogoliubov equations. It suffices to remark that the leading order term in transmission coefficients are independent of the energy of incident quasiparticles. Then, we calculate the heat conductivity across the interface at sufficiently low temperatures. We describe how the heat conductivity is affected by the Meissner effect and the species imbalance. Also the relation between the additional heat conductivity and penetration depth is obtained. The corresponding graphs is also plotted and discussed. Furthermore, we show the dependence of the heat conductivity on average chemical potential.

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### **SB08**

### Magnetic properties and structure evolution along R<sub>2</sub>RhIn<sub>8</sub> series

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Intermetallic compounds R<sub>2</sub>Rhln<sub>8</sub> (R = rare earth) are structurally related to a class of Ce-based heavy-fermion superconductors. Study of their magnetic structure and behavior is important for understanding mechanism of the unconventional heavy-fermion superconductivity. We have successfully grown single crystals of these compounds with R = Nd, Tb, Dy, Ho, Er and Tm from the melt in the In flux. We report the evolution of the magnetic properties of this series. All studied compounds exhibit antiferromagnetic behavior with the Neel temperatures up to 43 K for Tb compound. Magnetization and specific heat measurements have revealed metamagnetic transitions to another antiferromagnetic phase for Tb, Dy and Ho compounds. We show detailed magnetic phase diagrams for all these compounds and compare with results reported for Nd,Rhln<sub>8</sub> [1]. Low temperature magnetic structures were studied by neutron diffraction at the Institute Laue-Langevin on VIVALDI Laue diffractometer and four-circle D<sub>10</sub> diffractometer. From these experiments, propagation vector k = (1/2, 1/2, 1/2) and magnetic structure of studied compounds were determined. We present comparison of our results with the magnetic structure of Ce<sub>2</sub>Rhln<sub>8</sub> and other isostructural compounds from RRhln5 [2] and R<sub>2</sub>CoGa<sub>8</sub>[3] series.

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#### **SB09**

# Second magnetization peak and magnetic field distribution in superconductor

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The peak effect on M-H loop of superconductors is analyzed within the framework of the extended critical state model (ECSM) [D. M. Gokhfeld, D. A. Balaev, S. I. Popkov, K. A. Shaykhutdinov, M. I. Petrov, Physica C 434, 135 (2006)]. The second magnetization peak is referred to the vortex lattice transition or the normal phase rise occurred at a certain magnetic field value. According to ECSM, these processes extend from the surface of a sample (or a granule) to inside due to the magnetic field gradient. Any experimental M-H loops are fitted by assigning the appropriate distribution of the magnetic field in the sample and the field dependence of the critical current density. ECSM predicts that the asymmetry of M-H loop produces the mismatch of the peak positions in the field increasing and decreasing branches. The peak in the field increasing cycle is situated at lower H value than one in the field decreasing cycle.

#### **SB10**

#### Superconductivity in LuGe<sub>2</sub> single crystals

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Superconducting binary germanide LuGe<sub>2</sub> single crystals were synthesized using a high-temperature metal flux method. Several shiny plate-like single crystals were obtained and the compound has the orthorhombic ZrSi<sub>2</sub>-type structure (a = 3.9817 Å, b = 15.5573 Å, c = 3.8477 Å). Superconducting phase transition temperature at  $T_c = 2.35$  K was confirmed by temperature- and magnetic-field-dependent magnetization, and electrical resistivity measurement. In addition, we observed anomalous phase transition at around T = 4 K, and confirmed the conventional superconductivity by heat capacity measurement. Including these results, we will discuss the nature of LuGe<sub>2</sub> single crystals in detail.

#### **SB11**

#### Mutual interplay of magnon BEC and superconductivity Zvømunt Bak\*

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We present a theoretical study of the mutual interplay of magnon Bose-Einstein condensate and superconductivity /SC/, provided that they are simultaneously present. As a rule these phenomena are treated separately but it is evident that, at least in the case of triplet SC, there arises coherence in the electron as well as in the magnetic system. Therefore it is reasonable to assess whether magnon BEC can facilitate transition to the superconducting state. Such situation should arise when the electron pairing arises due to the magnetic interactions . In this case electron pair is accompanied by a cloud of magnons and below critical temperature there arises composed electronmagnon condensate. We point out that there are experimental evidences that magnetic coherence reduces or even switches off the destructive effect of magnetic excitations on SC. We prove that in the case of triplet superconductivity there should appear the quadrupolar coherence and what's more quadrupolar interactions contribute to the electron pairing. We give arguments that Josephson tunneling and "second sound" experiments in magnon and electron condenstes with various junction and setups can serve as a cross-tests verifying/falsifying hypothesis that in some systems SC is based on magnetic interactions.

### SB12

# Studies of the absolute value of lambda in unconventional superconductors

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The superconducting coherence length (\$\xi\$), magnetic penetration depth (\$\lambda\$), and their anisotropy are fundamental parameters that characterize superconducting materials. The value of \$\xi\$ can be estimated from the superconducting upper critical field (\$H {c2}\$) using the Ginzburg-Landau theory. In contrast to \$\xi\$, precise determination of the absolute \$\lambda\$ value is notoriously difficult. In this talk we present a simple, robust, and modelindependent technique to measure \$\lambda\$ using magnetic force microscopy (MFM), and discuss our recent results of the measurements of the absolute value of \$lambda\$ in a variety of unconventional superconductors, such as MgB2, borocarbides, and pnictides. In MgB<sub>2</sub> we will discuss the influence of the intra-band contributions on the superconducting properties and scattering effect between the two bands. In a recently discovered Ca10(Pt<sub>3</sub>As<sub>8</sub>) [(Fe1, Pt,),As2]s (x=0.097) single crystal, we observed large \$\lambda\$. Large \$\lambda\$ and high anisotropy in this system enhance thermal fluctuation, which leads to a liquid vortex phase in this low-temperature (T\$ c\$=11 K) superconductor. Also, we will discuss the effect of irradiation on the value of \$\lambda\$ in Ca<sub>05</sub>Na<sub>05</sub>Fe<sub>2</sub>As<sub>2</sub>. Finally, we will present our results on the anisotropy of \$\lambda\$ in a pnictide Ba-122 system, which shows a different temperature dependence from that of  $H_{c2}$ .

#### SB13

# Self-consistent calculations of the effects of disorder in d- and s-wave superconductors

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The kernel polynomial method (KPM) [1], as an approach for deriving highly accurate spectral properties of large sparse matrices, is very suitable for the numerical study on the Anderson localization problem. We investigate the effects of disorder in the two-dimensional inhomogeneous superconductors within the KPM by solving self-consistently the Bogoliubov-de Gennes equations [2, 3]. Comparing the behaviors of optical conductivity and the time evolutions of single particle Green's function in superconductor genese with different pairing symmetries, we find that the delocalized screening effect on the localized quasiparticle states is sufficiently larger in the d-wave superconductor. Within a considerable large range of disorder strength, the disorder-induced pseudogap state is derived close to the s-wave superconductor-insulator transition.

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### SB14

Superconductivity and structural transition of RPt<sub>2</sub>Si<sub>2</sub> (R = Y, La, Lu) Yutaro Nagano<sup>1\*</sup>, Nobutaka Araoka<sup>1</sup>, Akihiro Mitsuda<sup>1</sup>, Hirofumi Wada<sup>1</sup>, Masaki Ichihara<sup>2</sup>, Masahiko Isobe<sup>2</sup> and Yutaka Ueda<sup>2</sup> <sup>1</sup> Department of Physics, Kyushu University, Japan <sup>2</sup> Institute for Solid State Physics, University of Tokyo, Japan

The RPt<sub>2</sub>Si<sub>2</sub> compounds, where R stands for a rare-earth element, have a tetragonal CaBe<sub>2</sub>Ge<sub>2</sub>-type structure. LaPt<sub>2</sub>Si<sub>2</sub> and YPt<sub>2</sub>Si<sub>2</sub> are superconductors with Tc~2 K[1]. On the other hand, the compounds with R = La and Pr show anomalies in the temperature dependence of electrical resistivity[2,3]. These compounds are metallic but the resistivity shows a step at T\* ~110 K, which is attributable to a charge density wave (CDW) transition. In order to detect a periodic lattice distortion due to CDW, we performed selected area electron diffraction (SAED) and X-ray diffraction (XRD) measurements at low temperatures. It was found that (2 2 0) reflection peak was split into two sub-peaks with equal intensity below T\* in XRD measurements. This fact indicates that LaPt<sub>2</sub>Si<sub>2</sub> has an orthorhombic symmetry at low temperatures. The SAED patterns show (n/3 0 0) superlattice reflections with n = 1 and 2 below T\*. These results strongly support the formation of CDW in LaPt<sub>2</sub>Si<sub>2</sub>. In contrast, we did not observe any anomalies in the temperature dependence of resistivity except the superconducting transition for YPt<sub>2</sub>Si<sub>2</sub> and LuPt<sub>2</sub>Si<sub>2</sub>. These results are suggestive of no significant relation between CDW and superconductivity in the RPt<sub>2</sub>Si<sub>2</sub> compounds.

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#### **SB15**

# Magnetic property in ferromagnetic superconductor UGe<sub>2</sub> above ferromagnetic critical pressure

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UGe\_2 is a ferromagnet with the Curie temperature of 54 K at ambient pressure. The ferromagnetic transition temperature decreases with increasing pressure. The superconductivity appears in the pressure region of 1.0-1.6 GPa where UGe\_2 is in the ferromagnetic state. Ferromagnetic state disappears above  $P_c = 1.6$  GPa. Previous studies up to 2 GPa showed that the broad maximum appears in the temperature dependence of the ac magnetic susceptibility and that the ferromagnetic state is induced at metamagnetic field H\_m above  $P_c$ . These phenomena are resemblance to weak ferromagnetic compound in the 3d electron system such as  $ZrZn_2$ . In this study, we have measured the dc-magnetization up to 4.1 GPa using the ceramic anvil high pressure cell. We confirmed the maximum at T\_max in the temperature dependence of the dc magnetic susceptibility. With increasing pressure, T\_max is increased from 24 K at 1.9 GPa to 32 K at 2.6 GPa. The peak structure around T\_max becomes broader with increasing pressure and disappears above 4 GPa. The metamagnetic transition field H\_m increasing pressure. We discuss the experimental results with the phenomenological spin-fluctuation theory.

#### SB16

# Electron and hole transmission through superconductor - normal metal interfaces

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We have investigated the transmission of electrons and holes through interfaces between superconducting aluminum (T\_c = 1.2 K) and various normal metals (copper, gold, palladium, platinum, silver) using Andreev-reflection spectroscopy at T = 0.1 K. Most of the contacts had the typical Andreev reflection spectra and could be analyzed with the modified BTK theory. We included Dynes' lifetime as a fitting parameter  $\Gamma$  in addition to superconducting energy gap 2 $\Delta$  and normal reflection described by Z. For contact areas from 1 nm<sup>2</sup> to 10 000 nm<sup>2</sup> the BTK Z parameter was 0.5, corresponding to transmission coefficients of about 80%. This is considerably larger than estimated based on reports of other experimental methods (proximity effect and CPP resistance of multilayers) or expected theoretically. The small variation of Z indicates that the interfaces have a negligible dielectric tunneling barrier and that Fermi surface mismatch does not depend on the normal metal. Our results contradict standard expectations and call for a revision of the interpretation of the BTK theory and normal reflection. This is an especially important topic for applying Andreevreflection spectroscopy to unconventional superconductors or for measuring the local spin polarization of ferromagnets.

# SB17

# Magnetic field-induced odd-frequency superconductivity in s-wave superconductors

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The wavefunction of the Cooper pair of the conventional BCS state is characterized by the even parity (s-wave) and spin singlet state. It is well known that an external magnetic field modifies the wavefunction of the Cooper pair and breaks the conventional BCS superconducting state. We find that the wavefunction of the modified Cooper pair contains a spin triplet component with the s-wave orbital symmetry. While the spin triplet s-wave Cooper pair is forbidden owing to the fermion property, it is permitted in the generalized theoretical framework containing the oddfrequency dependence of the order parameter. We investigate this point on the basis of the self-consistent calculation in the Eliashberg theory and find that the odd-frequency order parameter coexists with the even-frequency one [1]. We also report how physical properties are affected by the induced odd-frequency superconducting order parameter [2].

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#### SB18

#### Enhanced pinning properties and the zero-resistance state in melttextured high-Tc superconductors probed by pulsed magnetic fields Fabio Teixeira Dias<sup>1\*</sup>, Valdemar Das Neves Vieira<sup>1</sup>, Douglas Langie Da Silva<sup>1</sup>, Sabrina Esperanca Nunes<sup>1</sup>, Frederik Wolff-Fabris<sup>2</sup>, Erik Kampert<sup>2</sup>, Jacob Schaf<sup>2</sup> and Joan Josep Roa Rovira<sup>4</sup> <sup>1</sup> Department of Physics, Universidade Federal de Pelotas, 96010-900, Pelotas, Brazil <sup>2</sup> Dresden High Magnetic Field Laboratory, HZ Dresden-Rossendorf, 01314, Dresden, Germany <sup>3</sup> Universidade Federal do Sul, 91501-970, Porto Alegre, Brazil <sup>4</sup> Universite de Poitiers, 86962, Poitiers, France

Among several techniques to grow high-temperature superconductors, melt-textured techniques have emerged as an alternative way to produce materials with physical properties needed for some technological applications. In this work we have studied the correlation between the magnetic inreversibility line and the zero-resistance state in melt-textured YBaCuO samples (Y123) with inclusions of Y211 phase particles by performing magnetization and magnetotransport measurements in DC and pulsed magnetic fields up to 60 r and temperatures down to 1.5 K [1]. Our results show for inplane magnetic fields an irreversibility field higher than 40 T at 77 K which evidences a strong pirming effect due to the interface of the Y211 particles embedded in the Y123 phase. We have observed effects of the intrinsic pinning due to the anisotropic superconductivity of the Y123 phase. Access to pulsed high magnetic fields reveals a suppression of the low temperature magnetic inversibility field relative to BC<sub>2</sub>. These results show experimental evidences for field-induced microscopic quantum fluctuations as observed in different high-Tc superconductors. Our findings in melt-textured YBCO samples exclude the possibility that the zero-resistance state and the magnetic inversibility merge in this field range [23,4], showing evidences of weak granularity and providing a good candidate for technological applications.

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### SB19

# **Phenomena of vortex pinning by composite pinning array on Nb films** Sheng Hao Wang<sup>1</sup>\*, Lance Horng<sup>1</sup>, T. C. Wu<sup>2</sup>, Chien-miao Chen<sup>1</sup>, R Cao<sup>1</sup> and J. C. Wu<sup>1</sup>

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In type II superconductor, magnetic fields can penetrate into the superconductor inhomogeneously as an array of vortices. The interaction of these vortices and vortexpin interactions in the superconductor give rise to a rich variety of static and dynamical phases. In our study, smaller pinning sites at the center of every hexagonal cell were added when fabricating film with honeycomb array to obtain films with composite pinning array. It has been prepared on Si<sub>3</sub>N<sub>4</sub>-coated Si wafers using electron-beam lithography in conjunction with reactive ion etching. Then a dc sputtering completed the four-terminal geometry niobium films over the composite array. MR measurements were carried out by a four-probe technique. Special matching peaks of critical currents are exhibited at different temperatures in the measurements. The positions and structures of the matching peaks seem irregular and are very different from the structures of the matching parays. Considering the multiple-vortex filling of different pinning sites, we can give a reasonable explanation to this interesting phenomenon.

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### **SB20**

#### Superconducting fluctuations near the Mott critical point

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The Nernst effect, the voltage that occurs transverse to crossed temperature gradient and magnetic field applied to a conductor, is a sensitive probe of superconducting fluctuations in unconventional superconductors such as the high-Tc cuprates [1] and organic superconductors [2]. The isostructural kappa-phase organic compounds (BEDT-TTF)2X offer a series of materials ranging from weakly correlated superconductors ( $X = Cu(NCS)_2$ ) through strongly correlated superconductors ( $X = Cu[(N(CN)_2]Br)$  to Mott insulators ( $X = Cu[(N(CN)_2]C)$ ) as the bandwidth decreases and the importance of correlations increases. We previously showed that while there are no fluctuations significantly above Tc in the weakly correlated superconductor [2], suggesting that the proximity of the Mott state drives the fluctuating state. Here, by alloying the strongly correlated and the Mott insulating materials, we tune the system towards the Mott critical point. Superconductors in the vicinity of the critical point exhibit a large and magnetic-field-dependent Nernst effect, characteristic of superconducting fluctuations, at even higher temperatures. The effect is apparent over the whole temperature range for which the conductivity is metallic, up to 50K, about six times Tc.

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#### **SB21**

#### Anomalous hall effect in superconductors with spin-orbit interaction

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### SB22

Analysis of the local current density in HTS coated conductors using low-temperature scanning laser and hall probe microscopy Sang Kook Park, Bo Ram Cho, Hee Yeon Park and Hyeong-cheol Ri\* Department of Physics, Kyungpook National University, Korea

Local current density in HTS coated conductors was investigated using Lowtemperature Scanning Laser and Hall Probe Microscopy (LTSLHPM). We prepared GdBCO and YBCO coated conductors to study the spatial distribution of the current density in a single bridge. Inhomogeneity of the local critical temperature in the bridge was analyzed from experimental results of scanning laser microscopy near the superconducting transition. The local transport and screening current in the bridge were also investigated using scanning Hall probe microscopy. From experimental results of scanning laser and Hall probe microscopy, we have observed the redistribution of current density and the critical temperature caused by defects of the HTS layer.

# **SB23**

#### Cooper pairing between conduction and localized electrons in heavyfermion systems

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We study the Cooper pairing between a conduction electron (c electron) and a localized f electron, referred to as the c-f paring, as a mechanism of s-wave superconductivity in heavy-fermion systems. To describe the physics of heavy-fermion systems, we start from the periodic Anderson model. We apply the Schrieffer-Wolff transformation to this model considering the situation where the f level is sufficiently deep and the Coulomb repulsion is finite but large. The resulting effective Hamiltonian includes direct and spin-exchange interactions between c and f electrons, which can be the driving force for the formation of the Cooper pairs. Within the mean-field approximation, we show that the c-f pairing state with anisotropic s-wave symmetry appears in a large region of the phase diagram. Such a pairing state can be realized when the quasiparticle band of the f electrons, which is formed by the strong Coulomb repulsion, is located near the energy level of the conduction band. We also discuss the relationship between the c-f pairing proposed here and the s-wave superconductivity found in the experiments on CeRu, and CeCo,

#### **SB24**

#### Low-temperature thermoelectric properties of the electron-doped Perovskites Sr<sub>1x</sub>CaxTi<sub>1x</sub>NbyO<sub>3</sub>

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Electron-doped perovskite SrTiO<sub>3</sub> is known to be one of candidates for an n-type thermoelectric (TE) oxide. [1] In this study, we tried to improve a TE property of a lightly-electron-doped SrTiO<sub>3</sub> single crystal below room temperature by substitutions of Ca and Nb for Sr and Ti. We found that SrTi<sub>0.99</sub>Nb<sub>0.01</sub>O<sub>3</sub> shows a large power factor of about 90 µW/K2 cm at 50 K and the largest dimensionless TE figure-of-merit (ZT ~ 0.07) below 40 K among the ever-reported materials. The good n-type TE response is due to a distinct electron-phonon interaction which could relate to the superconducting state [2], which enlarges a Seebeck coefficient not through an enhancement of effective mass, but through a change of relaxation time. On the other hand, a Ca2+ substitution for Sr<sup>2+</sup> increases the ZT at 300 K for Sr<sub>1x</sub>CaxTi<sub>0.97</sub>Nb<sub>0.03</sub>O<sub>3</sub> from 0.08 to 0.105. The enhancement of ZT originates not only in a large reduction of a thermal conductivity due to an introduced randomness into the crystal structure, but also in an unexpected enhancement of Seebeck coefficient by the Ca substitution.

[1] T. Okuda, K. Nakanishi, S. Miyasaka, and Y. Tokura, Phys. Rev. B 63, 113104 (2001). [2] H. Suzuki, H. Bando, Y. Ootuka, I.H. Inoue, T. Yamamoto, K. Takahashi, and Y. Nishihara, J. Phys. Soc. Jpn. 65, 1529 (1996). [3] J. Fukuyado, K. Narikiyo, M. Akaki, H. Kuwahara, and T. Okuda, Phys. Rev. B, to be published

#### **SB25**

#### Evolution of pairing potential in ladder materials under renormalization group transformations

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It is generally believed that electron-electron interactions are responsible for pair formation in unconventional superconductors. However, how electrons can form pairs in the presence of repulsive interactions remains counterintuitive and puzzling. Many studies suggest that spin fluctuations provides the "pairing glues" to account for the unconventional pairing. Here we take a different route: finding effective pairing potential at different length scales by renormalization group approach. We focus on the unconventional superconductivity in ladder materials and show how the profile of electron-electron interactions evolves under renormalization group transformation. Starting with repulsive short-ranged interaction, compelling evidence for spatially structured attractive interactions emerges in long-wavelength limit and explains the unconventional pairing symmetry naturally. Comparisons with other theoretical approaches and potential generalization to two dimensions are discussed at the end.

#### **SB26**

#### Evidence for multiband order parameters in the stong-coupling LaRu As., Skutterudite Superconductor

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Superconducting properties of high-purity LaRu4As12 single crystals with the highest critical temperature Tc = 10.45 K among the fully filled-skutterudite superconductors were investigated by dc magnetization, electrical resistivity, torque, and specific heat measurements both as a function of temperature and magnetic field. At B = 0 and below about Tc/2, the electronic specific heat exhibits two-band features, as concluded from an α-model analysis. Multi-hand effects are further inferred from a nonlinear magnetic-field dependence of the electronic specific heat coefficient in the zerotemperature limit as well as from a positive curvature of the upper critical field in the vicinity of Tc. Our findings, combined with results of previous de Haas-van Alphen experiments, suggest that enhanced superconductivity in LaRu<sub>4</sub>As<sub>12</sub>appears along with a spherical Fermi-surface sheet enclosing a small volume in the center of the Brillouin zone

#### **SB27**

NMR Study of magnetic order and the FFLO state in CeCoIn<sub>5</sub>

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A strongly-correlated superconductor, CeCoIn<sub>5</sub> is believed to host a Fulde-Ferrell-Larkin-Ovchinnkov (FFLO) state in a restricted region of high field and very low temperature. In addition to the field evolution of NMR results for H//a-axis [1], we will report recent results of the angle, $\theta$ , dependence of external fields with respect to the ab planes down to ~50mK. Detail phase diagram for  $\theta\text{-},$  B-, and T-parameters is obtained. The NMR spectra change dramatically upon entering the novel SC phase. A well-separated peak structure at the In(2b) site suggests the occurrence of the magnetic order which is emerging only in the newly-discovered SC state. The structure of the spin density wave (SDW) is modified with angle and field. The Knight shift of CeCoIn<sub>5</sub> provides a direct evidence for the emergence of the spatially-distributed normal quasiparticle region. The quantitative analysis for the field evolution of the paramagnetic magnetization and low-lying energy quasiparticle density of state is consistent with the nodal plane formation, which is characterized by an order parameter in the FFLO state

[1] K. Kumagai et al., Phys. Rev. Lett. 106, 137004 (2011)

### **SB28**

#### MgB2 coated conductors grown at various temperatures by hybrid physical-chemical vapor deposition

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For practical applications, MgB<sub>2</sub> superconductor ( $T_c \sim 39$  K) in the form of wires and tapes are required to be produced in long lengths to replace the conventional metallic superconductors, such as Nb-Ti and Nb<sub>3</sub>Sn. In this work, we have grown MgB<sub>2</sub> coated conductors by depositing MgB2 films directly on the flexible Hastelloy tapes at various temperatures ranging from 520 to 600 °C by using HPCVD. The MgB<sub>2</sub> coated conductors exhibit critical temperatures ranging from 37.5 to 38.5 K with superconducting transition width  $(T_c)$  of about 0.3?0.8 K. X-ray diffraction analysis revealed that the MgB<sub>2</sub> coated conductors are polycrystalline in nature. It was found that the MgB<sub>2</sub> coated conductor grown at 520 °C had dense microstructure with good grain connectivity. However, upon increasing the growth temperature from 520 to 600 °C deterioration of grain connectivity takes place and void formations were observed. The critical current density (Jc) of the order of Jc(5 K, 0 T)~107 and Jc(5 K, 3 T)~105 A/cm2 was obtained for the MgB, coated conductor fabricated at 520 °C. The fabrication process, crystallographic orientations, surface morphologies and superconducting properties [T<sub>c</sub>, Jc (H)] of MgB<sub>2</sub> coated conductors grown at various temperatures will be discussed in detail

# **SB29**

#### Low temperature properties of the weakly-coupled noncentrosymmetric superconductor LaNiC,

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In a non-centrosymmetric (NCS) superconductor, an asymmetric potential gradient yields an antisymmetric spin-orbit coupling (ASOC), which splits the degeneracy of conduction electron, and allows the admixture of spin-singlet and spin-triplet pairing states, leading to accidental nodes in the energy gap [1,2]. It was shown that the time reversal symmetry is broken in the weakly correlated intermetallic compound LaNiC, [3] but its pairing symmetry is still under debate. Here we present the low temperature specific heat C(T) and the magnetic penetration depth  $\lambda$ (T) for LaNiC<sub>2</sub>. It is found that both  $\lambda(T)$  and Ce(T)/T can be well described by a phenomenological two-gap model. The Sommerfeld parameter,  $\gamma(H)$ , increases steeply at low fields and eventually get saturated with increasing magnetic field, being also consistent with the behavior of two-gap superconductivity. We argue that the ASOC is weak in LaNiC2 and the spin-singlet state dominates in the pairing states, giving rise to two-gap-like BCS superconductivity.

[1] P. A. Frigeri, D. F. Agterberg, A. Koga, and M. Sigrist, Phys. Rev. Lett. 92, 097001 (2004). [2] H. Q. Yuan, D. F. Agterberg, N. Hayashi, P. Badica, D. Vandervelde, K. Togano, M. Sigrist, M. B. Salamon, Phy Rev. Lett. 97, 017006 (2006). [3] A. D. Hillier, J. Quintanilla, R Cywinski, Phys. Rev. Lett. 102, 117007 (2009).

#### **SB30**

#### Magnetic penetration depth and \$H\$-\$T\$ phase diagram in SrPd<sub>2</sub>Ge<sub>2</sub> H. Kim<sup>1</sup>, N. H. Sung<sup>2</sup>, M. A. Tanatar<sup>1</sup>, B. K. Cho<sup>2</sup> and R. Prozorov<sup>1</sup>\*

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In-plane magnetic penetration depth, \$\lambda(T)\$, was measured in single crystals of SrPd<sub>2</sub>Ge<sup>2</sup> superconductor down to 50 mK and in magnetic fields up to \$H\_{dc} = 1\$ T by using a self - oscillating tunnel diode resonator operating at 16 MHz with amplitude of \$H\_{ac} ≈ 20\$ mOe. In the pure Meissner state, at \$H {dc}=0\$, \$ \lambda(T)\$ saturates exponentially approaching \$T\rightarrow 0\$ indicating fully gapped superconductivity in SrPd\$\_2\$Ge\$\_2\$. In a magnetic field, the measured penetration depth is given by  $\lambda^2 = \lambda^2_L +$ \lambda^2 C\$, where \$\lambda^2 C ~ B/j^2 c\$ is the Campbell penetration depth and \$j c\$ is the critical current density. For \$H < 0.4\$ T, the shortest \$\lambda\_C\$ (strongest pinning) is achieved not at the lowest, but at some intermediated temperature. Additionally, \$\Delta \lambda(T)\$ shows hysteresis between zero field cooling and field cooling measurements which is an indication of non - parabolic pinning potential. Another interesting feature is a pronounced finite diamagnetic response above bulk superconducting transition, which could be related to superconducting fluctuations. Combining all measurements, the entire \$H-T\$ phase diagram of SrPd<sub>2</sub>Ge<sub>2</sub> has been constructed. Possible pairing mechanism and the superconducting gap structure, unusual vortex response and field - and temperature dependent critical current density will be discussed

### **SB31**

#### Eliashberg function of the overdoped Bi2212 superconductors deduced from the high resolution laser ARPES intensity

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We report the diagonal and off-diagonal Eliashberg functions a2F(+) and a2F(-) of overdoped Bi2212 supercondcutors deduced from the high resolution laser APRES experiments. By fitting the momentum distribution curves of the ARPES intensity employing the superconducting Green's function, the diagonal and off-diagonal selfenergies along several cuts in the Brillouin zone are extracted. Then, the diagonal and off-diagonal Eliashberg functions are deduced by inverting the d-wave Eliashberg equation using, respectively, the conrresponding self-energies. We will present the momentum and frequency dependences of the Eliashberg function a2F(+) and a2F(-). We then compare these results with other experiment and calculation results

### **SB32**

### The momentum and frequency dependences of the self-energy induced by the spin fluctuations for the cuprate superconductors

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We solve the momentum resolved d-wave Eliashberg equation to investigate the consistency of the spin fluctuation induced superconductivity theory for the cuprates. The effective interaction between electrons (Eliashberg function) is modeled in terms of the magnetic excitation spectrum measured by the inelastic neutron scattering on the LSCO superconductors reported by Vignolle et al[1]. The magnetic excitation spectrum consists of three parts: the resonance mode near 40-70 meV and the incommensurate parts above and below the resonance mode. The resonance mode has a peak at  $(\pi,\pi)$ . The low energy incommensurate part makes a nonmonotonic momentum dependence of the diagonal and off-diagonal self-energies. On the other hand, the resonance mode makes a monotonic momentum dependence of the self-energy. The angle dependence of the superconducting gap is influenced by the effects of these two energy scales. These results are compared with the ARPES experiments.

[1] B. Vignolle et al., Natrue Physics 3, 163 (2007)

### SC01

Novel non-centrosymmetric superconductors in 113 and 111 crystal structures Friedrich Kneidinger<sup>1</sup>\*, Ernst Bauer<sup>1</sup>, Herwig Michor<sup>1</sup>, Gerfried Hilscher<sup>1</sup>, Isolde Zeiringer<sup>2</sup>, Peter Rogl<sup>2</sup>, Nataliya Melnychenko<sup>3</sup>, Leonid Salamakha<sup>4</sup> and Adrian Hillier<sup>5</sup>

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Superconductivity (SC) in materials that don't possess a center of inversion due to their crystal structure has been recently found in materials like CePt\_3Si[1]. The lack of inversion-symmetry gives rise to an electrical field and therefore leads to an antisymmetric spin-orbit coupling (ASOC) of the Rashba-type[2]. This causes apditting of the Fermi surface, thereby leading to a superposition of spin-niglet and spin-triplet Cooper pairs in the SC condensate. Intriguing features like nodes in the superconducting gap, as well as H\_c2 exceeding the Pauli-Clogston limit, may occur. However, besides Li 2Pt 3B[3], spin-triplet pairing seems only to be dominant in heavy fermion systems. In order to study these features in further detail, new ternary compounds like BaPKis [34], SrPdSi 3, SrPdGe 3[5], SrPHSi 3 and SrPGe 3 of BaNiSn 3 type (space group 1/4mm) as well as promising 111-types like LaPSi[6,7] and LaIrSi[8] have been investigated, since there is a tuning possibility regarding the quadratic atomic number dependency ( $Z^2$ ) of the spin-orbit coupling. Measurements of the electrical resistivity and the specific heat have been conducted to analyze the superconducting ground state Additionally. uSR measurements have been performed to proof the microscopic pairing mechanism of the Cooper pairs. The results obtained so far indicate s-wave fully gapped BCS SC

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 K. Miliyanchuk et al., Journal of Physics: Conference Series 273, 012078 (2011) [6] K. Klepp & E. Parthe, Acta Crystallogr. B38, 1105 (1982) [7] J. Evers, G. Oehlinger and A. Weiss, Solid State Communications, Iol. 50, No. 1, pp. 61-62, (1984) [8] K. Klepp & E.Parthe, Acta Crystallogr. B38, 1541 (1982)

### SC02

#### Non-fermi liquid behavior of d-wave superconductor

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The d+id-density wave (Chiral DDW) order, at the anti-ferromagnetic wave vector  $Q = (\pi,\pi)$ , is assumed to represent the pseudo-gap (PG) state of a hole-doped cubrate superconductor. The d-wave superconductivity (DSC), driven by an assumed attractive interaction is discussed within the BCS framework together with the d+id ordering The single-particle excitation spectrum in the CDDW + DSC state is characterized by the Bogoluibov quasi-particle bands-a characteristic feature of SC state. The coupled gap equations are solved self-consistently together with the equation to determine the chemical potential (u). With the pinning of the van Hove-singularities close to  $\mu$ , one is able to calculate the thermodynamic properties of the under-doped cuprates in a consistent manner. Apart from the known facts that PG and DSC represent two competing orders as the former brings about a depletion of the spectral weight available for pairing in the anti-nodal region of momentum space and the depletion of the spectral weight below Tc at energies larger than the gap amplitude, we find a conspicuous feature that the electronic specific heat displays non-Fermi liquid behavior( anomalous temperature dependence) in the CDDW + DSC state

### SC03

#### Low temperature enhancement of the critical current in CeCoIn<sub>5</sub>. Possible signature of magnetic order

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We report on vortex dynamics, dc magnetization and RF penetration depth measurements in the heavy fermion compound CeCoIn<sub>5</sub> down to low temperatures, T=50 mK. No strong pining is observed as expected for a clean limit superconductor which does not break spontaneously any additional symmetries besides the U(1)-gauge. CeCoIn<sub>5</sub> display clear logarithmic relaxation curves as expected from Kim-Anderson theory. The temperature dependence of the relaxation rate, S, with a small but finite residual value indicate that quantum tunneling plays a significant role in the vortex creep only at very low temperatures. Remarkably, a new phase transition marked by a strong increase in the remnant magnetization, Mrem is observed around T = 0.3 K in low magnetic fields just high enough to put the sample in the Bean critical state. Moreover, this anomaly is corroborated by the DC magnetization and RF measurements at very low fields. We extended the vortex dynamics investigation to Pb irradiated CeCoIn<sub>5</sub>. While the defects created by irradiation have a clear effect on the relaxation rates the enhancement of Mrem still takes place at the same temperature. Our findings are consistent with the existence of a magnetically ordered phase, deep inside the superconducting state.

#### SC04

# Superconducting phase diagram in fcc phase of Cs<sub>3</sub>C<sub>60</sub>; A pressure dependence of resistivity

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 $\rm Cs_5C_{60}$  crystallizes three polymorphic phases;  $\rm A_{15},$  fcc and bco. Two polymorphs,  $\rm A_{15}$  and fcc phases, show pressure induced superconducting transitions around 38 K and 35 K, respectively. Here we report a resistivity for fcc phase of  $\rm Cs_5C_{60}$  as a function of pressure and present a dome-shaped superconducting phase diagram.

#### SC05

# The superconducting phases of URu<sub>2</sub>Si<sub>2</sub> from sound velocity measurements

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High resolution measurements of the changes in the longitudinal sound velocity in a magnetic field have been performed in the zero temperature limit in high quality single crystals of URu<sub>2</sub>Si<sub>2</sub>. These measurements reveal two distinct signatures attributable to multiple superconducting phases. A step change in the sound velocity, for propagation in the basal plane, is observed at the critical field, Bc<sub>2</sub>. This step broadens considerably as T tends to Tc with a concomitant decrease in magnitude. A second step is observed at a field -0.5 Bc<sub>2</sub> and it's magnitude remains constant at all temperatures. Inductive measurements of the transitions in a magnetic field, however, reveal only a single signature which coincides with the lower step change in the sound velocity. Measurements performed with B oriented at various angles between the a and c-axes reveal a weaker angular dependence of the lower step and confirm the rapid fall off of Bc<sub>2</sub> close to B||c-axis. A multiple superconducting phase diagram for URu<sub>2</sub>Si<sub>2</sub> is proposed.

#### SC06

#### Electronic structure of a superconducting boride, ZrB<sub>12</sub>

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Borides have attracted a great deal of attention followed by the finding of varied interesting properties. ZrB<sub>12</sub> is one such compound exhibiting highest superconducting critical temperature (Tc = 6 K) in MB12 family. In order to probe the electronic structure of single crystalline ZrB<sub>12</sub> using high resolution photoemission spectroscopy and ab initio band structure calculations. Zr 3d core level spectra exhibit signature of satellite features indicating finite correlation among the conduction electrons. The B 1s spectra appears to be sharp with unusual asymmetry, presumably due to surface contributions as found in hexaborides. In the valence band spectra, the spectral density of states reveal a dip at the Fermi level, which gradually increases with the decrease in temperature down to 10 K. The band structure results indicate dominant contribution from the B 2p electronic states in the valence band. However, the photon energy dependence in the experimental spectra indicate large contribution from the Zr 4d states near Fermi level suggesting their important role in the superconductivity.

#### **SC07**

# On the nature of an energy barrier between $(\pi, 0)$ and $(0, \pi)$ magnetic orders in Fe pnictides

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As temperature is lowered most of undoped Fe arsenides, parent compounds for recently discovered Fe based superconductors, undergo a transition into a collinear state with stripe-like magnetic order in which anti-ferromagnetic (AFM) Fe chains are ferromagnetically ordered along the direction perpendicular to the chains. Two such collinear magnetic structures, characterized by ordering vectors  $(\pi, 0)$  or  $(0, \pi)$ , are connected by infinite number of non-collinear states with two AFM sublattices of second Fe neighbors rotated by an arbitrary angle with respect to each other. In a classical Heisenberg model all these states are degenerate. Band structure calculation show, however, that the degeneracy is lifted already at the mean field LSDA level and that in Fe arsenides  $(\pi, 0)$  and  $(0, \pi)$  magnetic orders are separated by an energy barrier comparable to the energy difference between Neel and stripe AFM orders. We discuss a microscopic origin of the energy barrier and demonstrate that it is caused by closely related to underlying band structure. The results for Fe arsenides are compared to BaMn<sub>2</sub>As<sub>2</sub> and hypothetical KFe<sub>2</sub>Se<sub>2</sub> for which we found that a non-collinear 90-degree spin arrangement is more favorable than collinear ones. A doping dependence of the barrier is also discussed

#### **SC08**

#### Electronic structures and magnetic properties of LnFeAsO

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Recently, high transition-temperature superconductivity was discovered in the ironpnictide RFeAsO (R=rare-earths) family of materials. Among them, LaFeAsO undergoes a tetragonal-to-orthorhombic phase transition at ~160 K followed by an antiferromagnetic ordering at ~145 K. This material has also been reported to undergo a spin-density wave transition near 150 K. Through first-principles calculations we tried to understand the electronic structures and magnetic properties of LaFeAsO. We used full-potential linearized-augmented-planewave method in the antiferromagnetic ground state with orthorhombic phase. The physical origin of the SDW formation will be discussed.

# SC09

# Phase transition of a heavy fermion superconductor in a high magnetic field : entanglement analysis

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When the magnetic field is acting on the spin of electrons only, a transition from a normal to a modulated superconducting state or FFLO superconductor state may occur at low temperatures[1-3]. A FFLO superconducting state, accompany with an order parameter that oscillates spatially, may be stabilized by a high applied magnetic field or a molecular field. CeCoIn, is a sample of heavy-fermion superconductor[4]. Initially experiments on CeCoIn, indicate that in this substance a FFLO state in a exchange field is realized. Quantum multipartite entanglement is a new procedure for investigating quantum phase transition[5-6], which is one of the interesting topics in condensed matter. In this article, we deal with to the phase transition of FFLO state of CeCoIn, to the normal state by obtaining quantum multipartite entanglement of the system. For this purpose, using normal and anomalous Green functions and density matrix, concurrence, as a measure of bipartite entanglement, is obtained. Then, order parameter tandition is determined and the behavior of the system based on order parameter is discussed. Furthermore, the phase transition of both BCS and FFLO states to the normal state are commared.

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#### SC10

# Spin and charge excitations in antiferromagnetic metallic phase in multi-orbital systems: A case study of chromium

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Since the discovery of Fe-based superconductors, it has been recognized that the interplay of spin, charge, and orbital degrees of freedom is the key to understanding the physics of multiband itinerant systems. Antiferromagnetic (AF) phase seen in the parent compounds of Fe-based superconductors is a good example to study the multi-orbital physics in the metallic phase. The AF metallic phase similar to the parent compounds is realized in a simple chromium metal with AF order. To make a unified view in the AF metallic phase of multi-orbital systems, we study spin and charge excitations in AF Cr. We calculate the optical conductivity, neutron scattering spectra, and Cr L-edge resonant X-ray scattering (RIXS) spectra based on the mean-field theory for the multi-band Hubbard model and the random-phase approximation. The calculated results of the optical conductivity and neutron scattering spectra are compared with corresponding experimental data. The d etailed polarization and momentum dependence in RIXS is also discussed with predictions for the experiments.

### SC11

The dirty crossover - signature of a robust superfluid in the unitary regime. IIT Guwahati IIT Guwahati

We perform a systematic study of the BCS-BEC crossover in an ultracold fermionic gas in presence of a weak white noise-like random disorder which is incorporated in our mean-field treatment via Gaussian fluctuations. A careful investigation of different mean field quantities as a function of the interparticle interaction reveals a nonmonotonic behavior in the the condensate fraction data as disorder is included in our calculations. The inflexion point is seen to occur in the unitary regime, and afterwards there is a gradual depletion of the condensate density in the BEC side, thereby corroborating a robust paradigm of superfluidity in the unitary regime. Motivated by this result, we study the spectral gap and the density of states in the crossover region and their responses to the disorder. We further include a comparative discussion on the role of the pair and phase coherence lengths in this region in presence of disorder to assess how these results can be relevant to the ongoing discussion. Due to tremendous progress taking place in the experimental front in manipulating atoms on an optical lattice in presence of (speckle) disorder, a stable superfluid scenario in the unitary interpartice in the unitary regime.

Prof. S.S. Mandal, Indian Association for Cultivation of Science, Kolkata, India Prof. Y. Yanase, Niigata University, Japan

# SC12

# Spin-orbit coupling and the superconductivity in simple-cubic polonium

Chang-jong Kang, Kyoo Kim and B. I. Min\* Physics, POSTECH, Korea

Polonium is the only element which has the simple-cubic (SC) structure in the periodic table. We have studied its structural stability based on the phonon dispersion calculations using the first-principles all-electron full-potential band method. We have demonstrated that the strong spin-orbit coupling (SOC) in SC-Po suppresses the Peierls instability and makes the SC structure stable. Further, we have investigated the possible superconductivity in SC-Po, and predicted that it becomes a superconductor with Tc  $\sim$  4 K at ambient pressure. The transverse soft phonon mode at q  $\sim$  2/3 R, which is greatly affected by the SOC, plays an important role both in the structural stability and the superconductivity in SC-Po. We have explored effects of the SOC and the volume variation on the phonon dispersions and superconducting properties of SC-Po.

### SC13

#### Stability of FFLO states in optical lattices with layered structure

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Recently, ultracold fermions with spin-imbalanced populations have attracted much interest. One of the interesting questions is how the Fulde-Ferrell-Larkin-Ovchinnikov (FFLO) phase, where Cooper pairs are formed with a finite total momentum, is realized in the system. In the two dimensional optical lattice with a confining potential, it has been pointed out that two kinds of FFLO states are realized[1,2]. On the other hand, such FFLO states with a three-dimensional structure have not been studied so well although the interlayer coupling should be important for realistic optical lattice systems. Motivated by this, we investigate the stability of the superfluid state in a layered fermionic optical lattice system with a confining potential, using the Bogoliubov de-Gennes equations. It is clarified that in the imbalanced case, the introduction of the interlayer hopping stabilizes the radial FFLO state, while makes the angular FFLO state. It is clarified that in a certain ring region the A-FFLO state is indeed realized in a layered state. It is clarified that in a certain ring region the A-FFLO state is indeed realized in a large system.

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### SC14

# Flux quantization and its magnetic relaxation in a micrometer-sized superconducting ring of niobium

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We developed a high-sensitivity cantilever force magnetometry, which is capable of direct measurement of the magnetic moment of micron-scale samples at low magnetic fields and low temperatures, and studied magnetic properties of a 100 nm-thick Nb superconducting ring with 4 micrometer radius and 2 micrometer width. At a temperature of 3.5 K, we observed flux quantization of h/2e period using ac torque and dc force magnetometries with attonewton and sub-piconewton sensitivity, respectively, determining the magnetic moment of each flux quantum. First demonstration of magnetic relaxation at single flux quantum level is also presented, which provides information about the dynamics of thermally-activated single flux quantums depending on their number.

# SC15

# Measurement of rat biomagnetic signals by using a HTS-SQUID system

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We have measured biomagnetic signals in rats by using a high-TC SQUID magnetometer system employing a small magnetically shielded box. For the rat cardiac magnetic field measurements, healthy Wistar Kyoto rat and spontaneously hypertensive rat were anesthetized and fixed on a plastic xyz-stage in supine position inside the magnetically shielded box. Well defined P- , QRS- and T-waves were observed on the rat cardiac magnetic field measurements. We investigated circular measurements based on the hexaxial reference system, i.e., 12 measurement points were determined alongside the circumference of a circle centered at the heart position. These results were quite promising for use in biomedical applications of heart disease rat model study by using the high-TC SQUID magnetometer system.

# SC18

**SC19** 

Haigun Lee\*

#### Design, fabrication, and testing of a cooling system using solid nitrogen for a 3 T/60-mm RT bore superconducting HGMS

Jung-bin Song, Kwang Lok Kim, Dong Gyu Yang, Yoon Hyuck Choi, Jongseok Lee and Haigun Lee\*

Department of Materials Science and Engineering, Korea University, Korea



Purification of chemical mechanical polishing wastewater using a 2G

Dong Gyu Yang, Jung-bin Song, Young-gyun Kim, Jongseok Lee, Yeonjoo Park and

MOVED

Department of Materials Science and Engineering, Korea University, Korea

HTS high gradient magnetic separation system

### SC16

# Evolution of the effective mass approaching the quantum critical point in the heavy fermion superconductor $CePt_2In_7$

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We report Shubnikov-de Haas measurements on micro-structured CePt<sub>2</sub>In<sub>7</sub> single crystals under high pressures. CePt<sub>2</sub>In<sub>7</sub> belongs to the Ce<sub>m</sub>MnIn<sub>3m+2n</sub> heavy fermion family and orders antiferromagnetically below 5.5 K. Applying pressure induces a superconducting phase and suppresses the AFM order with the Neel temperature extrapolating to a quantum critical point. The magnetic fluctuations associated with the QCP are thought to stabilize the unconventional superconducting phase. The evolution of the effective electronic masses approaching the quantum critical point is essential to the understanding of the electronic correlation leading to unconventional superconductivity. To this end we have developed a method to contact micro-structured samples inside a diamond anvil cell allowing for complex sample shapes and several electric contacts.

### SC17

# Thermal stability of an epoxy-impregnated HTS racetrack coil without turn-to-turn insulation for rotating machines

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### SC20

# Effect of liquid cryogen on a 2G HTS magnet using a mixed cryogen cooling system

Kwang Lok Kim, Jung-bin Song, Yoon Hyuck Choi, Dong Gyu Yang and Haigun Lee\*

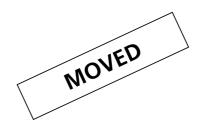
Department of Materials Science and Engineering, Korea University, Korea



# SC21

# Removal of silica and copper ions from CMP wastewater via magnetic seeding aggregation using superconducting HGMS

Jongseok Lee, Jung-bin Song, Dong Gyu Yang, Yeonjoo Park and Haigun Lee\* Department of Materials Science and Engineering, Korea University, Korea



### SC22

#### Removal of silica and copper ions from CMP wastewater via magnetic seeding aggregation using superconducting HGMS Jongseok Lee, Jung-bin Song, Dong Gyu Yang, Yeonjoo Park and Haigun Lee\*

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### SC23

#### Powder neutron diffraction study of HoCoGa<sub>5</sub>

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Ternary compound ACoGa<sub>5</sub>(A=rare earth or actinide) attracts a much attention due to high temperature superconductivity in PuCoGa<sub>5</sub> and related compounds. In addition, recent study on TbCoGa, reveals another interesting aspects of this system ; successive transitions in TbCoGa, could be "components-separated orders" which may arise from a novel type of frustration between multipole. HoCoGa5 also exhibits successive transitions at TN1 = 9.6 K and TN2 = 7.5 K. In order to investigate nature of these transitions in HoCoGa<sub>5</sub>, neutron powder diffraction experiment were carried out. Appearance of superlattice refractions was clearly observed below TN1. In addition, we observed that both peak position and intensity of superlattice peaks abruptly change at TN2 on cooling. The observed peak in the phase I(T < TN2) and II(TN2 < T < TN1) can be described with q1=(1/2 0 1/2) and q2=(1/2 0  $\tau$ ) with  $\tau$ =0.35, respectively. The determined propagation vector of the phase I is the same as that of TbCoGa5, whereas that of phase I I is altered. Detailed magnetic structure analysis for both phases are in progress.

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# SC24

### Evidence of two-band gap superconductivity in LaRu<sub>2</sub>P<sub>2</sub>

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We investigate the electronic transport properties and the superconducting gap characteristics of LaRu<sub>2</sub>P<sub>2</sub>. Superconducting transition was recognized at Tc = 4.4 K in LaRu<sub>2</sub>P<sub>2</sub> single crystal as reported in ref [1]. Values of upper critical field Hc<sub>2</sub> estimated from the electrical resistivity under magnetic field are isotropic in this system. These are good accordance with the results by Ying et al.[2]. However, they, although, claimed that LaRu<sub>2</sub>P<sub>2</sub> is a conventional superconductor, so called BCS type, it was revealed that LaRu<sub>2</sub>P<sub>2</sub> single crystal obtained in this work shows a remarkably larger Hc<sub>2</sub> than that expected from GL theory at considerably low temperature. In addition, we observed two clear discontinuities of specific heat at Tc<sub>1</sub> = 4.4 K and Tc<sub>2</sub> = 4.0 K, respectively, in specific heat as a function of temperature. Each transition temperature vary obeying Hc<sub>2</sub>(t)=Hc<sub>2</sub>(0)(1+2)/(1+(2)), here, t = T/Tc, indicating that both discontinuities at Tc<sub>1</sub> and Tc<sub>2</sub> are attributed to superconducting transitions. So, we tried to analyze thermodynamic properties of LaRu<sub>2</sub>P<sub>2</sub> by a two-band gap model [3]. In consequence, the numerically calculated C(T) curve was well reproduced experimentally measured specific heat as a function of temperature assumed to be 1.5 and 3.0, respectively.

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#### **SC25**

# Proximity effect for asymmetrical three layered F/S structures in external magnetic field

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We study the critical temperature Tc of the three-layered ferromagnet/superconductor (F/S) structures at the external magnetic field H parallel to the film. For the  $F_1/S/F_2$  and  $F_1/F_2/S$  asymmetrical trilayers the triplet superconducting component is generated at noncollinear magnetizations of the F layers. Assuming that all S and F layers are dirty we solve boundary problem for the Usadel function. The results of numerical calculations for Tc as function thicknesses both  $F_1$  and  $F_2$  layers at various parameters F/S structure are resented. The application to the spin-switch problem is discussed. We found that asymmetry can essentially change the spin-switch observation condition.

### SC26

# Influence of proximity effect with Umklapp processes on the Josephson current in the SFS nanostructure

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We consider the Josephson effect in symmetric superconductor-ferromagneticsuperconductor (SFS) system. The Josephson current is calculated as a function of the ferromagnetic layer thickness d. The Fulde-Ferrell-Larkin-Ovchinnikov (FFLO) pairs in the F layer have a nonzero wave vector, and the transformation the BCS pairs to the FFLO pairs (and vice versa) at passing through SF (FS) borders may be proceeded through the Umklapp processes [1]. To estimate the influence of proximity effect we use the expansion of superconductor energy in the powers of an order parameter (near the critical temperature Tc). We take into account the dependence of the critical temperature Tc from the phase difference  $\varphi$  between the order parameter in the left and right S side of the SFS contact. Our results are compared with known data for the Nb/ CuNi/Nb nanostructures [2] where the critical Josephson current oscillations due to transitions between 0 and  $\pi$  phase state were observed. A good agreement is obtained by taking into account the Umklapp processes.

I. Yu. A. Izyumov, Yu. N. Proshin, and M. G. Khusainov, Physics/Uspekhi 45, 109 (2002) 2. V.V. Ryazanov et al., Rev. Low Temp. Phys. 136, 385 (2004)

### SC27

#### Superconducting characters under pressure in heavy fermion compounds CeIr(In<sub>1-x</sub>Cd<sub>x</sub>)<sub>5</sub> studied by In-NQR Mitsuhami Yashima

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The heavy fermion compounds CeMIn<sub>5</sub> (M = Co, Rh, Ir) have provided to study the relationship between antiferromagnetism and superconductivity. CeCoIn5 becomes superconducting with  $T_c = 2.3 \text{ K}$  at ambient pressure1. This value is the highest record in the Ce-based heavy fermion superconductors. Since the existence of antiferromagnetic spin fluctuations is confirmed, it is expected that superconductivity is induced by antiferromagnetic spin fluctuations in CeCoIn5. CeRhIn5 is an antiferromagnet with \$T N\$ = 3.8 K at ambient pressure. Taking into account the lattice parameters of CeIrIn<sub>5</sub>, it is speculated that an antiferromagnetic order is realized at ambient pressure, but it actually shows superconductivity with  $T_c = 0.4 \text{ K}$ . The value of \$T c\$ is very low, as compared with the other CeMIn, compounds. \$T\_c\$ increases with increasing pressure and reaches about 1 K around 3 GPa, but antiferromagnetic spin fluctuations are greatly suppressed with the application of pressure. These results suggest that the superconducting mechanism of CeIrIn<sub>5</sub> may be quite different from that of the other CeMIn<sub>5</sub> compounds. In order to investigate the unconventional superconducting characteristics in CeIrIn5, the nuclear-quadrupoleresonance (NQR) measurements were performed under pressure in pure and Cd-doped CeIrIn.

#### **SD01**

# Giant and twofold oscillations of magnetoresistance in topological insulators Sb<sub>2</sub>Te<sub>3</sub> and Bi<sub>2</sub>Te<sub>3</sub> single crystals

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Topological insulators are quantum materials with an insulating bulk state and a topologically protected two dimensional metallic surface states. 1,2 Unique electronic properties, such as the quantum spin Hall effect, magnetoelectric effects, magnetic monopoles, and elusive Majorana states, are expected from topological insulators. 3,4 Topological insulators have great potential applications in spintronics and quantum information processing, as well as magnetoelectric devices with higher energy efficiency.5,6 To understand the macroscopic properties of the topological surface states and to investigate the possibility of their device applications, transport and magnetic measurements of high-quality single crystals are indispensable. Anisotropic magneto-transport properties were studied in p-type Sb<sub>2</sub>Te<sub>3</sub> and n-type Bi<sub>2</sub>Te<sub>3</sub> topological insulators through angular dependent magnetoresistance (MR) measurements. Giant MR of up to 230% was observed, which exhibits linear at high fields without any trend towards saturation. The giant MR displays a strong anisotropy, up to 210%, and twofold symmetry at different temperatures and fields. The giant MR might be due to intravalley and intervalley scattering, and the strong anisotropy might result from anisotropics of Fermi surface. The observed giant anisotropic MR of the Sb<sub>2</sub>Te<sub>3</sub> and Bi<sub>2</sub>Te<sub>3</sub> topological insulators could find applications in magneto-electronic devices based on topological insulators, such as magnetic sensors.

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#### **SD02**

#### Engineering and manipulating topological qubits in 1D quantum wires Panagiotis Kotetes<sup>1</sup>, Alexander Shnirman<sup>2</sup> and Gerd Schon<sup>1</sup>

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### SD03

# Angle dependence of the Landau level spectrum in twisted bilayer graphene

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In the context of the low-energy effective theory, the exact Landau level spectrum of quasiparticles in twisted bilayer graphene with small twist angle is analytically obtained by spheroidal eigenvalues. We analyze the dependence of the Landau levels on the twist angle to find the points where the twofold degeneracy for twist angles is lifted in the nonzero modes and below (above) which massive (massless) fermion pictures become valid. In the perpendicular magnetic field of 10 tesla, the degeneracy is removed at approximately 3 degrees for a few low levels, specifically approximately 2.56 degrees for the first pair of nonzero levels and approximately 3.50 degrees for the next pair. Massive quasiparticle appears at the angle less than 1.17 degrees in 10 tesla which matches perfectly with the recent experimental results. Since our analysis is applicable to the cases of arbitrary constant magnetic fields, we make predictions for the same experiment performed in arbitrary constant magnetic fields. For example, for 40 tesla we get the critical angle of 2.34 degrees and the sequence of angles 5.11,7.01,8.42,... for the pairs of nonzero energy levels. The symmetry restoration mechanism behind the massive (massless) transition is conjectured to be a tunneling (instanton) in momentum space.

#### **SD04**

# Surface band structure study of Bismuth-based ternary topological insulators

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Utilization of topological surface states is expected to lead to new vistas in electronics and fundamental physics. However, most of the known topological insulators either do not feature necessary band structure conditions (location of Dirac point with respect to the bulk band) or lack topological invariants essential for certain class of applications. Using angle-resolved photoemission spectroscopy (ARPES), we reveal the electronic band structure topology of a family of ternary spin-orbit insulators some of which feature functional electronic structure with in-gap Dirac point while others feature novel topological invariants (weak  $Z_2$  invariants) in crystalline form. These Bi-based ternary topological insulators provide a natural access to lattice matched magnetically ordered states to be utilized in future devices.

#### SD05

# Topological aspects and transport properties of edge states in the multi-band superconductor $\rm Sr_2RuO_4$

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Motivated by the spin-triplet chiral p-wave superconductor  $Sr_2RuO_4$ , we investigate the edge state of multi-band superconductor by means of the three band tight-binding model including inter-orbital hybridization and spin-orbit coupling effects. These bands correspond to the alpha, beta and gamma bands of  $Sr_2RuO_4$  where alpha and beta bands have hole- and electron-like characters, respectively. In particular we focus on the topological aspects and transport properties in the chiral spin-triplet phase. Although a full quasiparticle excitation gap in the bulk system, gapless edge states appear and affect both spontaneous spin and charge currents which give rise to anomalous and spin Hall effect. In this context, we study the interplay between electron- and hole-like particles and the effect of the genuinely two-dimensional gamma band in the formation of these edge states and edge currents. The topological aspects and correlation effects on this system are also discussed.

# SD06

#### Influence of geometry on the edge states of Bi Nanoribbons

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Two-dimensional (2D) topological insulators or 2D quantum spin Hall (QSH) systems provide an intriguing physical phenomena in condensed matter physics. The 2D QSH phase is a band insulator but has gapless edge states. The salient feature of these edge states is that they are topologically protected. Here, using first-principles densityfunctional calculations, we investigate the effect of geometry on the edge states in zigzag Bi nanoribbon (ZBNR) reported as a 2D QSH system.[1] We consider various edge geometries involving the ideal termination, relaxation, and rebonded reconstruction. We find that the rebonded reconstruction geometry is thermodynamically more stable than other geometries. Since the transition from the reconstruction to the relaxation geometry is kinetically feasible, the edge structure of ZBNR is expected to exhibit the orderdisorder transition as a function of temperature. We also find that the metallic feature of edge states is preserved with respect to the shear distortion and bond alternation, but it is dramatically altered to produce a band-gap opening through the rebonded reconstruction. We discuss such an influence of geometry on the edge states of ZBNR in terms of the penetration depth of the edge states.

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#### SD07

# Magnetotransport measurements in pulsed magnetic fields: a case for Fermiology studies in superconductors and topological insulators

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Superconductors and Dirac materials, such as graphene and topological insulators, are intensively studied in high magnetic fields. In this work, we present the current status of magnetotransport measurements at the Dresden High Magnetic Field Laboratory (HLD) and discuss examples of phase-diagram and fermiology investigations in different compounds [1,2], as well as the new Dirac material [3] SrMnBi<sub>2</sub>. Typical experiments in magnetic fields up to 85 T with 25 to 150 ms pulse duration in samples spaces up to 16 mm are performed by discharging a 50 MJ capacitor bank. We use standard four-probe-technique measurements with AC currents up to 200 kHz and the longitudinal and transverse (Hall effect) voltages signals from the samples are digitized with a fast recording system. We observed quantum oscillations in magnetotransport measurements up to 63 T on crystals of SrMnBi<sub>2</sub> and, together with other techniques, this demonstrates that there is a Dirac dispersion in the electronic structure of the double-sized Bi square net. However, in contrast to graphene, the Dirac cone in SrMnBi2 is highly anisotropic Dirac fermions. Aknowledgements Work at the HLD was supported by EuroMagNET II (No. 228043).

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#### SD08

# First-principles study of spin texture in the multilayer graphene on Ni(111)

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330 | The 19<sup>th</sup> International Conference on Magnetism

### SD09

A full quantum study on gapless modes and Axion electrodynamics in topological insulator heterostructure systems

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Three dimensional topological insulator heterostructure systems show the various novel phenomena such as gapless modes localized at defects or the Axion response in three dimensional topological insulator-magnetic insulator heterostructures. The origin of these phenomena is the nontrivial topology in both momentum and real spaces. Dealing with the spatial coordinates as the adiabatic parameter in the heterostructure systems, one can introduce the semiclassical Hamiltonian, which describes topological properties of the heterostructure system. On the other hand, in some cases, it also happens that semiclassical approaches fail to describe correctly topological features. In this study, we present full quantum treatment for three dimensional topological insulator heterostructure systems without the semiclassical approximation. We, first, prove the index theorem for gapless modes in line defects of the topological insulator-magnetic insulator heterostructure systems. Secondly, we demonstrate a full quantum calculation for the surface quantum anomalous Hall effect at the interface of the topological insulator-magnetic insulator heterostructure without using the Axion effective action. We also show that gapless modes in line defects and the Axion response may be governed not only by helical Dirac fermions on the surface, but also by gapped bound states at the interface, the contribution of which was missed in previous studies.

Xiao-Liang Qi, Taylor L. Hughes, and Shou-Cheng Zhang, Phys. Rev. B 78, 195424 (2008). Jeffrey C. Y. Teo and C. L. Kane, Phys. Rev. B 82, 115120 (2010).

#### **SD10**

### Magnetic properties of rare earth doped Bi<sub>2</sub>Te<sub>3</sub>

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Bi<sub>2</sub>Te<sub>3</sub> is one of 3-dimensional topological insulators together with Bi<sub>2</sub>Se<sub>3</sub> and Bi<sub>2</sub>Sb<sub>3</sub>. [1] Recently, Binghai Yan et. al have theoretically reported that LaBiTe<sub>3</sub> is a strong topological insulator. [2] We could conjecture that rare-earth materials have fascinating magnetic properties, since it has filled outer shell (5x25p6) and simulataneously unfiled inner shell (4f n). [3] Therefore, we made both LaxBi<sub>2x</sub>Te<sub>3</sub> and CexBi<sub>2x</sub>Te<sub>3</sub>, and measured transport and magnetic properties. LaxBi<sub>2x</sub>Te<sub>3</sub> is diamagnetic, while CexBi<sub>2x</sub>Te<sub>3</sub> is paramagnetic at low temperature. From the modified curie-wiess fit, we could obtain the effective magnetic moment of  $2.56\mu$ B. This value demonstrates that the Ce ion is in trivalent state. However, the carrier density of CexBi<sub>2x</sub>Te<sub>3</sub> is observed to be on order smaller than that of Bi<sub>2</sub>Te<sub>3</sub>, giving rise to a strong increase of carrier mobility. These results support a scenario of strong topological insulator. Furthermore, a magnetoresistance ratio of CexBi<sub>2x</sub>Te<sub>3</sub> is 1600% at 70kOe, even though ferromagnetic FexBi<sub>2x</sub>Te<sub>4</sub> have 600% magnetoresistance ratio at 70kOe.

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#### SD11

#### Iron doping effect in topological insulator: Bi2Te3

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 $Bi_2Te_3$  is well-known topological insulator with anti-site defects which complicates the manipulation of the Fermi level. It has also been reported that surface band of topological insulators is open when magnetic ions are doped, because magnetic ions are able to break the time reversal symmetry. In this report, we doped magnetic ions of Fe into  $Bi_2Te_3$  and studied the properties and the corresponding electronic structure. We found that a small amount of Fe ions less than 0.4% do not open the surface gap, but reduce the anti-site defects. Thus, Fe-doped  $Bi_2Te_3$  has linear magnetoresistance with high mobility (15,074cm2v-1s-1) which is three times higher than  $Bi_2Te_3$  while  $Bi_2Te_3$  has a deviation from the linear dependence of magnetoresistance. The reduction of anti-site defects enhances the sample quality, so that we could observe quantum oscillations such as Shubnikov de Haas and de Haas van Alphen effect. In addition, the spin-dependent ARPES data also demonstrate both clear surface state without gap opening and canted spin components in the c axis.

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### **SD12**

#### Magneto transport properties of topological insulator nanoribbons of Bi<sub>2</sub>Te<sub>3</sub>

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We report on the magneto transport properties of topological insulator nanoribbons of Bi2Te3. The Bi2Te3 nanoribbons were grown using the electrodeposition processes and contacted with 100-nm-thick cobalt electrodes. At low temperature below T = 10 K, the magnetoreistance increases with a magnetic field B, which is attributed to a weak antilocalization effect. When the magnetic field (B) is applied parallel to the ferromagnetic electrodes on top of the nanoribbon, a switching behavior of the magnetoresistance is observed near B = 400 Oe. Temperature and gate voltage dependences of the spin valve effect is discussed as well.

### **SD13**

Structural investigations of the topological insulators Bi<sub>2</sub>Se<sub>3-x</sub>Te<sub>x</sub> Geetha Balakrishnan<sup>1</sup>\*, Ravi Singh<sup>1</sup>, Devashibhai Adroja<sup>2</sup>, Kevin Knight<sup>2</sup> and

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The binary phases Bi2Se3 and Bi2Te3 have been discovered to exhibit three dimensional topological insulating behaviour. However, these compounds are not insulators when investigated in the bulk. Bi2Se3 is generally known to exhibit n-type conductivity and Bi2Te3 exhibits n or p-type conductivity depending on the synthesis conditions. The bulk conductivity exhibited by these materials are seen as being detrimental to the investigation of their surface conducting properties. We have examined the intermediate compositions Bi2Se3-xTex in order to study the role played by vacancies and anti-site defects in determining the bulk conduction mechanisms. We report detailed structural refinements obtained by neutron diffraction experiments on several of the intermediary phases and draw correlations between the defect structures and their conducting / semiconducting / insulating behaviour

### **SD14**

### Josephson effects in Bi<sub>2</sub>Se<sub>3</sub> topological insulator nanoribbons

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Topological insulators are exotic materials with bulk band gap and metallic edge states which are protected on their own boundary topologically. Here, we report on the fabrication and measurement results of the superconducting proximity junctions of topological insulator nanoribbons of Bi<sub>2</sub>Se<sub>2</sub>. Single-crystalline Bi<sub>2</sub>Se<sub>2</sub> nanoribbons are synthesized using the vapor-liquid-solid method, while the superconducting Al electrodes are formed on top of the nanowire. When a magnetic field (H) is applied along the axial direction the magneto-resistance data exhibit quasi-periodic oscillations with an average periodicity of  $H^* \sim 0.4$  T, which is consistent with the Aharonov-Bohm oscillations. In the superconducting state, the supercurrent branch with a critical current of Ic ~ 90 nA is clearly observed in the current-voltage curve as a result of the superconducting proximity effect in Bi-Se, nanoribbon. Quantized voltage steps of the Bi<sub>2</sub>Se<sub>2</sub> nanoribbon Josephson junction under the microwave irradiation satisfy the ac Josephson relation.

### SD15

#### Electronic structure and transport properties of pt based heusler compounds with C1b structure for topological quantum phenomena

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Many of the Heusler commounds based on heavy elements with 1:1:1 commosition and C1b structure are zero. band gap insulators and exhibit partially an inverted band structure [1]. Based on a topologically protected electronic surface states, this class of materials is supposed to open up innovative directions for future technological applications in spintronics quantum computing and thermoelectrics. The density of states of Pt based Heusler compounds PtMZ (M = Y, Gd, Lu and Z= Sb, Bi) were investigated by bulk sensitive hard X-ray photoelectron spectroscopy. The measured valence band spectra are clearly resolved and in well agreement with the first-principles calculations of the electronic structure of the compounds. Close to the Fermi energy, the linear behavior of the measured spectra at high excitation energy where influences of the surface can be neglected improve bulk origin of the Dirac-cone type density [2] Furthermore, the temperature dependence of electrical conductivity, Hall mobility, Seebeck coefficient and thermal conductivity were investigated. The promising properties of those compounds open a high potential for new technological applications.

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#### **SD16**

#### Gapless interface states in topological insulator/semiconductor heterostructures

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Based on first principles calculations we investigate the existence of topologicallynontrivial states at the interface between a topological insulator (TI) and a semiconductor (S) in TI/S heterostructures. We find that the emergence of gapless states depends on the symmetry of the junction, being protected by mirror symmetry. We discuss their characteristic dispersion and spatial distribution and analyze the quantum interference between helical states and its influence on the spin texture of the quasi two-dimensional states as a function of the semiconductor thicknesses. From the firstprinciples calculations, we derive an effective Hamiltonian, which captures the salient topological features. The effective model combined with a Green function matching method [1] allow us to investigate the evolution of the interaction between interface helical Dirac fermions from finite heterostructures to actual seminfinite interfaces and therefore to analyze the novel behavior of the spin-charge coupling in two-dimensional layered systems with strong spin-orbit coupling.

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#### **SD17**

#### Thermoelectric transport in topological insulators

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We study the thermoelectric properties of three-dimensional topological insulators with many holes (or pores) in the bulk. We show that at high density of these holes the thermoelectric figure of merit, ZT, can be large due to the contribution of the conducting surfaces and the suppressed phonon thermal conductivity. The maximum efficiency can be tuned by an induced gap in the surface states dispersion through tunneling, magnetic impurities, or external magnetic fields. The large values of ZT, much higher than unity for reasonable parameters make this system a strong candidate for applications in heat management of nanodevices, especially at low temperatures.

O. A. Tretiakov, Ar. Abanov, and Jairo Sinova, Applied Physics Letters, 99, 113110 (2011)

# SD18

#### Topological phase in a one-dimensional interacting fermion system Huaiming Guo

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We study a one-dimensional (1D) interacting topological model by means of the exact diagonalization method. The topological properties are first examined with the existence of the edge states at half-filling. We find that the topological phases are not only robust to small repulsive interactions, but also are stabilized by small attractive interactions, and also finite repulsive interaction can drive a topological nontrivial phase into a trivial one while the attractive interaction can drive a trivial phase into a nontrivial one. Next we calculate the Berry phase and parity of the bulk system and find that they are equivalent in characterizing the topological phases. With them we obtain the critical interaction strengths and construct part of the phase diagram in the parameters' space. Finally we discuss the effective Hamiltonian at the large-U limit and provide an additional understanding of the numerical results. These results could be realized experimentally using cold atoms trapped in the 1D optical lattice.

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#### **SD19**

#### Robustness of 1D topological superconductors with Majorana edge Evidence of a spin gap above the magnetic ordering temperature and states against lattice modulation crystal field excitations in CeOs<sub>2</sub>Al<sub>10</sub>

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#### SE01

#### Temperature-pressure phase diagram of quadrupolar order in PrTr<sub>2</sub>Al<sub>20</sub> (Tr=Ti,V)

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PrTr<sub>2</sub>Al<sub>20</sub> (Tr=Ti,V) crystallizes in a cubic CeCr<sub>2</sub>Al<sub>20</sub> type structure with the space group Fd<sub>3</sub> ?m. According to neutron experiment, crystal electric field (CEF) ground state of Pr ion is the non-Kramers  $\Gamma_3$  doublet [1]. Reflecting the non-magnetic ground state,  $PrTi_2Al_{20}$  and  $PrV_2Al_{20}$ exhibit quadrupolar ordering at TQ ~ 2.0 K and 0.6 K, respectively. It is notable that the low temperature electrical resistivity of PrV2Al20 shows VT dependence, suggesting the quadrupolar Kondo effect[2]. In the present study, we have measured electrical resistivity under high pressure on PrTr2Al20 (Tr=Ti,V) single crystals in order to clarity the pressure dependence of TQ and Kondo effect. Electrical resistivity of PrTi2Al20 shows the broad maximum centered at ~50 K due to incoherent Kondo-scattering processes on the ground state and the excited CEF level. With increasing pressure, the magnitude of this anomaly increases while the peak position is unchanged. At lower temperatures, it is found that a clear resistive drop at TQ  $\sim$  2 K at ambient pressure. TO first increases and shows the maximum under pressure. We will present the phase diagram of PrTi<sub>2</sub>Al<sub>20</sub> together with that of PrV<sub>2</sub>Al<sub>20</sub>.

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# SE02

### Soft point contact spectroscopy in the antiferromagnet Ce<sub>2</sub>RhIn<sub>8</sub>

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We performed soft point contact spectroscopy on the Ce-based heavy fermion systems (CeIn<sub>3</sub>, CeCoIn<sub>5</sub>, CeRhIn<sub>5</sub>, CeIrIn<sub>5</sub>, and Ce<sub>2</sub>RhIn<sub>8</sub>) with silver epoxy replacing a conventional tip. Unlike CeIn3 where the conductance dI/dV is enhanced below TN (=10K), the conductance of CeRhIn, at zero bias voltage is rapidly enhanced below 8K, which is higher than TN of CeRhIn<sub>5</sub> (=3.8K). Heavy fermion compounds with a magnetic order often show intermediate states that are characterized by a broad enhancement of conductance at temperatures higher than TN, while heavy fermion compounds without magnetic ordering reveal asymmetric conductance line-shape below characteristic temperature T\*. The antiferromagnet Ce2RhIn8 is unique in that the intermediate state shows a V-shaped dip around zero biased voltage and is enhanced with decreasing temperature below 25K, which is much higher than T\*=5K. Below TN=2.8K, a symmetric peak at zero bias voltage in the conductance spectra is enhanced within the V-shaped dip with decreasing temperature. We interpret the symmetric peak and V-shaped dip around zero bias as manifestation of the antiferromagnetic (AFM) order and Kondo dip in the Kondo lattice system, respectively, indicating that Ce2RhIn8 is at a boundary between AFM and Kondo states.

#### SE03

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Recently, the opening of a spin gap has been reported in the intermetallic compounds  $CeT^2AI_{10}$  (T = Ru and Os) below TN = 27-29K through heat capacity and magnetic susceptibility measurements. Various theoretical models have been proposed to explain the spin gap formation, for example a spin-Peierls and RVB models of Hanzawa for CeRu<sub>2</sub>Al<sub>10</sub>. To understand the nature of the spin gap formation in CeOs<sub>2</sub>Al<sub>10</sub>, we have carried out low and high energies inelastic neutron scattering (INS) measurements between 4.5 and 65K. Our INS studies reveal a clear sign of a spin gap of 11meV at 4.5K and of 4meV at 35K, while above 45K the observed response transformed into a quasielastic line with linewdith of ~ 8 meV. The observation of a spin gap above TN suggests that its origin is not only due to the gapped spin wave, but also related to the hybridization between 4f- and conduction electrons. Moreover the high energy INS study reveals two crystal field (CEF) excitations at 37 and 57meV. We have carried out an analysis of the INS data along with the single crystal susceptibility based on the CEF model, which gives good fit to both data sets.

### **SE04**

#### Role of quantum fluctuations in forming heavy-fermions for Ca<sub>2-</sub> <sub>x</sub>SrxRuO<sub>4</sub>

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In order to understand a role of quantum fluctuations in forming heavy-fermions for Ca2-xSrxRuO4, we investigated the effect of quantum fluctuations on the massenhancement on the basis of the three-orbital Hubbard model for the Ru-t2g-orbitals. In this presentation, I show the momentum-dependence of both the spin and charge susceptibilities for the tight-binding models at x = 2.0 and 0.5 within the random phase approximation (RPA). I also present the variation of  $\gamma$ e, obtained by using the Hellmann-Feynman theorem and the RPA, between x=2.0 and 0.5 and address a role of both the spin and orbital fluctuations in enhancing ye for Ca2-xSrxRuO4

### SE05

Study of long range magnetic ordering and spin gap formation in Ce(Ru1-xFex)2Al10 through muSR and neutron scattering measurements

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We have carried out muon spin relaxation (muSR) and inelastic neutron scattering (INS) investigations on Ce(Ru1-xFex)2Al10 (x=0, 0.3, 0.5, 0.8 and 1) compounds to investigate the origin of the spin gap formation. Our muSR study confirms the long-range magnetic ordering below 26K in x=0, 0.3 and 0.5, while x=0.8 and 1.0 remains paramagnetic down to 1.2K. Furthermore, INS measurements of x=0 clearly reveal the presence of a spin gap of 8meV at 5K. Interestingly at 5K the spin gap excitation broadens in x=0.3 and exhibits two clear INS peaks at 7.6(0.3) and 11.1(2.0)meV in x=0.5. Moreover the x=0.8 sample, which remains paramagnetic down to 1.2K, reveals a clear sign of a spin gap of 10-12 meV, with strong Q-dependent intensity that follows a spin-dimer structure factor. The observation of a spin gap in the paramagnetic sample (x=0.8) opens a question, what is the origin of the spin gap in CeT2Al10 (T=Ru and Os) compounds? Possibilities include gapped spin wave, hybridization gap, spin-dimer or a gap opening on a small part of the Fermi surface. Our analysis of the Q-dependent intensity reveals that the mechanism of the spin gap in the x=0.8 sample is possibly due to spin-dimer formation.

#### SE06

#### Specific heat of structure-disordered heavy-fermion CexY<sub>80</sub>-Mn<sub>20</sub> allovs

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Recently, we have reported physical properties of structure-disordered (a-)CexY<sub>80</sub>. <sub>x</sub>Mn<sub>20</sub> alloys which show a heavy fermion state in the non-Bloch system without translation symmetry[1]. The specific heat Cp for present alloys follows  $\gamma T + \beta T3$ down to about 10 K. The T-linear coefficient y shows a large value (~200 mJ/CemolK<sub>2</sub>) and almost independent of the Ce-concentration. However, the Cp/T vs. T2 plot of T < 10 K deviates from the T-linear relation with decreasing temperature. In this work, in order to clarify the factor of the deviation of Cp/T for a-CexY<sub>80-x</sub>Mn<sub>20</sub> alloys, we have measured the low-temperature specific heat. We have estimated the temperature TA of the deviation from the Cp/T vs. T<sub>2</sub> plot for present alloys. The TA is about 10 K, and almost independent of the Ce-concentration. Therefore, the deviation from the Cp/T for the present alloys would be considered the crystalline-electric-field splitting of the six-fold degenerate level of the Ce3+ ion. Detailed discussions will be presented at the conference

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#### SE07

#### Synchrotron X-ray diffraction study on crystal structure of URu<sub>2</sub>Si<sub>2</sub> Chihiro Tabata<sup>1</sup>\*, Reiji Kumai<sup>2</sup>, Kensuke Kobavashi<sup>2</sup>, Hironori Nakao<sup>2</sup>, Yoichi

Murakami<sup>2</sup>, Makoto Yokoyama<sup>3</sup>, Hiroyuki Hidaka<sup>1</sup>, Tatsuya Yanagisawa<sup>1</sup> and Hiroshi Amitsuka<sup>1</sup>

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URu<sub>2</sub>Si<sub>2</sub> (the tetragonal ThCr<sub>2</sub>Si<sub>2</sub> type crystal structure; space group I4/mmm) has attracted much interest because of its peculiar phase transition at 17.5 K (≡To) [1-3]. Despite more than 25 years of intense research trying to identify the order parameter, the intrinsic nature of the phase transition has not been elucidated [4]. The ThCr<sub>2</sub>Si<sub>2</sub> structure includes a free position parameter, z, of the Si 4(e) site. Since the Si ions are located at the second nearest-neighbor position of the U ions, the precise determination of the z parameter might provide useful information for evaluation of hybridization effects of 5f orbits with Si s. p states, calculations of a band structure. analyses for 29Si-NMR spectra, and so on. These pieces of knowledge might offer a clue to solving the issue of hidden order in URu2Si2. We will report our latest study of structural analyses of a single crystal of this material based on the X-ray diffraction measurements performed at BL-8A and B in the Photon Factory KEK

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#### **SE08**

#### Spin-density wave order in the 2D heavy fermion system CePt<sub>2</sub>In-

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The title compound is a recently discovered heavy fermion material where the spacing between Ce-In planes is drastically increased. Consequently, CePt<sub>2</sub>In, inherit a truly 2D electronic structure. It was recently discovered that CePt In, is not only AF (TN=5.3 K) at ambient pressure but also becomes SC under pressure with a maximum transition temperature Tc=2.1 K at P=3.12 GPa. At lower pressures an intriguing coexistence of the AF order and SC phase is found, and that with increasing pressure Tc is increasing while TN decreases. The growth of the SC on the expense of the AF order suggests a crossover behavior of the Ce-4f electrons from localized to itinerant, similar to what is considered for the well-known CeRhIn, compound In the presented work the low-temperature microscopic magnetic properties of the quasi-2D heavy fermion compound, CePt<sub>2</sub>In<sub>7</sub> are investigated by using µSR. Clear evidence for the formation of a SDW order is presented. The magnetic order parameter fit well to a modified BSC gap-energy function in a strong-coupling scenario, possibly predicting the evolution of unconventional pairing in the pressure induced superconducting phase of this compound. Finally, we also present recent pressure dependent µSR data.

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#### **SE09**

#### Metal-nonmetal transition in Cr partial substituted Ni0.96S Masanori Matoba\*, Yoichi Kamihara and Shuichiro Anzai

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Hexagonal NiAs-type NiS exhibits a first-order metal-nonmetal transition (MNT) at Tt=263K, resulting from the opening of narrow charge-transfer (CT) gap between the S 3p band and the upper Hubbard band [1,2]. The low temperature phase is the nonmetallic antiferromagnetic state, while the high temperature phase shows metallic and Pauli-paramagnetic behavior [3,4]. The substitution of vacancies for Ni atoms in NiS lowers Tt to be suppressed to be zero at d=0.035 in Ni1-4S, so that metallic Ni096S do not exhibit the MNT [5,6]. In this study, we synthesized Cr-partial substituted Ni<sub>0.96</sub>S or (Ni<sub>1-x</sub>Crx)<sub>0.96</sub>S by a solid state reaction in an evacuated silica tube to reveal the Cr-partial substitution effect on the physical properties of metallic Nings by the measurements of X-ray diffraction, electrical resistivity, magnetic susceptibility, and thermopower. Surprisingly, (Ni1-xCrx)096S clearly exhibits the MNT, while metallic Ni<sub>1005</sub>S do not exhibit the MNT. The Tt increases with x sharply from x=0.025 (Tt=50 K) to x=0.09 (Tt=260 K) by Cr-partial substitution for Ni in Ni {0.96} S At the conference. electronic nature of Cr-partial substituted Ni096S or (Ni12Crx)096S with ability for the opening of CT-gap will be discussed from the viewpoint of the strongly correlated electron systems.

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#### SE10

#### A study of ni-substitution and pressure effects on the heavy-fermion Superconductor CeCu<sub>2</sub>Si<sub>2</sub>

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The heavy-fermion superconductor CeCu2Si2 shows an attractive enhancement of T\_c around a critical pressure (~ 4.5 GPa), at which the T-linear variation of resistivity, the enhancement of the residual resistivity, the rapid decrease of T2-coefficient and a sign of a valence crossover have been observed [1-3]. These characteristics have been believed to be signatures of the critical valence fluctuations associated with a valence transition around its critical-end point [4, 5], while we have been requiring more signatures for clarification of the fascinating phenomena. Effects of the Ni substitution on CeCu2Si2 are investigated by means of specific heat and electrical resistivity measurements. A characteristic T-linear dependence of the resistivity is observed around  $x \sim 0.1$  with a fractional resistivity drop due to filamentary superconductivity. In addition, a variation of the power index n of a fit to  $\rho$ - $\rho_0 = aT^n$  at low temperatures against T\_1<sup>max</sup> at which the resistivity shows a maximum is almost fully identical to the result of the high pressure experiment [1]. We would like to discuss for the similarities between Ni doping and pressure and the origin of the anomalous Fermi liquid state.

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# SE11

#### Electrical resistivity measurement under pressure in the heavy fermion antiferromagnetic compound Ce<sub>2</sub>PtGa<sub>12</sub>

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Ce-based heavy-fermion ternary compounds have been of interest because of their exotic physical properties under pressure, such as magnetic ordering and superconductivity.[1] In these systems, there is a competition between magnetic order of f-electrons (RKKY effect) and the screening of magnetic moments (Kondo effect) producing large effective electron masses. Ce<sub>2</sub>PtGa<sub>12</sub> orders antiferromagnetically at 7 K, and specific heat measurements suggest that it is a moderate heavy-fermion system with  $\gamma = 191 \text{ mJ/mol K}_2$ . For understand of competition between RKKY and Kondo effect. We have measured the electrical resistivity under pressure on Ce-PtGa<sub>12</sub> single crystal. In this paper, we will report the suppression of the antiferromagnetic order under higher pressure and discuss the electronic state in the vicinity of the critical pressure.

[1] R. T. Macaluso et al., J. Solid State Chem. 178 (2005) 3547.

### SE12

#### Anisotropy of URhGe

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The heavy fermion material URhGe has garnered interest in recent years as an unconventional superconductor. It shows coexistence of ferromagnetism and p-wave superconductivity[1], and superconductivity which is re-entrant in magnetic field[2]. We now believe the latter effect to be the result of proximity to a topological change in the Fermi-surface[3]. In this poster I report recent resistivity measurements or samples of URhGe with well-defined current direction. Most previous works do not consider this when reporting resistivity, however we find that some features can vary significantly as a function of current direction. We discuss how these findings constrain the anisotropy of the Fermi surface, and the possibility of a field dependent and/or anisotropic scattering rate. I also present data taken on high quality samples close to the quantum critical point. Interestingly, there is no enhancement of the superconducting Tc near the QCP compared to that observed with the field applied entirely along the crystalline b-axis, where a low temperature tricritical point is present.

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### **SE13**

#### Study of spin wave and spin gap in single crystals of CeRu<sub>2</sub>Al<sub>10</sub> using inelastic neutron scattering measurements

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We have carried out inelastic neutron scattering (INS) investigations on single crystals of CeRu<sub>2</sub>Al<sub>10</sub> using the high neutron flux spectrometer MERLIN at ISIS facility. The crystals (total mass of 2.4g) were mounted in the b-c scattering plane. We carried out measurements using neutron incident energy Ei=20 meV at 4.5 K for b // ki and c // ki. We have seen a clear dispersive spin wave type excitations coming out from the [1 0 0] and [0 2 0] antiferromagnetic zone centres with a single ion anisotropic gap of -4.5 meV. On the other hand along [0 0 L] direction the spin excitations exhibit weak dispersion with almost Q-independent gap of 8 meV. The intensity of the peak along [0 0 L] is maximum at L = 0 and decreases with increasing L value, like Ce<sup>3+</sup> magnetic form factor. These results may suggest that excitations in the ab-plane can be explained by the gaped spin wave, but along [0 0 L] is dominated by the hybridization effect. We will present a detail analysis of the observed excitations

# SE14

#### Universal behavior in the nonlinear magnetic response of strongly correlated metal

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A diverse class of strongly correlated electronic materials irrespective of whether they order magnetically or not, often exhibit a peak in the temperature dependence of the linear magnetic susceptibility,  $\gamma^1$ . In non-linear susceptibility measurements on a number of heavy electron materials including in single crystals of the prototypical strongly correlated metal. UPt, we find that there exists a peak in the third order susceptibility,  $\gamma^3$  as well. This peak in  $\gamma^3$  occurs at a temperature T<sub>2</sub>, which is roughly half the temperature  $T_1$ , at which the peak in  $\chi^1$  appears. This proportionality between T<sub>2</sub> and T<sub>2</sub> also reported previously in other heavy fermion systems implies a universal behavior in the nonlinear magnetic response of correlated metals. The observed proportionality between  $T_2$  and  $T_1$  can is explained by a model which assumes a singlet ground state for the spins present in a crystalline electric field with a concomitant doublet excited state, separated from the ground state by an energy  $\Delta$ . This model also vields metamagnetic behavior at a magnetic field corresponding to  $\Delta$ .

### SE15

#### X-ray absorption studies of the Ce<sub>2</sub>Rh<sub>1-x</sub>Ir<sub>x</sub>In<sub>8</sub> intermetallic compounds

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Within the series of heavy fermions Ce2Rh1xIrxIn8 compounds it has been recently observed the occurrence of two low-temperature superconducting (SC) phases as a function of temperature both at ambient pressure and under applied pressure for intermediate Ir-content (x~0.25-0.7) with remarkable different behavior of both states as a function of pressure and x [1]. The details of the observed results seem to be reminiscent of the two SC phases found in the CeRh, Jr.In, system, thus related to the higher dimensionality and structural disorder for the Ce2-1-8 systems. Here, we report the results of extended x-ray absorption fine structure measurements performed on  $Ce_2Rh_{1-x}Ir_xIn_8$  (x = 0.0, 0.25, 0.75 and 1.0) at T = 10, 150 and 300 K at the Rh K-edge and at the Ce and Ir-L<sub>3</sub> absorption edges in order to check the extent of the structural disorder at the Rh-Ir atomic sites and their relevance for the SC phases observed in the resistivity and specific heat measurements in these compounds. Our results point to the trivalent state of the Ce ions, and that there is no evidence of phase separation in the lowest measured temperature.

[1] E. N. Hering, H. A. Borges et al. Phys. Rev. B 82 184517 (2010)

#### SE16

#### Polarized neutron diffraction study on the magnetic ordering in UMn<sub>2</sub>Al<sub>20</sub>

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The ternary intermetallic compounds UX2Al20 (X = Ti, V, Nb, Ta, and Cr), crystallizing with the cubic CeCr<sub>2</sub>Al<sub>20</sub> type crystal structure, were characterized in the literature as weak Pauli paramagnets [1-3]. In striking contrast, the isostructural phase UMn2Al20 was reported to order ferromagnetically due to the presence of magnetic moments on the uranium atom sites [4]. In this contribution, we present the results of our polarized neutron diffraction study on a high-quality single crystal of UMn2Al20, which was also carefully characterized by means of X-ray diffraction, microprobe, magnetic, electrical transport and heat capacity measurements. The bulk properties were found very similar to those published in Ref. 4. In particular, the ferromagnetic transition at TC = 17 K was clearly revealed from the magnetic data, while the temperature-dependent specific heat and electrical resistivity were found nearly featureless. The magnetization density maps derived from the neutron diffraction data, showed that the ferromagnetism in UMn<sub>2</sub>Al<sub>20</sub> arises due to the magnetic moments carried on the manganese atoms, while the uranium sublattice remains nonmagnetic. This result appears perfectly consistent with the behavior of the UX2Al20 (X = Ti, V, Nb, Ta, and Cr) compounds, in which the U atoms do not exhibit any magnetic moments

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### SE17

### Antiferromagnetic ordering in single-crystalline Ce2IrSi3

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The novel cerium-based ternary silicide Ce<sub>2</sub>IrSi<sub>3</sub>, crystallizing with fully ordered derivative of the hexagonal AlB<sub>2</sub>-type unit cell, has recently been characterized as a Kondo lattice that orders antiferromagnetically at TN = 1.3 K [1]. An antiferromagnetic character of the ordering was inferred merely from the specific heat and electrical resistivity data, collected in applied magnetic fields. With the aim at verification of the presumed nature of the ground state, we have performed magnetic studies down to 0.46 K and in external magnetic fields up to 7 T using a SQUID magnetometer. The results unambiguously corroborated the antiferromagnetic phase transition at 1.3 K, and additionally revealed pronounced magnetic anisotropy in the ordered state. The magnetization component measured within the hexagonal ab plane was found much larger than that taken along the c axis. This finding indicates that the magnetic mements in Ce<sub>2</sub>IrSi<sub>3</sub> are confined to the basal plane. In a magnetic-like phase transition. In turn, upon applying magnetic field along the hexagonal c axis some complex, yet more subtle changes in the magnetic structure seem to occur.

[1] M. Szlawska and D. Kaczorowski, Phys. Rev. B 84 (2011) 094430

#### **SE18**

Berezinskii-kosterlitz-thouless transition in heavy fermion superlattices Jian-huang She\* and Alexander V Balatsky

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Recently unconventional superconductivity has been observed in epitaxially grown heavy fermion superlattices, where Cooper pairs are confined within two-dimensional Kondo lattices (Nature Physics 7, 849(2011)). Here we propose to understand these experiments within the framework of Berezinskii-Kosterlitz-Thouless Transition. The observed behavior of the superconducting transition temperature can be understood in terms of phase fluctuations controlled by vortex-antivortex (un)binding. Reasonable agreement is found with the experiments. Furthermore, we predict a large renormalization of the penetration depth, which can be measured by future experiments. Evolution of the gap is explained by invoking effects of the interface.

**SE19** 

# Basic properties of the intermetallics APd<sub>5</sub>Al<sub>2</sub> (An=Ce, Th, U, Np, Pu, Am)

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### SE20

# Study of vibron quasibound state in CeAg<sub>1-x</sub>CuxAl<sub>3</sub>, 0<x<1

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The Ce-systems have attracted considerable interest due to the duality between the itinerant and the localized nature of 4f-electrons and the presence of strong electron-electron correlations. On the other hand, it is well known that the electron-phonon interaction drives the BCS-superconductivity. The nature of electron-phonon interaction in Ce-based compounds has its origin in an effective strong magnetoelastic coupling between some phonon modes and the Ce-localized states arising from crystal-field influence at some reciprocal space regions with high DOS, as in CeAl<sub>2</sub> [1], and recently CeCuAl<sub>3</sub> [2]. We have studied of crystal-field electron-phonon excitations along the series CeAgxCu<sub>1-x</sub>Al<sub>3</sub> (0<x<1), where the magnetoelastic coupling changes and disappears at critical concentration ? 0.65 under the hybridization of Cu and Ag d-bands with 4f-electrons. A screening model of dielectric permittivity was developed to expling this behaviour from a 3D-antiferromagnet in CeCuAl<sub>3</sub> to a in-basal-plane 2D-weakly ferromagnet for CeAgAl<sub>3</sub> [3]; both with a Kondo temperature inferior to 10 K

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#### SE21

### Magnetic properties of cubic GdTi<sub>2</sub>Al<sub>20</sub> single crystal

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The compounds crystallizing the RCr<sub>2</sub>Al<sub>20</sub>-type cubic crystal structure having 184 atoms in the unit cell with 8 formula units per unit cell, have attracted a great deal of interest as they exhibit a variety of interesting magnetic properties. For example, PrTi<sub>2</sub>Al<sub>20</sub> is the first example of a non-magnetic doublet ground state system, which exhibits a ln T behaviour in the electrical resistivity, indicating the Kondo effect [1]. Similarly, SmTi<sub>2</sub>Al<sub>20</sub> also exhibits Kondo effect [2]. Here we report on the single crystal growth and the transport and magnetic properties of an iso-structural compound: GdTi<sub>2</sub>Al<sub>20</sub>. Single crystals of GdTi<sub>2</sub>Al<sub>20</sub> were grown by high temperature solution growth using molten aluminum as flux. Large size single crystals with well-defined triangular shaped facets corresponding to (111) plane were obtained. The stoichiometry and the phase purity of the single crystal were confirmed by means of energy dispersive analysis by x-ray diffraction analysis. The estimated lattice constant from the x-ray diffraction was 14.682 A. It was found that GdTi<sub>2</sub>Al<sub>20</sub> undergoes an antiferromagnetic ordering at TN = 2.6 K. We will present our results on the magnetic and transport properties of GdTi<sub>2</sub>Al<sub>20</sub>.

[1] A. Sakai and S. Nakatsuji, J. Phys. Soc. Jpn. 80 (2011) 063701. [2] R. Higashinaka et al, J. Phys. Soc. Jpn. 80 (2011) 093703.

#### **SE22**

# Investigation of the heavy fermion Ce<sub>3</sub>Ir<sub>4</sub>Sn<sub>13</sub> by electrical resistivity under pressure

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The Ce<sub>3</sub>Ir<sub>4</sub>Sn<sub>13</sub> is a heavy fermion compound with an electronic specific heat  $\gamma = 670mJ/mol-CeK_2$  [1,2] for which two anomalies, at 0.6K and 2.1K, can be found on specific heat. Magnetic susceptibility data indicate that the anomaly at 0.6K is due to an antiferromagnetic ordering, while the transition at 2K is non-magnetic. The later one could be related to a change of the band structure accompanied by a lattice expansion[3] as was verified by XRD. This compound crystallizes in a cubic Yb<sub>3</sub>Rh<sub>2</sub>Sn<sub>3</sub> perovskite-like arrangement, with 40 atoms per unit cell [4]. In this work, we are studying the electrical resistivity (100mK < T < 300K and P<25kbar) performed on single crystals of Ce31r4Sn13, grown using a Sn self-flux technique. Our main motivation is to clarify the origin of the phase transition that occurs at 2K and to investigate the possibility of a vanishing of the Neel temperature over a quantum critical point. Our results show a Kondo-like increase of the resistivity at higher temperatures and a very sharp peak at 2K that moves to lower temperatures with applied pressure. Magnetic susceptibility measurements will be performed to investigate the behavior of the Neel temperature and pressure.

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# SE23

# Electronic structure studies of UPt<sub>3</sub> using soft x-ray angle-resolved photoemission spectroscopy and band calculation

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Electronic structure and Fermi surface of heavy fermion superconductor UPt3 have been studied by de Haas-van Alphen experiments and various theories [1], and it is suggested that USf itinerant band model is more comparable with the experiments. In order to observe the valence band structure of UPt3 directly, angle-angle photoemission spectroscopy had been done using a discharge lamp [2]. However a 3-dimensional band structure of UPt3 has been not understood well. We have carried out soft x-ray angle-resolved photoemission spectroscopy experiment with the twin-helical undulator beamline BL23SU of SPring-8, and as a result the USf itinerant bands of UPt3 can be detected together with the Pt 5d bands. In the paper we will discuss the band structure of UPt3 with a LDA band calculation.

[1] G. J. McMullan et al., New J. Phys. 10, 053029 (2008). [2] T. Ito et al., Phys. Rev. B 59, 8923 (1999).

#### SF01

# Field-dependent instability of the candidate quantum spin liquid in EtMe<sub>3</sub>Sb[Pd(dmit)<sub>2</sub>]<sub>2</sub> as revealed by NMR.

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#### SF02

# Magnetic states in quasi-2-D iridium oxides with large spin-orbit coupling

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The discovery of the novel Mott insulating state in Sr<sub>2</sub>IrO<sub>4</sub> has developed a new research field on solid-state physics [1]. The strong spin-orbit coupling with the moderate on-site Coulomb interaction in the 5d electron system yields the unconventional Jeff = 1/2 Mott ground state. The 'spin-orbit Mott state' is similar to conventional Mott states in 3d or 4d electron systems. However, the Jeff = 1/2 state is a considerably complex state hybridizing spin and orbital degrees of freedom. Therefore, investigation of this state may give a new insight into d-electron physics, and it may lead to possible unusual cooperative phenomena such as anisotropic superconductivity. In this presentation, we review electronic and magnetic states in the novel K<sub>2</sub>NiF<sub>4</sub>-type iridate Ba<sub>2</sub>IrO<sub>4</sub> [2]. We found that Ba<sub>2</sub>IrO<sub>4</sub> is a quasi-2-D (square lattice) Heisenberg antiferromagnet (TN ~ 240 K) in which the magnetic moment (~0.34 mB/Ir-atom) is significantly reduced by the low-dimensional quantum spin fluctuation with a large magnetic correlation |J|. Its electronic state undergoes an M-1 transition from the Mott insulating state to a non-Fermi-liquid state under high pressure (PC ~ 13.8 GPa) [3]. These behaviors are very similar to those in low-dimensional quartes.

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# **SF03**

#### Quantum critical end point in ucoal proved by NMR measurements

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Quantum critical end point (QCEP) emerges in a pressure - temperature - magnetic field phase diagram in ferromagnetic (FM) materials with tricriticality[1-3] In UGe<sub>2</sub>[1,2] a first-order metamagnetic transition, which appears above a critical pressure of ~1.5 GPa, possesses a critical end point (CEP) at a finite T\_CEP, above which the transition changes to crossover. The T\_CEP decreases under pressure, and reaches 0 K at the QCEP (P\_QCEP ~ 3.6 GPa and H\_QCEP ~ 18 T). In case of UCoAI, the QCEP is located at P\_QCEP ~ 1.5 GPa and H\_QCEP ~ 7 T.[3] Both systems exhibit the similar wing structures of the first-order plane, but they are quantitatively different. It is an interesting issue what induces such difference and what kinds of quantum criticality is involved in the vicinity of QCEP with both instabilities of the magnetism and the Fermi surface. In this presentation, we focus on 59Co-NMR measurements for UCoAI at ambient pressure [4] and under pressure. A contrasting behavior between the nuclear spin relavation rates 1/T<sub>1</sub> and 1/T<sub>2</sub> demonstrate that strong Ising-type magnetic fluctuation develops at around CEP. The drastic pressure evolution of the fluctuations toward the QCEP is also reported.

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 D. Aoki et al., J. Phys. Soc. Jpn. 80, 094711 (2011). [4] H. Nohara et al., J. Phys. Soc. Jpn. 80, 093707 (2011).

#### **SF04**

# The metal-insulator transition in ferromagnetic chromium hollandite

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We discovered a metal to insulator transition (MIT), retaining ferromagnetic in the insulator phase, in ferromagnetic K<sub>2</sub>Cr<sub>8</sub>O<sub>16</sub> [1]. The MIT is accompanied by a structural change from a tetragonal to a monoclinic structures [2,3]. In the low temperature monoclinic insulator phase, there is no evidence of charge separation/order. Instead, the characteristic displacements of the Cr and O sites, resulting in Cr-O bond alternations, are observed in the rectangular columns formed by the four chains (the four-chain columns). This indicates a lattice-dimerization with the formation of Cr-tetramer in the four-chain columns. The observed structural characteristics and electronic structure calcualitons lead us a model of MIT caused by the Peierls instability in the quasi-one-dimensional four-chain column [2], namely the MIT is a "pure" Peierls transition of spinless fermions (fully spin polarized electrons). On the other hand, isostructural RNCr<sub>8</sub>O<sub>16</sub>, shows a ferromagnetic transition at Tc= 270 K but does not show any distinct MIT. Rb-deficiency could be a reason for the absence of MIT. In a solid solution system, K2- yRbyCr<sub>2</sub>O<sub>16</sub>, the lattice parameters increase with increase of y. Te also increases as Rb concentration increases. This is consistent with the decrease of Tc with increasing hydrostatic pressure in K<sub>2</sub>Cr<sub>8</sub>O<sub>16</sub>.

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### SF05

#### Composition and transverse field-tuned quantum criticality in NbFe<sub>2</sub>

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The low-temperature band magnet NbFe2 can be tuned by varying composition, magnetic field or pressure, providing an attractive candidate system for investigating quantum criticality in a transition metal compound. Near the composition-tuned quantum critical point, NbFe2 displays signatures of a logarithmic Fermi liquid breakdown: The electronic specific heat diverges logarithmically  $\Delta C$ -T lnT and the resistivity exhibits non-Fermi-liquid forms  $\Delta p$ -T3/2 at low temperature T [1]. This quantum critical behavior is linked to the suppression of an antiferromagnetic phase in slightly Nb-rich samples which is stabilized towards Ferich samples. In addition, NbFe2 with slight Fe excess features a ferromagnetic ground state underlying the antiferromagnetic field perpendicular to the direction of the ordered moments, promoting NbFe2 as an interesting candidate to study transverse field tuning in an Ising-like metallic system. We report the significant suppression of Fermi liquid behavior close to the transverse-field induced quantum phase transition. Moreover, we utilize succeptibility measurements to study the nature of this quantum phase transition.

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# **SF06**

#### Tuning ferromagnetism in Ce<sub>1-x</sub>LaxAuGe: A specific heat and magnetic suscentibility study

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### **SF07**

#### Anomalous hybridization effects in the cubic quadrupole systems PrTr<sub>2</sub>Al<sub>20</sub> (Tr=Ti, V)

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The Kondo effect, the screening process of magnetic dipole moments by conduction(c) electrons. is a well-known many-body phenomenon and recognized as a key mechanism which induces a lot of interesting phenomena. Interestingly, nonmagnetic version of the Kondo effect, so-called quadrupolar Kondo effect, is theoretically predicted for the  $f_2 \Gamma_3$  crystalline electric field ground doublet[1]. In this theory, the ground state is no longer Fermi liquid due to the overscreening of the quadrupole moment by c-electrons. However, it is still uncertain what will happen in a real system with lattice periodicity. We have revealed that Kondo effect is observed in cubic PrTr<sub>2</sub>Al<sub>20</sub> for the first time in materials which have nonmagnetic  $\Gamma_3$  ground doublet. Quadrupolar ordering is observed at TQ = 2.0 K (Ti) and 0.6 K (V) respectively, and in particular ferroquadrupolar ordering in PrTi<sub>2</sub>Al<sub>20</sub>[2, 3, 5, 6]. Interestingly, anomalous metallic behaviors were observed in PrV<sub>2</sub>Al<sub>20</sub> above TQ, most likely attributed to the quadrupolar Kondo effect. In this presentation, we will report the results of low temperature thermal and transport measurements of PrTr<sub>2</sub>Al<sub>20</sub> and discuss the possibility of the quadrupolar Kondo effect and quantum critical effect of quadrupolar ordering.

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### **SF08**

#### Non-fermi-liquid properties of the non-centrosymmetric heavyfermion compound CePtSi: a magnetic field study.

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The ternary intermetallic compound CePtSi and its nonmagnetic equivalent LaPtSi form in the non-centrosymmetric tetragonal space group I41/md [1]. The structure is an ordered variant of the ThSi2 type which provides dedicated atomic sites for each of the rare-earth atom, Pt, and Si. CePtSi is a heavy-fermion compound with a very large electronic specific heat coefficient [2] and apparently with no spontaneous magnetic ordering. Non-Fermi-liquid (nFL) behaviour found at low temperatures in a Pt-Si variational study of CePtSi has been reported [3], as well as aspects of nFL signatures at the verge of magnetic ordering in CePtSi doped with Ge [4]. In this work we focus on the low-temperature behaviour of magnetic susceptibility, electrical resistivity, and specific heat of CePtSi together with comparative properties in LaPtSi. The focus is on the influence of applied magnetic fields and the stability of nFL scaling in CePtSi as a function of field.

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### SF09

#### Vibron quasi-bound state in the non-centrosymmetric tetragonal heavy-fermion compound CeCuAl<sub>3</sub>

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We have investigated the non-centrosymmetric tetragonal heavy-fermion antiferromagnetic compound CeCuAl, (TN=2.5 K) using inelastic neutron scattering (INS) Our INS results unequivocally reveal the presence of three magnetic excitations centred at 1.3, 9.8, and 20.5 meV. These spectral features cannot be explained within the framework of crystal-electric-field (CEF) models and recourse to Kramers theorem for a 4f. Ce<sup>3+</sup> ion. To overcome these interpretational difficulties, we have extended the vibron model for cubic CeAl, of Thalmeier and Fulde to tetragonal pointgroup symmetry. This extension provides a satisfactory explanation for the position and intensity of the three observed magnetic excitations in CeCuAl3, as well as their dependence on momentum transfer and temperature. On the basis of our analysis, we attribute the observed series of magnetic excitations to the existence of a vibron quasibound state.

#### **SF10**

#### Pressure-induced quantum criticality in the heavy-fermion compound CeCoGe<sub>2.2</sub>Si<sub>0.8</sub>

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The pressure-tuned quantum critical point (QCP) of the antiferromagnetic heavy-fermion compound CeCoGe21Si09 was claimed to be dominated by different effects on both sides of the QCP: on the magnetic side spin fluctuations govern the criticality, while on the non-magnetic side the criticality is dominated by disorder that quenches the spin fluctuations [1]. Here we study high-quality CeCoGe<sub>22</sub>Si<sub>0.8</sub> samples [2], with residual resistance ratios four times larger than those of the previously investigated CeCoGe, Sing samples Interestingly, while DC magnetic susceptibility measurements show that the Neel temperature of TN = 4K at zero pressure is only slightly reduced by pressure up to 3.0 kbar, a much stronger decrease is observed for the specific heat anomaly. We will present electrical resistivity and specific heat measurements up to 15 kbar and down to 0.05 K, and establish the pressure - temperature phase diagram for CeCoGe<sub>2</sub>Si<sub>0.2</sub>. The critical behavior shall be compared to the one observed for both CeCoGe, Sing and the pure reference compound CeCoGe<sub>3</sub> [3]. We hope that this investigation will elucidate the role of disorder. Acknowledgement: Financial support by the European Research Council (ERC Advanced Grant No 227378) is acknowledged.'

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#### SF11

#### Anomalous thermoelectric effects in the heavy fermion superconductor Ce<sub>2</sub>PdIn<sub>8</sub>

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# SF12

#### Unusual normal-state magnetotransport in the heavy-fermion superconductor Ce<sub>2</sub>PdIn<sub>6</sub>

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The normal-state behavior in the heavy-fermion superconductor Ce-PdIn. [1,2] has been probed by means of Hall coefficient (RH) and transverse magnetoresistivity (MR) measurements. Anomalous non-Fermi-liquid-like features, observed below 8 K in both RH(T) and MR(T), are related to underlying quantum critical point, previously evidenced in the specific heat and the electrical resistivity data [1-5]. Intriguing magnetotransport properties observed in this compound exhibit close similarity to those reported before for CeCoIn, and can be explained in terms of the recently developed theory in which anisotropy of the Fermi surface and strong backflow scattering of charge carriers on critical AF spin fluctuations are taken into account. What's more the magnetotransport in Ce<sub>2</sub>PdIn<sub>8</sub> is shown to exhibit novel types of scaling that may appear universal for similar systems being at the verge of magnetic instability

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#### SF13

#### Specific heat and thermal conductivity studies of UCu<sub>tix</sub>Alex compounds

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We report the thermal-conductivity and specific-heat measurements for eight  $UCu_{4x}Al_{8x}$  compounds ( $0 \le x$ ≤2.0) as a function of temperature and magnetic field. Previous studies provided evidence of a transition from magnetic to a non-magnetic heavy fermion state at xcr = 1.15. Here we present complementary thermal-property studies as a function of temperature (55 mK - 300 K) and applied magnetic field (0 - 15 T). Specific heat data of non-magnetic compound at xer = 1.15 shows logarithmic divergence of low-temperature C/T and this non-Fermi liquid (NFL) behavior can be attributed to the proximity of a quantum critical point. Non-magnetic compounds with higher Cu content (x > xcr) exhibit unusual temperature scaling in the specific heat possibly due to an increase in disorder between Cu and Al. The complete study of magnetic field effect on NFL behavior was hampered by high contribution from nuclear Schottky anomaly at low temperatures. The thermal conductivity data show stark contrast between the magnetic (x = 0.5) and non-magnetic compound (x = 1.75) - as evidenced in Lorentz numbers for example. Our results confirm that a simple free-electron picture is inadequate for the description of the low-temperature thermal conductivity properties in non-magnetic UCu<sub>4:1</sub>Al<sub>8.8</sub> compounds.

#### SF14

#### Dynamical cluster approximation results of the two-orbital Hubbard model

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The orbital-selective phase transition (OSPT), which is induced by the interplay of a narrow band of localized electrons and a wide band of itinerant electrons, has been intensively explored in the multi-orbital systems. Using a dynamical cluster approximation including the short-range correlation as well as quantum fluctuations, we would like to analyze the behaviors of both degenerate two-orbital Hubbard model, where both orbitals have equal bandwidths and one orbital is constrained to be paramagnetic, while the second one is allowed to have an antiferromagnetic solution and two-orbital Hubbard model with different bandwidth at half-filling. The OSPT, Fermi-liquid (FL), non-FL, and insulator behaviors are observed in the both cases. Finally, we would discuss the implications of the results in the context of the Fe-based superconductors

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# **SF15**

Renormalized parameters and convergence of energy scales on the approach to local quantum critical points

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We consider a two impurity/dot Anderson model, where the impurities, hybridized with separate conduction baths, are coupled by antiferromagnetic exchange interaction J and a direct interaction U<sub>22</sub>. The model has two types of quantum critical points associated respectively with a transition between Kondo singlet to a local state and a transition to a local charge ordered state. Using exact expressions for the spin and charge susceptibilities in terms of renormalized parameters we predict the values of the renormalized parameters on the approach to both types of critical point in terms of a single energy scale  $T^*$ , such that at the critical point  $T^{*-->} 0$ . From numerical renormalization group calculations we give complete phase diagram, and from an analysis of the low energy fixed point deduce the renormalized parameters. The results confirm the predictions of the renormalized perturbation theory and determine the value of T\*. The emergence of a single energy scale which goes to zero at a quantum critical point may be a general feature which would lead to a natural explanation of E/T scaling which has been observed at the quantum critical points of some heavy fermion compounds.

### **SG01**

#### NMR study of magnetic properties of Eu<sub>1-x</sub>SrxMnO<sub>3</sub>

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Perovskite-type manganite EuMnO<sub>2</sub> is known to be a compound with narrow eg-band width due to small Eu<sup>3+</sup> ionic radius compared to La<sup>3+</sup>. The Eu<sub>1-x</sub>SrxMnO<sub>3</sub> shows complex magnetic properties, depending on Sr concentration x [1,2]. We have carried out 55Mn spin-echo NMR measurements on polycrystalline Eu1-xSrxMnO3 (0≤x≤0.5) samples at 4.2 or 1.4K, in order to clarify the microscopic local magnetic state of Mn ions and the inhomogeneity in the compounds. Typical NMR spectra have been observed in the frequency range from 250MHz to 400MHz. For the A-type antiferromagnetic EuMnO<sub>2</sub>. the resonance frequency is about 260MHz, which is much lower than that observed for  $Mn^{3+}$  in La<sub>1,x</sub>SrxMnO<sub>3</sub>. For  $0.1 \le x \le 0.4$ , the resonance lines, corresponding to  $Mn^{4+}$  and  $Mn^{3+}$  have been observed around 310MHz and 390MHz, respectively. Furthermore the spectra originated from Mn<sup>2+</sup> have been observed around 590MHz for 0.1<x<0.4 For Eu<sub>0.5</sub>Sr<sub>0.5</sub>MnO<sub>3</sub>, the NMR spectrum spreads widely with several distinct peaks in the frequency range from 250MHz to 410MHz. In Eu1-xSrxMnO3, the spin-glass like behavior has been observed around x = 0.5[1,3]. The present NMR spectrum for Eu<sub>0.5</sub>Sr<sub>0.5</sub>MnO<sub>3</sub> shows that the compound is in the inhomogeneous state due to competition between the antiferromagnetic and the ferromagnetic interactions.

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### SG02

#### NMR study of successive magnetic transitions in A-site ordered perovskite LaMn<sub>3</sub>Cr<sub>4</sub>O<sub>12</sub>

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In a simple ABO3 perovskite, when 75% of the A-site cations are substituted by another element, an A-site-ordered perovskite with a chemical formula of A'A3B4O12 can be formed. These materials have attracted much attention, because they show a rich variety of physical phenomena, such as heavy-fermion behavior in CaCu3Ru4O12 and giant magnetoresistance in CaCu<sub>3</sub>Mn<sub>4</sub>O<sub>12</sub>. These features may be related to the A-A and A-B exchange interactions in addition to the usual B-B exchange interaction seen in simple perovskite materials. Recently, new A-site ordered perovskite, LaMn<sub>3</sub>Cr<sub>4</sub>O<sub>12</sub>, was prepared by using high pressure synthesis. The measurements of magnetic susceptibility and specific heat suggest two antiferromagnetic transitions at 150 K and 50 K in LaMn<sub>3</sub>Cr4O<sub>12</sub>. In this conference, we report microscopic investigation on the magnetic properties of this material probed by La NMR. The resonance line moderately broadens below 150 K and disappears below 50 K. These results are consistent with the two antiferromagnetic transitions. The temperature dependence of Knight shift indicates that the transition at 150 K is associated with the B-site Cr spin ordering and the other at 50 K is due to the A-site Mn spin ordering.

### **SG03**

### Magnetic properties of single crystalline U<sub>2</sub>Fe<sub>3</sub>Ge

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A number of Laves phases exist in intermetallic compounds based on uranium and transition metals. The compact structure of such phases, which leads to small U-U spacing and high U coordination numbers, gives rise to the 5f band and strong hybridization with the non-f states. U<sub>2</sub>Fe<sub>3</sub>Ge is a new ternary Laves phase [1], crystallizing in the Mg<sub>2</sub>Cu<sub>3</sub>Si-type structure, an ordered variant of the hexagonal MgZn<sub>2</sub>-type (C14, P6/mmc). It exhibits very short U-U distances (below the Hill limit) along the c-axis (3.2 Å), still reaching a ferromagnetic order (T<sub>c</sub> = 55 K). Surprisingly the Fe sub-lattice does not carry any significant ordered moment. A single crystal of U<sub>2</sub>Fe<sub>3</sub>Ge was prepared by the Czochralski method (a = 5.182(2) Å and c = 7.850(1) Å). Magnetization isotherms measured along the principal axes indicate that the magnetic moments lie in the basal plane of the hexagonal lattice, with the spontaneous magnetic moments of 1 µ<sub>b</sub>/f.u. and no anisotropy within the basal plane. Magnetic moments are therefore perpendicular to the nearest U-U link, but the anisotropy values are anomalously low comparing e.g. to isostructural UNi<sub>2</sub>.

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#### SG04

# Magnetic properties of a 5d transition metal oxide AOsO<sub>4</sub> (A = K, Rb, Cs)

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We investigate the magnetic properties of AOsO<sub>4</sub> (A = K, Rb, Cs) which takes a 5d<sub>1</sub> electron configuration and would provide us with an interesting playground for spin-orbit physics. The compound crystallizes in the scheelite structure comprising a distorted diamond network of Os atoms. Magnetic susceptibility at high temperatures shows Curie-Weiss behavior with Weiss temperatures of -94, -29 and 12 K for the K, Rb and Cs compounds, respectively. For the K and Rb compounds, broad humps are observed at 55 and 35 K upon cooling and followed by antiferromagnetic transition is observed down to 2 K, but an anomaly appears at 151 K, below which the Weiss temperature is reduced to 3 K. Markedly, the magnetic effective moments estimated from the Curie constants are 0.85, 0.84, and 0.86 Bohr magneton for the K, Rb, and Cs compounds, which are much smaller than the spin-only value of 1.73 Bohr

#### **SG05**

# Enhancement of curie temperature due to the coupling between fe itinerant electrons and Dy localized electrons in $DyFe_2Zn_{20}$

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The cubic RX2Zn20 has been recently insensitively examined, where R is rare earth atom and X is Fe, Co, Ru, etc. The exchange interaction between R-R is weak because R atoms are diluted in this compound. In fact, for X = Co and Ru, Tc is less than 10 K for any R. In the case of Fe, however, Tc is significantly enhanced, and the origin of this high Tc is suspected to be involved the magnetism of Fe atoms. But the origin is not clear at present. We examined the magnetic properties and specific heat of DyFe2Zn20, and revealed experimentally unusual magnetic anisotropies, i.e., large magneto-crystalline anisotropy at 2 K, which nature, however, disappears in the temperatures above 30 K below Tc. We analyzed these anomalous properties based on the three assumptions: (1) crystalline electric field and exchange interaction between R and R, (2) simplified Moriya theory for Fe, and (3) exchange interaction between R and Fe. As a result, the enhancement of Tc, the strong magnetic anisotropy at low temperatures, the disappearance of the magnetic anisotropy below Tc down to 30 K, and no anomaly of magnetic specific heat at Tc, etc. were well reproduced by this calculation.

### SG06

# Anomalous increase of TC in UGa2 under pressure

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#### SG07

# Phase transition between paramagnetic and spin polarized states in MnSi

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The of high-field region of the magnetic phase diagram of MnSi is probed by magnetization, resistivity and magnetoresistance measurements carried out in the temperature range 1.8-300 K for magnetic fields up to 8 T [1]. It is shown that the phase boundary between paramagnetic (PM) phase and spin-polarized (SP) phase is well-defined, has no positive slope as it was suggested previously, and appears to be practically vertical corresponding to transition temperature To-30 K. We argue that broad maxima of the resistivity and magnetization derivatives, which develops in the diapason T >Tc, are not associated with a "diffuse" SP-PM transition and are a consequence of the specific form of functional dependences of these quantities in the paramagnetic phase. It is found that in the paramagnetic phase of MnSi a universal relation between magnetoresistance and magnetization,  $\Delta\rho/\rho=$  a0M<sup>2</sup>, holds in a wide range where magnetoresistance data favors the explanation of magnetic properties of MnSi by Heisenberg-type localized magnetic moments rather than by itinerant magnetism approach. A low temperature anomaly at T~15K corresponding to magnetoresistance and g factor changes is reported.

1. S.V.Demishev et al., Phys. Rev.B 85, 045131 (2012)

#### **SG08**

# ESR in mnsi: Heisenberg localized magnetic moments and spin polarons

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High frequency (60 GHz) electron spin resonance (ESR) has been studied in manganese monosilicide, MnSi, single crystals [1]. The measurements performed within the 4.2-300 K temperatures range at the applied magnetic field up to 70 kOe have demonstrated that the magnetic resonance in MnSi is due to localized magnetic moments of the Heisenberg type with the g factor depending only slightly on temperature,  $g \sim 1.9-2$ . At the same time, it has been found that the ESR linewidth is determined by spin fluctuations and can be quantitatively described in the wide temperature range (4.2 K < T < 60 K) in the framework of the Moriya theory using the SL(T) function. The revealed deviations from the model of weak titnerant electron magnetism commonly used for the description of the magnetic properties of MnSi indicate a possible spin-polaron nature of the unusual magnetic properties of this strongly correlated metal.

1. S.V.Demishev et al., JETP Letters, 93, 213 (2011)

# SG09

# Magnetic susceptibility measurements at high pressures down to T=0.5 K with SQUID magnetometer

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We develop a 3He-insert attached to MPMS SQUID magnetometer and a pressure cell in order to measure magnetic susceptibility and magnetization precisely at high pressures down to T=0.5 K[1]. The 3He insert is made by stainless steel pipe with the inner diameter of 6.2mm and outer diameter of 6.5 mm. The liquid 3He is condensed in the stainless pipe. The vacuum jacket made of a copper pipe with the outer diameter of 8.6 mm is soldered to the stainless pipe, which is also used for the heat exchange between 3He gas and 4He bath at T=1.7K. The temperature was measured by RuO<sub>2</sub> thermometer mounted on the 3He container. A main body of piston-cylinder pressure cell is made from Be-Cu. The outer diameter of the cell is 6 mm, while the inner diameter is 1.8mm which corresponds to the size of the teflon cell.

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#### SG10

# Extinction of photo-luminescence of Mn-doped ZnS nanocolloids in weak magnetic field

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The extinction of light passing through a nanocolloidal medium containing ferromagnetic oxide (such as  $Fe_3O_4$ ) putting inside a cavity of weak magnetic field has been observed [1-3]. Here for the first time we report the extinction of photo-luminescence of Mn-doped ZnS nanocolloid under the applied field of 270 Gauss. The Mn-doped ZnS nanopowders have been synthesized by using the hydrothermal method from zinc acetate and natri thiosulfat as the precursors (with SPAN-80 as the surfactant). The particle size was determined to be within 10-20 nm; the lattice constant a = 5.41(2) A. The exitation wavelength was 623.8 nm from a He-Ne laser. We have observed a systematic reduction of photo-luminescence of a given liquid nanocolloid with time at applied field of 270 Gauss, despite of the field direction and field development (upwards or downwards). The 50% extinction was achieved after 30 min and the saturation extinction after 1 hour. The memory effect was shown to preserve for more than 2 hours. We have measured the temperature of the sample during the experiment which was kept constant at 70 oC. We believe the PL extinction was associated with the magnetic ordering but the ordering inside a liquid medium is a surprise.

1. John Philip, J. M. Laskar and Baldev Raj. Appl. Phys. Lett. 92, 221911 (2008). 2. Rajesh Patel. Pure Appl. Opt. 11 (2009) 125004 (Spp). 3. R.V. Mehta, Rajesh Patel, Rucha Desai, R.V. Upadhyay and Kinnari Parekh. Phys. Rev. Lett. 96,127402 (2006).

#### SG11

# Electronic structure of A-site ordered perovskite CaCu<sub>3</sub>Ti<sub>4</sub>O<sub>12</sub> studied by angle-resolved photoemission spectroscopy

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We present angle-resolved photoemission spectroscopy (ARPES) results of A-site ordered perovskite CaCu<sub>3</sub>Ti<sub>4</sub>O<sub>12</sub> (CCTO), which shows extremely high-dielectric constant over a wide range of temperature from 100 to 600 K and has an insulator phase in contrast to a metal phase from the result of LDA band calculation. We have observed the clear band dispersions, located in the higher binding energy than that expected in the LDA calculation, and the negligible spectral weight at the Fermi level in agreement with the results of electrical resistivity measurements. From the ARPES results, we suggest that CCTO is Mott-insulator caused by the strong correlation effects of the electrons in Cu 3d - O 2p hybridization bands.

# SG12

# Evidence of rattling transition in caged compounds $LaRu_2Zn_{20}$ and $LaIr_2Zn_{20}{:}$ La-NMR studies

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Recently, RT\_ZT<sub>20</sub> (R = Pr, La; T = Ir, Ru) has attracted much attention because of the coexistence of superconductivity and a quadrupole ordered state observed in PrIr\_ZT<sub>20</sub>[1,2] Superconductivity is also observed in LaRu\_ZT<sub>20</sub> (Te=0.2 K) and LaIr2Zn20 (Te=0.6 K), but not in PrIRu\_ZT<sub>20</sub>[1]. PrRu\_ZT<sub>20</sub> and LaRu\_ZT<sub>20</sub> undergo first-order structural phase transition at Ts=138 and 150 K, respectively, whereas LaIr\_ZT<sub>20</sub> undergoes second-order structural phase transition at Ts=200 K. To investigate the structural phase transition in RT\_ZT<sub>20</sub>, 139La(I=7/2) nuclear magnetic resonance (NMR) measurements have been carried out for LaRu\_ZT<sub>20</sub> and LaIr\_ZT<sub>20</sub>. For both compounds, quite narrow La-NMR line-widths less than 5 kHz without nuclear quadrupole splittings above Ts changes their spectral shape with broad tails distributed over the frequency range of 0.8 MHz, evidencing the lowering of the symmetry at La site[3]. Around Ts, the nuclear spin-lattice relaxation rate (1/T1) show an unusual enhancement associated with the slowing down of the EFG fluctuations[3]. To gain further insights, we have measured the nuclear spin-spin relaxation rate (1/T2) for both compounds. We have found that the temperature dependence of 1/T2 has maximum around Ts associated with the structural phase transition. We will discuss low-energy spin dynamics in RT\_ZT<sub>20</sub> from microscopic points of view.

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### **SG13**

# Unconventional magnetic ordering in spin-orbit mott insulator with honeycomb lattice

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Iridates have recently attracted much attention due to a novel  $SS_{mem} = 1/2S$  Mott insulating state, driven by the interplay of moderate electronic correlations with strong spin-orbit coupling. We have synthesized A<sub>2</sub>IrO<sub>3</sub> (A=Na,Li) which is a layered system with Ir moments sitting on a Honeycomb lattice. Theoretically, this system has been proposed as solid-state realization of the Heisenberg-Kitaev (HK) model, treating the superposition of isotropic antiferromagnetic exchange and highly anisotropic Kitaev type ferromagnetic nearest neighbor exchange. HK model predicts gapless spin-liquid ground state with Majorana fermionic excitations in Kitaev Limit[1]. Na,IrO<sub>3</sub> shows a Mott insulating state of SS\_{mem} eff]=1/2S moments with predominant antiferromagnetic coupling, indicated by a Weiss temperature of S/Theta\_W=-116S K[2]. A bulk antiferromagnetic transition occurs at a much reduced temperature of S/Theta\_W=-116S K[2]. A bulk antiferromagnetic structure[3]. For Mott insulating Li,IrO<sub>3</sub> we observe a similar ordering temperature of 15 K, while the S/Theta\_WS is drastically reduced to S-33KS. These observations are compatible with an enhancement of the Kitaev contribution compared to the Na-system, suggesting that Li,IrO<sub>3</sub> is located closer to the Kitaev contribution compared to the Murdus EURINDIA project and the AvH foundation.

[1] J. Chaloupka et al, Phys. Rev. Lett. 105, 027204 (2010), [2] Y. Singh and P. Gegenwart, Phys. Rev. B 82, 064412 (2010), [3] X. Liu et al., YSingh et al., Phys. Rev. B (bbf83), 220403(R) (2011), [4] Y. Singh, S. Manni, P. Gegenwart, arXiv:1106.0429

### SG14

#### YCr<sub>6</sub>Ge<sub>6</sub>: A kagome metal?

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Itinerant electrons on a kagome lattice are known to generate a flat band without dispersion. It is expected that the flat band causes exotic phenomena induced by the kagome geometry such as a flat-band ferromagnetism. However, this has not been experimentally observed in real compounds, possibly because the flat band tends to be located far from the Fermi level. We have been looking for a candidate compound to realize "kagome metal" and recently focused on  $YCr_6Ge_6$ . It crystallizes in the HfFe<sub>6</sub>Ge<sub>6-type</sub> structure (space group P6/mmm) where Cr atoms form a kagome lattice. Band-structure calculations show that a flat band really exists slightly below EF near the  $\Gamma$  point toward the K and M points. This suggests that hole doping, such as the chemical substitution of Ga for Ge, would make the flat band emerge at EF. We successfully synthesized single crystals of  $YCr_6Ge_6$  and  $YCr_6Ge_6.ax$  by the flux method. A Curie-Weiss like behavior is observed in magnetic susceptibility for both compounds. We will show the magnetic and transport properties and try to reveal the characteristics of the kagome metal.

### SG15

#### Raman scattering spectra of PrRu<sub>2</sub>Zn<sub>20</sub>

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 $PrRu_2Zn_{20}$  has the cubic structure with Fd-3m, where Pr-ion is encapsulated in the Frank-Kasper cage by 16 Zn atoms. In analogy of the filled skutterudites, this compound is expected to be observed an anharmonic Pr vibration and exotic physical properties due to orbital degrees of freedom of the local f electrons. shows structural transition at Ts=138K and remains a normal state in electronic conductivity down to 0.04K. On the other hand, isostructural PrIr<sub>2</sub>Zn<sub>20</sub> shows an antiferro-quadrupolar ordering(Tq=0.11K) and a superconductivity(Tc=0.05K), without structural change above Tq. The difference of the low temperature properties between PrRu<sub>2</sub>Zn<sub>20</sub> and PrIr<sub>2</sub>Zn<sub>20</sub> might be originated from the structural transition, however, the origin of the structural transition and also the crystallographic structure below Ts are not clarified. Therefore, we have measured Raman scattering spectra of  $\text{PrRu}_2\text{Zn}_{20}$  . Among the Raman active phonons; 3Ag<sup>+4</sup>Eg<sup>+9</sup>T2g, 2Ag<sup>+4</sup>Eg<sup>+6</sup>T2g peaks are observed at room temperature. With decreasing temperature, 4 peaks are newly observed below Ts in T2g spectra. This suggests the structural change around Ts. In comparison with Eg spectra, the observation of the similar energy peaks in both spectra suggests the crystal structure below Ts seems to be lower than cubic structure.

#### **SG16**

# Superconducting state in KSn<sub>2</sub> with a MgZn<sub>2</sub>-type (C14) Laves phase structure

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We have been searching for a new superconductor including alkali metal elements using high pressure / high temperture technique. High-pressure synthesis technique has been successfully employed in the search for new materials and new superconductors. To use a closed system for high pressure synthesis is effective not only to stabilize a composition but also to expand a solid-solution range. We successfully synthesized the polycrystalline sample of KSn<sub>2</sub> as a main phase and discovered that KSn<sub>2</sub> was a superconductor with  $T_{\rm c}$  = 3.2 K. KSn<sub>2</sub> has a hexagonal MgZn<sub>2</sub>-type (C14) structure (space group P6<sub>3</sub>/mmc, No.194) called Laves phase. The magnetization versus magnetic field curve shows a typical type-ll superconducting behavior. We determined the lower critical field and penetration depth, to be about  $H_{\rm e}(0)$  = 150 Oe and  $\lambda_{\rm cal}(0)$  = 210 nm, respectively, using the Ginzburg-Landau equations. From the density of state calculations, we found that Sn p-orbital mainly contributes near the N(E\_{\rm F}) level and plays an important role for the superconducting state in KSn<sub>2</sub>.

#### **SG17**

# Layered nanosized structures on basis of diluted magnetic semiconductors and heusler alloys

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The laser and RF magnetron synthesis of nanosized layers of the diluted magnetic semiconductors (DMS) on the basis of compounds III-V and elementary semiconductors Ge or Si with Mn or Fe impurities and Co<sub>2</sub>MnSi Heusler alloys (HA) on single crystal substrates GaAs, Si or Al<sub>2</sub>O<sub>3</sub> with ferromagnetism till 500K was earlier demonstrated [1-3]. In this report our new researches of properties of DMS Si:Mn and magnetic tunnel junctions (MTJ) with ferromagnetic HA plates and dielectric interlayer are presented. The features of FMR of MTJ with ferromagnetic exchange interaction, anisotropic negative magnetoresistance, I-V nonlinearity and hysteresis of MTJ and single DMS and HA layers at small current density were observed at 77-300K. The high-resolution transmission electron microscopy (HRTEM) and selected area electron

diffraction (SAED), made on JEM-2100F of JEOL, show that the laser synthesis of epitaxial layers of DMS Si:15%Mn/GaAs with Curie point 500K, perfect diamond like crystal structure and self-organized superlattice is possible. Supported by RFBR (08-02-01222-a, 11-02-00855-a), the Ministry of Education of Russian Federation (projects 2.1.1/2833 and 2.1.1/12029) and the State contract № 02.740.11.0672 of the Federal purpose program <<Scientific and scientific-pedagogical cadre of innovative Russia>2009-2013 is acknowledged.

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### SG18

#### Metal-insulator and spin-state transition in polycrystalline (Pr<sub>1-y</sub>REy)<sub>1-</sub>, CaxCoO<sub>3</sub> (RE=rare earth elements) in magnetic fields

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We have measured the temperature dependences of the resistivity and the magnetic susceptibility in various magnetic fields up to 17 T and 5 T, respectively, for polycrystalline  $(Pr_{1,\gamma}REy)_{1,\zeta}CaxCoO_3$  (RE=rare earth elements). For  $(Pr_{0,S}Sm_{0,2})_{0,\gamma}Ca_{0,3}CoO_3$  sample, a metal-insulator transition (MIT) and spin-state transition (SST) between the intermediate (IS) and the low spin (LS) states almost simultaneously took place around 40 K at 0 T. The MI-SST temperature decreased with increasing magnetic field, and the MI-SST was suddenly suppressed above 10 T. The MI-SST temperatures were about 38 K at 5 T and 29 K at 9 T. In the samples with higher y, in which the MI-SST temperature at 0 T was higher, the reduction of the MI-SST temperature became small and the MI-SST survived up to the highest magnetic field in this study. The similar behavior was observed in other RE substituted samples. The obtained results suggest that the magnetic field stabilizes the IS state, which contrasts with the fact that the application of the external pressure induces the LS state. We discuss the origin of the magnetic field effect on the MI-SST.

#### SG19

# Magnetic field-induced lattice effects in a quasi-2D organic conductor close to the Mott metal-insulator transition

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Organic charge-transfer salts of the  $\kappa$ -phase (BEDT-TTF)§\_2SX family have been recognized as prototypes for studying electronic correlation phenomena [1]. Their ground states can be tuned either by varying the counter anion X or by applying external pressure. For the X = Cu[N(CN)\_2[C] salt, for example, the ground state is an antiferromagnetic-Mott insulator, whereas for X = Cu[N(CN)\_3]Br, it is a superconductor. Fully deuterated single crystals of the latter substance are located on the verge of the Mott metal-insulator (MI) transition and, for this reason, constitute a material of particular interst for exploring the critical behavior of the  $\pi$ -electron system around the Mott transition [2,3]. Here we present ultra-high-resolution dilatometry studies at varying magnetic fields on a fully deuterated X = Cu[N(CN)\_3]Br salt. Our thermal expansion data reveal two remarkable features: (i) the Mott MI transition temperature TS\_{MI}S = (13.6 \pm 0.6)K is insensitive to fields up to 10.7, (ii) for fields along the interlayer b-axis, a B-induced first-order transition at TS\_{SF}S = (9.5 \pm 0.5)K is observed. This transition manifests itself in a sharp negative spike in the expansivity, whose position is found to be field insensitive. Possible scenarios to describe this field-induced phase transition will be discussed.

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#### SG20

# Metallic transition of the colossal magnetoresistance material $FexMn_{1-x}S$ (x=0.18) under high pressure

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The discovery of the colossal magnetoresistance in FexMn<sub>1,x</sub>S solid solutions [1] aroused many researcher's interests in studying the physical properties of this system. We studied the pressure induced phase transition of FexMn<sub>1,x</sub>S (x=0.18) by infrared (IR) reflection and X-ray analysis up to 40 GPa at room temperature. It is shown by X-ray analysis that the structure of this sample is NaCl type at ambient pressure and a structural change starts around 17 GPa and the mixed phase between the NaCl type low pressure phase and the structure unknown high pressure phase continues up to around 25 GPa. On the other hand, the IR reflectivity increases from 15 GPa and becomes highly remarkable around 20 GPa. The spectra do not show any changes from 30 GPa. These results suggest that the phase transition of FexMn<sub>1,x</sub>S (x=0.18) at room temperature starts around 15 GPa and completes around 30 GPa and the high pressure phase is not a band overlapping semimetal but a true metallic.

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# SG21

# Optimal design of IPMSM having double barrier for minimizing cogging torque and torque ripple

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In this paper, a design method for shape optimization of stator and rotor of a double barrier in the form of IPMSM for NEV(Neighborhood Electric Vehicle) is proposed. IPMSM has two big problems which are high cogging torque and high torque ripple that causes performance degradation of the motor and noise, and vibration during the motor rotation. To solve the problem, ways of reducing the cogging torque and lorque ripple as a method for forming the shape of the motor stator and rotor factors were searched. Especially, the factors which greatly affected cogging torque and torque ripple werified their mutual connection. Moreover it was found that the most influential factors was the slot open and the slot teeth angle. The results of the experiment which reduced cogging torque and torque ripple verified, while the optimization proceeded on the basis of the factors that were searched. So the prototype of the optimized IPMSM was designed, its performance was evaluated. The optimization was verified by the performance evaluation.

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#### SG22

# Implementation of first-principle calculation in combination with a dynamical cluster approximation

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While the combination of the local density approximation (LDA) with dynamical mean field theory (DMFT) opened the door to a realistic description of correlated compounds, the improvement of the method is ongoing. In this work, we present a combination of LDA with the dynamical cluster approximation (LDA+DCA) in the framework of the full-potential linear augmented plane-wave method, and compare our results to LDA+DMFT calculations as well as experimental observations on SrVO<sub>3</sub>. We find a qualitative agreement of the momentum resolved spectral function with angle-resolved photoemission spectra (ARPES) and former LDA+DMFT results. As a correction to LDA+DMFT, we observe more pronounced coherent peaks below the Fermi level, as indicated by ARPES experiments. In addition, we resolve the spectral functions in the  $\theta = \theta = 0$ , 0,0,0 and  $\theta = \theta = 0$ , 0,0,0 sectors of DCA, where band insulating and metallic phases coexist. Our approach can be applied to correlated compounds where not only local quantum fluctuations but also spatial fluctuations are important.

arXiv:1111.0890

#### SH01

# Oscillatory transformative domain wall inner structure of depinning domain wall around notched ferromagnetic wire

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We have investigated domain wall (DW) depinning behavior around a geometric notch in ferromagnetic nanowires by means of micromagnetic simulation [1]. Nanowires with different notch geometries by changing a depth (d) and a half-width of the notch (s) have been considered A domain wall with a transverse wall structure is initially prepared to be positioned at the notch. Depinning field pulse of 1 ns duration is applied to trigger DW depinning field transformation of depinning field strengths. With increase of s, the domain wall depinning field transformation of depinning field strengths. With increase of s, the domain wall depinning field strengths. Change of an aspect ratio is also considered, where a relatively insensitive variation of depinning field strength, we have found that DW internal structure changes during the depinning process. At lower depinning field (< 4 mT), DW keeps the initial transverse wall structure on the inner structure with antivortex soon after DW has escaped from the noteh. The transformation of the DW inner structure is obviously related to the Walker breakdown phenomenon.

[1] M. J. Donahue and D. G. Porter, OOMMF User's Guide, http://math.nist.gov/oommf (2002).

# **SH02**

#### Unidirectional thermal effects in current-induced domain wall motion

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We show unidirectional motion of a domain wall (DW) due to thermal effects in permalloy nanostrips using magnetic force microscopy on a RF-contacted sample. A deterministic thermal driven motion has been found for current excitation in the range of few nanoseconds. The cooperation between thermal effect and spin transfer torque induces larger and more effective displacements towards the center of nanostrip, the hotter part. The thermal effects are generated by Joule heating, and amplified by poor heat diffusion through thick SiO<sub>2</sub> (100 nm). An important issue in thermal transport is the temperature gradient related to the positions of sinks (thick Au electrodes), responsible (partly at least) for the unidirectional DW motion towards hotter part. Such behaviour has been recently predicted by other thermoelectric effects. However, we anticipate the spin currents created by Spin Seebeck effect [2] or anomalous Nernst effect [3] to be small compared to the critical current density for DW motion.

[1] D. Hinzke et al; PRL 107, 027205 (2011). [2] K. Uchida et al; Nature 455, 778 (2008). [3] S. Y. Huang et al; PRL 107, 216604 (2011).

### SH03

# Effect of the oersted field on current-induced domain wall motion and domain wall chirality in multilayer nanostripes

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Current-induced domain wall motion is more efficient in Permalloy(Py)/Cu/Co trilayer nanostripes than in single Py nanostripes. We used time-resolved photoemission electron microscopy combined with x-ray magnetic circular dichroism (XMCD-PEEM) to image the magnetic configuration in the Py layer of these nanostripes, during the current pulses, with 50ps time resolution. During the current pulses the Py magnetization tilts towards the transverse direction. The large tilt angles (up to 70°) can be reproduced by micromagnetic simulations only when both the Oersted field and the magnetostatic interactions between Py and Co are taken into account. Magnetic Force Microscopy (MFM) measurements were performed on Py/ Metal bilayer nanostripes, with different metal thicknesses, varying the transverse field for a given current density. We show that the Oersted field stabilizes transverse of the Py layer for which vortex walls are more stable. Once the chirality of the wall is compatible with the Oersted field, the domain wall moves with current pulses without changing configuration, indicating that the Walker breakdown is suppressed. This phenomenon could therefore explain the large domain wall velocities that we have observed in Py/Cu/Co nanostripes.

#### SH04

#### Domain-wall motion in permalloy nanowires with magnetic soft spots

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Recent concepts of nonvolatile data-storage devices involve the controlled motion of magnetic domain walls (DWs) in nanowires [1]. To realize such a device, a manageable fabrication process of reproducible and reliable pinning sites for individual DWs is required. We use irradiation with chromium ions to create local confining potentials via an effective reduction of the saturation magnetization in lithographically predefined regions. Fielddriven pinning and depinning at the so-called magnetic soft spots is directly observed using x-ray microscopy [2]. The shape of the potential is characterized via micromagnetic simulations and electrical measurements of the anisotropic magnetoresistance. Moreover, we demonstrate reliable DW depinning by single current pulses in a permalloy nanowire containing a square-shaped magnetic soft spot [3]. A DW can be moved back and forth between different soft spots along a wire. Lower requirements on the resolution of the lithography equipment in comparison to geometric constrictions on the nanoscale, a small distribution of properties due to parallel processing of many pinning sites during implantation, and fine tunability of the strength of the pinning potential via the chromium ion fluence make the magnetic soft spots a promising candidate for applications.

[1] D. A. Allwood, Gang Xiang, M. D. Cooke, C. C. Faulkner, D. Atkinson, N. Vernier, R. P. Cowhum, Science 296, 2003 (2002); S. S. P. Parkin, U. S. Patent No. US 683 400 5 (2004), IJ A. Vogel, S. Wintz, J. Kimling, M. Bolte, T. Strache, M. Arinsche, M.-Y. Im, P. Fischer, G. Meier, and J. Fassbender, IEEE Trans. Magn. 46, 1708 (2010), IJ A. Vogel, S. Wintz, T. Gerhardt, L. Bocklage, T. Strache, M.-Y. Im, P. Fischer, J. Fassbender, J. McCord, and G. Meier, Appl. Phys. Lett. 89, 202501 (2011).

### SH05

#### Direct observation of nearly mass-less domain walls in nanostrines with perpendicular magnetic anisotropy

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Important inertial effects have been recently shown for domain walls in permallov nanostripes. These transient effects, giving rise to an effective domain wall (DW) mass, are caused by deformations of the DW internal structure when a current pulse is applied. We have used time-resolved photoemission electron microscopy combined with x-ray magnetic circular dichroism to study current-induced DW motion in Pt/Co/AlOx nanostripes with perpendicular anisotropy and Rashba spin-orbit coupling. We show that in these nanostripes the DWs move at constant speed during the current pulse and inertial effects are much smaller than in permalloy nanostripes, i.e. the delay of the linear motion with respect to the pulse is less than 1 ns and the transient motion is smaller than 30 nm. The transient displacement δq depends of the change of generalized DW angle φ: δq = -Δ/α δφ where Δ the domain wall width at rest and  $\alpha$  the damping parameter. The negligible inertia in our system is accounted for by the narrow DW width and the small DW angle due to the presence of a transverse magnetic field of Rashba origin. Such small inertial effects could be efficiently exploited in devices based on manipulation of DWs.

#### SH06

#### Modified phase diagram of domain walls in FeNi/Cu/Co nanostripes

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Magnetic domain walls (DWs) confined in nanostructured magnetic materials display a wealth of specific static and dynamic physical properties. For magnetic nanostripes with sub-micron widths and in-plane magnetization, head-to-head DWs are either of transverse (TW) or vortex (VW) type and are characterized by a chirality, asymmetry and/or polarity. Due to their small size and fast dynamics, magnetic domain walls are promising candidates in the area of information processing and storage. The dynamic properties of these DW under application of current pulses through the spin transfer torque (STT) is one of the exciting subjects of spintronics. We used numerical simulations and analytical modelling to predict the phase diagram of transverse versus vortex walls in Py/Cu/Co trilayer nanostripes, in which high current-induced DW velocities have been observed. The DW is located in the Pv laver while the Co laver is initialized as single domain. We have found that due to magnetostatic interactions between Py and Co magnetizations, the TWs are energetically significantly more favourable in the Pv laver of the trilaver nanostripes than in single Pv nanostripes. We have experimentally verified part of this modified phase diagram with different microscopy techniques, including Magnetic Force Microscopy and Photoelectron Emission Microscopy.

#### SH07

#### Effect of current on a threshold width for a dimensional transition of domain wall dynamics in Co/Ni

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The effect of current on the dimensional transition of domain wall (DW) motion is investigated in Co/Ni nanowires by accompanying magnetoresistance measurement. As reducing the nanowire width, we found that the domain wall (DW) motion is changed from stripe pattern expansion to purely wall motion, implying that the dimensional transition from 2D to 1D. The threshold width for this transition is found to strongly depend on the relative direction of magnetic field and current on DW: parallel (antiparallel) direction has much smaller (larger) transition width, which is mainly because of different forces mechanism by magnetic field and current. The present work provides the upper boundary for domain-wall-motion based devices applications.

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### **SH08**

#### Simulations of field driven domain wall motion in permalloy nanowires with difference dimension

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In recent years, the domain wall motion in permalloy nanowire has been intensively investigated because it could give promise of a new type memory with many advantages. In our researches, the domain wall motion in permallov nanowires driver by external magnetic field have been simulated by a commercial software LLG simulation. The simulated results presented that the velocity of DW motion increases below a critical field value which is called Walker breakdown field: above that field the velocity of domain wall obviously decrease. The velocity of the domain wall and the critical field is dependent on the dimension of the nanowire. The domain wall velocity increases as the wire width increases, and the critical field decreases as wire width increases with non-linear relation.

1. GEOFFREY S. D. BEACH\*, CORNELIU NISTOR, CARL KNUTSON, MAXIM TSOI AND JAMES L. ERSKINE, nature materials VOL 4 OCTOBER 2005 2, Andrew Kunz, IEEE Trans Mag. VOL. 42, NO. 10, OCTOBER 2006 3. Stuart S. P. Parkin, \* Masamitsu Hayashi, Luc Thomas, Science 320, 190 2008

#### **SH09**

#### Voltage control of magnetisation and magnetic domain configurations in magnetostrictive epitaxial Fe1-xGax thin films

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The control of ferromagnetic domains and domain walls by electrical means is a desirable objective for applications in many existing and proposed information storage[1,2] and logical processing[3] devices. One route to achieve this goal is via hybrid ferromagnet/piezoelectric transducer devices [4-6] in which a voltage applied to the transducer induces mechanical strain in the ferromagnet affecting the magnetic anisotropy through the inverse magnetostriction effect. We have implemented epitaxially grown thin films of Fe81Ga19 in such a device. Through a combination of magnetotransport measurements, magneto-optical Kerr effect (MOKE) and high resolution photoemission electron microscopy (PEEM) we demonstrate that our epitaxial films nossess an exceptionally large magnetostriction (enhanced over that previously observed in bulk samples[7]) and a strong cubic magnetocrystalline anisotropy. We exploit these properties to achieve voltage control of the magnetisation and magnetic domain walls, including the reconfiguration of ordered domain patterns in structured geometries, and non-volatile switching of the magnetisation in the absence of external magnetic fields. Such functionalities will be directly applicable to information storage and processing technologies.

[1] W.J. Gallagher and S.S.P. Parkin, IBM J. Res. & Dev. 1, 5 (2006). [2] S.S.P. Parkin, M. Hayashi, and L. Thomas, Science 320, 190 (2008), [3] D.A. Allwood, et al. Science 309, 1688 (2005). [4] Samg-Koog Kim et al., J. Mag. Mat. Res. 267, 127 (2003), [5] J-W Lee et al., Appl. Phys. Lett. 82, 2458 (2003), [6] A.W. Rushforth et al., Phys. Rev. B. 78, 085314 (2008), [7] A.E. Clark, et al. J. Appl. Phys. 93, 8621 (2003).

#### SH10

#### Restricted oscillation period effect in the domain wall propagation after walker breakdown

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Understanding and control of dynamic behaviors for magnetic domain wall (DW) in ferromagnetic nanowire has become one of important issues in realization of potential spintronic applications such as DW logic and memory devices. In this work, we have intensively explored antivortex DW dynamic behaviors in ferromagnetic nanowires after Walker breakdown by means of the mircomagnetic simulation [1]. The nanowire has 2000nm length. The wire width w varied from 25 to 100 nm, and the wire thickness t varied from 0.5 to 12 nm. In all simulation, the material parameter of Permalloy is used. The unit cell dimension is 2.5×2.5×t nm3 and the Gilbert damping constant  $\alpha = 0.01$ . We have found that the DW dynamic behavior in the ferromagnetic nanowire significantly depend on the wire thickness and width. Interestingly, the oscillation period of the DW propagation is not significantly change with variation of the wire thickness is only dependent the wire width. By detailed analysis, we propose that there could be a new mechanism to restrict the oscillation period of the DW propagation in the ferromagnetic nanowire after the Walker breakdown.

[1] M. J. Donahue and D. G. Porter, OOMMF User's Guide, from http://math.nist.gov/oommf (2002).

# SH11

#### Domain wall motion by thermal gradients in Fe/W(110)

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The Spin Seebeck effect recently observed [1] has opened new areas of research in magnetic materials. One consequence of the Spin Seebeck effect is the movement of domain walls in magnetic materials [2]. In our current research we have studied the behavior of a domain wall in the presence of a thermal gradient in a 2D system by using the atomistic spin dynamics approach present in the UppASD code [3, 5]. This is based in the Landau-Lifshitz-Gilbert equation, where the atomistic spins are described using a Heisenberg Hamiltonian constructed from parameters obtained from first principles calculations. Temperature effects are taken into account by Langevin dynamics [3-4]. The system chosen is a monolayer of Fe on W (110) which exhibits a large anisotropy while having a soft exchange [5], resulting in domain wall thickness of the order of~ 1 - 2 nanometers. By subjecting this material to a thermal gradient we are able to observe a temperature dependent movement of the domain wall as well as changes of the spatial magnetization profile of the system. We compare this temperature gradient driven motion with the spin transfer torque driven domain wall motion observed for the system when subjected to a current.

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#### SH12

#### Transverse domain wall motion in notched ferromagnetic nanowire by spin transfer torque

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We have investigated domain wall motion in notched permalloy nanowires driven by spin polarized current. We solve modified Landau-Lifshitz Gilbert equation including the effect of spin-transfer torque by micromagnetic simulation. The ferromagnetic ground states for the studied geometries show transverse domain walls located at different regions depending upon the position, shape, dimension and symmetry of the notch in the nanowire. We have studied the effect of the spin torque on domain-wall displacement and domain-wall velocity. We find that the domain wall velocity is higher for single notched nanowire as compared to multiple notches. For rectangular notches of dimension 10 nm x 5 nm, transverse domain wall is pinned at one edge of the notch but it gets depinned for current density  $\sim 10^8 \text{ A/m}^2$ . We observe the current induced domain wall oscillation followed by a damped periodic translational motion of the transverse wall along with a switching of transverse domain wall from '?' to 'V' like configuration and vice versa via head to head or tail to tail domain wall formation.

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### SH13

#### Domain wall configuration and magneto-transport properties in dual spin-valve with

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The effect of the direction of the applied magnetic field on the magneto-transport properties including domain wall (DW) configuration and magnetization reversal process has been studied with a synthetic antiferromagnet-based dual spin-valve (SAF-DSV) structure (i.e. the SV structure is doubled symmetrically with respect to the FM). We can tune the DW configuration and its reversal process from a vortex to a transverse type by changing the direction of applied magnetic field respect to the nano-scale constricted SAF-DSV. When the field is applied along the perpendicular direction to the nano-scale constricted SAF-DSV, the perpendicular magnetic moments are developed due to the transverse magnetization reversal process. This multi step switching process reflects the pinning and depinning of a DW at the nano-scale constriction. Our results also show an asymmetric depinning field. We demonstrate, if nano-scale constriction is asymmetric along its length, i.e. expansions from both sides of the neck into the two nanowires are not identical, and then an asymmetric energy barrier to domain wall propagation is formed. This is due to the difference in DW width, which leads to an asymmetry in the domain wall depinning forces.

# **SH14**

#### Dynamics of domain-wall oscillations in magnetic nanorings driven by circularly rotating fields

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The manipulation of magnetic domain walls in geometrical confinements such as nanostrips and nanorings has become essential in information device technologies [1]. In particular ring structures have attracted much attention because domain walls can be easily positioned and manipulated by the application of magnetic fields or currents [2-4]. In the present work, we studied possible transverse domain-wall (TW) oscillations with certain frequency ranges in magnetic nanorings driven by circularly rotating fields on the nanoring plane. We used a permalloy nanoring with an onion state that has a tail-to-tail and a head-to-head transverse wall. From the micromagnetic simulations, it was found that the frequency of the TW oscillations varies with the frequency and amplitude of the applied rotating field. Moreover, the frequency range of the persistent TW oscillations become wider for a stronger field amplitude, as far as the frequency range is below a certain critical frequency value. We also found a phase diagram of the characteristic TW oscillation dynamics on the plane of the field strength and frequency. We will present our interpretations of those simulation results based on analytical calculations.

Corresponding author: sangkoog@snu.ac.kr [1] S. S. P. Parkin et al., Science 320, 190 (2008). [2] J. Rothman et al., Phys. Rev. Lett. 86, 1098-1101 (2001). [3] M. Klaui et al., Phys. Rev. B 68, 134426 (2003). [4] Y. Hou et al., Appl. Phys. Lett 98. 042510 (2011). [5] This work was supported by the Basic Science Research Program through the National Research Foundation of Korea (NRF) funded by the Ministry of Education, Science and Technology (Grant No. 20110000441).

### SH15

#### Modeling field-induced transformations of domain walls in magnetic stripes

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# July 13 (Fri)

### **SH16**

#### Interaction between propagating spin-waves and domain walls on a ferromagnetic nanowire

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The recent discovery that propagating spin waves (SWs) move a domain wall has created a new possibility to manipulate magnetization [Dong-Soo, Han et al., Appl. Phys. Lett. 94, (2009) 112502]. Here, we numerically investigate the interaction between propagating spin waves and a transverse domain wall in a nanowire by using micromagnetic simulations. In order to understand the mechanisms that lead to domain wall motions, we calculate to domain wall velocity and the depinning fields for a pinned domain wall that is depinned in and against the direction of the spin wave propagation. We find that the physical origin of the spin wave induced domain wall motion strongly depends on the propagating spin wave frequency. At certain spin wave frequencies, transverse domain wall oscillations lead to transverse wall displacement by the spin waves, while at other frequencies, large reflection and effective momentum transfer are the main drivers of the spin wave induced domain wall motion.

### **SI01**

#### Spin oscillations in a free molecular magnet

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Applying an external acoustic wave, we study spin oscillations in a magnetic nanoparticle that is free to rotate about its anisotropy axis. Using Hamiltonian of a rotated two-state spin system, we have shown that superposition of spin and rotational states makes a crucial effect on spin oscillations which exhibit quantum beats of the magnetization. In order to study such a beat structure, we compute dynamics of the magnetization by employing a perturbative approach, and discuss conditions under which this novel quantum effect can be generated. The results are expected to be tested in existing experimental techniques.

### **SI02**

### Effects of nonlinear spin dynamics on spin pumping

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Spin pumping has garnered significant interest in recent times. Most of the studies of spin pumping have been on systems in which the magnetization dynamics was linear. In nonlinear dynamics, however, large-angle, nonlinear, chaotic dynamics are regularly produced. We explore the nature of spin pumping signals resulting from nonlinear magnetization dynamical phenomenon. The device geometry comprises of a permalloy (Py) layer in contact with a nonmagnetic Pt layer. Magnetization dynamics is produced in the system by an rf signal applied to a coplanar waveguide patterned on top. Spin injected into the Pt is converted, due to the inverse spin Hall effect (ISHE), to a charge current and has been measured across the Pt layer. We will show that the system under study shown some high peaks at certain frequencies. The frequencies at which the peaks occur have been shown to be non-stochastic, even though one would expect them to be stochastic. Also, a frequency shift in the precessional resonant frequency has been observed as a function of the applied power. A similar shift had been previously observed in spin wave nonlinear dynamics in a different configuration. The origin of the peaks and frequency shifts are discussed.

#### **SI03**

# Sharp spectral linewidth in spin torque oscillator with perpendicular magnetized Co/Pd free layer

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There has been some scientific and practical interest in spin torque oscillator (STO) using perpendicular materials [1]. For application of STO to information-communication nano RF devices, it is important to investigate linewidth properties, although comprehensive understanding of linewidth properties in STO using perpendicular materials is still lacking. In this study, we have measured spin torque oscillation in giant magnetoresistance devices having a Co/Pd extended free layer and a NiFe/CoFe nanomagnet fixed bilayer. The total thickness of Co/Pd is 2.4 nm, and it was confirmed that Co/Pd has perpendicular magnetic anisotropy by vibrating sample magnetometer. Spin torque oscillation was measured under the low in-plane magnetic field using spectral analyzer. Microwave signal induced by spin transfer torque was clearly observed with oscillating frequency at 3.3 GHz. This frequency increased with increasing applied current i.e. positive nonlinear frequency shift. Obtained minimum spectral linewidth was 13 MHz, which is very small, compared with that of previous work on perpendicular materials [2]. Our results suggest that applying the low magnetic field to the Co/Pd film plane is efficient to obtain sharp spectral linewidth. This work was partly supported by Strategic Japanese-German cooperative program (ASPIMATT) from JST and by Japan Society for Promotion of Science.

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# **SI04**

#### Clocking schemes for soliton propagation in a ferromagneticallycoupled quantum-dot chain

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Transmission of magnetic signals along the ferromagnetically-coupled quantum-dot chain is an essential property to realize the magnetic quantum cellular automata. In the dot chain, a magnetic signal is represented by a topological soliton with a head-to-head or tail-to-tail state. The soliton is exceptionally stable coming from the topology of the system potential-energy surface, and could be driven by an external magnetic field. In addition to the external field, a global clocking field is used to further control propagation speed. We investigate the effect of the clocking field along the hard axis to initialize the soliton propagation by utilizing micromagnetic simulation. Interestingly, it is found that the required static field to initiate the soliton motion is a little larger than the minimum field to maintain the soliton motion, and only local trigger field is required to initiate the soliton propagation. Furthermore, we reveal the role of the global clocking field in the propagation speed, unexpectedly, the clocking field impedes the soliton motion.

#### SI05

# Spin wave propagation in single crystal Au(001)/Fe(001)/MgO(001) waveguides

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The needs for the development of future computing and information storage devices have led the emerging field of spintronics. Since spin waves can propagate macroscopic distance in ferromagnetic material, they can be used as information carrier. Typically spin wave propagations were investigated using gamet films or permalloy films, both of which have small crystallographic anisotropy. Here we investigated spin wave propagations using spin wave spectroscopy technique [1] in single crystal Fe waveguides with cubic and uniaxial surface anisotropies. The sample, Au(001)-50nm/ Fe(001) / MgO(001)-10nm, were fabricated on MgO(001) substrate using MBE system. We varied thickness of Fe (0.5-20 nm), waveguide width (1-100 µm), and distance of two antennae (0-40 µm). The static in-plane magnetic field was applied perpendicular to waveguide. In this way, magneto-static surface waves seme propagated along the ferromagnetic waveguides. The group velocity of spin wave was estimated from the oscillation period in transmission spectra. It was about 8.5 µm/ns in 20 nm thick waveguide under 500 Ce external magnetic field. We found that this value linearly decreases as decreasing the thickness. The contribution of magnetic anisotropies as functions of thickness and width of ferromagnetic layer to the spin waves also investigated.

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#### **SI06**

# Microscopic theory on the spin relaxation in an inhomogeneous spin dynamics

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The mechanisms of Gilbert damping in materials with inhomogeneous spin dynamics are more complicate than that with homogeneous spin dynamics (Kittel mode). It is well known that the magnetic impurities contribute to Gilbert damping with Kittel mode, on the other hand, the nonmagnetic impurities don't affect that. Recently, it has become clear that in the presence of spin waves (inhomogeneous dynamics), the nonmagnetic impurities contribute to Gilbert damping constant,  $\alpha$  [1]. In this case,  $\alpha$ is proportional to square of wave vector q in the limit q $\rightarrow$ 0. The aim of our study is to clarify the q-dependence of  $\alpha$  in the entire q range from microscopic theory. Our model is described by s-d model which consists of localized spins, and conducting electrons. In addition to this, we consider the nonmagnetic and magnetic impurities scattering. We obtain the precise expression of  $\alpha$  from the linear response theory. Consequently, we find that  $\alpha$  is proportional to 1/q in  $k\uparrow-k\downarrow<q<k\uparrow+k\downarrow$  ( $k\uparrow$ , $\downarrow$  are Fermi vector each spin) due to the Stoner excitation, and is continuously changed into the form of  $\alpha\theta+Aq^2$ in  $0<q<k\uparrow-k\downarrow$ .  $\alpha0$  originates from the magnetic impurities scattering while  $Aq^2$  is attributed to both the magnetic impurities scattering.

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# **SI07**

#### **Ultrafast transfer of spin in a non-collinearly magnetized multilayer** Koen Kuiper\*, Sjors Schellekens and Bert Koopmans *Eindhoven University of Technology, Netherlands*

Femtosecond laser pulses can trigger intriguing dynamic phenomena in magnetic systems such as quenching of magnetization within 100 fs. An effect hardly explored is that due to the laser heating high densities of hot electrons are generated, which can lead to transfer of spin angular momentum across different nano-layers at a femtosecond timescale. By controlling these spin currents, we effectively combine the field of spin-caloritronics with ultrafast magnetization dynamics, providing access to spin-caloritronics at ultimate timescales and a unique opportunity for the development of new spintronic devices. Malinowski et al. [1] showed in a timeresolved MO-Kerr effect experiment that spin-dependent transfer of hot electrons can speed up the demagnetization process. Here, in order to maximize the torque applied by these electrons, we explore laser induced dynamics in a SiOx/Pt/Co/Al(~1 nm)/Co/Al sample with a non-collinear orientation of the magnetization, i.e. an inplane magnetized top Co- and an out-of-plane bottom Co-layer. We provide a proofof-principle demonstration of laser-induced spin transfer across the Al spacer layer. The absorbed spin transfer and torque on the magnetization of the receiving layer cause an ultrafast canting of its magnetization, which is experimentally observed as a successively induced GHz-precession of the in-plane magnetized Co layer.

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#### SI08

#### Investigation of spin wave interference circuit with metallic thin film Nana Sato<sup>1</sup>, Koji Sekiguchi<sup>2</sup>\* and Yukio Nozaki<sup>3</sup>

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Spin waves are promising phenomenon for the future spintronic devices with low power consumption. For the realization of the spin wave device, we focused on the interference effect of spin waves[1]. Magnetostatic surface wave (MSSW) was excited in a NiFe thin film by applying a continuous RF field. In order to excite and detect MSSW, a pair of asymmetric coplanar wave guides (ACPWs) was fabricated on the NiFe thin film. The MSSW was excited by one ACPW, and its signal was detected by the other ACPW as an induced voltage in a real-time oscilloscope. The advantage of continuous excitation lies in that the MSSW signal becomes stable in time and space, enabling us to sensitively detect the spin wave interference. For the observation of spin wave interference, another strip line was inserted in the middle position between the ACPWs. From ACPWs in both sides of the inserted strip line two different MSSWs having opposite wave vectors were excited. By changing the phase of RF fields, the phase difference between two MSSWs was controlled. Induced voltage signal at the inserted strip line was changed according to the phase difference, indicating that we can observe the interference of MSSWs.

[1] M. P. Kostylev et al., Appl. Phys. Lett. 87, 153501 (2005)

#### SI09

#### Current-induced magnetization dynamics of synthetic antiferromagnetic free layers

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The spin polarized current-induced magnetization dynamics is studied theoretically in the frame of the driftdiffusion model [1-3] for two different spin-valve (SV) structures: (i) single free layer (SSV) Pol/Cu/FM and (ii) synthetic anti-ferromagnetic free layer (SyAF-SV) Pol/Cu/FMRu/FM. In a first step we have analyzed the effect of the angular dependence of the spin transfer torque (STT) on the magnetization dynamics for SSVs. For this four SSVs were considered: SSV1 (Co(3-5)/Cu(3)/Co(3)), SSV2 (Co(3-5)/Cu(3)/Py(3)), and SSV4 (Co(3-5)/Ru(3)/Py(3)). The polarizing layer, Co(3-5), is fixed. We found that these four different structures have different magnitudes and angular dependences of STT leading to different current - magnetic field phase diagrams with different critical values and ranges for dynamics. In the case of SyAF-SVs the drift-diffusion model allows us to take account of not only STT at Cu/FM but also STTs at FM/Ru and Ru/FM interfaces. The corresponding phase diagram is distinct from the one where STT across Ru has not been considered. Hence, this reveals that the STT through Ru considerably affects the dynamics of a synthetic free layer. We discuss the results as a function the strength of the Ruderman-Kittel-Kasuya-Yosida (RKKY)-type exchange interactions through Ru.

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### **SI10**

#### Magnon excitation studies in strongly correlated electron systems

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### SI11

#### Magnetoplasmonic hybrid nanoparticles

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Magnetoplasmonics is a very promising branch of active plasmonics that uses a magnetic field to modulate plasmon resonance; [1] the field-mediated modulation of the plasmonic response of nanoparticles is expected to have strong implications in light guiding applications on plasmon resonance-based sensing. [2,3] In this paper we show how ordinary gold nanoparticles' optical properties are modified by the magnetic field and describe the origin of such magnetoplasmonic behaviour. Then, using a combined approach with the aid of magneto-optics, magnetometry and x-ray spectroscopy we investigate several routes to increase the response to the magnetic field by adding a magnetic moiety to the gold nanoparticle. We found that when a transition metal oxide is added to the gold core in core@shell and heterodimer geometries, the two materials (magnetic and plasmonic) behave independently, thus giving rise to bifunctional magnetic-plasmonic nanoparticle hybrids. True magnetoplasmonic behaviour is instead found when a higher degree of hybridisation exists between the two moieties, as in the case when metallic magnetic metals are added to the gold moiety, thus achieving an enhancement of the magnetic response of the optical properties and giving rise to a synergic, non-additive spectroscopic signature in the magneto-optical spectrum.

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### SI12

# Inhomogeneous standing spin wave excited by the patterned periodic electrode

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There has been a great deal of attention in the spin waves excited in ferromagnetic films because of their various potential applications such as spin-wave-based logic circuit and the interconnection device between the electric and spintronic devices. Efficient excitation and manipulation of the spin wave are important issues for the realizations of practical device applications. Spin wave interference in three dimensional spin structures may enhance the amplitude of the spin oscillation. For this purpose, we here systematically investigated the thickness dependence of resonant frequency of the standing spin wave. We prepared Permalloy films with various thicknesses and excite the standing spin wave by using a periodic nonmagnetic Cu electrode. The resonant frequency of the spin wave was found to increase with the Permalloy thickness. This can be understood by the increase of the non-excited layer around the bottom of Permalloy layer, which induces the effective magnetic field to the excited layer via the exchange interaction. We also perform similar studies of the standing spin wave in the other ferromagnetic structures exchange biased and ferromagnetic/nonmagnetic multi-layered film and present the optimized structure for increasing the output signal induced by the spin wave excitation

### **SI13**

# Observation of spin-waves by time-resolved magneto-optic kerr effect microscope

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Spin-wave is an ordering of the magnetization when the magnetization of a ferromagnet is tilted away from its equilibrium state. The spin-waves in metallic thin films have been recently studied as an emerging field of magnetism for magnetic recording and spintronics devices. The spin-waves can be detected by either electrical or optical method for either frequency or time domain, respectively. In the frequency domain the ferromagnetic resonance has been commonly employed to detect the spin-waves, whereas in the time domain, it is typical to detect damped oscillation of magnetization excited by a short magnetic pulse via an inductive or optical apparatus. Here we report the experimental results on the spin-waves by use of all optical pumpprobe method. 20-nm-thick Co thin films are prepared on a GaAs substrate by dc magnetron sputtering. The pump beam is focused onto a spot with a diameter 1 µm of a sample under an applied magnetic field 2 kOe at normal direction of the film surface. The spin-waves are then observed by controlling the tilt angle of the probe beam with a fixed pump beam. Using this method, a propagating spin-wave can be measured directly and further analysis will be discussed.

### **SI14**

# Nuclear magnetic resonance study of proton dynamics in ZnO

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Hydrogen acting as a dopant is widely known in compound semiconductors including many oxides. However, in order for such doping to be technologically relevant, the thermal stability of the H-donor must be enough to avoid degradation during device operation at and above room temperature. Here, we have carried out 1H NMR measurements in order to understand proton motions in the ZnO lattice at temperatures between 200 K and 400 K. In our post-annealed ZnO, there are two proton species, one of which can be assigned to protons at interstitial sites (H<sup>-</sup>) in ZnO. The other species shows no appreciable motion in the temperature range. The distinct activation energies obtained for the Hi<sup>+</sup> motion, 0.27 eV and 0.42 eV, are associated with the reorientation around the bonding oxygen atom and hopping between neighboring oxygen atoms, respectively. These activation energies for diffusion, obtained from the spin-spin relaxation measurements, are in close agreement with those obtained from ab initio calculations.

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### **SI15**

# Micromagnetic study of magnonic band gaps in waveguides with a periodic variation of the saturation magnetization

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Spin wave propagation in micro-sized magnonic crystals (MCs) is intensively studied due to their potential technological application for signal processing in spintronic devices. Here we report on micromagnetic simulations [1] of the spin wave propagation in a MC realized as a permalloy ( $N_{iso}Fe_{20}$ ) waveguide with a periodical variation of its saturation magnetization. In real structures the variation of magnetization can be achieved by using an ion implantation technique. The 2 µm-wide waveguide of 40 nm thickness is magnetized transversal to its long axis. The MC lattice constant is equal to 1 µm. The spin-wave transmission characteristics have been studied as a function of the width of the implanted areas and of the level of the magnetization variation M/ M0. Frequency band gaps were clearly observed in the spin-wave transmission spectra. The dependences of the depth, width and the position in frequency and space of the rejection band gaps on the above parameters are referred in our studies. The role of the higher order spin-wave width modes on the MC properties is discussed as well. Support from DFG (grant SE-1771/1-2) is gratefully acknowledged.

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### SI16

#### Spin-transfer induced spin waves of a magnetic point contact with a confined domain wall

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Spin-transfer induced microwave oscillation will be a powerful candidate for applications in future wireless telecommunication technologies. According to Ref. [1], the oscillation of a domain wall between the Bloch- and Neel-type magnetic configurations is induced by a direct current and the corresponding oscillation frequency is proportional to the current density. It is also known that the current injected through a point contact can excite a variety of spin wave modes [2-4]. It is intriguing to ask if such spin wave modes can be excited in a thin magnetic layer connected to a magnetic point contact containing a domain wall. From our calculation we found that three different kinds of spin wave modes can be excited in the thin film depending on the bias current density. The low frequency mode just above the threshold is the propagating spin wave mode. Then the oscillation frequency increases drastically with increasing the current density and reaches more than 100 GHz. At high current density the excited spin wave mode is a propagating mode with a large angle precession. We also found that applied field independence of the frequency for the second mode. We shall discuss the physics behind the excitation of spin wave modes.

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#### SI17

#### Spin-torque-nano-oscillator using the perpendicular magnetized CoFeB/MgO/CoFeB magnetic tunnel junctions

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Perpendicular magnetized free layer in the magnetic tunnel junctions (p-MTJs) has advantages for spintorque-nano-oscillator because coherent precession, high output power, and increasing precession frequency due to high magnetocrystalline anisotropy, can be expected. In this study, the microwave transmission properties were investigated for the p-MTJs using thin CoFeB. The multilayer was prepared by UHV sputtering system and patterned it into 40-nm-sized junctions with coplanar waveguide electrode. The microwave properties were measured under the in-plane magnetic field of 3 kOe to tilt the free layer. The tunnel magnetoresistance ratio was 70% and 0.9% for the out-of-plane and in-plane direction, respectively. Below -0.24 mA, the spin torque noise was observed with broad signal. The STO signal was clearly observed above -0.30 mA at 6 GHz, and the second harmonic signal was not observed due to the single-precessionmode by small junction size. The linewidth and output power was 650 MHz and 1. nW, respectively. Power has a potential to be more enhanced by stabilizing the pinned layer to increase the relative angle between free and pinned layer. This study was partly supported by ASPIMATT (JST), by Grant-in-Aid for Young Scientists A (No. 22686001), by FIRST program (JSPS), and by Maekawa-Houonkai foundation.

#### SI18

#### Planar approximation for spin-transfer devices with tilted polarizer Ya. B. Bazaliv\*

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Planar spin-transfer devices with dominating easy-plane anisotropy can be described by an effective one-dimensional equation for the in-plane azimuthal angle [1-3]. Such a description maps the Landau-Lifshitz-Gilbert equation on the Newton equation for the motion of classical particle in a one-dimensional potential and thus provides an intuitive understanding of the magnetic dynamics, allowing one to obtain qualitative results without performing detailed calculations. We apply the effective planar equation to describe magnetic switching and precession states in spin transfer devices with tilted polarizer [4]. The approach allows one to list the possible dynamic regimes, including the precession cycles, and sketch the switching diagram of the device.

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# SJ01

#### Fast SpinRAM simulation by GPU Kiyoaki Oomaru\* and Yoshinobu Nakatani

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Spin transfer torque driven magnetoresistive random access memory (SpinRAM) is one of the candidates of the next generation non volatile memory, and it is studied intensively. Micromagnetic simulation is one of the methods to study it. In the SpinRAM simulation, because the size of the MTJ element is small, the number of the calculation cells is small, and the calculation for one simulation takes only a few minutes. However in the real analysis, more than millions of simulations are repeated with different conditions, the total simulation time becomes huge. Speed up of the simulation is important. Recently, a Graphic Processing Unit (GPU) is used to accelerate the micromagnetic simulation. The acceleration rate strongly depends on the number of the calculation cells, and a huge number of the calculation cells is required to achieve high speed up ratio. Speed up can not be expected in SpinRAM simulation by GPU. In this paper, we propose the simulation method to calculate a number of the MTJ elements in parallel by GPU, and achieve 75 times of the speed up ratio in maximum. This study was supported by New Energy and Industrial Technology Development organization (NEDO) partly.

#### SJ02

#### Micromagnetic simulations for the spin dynamics and Gilbert damping constants in nano-dot with perpendicular magnetic anisotropy Jungbum Yoon<sup>1</sup>, Chun-yeol You<sup>1\*</sup> and Myung-hwa Jung<sup>2</sup> <sup>1</sup> Department of Physics, Inha Universirty, Korea <sup>2</sup> Department of Physics, Sogang Universirty, Korea

Spin dynamics and the magnetic damping of a nano-dot ellipse with perpendicular magnetic anisotropy (PMA) are presented by the micromagnetic simulations. While the in-plane static magnetic field is applied to the long axis of the ellipse, the small in-plane RF magnetic field is applied to the short axis of the ellipse to mimic the ferromagnetic resonance (FMR) technique. The susceptibility can be obtained by the power spectra of fast Fourier transform (FFT) from magnetizations dynamics. There are two regions in the external static field dependence resonance frequencies for typical PMA system. (I) With small external field, the magnetization is tilted to the static field direction with the finite angle from the film normal. (II) If the static field is strong enough, the magnetization is aligned to the static field, where the Zeeman energy overcomes the PMA energy. In this simulation, the extract Gilbert damping constants from the line widths of the resonance peaks from two different regions. Surprisingly enough, the obtained values ( $\alpha_{\perp} I = 0.0165$  and  $\alpha_{\perp} II = 0.0194$ ) are different from the input value ( $\alpha = 0.027$ ). It implies that the experimentally obtained  $\alpha$  can be different from the real values.

#### SJ03

#### Micromagnetic study on micro-structured ferromagnetic thin film for high frequency device applications

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The ferromagnetic metallic thin films of Fe has been considered as promising materials for on-chip microwave devices, owing to the much higher saturation magnetization Ms compared with a conventional material of yttrium iron garnet (YIG). Understanding of demagnetizing effects, related to Ms, is important for the design of micro-structured devices with the high Ms material[1]. In the present study, geometrical effects of micro-structured magnetic thin films on high frequency magnetic properties are numerically investigated with micromagnetic simulations. The standard material parameters of Fe with cubic anisotropy are assumed in the simulation. The values of  $\mu'$  (@ 1 GHz) and the resonance frequency fFMR for various aspect ratios w/t are compared. The increase of  $\mu'$  with the decrease of w/t is attributable to the size effect in fFMR, related to the demagnetizing coefficient along the film normal direction. It should be noticed that the  $\mu'$  for the same w/t with different w and t is identical. The results suggest that the demagnetizing effect on the high frequency magnetic properties becomes prominent for the wit smaller than 400 in the structured Fe thin film.

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### **SJ04**

#### Atomistic modelling of magnetization dynamics with spin torque

Phanwadee Chureemart, Richard F. L. Evans, Irene D'amico and Roy W. Chantrell Physics, University of York, United Kingdom

The dynamics of the magnetization following the introduction of a spin current is studied by using an atomistic model coupled with spin accumulation. The evolution of the magnetization and spin accumulation in a trilayer system are calculated selfconsistently, the interaction between magnetisation and spin current provided by an s-d exchange term. A spin-polarised current is introduced into the system containing a domain wall whose width is varied by changing the anisotropy constant. The magnetisation is calculated using an atomistic model based on the Heisenberg exchange energy and spin dynamics introduced via the Landau-Lifshitz-Gilbert equation. The net spin torque contributes two important components: the adiabatic (AST) and nonadiabatic torques (NAST). We show that both arise naturally when calculated directly from the spin accumulation. The maximum value of the AST and NAST are considered for each thickness of domain wall. Both torques decrease as the domain wall thickness increases. The nonadiabaticity factor, defined as the ratio of the NAST and AST is also shown. Increasing the domain wall thickness gives rise to a small angle between magnetisation in domain wall. Interestingly the degree of NAST tends to decay to zero as the thickness of DW increases relative to the spin diffusion length.

### SJ05

# Highly parallelized micromagnetic simulator using fast multipole method

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We have developed an alternative method for computing the magnetostatic field of arbitrary-shaped ferromagnetic systems instead of fast Fourier transform (FFT). The method is based on the Taylor expansion of the potential, as is required by the fast multipole method (FMM)[1][2], the most efficient method known for calculating mangnetostatic fields in very large systems. The FMM scales as O(N) in both time and space complexity compared to the direct method with O(N<sup>2</sup>), or FFT with O(N logN)[3] and even easier to be parallelized, which is a more valuable solution for larger ferromagnetic systems. Graphical processing units (GPU) are now increasingly viewed as data parallel compute coprocessors that can provide significant computational performance at low price[4]. We have developed a micromagnetic simulator using the parallelized method with GPU and have tested in various configurations. In that case speedups of over a factor 100X can easily be obtained compared to the CPU-based OOMMF[5] program developed at NIST. The tests are carried out with µMag Standard Problem[6] #1,#2,#3,#4.

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### **SJ06**

# Effect of calculation conditions on the numerical simulation of magnetic materials with random magnetic anisotropy

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Numerical simulation was performed on magnetic materials with random anisotropy. An assembly of magnetically interacting grains with randomly oriented uniaxial anisotropy was provided and magnetization was calculated using the Landau-Lifshitz-Gilbert equation. The magnetization in a particular grain is assumed to align in the same direction (single spin model).[1] Calculations were carried out for N × N × N three dimensional cells changing cell sizes from 5 to 40 nm. The interaction at the interface between grains was taken into account by the interaction energy between the unit vectors representing grain magnetization directions. N was changed from 10 to 64. In the case of N = 10, the relation between coercive forces and grain sizes was obtained to be  $\delta = 5.7$  in He  $\sim D^{-}\delta$ , and fits to the primitive theory of random anisotropy model (RAM)  $\delta = 6$ . As the N increased, the  $\delta$  tends to decrease slightly from 5.7 to 4.2. The grain sizes where the coercive force becomes less dependent on the grain size, suggesting the natural exchange length, did not changed much.

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### **SJ07**

#### Vortex and antivortex formation in magnetic rolled-up nanotubes

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Magnetic nanotubes have attracted considerable attention owing to both fundamental and technological interest1-2. Using a rolling-up fabrication technology3, patterned multilayer tubes can be prepared that exhibit significant modifications of magnetic behaviors compared to those of thin-film structures4-5. In order to telucidate the underlying physics, we conducted micromagnetic simulations of possible spin configurations in rolled-up nanotubes. When a permalloy platelet of dimensions 400 x 400 x 5 nm3 is rolled up, various spin configurations are found to be either C-state, vortex(V) state or anti-vortex(AV) state in the plane of thin films with more spins being forced to be aligned along the tube axis. Therefore, the region whose spins are aligned transversely to the tube axis is compressed. The modified spin alignment is supposed to increase the anisotropy magnetoresistance ratio. Although the reorientation of local magnetizations in the rolled-up tubes results in increases of the exchange energy, it reduces much more the magnetostatic energy, particularly for AV states of the rolled-up tubes show another interesting feature: the stray fields inside the tubes are cancelled out at the positions.

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#### **SJ08**

# Magneto resistance study using micro magnetic simulations in permalloy nano ladder

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Micromagnetic simulation is one of the tools to investigate magnetization reversal mechanism in nano-wires and nano structures. In the open source simulators like OOMMF[1], magneto transport studies viz., Anisotropic Magnetoresistance (AMR) is not available and is non-trivial since it involves solving the generalized Laplacian equation for getting the spatial distribution of current density[2]. In the case of nano wires, the current density is considered to be uniform[3]; in this quasi uniform state one can apply phenomenological expression in deriving the AMR. We apply this condition in OOMMF-simulations and is compared with the commercial software "LLG-simulator"[4]. For this purpose, we considered single wire of length(L) 1000nm, width(W) 100nm and thickness(t) 5nm. Both the simulators agree with each other within a maximum error of 5%. Further, we studied a ladder structure formed by joining these two wires, with interconnects of L=100nm, W=20nm and t=5nm. In this ladder structure, we varied number of steps by keeping the periodic separation as 50nm. It was found that the coercivity falls non-linearly from 312 to 255G when the number of steps increased from 1 to 13. The resistance measured at remanence states drops linearly, which indicates the decrease in remanence due to the local nucleation centers.

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#### **SJ09**

# Soft layer driven switching of microwave-assisted magnetic recording on segmented perpendicular media

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In this work, we use micromagnetic modeling to investigate microwave-assisted magnetic recording on segmented perpendicular media. Two types of microwave fields, sinusoidal and finite bandwidth square microwaves, have been chosen for our study. We show that microwave fields, when applied to the soft segment, can assist the magnetization reversal of the hard segment. The calculated assisted frequencies are insensitive to the variation of inter-segment exchange coupling of segmented perpendicular media. Our results suggest that sinusoidal microwave fields are practically more feasible than square microwave fields for the soft layer driven switching mechanism of segmented perpendicular media.

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### SK01

# Top composite free layer non-collinear spin valve for hysteresis-free GMR sensors

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Top spin valves Ta/[FeNi/CoFe]/Cu/CoFe/FeMn/Ta based on FeMn antiferromagnet with varying [Fe<sub>20</sub>Ni<sub>80</sub>/Co<sub>90</sub>Fe<sub>10</sub>] composite free layer parameters and Cu interlayer thickness were prepared by DC magnetron sputtering. GMR characteristics of spin valves (SV) were investigated under the conditions of non-collinear SV geometry when the applied magnetic field and pinning direction for bottom magnetic layer are not collinear. It was shown that the value of low field hysteresis caused by free layer magnetization reversal can be reduced down to 0.1 Oe keeping the GMR ratio about 10 % by using both layers thicknesses and non-collinearity angle variation. The magnetoresistive sensitivity of non-collinear is about 1 %/Oe. Low-hysteresis free layer magnetization reversal is mainly due to coherent rotation of magnetization. The dependences of low field hysteresis value and GMR ratio on the Cu spacer layer thickness and the angle between applied magnetic field and pinning direction are presented.

#### SK02

Strong <111> texture and low hysteresis in MnIr-based top spin valve Mikhail Milyaev, Larisa Naumova, Vyacheslav Proglyado, Tatiana Krinitsina, Nataliya Bannikova and Vladimir Ustinov Institute of Metal Physics UB RAS, Russia

"Spin valve" nanostructures with GMR and TMR are widely used in spintronics. In some devices the hysteresis of magnetoresistance must be negligible. The hysteresis is weak in the spin valves with <111> axial texture. In this work we study the dependence of the free layer hysteresis loop width on the texture strength that is characterized by the FWHM of rocking curve. The top spin valves [Ta, (Ni<sub>80</sub>Fe<sub>30</sub>/<sub>b0</sub>Cr<sub>40</sub>)]/Ni<sub>80</sub>Fe<sub>30</sub>/Co<sub>30</sub>Fe<sub>10</sub>/ Cu/Co<sub>30</sub>Fe<sub>10</sub>/Mn<sub>75</sub>Ir<sub>25</sub>/Ta with various thickness of the magnetic and nonmagnetic layers are investigated. On the base of magnetor, magnetoresistance and XRD data the nonlinear dependence between FWHM of rocking curves measured at the integrated (111) Bragg peak of Ni<sub>80</sub>Fe<sub>20</sub>, Co<sub>30</sub>Fe<sub>10</sub> and Cu layers and free layer coercivity is established. A correlation between GMR ratio and FWHM is not revealed. It is shown that in the spin valve Ta(50A)/NiFe(30A)/CoFe(20A)/Cu(22A)/CoFe(25A)/MnIr(50A)/Ta(20A) with sharp <111> texture, GMR ratio of 10.6 % and minimal coercivity of 8.7 Oe, the free layer hysteresis can be effectively reduced down to few tenth of Oersted [1]. This study is supported by RFBR, project No. 10-02-00590, by the Program of Presidium RAS, project No. 12-2-21051, and OFI-UB RAS, project No. 11-2-23-NPO.

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#### **SK03**

# Low temperature crystallization process in Co<sub>2</sub>FeSi Heusler alloy thin films

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One of the major challenges for Heusler alloys is to achieve a high  $L_{21}$  ordering at low annealing temperatures in order to limit interdiffusion in multilayer devices. We report a low temperature crystallization process of polycrystalline Co<sub>2</sub>FeSi films using in-situ imaging techniques. 20nm thick Co<sub>2</sub>FeSi films were annealed in-situ up to 600°C with both TEM and diffraction patterns recorded. We show that the grain crystallization starts to occur at 230°C with no additional nucleation occurring after 30 minutes but further annealing does improve the grain ordering. By varying the growth conditions it is possible to change the number of nucleation sites and also the median grain size. Through varying the initial film morphology it is possible to change the number of nucleation sites and also the median grain size from 20 to 250nm. Magnetic measurements show that the magnetic moment is over 80% of the theoretical value predicted [1]. This shows that a high degree of ordering can be created for polycrystalline films after a low temperature anneal of less than 300°C.

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# SK04

# Spin-polarized itinerant electrons in Co<sub>2</sub>MnAl and Co<sub>2</sub>MnSi studied by magnetic Compton scattering

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Investigating the spin-polarized itinerant electron states in the intermetallic alloys and compounds is indispensable for the development of potential spintronic materials. In this work, we have studied the relationship between the number of sp itinerant electrons and the spin moment in the compounds Co<sub>2</sub>MnZ (Z=Al or Si), using synchrotron-based Magnetic Compton scattering (MCS)[1], which provide the bulk spin-polarized electron momentum distributions, i.e. magnetic Compton profiles (MCPs). The shape of MCPs depends on the number of spin-polarized electrons of occupied orbitals, and the area of MCP corresponds to the spin moment. Co<sub>2</sub>MnZ (Z=Al or Si) is a good system to investigate the spin-polarized itinerant electron states since the number of sp itinerant electrons can be changed without affecting the number of 3d electrons. The following results were obtained . 1. The total spin moment increases by 1 µ<sub>b</sub>/L. from Co<sub>2</sub>MnA1 (4µ<sub>b</sub>/Lu) to Co2MnSi (Syu<sub>b</sub>/Lu), with one sp electron difference. 2: The proportion of the partial spin moment carried by sp itinerant electrons is same for Co<sub>2</sub>MnA1 and Co<sub>2</sub>MnSi. From these findings, additional sp electrons occupy the 3d electron states and increases the localized spin moment. This enhances the spin polarization of itinerant electrons through the indirect magnetic exchange interaction.

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#### SK05

# Evidence of inelastic tunneling in magnetic tunnel junctions via capacitance-voltage characteristics

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Magnetic tunnel junctions (MTJ) have played a key role in the development of HDDs as well as next generation non-volatile memories. For high speed applications, capacitance and RC time constant analysis of MTJs becomes very important. Also for metal-insulator-metal structures with extremely thin tunnel burniers as in MTJs, the capacitance is mainly determined by the interface properties of the structure as the electric field penetration becomes significant. Since in MTJs the interfaces are very critical in determining the tunneling magnetoresistance (TMR), study of capacitance becomes very relevant for better understanding of these devices. We report the temperature dependence of C-V characteristics in MTJs. All our MTJs show a negative tunneling magnetocapacitance (TMC) with a magnitude lower than the TMR [1]. At temperatures lower than 100 K, a zero-bias anomaly appears in the C-V profile. The behavior is similar to inelastic tunneling or zero-bias anomaly in MTJs. Zero-bias anomaly in the R-V characteristics of MTJs has been attributed to magnon and phonon assisted tunneling or thorough impurities inside the barrier. But since the capacitance is mainly contributed from the interfaces, the origin of these effects is from the interfaces rather than the bulk.

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### SK06

#### Spin-polarization measurements for Co<sub>2</sub>MnSi using Co<sub>2</sub>MnSi/MgO/ NbN epitaxial tunnel junctions

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Spin-polarization measurements were carried out for the Co-based Heusler alloy, Co<sub>2</sub>MnSi (CMS, for short), using CMS/MgO/NbN epitaxial tunnel junctions. The layer structure was as follows: (from the substrate side) MgO buffer (10 nm)/NbN (10 nm)/MgO barrier (1.4-3.2 nm)/CMS (3 nm)/Co<sub>38</sub>Fe<sub>50</sub> (1nm)/IrMn (10nm)/Ru cap (5 nm). All the samples were successively deposited on MgO(001) single-crystal substrates in an ultrahigh vacuum chamber. Spin-polarization measurements were carried out at 1.5 K using the conventional lock-in technique. The spin polarizations (P) were evaluated for three different film compositions of CMS:  $Co_2Mn_{0.7}Si_{0.81}$ ,  $Co_2Mn_{1.40}Si_{0.81}$ , and  $Co_2Mn_{1.40}Si_{0.81}$ . At H = 0 T, typical differential conductance curves were clearly observed, and were well-fitted by the modified-BTK theory [1]. The superconducting energy gap extracted from the analysis was about 2.5 meV [2]. The P values evaluated at H = 8 T were 0.41, 0.50, and 0.54 for respective CMS films. Thus, the P increase of the TMR ratios for CMS/MgO/CMS fully-epitaxial magnetic tunnel junctions, and can be qualitatively understood from the view point of the suppression of minority-spin in-gap states that resulted from the reduction of COMn anti-sites [3].

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# SK07

# Variation of point-contact and reev reflection spectra of ferromagnetically ordered metals

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Point-contact Andreev reflection (AR) spectra of ferromagnets are usually fitted with modified BTK models that take into account the spin polarization. On the other hand, AR spectra of non-magnetic metals are often described by a BTK model that includes Dynes' lifetime parameter. The two models yield rather similar spectra and can in fact be used to fit the spectra of both magnets and non-magnets. The differences between them are most pronounced at low temperatures but strongly suppressed by thermal smearing. We have applied the two models to point contacts between the classical band ferromagnets Co, Fe, and Ni and the conventional superconductors Nb, Ta, and AI as function of temperature. The spectra varied between an absent superconducting signal and an occasional Josephson-like anomaly with a large resistance drop at zero bias, while many spectra had the typical Andreev reflection double-minimum structure. However, we could reproduce the superconducting gap only in a few cases. This unexpected variation could be related to the interplay of ferromagnetism and superconductivity at the interfaces.

### **SK08**

#### Fabrication of Highly Sensitive Magnetic Tunnel Junctions for Biomagnetic Field Sensor Application

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The discovery of large tunnel magneto-resistance (TMR) effect of over 200% in MgO based magnetic tunnel junctions (MTJs)[1,2] enabled us to design highly sensitive magnetic field sensors[3,4], such as a bio-magnetic field sensor. For this purpose, we need MTJs with a high sensitivity and hysteresis-free responses. In this study, we fabricated MTJs with a NiFe<sub>n</sub>/ Ru<sub>0.9</sub>/ CoFeB<sub>3</sub> (in nm) bottom synthetic ferri-coupled free layer and an MgO barrier layer. This thick NiFe layer serves low magnetic anisotropy[5]. The MTJs were annealed at 325°C by applying a magnetic field. After the first annealing, we rotated direction of the magnetic field along the hard axis of the free layer, and we have successfully observed the very high sensitivity of 25.3%/Oe and hysteresis free response. Developed MTJs are promising for bio-magnetic sensor application. This work was partly supported by S-Innovation program, Japan Science and Technology Agency (JST).

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### SK09

Oscillatory exchange coupling and strong magnetoresistance effect in Fe/AgX/Fe (001) heterostructures with X=Cl and Br  $\,$ 

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AgCl and AgBr silver halides crystallize under ambient conditions in rock salt geometry having lattice parameters epitaxially compatible with Fe structure. AgCl and AgBr compounds are wide bang gap semiconductors with indirect band gaps of 3.2 eV and 2.69 eV, respectively. Thus Fe/AgX/Fe (001) magnetic tunnel junctions (MTJs) with X=Cl and Br may represent feasible heterostructures for comparative studies with sp based Fe/MgO/Fe and Fe/NaCl/Fe junctions as well as for spintronic applications. In present contribution structural, electronic and magnetic properties of Fe/AgX/Fe (001) MTJs are theoretically investigated by means of a self-consistent Green's function technique for surface and interfaces implemented within tight-binding linear muffin-tin orbital (TB-LMTO) method. The spin-dependent transport properties are determined by means of Kubo approach implemented within TB-LMTO formalism. Formation of sharp Fe/AgCl(AgBr) interfaces and enhancement of interfacial Fe magnetic moments are predicted. Metal induced gap states on both anion and cation sites are observed. Oscillatory exchange couplings between ferro- and antiferromagnetic states with decreasing amplitude are evidenced. Spin-dependent transport properties are determined by the complex band structures of AgCl and AgBr barriers. Tunneling magnetoresistance ratios as high as 3500 % are predicted for Fe/AgCl/ Fe MTJs in asymptotic regime and even higher values for Fe/AgBr/Fe junctions.

### SK10

# Transport properties in double MgO barrier magnetic tunnel junctions with Fe nano-particles

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Interplay between spin dependent electron tunneling and Coulomb blockade effect [1,2] is one of the most attractive subjects in spintronics field because it is expected to give rise to remarkable tunnel magnetoresistance (TMR) effect as well as exhibit fascinating phenomena such as Kondo effect[3]. And in order to realize this interplay, double barrier magnetic tunnel junctions (MTJs) with ferromagnetic nano-paricles are good candidates. In this research, double barrier MTJs with structure of MgO(001) substrate/MgO/Cr/Fe/MgO/Fe particles(t = 0.12 - 1.37)/MgO/Fe/Au (unit in nanometer) are fabricated. An average diameter of the Fe nano-particles setimated by Langevin function fitting reveals an almost linear relation with t for t  $\leq$  1.02 nm. Coulomb blockade effect was observed for t  $\leq$  1.02 nm, and the t dependence of the Coulomb energy of the Fe nano-particles was explained consistently by a simple capacitance model. Interestingly, electron tunneling spectra (dI(V)) and TMR show a dip around zero bias at low temperatures. A suppression of the TMR around zero bias implies a phenomenon like Kondo physics. The transport properties in the double barrier MTJs having Fe nano-particles were also compared with those in a double barrier MTJ with the continuous middle Fe layer and a single barrier MTJ.

Acknow ledgements: This work is supported by Grant-in-Aid for Scientific Research (S), MEXT, Japan. (No. 23226001). [1]. S. Mitani et al. Phys. Rev. Lett. 81, 2799 (1998). [2]. H. Sukegawa et al. Phys. Rev. Lett 94, 068304 (2005). [3]. K. I. Lee et al., Phys. Rev. Lett. 98, 107202 (2007).

#### SK11

Characterisation of Epitaxial and Polycrystalline Co<sub>2</sub>FeSi thin films James Sagar<sup>1</sup>, Hiroaki Sukegawa<sup>2</sup>, Leonardo Lari<sup>3</sup>, Vlado K Lazarov<sup>3</sup>, Seiji Mitani<sup>4</sup> and Atsufumi Hirohata<sup>5</sup>\*

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We have grown and characterised both epitaxial and polycrystalline Co<sub>2</sub>FeSi thin films using UHV magnetron and HiTUS sputtering. The structure of the films has been studied using XRD, TEM and STEM. The combination of these techniques allows analysis of the structure and defects within the films. Magnetic characterisation has been carried out using AGFM and VSM. UHV magnetron sputtered films have been found to become increasingly L<sub>21</sub> ordered after annealing above 500°C. HiTUS deposited samples form a polycrystalline film from and amorphous Co, Fe and Si matrix after annealing at 500°C for 3, 6 and 9 hours with large (50nm-250nm) L<sub>21</sub> and B<sub>2</sub> ordered grains. This growth process results in films with large Hc (>500Oe) these values are two orders of magnitude greater than those found to improve with the use of a form Cr or Ag buffer layer. This is shown by a decrease is Hc and an increase in Ms to greater than 85% of the theoretical value.

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#### SK12

#### Inverse spin polarization at benzene/iron interface

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Organic spin-valves involving an organic spacer inserted between two magnetic electrodes have been successfully realized and led to sizable magnetoresistance ratios [1-3]. The accurate design of the interfacial electronic structure controlling the spin injection is critical to obtain high performance devices. Intriguingly, recent reports claimed that interfacial spin polarization may be reversed in some cases [4, 5]. We address the nature of the spin polarization Benzene/ Fe (100) interface using ab-initio approach. The surface is modeled by 6 layers of Fe (100) and the position of the Benzene molecule on Fe (100) surface is determined by a combination of the chemisorption of the  $\pi$ -molecule on the reactive surface and the strong interfacial spin polarization. While the Benzene molecule dramatically modifies the electronic structure of the interface and can lead to drastic change in the interfacial spin polarization. While the Benzene molecule becomes slightly magnetic (0.032  $\mu_{\rm B}$  by C atom), it is found that the polarization of the interface). This polarization reversal is analyzed in terms of orbital overlaps between the Benzene protitals and Fe-d orbitals .

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### SK13

#### Local atomic structure analysis of ferromagnetic semiconductor GeMnTe by atomic resolution holography

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The IV-VI diluted magnetic semiconductor GeMnTe is expected as a spintronics material, because it shows a ferromagnetic order below 190 K at x = 0.08 [1]. It is believed that its ferromagnetism is a carrier induced type. However, another possibility is a cluster formation of a ferromagnetic MrGe binary alloy. The solution can be obtained by an investigation of the crystal structure around the each element. Recently, we have investigated the atomic arrangement around the Ge and Mn atoms in the GeMnTe by an X-ray fluorescence holography (XFH) [2, 3]. The obtained atomic images from the XFH suggest that the Mn atoms replace with the Ge atoms in the host GeTe, the Mn position is stable in the exact positions of the anion fcc sub lattice [2], and the Ge position is fluctuated [3]. We have also performed the Te 4d photoelectron holography (PEH) in order to investigate the local structure around the Te atoms in the GeMnTe thin film single crystal. The obtained atomic image from the XFH three the three the obtained atomic image from the reat opticate the local structure around the the atoms in the GeMnTe thin film single crystal. The obtained atomic image from the PEH suggests that the Te atoms are not located at the exact positions of the fcc anion sublattice.

[1] Y. Fukuma, et al., Appl. Phys. Lett. 93 (2008) 252502. [2] N. Happo, et al., Jpn. J. Appl. Phys. 50 (2011) 05FC11. [3] N. Happo, et al., e-Journal of Surface and Nano Technology 9 (2011) 247.

#### SK14

# Dependence of the tunneling magnetoresistance on the inserted nonmagnetic layer

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When a crystalline nonmagnetic (NM) layer was inserted between the tunnel barrier and a magnetic layer in magnetic tunnel junctions (MTJs), it was observed that the tunneling magnetoresistance (TMR) decayed and oscillated as a function of the thickness of the nonmagnetic layer. In this paper, the TMR dependence on the NM layer is investigated theoretically. The conductance was calculated with the Landauer-Buttiker formalism, and the conductance difference between the parallel and antiparallel magnetizations was given in an analytical form with the spin-dependent reflection amplitudes. In order to check the validity of our model, numerical calculations were carried out. An effective mass theory was adopted for the electronic structure of each material, and several cases were considered. It was shown that our simple model agreed well with the results of the numerical calculations. Our model can be applied also to more realistic cases with the full-band structures. The oscillation period was determined from the extremal vectors of the Fermi surface of the NM material, and the selection rules for the oscillation period were obtained. Comparisons with the experimental data and the previous theoretical results will be presented.

#### SK15

# The bulk Fe-Mo double perovskite analyzed from a small clusters perspective

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Being the  $Sr_2FeMoO_6$  double perovskite among the promising materials for spintronic applications, the bulk case was studied in this work; but attention was focused, also, on the octahedral arrangements  $FeO_6$  and  $MoO_6$  as well as on  $FeO_4$  and  $FeMoO_{11}$  clusters. Calculations were made using Density Functional Theory (DFT) and the Perdew-Burke-Emzerhof (PBE) functional within the Generalized Gradient Approximation (GGA) scheme. In order to understand the differences in the behavior among up- and down-spin electrons observed in the half-metallic Sr\_FeMoO<sub>6</sub> double perovskite, the orbitals and density of states were analyzed. Our results reveal that the half-metallic character is present ever since the isolated (FeO<sub>6</sub>)4- cluster and is a consequence of the non-zero spin at Fe atom coming from t2g orbitals. The (MoO<sub>6</sub>)6octahedra does not inhibit the perovskite half-metallic character because this Mo-containing cluster has a zero total spin. In addition, Sr atoms were aggregated to the studied cluster, not just because of the searched relation between the bulk material and the cluster units properties, because they give the impression that play a relevant role to become stable the charge even though the energy deepest of the states contributed by them. Acknowledgments: Work supported by Multidisciplinary Project SIP-IPN: 2012-1439, PAPIIT-IN108710 UNAM.

### SK16

Tunneling magnetoresistance effect in magnetic tunneling junctions with a high resistance ferromagnetic oxide  $Fe_{2.5}M_{0.5}O_4(M = Mn, Zn)$  electrode

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Recently, molecular spintronics has attracted much attention. To improve the spin injection efficiency from ferromagnetic electrodes into molecular materials, by avoiding an impedance mismatch problem, the use of high resistance ferromagnetic electrodes is more accessible than other methods. In this study, a high resistance ferromagnetic electrodes is more accessible than other methods. In this study, a high resistance ferromagnetic oxide  $F_{c_2}M_{0.5}O_4(FMO: M = M, Z)$  was focused and observation of a tunneling magnetoresistance (TMR) effect in magnetic tunneling junctions (MTJs) with an FMO electrode was implemented. MTJs composed of MgO(100)-sub. / FMO(thickness, 25-30 nm) / AI-O(1.5-2) / Ni<sub>80</sub>Fe<sub>20</sub>(20) were fabricated by using pulsed laser deposition, photolithography, Ar ion milling, and sputtering. Magnetic properties were investigated with magnetic Kerr effect measurements. Magnetoresistance properties were measured with a 4-terminal-method. A TMR effect reflecting the magnetization processes of the both ferromagnetic electrodes in the MTJs with M = Mn was successfully observed, and the TMR ratio was approximately 0.85% at room temperature (RT) [1]. The spin-polarization of FMO with M = Mn at RT was estimated to be at least 0.94%[1]. Higher spin-polarization is expected by improving the interface and the barrier qualities between the ferromagnetic electrodes.

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#### SK17

#### Monitoring of gamma radiation interaction in PHR sensor D.g. Park and Hoon Song

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The Planar Hall Resistance (PHR) sensors have been found to be capable of magnetic micro and nano bead detection as the signal to noise ratio in these PHR sensors are higher by several orders compared to others magnetoresistive biosensors. In addition, the spin valve multilayer are considered as novel structures for a high sensitivity and uniform magnetization rotation compared to exchange biased bilayer structure for planar Hall resistance sensors.(1). In this work, we discussed a method for immobilization of magnetic labels on the PHR sensor for sensitive detection of magnetic labels. We used photolithography to generate the functionalized sensor surface to demonstrate the capacity of monitoring the immobilization of biomolecule. The PHR sensor was irradiated by Co source up to 80 Gy and the sensor characteristic was compared with non irradiated one. The field sensitivity obtained by the slope ( $\Delta V$ / AH) of the field-voltage curve was not changed by high dose irradiation up to 80 Gy The sensitivity of the PHR sensor was identified detecting the different concentration of magnetic beads. We have successfully demonstrated the immobilization of strentavidin coated magnetic bead to the single stranded oligonucleotide on the Au coated PHR sensor surface for the detection of magnetic labels.

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### SK18

#### Reactively sputtered MgAl2O4 barrier layers for Heusler tunnel iunctions

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Magnetic tunnel junctions (MTJs) using epitaxial MgO barriers have attracted great interest because they exhibit large tunnel magnetoresistance (TMR) ratios. However, the large lattice mismatch (about 4%) between MgO and half-metallic Heusler alloy induces misfit dislocations and antisite disorder at the surface of ferromagnetic electrode[1]. Thus, coherent growth of MTJ is important for further improvement of device performance. This study focuses on spinel MgAl<sub>2</sub>O<sub>4</sub> as a tunnel barrier lattice-matched with half-metallic Heusler alloy. We fabricate a Fe<sub>2</sub>CrSi/MgAl<sub>2</sub>O<sub>4</sub> junction by reactive sputtering, which can control the crystallinity and degree of oxidation of the films strictly, and investigate its crystal structure and electrical transport properties. MgAl<sub>2</sub>O<sub>4</sub> junctions were obtained by inserting an ultrathin MgAl<sub>2</sub> interlayer, which worked as a protection layer for oxidization at the surface of Fe<sub>2</sub>CrSi. The rocking width of MgAl<sub>2</sub>O<sub>4</sub> (004) was around 1°, which is small value compared to that of the electron-barrier nlight  $\phi = 0.860.95$  eV for the Fe<sub>2</sub>CrSi/MgAl<sub>2</sub>O<sub>4</sub>/Or Fe MTJ. Present epitaxial MgAl<sub>2</sub>O<sub>4</sub> barrier deposited by reactive sputtering theory, and we obtained barrier hight  $\phi = 0.860.95$  eV for the Fe<sub>2</sub>CrSi/MgAl<sub>2</sub>O<sub>4</sub>/Or Fe MTJ. Present epitaxial MgAl<sub>2</sub>O<sub>4</sub> barrier deposited by reactive sputtering is expected to realize high performance spintronic devices.

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### SK19

#### Half-metallic properties of (001) surfaces of the Cr substituted rocksalt GeTe-based compounds

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Much effort has been devoted to the development of new magnetic materials for the spintronic applications. Recently, new half-metallic ferromagnets based on the rock-salt semiconductor GeTe were found [1]. In the bulk structure, GeTe becomes half-metal when some of Ge atoms are substituted by Cr or V. We report the properties of (001) surfaces of Cubic CrGe<sub>3</sub>Te<sub>4</sub> and tetragonal CrGeTe<sub>2</sub> with two different terminations, i.e. GeTe and CrTe terminations, as obtained by using the FLAPW [2] within GGA. We found that both CrTe- and GeTe-terminated CrGe<sub>3</sub>Te<sub>4</sub> and CrGeTe<sub>2</sub> (001) surfaces conserve the half-metallic properties of their bulk structures by analyzing the DOS and confirmed from the integer value of total magnetic moment (MM) per unit cell. The center layer energy gap is much larger as compared to its bulk. The energy gap and half metallic gap became wider with the larger substitution of Ge atom by Cr atoms. In the CrTe-terminated surface the MM of the Cr atom at the surface layer has been enhanced as compared to the center layer of the surfaces is 3.7  $\mu_{\rm a}$  in the both system.

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#### SK20

#### Electronic structure and spin polarization of Co<sub>2-x</sub>Fe<sub>1+x</sub>Si Heusler alloy Hirovoshi Itoh<sup>1</sup>\* and Svuta Honda<sup>2</sup>

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We study the electronic structure and spin polarization for Co-based Heusler alloys of Co<sub>2x</sub>Fe<sub>1+x</sub>Si using the first principles calculations to discuss the spin injection into semiconductors and paramagnetic metals. Half-metallic Heusler alloys such as Co<sub>2</sub>MnSi, Co<sub>2</sub>FeSi, and Fe<sub>2</sub>MnSi are promising materials for spin generator, injector, and detector. Recently, epitaxial films of Fe<sub>2</sub>MnSi [1] and Co<sub>2</sub>FeSi [2] have been successfully fabricated on Ge or Si substrates and the spin-dependent transport properties have been measured. Obtained results seem sensitive to the degree of disorder and the composition of the alloy. In this contribution, we examine composition dependence of the electronic structure for Co<sub>2x</sub>Fe<sub>1+x</sub>Si by performing band calculations based on the density-functional theory (DFT) in the generalized gradient approximation (GGA) using the projector-augmented wave (PAW) method. It is shown that the spin polarization defined by the total DOS decreases from +100% at x=0 and changes the sign with increasing Fe composition x because the half-metallicity is lost. However, the spin polarization defined by the s-component of DOS (partial DOS) keeps positive although it decreases with increasing x from x=0 to x=2.

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### SK21

# Large enhancement of Kerr rotation of GMR periodic patterns using Pt / Co free layer

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Magneto-optical spatial light modulator (MO-SLM) driven by spin transfer switching is a promising device for holographic 3D display. However, a small Kerr rotation angle is an issue to be solved to realize the MO-SLM. In this paper, we report an enhancement of Diffraction Magneto-Optic Kerr Effect (DMOKE) obtained for periodic GMR patterns. A GMR structures with Pt/Co free layer were used in this experiment. The samples were patterned into periodic structures with sizes of 5-50 µm and ptiches of 10-100 µm. SiO<sub>2</sub> was deposited on the etched area after the patterning to obtain flat surfaces. DMOKE of the GMR patterns were measured by using a He-Ne laser with a wavelength of 632.8 nm and an incidence of 30°. DMOKE hysteresis loops were clearly observed, corresponding to the magnetization reversal of the free layers, and Kerr rotation angles of (1,0) and (1,1) diffraction spots were measured to be 0.5°, which is six times larger than a polar Kerr rotation angle measured before the patterning. We consider that the large enhancement of DMOKE is due to an interference effect occurred in the periodic GMR structures embedded in SiO2. This work was supported by the NICT.

K. Aoshima et al, Appl Phys Let, 44, 052507 (2007)

# SK22

# Ab-initio and tight-binding calculations of magnetic anisotropy phenomena in CoPt

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Ordered CoPt alloys hold potential for applications in high density magnetic recording due to the combination of exchange and spin-orbit interactions giving rise to large magnetic anisotropies. Tunneling Magnetoresistance [1] and Tunneling Anisotropic Magnetoresistance [2] have been demonstrated in tunneling devices with CoPt electrodes. Controlling the spin-orbit coupling (SOC) phenomena is essential as the same anisotropy that ensures thermal stability of a memory element makes the writing by external magnetic fields or electric currents challenging. An alternative to the spin transfer torque (STT) effect in multi-domain systems is based on current induced polarization of the conduction electron spins in structures possessing SOC and lacking inversion symmetry. Switching of a ferromagnetic layer induced by in-plane current has been shown recently in a Co-Pt bilayer [3]. Another promising way of controlling the magnetic anisotropy via SOC utilizes lattice strains induced by piezo stressors. We study SOC effects in strained CoPt multilayer structures using a relativistic full-potential linearized augmented plane-wave method. We compare the variations of anisotropies of magnetic total energies (MAE) and density of states (ADOS) also to ur tight-binding calculations using a realistic Slater-Koster parametrization which provide more practical basis for further research of magnetotransport anisotropies in CoPt based nanostructures.

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#### SK23

# Spin polarization of half-metallic Heusler alloy Co<sub>2</sub>MnSi by Andreev reflection measurements

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We report spin polarization measurements of half-metallic Heusler alloy Co<sub>2</sub>MnSi. Heusler alloys are regarded as the most promising half-metallic materials for spintronics applications because of their high Curie temperature and low coercivity. Recently, large tunneling magnetoresistance (TMR) value was experimentally reported from a Co<sub>2</sub>MnSi/Al-O/Co<sub>2</sub>MnSi magnetic tunnel junction (MTJ), which was attributed to the half-metallicity of Co<sub>2</sub>MnSi. Therefore, we have determined the spin polarization of Co<sub>2</sub>MnSi by the Andreev reflection technique. Polycrystalline Co<sub>2</sub>MnSi samples were fabricated by using a high-frequency induction furnace. We have made planar-type junctions by depositing Pb on the polished surface of the polycrystalline Co<sub>2</sub>MnSi samples. The differential conductance of Co<sub>2</sub>MnSi/Pb planar-type junctions was measured between 1.5 K and 10K by the lock-in amplifier technique. The obtained differential conductance was able to be fitted very well by the modified Blonder-Tinkham-Klapwijk (BTK) model. We have found the spin polarization of approximately 50% from the analysis of Andreev reflection measurements of half-metallic Heusler alloy Co<sub>2</sub>MnSi.

### SK24

# Electronic properties of Co2Fe\_xMn\_{\rm Lx}Si Heusler alloys studied by hard X-ray photoelectron spectroscopy

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One major research area where Co<sub>2</sub> based Heusler compounds are investigated is on TMR junctions. Fully epitaxial magnetic tunnel junctions (MTJs) with a Co<sub>2</sub>YZ thin film as a base electrode and a MgO tunnel barrier exhibit high TMR ratios. Quaternary Co2Fe<sub>2</sub>Mn<sub>1,x</sub>Si (CFMS) Heusler alloys are even more promising because the substitution of Mn by Fe increases the number of valence electrons and therefore leads to higher Curie temperatures above 1000 K. Kubota et al. found a maximum TMR ratio for Co2Fe<sub>2</sub>Mn<sub>1,x</sub>Si (x = 0. . . .)-CoFe MTJs [1] for Fe fractions of between x=0.4 and 0.6. The halfmetallic character disappeared for samples with x≥0.8. The reason of this behaviour is still unknown. A B<sub>2</sub> and L<sub>21</sub> site ordered stack of MgO substrate/ Cr [40 nm] buffer layer/ Co2Fe<sub>2</sub>Mn<sub>1,x</sub>Si [30 nm]/ MgO [2 nm]/AlO<sub>x</sub> [1.3 nm] (x = 0, 0.2...1.0) films was studied at beamline P09 at PETRA III and beamline ID32 at the ESRF. The hard X-Ray photoelectron spectroscopy studies were performed with an excitation energy of 6 keV. It was found that the position of Co 2p core levels and the density of states in the valence band region strongly depend on the Fe/Mn ratio.

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### SK25

#### Ab-initio calculation of the magnetic properties of bn nanoribbon Jeffrev Rufinus

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The emerging field of spin electronics (spintronics) has been continuously attracting researchers. Substantial theoretical and experimental efforts have been made in the quest to find the candidates for future spintronics devices. Recently, the search for new spintronics materials has also included graphene-based materials due to the theoretical prediction that this type of material may show the half-metallic property. Here, we present the results of a density functional theory within a generalized gradient approximation study of Boron Nitride (BN) nanoribbon. The objective of this study is to determine whether this type of material will be ferromagnetic or antiferromagnetic if the width is varied. Our results show that the narrow zigzag shaped BN nanoribbon prefers ferromagnetic or antiferromagnetic state depending on the width. Thus, by controlling the width (and growth direction) it is possible to manipulate the magnetic property of this material. These results are of the scientific interest in exploring the magnetic properties of BN-based nanoribbon for future spintronics.

#### SK26

# Synthesis and functional properties of polycrystalline Fe<sub>3</sub>Si-based magnetic tunnel junctions

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Magnetic tunnel junctions (MTJs), consisting of two ferromagnetic (FM) layers separated by ultrathin tunnel barrier are key elements in a wide range of spintronic devices, particularly, as bit elements in magnetic memories and as reading heads in racetrack memory concept. The improvement of the MTJs tunnel magnetoresistance (TMR) might be attained by the selection of the materials with high spin-polarization and the appropriate combination of FM electrodes and tunnel barrier [1]. In this work, we discuss the applicability of a half-metal Heusler alloy Fe<sub>5</sub>Si as an top and bottom electrodes in MTJ, the material which has been poorly adopted so far [2,3]. MTJs based on polycrystalline Fe<sub>5</sub>Si ferromagnetic (FM) electrodes and both SiO<sub>2</sub> and MgO tunnel (~1~3 nm) insulators (TI) have been grown by the combination of pulsed laser deposition (using two-laser co-deposition approach) and ion beam assisted deposition techniques. The morphology, phase composition and structural properties of individual layers as well as the intermixing and chemican at reactions at the FM/TI interface have been investigated. The magnetic properties of the trilayers and preliminary TMR measurements on Fe<sub>5</sub>Si-based MTJs obtained at wide range of temperatures (2-400 K) will be presented, aiming at the evaluation of potential of Fe<sub>5</sub>Si-based MTJs for spintronics.

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#### SK27

# Half-metallic properties of the (001) surfaces of the half-heusler compounds GeKCa and SnKCa

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The sp or d0 half-metallic materials have been attracted more and more increasing attention. A recent first-principles calculation found that the half-Heusler compounds, GeKCa and SnKCa, are half-mtallic [1]. We investigated the magnetic and half-metallic properties of the (001) surfaces of the GeKCa and the SnKCa by using the all electron full-potential linearized augmented plane wave method [2] within generalized gradient approximation. We considered two different terminations for each compound, i. e. Ge (or Sn) and KCa terminations. From the calculated density of states, we found that the Ge or the Sn termination reserves the half-metallicity, while the KCa termination for both compounds is not half-metallic. The magnetic moments for the surface Ge and Sn atoms are much enhanced compared to the bulk values, but the KCa surface for both compounds is almost magnetically dead.

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# SK28

### Half-metallicity in hydrogenated carbon nanotubes

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Zigzag-edge graphene nanoribbons (ZGNRs) show half-metallicity arising from edgelocalized electron states under an electric field, enabling fully spin-polarized electrical current and thus providing a basis for spintronics. While carbon nanotubes (CNTs), obtained by rolling up the ZGNRs, have themselves no edges, edge states have been identified in hydrogenated CNTs. Thus, whether or not the CNTs can be half-metallic under an electric field becomes an issue of great interest. Here, our density functional theory (DFT) calculations show that the edge states in the CNTs are unstable under an electric field due to the spin-conserving electron transfer between the edges, but a large enough transfer barrier between the edge states, obtained by controlling the adsorption patterns, can render the CNTs half-metallic. Our results can be extended to most carbon systems thus opening a new route to materials engineering and providing a deeper understanding of the electronic properties of  $\pi$ -bond network edges.

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#### SK29

# Enhanced perpendicular magnetic anisotropy in Fe/(MgAl<sub>2</sub>)OX bilayer structures with interface optimization processes

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Ferromagnetic metals (FM) electrodes with large perpendicular magnetic anisotropy (PMA) are actively being investigated. Recently PMA arising at the interface between a FM and an oxide layer such as CoFeB/MgO [1] and Co<sub>2</sub>FeAl/MgO [2] were reported. The PMA at which between FMs and oxide barrier materials has attracted much attention not only for the STT-MRAM applications but also for spintronic devices with electric field control of magnetization. The detailed mechanism of the interface PMA is still an open question since the experimental reports, in particular the materials examined, are much limited. Here, we report the enhanced PMA in Fe/MgAl<sub>2</sub>O<sub>x</sub> structures with the effective perpendicular magnetic anisotropy energy density (Keff)  $\sim$  3 erg/cm<sup>3</sup>. In the case of an Fe layer with a crystalline (MgAl<sub>2</sub>O<sub>x</sub> layer, it exhibited perpendicular magnetization characteristics, and the value of Keff was increased from  $\sim$  0 to  $\sim$ 3 erg/cm<sup>3</sup> with increasing annealing temperature. Furthermore, even in the Fe layer with amorphous (MgAl<sub>2</sub>)O<sub>x</sub>, a large interface magnetic anisotropy was obtained.

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#### SK30

#### Graphene nano-ribbon and the ripple effect Hsin-han Lee, Kuo-chin Chen and Ching-ray Chang\* Physics. National Taiwan University. Taiwan

Graphene is a possible nano-device in the future. There are several experiment and theory exhibit that the graphene would generate the ripple in some conditions[1-3], and there is a "geometrical effect" on the curved surface with Rashba spin-orbit effect[4]. We investigate the transport property and the spin density accumulation on the graphene nano-ribbon with the ripple. This research used the non-equilibrium Green function(NEGF) of the Landauer formula[5.6] to compute the numerical results for the system.

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### SK31

#### Enhancement of spin signal in all-metallic lateral spin valves with halfmetallic heusler alloy

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A lateral spin valve (LSV), consisting of two ferromagnetic wires connected with one non-magnetic wire, is one of the promising devices to convert charge currents to pure spin currents, and vice versa. The conversion ratio is, however, still too small, because of the large spin absorption in ferromagnetic wires. Recently, Fukuma et al. have demonstrated that inserting the MgO barrier between ferromagnetic and non-magnetic wires suppresses the spin absorption, resulting in the enhancement of the spin signal[1]. Here, we demonstrate another way, using the half-metallic Heusler alloy as ferromagnetic electodes. The devices were fabricated from a film stack of SiO<sub>2</sub>(2)/Cu(10)/Co.<sub>2</sub>FeGa.<sub>3</sub>Ge.<sub>0</sub>(25)/Ag(10)/Cr(1) on MgO substrate (unit in nm). The non-magnetic wire was made of 100 nm-thick Cu. The typical wire width was 150 nm, while the distance between two ferromagnetic wires was 350 nm. The non-local resistance observed is 12.8 m\Omega at room temperature that increases up to 44.8 m\Omega below 10 K[2]. These values are much larger than those in typical LSVs made of Permalloy and Cu in previous reports, indicating the advantage of the half-metallic Heusler alloy as a ferromagnetic electodes in LSV.

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### SK32

Effect of electron beam rapid thermal annealing on the TMR of CoFeB/MgO/CoFeB magnetic tunnel junctions

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In this paper, we present our studies on the effect of Electron Beam Rapid Thermal Annealing (ERTA) on the structural evolution and related transport properties of CoFeB/MgO/NiFe MTJ. The MgO based tunneling junctions were fabricated in a DC-RF magnetron sputtering system with a base pressure 2 X 10<sup>-7</sup> mBar. An e-beam gun with a 3kV power supply was used to facilitate electron bombardment on the MTJ for the purpose of rapid annealing .We observed the evolution of TMR in CoFeB/MgO/NiFe MTJ when subjected to ERTA. The main factor that distinguishes this study from other rapid annealing time (20 seconds). We observed an ultrafast crystallization of the FM layers of the MTJ which resulted in the evolution of TMR.

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### SK33

#### Preparation of Ti-N films for a capping layer of a CoFeB/MgOmagnetic tunnel junction

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CoFeB/MgO-MTJ is the most promising system for Spin-RAM because of its high tunnel magneto-resistance (TMR) [1]. Recently, perpendicular magnetic anisotropy (PMA) which was expected to reduce current density for magnetic switching by spin-transfer torque (Jc0) down to applicable value (<10a A/cm<sup>2</sup>) was reported at reduced CoFeB layer thicknesses (<1.5 nm) [2]. However, drastic increase of damping constant (a) was observed as the reduction of the CoFeB thickness, which consequently lead to high J<sub>c0</sub> (3.9 x 10<sup>6</sup> A/cm<sup>2</sup>). The possible cause of the increase of a is interdiffusion between Ta capping layer and CoFeB layer[3], [4]. Therefore, a capping layer material whose thermal stability is high than that of Ta is necessary. Our study is to observe a better capping layer material for CoFeB/MgO-MTJ than Ta. We focused on Ti-N and prepared CoFeB layers capped by Ta and Ti-N on MgO substrates and annealed them at various temperatures (T<sub>a</sub>) up to 800 °C. Clearly different behavior at over T<sub>a</sub>=600 °C in Ta dependence of saturation magnetization of CoFeB was indicated that Ti-N was a promising capping material for reduction of J<sub>c0</sub> in CoFeB/MgO-MTJ.

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# SK34

#### Half-Metallic Molecular Wire on Silicon Surface

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Based on DFT calculations including the molecular dynamic simulations, at first time, we describe a feasible reaction path to make a robust half metallic Mo-borine sandwich molecular wire on hydrogen-terminated Si(001) surface along [100] direction. The wire is chemically bonded to the surface, therefore, it is solid enough against thermoperturbation, at least up to room temperature. Meanwhile, the electronic and magnetic properties of this molecular wire are tunable via an external electric field (EEF), and the wire exhibits multifunctional electronic and magnetic states, which allows it to work as a multistep molecular switch or/and multistep information storage component. Fantastically, this surface wire seems independent from the underplayed silicon bulk even under EEF situations since bulk properties have not been affected. This is because the charge transfers are constrained within the wire and the upmost layer Si atoms only. All of these features should be meaningful in the applications of molecular electronics or spintronics in near future.

### **SL01**

# Structural and magnetic studies of sol-gel prepared hexagonal BaFe<sub>12</sub>O<sub>19</sub>

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Barium hexaferrite particles prepared by sol-gel method was characterised using X-ray diffraction spectroscopy (XRD), Field emission scanning electron microscopy (FeSEM), Superconducting Quantum Interference Device (SQUID) and Thermogravimetric analysis (TGA). Nano-sized particles  $BaFe_{12}O_{19}$  and  $Fe_{20}$  were obtained after calcination at 8000C. The  $BaFe_{12}O_{19}$  microstructure showed an average grain size of 200nm and recorded a saturation and remanent magnetisation of 60 and 35 emu/g, respectively. In addition, a commercially available  $BaFe_{12}O_{19}$  was also studied. The feature observed was quite distinctive and was mostly in platelets form and much larger in size. It has a saturation and remanent magnetisation of 70 and 5 emu/g, respectively. The poor remanent magnetism of this commercially available  $BaFe_{12}O_{19}$  was due to large particle size and accompanied by an increased in surface area. TGA performed at 1000°C recorded about 1.8% wt loss in the sol-gel prepared barium hexaferrite compared with 0.2% as observed in the commercial product.

#### **SL02**

#### Magnetization reversal process in antiferromagnetically coupled (Co/ Pd)/Ru/(Co/Pd) multilayer dot pattern

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Magnetization reversal process, exchange interaction Hex between layers and coercivity in the antiferromagnetically coupled (Co/Pd)/Ru/(Co/Pd) dot patterns were studied as a function of dot size. Pd<sub>10.5</sub>/[Pd<sub>0.7</sub>/Co<sub>0.4</sub>]<sub>4</sub>/Ru<sub>0.4</sub>/[Co<sub>0.4</sub>/Pd(1.5 nm)]6 were sputtered onto pre-patterned Si substrates. Dot shape was designed to be 100 - 3000 nm square. Magnetization measurements revealed that top and bottom layers have a perpendicular magnetic anisotropy and couple each other antiferromagnetically at remanent state. MFM measurement indicated that the magnetic state of each layer becomes single domain below side length d=200 nm. Complicated magnetization reversal process was observed. The magnetization reversal starts from a nucleation of reverse domain at dot edge in the reversal from the  $\uparrow/\downarrow$  ( $\downarrow/\uparrow$ ) state to the  $\uparrow/\uparrow$  ( $\downarrow/\downarrow$ ) state, while the coherent rotation type behaviour was found in the  $\uparrow/\uparrow(\downarrow/\downarrow)$  to  $\uparrow/\downarrow$  $(\downarrow\uparrow\uparrow)$  reversal. Hex was about 2.2 kOe in the film and decreases to less than 1 kOe at d=100 nm. Coercivity is about 1 kOe for the film and is almost constant independently on the dot size. The magnetization reversal process and the dependence of Hex and coercivity on dot size are discussed in terms of the structural and magnetic property of intermediate Ru layer in patterned dots.

### **SL03**

#### Magnetic properties of mechanically alloyed Fe-Cu particles Osamu Kohmoto, Masakazu Uchida and Yasushi Matsushima Okayama University, Japan

Mechanically alloyed powders have metastable structure. In the Fe-Cu phase diagram, both the Fe and Cu are known to have quite small solubility in to Cu and Fe, respectively, at room temperature. We have prepared x at% Cu-Fe particles by mechanical alloying. The starting materials were pure Fe and Cu powers. The crystallite sizes of Fe with 700 A for bcc(110) and Cu with 45 nm for fcc(111) decreases below 20 nm at milling time of 500 h. The 500-h-milled alloys with x=0-30 % Cu show a single bcc phase, and the alloys with x=35-100 % Cu show a single fcc phase. Co-existence of fcc and bcc phases is not observed at x=30 and 35 %Cu. However, at x=30 at %, a slight increase in the mass density deduced by lattice parameter is observed and the density jumps by 0.85 g/cm3 at 35 at%. The magnetization at 16 kOe decreases from 200 emu/g with increasing Cu content, however, it has a sudden decrease at 35at% Cu of the phase boundary. The magnetization reaches zero at x=90 % Cu. Coercive force increases initially with the Cu content, however, it has a sudden jump at the phase boundary.

#### **SL04**

#### Temperature dependence of the coercive force of ferromagnetic TM-Al-O (TM=Fe, Co) granular films

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Temperature development of the coercive force (Hc) of ferromagnetic TM-Al-O granular films has been investigated in a temperature range from 290 K down to 5 K. Each film is prepared using the magnetron sputtering method. During the depositing process, a dc magnetic field was applied parallel to the film surface to induce in-plane uniaxial anisotropy. An Fe-Al-O granular film shows no clear in-plane anisotropy at 290 K. On the other hand, Co-Al-O films show clear in-plane anisotropy at 290 K. The Hc of the Fe-Al-O films strongly depends on temperature. The Hc of Co-Al-O films along the easy axis increases rapidly with decreasing temperature. This behavior in the Hc also depends on the Co-concentration. For films with lower Co-concentration and larger electrical resistivity, enhancement of the Hc becomes larger at the lowest temperature. The Hc of Co-Al-O films shows characteristic behavior at low temperatures. It increases with decreasing temperature, following a logarithmic law below about 50 K.

#### SL05

# Microscopic dipole-exchange theory for magnonic crystal arrays of interacting ferromagnetic nanorings

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A microscopic theory is employed to study magnonic crystals represented by onedimensional arrays of ferromagnetic nanorings. The Hamiltonian-based formalism includes exchange and dipole-dipole coupling within each ring, the long-range dipolar coupling between rings, and an external applied magnetic field. Depending on the applied field, as well as the size of the rings and the inter-ring spacing, the array elements may be in a vortex, onion (bidomain), or other inhomogeneous magnetization state. Numerical results are reported for the magnonic bands and gaps in Permalloy nanorings with outer diameter typically 80 to 120 nm, thickness typically 15 to 30 nm, and spacing varied between 10 and 100 nm. For low values of the field (e.g., 0.02 T) along the length of the array, the vortex state is favored energetically and it is found that the magnonic bands are fairly broad and the gaps are relatively small, particularly when the spacing is reduced. Conversely, at larger fields (e.g., 0.2 T) favoring the onion state, the magnonic bands are narrower and the gaps larger. Varying the inner diameter of the rings, as well as the spacing, can control the transition fields between the different states, which may be advantageous for device applications.

# **SL06**

# Synthesis and magnetic properties of zinc ferrite nanocrystals and their applications

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We studied magnetic properties of zinc ferrite nanoparticles via thermal decomposition whereas zinc ions partially replace iron ions at A sites with a nominal composition of  $(Zn_{0}Fe_{0})Fe_{0}O_{4}$ . Our first-principle calculation has shown that Zn ion has significantly lower formation energy at A-site compared to the substitution of B-site and an enhancement of saturation magnetization of 10-20% is obtained due to the reduction of the total magnetic moment of A-site which is antiferromagnetically coupled with B-site. The iron and zinc precursors [Fe(acac)3, Zn(acac)2,xH<sub>2</sub>O] were used in the synthesis of zinc ferrite particles. We have found that a molar ratio of zinc precursor to iron precursor of 1.2 is required for the end product of  $(Zn_{0},Fe_{0})Fe_{2}O_{4}$ . It has been found that magnetization measured at room temperature is dependent sensitively on particles. Room temperature magnetization reaches 100.9 emu/g for zinc ferrite nanoparticles with an average particle size of 85 nm. The value is significantly larger than that of bulk Fe<sub>2</sub>O<sub>4</sub> (80-90 emu/g). The as-synthesized zinc ferrite particles with an average particle size of 26 nm were selected for magnetic fluid hyperthermia. After mPEG surface coating, the particles could be well dispersed in water. A comparative high SAR value of 630 W/g was obtained.

#### **SL07**

#### Nanocrystallite size-induced changes in the magnetic and transport properties of La<sub>1-x</sub>Ca<sub>x</sub>MnO<sub>3</sub> (x = 1/8, 3/8, 5/8) manganite Yugandhar Bitla and S. N. Kaul\*

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Magnetization, magnetic hysteresis loops/isotherms, linear and non-linear ac susceptibility, electrical resistivity and magnetoresistance measurements have been performed on 40 nm nanocrystalline La<sub>1-x</sub>Ca<sub>x</sub>MnO<sub>3</sub> (x = 1/8, 3/8, 5/8). Compared to the bulk, reducing the crystallite size down to 40 nm results in the following changes. For x = 1/8 [1], lowers paramagnetic insulating (PMI)- ferromagnetic insulating (FMI) transition temperature T\_(c) to 143 K, induces the metal-insulator (MI) transition at T\_(MI) = 77 K and charge ordering (CO)/orbital ordering (OO) transition at T\_(CO/OO) = 35 K but leaves the insulating cluster spin glass transition temperature T\_(g), the temperature T\_(TT) below which short-range antiferromagnetic (AF) order coexists with OO and the spin-wave stiffness at 0 K, D(0) = 35(2) meVA<sup>2</sup>, essentially unaltered. For x = 3/8, lowers T\_(MI) to 234 K, a second-order (first-order) PMI-FMI phase transition occurs at T\_(c) = 273.5 K when the applied magnetic field, H = 0 (H  $\neq$  0), D(0) retains its bulk value [2] 145(5) meVA<sup>2</sup> and T\_(CO) reduces to 30 K. For x = 5/8, FM and CO ordering in the insulating state sets in at T\_(c) = T\_(CO) = 270 K and transition to the AF/OO phase occurs at T (AF) = 50 K.

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### **SL08**

# Magnetic and magneto-optical properties of bilayered Co/Ni anti-dot arrays

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In this study, we have successfully fabricated bilayered Co/Ni anti-dot arrays using photolithography and controlled wet-etching processes. Such a ferromagnetic structure was made by sequentially depositing Co (40 nm)/Ni (5 nm) bilayer on a Si substrate. The anti-dot patterning with a hole diameter of 1 -m was done only on the upper Co layer, while the Ni underlayer was kept uniform. The magnetic properties of anti-dot arrays were investigated by using a superconducting-quantum-interference-device magnetometer and by magnetic-force microscopy (MFM). The longitudinal Kerr rotation (LKR) of the zeroth- and the first-order diffracted beams were measured at an incidence angle of 30°. Significantly, it was observed that the LKR of the first-order diffracted beam is nearly 4 times larger than that of the zeroth-order beam. The simulated results for the hysteresis loops matched qualitatively well with the experimentally obtained ones. The MFM images revealed well-defined periodic domain structures which can be ascribed to the anisotropis such as magnetic uniaxial anisotropy, configurational anisotropy, etc. We observed that the magnetization reversal of such a system proceeds through the formation and the annihilation of domains. The observed changes in the magnetic properties are closely related to the patterning that hinders the domain-wall motion.

# **SL09**

# Diffracted magneto-optical Kerr effect of Co anti-dot structure in different arrangements

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To fullfill the increasing requirement for the miniaturization of magnetic materials, the effective extraction of information from the nano-scale structures becomes hot issue. Some of techniques, such as diffracted magneto-optical Kerr effect (DMOKE), magnetic-force microscopy and Lorentz scanning transmission electron microscopy, have been utilized to read the magnetic information. As a derivative technique of MOKE, which can be applied in describing the magnetic hysteresis loops, the DMOKE provides us more unique performance, since the diffracted light is more sensitive to the change of the magnetization distribution in magnetic material. Micromagnetic simulation is also a useful method to study the process of magnetization reversal. In order to get the distribution of magnetization, the Object Oriented Micromagnetic Framework (OOMMF) program with periodic boundary condition was employed. By combining DMOKE and the micromagnetic simulation, we can theoretically and experimentally get more exact magnetic information. We investigated the transverse DMOKE and the magnetization-reversal process in rhombic and square lattice arrangements of Co anti-dot structures. The hysteresis loops of the diffracted beams are described by the magnetic form factor. The shape of hysteresis loops of the diffractions, especially, at the points of inflection of the loop, is expected to indicate a special distribution of magnetization.

### **SL10**

# Monodisperse magnetic nanoparticles: effects of surfactant concentration on the reaction between Fe(acac)3 and Pt(acac)2

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Magnetic properties of monodisperse nanoparticles for ultrahigh density recording and biomedical applications are sensitive to their shape and size distributions. These attributes are, in turns, dictated by several parameters during the synthesis and heat treatments. In this work, monodisperse magnetic nanoparticles with size of around 4-5 nm were synthesized by the modified polyol process with 1.0 mmol Fe(acac)3 and 0.5 mmol Pt(acac)2 as starting materials. The as-synthesized nanoparticles were transformed from their superparamagnetic phase to the ferromagnetic state by annealing in argon atmosphere at 570 C for 1 hour. Without any reducing agents, concentration of surfactants (oleic acid and oleylamine) of 1.5 - 4.5 mmol affected morphology and magnetic properties of the nanoparticles. When 2.5 mmol of each surfactant was used, the as-synthesized nanoparticles have the smallest variations in size and shape. However, the higher concentration of oleic acid and oleylamine led to a higher corecrivity in annealed nanoparticles which are related to the size distribution. These results indicated that the appropriate amount of surfactants must be used in order to optimize the size distribution and magnetic hysteresis of nanoparticles.

### SL11

#### Hollow MnCO<sub>3</sub> and MnSiO<sub>3</sub> nanosphares

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Magnetic nanoparticles can be used in many different applications ranging from medicine, electromagnetic wave detectors to digital information storage. Hollow nanosphares offers additional functionality. For example they show improved electrochemical performance for batteries or can be used for drug delivery. We synthesised MnCO<sub>3</sub> and MnSiO<sub>3</sub> hollow nanospheres with a diameter of few nanometers and thickness of only few atomic layers. In the MnSiO<sub>3</sub> hollow nanospheres a small concentration of  $Mn_3O_4$  was detected from magnetic measurements. Its strong ferrimagnetic signal at TN = 42 K suppressed the signal of intrinsic magnetic properties of MnSiO<sub>3</sub> nanospheres. On the other hand in the MnCO<sub>3</sub> nanospheres no  $Mn_3O_4$  was detected which enable us to study their dc and ac magnetic properties. The Neel temperature at around 33.5 K is between the reported values for MnCO<sub>3</sub> single crystal (34.3 K) and powder sample (31.0 K)[1]. A coercive field and remanent magnetization of MnCO<sub>3</sub> nanospheres.

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#### SL12

#### **Shape-controlled synthesis and magnetic properties of FePt nanocubes** Mingge Zhou<sup>1</sup>\*, Wei Li<sup>2</sup>, Minggang Zhu<sup>2</sup>, Dong Zhou<sup>2</sup> and Yanglong Hou<sup>3</sup>

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In this article, synthesis of monodisperse iron-platinum (FePt) nanoparticles by reduction of platinum acetylacetonate followed by thermal decomposition of iron pentacarbonyl in the presence of oleic acid and oleylamine stabilizers is reported, Particle shape is controlled by the addition of Mo(CO)6, our synthesis yields particles with shapes that are close to cubic and their size can be controlled from 7 to 10 nm, In order to form FePt film, Then the self-assembly FePt nanocubes were annealed at 700 °C for an hour did the FePt nanoparticle transform to the fct phase. The phase analysis, structure, and magnetic properties were determined by X-ray diffraction (XRD), High resolution transmission electron microscope (HRTEM), Scanning electron microscope and vibrating sample magnetometer (VSM) techniques.

#### SL13

#### Magnetic properties of NiO/Ni(OH)2 core-shell nanostructures Mangesh B. Mahajan and P. A. Joy\*

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Antiferromagnetic materials show interesting and anomalous properties when their size is reduced to nanodimensions. The magnetic susceptibility is increased with decreasing particle size(1). Bulk NiO is antiferromagnetic and recent studies showed that NiO nanoparticles exhibit aging and memory effects similar to that of spin glasses(2). The finite size and surface effects play major role in determining the magnetic behaviour of these nanoparticles(3). We have synthesized NiO nanoparticles by a simple chemical route followed by calcination. The freshly prepared sample was dispersed in water, then filtered and dried in air to get NiO/Ni(OH)2 core-shell type nanostructures. These nanostructures were characterized by using different techniques and magnetic properties using SQUID-VSM. Magnetic properties of nanoparticles of NiO, Ni(OH)2, and NiO/Ni(OH), have been compared to study the effect of Ni(OH), layer on NiO. Low temperature M-H curves showed considerable difference in the coercivities and saturation magnetization values of the core-shell structure. From ZFC and FC measurements it is observed that magnetic behaviour of the NiO/Ni(OH), core-shell nanoparticles differs substantially from that of bare NiO and Ni(OH), nanoparticles. The observed changes are explained in terms of lower size of NiO core and the effect of different inter-particle interactions with and without the hydroxide layer.

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#### SL14

## Structual and magnetic properties of $MFe_2O_4$ (M=Ni, Mg) nano hollow spheres

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Hollow sphere structures have recently received considerable attention because of their novel interior geometry and surface functionality. In this report, sub-micrometersized hollow spheres assembled from ferrite nanoparticles have been prepared by a wet chemial process. Measurements of dc-susceptibility and magnetization were carried out on samples of MFe<sub>2</sub>O<sub>4</sub> (M = Ni, Mg) nanoparticles with average diameter 4-5 nm, and hollow spheres (500 nm in diameter) with about 12-17nm thickness shell. A superparamagnetic-type blocking process is observed in the ZFC and FC magnetization and magnentic susceptibility for all samples. Furthermore, a behavior due to surface spin freezing of nanoparticles is found in Ni-ferrite. The surface spin freezing temperatures Tf are quite lower than the blocking temperature, and the magnetic field dependence is different between nanoparticle and hollow sphere. The spin-glass nature of Ni-ferrite nanoparticle is demonstrated by the magnetic field dependence of Tf following the well known Almeida-Thouless line. In contrast, Tf (H) of Ni-ferrite hollow sphere hints at failure of Gabay-Toulouse theory as well as the AT type.

#### **SL15**

#### **Magnetization reversal in patterned arrays of (001)Fe particles** Maj Hanson<sup>1</sup>\*, Rimantas Brucas<sup>2</sup> and Erik Wahlstrom<sup>3</sup>

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Arrays of elliptical Fe particles prepared by patterning [1] of epitaxial (001)Fe films may serve as model systems for studies of the processes of magnetization reversal in sub-micron size particles. In this work we studied such systems with film thicknesses in the range between 10 nm and 50 nm. Ellipses with aspect ratio 1:3 were prepared with their short axes being 50, 100 and 150 nm and varying inter-particle distances. The magnetic domain state and the magnetization reversal of individual particles were investigated by magnetic force microscopy (MFM) and micromagnetic calculations. Magnetization curves of the arrays, comprising about 106 particles, were measured with an alternating gradient force magnetometer (AGM). Within the range of actual sizes we find that the largest particles may be left in a multi- or bidomain (MD) state after demagnetization whereas the smallest show single-domain (SD) properties. In particular we analyze the influence of interactions and the transition from MD to SD behaviour as displayed in the shape of the hysteresis curves and in the MFM images taken after remanent demagnetization. The single particle properties are compared with results from micromagnetic calculations.

1. M. Hanson et al. JAP 85 (1999) 2793

#### SL16

## Controllable structure and magnetic properties of cobalt nanowires by tuning deposition voltage

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Magnetic properties of Co nanowires have been shown to be controlled by various growth parameters. The deposition voltage shows great convenience in controlling microstructure. Reported results were mostly focused on low voltages [1-4]. In this work, Co nanowires were fabricated into ordered AAO templates by dc electrodeposition in a wide range of deposition voltage at 20°C. With deposition voltages of 2.0 and 4.5 V, Co nanowires are of hcp(100) and fcc(220) along the nanowire-axis, respectively. Two structures coexist for medium voltages. For all samples, the easy axes of magnetocrystalline (Kmc) and shape (Ksh) anisotropy make angles to each other. For low voltages, Kmc and Ksh are close to each other, resulting in small coercivity Hc(l) and Hc( $\frac{1}{4}$ ) for H parallel and perpendicular to nanowire-axis. Since Kmc is one order smaller than Ksh and thus the easy-axis lies along the nanowire-axis for high voltages, large Hc(l) is obtained at room temperature. For low and high voltages, Hc(l) first decreases and then increases with a minimum at 200 K, and increases monotonically with T increasing from 50 to 390 K[5], respectively. The angular dependence of Hc shows a combined magnetic reversal process which can be described by curling and/or transverse domain wall modes[6].

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#### SL17

#### Magnetotransmission effect in Nd<sub>0.5</sub>Sr<sub>0.5</sub>MnO<sub>3</sub> nano-composites

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Magnetic and optical properties of composites on the base of Nd<sub>0</sub>, Sr<sub>0</sub>, MnO<sub>3</sub> powders diluted in dielectric

matrix have been studied. Nd<sub>0.</sub>Sr<sub>0.</sub>MnO<sub>3</sub> manganite is interesting due to existence of two different type of phase transitions: the first one is from paramagnetic semiconductor to the FM metallic state near TC-250 K, the second one is from the FM metallic state to CE-type AFM insulating state with charge ordering near TCO-160 K [1]. The giant negative magnetotransmission effect (relative change of natural light transmission under the magnetic field) up to 14% in the field H = 0.8 T was revealed for Nd<sub>0.5</sub>Sr<sub>0.3</sub>MnO<sub>3</sub> composites near TC in the infrared range. The charge ordering and AFM state may by destroyed by decreasing of grains to nanoscale size [2] resulting in appearing of new effects. The remarkable feature of Nd<sub>0.5</sub>Sr<sub>0.3</sub>MnO<sub>3</sub> nano-composite with average grain size about 100 nm is an existence of significant magnetotransmission about 7% in the wide temperature range near and far below TC. It was explained by the significant contribution of nano-particle shells with disordered spins to the magnetic properties of nano-composite. Finally, such a composite material is a promising candidate for future magneto-optical applications [3].

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#### SL18

#### Influence of asymmetric permalloy ring on magnetization configuration and switching behavior

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By introducing an inner circle shifted along the diameter of outer circle, an asymmetric ring and the asymmetric ratio (R) are defined to control over the vortex state (flux-closure state) nucleation and annihilation field in asymmetric Permalloy (Py) rings. Series of asymmetric Permalloy rings with different diameter and asymmetric ratio were prepared to observe the magnetic switching behavior and the magnetization configuration. The asymmetry of Py rings with thickness of 20 nm, outer (Do) of 300, 500, 800 nm; inner diameter (d) of 100, 150, 200, 250, 300 nm, and shift length (S) of 30, 50, 70, 100, 120, 150 nm. Numerical micromagnetic simulations are carried out using the object-oriented micromagnetic firamework (OOMMF) software. With applied field from the positive to negative magnetic field, the magnetic field at which the vortex nucleate in the Py ring is defined as the nucleation field (Hn). The annihilation field (Han) is the magnetic field where the vortex annihilate. Linear relations of vortex state nucleation and annihilation fields to asymmetric ratio were found. and are useful mainty for controlline the nucleation and annihilation fields.

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#### SL19

#### Fabrication of high aspect ratio nanoscale magnetic tunnel junction etch mask by oxygen plasma assisted resist trimming

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The development of large-area, nanostructure patterning with high aspect ratios is a challenging problem that must be addressed for potential applications in high performance nanoscale devices[1]. When negative resist pillar patterns are formed using electron beam lithography, the forward and backward scattering of the electrons limit the resolution therefore, the resist layer should be thinner. When the patterned resist pillars were used as dry etching masks, the resist polymer during the process requires the pillars to have high aspect ratio. Also, due to chemical processing, high aspect ratio resist pillars have a high probability of falling over due to capillary forces on the pillars[2]. We report on a simple method to overcome the low aspect ratio and the collapse of resist by using oxygen plasma assisted resist trimming. Firstly, We used 170 nm thick negative electron beam resist and patterned 55 nm dot arrays using 80 kV electron-beam.Next, resist trimming process resulted in 5 nm trimming in all direction. Therefore the 70 s ashing resulted in a 30 nm diameter and 145 nm tall resist pillar increasing the aspect ratio from 3 to 5, which is about 70 % increase.

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#### SL20

### Magnetic properties of iron(III) oxide nanostructures by hydrothermal synthesis

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The hydrothermal synthesis of nanoparticles enables functional control on the distribution of particle size, morphology and crystallinity by controlling the synthesis parameters. The synthesis of nanoparticles of iron oxide by this method was performed by raising the temperature of a autoclavable bottle containing a solution of iron(III) chloride, sodium nitrate and hydrochloric acid in appropriate quantities. Iron oxide powders synthesized under conditions previously published in the literature were treated in a box furnace at temperatures of 390, 450, 550 and 800 °C. For these samples we performed structural, morphological and magnetical characterization. The images of transmission electron microscopy provide evidence that the 1D nanostructures of iron(III) oxide, are joined by surfaces with non-oriented crystallographic planes, suggesting that the growth of these particles occurs by oriented coalescence. After heat treatment a phase transition occurs from akaganeite to hematite. Magnetization measurements show that akaganeite presents weak magnetic behavior, whereas the hematite phase presents a larger remnant magnetization and coercive field. We thank the support of Brazilian agencies FAPESP and CNPq.

#### SL21

#### Magnetic dot-antidot lattice for control of magnetic anisotropy

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The magnetic anisotropy of epitaxial Fe/GaAs(100) films varies with thickness of the film.[1, 2] We have prepared rectangular lattices of antidots using epitaxial 35 monolayer (ML) Fe/GaAs(100) films deposited by Pulsed laser deposition in an ultrahigh vacuum chamber whose base pressure was less than 2exp(-10) Torr. 10ML MgO/35ML Fe/GaAs(100) film was controllably etched to reduce the thickness of lithographically defined regions of the Fe film to 19ML and capped with MgO before liftoff. Thus square antidot regions of size 200nm containing Fe of thickness lower than the surrounding matrix were created. Thus we have created a 2 dimensional magnetic system having a uniform material composition but controlled lateral variation of local magnetic anisotropy. The spacing between the antidot regions is varied from ~ 200nm to 400nm. Investigations of magnetic anisotropy on the dot-antidot bilayer were performed using Magneto-Optical Kerr effect in longitudinal configuration and Ferromagnetic Resonance. Anisotropy constants are evaluated and it shows that the ratio of the magnitude of the uniaxial anisotropy to the magnitude of the four fold anisotropy is much higher (0.96) in the dot-antidot bilayer even with a square lattice. This is discussed in terms of exchange coupling between the antidot region and the surrounding matrix

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#### SL22

## Effect of diameter of the wires on magnetic properties of electrodeposited CoNiP hard magnetic nanowires

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The magnetic properties in CoNiP nanowires arrays with diameter of nanowires in range of 50-800nm were investigated. All the samples were prepared by electrodeposition method with pH of 3.1 and room temperature. The electrochemical potential of CoNiP were determined by cycle voltammetary. The crystalline structure and morphology of the samples were characterized by X-ray diffraction (XRD) and Scanning Electron Miroscopy (SEM), respectively. The hysteresis loops were measured at room temperature using vibrating sample magnetometry (VSM) and alternating gradient magnetometry (AGM) with magnetic field applied perpendicular and parallel to the wire axis. The mixture of fcc and hcp polycrystalline phases of the CoNiP based nanowires has been indicated by the XRD pattern. The obtained results show that with 800 nm of diameter of the nanowires we can obtain maximum coercivity value (2070 Oe). The uniaxial anisotropy values dependent on diameter of the wires are calculated. The anisotropy of the samples are quasi-onedimensional anisotropic is also found when diameter of the wires are less than 500 nm.

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#### **SL23**

## Switching behavior of lithographically defined grid of permalloy nanowires studied with magnetoresistance

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We report the switching behavior of a grid [1] of permalloy nanowires fabricated using e-beam lithography and DC-magnetron sputtering followed by lift-off process. The wires are of 20 nm thickness and about 120 nm wide. The periodicity between the wires is maintained at one micron throughout the grid. The grid is made as a Hall bar geometry which consists of few hundreds of junction regions. The magnetoresistance study indicate that the node regions in the grid acts as defects in the horizontal wires in longitudinal magnetic field, which in turn leads to the nucleation of domains. Thus, the switching process in the horizontal wires is governed by these defect centers [2]. The switching field varies from 80 Oe to 140 Oe for transverse magnetoresistance is due to the presence of nodes. These nodes possess four-fold anisotropy with easy axes along the diagonal directions. Further, the micro-magnetic simulations using Object Oriented Micro Magnetic Framework (OOMMF) with the two dimensional periodic boundary condition, gives the information about the domain configuration which supports our experimental findings.

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### POSTER PRESENTATION

#### SL24

## Magnetization reversal modes in narrow FePt nanowires with high perpendicular anisotropy

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In this contribution, we present the dependence of magnetization reversal modes in FePt nanowires with the width ranging from 2000 nm to 30nm using magneto-transport and magnetic force microscopy (MFM) measurements. Magneto-transport measurements showed an enhancement of magnetic corecivity when reducing the width. Theoretical calculations suggest that this coercivity increase is not only due to the suppression of available propagation paths but also to the contribution of the roughness. Further study of the magnetization reversal process using MFM shows that three different reversal modes can appear. Indeed, for widths above 500 nm, the structure of the reversed magnetic domains appears to be similar to those of the continuous FePt film: the reversed magnetic domain grows and expands within a dendritic structure. For widths below the characteristic dendrite width (~300 nm), the reversal takes place by propagation of a single DW that reverse the whole wire. Finally, we show that another behavior appears at very low widths: when reducing the widths below 50 nm, the propagation field becomes larger than the nucleation field. Nucleation thus occurs randomly, the reversal consisting in a mix of nucleation and propagation. Such behavior could prevent the use of ultra-narrow wires for DW-based spintronic devices.

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#### SL25

#### Structural and magnetic behaviour of nanocrystalline CaFe<sub>2</sub>O<sub>4</sub>

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Abstract. Study of the structural and magnetic properties of the nanoparticles of large cation substituted ferrite, CaFe2O4 have been reported here. The cubic spinel phase of CaFe<sub>2</sub>O<sub>4</sub> could be synthesized effectively in nanocrystalline form of size 5.9 nm using advanced sol-gel method. The sample was then heated at 1000 °C leading to bulk size of orthorhombic structure. Very interestingly, the nanoparticles of CaFe2O4 get transformed from cubic spinel structure to its characteristic orthorhombic structure on annealing at 1000 0C. Rietveld profile refinement of the XRD patterns was performed to study the nature of crystal structure. The SQUID magnetization measurements divulged that the nanocrystalline sample is superparamagnetic above the blocking temperature of 150 K whereas the bulk sample is ferromagnetic even at room temperature. The reduction in saturation magnetization in the case of nanoparticles as compared to its bulk counterpart has been explained on the basis that the magnetic moments in the surface layers of a nanoparticle are in a state of frozen disorder. The departure of the field cooled curve from the zero field cooled curves in the momentversus-temperature plot, further confirmed the room temperature superparamagnetic behavior of the nanocrystalline CaFe2O4.

#### **SL26**

### Magnetoresistance of helimagnetic ordering in single crystal FeGe nanowires

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We report on the synthesis, structural characterization, and magneto-transport of single crystalline FeGe NWs. Transition metal silicides and germanides have attracted a lot of attention because of their chiral cubic helimagnetism as a consequence of the Dzyaloshinskii-Moriya interaction. FeGe is known to have a high Neel temperature and helical spin order with a relatively long period. In addition, one dimensional-confinement of the spin structure might enhance the helical magnetic ordering of FeGe NWs as the nano-confinement often alter or improve the properties. We have synthesized single crystalline FeGe NWs by CVD process without any catalyst. We have investigated the helimagnetic ordering in FeGe NWs by using electrical and magneto-transport measurements. When the magnetic field was applied longitudinal to the NW axis, the signature of the helimagnetism in the FeGe NWs was observed up to near room temperature. The magneto-transport measurements reveal three magnetic states of the FeGe NWs: the helimagnetic state, conical helimagnetic state, and ferromagnetic state. The magnetic transitions from the conical helimagnetic state to ferromagnetic state in FeGe NWs was clearly observed and occurred at much higher field those observed in the bulk FeGe. The relationship between nano-sized confinement and the conical helimagnetic state will be discussed in detail

#### SL27

#### Arrays of interacting ferromagnetic nanofilaments: small-angle neutron diffraction study

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Magnetic properties of spatially ordered arrays of interacting nanofilaments have been studied by means of small-angle diffraction of polarized neutrons. Several diffraction maxima or rings that correspond to the scattering of the highly ordered structure of pores/filaments with hexagonal packing have been observed in neutron scattering intensity maps. The interference (nuclear-magnetic) and pure magnetic contributions to the scattering have been analyzed during the magnetic reversal of the nanofilament array in a field applied perpendicular to the nanofilament axis. The average magnetization and the interference contribution proportional to it increase with the field and are saturated at H = Hs. The magnetic reversal process occurs almost without hysteresis. The intensity of the magnetic contribution has hysteresis behavior in the magnetic reversal process for both the positive and negative fields that from the field dependence of the intensity in a butterfly shape. It has been shown that this dependence is due to the magneto is troposed.

#### **SL28**

#### Magnetic properties of nanometer scale FeCr antidot array system

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In this work, we have fabricated periodic nanometer scale antidot array of FeCr. The deposition was done using electron beam evaporation. The substrate was first coated with 100nm of a suitable resist (ZEP-520A) using spin coating. Then, the mask was patterned on the substrate using an electron beam attached to field emission scanning electron microscope to generate the necessary pattern. After the development, the FeCr (40nm) was coated onto the resist at 0.1A/s using electron beam evaporation. The antidot pitch was maintained at 400nm and the antidot diameter was found to be 300nm. The magnetic nature of the sample was characterized using vibrating sample magnetometer and the in-plane magnetic anisotropy of the system was investigated using torque magnetometer. Further, the magnetoresistance measurements were taken using an inline four point probe. The results indicate that the antidot array system has a predominant ferromagnetic nature, has an n-fold rotational symmetry with relation to the magnetic anisotropy. The easy axis, the hard axis of magnetization along with the net anisotropic energy has also been determined. All data have been appropriate discussed in the work. The use of antidot array systems for various applications is an intriguing new area of research.

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#### SL29

## Fabrication of Al-Ni core-shell structured particles via electroless ni plating

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Al alloy powders have been of much interest in recent years due to their various applications such as metallurgical, chemical, paint & pigment industries, aerospace applications due to their lightness in weight and processability. Major research efforts have been made for improving their functionalities by addition or deposition of alloying elements such as Ni and by controlling the structure. However, it is difficult to fabricate these alloy powders by traditional metallurgical process and it is necessary to evolve new methods. Core-shell structured particles are attracting considerable attention as a consequence of their potential application in different areas of science and technology. Functionalizing the Al core with suitable shell materials indeed opens the door to specific physical, chemical and optical performances of Al. However, due to the oxide layer present on Al surface, it is difficult to plate Al with any metal or metallic base surface coating by post wet deposition techniques. Thus, we fabricated Al-Ni core-shell structured powders using electroless Ni plating. The pre-treatment was conducted for the formation of Ni shell as well as removal of the oxide layer The formation of Ni shell on Al core powder was confirmed by SEM, XRD and VSM results

#### SM01

## Breakdown of barkhausen critical scaling behavior with increasing domain wall pinning in fe thin films

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We report a breakdown of Barkhausen critical scaling behavior in Fe films with increasing domain wall pinning by means of a Kerr microscope, capable of direct domain observation. From the the domain reversal patterns in the films, we find that the Barkhausen jump size generally decreases with increasing Fe thickness, showing the increased domain wall pinning density. Surprisingly, the power-law scaling behavior of Barkhausen jump size distribution gradually disappears as pinning of domain walls in the Fe layer increases. This is due to the fact that magnetization reversal mechanism is changed from a random Barkhausen avalanche dominant mode to thermally activated domain wall creep dominant mode.

#### SM02

#### Enhanced critical fields in MnSi thin films

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The magnetic properties of the cubic helimagnet MnSi are governed by the Dzyaloshinskii- Moriya interaction which in the bulk stabilizes a helical ground state below 29 K with a spin-helix wavelength of 18 nm and a propagation vector along the [111] direction. Recently, MnSi has attracted a lot of attention after a skyrmion lattice had been discovered in the so called A-phase. We have prepared thin MnSi films on Si(111) substrate by alloying a single Mn layer into the substrate as well as by codeposition of Mn and Si. RHEED analysis in combination with TEM investigation reveals that codeposition leads to smooth films whereas the alloying process generates rather three dimensional MnSi islands. In comparison to bulk material, the films show an enhanced magnetic ordering temperature of 45 K, which has been determined using SQUID magnetometry. Magnetization, resistivity, magnetoresistance, and Hall effect measurements have been performed on the films. They show that the critical fields describing the transition from the helical to a conical spin phase and from the conical phase to a parallel spin alignment are dependent on the film thickness and enhanced in comparison to single crystalline MnSi.

#### SM03

#### Magnetism and Cr<sub>2</sub>O<sub>3</sub>-Fe<sub>2</sub>O<sub>3</sub> structure of CoFe/Cr-NOL surface Naoki Shimomura<sup>1</sup>\*, Kazuya Sawada<sup>2</sup>, Tomohiro Nozaki<sup>1</sup>, Masaaki Doi<sup>3</sup> and Masashi

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Magnetoelectric (ME) effect has been investigated for realizing electric controlled spintronic device, such as MERAM.  $Cr_2O_3$  is the most promising material for the application due to high Neel temperature (TN : ~307K). Switching exchange bias field by applying electric field using bulk  $Cr_2O_3$  (0.6mm) has been reported [1]. Recently, we have observed high blocking temperature (TB : ~250K) of CoFe/Cr-NOL (Nano-Oxide-Layer) containing  $Cr_2O_3$  and Fe<sub>2</sub>O<sub>3</sub> (TN ~ 950K) in its surface, in spite of the ultrathin thickness (less than Inm). The NOL is prepared by Natural Oxidation (N.O.) of CoFe and ultrathin (1ML) Cr layer [2]. To clarify the reason of high TB of the NOL is important for achieving thin  $Cr_2O_3$  film that shows ME effect higher than room temperature. In this study, we have investigated the structure and magnetic property of CoFe/Cr-NOL by varying constitution condition, such as deposition system and N.O. intensity. It was found that mixing of Cr into CoFe layer before N.O. and high N.O. intensity cause the formation of  $Cr_2O_3$ -Fe<sub>2</sub>O<sub>3</sub> solid solution in NOL surface, which lower the TB. It is indicated that the presence of not Cr-O<sub>3</sub>-Fe<sub>3</sub>O<sub>3</sub> solution in the existence of both Cr-O<sub>3</sub> and Fe<sub>3</sub>O<sub>4</sub> is needed for high TB.

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#### **SM04**

#### Magnetic phase diagram for non-epitaxial Cr/Gd/Cr-multilayers

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In this work, the low-temperature phase and critical behavior at high temperature of inhomogeneous gadolinium-chromium samples were investigated using DC magnetization measurements. The samples were prepared as Cr(50nm)/Gd(100nm)/Cr(50nm) multilayers grown onto Si(110) single crystal substrates by means of sputter deposition. The samples' inhomogeneity was controlled by annealing at different temperatures Tan = 200, 400, and 500°C during 10 minutes, and their magnetization was measured using zero field cooling (ZFC) and field cooling (FC) procedures at several applied fields. In the vicinity of room temperature, the magnetic behavior of the samples is compatible with a ferromagnetic ground state, which undergoes a rounded ferromagnetic (F)-paramagnetic (P) phase transition at a critical temperature, Tc. At temperatures below Tc, however, differences between the magnetization measured in the zero field cooling (Mzfc) and field cooling (Mfc) procedures are observed. Mfc decreases with increasing temperature. Mzfc increases with increasing temperature and, in addition to the F-P phase transition at Tc, shows a peak at a certain temperature Tg < Tc. It was also found that Mzfc is time dependent at low field and irreversible below an irreversible temperature Ti(H). A (H,T) phase diagram that contains Tc, Ti(H) and Tg(H) is presented for each of the samples.

#### SM05

## Revealing the volume magnetic anisotropy of Fe films epitaxied on GaAs(001) surface

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Fe film grown on GaAs(001) is a model system for studying the spin injection from a ferromagnetic layer into the semiconductor substrate. The Fe/GaAs(001) system exhibits a remarkable in-plane uniaxial magnetic anisotropy (UMA) with an easy axis (EA) parallel to the GaAs[110] direction, and such UMA is usually believed to originate entirely from the Fe/GaAs interface. In this contribution, we will show our quantitative studies on the thickness dependent magnetic anisotropy in Au/Fe/ GaAs(001) system using the magneto-optic Kerr effect with a rotating magnetic field. Through a clear 1/dFe relation of the UMA, we found a UMA component with the EA parallel to the GaAs direction which originates from the volume contribution. Such volume UMA is sensitive to the growth temperature and also strongly correlate with the interface anisotropy. We can conclude that both the volume anisotropy and the interfacial anisotropy are related to the ordered atomic structure at Fe/GaAs interface. Our results may introduce a new aspect for further understanding the origin of UMA in Fe/GaAs(001) system.

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#### SM06

### Structures and magnetic properties of ultrathin Ni/Cu(100) in hydrochloric acid

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Microstructure and magnetic properties of ultrathin Ni/Cu(100) prepared by way of electrochemical approaches are investigated. In a pure supporting electrolyte, the STM image of the chloride covered Cu(100) surfaces show the formation of 90-degree step edges of the terraces as a result of the electrochemical annealing [1]. After adding NiCl<sub>2</sub>, the hydrogen evolution reaction is advanced to a more positive potential. Nickel atoms attach onto the steps and the surface shows single atomic steps corresponding to a layer-by-layer growth. As the coverage of Ni increases, large amount of clusters form on the surface. The adsorption of chloride anions occurs on the top of the films [2]. For thin Ni layers, no magnetic hystsresis is observed due to the lowered Curie temperature for ultrathin overlayers. As the Ni coverage increases, hystsresis loop is observed with in-plane anisotropy of the films. For thicker films spin reorientation transition occurs that may be due to the strain relaxation in the Ni overlayer. The high squareness of the magnetic hysteresis loops confirms the smooth interfaces of the Ni films.

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### POSTER PRESENTATION

#### **SM07**

## Improvement in structural and magnetic properties of laser ablated Ni-Zn Ferrite thin films

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Zinc, as known as non-magnetic element, has been doped into Nickel-Ferrites under thin film form, with a hope to modify its structural and magnetic properties, in order to have the appropriate requirements for high frequency applications. Laser ablated Ni<sub>1</sub>, x<sub>2</sub>n<sub>4</sub>Fe<sub>2</sub>O<sub>4</sub> (0 ≤ x ≤1) thin films grown on R-cut Al<sub>2</sub>O<sub>3</sub> (0001) substrates using pulsed laser deposition (PLD) technique are single phase with (1 1) orientation, and they are room temperature strongly ferromagnetic, while others' similar works reported ferromagnetism only at low temperatures. Additionally, the higher HC value in our films compared to what others reported indicates that somehow we have found an appropriate way to make Ni-Zn ferrites become potential candidates for modern miniaturization of electronic devices. The zero field cooled ? field cooled (ZFC-FC) magnetization curves for all the films shows the blocking temperature (TB) suggesting a strong anisotropy. Research on Ni<sub>1,x</sub>Zn<sub>x</sub>Fe<sub>2</sub>O<sub>4</sub> (0 ≤ x ≤1) thin films would be very promising for device market.

#### **SM08**

### A basic study of magnetic anisotropy strength control using FeSiB magnetostrictive thin film

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The Magneto-Impedance (MI) sensors are used for various applications [1, 2]. A magnetic anisotropy controlling is important to improve sensitivity of MI sensor. Therefore, our research group proposed a new and simpler method to control magnetic anisotropy using inverse-magnetosriction and difference of thermal expansion coefficient in bimetal thin film[3]. To exploit the proposed method, investigation of a fabrication ratio between two layers (thickness and length) requires due to the sensitivity of proposed MI sensor is determined by changes in fabrication ratio. In this paper, we introduce the changes of magnetic anisotropy field (Hk) strength to develop high and adjustable sensitive MI sensor according to changes in geometrical properties of layers. The thin films were deposited by RF magnetron sputter. Fer<sub>2</sub>Si<sub>44</sub>B<sub>44</sub> and Molybdenum were used for the magnetic and conductive layer, respectively. The fabricated thin films annealed to release the local stress. We carried out magnetic domain observation by Kerr-microscope. Moreover, Hk value was obtained by VSM. Based on the experiments, we could confirm that Hk value was controlled by the geometrical properties because the geometrical properties can determine the generated stress in the magnetostrictive magnetic layer. This research provides the fabrication ratio and expectation of the sensitivity of MI sensor.

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#### **SM09**

#### Magnetic properties of co thin films on polyethylene naphthalate organic substrates

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Recently we have proposed spin quantum cross (SQC) devices, in which molecules are sandwiched between two edges of magnetic thin films deposited on organic substrates with their edges crossed. According to our recent calculation, SQC devices exhibit large magnetoresistance effect at room temperature. In this study, towards the fabrication of SQC devices, we have investigated magnetic properties and surface morphologies of Co thin films on polyethylene naphthalate (PEN) organic substrates. Co thin films were thermally evaporated on PEN substrates under a magnetic field in a high vacuum chamber. The magnetization curves were measured by focused magnetooptic Kerr effect techniques at room temperature. The surface morphologies were analyzed by atomic force microscope. As a result, the surface roughness decreases from 1.3 to 0.55 nm with increasing the Co thickness up to 55 nm, where a two-step smoothing phenomenon can be seen. As for magnetic properties, the coercive force and the squareness of the hysteresis loop show the maximum values at a Co thickness of 5.3 nm. This peak is due to the rotation of the uniaxial magnetic anisotropy formed from both the induced magnetic anisotropy and the shape magnetic anisotropy affected by the surface morphologies of Co/PEN and PEN substrates.

#### SM10

#### Effects of dimensionality on magnetization of Ho and Sm-doped BiFeO, thin films

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BiFeO3 is promising multiferroic material due to high ferroelectric (1100 K) and antiferromagnetic (650 K) ordering temperatures. Reports of other groups have suggested that by reducing the dimensionality, the spiral magnetic ordering could be suppressed [1]. Thus, multiferroics in thin films and nanocrystalline forms could enhance the magnetic properties [2]. In BiFeO<sub>3</sub>, if Bi is substituted by a small amount of divalent or trivalent metal ions, or Fe is substituted by transition metal ion, a significant enhancement in magnetization can be achieved. We have prepared the nanocrystalline RE-BiFeO<sub>3</sub> (RE = Sm and HO) by using sol-gel method to be target materials for thin films fabrication. RE-BiFeO<sub>3</sub> (RE = Sm and HO) by using sol-gel method to be target materials for thin films fabrication. RE-BiFeO<sub>3</sub> films were grown on LaAIO<sub>3</sub> substrates using pulsed laser deposition technique. All the films show a single phase of rhombohedral structure with space group R3c. Significant changes in the structural properties were observed as the thickness of the films increases. The magnetization observed in the case of HO is much higher than those reported in literature and also Sm-doping shows better magnetization values. The magnetization increases as the film becomes thinmer. This is a strong evidence for the origin of surface maanetism that exist in the films.

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#### SM11

#### Substrate-dependent electronic anisotropy of epitaxial multiferroic DyMnO<sub>3</sub> and Dy<sub>0.8</sub>Ca<sub>0.2</sub>MnO<sub>3</sub>thin films

Kueih-tzu Lu, S. C. Haw, T. L. Chou, J. M. Lee, S. A. Chen and J. M. Chen National Synchrotron Radiation Research Center, Taiwan

We investigated the substrate-dependent electronic structure and anisotropic bonding of the Mn 3d states in DyMnO<sub>3</sub> and Dy<sub>0.8</sub>Ca<sub>0.2</sub>MnO<sub>3</sub> thin films on SrTiO<sub>3</sub>(001) and LaAlO<sub>3</sub>(110) substrates using polarization-dependent x-ray absorption spectroscopy (XAS) at O K-, Mn L- and Mn K-edges for three polarizations, E || a, E || b and E || c. Polarization-dependent x-ray absorption spectra at O K-, Mn L23- and Mn K-edges of orthorhombic DyMnO<sub>3</sub>/LaAlO<sub>3</sub>(110) thin films show a strong polarization dependence, whereas orthorhombic DyMnO<sub>3</sub> /SrTiO<sub>3</sub>(001) thin films show nearly isotropic spectral structure. The main peak in polarized Mn L2,3-edge XAS spectra of DyMnO<sub>3</sub>/LaAlO<sub>3</sub>(110) thin films for the E || b polarization lies at a lower energy than for polarizations E || a and E || c. This indicates a great anisotropy in Mn 3d-O 2p hybridization, reflecting an orbital ordering and a highly anisotropic coplanar Mn-O bonding in DyMnO<sub>3</sub>/LaAlO<sub>3</sub>(110) thin films. Orbital ordering of eg-orbital and the highly anisotropic in-plane Mn-O bonding is an indispensable factor to the formation of complicated incommensurate modulated magnetic structures observed in orthorhombic DyMnO3. The present results provide important implications for the microscopic understanding of the multiferroic DyMnO3.

#### SM12

#### Enhanced magnetization by substitution of Zn<sup>2+</sup> in Fe<sub>3</sub>O<sub>4</sub> films

Chang-yup Park<sup>1</sup>, Jong-hyun Song<sup>2</sup>, Chun-yeol You<sup>3</sup> and Sung-chul Shin<sup>4</sup>\* <sup>1</sup> Department of Physics and Center for Nanospinics of Spintronic Materials, Korea Advanced Institute of Science and Technology, Daejeon 305-701, Korea <sup>2</sup> Department of Physics, Chungnam National University, Daejeon 305-764, Korea <sup>3</sup> Department of Physics, Inha University, Incheon 402-751, Korea <sup>4</sup> Department of Physics and Center for Nanospinics of Spintronic Materials, KAIST, Daejeon, Department of Emerging Materials Science, DGIST, Daegu, Korea

Fe<sub>3</sub>O<sub>4</sub> is one of the key materials as efficient spin injector for spintronic device. Thereby considerable researches have been carried out to obtain high magneto resistance ratio (MR), however no report show enough MR because of interface concerned problems and uncertainty of half metallic identity itself. Hence we choose Fe<sub>3</sub>O<sub>4</sub> to investigate its nature and enhancement of magnetic and electric properties by substitution of Zn<sup>2+</sup> ion. Firstly we carried out magnetization hysteresis loop measurement by means of alternating gradient magnetometer. To find out a role of Zn substitution, we fabricated Fe<sub>1.8</sub>Zn<sub>1.2</sub>O<sub>4</sub> film and the saturation magnetization strikingly enhanced to 283 emu/cc from 79 emu/cc of the epitaxial Fe<sub>5</sub>O<sub>4</sub> film. We believe that the 3.6 times enhancement is astribed to the substitution of Zn<sup>2+</sup> ion instead of Fe ion at tetrahedral A site which spin is antiparallel to that of Fe ion at octahedral B site in inverse spinel structure of Fe,O<sub>4</sub>.

#### SM13

## Growth temperature dependence of crystal orientation and magnetic properties of CoMn<sub>2</sub>O<sub>4</sub> thin films

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 $Co_{z}Mn_{z_{n}}O_{d}(0\leq\times3)$ , which crystallize in the spinel structure, are strong candidates of multiferoics which can be used in various technical applications due to their thermal, electrical and magnetic properties. In the present work, we report on the growth of epitaxial CoMn<sub>2</sub>O<sub>4</sub> thin films on Nb-doped SrTiO<sub>4</sub>(001) substrate using Pulsed Laser Deposition and their magnetic properties. With all other growth parameters fixed, the growth temperature (Tg) was changed from 720oC to 9200C. The sample with Tg =720oC shows a strongly tensile strained epitaxial crystal structure with the orientation of (400). The M-H loops shows a ferrimagnetic behavior with the corecive field of ~1.7 T. However, with increasing Tg up to 870oC, the tensile strained crystal domain becomes weaker whereas the domain of mixed orientations of relaxed (400) and (004) appear and has gained predominance. The M-H loops shows a weakened ferrimagnetism compared to the sample with Tg =720oC. With further increasing Tg to 920oC, interestingly, the crystal structure turns to be oriented (004) and the magnetic properties schibt enhanced saturated magnetization with negligible corective field. These observations indicate that the growth temperature plays a crucial role in determining the crystal structure and magnetic properties of CoMn<sub>2</sub>O<sub>4</sub> thin film.

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#### SM14

#### Structure and magnetic properties of Fe<sub>3</sub>O<sub>4</sub> thin films on different substrates by Pulsed Laser Deposition Xuelian Huang, Yang Yang and Jun Ding\*

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Magnetite is an attractive magnetic oxide to be used in magneto-transport applications.[1]Many researchers have studied how to optimize the deposition condition for high quality magnetite thin films.[2]In this work,magnetite thin films were deposited on different substrates at temperatures varying from 100°C to 450°C,and their structure,magnetic and magneto-transport properties have been studied. XPS and Raman spectra reveal single phase of Fe<sub>3</sub>O<sub>4</sub> on all the substrates. When film thickness is small and temperature is relatively high(>300°C),magnetite thin films on single crystal SiO<sub>2</sub>Si and amorphous glass exhibit strong (111) texture on matter they have either a huge lattice mismatch or none matching with Fe<sub>3</sub>O<sub>4</sub>. The (111)texture can be formed at lower temperature(100°C) on (002)-ZnO due to the minimization of both surface and interfacial energy. The films on (0001)-Al<sub>2</sub>O<sub>3</sub> show (111)-epitaxial structure in a large temperature range(100-450°C). The saturated magnetization for all the (111)oriented magnetite films are in the reported thin-film range.Magneto resistance up to 4% has been obtained. The out-of-plane MR exhibit quadratic behavior up to 6000 Oe, whereas the in-plane MR is linear in this regime, which suggests the presence of anti-phase boundaries.[3] Large substrate-induced strain was found in the films grown on Al<sub>2</sub>O<sub>3</sub> due to the lattice

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#### SM15

## Influence of crystallographic orientation on the magnetic properties of NiFe, Ni, and Co epitaxial fcc films grown on single-crystal substrates Taiki Ohtani\*, Tetsuroh Kawai, Mitsuru Ohtake and Masaaki Futamoto *Chuo Univercity, Japan*

NiFe, Ni, and Co crystals with fcc structure are soft magnetic materials and have been used in magnetic thin film devices like magnetic heads and sensors. Magnetic property varies depending on the film texture [1,2]. In the present study, Ni<sub>40</sub>Fe<sub>20</sub> (at %), Ni, and Co films with fcc (100) and (111) orientations were prepared by hetero-epitaxial growth on single-crystal substrates of GaAs(100) and Al<sub>2</sub>O<sub>4</sub>(0001). The magnetic anisotropy, the magnetization structure, and the magnetostriction were investigated by employing RHEED, XRD, Bitter, and magnetostriction measurements. NiFe, Ni, and Co films of fcc(100) orientation showed four-fold symmetries in-plane magnetic anisotropy. The easy magnetization axes were parallel to [011] and [01-1] directions. 90° magnetic domain walls were observed for the epitaxial films. The in-plane magnetization properties were reflecting the magnetostriction behavior of respective bulk fcc crystals. On the contrary, the in-plane magnetization properties were isotropic for NiFe, Ni, and Co films with fcc(111) orientation. In magnetostriction measurement, the magnetostriction behavior of epitaxial magnetic thin film depended delicately on the magnetic domain structure and the domain wall motion under an influence of magnetic field. At the conference, the inter-relationships of magnetic anisotropy, magnetization structure, and magnetization structure, and the formain walls were structure and the domain walls were structure and the domain walls were structure.

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#### SM16

## Magnetic coupling in manganite-based thin film heterostructures studied by Electron Holography

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Advances in nanofabrication of magnetic nanomaterials leads to search for new tools to measure physical properties at the nanoscale, such as magnetization. Accurate characterization of the magnetization states in nanostructures is of the utmost importance in the development new devices such as magnetic tunnel junctions (MTJs) [1]. Transmission electron microscopy (TEM) techniques such as Electron Holography (EH) allows the quantitative imaging of the magnetization states of ferromagnetic materials with unprecedented spatial resolution at the nanometer level [2]. Furthermore, EH can be combined with the in situ variation of external parameters such as magnetic and electric fields, temperature, etc. In particular, we use a TEM cryo-holder to image magnetization states of ferromagnets with low TC (down to 100 K) while varying in situ the magnetic field applied by the objective lens of the TEM. In this work, we present the EH study of the magnetization svitching of La-Sr and La-Ca manganite thin films (TC = 180 and 300 K, respectively) and manganite-based MTJs [3]. The (de)coupling of magnetic electrodes in MTJs is analyzed by performing hysteresis loops at 100 K to determine the switching fields from the direct observation of the magnetization orientation of the electrodes upon magnetic field [4].

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#### SM17

Properties of hybrid superconductor/ferromagnet (SC/FM) multilayers U. D. Chacon Hernandez<sup>1</sup>, Y. T. Xing<sup>2</sup>, William E. Alayo<sup>1</sup>, Magda B. Fontes<sup>1</sup>, Jorge L. Gonzalez<sup>3</sup>, Liying

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Recent research on systems that contain two types of materials (hybrid systems) shows that new phenomena can be observed. One of such hybrid systems is superconductor/ferromagnet (SCFM) nanocomposites, which shows coexistence of superconductivity and ferromagnetism with unusual properties. We use niobium (Nb) as the superconducting layers with fixed thickness (50 nm) and cobalt or permalloy (NiFe) as the ferromagnetic layers, with different thickness (5, 10 and 20 nm). The films prepared by magnetron sputtering deposition technique, have been characterized by x-rays diffraction and in some cases the cross section at interface were studied by TEM. Characterization of SC and FM transitions was done by resistivity and magnetization measurements. The properties of superconductor changed at the SC/metal interface because of the proximity effect, for example, decrease of Tc. Moreover, the magnetic layers between superconducting layers will give have more influence on the superconductivity due to its stray field. Our results show that when the external magnetic field is close to the upper critical field Hc2, the SC/FM multilayers can have a vortex solid to a liquid phase transition. We believe that the stray field of the magnetic layers contributes to the melting of the vortex solid state.

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#### SM18

## Effect of substrate on the anisotropic magnetotransport in $Sm_{0.45}Nd_{0.10}Sr_{0.45}MnO_3$ thin films

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 $Sm_{0.45}Nd_{0.1}Osr_{0.45}MnO_3$  thin films on single crystal LAO and STO substrates were prepared by ultrasonic nebulized spray pyrolysis. The PM-FM transition was observed at TC - 153 K and 150 K in film on LAO and STO, respectively. The magnetic moment was found to be smaller in film on STO and it showed stronger bifurcation between the ZFC-FC magnetization. The isothermal M-H loops measured with H applied parallel and normal to the film surface showed that the easy axis was in the film plane, while the hard axis was along the normal. The dominance of the shape anisotropy explains the in-plane magnetic easy axis. Large low field magnetoresistance (MR) is observed in both films. The occurrence of significantly smaller peak anisotropic magnetoresistance (AMR) at a temperature lower than the TIM and the broadening of the AMR-T curve suggest that films on STO have higher fraction of the AFM-COI phase. The difference in the magnetotransport anisotropy in the set of films has been explained in terms of substrate induced strain, e.g., STO provides tensile strain, which is known to favor the AFM-SE and CO-OD phases.

#### **SM19**

#### Magnetism of multilayer (CoNiPsoft/CoPhard)n films

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The (CoNiPsoft/CoPhard)n films were prepared by chemical deposition method. The content of phosphorus was 8 % at. in each layer. In hard magnetic layer the CoP was in hexagonal state and in soft magnetic layer the CoNiP was in amorphous state. The thickness of each layer was t = 4 nm. Magnetic measurements were made with vibrating-coil magnetometer in temperature range T = 77 ÷ 400 K and in magnetic fields H < 10 kOe. We investigated changes of magnetic parameters in dependence on number of layers in multilayer structure. In our experiments number of layer pairs was n ≤ 15. Earlier [1] it was established that the coercive force (HC) of CoP layers depends on layer thickness and HC rises when layer thickness increases. In given case at nitrogenous temperatures the coercive forces of soft magnetic and hard magnetic layers were ~ 20 Oe and ~ 1000 Oe, respectively. The combination of these layers into structure leads to substantial change of magnetic behavior of structure in whole [2]. It attracts attention that the soft magnetic layer determines the behavior of magnetization of film structure decreasing the coercive force of system.

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#### SM20

Significant change in the antiferromagnetic-to-ferromagnetic phase transition temperature of epitaxial FeRh thin films by Ga substitution Ippei Suzuki<sup>1\*</sup>, Mitsuru Itoh<sup>1</sup>, Tetsuya Sato<sup>2</sup> and Tomoyasu Taniyama<sup>1</sup>

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 $B_2$  ordered FeRh alloys exhibit fascinating magnetic properties such as the firstorder magnetic phase transition from the antiferromagnetic (AF) to ferromagnetic (FM) states at around 400 K. For a step toward the incorporation of this material into devices, the control of the transition temperature in a thin film form is one of the most critical issues. In this study, we find that Ga substitution in FeRh efficiently reduces the transition temperature down to room temperature, making the use of this material a reality. A 30-nm-thick FeRh thin film and a Ga-substituted FeRh(Ga-FeRh) were epitaxially grown on MgO(001) substrates by molecular beam epitaxy. The temperature dependent magnetization of the epitaxial films clearly demonstrates the reduction of the transition temperature down to room temperature. The change in the transition temperature can be understood in terms of a possible change in the electronic states near the Fermi level due to the lattice expansion, leading to a more stable FM state in Ga-FeRh. In fact, the out-of-plane lattice parameters calculated from the XRD patterns are found to expand from 2.990 to 3.003 Å by the Ga substitution. Work supported in part by Industrial Technology Research Grant Program in 2009 from NEDO, Japan.

#### **SN01**

### Structural and magnetic properties of Sm<sub>2</sub>Co<sub>17</sub> nanoflakes prepared by surfactant-assisted ball milling

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Microstructure, crystal structure and magnetic properties were studied for Sm2Co17 nanoflakes prepared by surfactant-assisted high energy ball milling with heptane and oleic acid as the milling medium. Effect of ball milling time on the crystallographic alignment evolution and coercivity of the nanoflakes were systematically investigated. Microstructure observation shows that the Sm2Co17 nanoflakes have an average thickness of less than 100 nanometers with an aspect ratio as high as 100. For the 2 hours ball milled nanoflakes, the intensity ratio between (004) and (302) reflection peaks, which indicates the degree of c-axis crystal texture of the Sm<sub>2</sub>Co<sub>17</sub> phase, reaches the maximum value among all the samples, revealing that the strongest c-axis crystal texture was obtained in this nanoflakes. As the ball milling time increases, the intensity ratio drops gradually indicating that the long time milling process undermines the c-axis crystallographic alignment of Sm2Co17 phase in the nanoflakes. On the other hand, it is found that as the ball milling time increases from 2h to 10h, the coercivity of the nanoflakes firstly increases, peaks at 2.8kOe for 8h, and then decreases again. Further investigation indicates that the microstructure evolution plays an important role in the coercivity variance of the Sm2Co12 nanoflakes.

#### **SN02**

### Rotor structure optimization of interior permanent magnet by using response surface method

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Due to limited energy sources and environmental pollution issues recently, the needs of high efficiency electricity equipments are steadily increasing. Among them, an IPM is being developed for the application of eco-friendly vehicle. An Interior Permanent Magnet has a structure of inserted permanent magnet in the inner rotor. These motor requires high torque and low torque ripple. It was designed with 8 pole 36 slot of stator and rotor, alternating N pole and S pole of permanent magnet and triangle type permanent magnet. In order to find the optimized condition, Central Composite Designs method and Box-Behnken method of Response Surface Method were used for finding design variable and minimizing variation of design variable. Comparing two Surface Reaction Methods, electromagnetic finite element method analysis based on inserting variation of design variables, torque, torque ripple, harmonics were obtained and optimized final value was obtained by safety factor through mechanical stiffness analysis.

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#### SN03

Magnetization reversal behavior of FePt/MgO/FePt thin film

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L10 ordered FePt allow has attracted much attention since they are believed to be good candidates for future magnetic devises. It is also thought to overcome the problem of thermal fluctuation in nano-meters scale. A lot of studies such as thin films and self-assembled nano particles have been investigated [1-3]. However, the magnetic properties of the multilayered FePt alloy have not been completely clarified. In order to understand the magnetization reversal, FePt/MgO/FePt trilayer films have been fabricated and detailed magnetic domain observation have been investigated. FePt bottom layer was co-deposited with Fe and Pt on Au buffer layer by using an UHV compatible magnetron sputtering system. Then, the intermediate MgO layer and FePt top layer were deposited. Finally, the samples were annealed at Ta of 600 °C. The magnetic domain structure was observed by magnetic force microscope (MFM) in applied fields up to ±6 kOe. Remarkable steps of the magnetization in the second quadrant were observed for all the samples. This is thought to arise from the difference in the chemical ordering of top and bottom FePt layers. Independent magnetization reversal of the top and bottom FePt layers was clearly observed at the FePt/MgO/FePt circular dots by the MFM measurements

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 1050 (2002). [3] S. Matsumoto et al., J. Phys.: Conf. Ser. 266. 012038 (2011).

#### SN04

## Efficiency and torque density improvement of interor permanant magnet synchronous motor

Mi Jung Kim, Ki Doek Lee, Jae-jun Lee, Jeong-ho Han, Tae-chul Jeong, Woong-chan Chae and Ju Lee\*

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The motor should be designed with a minimum of because motor weight has quite large portion in totality mass; because there is no generator as electrical vehicle, the high efficiency is demanded in order to increase the 1 charging distance. In this paper, the technique is presented for progression of the torque density and the efficiency through the relative electric loading and the shape of rotor magnet. In case of the established model, the required torque satisfied at the rated speed and maximum speed, however, the torque density is low. When the mass of occupied by the motor is going to be bulky, 1 charging distance is decrease due to the increase in the driving load. The current density was increased to boost the torque density. The motor size could be decreased owing to change the current density, but there is a weakness at efficiency and power density, the relative magnetic loading also should increase according to increase the relative electric loading. To increase the relative magnetic loading. To increase the relative magnetic loading.

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#### SN05

#### An effective skew method for torque ripple reduction in surfacemounted permanent magnet motor

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This paper deals with an effective skew method for surface-mounted permanent magnet motor (SPMSM). Skew is a common method to reduce cogging torque and torque ripple as well, and it can be done by continuous way in small machines or multi-stepwise method in big machines [1-2]. Generally, these conventional ways assume that the magnetic flux density distributions are same in axial direction. However, the air-gap flux density is not uniform in axial direction, and it is rapidly reduced at both ends of the motor. Especially, in the pancake type motor, the influence of end sides is significant. Thus, for skew to be effective, non-uniform air gap flux density distribution in axial direction should be considered. In this paper, one method for an effective skew is suggested by considering flux density distribution in axial direction, and the results will be confirmed by 3D finite element analysis.

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#### **SN06**

## Effect of carbon additive on the $TbCu_7$ -type melt spun Sm(Co, M)7 (M=Ti, Zr, Hf, V, Nb and Ge) ribbons

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Recently, Sm-Co compounds with TbCu7-type structure have received intensive attention due to their excellent intrinsic magnetic properties, which have potential to be used as the starting materials for making bonded magnets for above 150 oC applications. Our previous studies indicated that pure TbCu7-type structure could be obtained easily in melt spun Sm(Co,M)7 (M=Ti, Zr, Hf, V, Nb and Ge) ribbons. Based on the purpose of developing Sm(Co,M)7 alloys with TbCu<sub>7</sub>-type structure for practical usage, effect of carbon on the microstructure and magnetic properties of melt spun SmCo7, M, C0.1 (M=Ti, Zr, Hf, V, Nb and Ge; x=0-0.3) ribbons were studied. Based on the XRD and TEM analysis, a pure 1:7 phase could be kept for the ribbons with M= Nb, and Ge, but carbide phases, eg. TiC for M=Ti, ZrC for M=Zr, Sm<sub>2</sub>C<sub>3</sub> for M=Hf, and SmCoC<sub>2</sub> for M=V, respectively, were found for other ribbons. Nevertheless, a slight C addition may effectively refine the microstructure and improve both the intrinsic coercvity and the magnetic energy product in all the studied ribbons. Among them, the SmCo<sub>69</sub>V<sub>0.1</sub>C<sub>0.1</sub> ribbons with or=58.7 emu/g, iHc=13.5 kOe and (BH) max=9.3 MGOe, and SmCo69Hf01C01 ribbons with or=61.6 emu/g, iHc=11.8 kOe and (BH)max=10.3 MGOe are suitable for the bonded magnet applications

#### SN07

#### **Study of designed induction motor on cryogenic LNG pump system** Jinsung Kim and Gwansoo Park\*

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In this Paper, I designed LNGs main cargo pump's motor using LNGs ship. LNGs ship has made a high value product and won steady demand order in every year. However, in Korea, when we make a LNGs ship, we import the key products from other countries(It's almost 60% of order price). So, we need to develop original skill about key product. LNGs main cargo pump's motor is submerged, and induction motor type. In this paper, first of all, design the induction motor using constants at room temperature. And then, we re-design the motor considering an extremely low temperature. This motor's working operation temperature is -163°C. Because of that, it can make a special design. In other words, we don't need to think about heat loss as like as room temperature. In this reason, we consider about losses by current density and flux density. First, coil loss is generated a heat. So, it can be cooling the coil heat loss immediately. On the other hands, magnetic saturation has effect on characteristics of material (specially, B-H curve). So, we think about magnetic saturation, even it works at low temperature.

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#### **SN08**

## Structural and magnetic properties of nanocrystalline $BaFe_{12}O_{19}$ synthesized by microwave-hydrothermal method

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The nanocrystalline BaFe<sub>12</sub>O<sub>19</sub> powders were prepared by Microwave-Hydrothermal (M-H) method at 200°C/45 min. The as synthesized powders were characterized by using X-ray diffraction (XRD). The present were densified at different temperatures i.e., 750°C, 850°C, 900°C and 950°C for 1h using microwave sintering method. The phase formation and morphology studies were carried out using XRD and field emission scanning electron microscopy (FE-SEM). The average grain sizes of the sintered samples were found to be in the range of 185 to 490 nm. The magnetic properties such as saturation magnetization and coercive field of sintered samples were claulated based on magnetization curves. A possible relation between the magnetic hysteresis curves and the microstructure of the sintered samples was investigated.

#### **SN09**

## Research magnetic properties of Fe-O alloys with different texture degrees and ratios of phases using simulating

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The present article is concerned with investigation of magnetic properties of Fe-O alloys, simulating their magnetic behaviour and causation of the high-coercivity state of these alloys. Initial materials were powder so-called "Blue dust" and mixture of this powder with Fe powder with different rations of these phases. "Blue dust" is a powder of Fe<sub>2</sub>O<sub>3</sub>, it's a natural Indian raw. Reserves of this material sufficient large, so a production material with high magnetic properties from it is an urgent problem. Initial samples were powdered in high-energy ball mill for different time periods. This treatment leads to phase transformation. The phase compositions were investigated by X-ray analysis. The impact of time of milling and initial composition on magnetic properties was explored. Milled nanocrystalline powders were subjected to a low-temperature crystallization annealing in magnetic field. It leaded to phase transformation. Powder was pressed into cubic shape samples. The dependence of magnetic properties on composition after annealing was also investigated. The computer simulating was used to understand which phase is cause og forming the highcoercitive state of material. Initial parameters were set so that result corresponds to experimental data. The parameter changing showed quite how each of them influence on magnetic properties

#### SN10

#### Study on FePt/Fe exchange coupling nanocomposite thin films Wenli Pei

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Granular FePt films with island microstructures have been deposited by DC magnetron sputtering. The FePt grain size and distance among FePt grains could be controlled by tuning sputtering condition. An iron layer was deposited on the controllable island FePt film to produce FePt/Fe nanocomposite thin film. The composite films show a hard/soft magnetic exchange coupling after annealing, because the FePt grains in the composite film became L10 structure with good hard magnetic properties after annealing. The exchange coupling could be modified by tuning granular FePt structure and thickness of Fe layer, which means it's promising to produce FePt/Fe nanocomposite magnet with good magnetic properties by this method.

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#### **SN11**

Effects of Sm content on thermal stability of sintered Sm2Co<sub>17</sub> magnets Minggang Zhu, Haibo Feng, Wei Li, Yikun Fang, Wenchen Zhang and Wei Pan China Iron & Steel Research Institute Group, China

Investigations on environment stability of the Sm2Co17 type magnets have been one of hot pots on applications of Sm-Co based magnets. Effects of heavy rare-earth elements substitute Sm on temperature coefficient of remanence (?) of the magnets have been investigated systematically [1]. Sm2Co17 magnet with high coercivity has the best thermal stability among all permanent magnets [2]. The reversible temperature coefficient ? was zero in the (Sm1-xGdx)(CoFeCuZr)z magnet when x=0.55[3]. In this paper, the dependence of Sm content on temperature coefficient of 2:17 type Sm-Co magnets at different temperature were investigated. As shown in Fig.1, the temperature coefficient increased with temperature when z value of the magnets is the same. While, the temperature coefficients of remanence firstly decreased and then increased with z value at the same temperature. When z=7.87, magnet with the best temperature stability, the highest coercivity and the lowest temperature coefficient of remanence were obtained.

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#### **SN12**

#### Simulation of die-upsetting process of hot-deformed magnets

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A plastic deformation simulation was performed by using a three dimensional finite element method based deformation modeling software in order to clearly understand the die-upsetting process in detail. It was found by the plastic deformation simulation that the effective strains of Nd-Fe-B die-upset magnets reduced from the middle to the ends of magnetic specimens along the c-axis, or the press direction, so the crystallographic orientation deteriorated for the grains approached to both ends of die-upset magnets. The magnetic properties were improved with increase of effective strains of die-upset magnets. The remanence Br and maximum energy product (BH) m were high for the specimens, machined from the central region of magnets, because the microstructure of die-upset magnets was homogeneous and grains were aligned well in the central region of deformed magnets according to the simulation. Experimentally, the magnetic properties, X-ray diffraction patterns and microstructure of hot-deformed Nd-Fe-B magnets were studied. The magnetic properties were improved with decreased height of magnet specimens which were obtained by symmetrically removing both ends of die-upset magnets. The grain misalignment was as enhanced from the middle to both ends of die-upset magnets along c-axis. The microstructural characterization was verified by experiment

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#### **SN13**

#### Synthesis of high magnetic moment nanowires for encoding and decoding of barcode segments for multiplexing bio- applications

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Abstract: Sequential template electrochemical synthesis of multisegment nanowires with magnetic and non-magnetic segments have been extremely significant for barcoding, multiplexing and biosensing applications [1,2]. Especially, Hard magnetic segment with high saturation magnetization and remenance materials are more convenient for multiplexing biological applications. In the present article we have synthesized three different types of magnetic nanowires ( CoNiP, CoPtP and CoFeP) for a magnetic segment of barcode nanowires by electrochemical deposition technique using a polycarbonate membrane with a diameter of 50 nm. The length of the nanowires was found to be around 6 µm. We compared the magnetic properties of all the synthesized materials and found to be the CoPtP magnetic nanowires have higher saturation and remenence. In order to demonstrate the decoding of barcode nanowires using the magnetoresistance sensor in flow cytometry, we calculated the spatial distribution of the stray magnetic field produced by the barcode nanowire by means of finite element method (FEM) using the commercial Maxwell software. The CoPtP shows most higher spatial variation compared to the CoFeP and CoNiP, which means CoPrP is most advantageous composition for the hard magnetic segment

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#### A new mechanism of electromagnetic linear-actuator using a magnetic silicone rubber Takanori Fukushi, Sung Hoon Kim\*, Shuichiro Hashi and Kazushi Ishiyama

force, and the flexibility as soft elastmer, so it can be used for various biomedical applications. The details will be introduced in the conference [1] B.Kim, M.G.Lee, Y.P.Lee, Y.Kim, and G.Lee, "An earthworm-like micro robot using shape memory alloy actuator," Sens. Actuators A., 125, 429-437, 2006. [2] H.Lu, J.Zhu, Z.Lin, and Y.Guo, "An inchworm mobile robot using electromagnetic linear actuator," Mechatronics 1-10, 2009. [3] T.Niino, S.Egawa, H.Kimura, and T.Higuchi, "Electrostatic Artificial Muscle: Compact, High-Power Linear Actuators with Multiple-Laver Structures," Proc. IEEE Micro Electro Mechanical Systems Workshop, 94, 130-135, 1994

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We propose a new mechanism of electromagnetic linear-actuator for biomedical applications such

as an artificial muscle and medical micro-robots [1-3]. The proposed mechanism requires two

armatures and magnetic rubber in the DC-solenoid. The proposed structure is series connection:

armature\_1, magnetic rubber, and armarture\_2. Therefore, D.C magnetic field produces attractive

force between the two armatures. At that time, the magnetic nubber causes increase of attractive

force because the magnetic nibber was fabricated by iron powders with silicone elastormer (36 · 64

wt %), as a cylindrical shape (diameter: 12 mm and height: up to 8.9 mm). Its magnetic properties

were measured by VSM: magnetization and coercive force are 50 emu/g and 40 Oe, respectively,

We compared magnetic force between the two armatures with the magnetic rubber and with a non-

magnetic rubber in magnetic field of 600 Oe. We found the magnetic rubber made the attractive force

1.7 times larger. This also means the use of magnetic rubber makes the low power consumption. The

largest advantages of this magnetic rubber are the increase of attractive force, generation of reaction

#### **SN15**

**SN14** 

#### Novel microcrystalline Co-Zr-B RE-free hard magnetic alloys Sofoklis Makridis and Evagelos Gkanas

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The investigation of rapidly solidified alloys is important both for the basic information it gives about the structural and compositional limits of magnetism and for the technical potential of new alloys. Co-based amorphous alloys have the advantage of relatively high saturation flux density and low magnetostriction. The glass formability and technical magnetic properties of rapidly solidified Co(rich)-Zr-B alloys were investigated in the composition range Co82ZrxB20-x. in order to determine the glass forming range and explore the magnetic characteristics as a function of composition and guenching rate. The allovs were found to be completely glassy for low x whereas all the other compositions studied were completely or partially crystalline. The typical glassy alloys showed good soft magnetic properties. The crystalline alloys with Zr content higher than 80 % in Boron were characterized by a coercive force in excess of 4 kOe. Slow-cooled samples of the same composition showed coercivity in the order of 120 Oe.

#### SO01

#### First-principles study on the half-metallicity of full-Heusler alloy Co<sub>2</sub>VGa (111) surface

Hongpei Hong, Kailun Yao\* and Guoying Gao physics, Huazhong university of science and technology, China

A recent experimental study indicates that full-Heusler alloy Co2VGa exhibits halfmetallic ferromagnetism and the magnetic moment measured at 5 K is independent of high pressure up to 12.1 kbar (T. Kanomata et al Phys. Rev. B 82, 144415 (2010)). Using the full potential linearized augmented plane-wave method, we here further investigate the structural, electronic and magnetic properties of all possible (111) surfaces of Co2VGa. Both structural relaxation and calculated surface energy reveal that Ga-terminated surface is more stable than other terminations. From the analysis of the spin-polarized density of states, we find that the half-metallicity confirmed in bulk Co<sub>2</sub>VGa is lost at Co-terminated surfaces but still maintained at V-terminated and Gaerminated surfaces. Moreover, the obtained magnetic moments show that the atomic magnetic moments of the (111) surfaces are greatly different from the bulk values.

#### SO02

#### Magnetic and resonance properties of Bi24(CoBi)O40

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The cubic bismuth oxide is characterized by charge ordering of bismuth ions. A study of magnetic moment of cobalt ions is used to establish the valence state of bismuth ions. The aim of study is to establish the interrelation between spin and elastic subsystems. Measurements of the magnetization was carried out over the temperature range 4K-300 K in the field H=50 kOe and the magnetization curve M(H) was determined at the temperatures T=5 K and T=300 K. The temperature dependence of magnetic susceptibility is well fitted to the Curie-Weiss law with the Curie-Weiss temperature  $\theta$ =- 12.3 K and the effective magnetic moment is u=5.08uB. The bismuth cobaltit is paramagnetic at T >4 K. The electron paramagnetic resonance (EPR) spectra were recorded from 77 K to 300 K at 9.3 GHz. The g- factor increases versus temperature that can be qualitatively explained by local deformation of structure at variation of bismuth ion valence. As a result, the magnetic anisotropy field is enhanced at the heating. EPR linewidth varies linearly with temperature that is explained by spinphonon interaction. Relaxation of magnetic moment is due to strong interaction with the elastic lattice.

#### SO03

#### Field induced phase transitions and magnetocaloric properties in Er1. "Lu, Fe<sub>2</sub>O<sub>4</sub> compound

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Rare earth (RE)-ferrites have attracted attention in recent years due to their various phase transitions (PTs) near room temperature. Their magnetic phase is determined by competition between antiferromagnetic and ferrimagnetic coupling of the localized Fe ions. Magnetocaloric effect (MCE) reaches the maximum near the magnetic PT temperature. Despite of many works of the PT, however, there are very limited works on MCE in the RE-ferrites with two RE elements. In this work one has investigated field induced magnetic PT properties and MCE in the vicinity of the PT point. Stoichiometric polycrystalline samples of  $Er_{1,x}Lu_xFe_2O_4$  (0.1  $\leq x \leq 0.8$ ) were prepared by solid-state reaction method. Magnetic measurements were performed using a SQUID magnetometer. Fig.1 gives isothermal M-H curves of Er<sub>0.9</sub>Lu<sub>0.1</sub>Fe<sub>2</sub>O<sub>4</sub>, from which the field induced PT can be seen clearly from the magnetization curves in the range of 208 K-240 K indicating a first order structural type PT[1]. Fig. 2 shows magnetic entropy change ( $\Delta$ SM) curves. With a maximum field of 1 T and 3 T, the maximum value of ∆SM are 0.49 J•kg-1K-1 and 1.04 J•kg-1K-1 at 249 K, respectively. The maximum  $\Delta$ SM were appeared near the transition temperature and their MCE will be explained in terms of PTs.

#### SO04

#### Magnetic properties of heusler-type Ni-Mn-Ga glass-coated microwires Valeria Rodionova<sup>1</sup>, Maxim Ilvn<sup>2</sup>, Valentina Zhukova<sup>2</sup>\*, Alexander Granovsky<sup>3</sup>, Alexander

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Studies of magnetic shape-memory alloys (MSMAs) attracted growing attention within last few years owing to significant magnetic-field-induced strain (MFIS) useful for applications in actuators, sensors and for energy harvesting. From point of view of applications, miniaturizing of MSMA-based devices based on small-size MSMA materials is quite important. Main advantage of Ni-Mn-Ga Heusler-type microwires is related with composite character, allowing production of long microwire from brittle Ni-Mn-Ga alloy. Consequently the aim of this paper is to present results on Heusler-type Ni495Mn254Ga251 glass-coated microwires. We fabricated few metres long microwire and measured magnetization, M, versus magnetic field, H, curves in as-prepared and annealed up to 790 K microwires. After annealing M(H) curves drastically changed: in as-prepared microwires magnetic transformation appears around 175 K. The sample seems contain the second magnetic phase which orders at low temperature giving rise to fast increase of the magnetization below 30 K. Meanwhile common ferromagnetic behaviour was observed at around 325 K for the annealed sample without any evidence of existence of second magnetic phase. X-ray diffraction allows to detect significant structural changes after annealing. Obtained results are discussed taking into account strong internal stresses induced by fabrication of composite microwires and high quenching rate.

#### **SO05**

#### Compensated magnetism in double perovskites A<sub>2</sub>CrFeO<sub>6</sub> (A=La,Sr)

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Precise compensation of magnetic moment in solid state compounds is very rare. An example is the so-called half-metallic antiferromagnet (HMAFM), which has one metallic spin channel and the other insulating spin channel with the zero net moment. Double perovskite compounds of A<sub>2</sub>BB'O<sub>6</sub> have been theoretically investigated as the most promising candidates for the last 15 years. Very recently, nearly wellordered epitaxial La<sub>2</sub>CrFeO<sub>6</sub> films, theoretically predicted to be a HMAFM, were synthesized using pulsed laser deposition. This sample shows the saturated moment of  $\sim 2 \mu B$ , implying the ferrimagnetically ordered S=3/2 Cr<sup>3+</sup> and S=5/2 Fe<sup>3+</sup>. This results are controversial to the previous experimental results on artificial LaCrO<sub>3</sub>/LaFeO<sub>3</sub> superlattices, which proposed a ferromagnetic order, as well as the existing theoretical predictions. In this presentation, we will revisit this system with first principles calculations using LDA+U, the modified Becker-Johnson functional, and fixed spin moment. Our results are consistent with the recent experimental results, showing a ferrimagnetic insulating La2CrFeO6. Furthermore, we have investigated the Sr-analog Sr<sub>2</sub>CrFeO<sub>6</sub>, which is as yet unsynthesized. Contrary to La<sub>2</sub>CrFeO<sub>6</sub>, this system has the tetravalent Cr and Fe ions with antialigned S=1 moments, resulting in zero net moment. Our findings suggest a precise HMAFM Sr2CrFeO6.

#### **SO06**

#### Possible half-metal antiferromangetism in double perovskites A<sub>2</sub>VMnO<sub>6</sub> (A=La,Sr)

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A half-metallic antiferromagnet (HMAFM) is half-metallic without macroscopic magnetic moment. HMAFM has been considered a promising candidate of a spin injector in spintronics, though no true HMAFM has been established yet. Very recently, A well-ordered double perovskite La<sub>2</sub>VMnO<sub>6</sub>, which was theoretically predicted to be HMAFM 15 years ago, has been synthesized by pulsed-laser deposition. Contrary to the theoretical predictions, the saturated magnetic moment of 2 µB/f.u. implies an ferrimagnetically ordered S=1 V3+ and S=2 Mn3+ ionic state, below TN=20 K. Besides, the synchrotron radiation photoemission measurement indicates an insulating La<sub>2</sub>VMnO<sub>6</sub> with a gap of 0.9 eV, instead of being half-metallic. In this presentation, we will revisit this system using various first principles calculations to unravel these discrepancies between the previous theoretical suggestions and the experimental results. Furthermore, our calculations suggest that the Sr-analog Sr2VMnO6 is a ferrimagnetic half-metal, contrary to the ferrimagnetic insulator La<sub>2</sub>VMnO<sub>6</sub>.

#### SO07

#### Focused magneto-optic kerr effect spectroscopy in Ni<sub>75</sub>Fe<sub>25</sub> and Fe ferromagnetic thin films on organic substrates

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We have developed the theory of focused magneto-optic Kerr effect (MOKE) for multilayer thin films and applied it to the surface magnetic properties, which have been observed by focused MOKE, for Ni75Fe25 and Fe thin films evaporated on polyethylene naphthalate (PEN) organic substrates. We have experimentally obtained the Ni75Fe25 thickness dependence of Kerr rotations for as-deposited thin films and exposed thin films at an atmosphere for 1-year, respectively. The theoretical thickness dependences of Kerr rotations have been calculated using magneto-optic constants Q's of 0.01\*exp (- i  $48\pi/180$ ) and 0.025\*exp (- i  $57\pi/180$ ), respectively. The calculated results are in good agreement with the experimental data. Also, the Ni<sub>75</sub>Fe<sub>25</sub> thin films exposed for 1-year show larger magnetization than the as-deposited  $Ni_{75}Fe_{25}$  thin films in the range of 1-20 nm. We consider that the reason is that the surface of Ni<sub>75</sub>Fe<sub>25</sub> has transformed to locally Fe-rich surface because of the segregation of Fe. We have also measured the Kerr rotations of Fe deposited on PEN. The Q of Fe/PEN was 0.025\*exp (- i  $47\pi/180$ ). To our knowledge, the magneto-optic constants have been estimated for ferromagnetic thin films on PEN for the first time.

## POSTER PRESENTATION

#### **SO08**

The structural and magnetic properties of the magneto-caloric compounds  $Mn_{0.66}Fe_{1.29}P_{1.4}Si_x$  (x = 0.34, 0.37 and 0.42)

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The magneto-caloric effect (MCE) is the thermal response of a magnetic material to an applied magnetic field, results in a change of temperature [1,2] (Mn,Fe)2(P,Si) compounds are known for their potential for room-temperature magnetic refrigeration, due to large MCE and tunable working temperatures, cheap and abundant raw materials [3] In this study, the single phase compounds  $Mn_{0,68}$ Fe<sub>1,39</sub>P<sub>1,4</sub>Si<sub>4</sub> (x = 0.34, 0.37 and 0.42) have been synthesized. The crystal and magnetic structure have been investigated by means of neutron and X-ray powder diffraction. The refinement shows the compounds crystallize in the hexagonal Fe<sub>2</sub>P-type structure (P-62m). The 3f site is nearly completely occupied by Fe atoms, the 3g site is occupied by Mn atoms and the excess Fe atoms, while the P and Si are randomly mixed. It is shown that the average magnetic moment in the 3g site is larger than that in the 3f site 2.64 versus 1.81 µB for the compound with x = 0.34 at 10 K. The alignment of the magnetic moment has changed from almost e-direction to a-b plane with increasing Si content from x = 0.34 to x = 0.42. The compounds show a strong magnetic field induced transition, small magnetic hysteresis and large magnetic entropy changes.

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#### **SO09**

## Faraday rotation characteristics in wide-gap magnetic semiconductor ZnMnTe and ZnMnSe films

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II-VI based magnetic semiconductors with a wide optical band gap are expected to show high potential for optical applications utilizing short wavelength laser diodes, such as 532-nm green and 475-nm blue LDs [1]. ZnMnTe and ZnMnSe exhibit their absorption edges at 428-544 and 428-458 nm, respectively [2]. The edge is not so influenced by the Mn concentration, as is typically observed in CdMnTe. We have confirmed that the Faraday rotation in the ZnMnTe films deposited on quartz glass substrates is large near the absorption edge. Crystallinity was evaluated using an x-ray diffractometer. The preferred (111) growth reported previously for CdMnTe films on QG substrates was also observed in the ZnMnTe and ZnMnSe films with weaker peak intensities. This paper reports the Faraday rotation angle of those films synthesised on QG substrates by using molecular beam epitaxy with a thickness of 2 µm. A Faradayeffect signal observed for the ZnMnTe film using a 532-nm green LD has shown that ZnMnTe films are useful for green lights.We developed equipment for observing the Faraday effect directly under ac magnetic fields generated by a ring magnet. The results of a direct Faraday rotation observation sucessfully made for the ZnMnTe films under ac fields are shown.

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#### **SO10**

#### Magnetism and multiplets in Fe-phthalocyanine molecules

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A challenge to miniaturize devices for novel magnetic application now extends to treating the extreme limit of a single atom or molecule. Moreover, for molecules with transition-metals, multiplets are essential aspects of their electronic structure. It is often difficult, however, to evaluate such multiplets within a given ligand (or crystal) symmetry. Here, we carry out first-principles calculations with the FLAPW method[1] for Fe-phthalocyanine (FePc), in which the multiplets are treated by imposing a density matrix constraint on the d-orbital occupation numbers. Results predict that there are three stationary multiplets - <sup>3</sup>Eg (d<sup>2</sup>xy, d<sup>2</sup>z2, d<sup>4</sup>xz+yz), <sup>3</sup>B2g (dxy, dz2, d<sup>4</sup>xz+yz), and <sup>3</sup>A2g (d<sup>2</sup>xy, d<sup>2</sup>z2, d<sup>2</sup>xz+yz) - in a single FePc molecule, and total energy calculations demonstrate that the ground state is the <sup>3</sup>A2g. Furthermore, a columnar stacking of the FePc molecules (as seen in α-FePc) changes the ground state to <sup>3</sup>Eg due to hybridization between adjacent molecules. In qualitative agreement with recent XMCD experiments[2], the magnetic anisotropy energy for the columnar stacking structure has a large negative value of 0.6 meV/molecule, indicating that the magnetization favors pointing in the planar direction, and a large orbital moment of 0.14 uB is induced.

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#### SO11

#### Elastic anisotropy in ferromagnetic shape memory alloy of nonstoichiometric Ni,MnGa<sub>1,x</sub>In<sub>x</sub>

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In magnetic shape memory (MSM) effect, the elastic and magnetocrystalline anisotropies play an important role. Moreover, experiments reported an enhancement of the MSE effect in a non-stoichiometric metallic compounds.[1] Here, in order to discuss a role of the non-stoichiometry, we developed a method for determining elastic property of non-stoichiometric compounds from first principles calculations and cluster expansion method,[2] and applied to non-stoichiometric Ni2MnGa1.xInx. For calculations, first, the formation energies of ordered structures as functions of volume V and c/a,  $\Delta E$ o(V,c/a), were calculated by using the FLAPW method.[3] Then, by using the cluster expansion method, the formation energy of a disordered state with non-stoichiometric compositions AE o(Vc/a) were obtained in terms of many-body correlation functions Elastic properties of the bulk modulus. B. and the elastic tetragonal anisotropy. C'=(C 11-C 12)/2, were numerically obtained from the  $\Delta E d(V,c/a)$  of the disordered state. Results indicate that the B has an almost linear relation with respect to x. However, a deviation from stoichiometric composition is found to enhance the elastic tetragonal anisotropy; e.g., for a non-stoichiometric composition of x=0.125, the C' decreases by 0.4 GPa compared to that in x=0, and the C' increases when the x increases over about 0.125.

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#### **SO12**

### Structure and properties of double perovskite system La<sub>2</sub>Co<sub>1-x</sub>Fe<sub>x</sub>MnO<sub>6</sub> (x=0, 0.1, 0.2, 0.3)

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Insulating ferromagnets of type RE<sub>2</sub>AMnO<sub>6</sub> where A is transition metal and RE rare earth have received a renewed interest. Although the bulk La<sub>2</sub>CoMnO<sub>6</sub> has been extensively studied, its Fe-doped possed several questions that should be addressed. The substitution of Fe into the B site induced the increase in ferromagnetism at the cost of TC. It was also accompanied by a structure transition from orthorhombic to monoclinic. The monoclinic crystal symmetry occurred in the space group P21/n, with cell parameters a = 5.532 A, b=5.492 A, c=7.784 A, beta = 89.92°. The magnetic measurement showed that the magnetization reached maximum value of 4.7uB/ u.c. This value agrees well with the theoretically expected value.

#### SO13

### Density-functional study on spin-crossover in several Fe-based molecules

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Spin crossover (SCO) complexes are now very potential candidates for applications such as molecular switches, display and memory devices [1]. The SCO phenomenon can be qualitatively explained by the ligand field model, however, designing transition metal complexes with expected SCO behavior is still a big challenge in the field of materials science. In this paper, in order to explore more about the way to control SCO behavior of transition metal complexes, we present a density functional study on the geometric structure, electronic structure and SCO of a series of eight Fe-based molecules [2-7]. Our calculated results show that SCO in these Fe-based molecules is accompanied with charge transfer between the Fe atom and ligands. This causes change in the electrostatic energy ( $\Delta U$ ) as well as the total electronic energy of these molecules. Moreover, our calculated results demonstrate an important contribution of the interionic interactions to  $\Delta U$ , and there is the relation between  $\Delta U$  and the thermal hysteresis behavior of SCO in these molecules. These results should be helpful for developing new SCO molecules.

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#### SO14

## Magnetic properties of Cu<sub>70.9</sub>Al<sub>18.1</sub>Mn<sub>11</sub> ferromagnetic shape memory alloy

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The ferromagnetic shape memory alloy (FSMA) of nominal composition  $Cu_{70.9}Mn_{11}Al_{18.1}$  has been studied through electronic transport, dc magnetization and ac susceptibility measurements. Unlike other Cu containing shape memory alloys, Cu-Mn-Al systems of alloys show interesting magnetic behaviour. Ferromagnetism in Cu-Mn-Al alloys come from the Rudermann-Kittel-Kasuya-Yosida (RKKY) type ferromagnetic to glassy transition below martensitic transition (MT). Clear frequency shift in ac susceptibility measurement is observed near the step like anomaly present in the zero field cooled dc magnetization data, which actually indicates the onset of spin glass freezing in the sample. Isothermal magnetization (M) measurement as a function of magnetic field (H) at 5K with field cooled condition (sample cooled in presence of different magnetic field) has been performed but no shift in M(H) loop is observed. It indicates that no exchange pinning is originated in the sample due to field cooling.

#### SO15

### Magnetic properties of dioxygen molecules confined in single-walled carbon nanotubes

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Materials confined in quasi-1D single-walled carbon nanotubes (SWCNTs) exhibits unique properties that differ from the bulk counterparts [1]. Among many kinds of intercalates, a dioxygen molecule (O<sub>2</sub>), which has a spin S = 1 in the ground state, confined in SWCNTs may offer well-defined environment to investigate the physics of a quasi-1D spin-1 magnet. We present magnetic susceptibility and x-ray diffraction measurements as well as theoretical simulations on the dioxygen molecules confirmed in SWCNTs. Our XRD measurements confirmed that the O<sub>2</sub> molecules are absorbed inside the SWCNTs at low temperatures. Our magnetic susceptibility data shows the signature of the opening of a gap in spin excitations at low temperatures, which differs from bulk O<sub>2</sub> molecules. We interpret this result in terms of 1D antiferromagnets with S = 1, which is supported by our first-principles calculation suggesting that an isolated O<sub>2</sub> molecule inside a SWCNT of over 0.8 nm in diameter possess a spin S = 1. We will also discuss the condensed structure of O<sub>2</sub> inside SWCNTs on the basis of our XRD data and molecular dynamics simulations [2].

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#### SO16

## Magneto-elastic coupling and magnetocaloric effect in $Fe_2P\mbox{-}based$ Mn-Fe-P-Si compounds

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Structural, magnetic and magnetocaloric properties of Fe<sub>2</sub>P-based Mn-Fe-P-Si compounds were investigated. The study reveals a strong magneto-elastic coupling that starts to develop in the paramagnetic state and grows when the ferromagnetic transition temperature is approached. This magneto-elastic coupling over a large temperature range around the transition temperature results in a first-order phase transition with abrupt changes in lattice parameters. Moreover, a correlation is observed between the hysteresis and the discontinuous changes in lattice parameters at the transition temperature, which is attributed to competition between moment formation and bonding [1], contributes to a reduced hysteresis. We also show the thermal evolution of the magnetic moments. Interestingly, the magnetic moments on the 3f site are formed and gradually increase in the paramagnetic state, and suddenly jump to a larger value when magnetic order is established.

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#### SO17

#### Magnetorefractive effect in manganites

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We studied a magnetic-field-induced change in reflectivity (magnetoreflection) and transmission (magnetotransmission) of natural light in La<sub>1-x</sub>A<sub>x</sub>MnO<sub>3</sub> (A=Ca,Sr,K,Ag) manganites possessed the colossal magnetoresistance effect (CMR)[1-3]. It was showed the magnetoreflection and magnetotransmission effects in manganites are an optical response to the CMR in the IR-region. Effects can reach tens of percents in the field of 0.3-1 T near the magnetic phase transition temperature. The observed phenomena are connected with the change of ratio between the localized and delocalized charge carriers under the magnetic field applied [2]. At the same time, there is no strict correlation between these effects and CMR in the visible range. The observed phenomena are connected with the alteration of the optical density under the magnetic field in the range of interband transitions [3]. Magnitude of effects is one order as less as that in the IR-region. The nature of observed effects in manganites can be satisfactory explained within the framework of the theory of magnetorefractive effect [4]. The results of study of magnetoreflection and magnetotransmission in manganites may be proposed for creation of different magnetic and electronic sensors and light modulators.

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#### **SO18**

#### Electroresistance and joule heating effects in manganite thin films

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Electroresistance (ER), i.e. electric field- and/or current-induced resistance switching, has attracted much attention recently because of the possibility of using it for the implementation of resistance random access memories (ReRAM). Although ER is a quite common phenomenon in transition metal oxides, that has been extensively studied both theoretically and experimentally, the precise mechanism involved is not clear yet. In this work we report on the ER measurements in patterned La<sub>23</sub>Sr<sub>23</sub>MnO<sub>3</sub> (LSMO) thin films prepared by sputtering. In order to analyze Joule heating effects we have evaporated a Pt layer on top of the LSMO path to have access to the actual temperature of the sample while measuring resistance of LSMO path or I(V) characteristic curves. I(V) curves have been measured at different temperatures and the corresponding resistance values are compared with that of the R(T) curve taking into account the actual temperature of the sample in order to clarify the role of Joule heating in the observed change of the resistance.

#### SO19

#### Magnetocaloric effect and other properties of cold rolled Gd ribbons

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The number of papers devoted to magnetic refrigeration, magnetocaloric materials and high-field magnetic systems has grown at an exponential rate in the past 10 years (K.A. Gschneidner et al., 2008) [1]. Since First International Conference on Magnetic Refrigeration at Room Temperature many researches in different tourhies mobilize their efforts in the field of constructing commercial magnetic refrigerator. There are numerous different magnetocaloric alloys such as La-Fe-Si, Gd-Si-Ge, Ni-Mn-X (X = Ga, In, Sn, Sb) [2-3]. Usually the magnetic refrigerant materials are either packed in a magnetocaloric bed as small spheres attached around a disk as a ribbon or separated in a pile of thin, equally spaced sheets. Thus one the ways of engineering MCE materials is tightly connected with preparing very thin (a few microns) ribbons of high value MCE alloys with good mechanical properties. At present rapid solidification is the main technique for producing this kind of materials [4]. In our work we investigate magnetocaloric effect, magnetic and mechanical properties of Gd cold rolled samples of different thickness. Both theoretical and experimental aspects of the investigation are discussed.

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#### **SO20**

**Fabrication and properties of double perovskite SrLaVRuO**<sub>6</sub> Ryosuke Zenzai\*, Tetsuya Miyawaki, Kenji Ueda and Hidefumi Asano *Crystalline Materials Science, Nagoya University, Japan* 

Half-metallic antiferromagnets (HMAF) exhibiting both fully spin-polarized charge transport and zero net magnetic moment are expected to provide not only a fertile playground for fundamental research but also new spintronic devices. Although the theoretical calculations predicted candidates for HMAFs, experimental searches for HMAFs have so far been unsuccessful. This paper reports on the fabrication and properties of SrLaVRuO<sub>6</sub> polycrystalline bulks and epitaxial films. SrLaVRuO<sub>6</sub> bulks were synthesized by a solid-state reaction at 1100~1350 °C in Ar + (0~3) % H<sub>2</sub> gas. Structural characterization revealed that the main phase of the bulk samples was of a double perovskite crystal structure with lattice parameters of 0.7888~0.7905 nm. However, all the bulks contained the impurity phase (Sr<sub>3</sub>V<sub>2</sub>O<sub>8</sub>), which is a nonmagnetic semiconductor. The bulks exhibited magnetization values of 10-3 µB per formula unit at 77 K, negative paramagnetic Curie temperature values and metallic temperature dependence of the resistivity. These properties are assumed to be intrinsic to double perovskite SrLaVRuO6 by considering the nonmagnetic semiconducting nature of Sr<sub>3</sub>V<sub>2</sub>O<sub>8</sub>. Moreover, SrLaVRuO<sub>6</sub> thin films were prepared on LaAlO<sub>3</sub>, SrTiO<sub>3</sub> (001) substrates by magnetron sputtering. Structural characterization revealed that these thin films were grown epitaxially on these substrates, and phase pure double perovskite.

J. H. Park, S. K. Kwon, and B. I. Min, Phys. Rev. B 65 (2002) 174401.

#### SO21

Korea

#### Giant magnetoresistive effect in non-magnetic silicon

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Recently magnetoresistive effects in non-magnetic semiconductors, particularly silicon, have gained considerable attention owing to the large magnitude of the effect, which is comparable, or even larger than, that of the performance of commercial giant-magnetoresistance devices. The underlying mechanism responsible for the giant magnetoresistance in silicon, however, has not yet been fully explored. Here we report that a simple device, based on a lightly doped silicon substrate contacted with two indium contacts, shows a positive magnetoresistance over 10,000 per cent for magnetic fields between 0 and 1 T at low temperature of 3 K. Current-voltage characteristics exhibit a nonlinear behavior, which is highly sensitive to temperature and magnetic field. We have extensively studied magnetoresistive effect in silicon with varying bias voltage, temperature, electrode spacing, and the relative angle between the substrate and the applied magnetic field. Since our device is based on a conventional silicon platform and is highly sensitive to low magnetic field, it could be used to develop new devices of silicon-based magnetoelectronics.

#### **SO22**

## Effect of Fe substitution on Ni-Mn-In shape memory alloys: Magnetic and magneto-structural properties

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The Ni-Mn based ferromagnetic shape memory alloys have been extensively investigated for their multifunctional properties emerging by suitable chemical substitutions. The off-stoichiometric  $N_{ig}Mn_{2:h_s}I_{n_{2:h_s}}$  exhibit martensitic transition for In concentration in the range 13-16 at%. In the  $N_{ig}Mn_{2:h_s}I_{n_{2:h_s}}$  alloy, the ferromagnetic transition occurs at Curie temperature TC = 305 K with a concommittant austenite to martensitic transition. A large inverse magnetocaloric effect and anomalous magnetic behavior have earlier been reported for this alloy [1]. The present work reports the effect of F-substitution (05 to 2 at%) on the magneto-structural transition and magnetic behavior of the  $N_{ig}Mn_{3:h_1}$  alloy. The martensitic transition goes down to 255 K for 1at% Fe alloy with strong hysteresis between zero field-cooled (ZFC) and field-cooled (FC) data while it nearly disappears for the 2at% Fe alloy. The 0.5 at% Fe substituted alloy shows a temperature dependent field-induced metamagnetic transition between 250-220K while this transition is not observed up to fields of 6T in 1 at% Fe substituted alloy. A thermal hysteresis in the temperature dependence of resubstitution (05 a 2 at%) for substitution 07 a at% Fe substitution of 2 at% Fe substituted alloy.

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#### SO23

### Magnetic properties and magnetocaloric effect in shape memory alloys Ni-Mn-Ga

Rafael Fayzullin, Vasiliy Buchelnikov, Sergey Taskaev and Mikhail Drobosyuk Chelyabinsk State University, Russia

The magnetocaloric effect (MCE) is the ability of magnetic materials to heat up or cool down when placed in or removed from an external magnetic field. The magnetic materials with large values of MCE can be applied in the magnetic refrigeration technique [1]. Recent experimental studies have shown that Ni-Mn-Ga Heusler alloys are also attractive for the application in magnetic refrigeration [2]. In this work the magnetic and magnetocaloric properties of the ferromagnetic shape memory alloys Ni<sub>21-x</sub>Mn<sub>1</sub>.  $_{x}$ Ga (x=0.07; 0.08; 0.09) were experimentally studied. The MCE measurements were performed by the setup produced by AMT&C [3]. In this setup, the adiabatic temperature change  $\Delta$ Tad was measured by a direct method with help of a thermocouple. Magnetic field up to 2 T was created by Halbach permanent magnet. The magnetic field strength was measured by Hall probe. Signals from the thermocouple and Hall probe were recorded simultaneously what allowed us to measure  $\Delta$ Tad as a function of magnetic field H. All alloys exhibit the positive MCE typical near the second-order phase. Our MCE measurements of Ni<sub>21-x</sub>Mn<sub>1-x</sub>Ga have shown that the maximal MCE is observed near the Curie point. The maximal values of MCE are 1.2 - 1.6 K near the room temperature.

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#### **SO24**

## Magnetic susceptibility avalanches in thermally-induced first-order phase transition of La(Fe\_{0.88}Si\_{0.12})\_{13} magnetocaloric compound

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Magnetic materials with large magnetocaloric effects due to the first-order phase transition attract attention as candidates for room temperature magnetic refrigerants. NaZn13type La(Fe,Si1,,)13 compounds exhibit thermally-induced first-order ferromagnetic (F) paramagnetic (P) transition at the Curie temperature Tc and the itinerant-electron metamagnetic transition takes place above Tc [1]. Evolution of the transition determines the operating speed of magnetic refrigeration, therefore, a detailed observation is necessary. In the present study, appearance of the avalanche behavior [2] in the progress of thermally-induced transition has been examined by measuring the magnetic susceptibility  $\chi$ . The value of  $\chi$  upon cooling at 0.4 K / min. exhibits a smooth variation across the transition and its first derivative, dg/dT, shows no jerky change. Meanwhile, the avalanche-like behavior is observed in dg/dT around Tc upon heating. Furthermore, the number and height of avalanches are dramatically changed by external magnetic fields. Generally, the avalanche behavior is induced by the quenched disorders, however, such a mechanism can not explain the influence of magnetic field. Considering the fact that the kinetics of the transition is influenced by the demagnetization effect [3], these results indicate that the nucleation- growth behavior is affected by the magnetic dipole interaction.

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#### SO25

#### Phase coexistence and magnetocaloric effect on martensitic transition in Ni-Co-Mn-Sn metamagnetic shape memory alloy

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Since the discovery of the metamagnetic shape memory effect in Ni<sub>ex</sub>Co<sub>3</sub>Mn<sub>36</sub>, In<sub>133</sub> Heusler alloys [1], a member of studies on magnetocaloric effect (MCE) have been made in metamagnetic shape memory alloys. It is found that metamagnetic shape memory alloys exhibit large MCE in the vicinity of martensitic transformation. But for the first-order transition, whether and to what extent the Maxwell relation is applicable are problems requiring further study [2]. In this paper, we will discuss the appropriate approach to evaluate the MCE in Ni-Co-Mn-Sn alloys. The energy dispersive spectrometer shows that the sample composition is Ni<sub>6.2</sub>Co<sub>1</sub>,Mn<sub>6.8</sub>Sn<sub>119</sub>. The thermo-magnetic curves indicate the martensitic transformation can be driven by the magnetic field. It is observed that metamagnetic transition takes place in a wide temperature range from the magnetization isotherms. The entropy change calculated by Maxwell relation is up to 301/kgK under a field change of 0-5T. Meanwhile, we calculated the entropy change based on heat capacity measured under 0 and ST, and an entropy change of ~31/ kgK is obtained for a field change of 0-5T. According to the different results from magnetic and calorimetric measurements, we find that the Maxwell relation should be used with caution on martensitic transformation including phase coexistence.

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#### SO26

Balance between the growth rate of ferromagnetic phase and demagnetizing fields in itinerant electron metamagnetic transition of La(Fe<sub>0.88</sub>Si<sub>0.12</sub>)<sub>13</sub> Hitomi Yako\* and Asaya Fujita *Tohoku university. Japan* 

Strong interplay between an itinerant electron metamagnetic (IEM) transition and magnetocaloric effects in NaZn<sub>13</sub>-type La(FexSi1-x)13 has been investigated.[1-3] Recently, we have demonstrated that progress of the IEM transition is influenced by the demagnetization effect in La(Fe<sub>0.88</sub>Si<sub>0.12</sub>)<sub>13</sub>-[4] Especially, the progress is arrested in midstream even under the magnetic fields, µH larger than the critical value, µHe by setting longitudinal direction of the sample perpendicular to the field direction ( $\theta = 90$  degrees). In the present study, the angular dependence of the transition progress is examined and the influence of demagnetization effect on the IEM transition is discussed. The time-dependent magnetization measurements were carried out at 200 K under the constant value of µH by changing  $\theta$  from 0 to 90 degrees. The transition starts to proceed under the magnetic field larger than µHs. However, the transition is arested in midstream, and the completion of transition emerges only after µHc exceeds the threshold value, µHt. The values of µHs and µHt are evaluated to be 0.60 and 0.75 T for  $\theta = 30$  degrees and 0.70 and 1.00 T for 90 degrees, respectively. The growth rate of ferromagnetic phase is dominated by keeping the balance with demagnetizing fields.

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#### SO27

#### Temperature dependent structural and magnetic properties of Ni-Mn-In Heusler alloy glass-coated microwires

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Magnetic shape-memory alloys (MSMAs) attracted growing attention owing to shape memory effect, originated from the coupling between magnetic and structural ordering. From point of view of technological applications, miniaturizing of MSMA-based devices using MSMA particles, wires, nbbors, films, bi- and multilayers, pillars is quite important. Here we report on studies of magnetic properties of non-stoichiometric Heuslertype Ni<sub>8</sub>Mn<sub>34</sub>, ln<sub>155</sub>, Ni<sub>68</sub>Co<sub>5</sub>Mn<sub>34</sub>, ln<sub>155</sub>, Ni<sub>68</sub>Co<sub>5</sub>Mn<sub>54</sub>, ln<sub>155</sub>, Ni<sub>68</sub>Co<sub>54</sub>, ln<sub>156</sub>, ln<sub>156</sub>, Ni<sub>68</sub>Co<sub>54</sub>, ln<sub>156</sub>, ln<sub>156</sub>, Ni<sub>68</sub>Co<sub>54</sub>, ln<sub>156</sub>, ln<sub>156</sub>, Ni<sub>68</sub>Co<sub>54</sub>, ln<sub></sub>

#### SO28

#### Magnetooptical Kerr effect enhancement in Co/TiO<sub>2</sub> layered films Victor Polyakov<sup>1</sup>, Konstantin Patrin<sup>1</sup>, Klaudia Polyakova<sup>1\*</sup>, Vitaly Seredkin<sup>1</sup> and

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As is known, the magneto-optical properties of the composite metal-dielectric systems are determined by the volume of a magnetic phase and the properties of a dielectric matrix or dielectric layers. Of special interest are the magneto-optical properties of the composite films containing TiO<sub>2</sub> whose permittivity is higher as compared to SiO<sub>2</sub> and Al<sub>2</sub>O<sub>3</sub>. Here, we report the results of the magneto-optical study of Co/TiO<sub>2</sub> multilayers. Particularly we studied spectral dependence of the polar Kerr effect of Co / TiO<sub>2</sub> layered structure of with various number of layers. The composite Co/TiO<sub>2</sub> periodic structures consisting of 2–12 pairs of layers (n) were prepared by sequential deposition of Co and titanium oxide by ion-plasma sputtering and reactive ion-plasma sputtering, respectively, onto glass substrates. The respective thicknesses of magnetic and titanium oxide layers were 5 and 17 nm. Study of spectral dependence of polar Kerr effect showed that the spectra of the Co/TiO<sub>2</sub> layered structures are resonance, with the large growth of the angle of rotation near the resonance. The maximum angle of rotation nonmonotonically depends on the number of layers. The maximum Kerr rotation angle (20k=7.3 deg) observed at a wavelength of 540 nm corresponds to n = 8.

#### SO29

## Electronic structure calculations and the magnetic properties of $Ru_2FeZ$ (Z = Al and Ga)

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After the discovery of half-metallic character in NiMnSb by de Groot et al. Heusler alloys have drawn much attention due to the possible application in magneto-electronic devices [1,2]. In the quest of finding half-metallic character in 4d transition element based Heusler alloy we present the electronic structure and magnetic properties of Ru<sub>2</sub>FeGa/Al. The band structure calculations were done using the state-of-theart density functional technique with the PWSCF code [3]. The equilibrium lattice parameter was found by fitting energy vs volume data to the Murghnan equation of state. At the equilibrium lattice parameter the total energy difference between the nonmagnetic state and ferromagnetic state was observed to be of the order of 5 meV which suggest that these systems are having stable magnetic ground state. In ferromagnetic ground state the lattice parameter was observed to be 5.988 A for Ru<sub>2</sub>FeGa and 5.982 A for Ru<sub>2</sub>FeAl. From the DOS Vs Energy diagram it is found that the systems are not half-metallic in nature. However, the calculation showed very high spin polarization at the Fermi level. The spin polarization (P) value was found to be 83 % in both the cases.

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#### **SO30**

#### Disorder effects in giant magnetocaloric materials

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In striving for mass production of magnetocaloric materials for refrigeration applications, one has to face the fact that these will not be laboratory-grade materials, and will necessarily present composition gradients due to lower quality reagents. Some consequences are expected: the maximum value of magnetic entropy change ( $\Delta$ SM) should decrease compared to a "pure" material, and broadening of the  $\Delta$ SM(T) curves, as observed in elementary ferromagnets [1]. Some theoretical work has focused on this topic, for second-order phase transition systems [2-5]. Still, these theoretical considerations do not directly apply to giant magnetocaloric materials. We present a study on the effect of disorder on the magnetic and magnetocaloric materials of first-order phase transition systems, via the use of the Bean-Rodbell [6] model. Disorder is shown to "smooth" the discontinuities on magnetization and  $\Delta$ SM, and also affect magnetic hysteresis. For sufficiently large disorder, the  $\Delta$ SM(T) curves resemble the distribution function. We discuss how the magnetic field dependence of  $\Delta$ SM is affected by disorder, in light of a "second-order like" dependence of  $\Delta$ SM on applied magnetic field, for disordered LaFeSi samples [5]. Aging effects on hydrogenated LaFeSi samples, that degrade even in storage [7], are discussed taking into account our theoretical simulations.

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#### SO31

#### On the Curie temperature dependency of the Magnetocaloric effect

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A theoretical study of the magnetcaloric effect dependency on simple ferromagnetic microscopic parameters is here presented. Considering the Bean-Rodbell model [1] of magnetovolume coupling within the Weiss mean field theory, the magnetic entropy change of systems undergoing first and second order magnetic phase transition was numerically simulated. Through these simulations, the magnetic entropy change behavior was studied as a function of: Tc , spin value, applied field changes and the nature of the transition considered. The main result found for both first and second order magnetic transitions systems is the linear dependency of the  $\Delta$ SM(T) maximum value on Tc<sup>23</sup>. By starting from the magnetization mean field state equation an approximated expression for the magnetic entropy change as a function of microscopic magnetic parameters was reached, confirming the Tc<sup>23</sup> dependency. We have found that many ferromagnetic systems, both undergoing second (GdTbCo [2], RR'AI [3] (R = rare earth), RCoMn [4]) and first-order (MnFePAs [5], GdSiGe [7], RCo [7]) phase transitions, obey this relation between the magnetocaloric effect and Tc.

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•			Akama, Takamitsu	PG07	143 Amaral, Vitor S.	KB06
Α			Akatsu, Mitsuhiro	RB05, RD07	254, 263 Amarande, L.	SA11
Abbamonte, P.	PH10	148	Akatsuka, Hiroyuki	PB34	125 Amato, Alex	IA02, PI06, QB
Abbasian Arani, A. A.	PO30	182	Akbari, Alireza	PD06	132 Amemiya, Kenta	BH06
Abbasian Arani, Ali Akbar	PO28	181	Akdogan, Ozan	CB05, KB04	49, 108 Amitsuka, Hiroshi	PD07, QC25, S
Abdel - Hafiez, M.	QF22	207	Akerman, Johan	CE01, CE02, PK13	51, 162 Amyan, Adham	QA11, QJ11
Abdul Ahad, Faris B.	RK07	287	Akerman, Johanna	KE01	110 Analytis, James	HA05
Abdullah, Nor Hapishah	BJ06	48	Akhavan, M.	QP23	249 Anand, Vivek Kumar	QB14
Abe, Kazutaka	IF02	94	Akimitsu, Jun	DF03, PB06, PB07, PC21,	60, 121, 130, 143, Anane, Abdelmadjid	IE04
Abe, Nobuyuki	QA10, QA13, QA17	184, 175, 185		PG09, QG22, QG28,	211, 212, 276, 337, Andersen, Ole Krogh	SC07
Abe, Taro	SM09	362		RH13, SF02, SG16	342 Andersson, Gabriella	EI03
Abid, Mohamed	PJ14, SH13	157, 345	Akiyama, T.	RM04	295 Andersson, Yvonne	PM10
Abo, Gavin Sky	JH04	104	Akiyama, Toru	RM05, SO10, SO11	296, 368 Ando, K.	CE03, ID02, RJ
Abramova, G. M.	SG20	342	Akutsu, Yasuhiro	RI01	279 Ando, Kazuya	AI02
Abrosimova, Galina	SO04	367	Alameda, Jose Maria	GG03	80 Ando, Ken	KF01
Abrudan, Radu	EC04	66	Alaria, Jonathan	PI04	151 Ando, Koji	BI01
Acquarone, Marcello	SI10	347	Alayo, William E.	SM17	363 Ando, Y.	SK24
Aczel, A. A.	GF01	79	Alayo, Willian	PL09, QM14	165, 233 Ando, Yasuo	CH02, El02, EJ
Adachi, Hiroto	FA03	72	Alba Venero, D.	QM06	232	RL13, RO05, R
Adachi, Nobuyasu	KC02	108	Albert, Barbara	PI05	151	SI17, SK08
Adachi, Yoshiya	PG04, QI05	143, 218	Albertini, Franca	AG02	37 Ando, Yoichi	HI01, RC11
Adam, Philip	CH01	54	Albrecht, Michele	RK05	286 Ando, Yuhki	SA03
Adam, Sorin	PM06	168	Albrecht, Ole	FD02	74 Andrade, Bruna	QL01
Adamo, Carolina	SA08	317	Aldus, Robert	PO12	179 Andre, Gilles	JB05, RA11, RI
Adams, M	SE05	334	Aleshkevych, Pavlo	QM17	234 Andreev, A. V.	QG07, QI03, Q
Adams, T.	PI09	152	Alfadhel, Ahmed	PM27	171 Andreev, A.v.	GJ03
Adams, Tim	DF02	60	Alfahad, S.	RN18	299 Andreev, Alexander	RE09, SG03
Adenot Engelvin, Anne Lise	AJ01	39	Alff, Lambert	PI05	151 Andreev, Alexander V.	QI19
Adiga, Shilpa	QA23	186	Algarabel, P. A.	KC05, PJ23	109, 158 Andrei, Gloskovskii	SD15
Adriano, Cris	SE15	335	Algarabel, Pedro Antonio	SM16	363 Andreica, Daniel	CD04, SE08
Adroja, D T	KA02, SE03, SE05, SE13,	107, 333, 334, 335,	Ali, Md. Ehesan	BC02	42 Angelakeris, Makis	RP15
	SF09	338	Ali, Naushad	CH01	54 Angst, Manuel	QA23
Adroja, Devashibhai	PC12, SD13	128, 332	Aliakbari Miyan Mahaleh, M.	PO30	182 Anh Nguyen, T. N.	PK13
Adroja, Devashibhai T.	SE20	336	Alidoust, N.	SD04	330 Anh Tuan, Duong	PJ15
Adschiri, Tadafumi	RN11	298	Aliouane, Nadir	PA10	117 Anh, D. T. K.	QG11
Aeppli, Gabriel	PJ17	157	Alireza, Patricia	PC02, RQ09	127, 308 Anil Kumar, P. S.	PH04
Afzali, Reza	SC09	327	Alireza, Patricia L	SC16	328 Anisimov, M. A.	QC22, RH07
Agbelele, Arsene	QA27	187	Allred, Jared M	QB13	189 Annal Therese, Helen	QA25
Agrestini, Stefano	BF03	45	Al-mahdawi, Muftah	PJ11	156 Annett, James	JA05
Agterberg, Daniel	RB26	258	Alnassar, Mohammed	SA02	316 Antohe, Stefan	QK20
Agui, Akane	PM02	167	Alonso, J	QM06	232 Antonelli, Alex	RB18
Ah, Chil Seong	RL21	293	Alotibi, M.	RN18	299 Antos, Roman	KD03
Ahmad, Dawood	QB25	191	Alqahtani, Mohammed S.	RN18	299 Antropov, Vladimir	SC07
Ahmad, Ehsan	EH04	70	Al-senany, Nourh A.	PG01	142 Anwar, M S	PN15
Ahmed, Faheem	PM14, PN15	169, 174	Alshammari, Marzook S.	RN18	299 Anzai, Shuichiro	SE09
Ahmed, Mohamed Abbas Ali	RN21	300	Altarawneh, M. M.	SC03	326 Aoki, Dai	HP11, JC02, PI
Ahn, Chan-woo	RR31	316	Altarawneh, Moaz	BA01, EA03	40, 64	SF03
Ahn, Cheol-woo	DJ02	63	Althoyaib, S	QL21	231 Aoki, Haruyoshi	QC10, QC11
Ahn, Ho-seong	PM19, PM21	170	Altshuler, Tatiana Semenovna	PE09	137 Aoki, Hideo	EB02, RI14
Ahn, Jae Young	QP18	248	Alvarez, Lorena	Q026	245 Aoki, Kento	QA05
Ahn, Kyo-hoon	SO05	367	Alvarez, Pablo	JJ04	106 Aoki, Yuji	RD17
Ahn, San	SC15	328	Alves, E.	QM13	233 Aoshima, Ken-ichi	SK21
Ahn, Sora	RL17	292	Alyamani, A.	RN18	299 Aoyagi, Mitsuharu	RP17
Aichhorn, Markus	SG22	343	Alzayed, Nasser	RB29	258 Aoyama, Satoshi	PH20
Ain, Qurat-ul	PB13	122	Amakai, Yusuke	PB08, QA19, QJ04,	121, 185, 221, 263, Aparnadevi, M	RA08
Aizawa, Kazuya	RH12	276		RD06, RD22, SE06	265, 334 Apinaniz, Estibaliz	QH16
Akagi, Yutaka	PG23	146	Amako, Yasushi	QH04, RD19	213, 265 Aplesnin, Sergey	QG05
Akahoshi, Daisuke	QE20, QF21	202, 207	Amaral, Joao S	RI18, SO30, SO31	281, 371 Aplesnin, Sergey Stepanovic	
Akaki, Mitsuru	PA06, PA19, QA05,	116, 119, 183, 250,	Amaral, V. S.	RH02	275 Aprilia, Ely	RC02
	RA06, SB24	324	Amaral, Vitor S	RI18, SO30, SO31	281, 371 Arac, Erhan	Q024

Arai, Hiroko Arai, Masatoshi Arakawa, Akira Arakawa, Naoya Araki, Koji Araki, Shingo Araki, Yusuke Aramberri, Hugo Arango, Y. C. Araoka, Nobutaka Araujo, Joao P Araujo, M. A. N. Araujo, Miguel Arbelo Jorge, Elena Ardavan, Arzhang Arenholz, Elke Arepalli, Sivaram Arguello, C. Argyriou, D. Argyriou, Dimitri Argyriou, Dimitri N. Argyriou, Dimitrios Arhan, H. Z. Arikan, Mustafa Arima, Takahisa Arima, Taka-hisa Arinicheva, Olga Arita, Masashi Arita, Ryotaro Arnold, Z. Arnold, Zdenek Aronin, Alexander Arora, Sunil Kumar Arredondo, Yesenia Arshi, Nishat Artyukhin, Sergey Arvanitis, D. Arzhnikov, Anatoly Asada, Hironori Asai, Shinichiro Asai, Takuya Asai, Yugo

JB05, RA11, RD14

CE03, ID02, RJ07

CH02, EI02, EJ01, KD04,

IA02, PI06, QB22

PD07, QC25, SE07

108

318

47

84

189

94

326

71

169

39

110

47

354

89, 260

100, 251, 264

316

228

82

220

332

305

186

162

157

210

147

186

99

227

257

109

326

174

334

337

193, 194

65, 281

183

264

353

306

150

250

215

209

259

244

139, 367

HP11, JC02, PD17, RE10, 30, 100, 134, 268,

195, 275

51, 334

267, 340

RL13, RO05, RO06, SI03, 301, 346, 348, 351

QG07, QI03, QI09, RD26 209, 218, 219, 266

51, 93, 283

54, 70, 71, 110, 292,

90, 151, 190

133, 196, 334

184, 222

Asaka, Souta Asaki, Kenji Asano, Hidefumi

Asano, Hidehumi Asano, Takayuki Asaoka, Rui Aschenbach, Konrad Hsu Ashida, Takuya Aslibeiki, B. Aso, Naofumi Asokan, K Asokan, K. Assmus, Wolf Asthana, Saket Atkinson, Del

RJ06, SI16 283, 348 DA01, RH12 56, 276 SL02 356 333 SE04 254 RB05 HC02, PE15, QC18, SE10 84, 138, 195, 334 222 QJ12 SD16 332 QF22 207 SB14 322 371 SO31 323 SB21 258, 271 RB25, RF13 54 CH04 SB20 323 62 DH01 PJ32 160 GF01 79 RB17 256 RB01, RB07 254, 255 PA10 117 CI01 54 GA01 75 GH03 81 145, 196, 242 PG20, QD02, QO13 IA04, PA12, PG14, 91, 117, 144, 184, QA10, QA13, QA17 175, 185 223 QJ13 265 RD20 AE05, QC04 36, 192 158 PJ23 169 PM11 367 SO04 149 PH16 211 QG21 174 PN15 54 CI01 58 DD02 45 BF04 PG06, QM10, RP12, 143, 233, 305, 352 SK13 PH12 148 RD08 263 PM18 170 BG04 46 SG12 341 PJ12, PL12, RC18, SA07, 156, 165, 261, 317, SA17, SO20 319, 370 PB34, SK18 125, 353 PG17, QF11, QF14 145, 205, 206 RF11 270 EG01 69 PB25, PL10 124, 165 PI23 154 PD14, QC11 134, 194 PG27 146 Q027 245 201 QE14 PN18 175 Q024 244

ex	
Author Index	

Attane, Jean-philippe	Al04, GE01, QK15, QK16, QK17, RL24, SL24	39, 78, 226, 293, 360	Bak, Zygmunt Baker, Peter J	SB11 PN23	322 175
Atxitia, U.	GD03	78	Baker, Peter. J.	QG19	211
Atxitia2, U	FC01	73	Bakr, Mohammed	CA03	49
Au, Yat-yin	EH04, HD03	70, 85	Balaev, Alexander	PI11	152
Augendre, Emmanuel	AI04	39	Balaev, Dmitry	SM19	364
Augstine, Saji	Q027	245	Balagurov, Leonid A	PN22	175
	QL02	228	-		
Augustine, Tan Lee			Balakhonov, Sergey	ED06	67
Augustine, Tan Tuck Lee	PN03	172	Balakirev, Fedor	RB21	257
Avanesyan, Grisha	ID03	93	Balakrishnan, G.	CC04, II02	50, 97
Avdeev, M.	SO08	368	Balakrishnan, Geetha	GF03, SC06, SD13	79, 326, 332
Avdeev, Maxim	IJ03	97	Balatsky, A. V.	QB08	188
Avdeev, Maxim V.	SC25	329	Balatsky, Alexander	KG01	111
Avella, Adolfo	RG18	274	Balatsky, Alexander V	EB04, SE18	65, 336
Aviani, Ivica	IA02	90	Balbashov, A. M.	PA15	118
Avignon, Michel	QG21, QO26	211, 245	Balcells, Lluis	IH01, SO18	95, 369
Avila, M A	QJ05	221	Balcerzak, Tadeusz	RI03, RI09	279, 280
Avila, Marcos A.	Q003	241	Baldassarre, Leonetta	SA08	317
Avramov, Pavel V	BH05, BH06	46, 47	Baldes Bango, Fernando	GG03	80
Awan, Muhammad Saifullah	SA15	318	Baldovi, J.	BC01	42
Awana, V. P. S	QM24	235	Balents, Leon	JD03, PL08	101, 165
Awano, Hiroyuki	GE04, JE04, PJ09, QP20	78, 102, 156, 248	Balicas, Luis	CC02, EA03, JA02, RB19	50, 64, 98, 257
Awschalom, David D.	Al01	39	Balke, Benjamin	QN19	239
Ayala, Oscar	QB13	189	Ballou, Rafik	AC04	35
Ayala-valenzuela, Oscar	GC03	77	Baltz, Vincent	AI04, BD02, DD01	39, 43, 58
Azarevich, Andrey	QC20	195	Bance, Simon	DE01	59
Azuma, Masaki	QE10	201	Bando, Hiroshi	PM13	169
			Bandou, Hisamitu	PC05	127
В			Banerjee, Animesh	CG05	53
3., Parvatheeswara Rao	BJ05	48	Bang, Do	JE04, SK10	102, 352
Baba, Ryo	RR05	311	Bang, Dong H.	CJ03	55
Babaei Brojeny, Aliakbar	PC16	129	Bang, Ngac An	AJ06	40
Babic, Michal	вноз	46	Banholzer, Anja	RJ11	284
	01105	40	barinoizer, Arija		350
	DNOF	207	Describer on Matalian		
Bach, Long Giang	RN05	297	Bannikova, Nataliya	SK02	
Badini Confalonieri, Giovanni A.	GG04	80	Bansil, A.	HI02	89
Badini Confalonieri, Giovanni A. Bae, Dong Sik	GG04 PH21	80 150	Bansil, A. Bansil, Arun	HI02 BA04	89 41
Badini Confalonieri, Giovanni A. Bae, Dong Sik Bae, Joohyung	GG04 PH21 QK12	80 150 226	Bansil, A. Bansil, Arun Bansil, B.	HIO2 BAO4 SDO4	89 41 330
3adini Confalonieri, Giovanni A. 3ae, Dong Sik 3ae, Joohyung 3ae, Jungchan	GG04 PH21 QK12 PM17	80 150 226 170	Bansil, A. Bansil, Arun Bansil, B. Bao, Guizhi	HI02 BA04 SD04 PC20, QJ07	89 41 330 130, 222
3adini Confalonieri, Giovanni A. 3ae, Dong Sik 3ae, Joohyung 3ae, Jungchan 3ae, Kyoung-hoon	GG04 PH21 QK12 PM17 RO02	80 150 226 170 301	Bansil, A. Bansil, Arun Bansil, B. Bao, Guizhi Bao, Jinke	HI02 BA04 SD04 PC20, QJ07 RB13	89 41 330 130, 222 256
3adini Confalonieri, Giovanni A. 3ae, Dong Sik 3ae, Joohyung 3ae, Jungchan 3ae, Kyoung-hoon 3ae, Moon Seob	GG04 PH21 QK12 PM17 RO02 RP23	80 150 226 170 301 307	Bansil, A. Bansil, Arun Bansil, B. Bao, Guizhi Bao, Jinke Bao, Nina	HI02 BA04 SD04 PC20, QJ07 RB13 DD03	89 41 330 130, 222 256 58
Badini Confalonieri, Giovanni A. Bae, Dong Sik Bae, Joohyung Bae, Jungchan Bae, Kyoung-hoon Bae, Moon Seob Bae, Seok	GG04 PH21 QK12 PM17 RO02 RP23 RH08	80 150 226 170 301 307 276	Bansil, A. Bansil, Arun Bansil, B. Bao, Guizhi Bao, Jinke Bao, Nina Bao, Wei	HI02 BA04 SD04 PC20, QJ07 RB13 DD03 RB07	89 41 330 130, 222 256 58 255
3adini Confalonieri, Giovanni A. 3ae, Dong Sik 3ae, Joohyung 3ae, Jungchan 3ae, Kyoung-hoon 3ae, Moon Seob	GG04 PH21 QK12 PM17 RO02 RP23	80 150 226 170 301 307 276 105	Bansil, A. Bansil, Arun Bansil, B. Bao, Guizhi Bao, Jinke Bao, Nina Bao, Wei Bar-ad, S.	HI02 BA04 SD04 PC20, QJ07 RB13 DD03	89 41 330 130, 222 256 58 255 43
Badini Confalonieri, Giovanni A. Bae, Dong Sik Bae, Joohyung Bae, Jungchan Bae, Kyoung-hoon Bae, Kyoung-hoon Bae, Moon Seob Bae, Seok Bae, Taejin Bae, Yu Jeong	GG04 PH21 QK12 PM17 RO02 RP23 RH08 JI02 CB01, PJ17, QP18	80 150 226 170 301 307 276	Bansil, A. Bansil, Arun Bansil, B. Bao, Guizhi Bao, Jinke Bao, Nina Bao, Wei Bar-ad, S. Baranov, Denis G.	HI02 BA04 SD04 PC20, QJ07 RB13 DD03 RB07 BD01 HD05	89 41 330 130, 222 256 58 255 43 86
Badini Confalonieri, Giovanni A. Bae, Dong Sik Bae, Joohyung Bae, Jungchan Bae, Kyoung-hoon Bae, Moon Seob Bae, Seok Bae, Taejin	GG04 PH21 QK12 PM17 RO02 RP23 RH08 JI02	80 150 226 170 301 307 276 105	Bansil, A. Bansil, Arun Bansil, B. Bao, Guizhi Bao, Jinke Bao, Nina Bao, Wei Bar-ad, S. Baranov, Denis G. Baranov, Nikolai V.	HI02 BA04 SD04 PC20, QJ07 RB13 DD03 RB07 BD01	89 41 330 130, 222 256 58 255 43
Badini Confalonieri, Giovanni A. Bae, Dong Sik Bae, Joohyung Bae, Jungchan Bae, Kyoung-hoon Bae, Kyoung-hoon Bae, Moon Seob Bae, Seok Bae, Taejin Bae, Yu Jeong	GG04 PH21 QK12 PM17 RO02 RP23 RH08 JI02 CB01, PJ17, QP18	80 150 226 170 301 307 276 105 49, 157, 248	Bansil, A. Bansil, Arun Bansil, B. Bao, Guizhi Bao, Jinke Bao, Nina Bao, Wei Bar-ad, S. Baranov, Denis G.	HI02 BA04 SD04 PC20, QJ07 RB13 DD03 RB07 BD01 HD05	89 41 330 130, 222 256 58 255 43 86
Badini Confalonieri, Giovanni A. Bae, Dong Sik Bae, Joohyung Bae, Jungchan Bae, Kyoung-hoon Bae, Kyoung-hoon Bae, Moon Seob Bae, Seok Bae, Taejin Bae, Yu Jeong Baek, In-bok	GG04 PH21 QK12 PM17 RO02 RP23 RH08 JI02 CB01, PJ17, QP18 RL21	80 150 226 170 301 307 276 105 49, 157, 248 293	Bansil, A. Bansil, Arun Bansil, B. Bao, Guizhi Bao, Jinke Bao, Nina Bao, Wei Bar-ad, S. Baranov, Denis G. Baranov, Nikolai V.	HI02 BA04 SD04 PC20, QJ07 RB13 DD03 RB07 BD01 HD05 QN10	89 41 330 130, 222 256 58 255 43 86 238
Badini Confalonieri, Giovanni A. Bae, Dong Sik Bae, Joohyung Bae, Jungchan Bae, Kyoung-hoon Bae, Kyoung-hoon Bae, Moon Seob Bae, Seok Bae, Taejin Bae, Yu Jeong Baek, In-bok Baek, Seung Ki	GG04 PH21 QK12 PM17 RO02 RP23 RH08 JI02 CB01, PJ17, QP18 RL21 R119, RI20	80 150 226 170 301 307 276 105 49, 157, 248 293 282	Bansil, A. Bansil, Arun Bansil, B. Bao, Guizhi Bao, Jinke Bao, Nina Bao, Wei Bar-ad, S. Baranov, Denis G. Baranov, Nikolai V. Baranov, Nikolay	HI02 BA04 SD04 PC20, QJ07 RB13 DD03 RB07 BD01 HD05 QN10 QG04	89 41 330 130, 222 256 58 255 43 86 238 208
Badini Confalonieri, Giovanni A. Bae, Dong Sik Bae, Joohyung Bae, Jungchan Bae, Kyoung-hoon Bae, Kyoung-hoon Bae, Moon Seob Bae, Seok Bae, Taejin Baek, Taejin Baek, In-bok Baek, Seung Ki Baek, Seung Ki Baek, Seung-heon Chris	GG04 PH21 QK12 PM17 RO02 RP23 RH08 JI02 CB01, PJ17, QP18 RL21 R119, RI20 RL30	80 150 226 170 301 307 276 105 49, 157, 248 293 282 294	Bansil, A. Bansil, Arun Bansil, B. Bao, Guizhi Bao, Jinke Bao, Nina Bao, Wei Bar-ad, S. Baranov, Denis G. Baranov, Nikolai V. Baranov, Nikolay Baranov, Nikolay V.	HI02 BA04 SD04 PC20, QJ07 RB13 DD03 RB07 BD01 HD05 QN10 QG04 QI08	89 41 330 130, 222 256 58 255 43 86 238 208 219
Badini Confalonieri, Giovanni A. Bae, Dong Sik Bae, Joohyung Bae, Jungchan Bae, Kyoung-hoon Bae, Kyoung-hoon Bae, Moon Seob Bae, Seok Bae, Taejin Bae, Taejin Baek, In-bok Baek, Seung Ki Baek, Seung Ki Baek, Seung min	GG04 PH21 QK12 PM17 RO02 RP23 RH08 JI02 CB01, PJ17, QP18 RL21 R119, R120 RL30 QP15, QP16	80 150 226 170 301 307 276 105 49, 157, 248 293 282 294 294 247, 248	Bansil, A. Bansil, Arun Bansil, B. Bao, Guizhi Bao, Jinke Bao, Nina Bao, Wei Bar-ad, S. Baranov, Denis G. Baranov, Nikolai V. Baranov, Nikolay V. Baranov, Nikolay V. Barbar, S. K.	HI02 BA04 SD04 PC20, QJ07 RB13 DD03 RB07 BD01 HD05 QN10 QG04 QI08 RA03	89 41 330 130, 222 256 58 255 43 86 238 208 219 249
Badini Confalonieri, Giovanni A. Bae, Dong Sik Bae, Joohyung Bae, Jungchan Bae, Kyoung-hoon Bae, Kyoung-hoon Bae, Moon Seob Bae, Seok Bae, Taejin Bae, Yu Jeong Baek, In-bok Baek, Seung Ki Baek, Seung Ki Baek, Seung min Baek, Younkyoung	GG04 PH21 QK12 PM17 RO02 RP23 RH08 JI02 CB01, PJ17, QP18 RL21 R119, R120 RL30 QP15, QP16 PM24	80 150 226 170 301 307 276 105 49, 157, 248 293 282 294 247, 248 171	Bansil, A. Bansil, Arun Bansil, B. Bao, Guizhi Bao, Jinke Bao, Nina Bao, Wei Bar-ad, S. Baranov, Denis G. Baranov, Nikolai V. Baranov, Nikolay V. Baranov, Nikolay V. Barbar, S. K. Barberis, Gaston Eduardo	HI02 BA04 SD04 PC20, QJ07 RB13 DD03 RB07 BD01 HD05 QN10 QG04 QI08 RA03 RB18	89 41 330 130, 222 256 58 255 43 86 238 208 219 249 257
Badini Confalonieri, Giovanni A. Bae, Dong Sik Bae, Joohyung Bae, Jungchan Bae, Kyoung-hoon Bae, Kyoung-hoon Bae, Kyoung-hoon Bae, Taejin Bae, Taejin Baek, In-bok Baek, Seung Ki Baek, Seung-heon Chris Baek, Seungmin Baek, Younkyoung Baek, Younkyoung Baek, Youn-kyoung	GG04 PH21 QK12 PM17 RO02 RP23 RH08 JI02 CB01, PJ17, QP18 RL21 R119, RI20 RL30 QP15, QP16 PM24 RN06, SL29	80 150 226 170 301 307 276 105 49, 157, 248 293 282 294 247, 248 171 297, 360	Bansil, A. Bansil, Arun Bansil, B. Bao, Guizhi Bao, Jinke Bao, Nina Bao, Wei Bar-ad, S. Baranov, Denis G. Baranov, Nikolai V. Baranov, Nikolay V. Baranov, Nikolay V. Barbar, S. K. Barberis, Gaston Eduardo Barbiellini, Bernardo	HI02 BA04 SD04 PC20, QJ07 RB13 DD03 RB07 BD01 HD05 QN10 QG04 QI08 RA03 RB18 BA04	89         41         330         130, 222         256         58         255         43         86         238         208         219         257         41
Badini Confalonieri, Giovanni A. Bae, Dong Sik Bae, Joohyung Bae, Jungchan Bae, Kyoung-hoon Bae, Kyoung-hoon Bae, Kyoung-hoon Bae, Seok Bae, Taejin Baek, Taejin Baek, In-bok Baek, Seung Ki Baek, Seung Ki Baek, Seung-heon Chris Baek, Seungmin Baek, Younkyoung Baek, Younkyoung Baek, Younkyoung Baek, Younkyoung Baegio-saitovitch, E.	GG04 PH21 QK12 PM17 RO02 RP23 RH08 JI02 CB01, PJ17, QP18 RL21 RI19, RI20 RL30 QP15, QP16 PM24 RN06, SL29 SM17	80 150 226 170 301 307 276 105 49, 157, 248 293 282 294 247, 248 171 297, 360 363	Bansil, A. Bansil, Arun Bansil, B. Bao, Guizhi Bao, Jinke Bao, Nina Bao, Wei Bar-ad, S. Baranov, Denis G. Baranov, Nikolai V. Baranov, Nikolay Baranov, Nikolay V. Barbar, S. K. Barberis, Gaston Eduardo Barbiellini, Bernardo Barco, Reginaldo	HI02 BA04 SD04 PC20, QJ07 RB13 DD03 RB07 BD01 HD05 QN10 QG04 QI08 RA03 RB18 BA04 PG03	89         41         330         130, 222         256         58         255         43         86         238         208         219         249         257         41         142
Badini Confalonieri, Giovanni A. Bae, Dong Sik Bae, Joohyung Bae, Jungchan Bae, Kyoung-hoon Bae, Kyoung-hoon Bae, Kyoung-hoon Bae, Seok Bae, Taejin Baek, Taejin Baek, In-bok Baek, Seung Ki Baek, Seung Ki Baek, Seung-heon Chris Baek, Seungmin Baek, Younkyoung Baek, Younkyoung Baek, Younkyoung Baek, Younkyoung Baegio-saitovitch, E.	GG04 PH21 QK12 PM17 RO02 RP23 RH08 JI02 CB01, PJ17, QP18 RL21 R19, RI20 RL30 QP15, QP16 PM24 RN06, SL29 SM17 PL09, QM14, QM15,	80 150 226 170 301 307 276 105 49, 157, 248 293 282 294 247, 248 171 297, 360 363 165, 233, 234, 255,	Bansil, A. Bansil, Arun Bansil, B. Bao, Guizhi Bao, Jinke Bao, Nina Bao, Wei Bar-ad, S. Baranov, Denis G. Baranov, Nikolai V. Baranov, Nikolai V. Baranov, Nikolay Baranov, Nikolay V. Barbar, S. K. Barberis, Gaston Eduardo Barbeilini, Bernardo Barco, Reginaldo Barcones, Beatriz	HI02 BA04 SD04 PC20, QJ07 RB13 DD03 RB07 BD01 HD05 QN10 QG04 QI08 RA03 RB18 BA04 PG03 JE01	89         41         330         130, 222         256         58         255         43         86         238         208         219         249         257         41         142         101
3adini Confalonieri, Giovanni A. 3ae, Dong Sik 3ae, Joohyung 3ae, Jungchan 3ae, Kyoung-hoon 3ae, Kyoung-hoon 3ae, Kyoung-hoon 3ae, Seok 3ae, Taejin 3ae, Taejin 3ae, Yu Jeong 3aek, Taejin 3aek, Seung Ki 3aek, Seung Ki 3aek, Seung-heon Chris 3aek, Seungmin 3aek, Youn-kyoung 3aek, Youn-kyoung 3aggio-saitovitch, E. 3aggio-saitovitch, Elisa 3agheri, Mahdiyar	GG04 PH21 QK12 PM17 RO02 RP23 RH08 JI02 CB01, PJ17, QP18 RL21 RI19, RI20 RL30 QP15, QP16 PM24 RN06, SL29 SM17 PL09, QM14, QM15, RB10, RB11, SE22 QA25, SA22	80 150 226 170 301 307 276 105 49, 157, 248 293 282 294 247, 248 171 297, 360 363 165, 233, 234, 255, 336 186, 319	Bansil, A. Bansil, Arun Bansil, B. Bao, Guizhi Bao, Jinke Bao, Nina Bao, Wei Bar-ad, S. Baranov, Denis G. Baranov, Nikolai V. Baranov, Nikolay V. Barbar, S. K. Barberis, Gaston Eduardo Barbiellini, Bernardo Barco, Reginaldo Barcones, Beatriz Barik, Sujit Kumar Barisic, Neven	HI02 BA04 SD04 PC20, QJ07 RB13 DD03 RB07 BD01 HD05 QN10 QG04 QI08 RA03 RB18 BA04 PG03 JE01 PI01, RA08 CA03	89         41         330         130, 222         256         58         255         43         86         238         208         219         249         257         41         101         150, 250         49
Badini Confalonieri, Giovanni A. Bae, Dong Sik Bae, Joohyung Bae, Jungchan Bae, Kyoung-hoon Bae, Kyoung-hoon Bae, Kyoung-hoon Bae, Moon Seob Bae, Seok Bae, Taejin Baek, Taejin Baek, Taejin Baek, Young Baek, Seung Ki Baek, Seung Ki Baek, Seung Ki Baek, Seung Ki Baek, Seung Ki Baek, Youn-kyoung Baek, Youn-kyoung Baggio-saitovitch, E. Baggio-saitovitch, Elisa Bagheri, Mahdiyar Bahmanrokh, Ghazaleh	GG04 PH21 QK12 PM17 RO02 RP23 RH08 JI02 CB01, PJ17, QP18 RL21 R19, RI20 RL30 QP15, QP16 PM24 RN06, SL29 SM17 PL09, QM14, QM15, RB10, RB11, SE22 QA25, SA22 RQ22	80 150 226 170 301 307 276 105 49, 157, 248 293 282 294 247, 248 171 297, 360 363 165, 233, 234, 255, 336 186, 319 310	Bansil, A. Bansil, Arun Bansil, B. Bao, Guizhi Bao, Jinke Bao, Nina Bao, Wei Bar-ad, S. Baranov, Denis G. Baranov, Nikolai V. Baranov, Nikolay V. Barbar, S. K. Barberis, Gaston Eduardo Barbiellini, Bernardo Barbiellini, Bernardo Barcones, Beatriz Barik, Sujit Kumar Barisic, Neven Barker, J	HI02 BA04 SD04 PC20, QJ07 RB13 DD03 RB07 BD01 HD05 QN10 QG04 QI08 RA03 RB18 BA04 PG03 JE01 PI01, RA08 CA03 FC01	89         41         330         130, 222         256         58         255         43         86         238         208         219         249         257         41         142         101         150, 250         49         73
Badini Confalonieri, Giovanni A. Bae, Dong Sik Bae, Joohyung Bae, Jungchan Bae, Kyoung-hoon Bae, Kyoung-hoon Bae, Kyoung-hoon Bae, Seok Bae, Taejin Baek, Taejin Baek, Taejin Baek, Yu Jeong Baek, In-bok Baek, Seung Ki Baek, Seung Ki Baek, Seung Ki Baek, Seung Ki Baek, Seung Ki Baek, Seung Ki Baek, Youn-kyoung Baek, Youn-kyoung Baggio-saitovitch, E. Baggio-saitovitch, Elisa Bagheri, Mahdiyar Bahmanrokh, Ghazaleh Bahramy, M.s.	GG04 PH21 QK12 PM17 RO02 RP23 RH08 JI02 CB01, PJ17, QP18 RL21 R119, R120 RL30 QP15, QP16 PM24 RN06, SL29 SM17 PL09, QM14, QM15, RB10, RB11, SE22 QA25, SA22 RQ22 II02	80 150 226 170 301 307 276 105 49, 157, 248 293 282 294 247, 248 171 297, 360 363 165, 233, 234, 255, 336 186, 319 310 97	Bansil, A. Bansil, Arun Bansil, B. Bao, Guizhi Bao, Jinke Bao, Nina Bao, Wei Bar-ad, S. Baranov, Denis G. Baranov, Nikolai V. Baranov, Nikolai V. Baranov, Nikolay V. Barbar, S. K. Barberis, Gaston Eduardo Barbiellini, Bernardo Barbiellini, Bernardo Barco, Reginaldo Barcones, Beatriz Barik, Sujit Kumar Barisic, Neven Barker, J Barker, Je	HI02 BA04 SD04 PC20, QJ07 RB13 DD03 RB07 BD01 HD05 QN10 QG04 QI08 RA03 RB18 BA04 PG03 JE01 PI01, RA08 CA03 FC01 RK17	89         41         330         130, 222         256         58         255         43         86         238         208         219         249         257         41         142         101         150, 250         49         73         288
Badini Confalonieri, Giovanni A. Bae, Dong Sik Bae, Joohyung Bae, Jungchan Bae, Kyoung-hoon Bae, Kyoung-hoon Bae, Kyoung-hoon Bae, Seok Bae, Seok Bae, Taejin Baek, Seung Baek, Juleong Baek, Juleong Baek, Seung Ki Baek, Seung Ki Baek, Seung Ki Baek, Seung Ki Baek, Seung Ki Baek, Seung Ki Baek, Youn-kyoung Baek, Youn-kyoung Baggio-saitovitch, E. Baggio-saitovitch, Elisa Bagheri, Mahdiyar Bahmanrokh, Ghazaleh Bahramy, M.s. Bahramy, Mohammad Saeed	GG04 PH21 QK12 PM17 RO02 RP23 RH08 JI02 CB01, PJ17, QP18 RL21 R119, RI20 RL30 QP15, QP16 PM24 RN06, SL29 SM17 PL09, QM14, QM15, RB10, RB11, SE22 QA25, SA22 RQ22 II02 AE05	80 150 226 170 301 307 276 105 49, 157, 248 293 282 294 247, 248 171 297, 360 363 165, 233, 234, 255, 336 186, 319 310 97 36	Bansil, A. Bansil, Arun Bansil, B. Bao, Guizhi Bao, Jinke Bao, Nina Bao, Wei Bar-ad, S. Baranov, Denis G. Baranov, Nikolai V. Baranov, Nikolay V. Barbar, S. K. Barberis, Gaston Eduardo Barbiellini, Bernardo Barbiellini, Bernardo Barco, Reginaldo Barcones, Beatriz Barik, Sujit Kumar Barisic, Neven Barker, J Barker, Joe Barman, Anjan	HI02 BA04 SD04 PC20, QJ07 RB13 DD03 RB07 BD01 HD05 QN10 QG04 QI08 RA03 RB18 BA04 PG03 JE01 PI01, RA08 CA03 FC01 RK17 HD04, RK11, SH12	89         41         330         130, 222         256         58         255         43         86         238         208         219         257         41         142         101         150, 250         49         73         288         85, 287, 345
Badini Confalonieri, Giovanni A. Bae, Dong Sik Bae, Joohyung Bae, Jungchan Bae, Kyoung-hoon Bae, Kyoung-hoon Bae, Kyoung-hoon Bae, Seok Bae, Seok Bae, Taejin Baek, Seung Ki Baek, Youn-kyoung Baek, Youn-kyoung Baggio-saitovitch, E. Baggio-saitovitch, Elisa Bagheri, Mahdiyar Bahmanrokh, Ghazaleh Bahramy, M.s. Bahramy, Mohammad Saeed Bai, X. J.	GG04 PH21 QK12 PM17 RO02 RP23 RH08 JI02 CB01, PJ17, QP18 RL21 R19, RI20 RL30 QP15, QP16 PM24 RN06, SL29 SM17 PL09, QM14, QM15, RB10, RB11, SE22 QA25, SA22 RQ22 II02 AE05 PL20	80 150 226 170 301 307 276 105 49, 157, 248 293 282 294 247, 248 171 297, 360 363 165, 233, 234, 255, 336 186, 319 310 97 36 167	Bansil, A. Bansil, Arun Bansil, B. Bao, Guizhi Bao, Jinke Bao, Nina Bao, Wei Bar-ad, S. Baranov, Denis G. Baranov, Nikolai V. Baranov, Nikolay V. Barbar, S. K. Barberis, Gaston Eduardo Barbiellini, Bernardo Barberis, Gaston Eduardo Barto, Reginaldo Barcones, Beatriz Barik, Sujit Kumar Barisic, Neven Barker, J Barker, Joe Barman, Anjan Barman, Saswati	HI02 BA04 SD04 PC20, QJ07 RB13 DD03 RB07 BD01 HD05 QN10 QG04 QI08 RA03 RB18 BA04 PG03 JE01 PI01, RA08 CA03 FC01 RK17 HD04, RK11, SH12 HD04, RK11, SH12	89         41         330         130, 222         256         58         255         43         86         238         208         219         249         257         41         142         101         150, 250         49         73         288         85, 287, 345         85, 287, 345
Badini Confalonieri, Giovanni A. Bae, Dong Sik Bae, Joohyung Bae, Jungchan Bae, Kyoung-hoon Bae, Kyoung-hoon Bae, Kyoung-hoon Bae, Seok Bae, Seok Bae, Taejin Baek, Seung Baek, Juleong Baek, Juleong Baek, Seung Ki Baek, Seung Ki Baek, Seung Ki Baek, Seung Ki Baek, Seung Ki Baek, Seung Ki Baek, Youn-kyoung Baek, Youn-kyoung Baggio-saitovitch, E. Baggio-saitovitch, Elisa Bagheri, Mahdiyar Bahmanrokh, Ghazaleh Bahramy, M.s. Bahramy, Mohammad Saeed	GG04 PH21 QK12 PM17 RO02 RP23 RH08 JI02 CB01, PJ17, QP18 RL21 R119, RI20 RL30 QP15, QP16 PM24 RN06, SL29 SM17 PL09, QM14, QM15, RB10, RB11, SE22 QA25, SA22 RQ22 II02 AE05	80 150 226 170 301 307 276 105 49, 157, 248 293 282 294 247, 248 171 297, 360 363 165, 233, 234, 255, 336 186, 319 310 97 36	Bansil, A. Bansil, Arun Bansil, B. Bao, Guizhi Bao, Jinke Bao, Nina Bao, Wei Bar-ad, S. Baranov, Denis G. Baranov, Nikolai V. Baranov, Nikolay V. Barbar, S. K. Barberis, Gaston Eduardo Barbiellini, Bernardo Barbiellini, Bernardo Barco, Reginaldo Barcones, Beatriz Barik, Sujit Kumar Barisic, Neven Barker, J Barker, Joe Barman, Anjan	HI02 BA04 SD04 PC20, QJ07 RB13 DD03 RB07 BD01 HD05 QN10 QG04 QI08 RA03 RB18 BA04 PG03 JE01 PI01, RA08 CA03 FC01 RK17 HD04, RK11, SH12	89         41         330         130, 222         256         58         255         43         86         238         208         219         257         41         142         101         150, 250         49         73         288         85, 287, 345

Barnes, Crispin	BH01	46
Barnes, Crispin H.w.	PN23	175
Baron-gonzalez, Aura Janeth Barraclough, J. M.	EC04 GC01	66 77
Barraclough, Jack M	SE12	335
Barrera, Gabriele	AG02	37
Barshilia, Harish C	PK16	163
Barski, Andre	Al04, DD01	39, 58
Barthelemy, Agnes	AE03, HE04, QA27	36, 86, 187
Barthelemy, M.	PO23	181
Barthelemy, Marie	FC02	73
Bartolome, Fernando	EF03, GG03, GG05, QI14	68, 80, 220
Bartolome, Juan	EF03, GG03, GG05	68, 80
Basak, Susmita	BA04	41
Bascones, Elena	QB18	190
Bascones, Elena .	QB19	190
Basso, Vittorio	HH01	88
Batail, Patrick	SB20	323
Bataille, Alexandre	DD05	59
Batista, Cristian D	JD01	100
Batko, Ivan	RE13	268
Batkova, Marianna	RE13	268
Batlle, X. Batlle, Xavier	BD01, ED03 BD04, RN14	43, 66 43, 299
Batlogg, B.	PC29, PI15	131, 153
Batlogg, Bertram	RB21, SC16	257, 328
Battiato, Marco	RK12	287
Bauer, Andreas	DF02, QC17, QH24	60, 195, 216
Bauer, E D	EA03	64
Bauer, E.	QB31, QD10	192, 198
Bauer, E. D.	FB01, QB12	73, 189
Bauer, E.d.	PC19	129
Bauer, Eric D	FB03	73
Bauer, Eric D.	GC03, PC03, RD09,	77, 127, 263, 328,
	SC16, SE02, SE08	333, 334
Bauer, Ernst	IA02, SC01	90, 325
Bauer, Thomas	SF05	337
Bauer, Uwe	AD02, AE01	35
Baum, Max	AA05	33 192
Baumbach, R. Baumberger, Felix	QB31 II02	97
Baumgarten, Martin	HC04	85
Bayle-guillemaud, Pascale	Al04, DD01	39, 58
Bazaliy, Ya. B.	SI18	348
Beach, Geoffrey	AE01	35
Beard, Jerome	RQ16	309
Beatrice, Cinzia	HH01	88
Beaurepaire, Eric	GI03	82
Beck, Adli	QH22	216
Becker, Petra	AA05, QA22, RI12	33, 186, 280
Behnia, Kamran	KJ04	114
Beigne, Cyrille	DD01, El04, GE01, QK15, QK16, RL24	58, 71, 78, 226, 293
Bekyarova, Elena	PN16	174
Belinsky, Moisey	RA25	253
Bellaiche, Laurent	QA27	187
Belo, Joao H	SO31	371
Ben Youssef, Jamal	SA24	320
Benatmane, N. Repock Pater	PK13	162
Bencok, Peter	PA15	118

## AUTHOR INDEX

Bending, Simon	Q024	244
Bendix, Jesper	BC02	42
Bera, A. K.	DH05	62
Beran, Premysl	BE05	44
Berardan, David	RB15	256
Berger, Andreas	SM04	361
Berger, Claire	PN16	174
Berger, H.	PI09	152
Berger, Helmuth	QA20	185
Bergman, Anders	SH11	345
Bergmann, Christoph	AB02	34
Bergqvist, Lars	SH11	345
Berkowski, Marek	QB30	192
Bernhard, Christian	HA02	83
Berntsen, Magnus	HA03	83
Bersweiler, Mathias	KF02	110
Bertacco, Riccardo	AI03, DE03	39, 59
Berthet, Patrick	GF05	79
Berthier, Claude	IC02	92
Bertram, Batlogg	RB19	257
Besnard, Celine	QC13	194
Bessais, Lotfi	PM06	168
Bessonov, Vladimir	SO17	369
Beyer, Markus	HA02	83
Bezmaternykh, Leonard N	CI04	55
Bezmaternykh, Leonard N.	PA02	116
Bhagyashree, Karki S	SL21	359
Bharathi, A.	QA14	184
Bhat, Shwetha G.	CG04	53
Bhat, Subray V	SL21	359
Bhatia, Charanjit S	SK05	351
Bhatia, Charanjit Singh	PJ08	156
Bhatia, Pushpinder Gupta	PN06	173
Bhatt, Pramod	PN17	174
Bhattacharya, Pallab	CG05	53
Bi, Ke	RA12	251
Bialek, Beata	SK19, SK27	353, 354
Bianchi, Andrea	JC01	100
Bianchi, Andrea D.	CD04, PD05	51, 132
Bianchi, M.	1102	97
Bibes, Manuel	AE03, HE04, QA27	36, 86, 187
Bickel, Jessica E.	IH03	96
Bigot, Jy.	PO23	181
Bigot, Jean-yves	FC02, RK05, RK13, RK15	73, 286, 288
Bihlmayer, Gustav	SD06	331
Billette, Julien	RQ16	309
Biner, D	RI05	279
Biondo, Armando	QM15	234
Biswas, Abhijit	PN28	176
Biswas, Deep Narayan Biswas	SC06	326
Biswas, S.	PL07	165
Bitla, Yugandhar	SL07	357
Bittar, Eduardo	SE22	336
Bjorck, Matts	EI03	71
Blackburn, S.	PD05	132
Blaha, Peter	RB21	257
Blanco, J. A.	GH05	81
Blanco, Jesus A.	JJ04	106
Blanco, Juan Maria	AJ02	40
Blanco-canosa, Santiago	BA03	41

Author Index

Blank, Dave	DH02, SA10	62, 317	Bourdarot, Frederic	JC02	100	Bunce, Christopher	RJ11	284
Blasco, Javier	EC04	66	Bourgeois, Julie	RA11	251	Bunkov, Yury	IC04	92
Blavette, Didier	BD02	43	Bourges, Philippe	DA01	56	Burdin, Sebastien	PD04	132
Blinova, Julia	PB22	123	Bouzehouane, Karim	AE03	36	Burkov, Alexander T	QI12, QI13	219
Bluegel, Stefan	SD06	331	Boyn, Soeren	HF01	87	Burzo, Emil	SK09	351
Blugel, Stefan	IH03	96	Bozzo, Bernat	SO18	369	Bussetti, Gianlorenzo	PL21	167
Blum, Christian	QH24	216	Braden, Markus	AA05	33	Bustamante, Rodney	KB06	108
Blundell, Stephen. J.	QG19	211	Braicovich, Lucio	BA03	41	Buzzi, Michele	RK17	288
Blythe, H. J.	RN18	299	Brambilla, Alberto	PL21	167	Bychkov, Igor V.	SO19	369
Boarino, Luca	AG02	37	Bramwell, Steven	CC01	50	Bychkov, Igor Valer'evich	QH17	215
Bobade, Santosh M	Q011	242	Bran, Christina	DG01	61	Byeon, Hyehyeon	PN08, PN09, QM04	173, 232
Bobak, Andrej	QE01, QE02	199	Bran, Cristina	DG03, EI03	61, 71	Byzov, Ilya	QL07	229
Bobrov, Aleksandr	SG17	342	Bran, Julien	DG02, PJ22	61, 158	byzov, nyu	QLUI	225
Bochenek, Lukasz	SB26	324	Brandenburg, Jens	QA11	184	С		
	PE19	139	Brandl, G.	PIO9	152	Cabra Danial C	QE26	203
Bochenek, Łukasz						Cabra, Daniel C.		
Bocher, Laura	AE03	36	Brando, Manuel	AB02, QH12, SF05	34, 214, 337	Cai, Jinzhu	PN05, QL04	172, 228
Bocklage, Lars	SH04	343	Braun, Anke	PE07	137	Cai, Yiwei	RE11	268
Bode, Matthias	AD03	35	Braun, H. B.	PJ26	159	Cai, Yong	QC09	193
Bodensiek, Oliver	PE03	136	Braun, Hans - Benjamin	GF02	79	Cai, Yongfu	RP17	306
Boehm, Helga M	RG15	274	Breard, Yohann	BB03, RH22	41, 278	Cai, Yongqing	QP06	246
Boehm, M.	QE19	202	Brede, Jens	GH02	81	Caizer, Costica	PO07	178
Boehnert, Tim	RP04	303	Brenac, Ariel	EI04	71	Caldas, Heron	KG04	112
Boeni, Peter	DG01	61	Brenig, Wolfram	EC05	66	Calderon Filho, Cesar Jose	RB18	257
Boeri, Lilia	SC07	326	Breunig, Oliver	RI12	280	Calderon, Maria Jose	QB18, QB19	190
Bof Bufon, Carlos Cesar	SJ07	350	Bridges, F.	CF01	52	Callaudin, Aurelie	KJ04	114
Bogach, A. V.	QC22, RH07	195, 275	Bridoux, G.	DI01	62	Calloni, Alberto	PL21	167
Bogach, Alexey	QC20, RE13	195, 268	Brinkman, Alexander	DH02, SA10	62, 317	Calvani, Paolo	SA08	317
Bogach, Alexey V.	QC19	195	Broholm, C.	CF01	52	Calvo, Irene	EF03	68
Bogdanov, Savva	QG02	208	Brookes, Nicholas	PI08	152	Camarero, Julio	SH03	343
Bohaty, Ladislav	AA05, RI12	33, 280	Brooks, James S.	SB05	321	Camilo, Nilmar Silva	SE15	335
Bok, Jin Mo	SB31	325	Brown, Stuart E.	GB03, SF01	76, 337	Campbell, Stewart	EF02	68
Boldt, Regine	PC27	131	Brucas, Rimantas	SL15	358	Campbell, Stewart J	QG26	212
Boldyrev, Kirill N.	PA20, RA24	119, 253	Bruck, E.	RQ12, SO08, SO16	309, 368, 369	Campbell, Stewart James	QG09	209
Bombardi, Alessandro	BF03	45	Bruck, Ekkes	PM10, RI18	169, 281	Campion, Richard	IE03	93
Bonanni, Alberta	BE03	44	Brueckl, Hubert	AH01	38	Campion, Richard P	SH09	344
Bonell, Frederic	PK20	163	Bruel, Andreas	SG19	342	Campo, Giulio	CB04, SI11	49, 347
Bonfim, Marlio	SH05	344	Bruhl, Andreas	RI13	281	Campo, Javier	GF04, KB06, PG22,	79, 108, 146, 24
Boni, Peter	DF02	60	Brune, Harald	BC02	42	cumpo, sanci	Q016	, , , , , , , , , , , , , , , , , , , ,
Bonilla, Marcela	EF03	68	Buchanan, Kristen	RJ20	285	Canfield, P. C.	GA01, IA05, RD02	75, 91, 262
Bonn, Doug	BA01	40	Buchelnikov, Vacily	PI20	154	Cantoni, Matteo	AI03	39
Booth, John Graham	QG01	208	Buchelnikov, Vacily	PM22, QH05, SO23	171, 213, 370			167
						Cao, C. D.	PL20	
Boothroyd, A. T. Boothroyd, Androw	PC22, PI03, QN01	130, 151, 236	Buchelnikov, Vasiliy D.	Q012, S019	242, 369	Cao, Guanghan	AB04, IJ04, RB13	34, 98, 256
Boothroyd, Andrew	AF01, DH03	36, 62	Buchelnikov, Vasily Dmitrievich		215	Cao, Helin	HI02	89
Boothroyd, Andrew T	HA04	83	Buchner, B.	QF18	206	Cao, R	SB19	323
Borchers, Julie	KD02	109	Buckow, Alexander	PI05	151	Cao, Rong	PC30	131
Borchmann, Jan	SL05	356	Buda-prejbeanu, Liliana	SI09	347	Cao, Shixun	SG01	339
Borges, Hortencio Alves	SE22	336	Bud'ko, S. L.	GA01, IA05, RD02	75, 91, 262	Cao, W.n.	PL02	164
Borges, R. P.	QM13	233	Budnick, Joseph I.	HA03	83	Capelli, Silvia	RE12	268
Boris, A. V.	GA03	75	Budrikis, Zoe	KE01	110	Capik, Marek	IB01	91
Borovsky, Michal	QE02	199	Buechner, B.	QE23, QF22, RB17	203, 207, 256	Caprile, Ambra	HH01	88
Borozdina, Yulia	HC04	85	Buechner, Bernd	QH24	216	Cardona, S.	BC01	42
Borrell, A	RQ10	308	Bufaical, Leandro F	SE15	335	Carlos, Luis D.	KB06	108
Borrmann, Horst	QD07	197	Buhot, Jonathan	PD17	134	Carlotti, Giovanni	GG02	80
Bortolotti, P.	ID02	93	Bui, Hong-van	SG10	341	Carman, Gregory P.	QA15	185
Bortolotti, Paolo	CE03, ID03, RJ07	51, 93, 283	Buitrago, Ivon R.	SI10	347	Carmelo, Jose	RF13	271
Bosak, Alexei	PA08	117	Bukowski, Z	QB06	188	Caron, Luana	PM10, RN07	169, 297
Boseggia, S.	QN01	236	Bukowski, Z.	PC01	126	Carpenter, Michael A	EF03	68
Bouchenoire, L.	QN01	236	Bukowski, Zbigniew	RB21	257	Carretero, Cecile	QA27	187
Bouillot, Pierre	CC03	50	Buldmann, Benjamin	RI12	280	Carta, Daniela	CB04	49
			· · ·					

## AUTHOR INDEX

Carva, Karel	RK12	287
Carvajal Quiroz, Eliel	SK15	352
Casa, Diego	DC03	58
Casanova, Felix	HE02	86
Casanove, Marie-jose	IH01	95
Casola, Francesco	IC02	92
Casoli, Francesca	AG02	37
Casper, Frederick	RC10	260
Castan, Celia	EF03	68
Castan-guerrero, Celia	GG03, GG05	80
Castano, F.	ED03	66
Castellano-hernandez, E.	ED03	66
Castrillon, Mariana	RN09	298
Caux, Jean-sebastien	IC05	92
Cava, R. J.	SD04	330
Cava, Robert J	QB13	189
Cavill, Stuart A	SH09	344
Cazayous, Maximilien	PD17, QA27	134, 187
Celegato, Federica	AG02, DG04	37, 61
Cermak, Petr	SB08	321
Cezar, Julio	PI08	152
Cezar, Julio C.	SH03, SH05	343, 344
Cha, Hyun Rok	QO20, SG21	244, 343
Cha, Yun-yeop	PO06, PO10	178
Chacon Hernandez, U. D.	SM17	363
Chacon, A.	PI09	152
Chadov, Stanislav	CH04, II03, RC10	54, 97, 260
Chadova, Kristina	1103	97
Chae, Keun Hwa	PG27, QO27	146, 245
Chae, Woong-chan	SN04	364
Chai, Yi Sheng	PA09	117
Chai, Yisheng	JB02	99
Chakraborty, Monodeep	RG02	272
Chaloupka, Jiri Chan. Ho Bun	BA03 BC05	41 42
Chan, Mun K.	CA03	42
Chan, Wen-yuan	PK02	160
Chandarak, Sujittra	SL10	357
Chandra, Shekhar	SD15	332
Chandra, Sumal	PB24, QG03	124, 208
Chandrasekaran, Gopalakrishnan	PO21, QA25, QN26	180, 186, 240
Chang, Bin	НЈО5	90
Chang, Cheng-hsun	PK02	160
Chang, Chia-chi	SH08	344
Chang, Ching Hao	KI04	113
Chang, Ching-ray	QK22, RK14, SK30	227, 288, 355
Chang, H.w.	SN06	365
Chang, Hun	CI03, PA01, PA03	55, 116
Chang, J.	PC01, PI15	126, 153
Chang, Jae Kyung	QH29, SC08	217, 326
Chang, Johan	HA03	83
Chang, Joonyeon	KC04, PJ30, PL16, QK10, QK11, QK12, QK13, RL28, RL31, SL26	109, 159, 166, 225, 226, 294, 295, 360
Chang, Jui-hang	RK14	288
Chang, Junghwan	RR16, RR17, SN05	313, 365
Chang, Liang-juan	RL18	292
Chang, P. C.	PE02	136
Chang, P. H.	PC13	128
Chang, Sung	QA20	185

Chang, Tr	SD04	330	Cheng, Chih-wei	PK11, PK15, PK19	162, 163
Chang, W.c.	SN06	365	Cheng, Hong-wen	QM09	233
Chanthbouala, Andre	IE04	94	Cheng, Nian-jia	SL18	359
Chantrell, Roy	RK17	288	Cheng, Po-wen	BJ01	47
Chantrell, Roy W	FC01	73	Cheng, Tsung-i	PK19	163
Chantrell, Roy W.	QK09, SJ04	225, 349	Cheon, Jinwoo	CB01, PO15, PO16	49, 179
Chantrell, Rw	FC01	73	Cheong, S. W.	AA04	33
Chao, C. H.	PK01	160	Cheong, S.w.	PH10	148
Chao, Wu	Q015	243	Cheong, Sw.	QA04	183
Chapon, Laurent C	BF03, CI02	45, 55	Cheong, Sang-wook	PP06, QF20	28, 207
Chapon, Laurent C.	GF03	79	Cheong, S-w.	RA22	253
Charnukha, Aliaksei	GA03	75	Cheptiakov, Denis	PI21	154
Chatterjee, S	SO14	369	Chern, G.	PK11, PK15, PK19	162, 163
Chattopadhyay, S	QL14	230	Chern, Gung	QM11, QM25	233, 235
Chaturvedi, Anurag	DJ01	63	Chernets, Ivan A.	Q012	242
Chaudhary, S.	PG27	146	Chernobrovkin, Alexey	ED06, KI05	67, 113
Che, Sun Chang	QG18	211	Chernyshov, Dmitry	PA14	118
Chebotkevich, Ludmila	AG05, QM05	38, 232	Chi, Sonqxue	SC23	329
Cheikhrouhou, Abdelwaheb	DJ05	64	Chiba, D.	BE02, RL07, SH07	44, 291, 344
Cheikhrouhou-koubaa, Wissen		64	Chiba, Daichi	RL10	291
Chen, Bin	PC02, PC09	127, 128	Chicinas, Ionel	EJ03	72
Chen, C T.	QB28	191	Chico Carpio, Jonathan Philippe		345
Chen, Chien-miao	SB19	323	Chien, Chia-ling	FA04	72
Chen, Chih_yung	РК09	161	Chikako, Moriyoshi	QD14	198
Chen, Chih-wen	QN08	237	Chikoidze, Ekaterina	SM07	362
Chen, Chih-yung	QM01	231	Chin, Yi-ying	PIO8	152
Chen, G F	QB11	188	Chioran, Viorica	RR02, RR03	311
Chen, G.	SM05	361	Chiu, Yu-che	RK07	287
Chen, Genfu	RB11	255	Chizhik, Alexander	CJ04	56
Chen, H.m.	PK11, PK15	162	Chizhik, Alexandr	AJ02	40
Chen, Hao-hsuan	RK14	288	Cho, B. K.	PK04, QI15, QK03, SB10,	161, 220, 224, 321,
Chen, Hsin-hsien	QM27	236		SB30	325
		320	Cho. Bo Ram	SB22	323
Chen, I. A.	SB04	320 189	Cho, Bo Ram	SB22 PN21_OL18	323 175 231
Chen, I. A. Chen, J.	SB04 QB12	189	Cho, Chae Ryong	PN21, QL18	175, 231
Chen, I. A. Chen, J. Chen, J. M.	SB04 QB12 SM11	189 362	Cho, Chae Ryong Cho, Daeheum	PN21, QL18 QG17	175, 231 211
Chen, I. A. Chen, J. Chen, J. M. Chen, Jian	SB04 QB12 SM11 SB29	189 362 325	Cho, Chae Ryong Cho, Daeheum Cho, Deo-kyong	PN21, QL18 QG17 PG24	175, 231 211 146
Chen, I. A. Chen, J. Chen, J. M. Chen, Jian Chen, Jin-ming	SB04 QB12 SM11 SB29 QB21	189 362 325 190	Cho, Chae Ryong Cho, Daeheum Cho, Deo-kyong Cho, Hyun Jin	PN21, QL18 QG17 PG24 PG25	175, 231 211 146 146
Chen, I. A. Chen, J. Chen, J. M. Chen, Jian Chen, Jin-ming Chen, K.I.	SB04 QB12 SM11 SB29 QB21 QM27	189 362 325 190 236	Cho, Chae Ryong Cho, Daeheum Cho, Deo-kyong Cho, Hyun Jin Cho, Hyunduck	PN21, QL18 QG17 PG24 PG25 PJ17	175, 231 211 146 146 157
Chen, I. A. Chen, J. Chen, J. M. Chen, Jian Chen, Jin-ming Chen, K.I. Chen, Kuo-chin	SB04 QB12 SM11 SB29 QB21 QM27 QK22, SK30	189 362 325 190 236 227, 355	Cho, Chae Ryong Cho, Daeheum Cho, Deo-kyong Cho, Hyun Jin Cho, Hyunduck Cho, Hyun-sung	PN21, QL18 QG17 PG24 PG25 PJ17 RR15	175, 231 211 146 146 157 313
Chen, I. A. Chen, J. Chen, J. M. Chen, Jian Chen, Jin-ming Chen, K.I. Chen, Kuo-chin Chen, L. M.	SB04 QB12 SM11 SB29 QB21 QM27 QK22, SK30 IC03	189 362 325 190 236 227, 355 92	Cho, Chae Ryong Cho, Daeheum Cho, Deo-kyong Cho, Hyun Jin Cho, Hyunduck Cho, Hyun-sung Cho, Jun-hyung	PN21, QL18 QG17 PG24 PG25 PJ17 RR15 SD06	175, 231 211 146 146 157 313 331
Chen, I. A. Chen, J. Chen, J. M. Chen, Jian Chen, Jin-ming Chen, K.I. Chen, Kuo-chin Chen, L. M. Chen, Lin	SB04 QB12 SM11 SB29 QB21 QM27 QK22, SK30 IC03 QK23	189 362 325 190 236 227, 355 92 227	Cho, Chae Ryong Cho, Daeheum Cho, Deo-kyong Cho, Hyun Jin Cho, Hyunduck Cho, Hyun-sung Cho, Jun-hyung Cho, K.	PN21, QL18 QG17 PG24 PG25 PJ17 RR15 SD06 QC23	175, 231 211 146 146 157 313 331 196
Chen, I. A. Chen, J. Chen, J. M. Chen, Jian Chen, Jin-ming Chen, K.I. Chen, Kuo-chin Chen, L. M. Chen, Lin Chen, M.j.	SB04 QB12 SM11 SB29 QB21 QM27 QK22, SK30 IC03 QK23 QM27	189 362 325 190 236 227, 355 92 227 236	Cho, Chae Ryong Cho, Daeheum Cho, Deo-kyong Cho, Hyun Jin Cho, Hyunduck Cho, Hyun-sung Cho, Jun-hyung Cho, K. Cho, K. H.	PN21, QL18 QG17 PG24 PG25 PJ17 RR15 SD06 QC23 SB28	175, 231 211 146 146 157 313 331 196 324
Chen, I. A. Chen, J. Chen, J. M. Chen, Jian Chen, Jin-ming Chen, K.I. Chen, Kuo-chin Chen, L. M. Chen, Lin Chen, Lin Chen, P.	SB04 QB12 SM11 SB29 QB21 QM27 QK22, SK30 IC03 QK23 QM27 AA04	189 362 325 190 236 227, 355 92 227 236 33	Cho, Chae Ryong Cho, Daeheum Cho, Deo-kyong Cho, Hyun Jin Cho, Hyunduck Cho, Hyun-sung Cho, Jun-hyung Cho, K. Cho, K. H. Cho, K. Wang Lae	PN21, QL18 QG17 PG24 PG25 PJ17 RR15 SD06 QC23 SB28 P113	175, 231 211 146 146 157 313 331 196 324 152
Chen, I. A. Chen, J. Chen, J. M. Chen, Jian Chen, Jin-ming Chen, K.I. Chen, Kuo-chin Chen, L. M. Chen, Lin Chen, Lin Chen, P. Chen, Qianping	SB04 QB12 SM11 SB29 QB21 QM27 QK22, SK30 IC03 QK23 QM27 AA04 SA19, SA20	189 362 325 190 236 227, 355 92 227 236 33 319	Cho, Chae Ryong Cho, Daeheum Cho, Deo-kyong Cho, Hyun Jin Cho, Hyunduck Cho, Hyun-sung Cho, Jun-hyung Cho, K. Cho, K. H. Cho, K. H. Cho, Kwang Lae Cho, Lee-hyun	PN21, QL18 QG17 PG24 PG25 PJ17 RR15 SD06 QC23 SB28 PI13 SK27	175, 231 211 146 146 157 313 331 196 324 152 354
Chen, I. A. Chen, J. Chen, J. M. Chen, Jian Chen, Jin-ming Chen, K.I. Chen, Kuo-chin Chen, L. M. Chen, Lin Chen, M.j. Chen, P. Chen, Qianping Chen, S. A.	SB04 QB12 SM11 SB29 QB21 QM27 QK22, SK30 IC03 QK23 QM27 AA04 SA19, SA20 QB21, SM11	189 362 325 190 236 227, 355 92 227 236 33 319 190, 362	Cho, Chae Ryong Cho, Daeheum Cho, Deo-kyong Cho, Hyun Jin Cho, Hyunduck Cho, Hyun-sung Cho, Jun-hyung Cho, K. Cho, K. H. Cho, K. Wang Lae Cho, Lee-hyun Cho, Myeong Woo	PN21, QL18 QG17 PG24 PG25 PJ17 RR15 SD06 QC23 SB28 Pl13 SK27 QN09	175, 231 211 146 146 157 313 331 196 324 152 354 237
Chen, I. A. Chen, J. Chen, J. M. Chen, Jian Chen, Jin-ming Chen, K.I. Chen, Kuo-chin Chen, Kuo-chin Chen, L. M. Chen, Lin Chen, M.j. Chen, P. Chen, Qianping Chen, S. A. Chen, S. C.	SB04 QB12 SM11 SB29 QB21 QM27 QK22, SK30 IC03 QK23 QM27 AA04 SA19, SA20 QB21, SM11 PH01	189 362 325 190 236 227, 355 92 227 236 33 319 190, 362 147	Cho, Chae Ryong Cho, Daeheum Cho, Deo-kyong Cho, Hyun Jin Cho, Hyunduck Cho, Hyun-sung Cho, Jun-hyung Cho, K. Cho, K. H. Cho, Kwang Lae Cho, Lee-hyun Cho, Myeong Woo Cho, Sang-uk	PN21, QL18 QG17 PG24 PG25 PJ17 RR15 SD06 QC23 SB28 PI13 SK27 QN09 PH22	175, 231 211 146 146 157 313 331 196 324 152 354 237 150
Chen, I. A. Chen, J. Chen, J. M. Chen, Jian Chen, Jin-ming Chen, K.I. Chen, Kuo-chin Chen, Kuo-chin Chen, L. M. Chen, Lin Chen, M.J. Chen, P. Chen, Qianping Chen, S. A. Chen, S. C. Chen, Shih-jia	SB04 QB12 SM11 SB29 QB21 QM27 QK22, SK30 IC03 QK23 QM27 AA04 SA19, SA20 QB21, SM11 PH01 QM09	189 362 325 190 236 227, 355 92 227 236 33 319 190, 362 147 233	Cho, Chae Ryong Cho, Daeheum Cho, Deo-kyong Cho, Hyun Jin Cho, Hyunduck Cho, Hyun-sung Cho, Jun-hyung Cho, K. Cho, K. H. Cho, K. Wang Lae Cho, Lee-hyun Cho, Myeong Woo	PN21, QL18 QG17 PG24 PG25 PJ17 RR15 SD06 QC23 SB28 PI13 SK27 QN09 PH22 EH05, PN29, PN30,	175, 231 211 146 146 157 313 331 196 324 152 354 237
Chen, I. A. Chen, J. Chen, J. M. Chen, Jian Chen, Jin-ming Chen, K.I. Chen, Kuo-chin Chen, L. M. Chen, Lin Chen, Lin Chen, M.J. Chen, P. Chen, Qianping Chen, S. A. Chen, S. C. Chen, Shih-jia Chen, Sui-pin	SB04 QB12 SM11 SB29 QB21 QM27 QK22, SK30 IC03 QK23 QM27 AA04 SA19, SA20 QB21, SM11 PH01 QM09 RL29	189 362 325 190 236 227, 355 92 227 236 33 319 190, 362 147 233 294	Cho, Chae Ryong Cho, Daeheum Cho, Deo-kyong Cho, Hyun Jin Cho, Hyunduck Cho, Hyun-sung Cho, Jun-hyung Cho, K. Cho, K. H. Cho, Kwang Lae Cho, Lee-hyun Cho, Myeong Woo Cho, Sang-uk Cho, Sunglae	PN21, QL18 QG17 PG24 PG25 PJ17 RR15 SD06 QC23 SB28 PI13 SK27 QN09 PH22 EH05, PN29, PN30, QO22, QO23	175, 231 211 146 146 157 313 331 196 324 152 354 237 150 70, 176, 177, 244
Chen, I. A. Chen, J. Chen, J. M. Chen, Jian Chen, Jin-ming Chen, K.I. Chen, Kuo-chin Chen, L. M. Chen, Lin Chen, Lin Chen, M.J. Chen, P. Chen, Qianping Chen, S. A. Chen, S. C. Chen, Shih-jia Chen, Sui-pin Chen, Szu-wei	SB04 QB12 SM11 SB29 QB21 QM27 QK22, SK30 IC03 QK23 QM27 AA04 SA19, SA20 QB21, SM11 PH01 QM09 RL29 AD01	189 362 325 190 236 227, 355 92 227 236 33 319 190, 362 147 233 294 35	Cho, Chae Ryong Cho, Daeheum Cho, Deo-kyong Cho, Hyun Jin Cho, Hyunduck Cho, Hyun-sung Cho, Jun-hyung Cho, K. Cho, K. H. Cho, Kwang Lae Cho, Lee-hyun Cho, Myeong Woo Cho, Sang-uk Cho, Sunglae Cho, Suyeon	PN21, QL18 QG17 PG24 PG25 PJ17 RR15 SD06 QC23 SB28 P113 SK27 QN09 PH22 EH05, PN29, PN30, QO22, QO23 GF05	175, 231 211 146 146 157 313 331 196 324 152 354 237 150 70, 176, 177, 244
Chen, I. A. Chen, J. Chen, J. M. Chen, Jian Chen, Jin-ming Chen, K.I. Chen, Kuo-chin Chen, L. M. Chen, Lin Chen, Lin Chen, M.J. Chen, M.J. Chen, P. Chen, Qianping Chen, S. A. Chen, S. C. Chen, Shih-jia Chen, Sui-pin Chen, Szu-wei Chen, Wei-hsiang	SB04 QB12 SM11 SB29 QB21 QM27 QK22, SK30 IC03 QK23 QM27 AA04 SA19, SA20 QB21, SM11 PH01 QM09 RL29 AD01 IH06, PL03, PL13	189 362 325 190 236 227, 355 92 227 236 33 319 190, 362 147 233 294 35 96, 164, 166	Cho, Chae Ryong Cho, Daeheum Cho, Deo-kyong Cho, Hyun Jin Cho, Hyunduck Cho, Hyun-sung Cho, Jun-hyung Cho, K. Cho, K. H. Cho, Kwang Lae Cho, Lee-hyun Cho, Sang-uk Cho, Sang-uk Cho, Suyeon Cho, Suyeon Cho, Yong Chan	PN21, QL18 QG17 PG24 PG25 PJ17 RR15 SD06 QC23 SB28 P113 SK27 QN09 PH22 EH05, PN29, PN30, QO22, QO23 GF05 PN20, PN21	175, 231 211 146 157 313 331 196 324 152 354 237 150 70, 176, 177, 244 79 175
Chen, I. A. Chen, J. Chen, J. M. Chen, Jian Chen, Jin-ming Chen, K.I. Chen, Kuo-chin Chen, L. M. Chen, Lin Chen, Lin Chen, M.J. Chen, P. Chen, Qianping Chen, S. A. Chen, S. C. Chen, Shih-jia Chen, Sui-pin Chen, Szu-wei Chen, Wei-hsiang Chen, Wen-chen	SB04 QB12 SM11 SB29 QB21 QM27 QK22, SK30 IC03 QK23 QM27 AA04 SA19, SA20 QB21, SM11 PH01 QM09 RL29 AD01 IH06, PL03, PL13 PL18	189 362 325 190 236 227, 355 92 227 236 33 319 190, 362 147 233 294 35 96, 164, 166 166	Cho, Chae Ryong Cho, Daeheum Cho, Deo-kyong Cho, Hyun Jin Cho, Hyunduck Cho, Hyun-sung Cho, Jun-hyung Cho, K. Cho, K. H. Cho, Kwang Lae Cho, Lee-hyun Cho, Myeong Woo Cho, Sang-uk Cho, Sunglae Cho, Suyeon Cho, Suyeon Cho, Yong Chan Cho, Yong Chan	PN21, QL18 QG17 PG24 PG25 PJ17 RR15 SD06 QC23 SB28 P113 SK27 QN09 PH22 EH05, PN29, PN30, QO22, QO23 GF05 PN20, PN21 PH22, QL18	175, 231 211 146 157 313 331 196 324 152 354 237 150 70, 176, 177, 244 79 175 150, 231
Chen, I. A. Chen, J. Chen, J. M. Chen, Jian Chen, Jin-ming Chen, K.I. Chen, Kuo-chin Chen, Kuo-chin Chen, L. M. Chen, L. M. Chen, Lin Chen, D. Chen, M.j. Chen, P. Chen, Qianping Chen, S. A. Chen, S. C. Chen, Shih-jia Chen, Sui-pin Chen, Sui-pin Chen, Sui-wei Chen, Wei-hsiang Chen, Wen-chen Chen, Xiang-bai	S804 QB12 SM11 SB29 QB21 QM27 QK22, SK30 IC03 QK23 QM27 AA04 SA19, SA20 QB21, SM11 PH01 QM09 RL29 AD01 IH06, PL03, PL13 PL18 SA16	189 362 325 190 236 227, 355 92 227 236 33 319 190, 362 147 233 294 35 96, 164, 166 166 318	Cho, Chae Ryong Cho, Daeheum Cho, Deo-kyong Cho, Hyun Jin Cho, Hyunduck Cho, Hyun-sung Cho, Jun-hyung Cho, K. Cho, K. H. Cho, Kwang Lae Cho, Lee-hyun Cho, Myeong Woo Cho, Sang-uk Cho, Sunglae Cho, Suyeon Cho, Suyeon Cho, Yong Chan Cho, Yong-chan Cho, Yun Hyun	PN21, QL18 QG17 PG24 PG25 PJ17 RR15 SD06 QC23 SB28 P113 SK27 QN09 PH22 EH05, PN29, PN30, QO22, QO23 GF05 PN20, PN21 PH22, QL18 RR19, RR24	175, 231 211 146 146 157 313 331 196 324 152 354 237 150 70, 176, 177, 244 79 175 150, 231 314
Chen, I. A. Chen, J. Chen, J. M. Chen, Jian Chen, Jin-ming Chen, K.I. Chen, Kuo-chin Chen, L. M. Chen, Lin Chen, Lin Chen, M.j. Chen, P. Chen, Qianping Chen, S. A. Chen, S. C. Chen, Shih-jia Chen, Sui-pin Chen, Szu-wei Chen, Wei-hsiang Chen, Wei-chen Chen, Xiang-bai Chen, Xing-yuan	SB04 QB12 SM11 SB29 QB21 QM27 QK22, SK30 IC03 QK23 QM27 AA04 SA19, SA20 QB21, SM11 PH01 QM09 RL29 AD01 IH06, PL03, PL13 PL18 SA16 QA03	189 362 325 190 236 227, 355 92 227 236 33 319 190, 362 147 233 294 35 96, 164, 166 166 318 183	Cho, Chae Ryong Cho, Daeheum Cho, Deo-kyong Cho, Hyun Jin Cho, Hyunduck Cho, Hyun-sung Cho, Jun-hyung Cho, K. Cho, K. H. Cho, Kwang Lae Cho, Lee-hyun Cho, Myeong Woo Cho, Sang-uk Cho, Sunglae Cho, Sunglae Cho, Yong Chan Cho, Yong Chan Cho, Yong Chan Cho, Yun Hyun Cho, Yun Hyun	PN21, QL18 QG17 PG24 PG25 PJ17 RR15 SD06 QC23 SB28 PI13 SK27 QN09 PH22 EH05, PN29, PN30, QO22, QO23 GF05 PN20, PN21 PH22, QL18 RR19, RR24 RR25	175, 231 211 146 146 157 313 331 196 324 152 354 237 150 70, 176, 177, 244 79 175 150, 231 314 315
Chen, I. A. Chen, J. Chen, J. M. Chen, Jian Chen, Jin-ming Chen, K.I. Chen, Kuo-chin Chen, Kuo-chin Chen, Lin Chen, Lin Chen, M.j. Chen, M.j. Chen, P. Chen, Qianping Chen, S. A. Chen, S. A. Chen, S. C. Chen, Shih-jia Chen, Sui-pin Chen, Sui-pin Chen, Sui-pin Chen, Sui-wei Chen, Wei-hsiang Chen, Wen-chen Chen, Xiang-bai Chen, Xing-yuan Chen, Xumin	SB04 QB12 SM11 SB29 QB21 QM27 QK22, SK30 IC03 QK23 QM27 AA04 SA19, SA20 QB21, SM11 PH01 QM09 RL29 AD01 IH06, PL03, PL13 PL18 SA16 QA03 AD03	189 362 325 190 236 227, 355 92 227 236 33 319 190, 362 147 233 294 35 96, 164, 166 166 318 183	Cho, Chae Ryong Cho, Daeheum Cho, Deo-kyong Cho, Hyun Jin Cho, Hyunduck Cho, Hyun-sung Cho, Jun-hyung Cho, K. Cho, K. H. Cho, Kwang Lae Cho, Lee-hyun Cho, Myeong Woo Cho, Sang-uk Cho, Sunglae Cho, Sunglae Cho, Yung Chan Cho, Yong Chan Cho, Yung Chan Cho, Yun Hyun Cho, Yun Hyun Cho, Yun Hyun Cho, Yunhyun	PN21, QL18 QG17 PG24 PG25 PJ17 RR15 SD06 QC23 SB28 PI13 SK27 QN09 PH22 EH05, PN29, PN30, QO22, QO23 GF05 PN20, PN21 PH22, QL18 RR19, RR24 RR25 KE03, SI13	175, 231 211 146 146 157 313 331 196 324 152 354 237 150 70, 176, 177, 244 79 175 150, 231 314 315 110, 348
Chen, I. A. Chen, J. Chen, J. M. Chen, Jian Chen, Jin-ming Chen, K.I. Chen, Kuo-chin Chen, Kuo-chin Chen, L. M. Chen, Lin Chen, D. Chen, J. Chen, P. Chen, Qianping Chen, S. A. Chen, S. C. Chen, Sui-pin Chen, Sui-pin Chen, Sui-pin Chen, Sui-pin Chen, Sui-pin Chen, Xiang-bai Chen, Xiang-bai Chen, Xing-yuan Chen, Xumin Chen, Xumin	SB04 QB12 SM11 SB29 QB21 QM27 QK22, SK30 IC03 QK23 QM27 AA04 SA19, SA20 QB21, SM11 PH01 QM09 RL29 AD01 IH06, PL03, PL13 PL18 SA16 QA03 AD03 QB12	189 362 325 190 236 227, 355 92 227 236 33 319 190, 362 147 233 294 35 96, 164, 166 166 318 183 35 189	Cho, Chae Ryong Cho, Daeheum Cho, Deo-kyong Cho, Hyun Jin Cho, Hyunduck Cho, Hyun-sung Cho, Jun-hyung Cho, K. Cho, K. H. Cho, Kwang Lae Cho, Lee-hyun Cho, Myeong Woo Cho, Sang-uk Cho, Sunglae Cho, Sunglae Cho, Yung Chan Cho, Yong Chan Cho, Yung Chan Cho, Yun Hyun Cho, Yun Hyun Cho, Yunhyun Choe, Sug-bong Choi, Chang Ho	PN21, QL18 QG17 PG24 PG25 PJ17 RR15 SD06 QC23 SB28 P113 SK27 QN09 PH22 EH05, PN29, PN30, QO22, QO23 GF05 PN20, PN21 PH22, QL18 RR19, RR24 RR19, RR24 RR25 KE03, S113 RK06	175, 231 211 146 146 157 313 331 196 324 152 354 237 150 70, 176, 177, 244 79 175 150, 231 314 315 110, 348 286
Chen, I. A. Chen, J. Chen, J. M. Chen, Jian Chen, Jin-ming Chen, K.I. Chen, Kuo-chin Chen, Kuo-chin Chen, Lin Chen, Lin Chen, M.j. Chen, P. Chen, Qianping Chen, S. A. Chen, S. C. Chen, S. C. Chen, Sui-pin Chen, Szu-wei Chen, Sui-pin Chen, Szu-wei Chen, Wei-hsiang Chen, Xiang-bai Chen, Xiang-bai Chen, Xing-yuan Chen, Xumin Chen, Y. Chen, Yao-jung	SB04 QB12 SM11 SB29 QB21 QM27 QK22, SK30 IC03 QK23 QM27 AA04 SA19, SA20 QB21, SM11 PH01 QM09 RL29 AD01 IH06, PL03, PL13 PL18 SA16 QA03 AD03 QB12 PL03	189 362 325 190 236 227, 355 92 227 236 33 319 190, 362 147 233 294 35 96, 164, 166 166 318 183 35 189 164	Cho, Chae Ryong Cho, Daeheum Cho, Deo-kyong Cho, Hyun Jin Cho, Hyunduck Cho, Hyun-sung Cho, Jun-hyung Cho, K. Cho, K. H. Cho, Kwang Lae Cho, Lee-hyun Cho, Myeong Woo Cho, Sang-uk Cho, Sunglae Cho, Sunglae Cho, Yung Chan Cho, Yong Chan Cho, Yong Chan Cho, Yun Hyun Cho, Yun Hyun Cho, Yun Hyun Cho, Yun Hyun Cho, Yunhyun Cho, Sug-bong Choi, Chang Ho Choi, Chang Ho	PN21, QL18 QG17 PG24 PG25 PJ17 RR15 SD06 QC23 SB28 P113 SK27 QN09 PH22 EH05, PN29, PN30, QO22, QO23 GF05 PN20, PN21 PH22, QL18 RR19, RR24 RR19, RR24 RR25 KE03, SI13 RK06 SK14	175, 231 211 146 146 157 313 331 196 324 152 354 237 150 70, 176, 177, 244 79 175 150, 231 314 315 110, 348 286 352
Chen, I. A. Chen, J. Chen, J. M. Chen, Jian Chen, Jin-ming Chen, K.I. Chen, Kuo-chin Chen, Kuo-chin Chen, Lin Chen, Lin Chen, M.j. Chen, P. Chen, Qianping Chen, S. A. Chen, S. A. Chen, S. C. Chen, Shih-jia Chen, Sui-pin Chen, Szu-wei Chen, Wei-hsiang Chen, Wei-shiang Chen, Xiang-bai Chen, Xiang-bai Chen, Xiang-bai Chen, Ximin Chen, Xumin Chen, Y. Chen, Yao-jung Chen, Ye	SB04 QB12 SM11 SB29 QB21 QM27 QK22, SK30 IC03 QK23 QM27 AA04 SA19, SA20 QB21, SM11 PH01 QM09 RL29 AD01 IH06, PL03, PL13 PL18 SA16 QA03 AD03 QB12 PL03 RD12	189 362 325 190 236 227, 355 92 227 236 33 319 190, 362 147 233 294 35 96, 164, 166 166 318 183 35 189 164 264	Cho, Chae Ryong Cho, Daeheum Cho, Deo-kyong Cho, Hyun Jin Cho, Hyunduck Cho, Hyun-sung Cho, Jun-hyung Cho, K. Cho, K. H. Cho, Kwang Lae Cho, Lee-hyun Cho, Myeong Woo Cho, Sang-uk Cho, Sunglae Cho, Sunglae Cho, Sunglae Cho, Yong Chan Cho, Yong Chan Cho, Yong Chan Cho, Yong Chan Cho, Yun Hyun Cho, Yun Hyun Cho, Yun Hyun Cho, Yunhyun Choe, Sug-bong Choi, Chang Ho Choi, Changsik Choi, Choljin	PN21, QL18 QG17 PG24 PG25 PJ17 RR15 SD06 QC23 SB28 P113 SK27 QN09 PH22 EH05, PN29, PN30, Q022, Q023 GF05 PN20, PN21 PH22, QL18 RR19, RR24 RR19, RR24 RR25 KE03, S113 RK06 SK14 RN23	175, 231 211 146 146 157 313 331 196 324 152 354 237 150 70, 176, 177, 244 79 175 150, 231 314 315 110, 348 286 352 300
Chen, I. A. Chen, J. Chen, J. M. Chen, Jian Chen, Jin-ming Chen, K.I. Chen, Kuo-chin Chen, Kuo-chin Chen, L. M. Chen, Lin Chen, M.j. Chen, M.j. Chen, Qianping Chen, Qianping Chen, S. A. Chen, S. A. Chen, S. C. Chen, Shih-jia Chen, Sui-pin Chen, Szu-wei Chen, Wei-hsiang Chen, Xiang-bai Chen, Xiang-bai Chen, Ximin Chen, Xumin Chen, Y. Chen, Yao-jung Chen, Ye	SB04 QB12 SM11 SB29 QB21 QM27 QK22, SK30 IC03 QK23 QM27 AA04 SA19, SA20 QB21, SM11 PH01 QM09 RL29 AD01 IH06, PL03, PL13 PL18 SA16 QA03 AD03 QB12 PL03 RD12 PC18	189 362 325 190 236 227, 355 92 227 236 33 319 190, 362 147 233 294 35 96, 164, 166 166 318 183 35 189 164 264 129	Cho, Chae Ryong Cho, Daeheum Cho, Deo-kyong Cho, Hyun Jin Cho, Hyunduck Cho, Hyun-sung Cho, Jun-hyung Cho, K. Cho, K. H. Cho, Kwang Lae Cho, Lee-hyun Cho, Myeong Woo Cho, Sang-uk Cho, Sunglae Cho, Sunglae Cho, Sunglae Cho, Sunglae Cho, Yong Chan Cho, Yong Chan Cho, Yong Chan Cho, Yong Chan Cho, Yun Hyun Cho, Yun Hyun Cho, Yun Hyun Cho, Sug-bong Choi, Chang Ho Choi, Changsik Choi, Choljin Choi, Choljin	PN21, QL18 QG17 PG24 PG25 PJ17 RR15 SD06 QC23 SB28 P113 SK27 QN09 PH22 EH05, PN29, PN30, QO22, QO23 GF05 PN20, PN21 PH22, QL18 RR19, RR24 RR25 KE03, S113 RK06 SK14 RN23 PM24	175, 231 211 146 146 157 313 331 196 324 152 354 237 150 70, 176, 177, 244 79 175 150, 231 314 315 110, 348 286 352 300 171
Chen, I. A. Chen, J. Chen, J. M. Chen, Jian Chen, Jin-ming Chen, K.I. Chen, Kuo-chin Chen, Kuo-chin Chen, Lin Chen, Lin Chen, M.j. Chen, P. Chen, Qianping Chen, S. A. Chen, S. A. Chen, S. C. Chen, Shih-jia Chen, Sui-pin Chen, Szu-wei Chen, Wei-hsiang Chen, Wei-shiang Chen, Xiang-bai Chen, Xiang-bai Chen, Xiang-bai Chen, Ximin Chen, Xumin Chen, Y. Chen, Yao-jung Chen, Ye	SB04 QB12 SM11 SB29 QB21 QM27 QK22, SK30 IC03 QK23 QM27 AA04 SA19, SA20 QB21, SM11 PH01 QM09 RL29 AD01 IH06, PL03, PL13 PL18 SA16 QA03 AD03 QB12 PL03 RD12	189 362 325 190 236 227, 355 92 227 236 33 319 190, 362 147 233 294 35 96, 164, 166 166 318 183 35 189 164 264	Cho, Chae Ryong Cho, Daeheum Cho, Deo-kyong Cho, Hyun Jin Cho, Hyunduck Cho, Hyun-sung Cho, Jun-hyung Cho, K. Cho, K. H. Cho, Kwang Lae Cho, Lee-hyun Cho, Myeong Woo Cho, Sang-uk Cho, Sunglae Cho, Sunglae Cho, Sunglae Cho, Yong Chan Cho, Yong Chan Cho, Yong Chan Cho, Yong Chan Cho, Yun Hyun Cho, Yun Hyun Cho, Yun Hyun Cho, Yunhyun Choe, Sug-bong Choi, Chang Ho Choi, Changsik Choi, Choljin	PN21, QL18 QG17 PG24 PG25 PJ17 RR15 SD06 QC23 SB28 P113 SK27 QN09 PH22 EH05, PN29, PN30, Q022, Q023 GF05 PN20, PN21 PH22, QL18 RR19, RR24 RR19, RR24 RR25 KE03, S113 RK06 SK14 RN23	175, 231 211 146 146 157 313 331 196 324 152 354 237 150 70, 176, 177, 244 79 175 150, 231 314 315 110, 348 286 352 300

Choi, Dong Hoon	EH02	70
Choi, Doo-hyun	RR23, RR27	314, 315
Choi, Eun Sang	GJ02, PA21	82, 119
Choi, Eun-sang	SB05	321
Choi, Han-yong	SB31, SB32	325
Choi, Heonhwa	RQ15	309
Choi, Heon-hwa	SC14	327
Choi, Heon-jin	QK11, SL26	225, 360
Choi, Hong Chul	PE17	138
Choi, Hyeok Cheol	QM18, QM20	234
Choi, Hyeok-cheol	RK02	286
Choi, Hyoung Jin	QN03, QN09	236, 237
Choi, Hyoung Joon	QK21, RB16, RC15	227, 256, 261
Choi, Jae-hak	RR31	316
Choi, Jae-hyuk	PB21, RQ15, SC14	123, 309, 327
Choi, Jong-gu	BH02, PL11, RR15	46, 165, 313
Choi, Jong-jin	DJ02	63
Choi, Joon-hwan	DJ02	63
Choi, Jun Woo	KC04	109
Choi, K. Y.	QF13	206
Choi, Ki-young	JA03	98
Choi, Kwang Yong	RP13	305
Choi, Kyong-hoon	RN20	300
Choi, Mahn-soo	SC14	327
Choi, Min-young	SD03	330
Choi, Moo Young	RC08, RE07	260, 267
Choi, Seokhwan	RQ11	308
Choi, Seonghoon	QL08	229
Choi, Seongil	PG24	146
Choi, Seonjun	QP19, RL22, SL19	248, 293, 359
Choi, Sul A	SA09	317
Choi, Sungil	PI10	152
Choi, Won Young	QK13	226
Choi, Y.j.	PH10	148
Choi, Yong Nam	PN20, QL18	175, 231
Choi, Yoon Hyuck	PB16, PB18, PB35, PB36, PB38	122, 123, 125, 126
Choi, Young Ha	RC07	260
Choi, Youngha	SD10, SD11	331
Choi, Youn-seok	DE05, RJ16, SH14	60, 285, 345
Chokprasombat, Komkrich	SL10	357
Choo, Seongmin	PN19	175
Choo, Seong-min	Q010, Q011	242
Chou, Chih-chieh	QA20	185
Chou, T. L.	SM11	362
Choudhary, R.j.	SA14	318
Chowdhury, Prasanta	PK16	163
Chu, Ching-wu	JA02	98
Chu, Yu-hsun	AD01	35
Chuan'an, Cui	Q021	244
Chubykalo-fesenko, O	FC01	73
Chubykalo-fesenko, O.	GD03	78
Chubykalo-fesenko, Oksana	RJ04	283
Chudo, Hiroyuki	PC03	127
Chumak, Andrii V.	SI15	348
Chumakov, Andrey	SL27	360
Chun, Byong Sun	IH04, PJ14, RK18, SH13	96, 157, 288, 345
Chun, Sae Hwan	JB02, PA09	99, 117
Chun, Sae-hwan	PA04	116
Chun, Sungwoo	QP19, RL22, SL19	248, 293, 359

Chung, Chung-hou	DB04, RE06
Chung, D. Y.	GA01
Chung, J H.	PH03
Chung, Jaeho	PA01
Chung, Jae-ho	CI03, JB02, PA03, PA
Chung, K. C.	SB28
Chung, Kookchae	RN23
Chung, Kook-chae	QM22
Chung, Sunjae	PK13
Chung, Sunjea	CE02
Churagulov, Bulat	ED06
Chureemart, Phanwadee	SJ04
Cibert, Joel	DD01
Ciccacci, Franco	PL21
Cichorek, Tomasz	PE19, SB26
Cieslak, Jakub	QH18
Ciftja, Orion	RG01
Cinal, Marek	AD02
Cioangher, M.	SA11
Cisarova, Ivana	BB04, PD19
Ciubotaru, Florin	SI15
Ciuta, G.	JH02
Civale, L.	QB31, SB12
Cizmar, E.	QF02, QF03
Cizmas, Corneliu Bazil	PM06
Clarke, Simon J.	HA04
Clarke, Simon. J.	QG19
Claudia, Felser	SD15
Claydon, J S	QO24
Clemente-juan, J. M.	BC01
Coey, Michael	PO32
Coisson, Marco	AG02, DG04
Colak, Levent	KB04
Coleman, Piers	HP22
Colineau, Eric	RE03, SE19
Colis, Silviu	PI04
Collin, Sophie	HE04
Collins, S. P.	QN01
Collins, Steve	DH03
Conder, Kazimierz	PI21, PO17, QF23
Cong, Pham Thanh	HC04, QE14, RI13
Cong, Rihong	PH02
Continentino, Mucio A.	KG04
Continenza, Alessandra	QG25
Coqblin, B.	QD10
Coqblin, Bernard	PE01
Cordoba, Rosa	DG05
Cornell, N.	QB12
Coronado, Eugenio	BC01
Corrias, Anna	CB04
Cossu, Fabrizio	QP01
Costache, M. V.	DI01
Cote, Michel	PD05
Cottam, Michael	QG12, SL05
Cowburn, Russell	KE02
Cowlam, Neil	PG01
Crassous, Arnaud	AE03
Crichton, Wilson	RB20
Cristiani, G.	ED05
Crooker, Scott A	JD01

Cros, V.	CE03, CE06, ID02, RJ07	51, 52, 93, 283
Cros, Vincent	HF01, ID03, IE04	87, 93, 94
Crow, D. J. G.	SF15	339
Crowell, Paul A	CG01	53
Cruz Irisson, Miguel	SK15	352
Cubukcu, Murat	Al04, GE01, QK15, QK16, QK17, SL24	39, 78, 226, 360
Cuello, Gabriel	JJ04	106
Curiale, Javier	SH02	343
Curlik, Ivan	QD08, RD21	197, 265
Custardoy, Laura	RN08	298
Custers, J.	CD01	50
Custers, Jeroen	BB04, PD18, PD19	42, 134, 135
Cywinski, Robert	JA05	99
D		
Da Silva, Douglas	PB33	125
Da Silva, Douglas Langie	PB09, PB11, RB09, SB18	121, 255, 323
Da Silva, Mathieu	KF02	110
Da Silva, R. C.	QM13	233
Daadmehr, Vahid	PB23	123
Dabkowska, H.	GF01	79
Dabkowska, Hanna A	JD01	100
Dabkowski, A.	GF01	79
Dabrowski, Maciej	AD02, El01	35, 70
Dadras, Sedigheh	PB23	123
Daegeun, Park	RR20	314
Dai, Bing	RL08	291
Dai, Jia	QB07	188
Dai, Jianhui	AB04, IJ04	34, 98
Dai, Jianming	SA01	316
Dai, Pengcheng	GA02, RB10, RB11	75, 255
Dai, Yuqiang Daimon, Hiroshi	SA01 SK13	316 352
Damay, Francoise		100, 251
D'amico, Irene	JB05, RA11 QK09, RL12, SJ04	225, 291, 349
Dan, Nguyen Huy	AJ06	40
Dang, Dung Duc	EH05	70
Dang, Hue Thi Minh	RN19	299
Dang, N. V.	RA01	249
Danis, Stanislav	QI01	217
Danzenbacher, Steffen	CD03, PD12	51, 133
Daou, Ramzi	JD01	100
Dariani, R. S.	QP23	249
Das, A.	PI12	152
Das, Kaustuv	HD04	85
Das, Pintu	QA11, QJ11	184, 222
Das, Pranab Kumar	QJ17	223
Das, S.	PL07, RH02	165, 275
Das, S. D.	GA01	75
Das, Soma	SO30, SO31	371
Das, Tanmoy	HI03, QB08	89, 188
Dasari, Venkateswarlu	SJ08, SL23	350, 359
Davesne, Vincent	GI03	82
Davis, J. C. Seamus	CA02	48
Davison, Toby	EH04	70
De Bergevin, Francois	AF01	36
De Brion, Sophie	KB06	108
De Julian Fernandez, Cesar	CB04, SI11	49, 347 226

De La Fuente, Jesus M	RN08	298
De La Torre, A	1102	97
De Leo, Natascia	AG02	37
De Loubens, Gregoire	ID03	93
De Matteis, Laura	RN08	298
De Pedro, I.	GH05	81
De Pedro, Imanol	JJ04	106
De Ranieri, Elisa	IE03	93
De Souza, Mariano	QJ11, SG19	222, 342
De Teresa, J. M.	KC05	109
De Teresa, Jose Maria	DG05, SM16	61, 363
Decker, Regis	GH02	81
Decourt, Rodolphe	EH03	70
Deen, P P	SE03, SE05	333, 334
Deguchi, Hiroyuki	PB25, PG09, PG10,	124, 143, 144, 211,
	QG19, RN03, SA03	297, 316
Deguchi, Kazuhiko	PD14, PD20, RD05	134, 135, 262
Deheer, Walt A.	PN16	174
Deisenhofer, Joachim	PA23, RH24	119, 278
Dejneka, Alexandr	AH02, BH03	38, 46
Dekadjevi, David Toyo	SA24	320
Del Moral, A	SF09	338
Del Moral, Agustin	SE20	336
Del Real, Rafael Perez	DG03	61
Del Val, Juanjo	PJ25	158
Delfanazari, Kaveh	KJ01	113
Delley, B.	PC01	126
Delmas, Claude	EH03	70
Delmo, Michael P.	PN12	174
Delval, Juanjose	DJ03	63
Demidenko, Olga	QG05	209
Demidov, Evgeny Sergeevich	SG17	342
Demishev, S. V.	PE12, QC22, RH07	138, 195, 275
Demishev, Sergey	ED06, KI05, QC20,	67, 113, 195, 340
bernisher, bergej	SG07, SG08	0,, 1,0, 1,0, 0,0,0,0
Demishev, Sergey V.	QC19	195
Demmel, F	SE03, SE05	333, 334
Dempsey, Nora	JH02	104
Demsar, Jure	HA02	83
Deng, Guochu	PI21, PO25	154, 181
Denisovsky, Andrey N.	Q012	242
Deorani, Praveen	EG03, PJ08, SI02	69, 156, 346
Deranlot, Cyril	HE04	86
Deranlot, Cyrile	CG03	53
Derlet, Peter	QF20	207
Deshpande, N. G.		357
Desilets-benoit, Alexandre	SL08, SL09 CD04	51
		354
Detlefs, B. Devereaux, T.	SK24 EB01	554 64
Devillers, Thibaut	BE03, DD01	44, 58
Devlin, Eamon	BJ02, RO08	47, 302
Dewhurst, Charles	DF04	60
Dey, G K	PM01	167
Dhar, Sudesh Kumar	QJ17, RB24, SE21	223, 258, 336
Dhawan, M. S.	SL25	360
Dho, Joonghoe	PL15, QM21	166, 235
Di Marco, Igor	RG06	272
Diaconu, Andrei	RB12	256
Dias, Fabio	PB10, PB33	121, 125
Dias, Fabio Teixeira	PB09, PB11, RB09, SB18	121, 255, 323

Didukh, Leonid	RG13	274
Dieny, Bernard	BD02, SI09	43, 347
Dietl, T.	BE02	44
Dietl, Tomasz	BE03, DD02	44, 58
Dilley, N.	IA05	91
Dilullo, Andrew	GH02	81
Ding, Jun	DD03, QP22, SL06, SM14	58, 249, 357, 363
Dinh, Van-chau	SG10	341
Dinia, Aziz	PI04	151
Dirks, Andreas	FE03	75
Divis, Martin	HJ04, QC21, QJ02	90, 195, 221
Djuhana, Dede	RA15, SH01	251, 343
Dkhil, Brahim	QA27	187
Dluzewski, Piotr	DD02, QB30	58, 192
Do, Anh Thi Kim	CJ05	56
Do, Bang	PJ09	156
Do, Thang Viet	CJ05	56
Dobkowska, Sylwia	BE03	44
Dobrynin, Alexey	PA15	118
Doenni, Andreas	JJ05, PH02	106, 147
Doerr, Mathias	PH04	147
Dogan, Fatih	CG05, IE01	53, 93
Doh, Jaewon	PO14	179
Doh, Yong-joo	SD12, SD14, SO21	332, 370
Doi, Masaaki	SM03, SN03	361, 364
Doi, Naohiro	QE13	201
Dolgov, O. V.	GA03	75
Dolia, S. N.	SL25	360
Dolinsek, Janez	SL11	357
Dominguez, Claribel	SA12	318
Doria, Mauro Melchiades	RB09	255
Dos Santos, Maria Elenice	QA06	183
Dou, Shixue	EF02, QG09, SD01	68, 209, 330
Drechsler, SI.	QF18, QF22	206, 207
Dreiser, Jan Gui-hyon	BC02	42
Drews, Andre	FD02	74
Driver, Sarah L	EF03	68
Drobosyuk, Mikhail	PM22, SO23	171, 370
Drube, W.	SK24	354
Drymiotis, Fivos R	QO09	242
Du Plessis, Paul De Villiers	RD23	265
Du, An	RK16	288
Du, Jun	PL01, RA13	164, 251
Du, M-h	GF01	79
Dua, P.	QH26	217
Duan, Truong Cong	HC02	84
Duanmu, Qingyong	PK03	160
Dubenko, Igor	CH01	54
Dubey, Sheshmani K.	PN01	172
Dubi, Yoni	KG01	111
Dubiel, Stanislaw M	PG03, QH18	142, 215
Dubowik, J.	QG27	212
Duc Dung, Dang	PJ15	157
Duc, Fabienne	RQ16	309
Duc, Nguyen Huu	AJ06	40
Ducruet, Clarisse	AI04	39
Dudzik, Esther	EC04	66
Duerr, George	GG02	80

SE20

336

De La Fuente, Cesar

## AUTHOR INDEX

Dufour, Catherine Dugaev, V K Dukhnenko, A. Dukhnenko, A. V. Dumas, R. K. Dumas-bouchiat, F. Dumesnil, Karine Dumont, Matthieu F. Dumont, Yves Duncan, William J Dung, Dang Duc

Dung, N. H. Dung, N.h. Dunn, J. Dunne, Peter Dunsiger, Sarah Ruth Duo', Lamberto Duong, Thiet Van Duong, Tuan Anh Dupuis, Veronique Duque, Jose G. S. Durakiewicz, Thomasz Durr, Hermann Andreas Dussaux, A. Dvornik, Mykola Dwevedi, Sandhya Dwivedi, R. K. Dyadkin, Vadim A. Dyadkina, Ekaterina Dyakina, Veronika

#### Ε

Ebels, Ursula Ebert, Hubert Ebihara, Takao Ebisu, Shuji Ebrahimian, Neda Echevarria-bonet, C. Eckerlebe, Helmut Eckert, James C Eckold, Gotz Eckstein, Martin Eddrief, Mahmoud Eder, Robert Edmonds, Kevin W Egashira, Yusaku Egawa, Genta Egetenmeyer, Nikola Eggert, Sebastian Eguchi, Gaku Eguchi, Ryo Ehlers, Georg Ehmler, H. Einaga, Yasuaki Eisaki, H. Eisaki, Hiroshi

El Hafid, Hassan

QC07

QB28

EH03

CA02, DA04, QA05,

QA07, QA09, QB02,

QB15, QB29, RH11

DD05, KD02, KF02 59, 109, 110 SB21 JC04 PE12 CE02, PK13 JH02 DD05, KD02, KF02 KB03 SM07 QH12, SF05 PN29, PN30, QO22, Q023 SO16 SO08 HA05 84 PO32 GF01 79 PL21 EH05 70 EH05 70 KB05, QN06, RK03 PI19 SE08 GD02 77 CE03, CE06, ID02 EH04, JG01 PG19 RA27, SM18 DF04 60 QM08 PC07 SI09 1103 97 RA19 QH14, QO07 SB07, SC09 QD10 SL27 PK09, QM01 AA05 33 EB02 65 QL06 PF14, RF03 SH09 PE10 RP06 CD04 51 HC04 85 PC20 JE04 CC03 50 RQ21

Author Index

El Mousaoui, Souliman         RK17         288         Fan, C.         ICO3         92           Elgazar, Saad         RG05         272         Fan, Raymond         PA15         11           El-hagary, M A         QL21         231         Fan, Xiu Xiu         SL16         35           El-hagary, Magdy         RB29         258         Fan, Zhu         SH10         36           Elisev, Andrey         SL27         360         Fang, Dong         IEO3         92           Elizabeth, Suja         PH04         147         Fang, Dong         IEO3         92           Elizabeth, Suja         PH04         147         Fang, Nakun         R007, SN11         30           Elimers, H J         QM12, RN12, RP20         233, 298, 306         Farle, M         R08         30           Elmers, H J         QM12, RN12, RP20         233, 298, 306         Farle, M         R08         30           Eliorid, Rachel         RE03         266         Fauque, Benoit         K104         11           Elsayed, H. M.         RN01         296         Fauzulin, Rafael         PM22, S023         17           Elsayed, H. M.         RN01         296         Fauzulin, Rafael         PM22, S023         17 <th>18 158 144 13 101, 366 19 117 104 184, 343 96</th>	18 158 144 13 101, 366 19 117 104 184, 343 96
Elgazar, Saad         RG05         272         Fan, Raymond         PA15         11           El-hagary, M A         QL21         231         Fan, Xiu Xiu         SL16         35           El-hagary, Maqdy         RE29         258         Fan, Zhe         SH10         36           Eliseev, Andrey         SL27         360         Fang, Dong         IEO3         93           Elizabeth, Suja         PH04         147         Fang, Yikun         R007, SN11         30           Elkaim, Erik         RA11         251         Fantechi, Elvira         CB04         46           Elisir, Tim         JE01         101         Fantechi, Elvira         CB04         46           Eliners, H J         QM12, RN12, RP20         233, 298, 306         Farle, M         RP08         30           Elmers, H J         QM12, RN12, RP20         233, 298, 306         Farle, M         R04         11           Eloiner, Rachel         RE03         266         Faugue, Benoit         KI04         11           Elsayed, H. M.         RN01         296         Fayzulin, Rafael         PM22, S023         17           Elwindari, Nastiti         B103         48         Fdez-gubieda, M. L.         QM06         23	18 158 144 13 101, 366 19 117 104 184, 343 96
Elhagan, M A         QL21         231         Fan, Xiu Xiu         SL16         35           El-hagan, Magdy         RB29         258         Fan, Zhe         SH10         34           Eliseev, Andrey         SL27         360         Fang, Dong         E03         99           Elizabeth, Suja         PH04         147         Fang, Yikun         R007, SN11         30           Elikabeth, Suja         PH04         147         Fang, Yikun         R007, SN11         30           Elikatim, Erik         RA11         251         Farlan, M. Arshad         QH27         21           Elmers, HJ         QM12, RN12, RP20         233, 298, 306         Farle, M         RP08         30           Elior, Itans-joachim         AG04, CH04         38, 54         Fassbender, Juergen         R111, SH04         28           Eloi, Jean-charles         KB01         107         Faulhaber, Enrico         QC24         19           Eloidi, Rachel         RE03         266         Faugue, Benoit         KI04         11           Elsayed, H. M.         RN01         296         Faugue, Banoti         KI04         13           Elsayed, H. M.         RU03         54         Fecher, G. H.         QG24, SK24 <td< td=""><td>158 144 13 101, 366 19 117 104 184, 343 96</td></td<>	158 144 13 101, 366 19 117 104 184, 343 96
El-hagar, Magdy         RB29         258         Fan, Zhe         SH10         34           Eliseev, Andrey         SL27         360         Fang, Dong         IE03         93           Elizabeth, Suja         PH04         147         Fang, Yikun         R007, SN11         36           Elizabeth, Suja         PH04         147         Fang, Yikun         R007, SN11         36           Elikarm, Erik         RA11         251         Fartechi, Kura         CB04         46           Ellis, Tim         JE01         101         Fartechi, M. Arshad         QH27         21           Elmers, Hans-joachim         AG04, CH04         38, 54         Fasbender, Juergen         R111, SH04         28           Eloi, Jean-charles         KB01         107         Fauluber, Enrico         QC24         15           Eloirdi, Rachel         RE03         266         Faugue, Benoit         KJ04         11           Elswed, H. M.         RN01         296         Fayzulin, Rafael         PM22, SO23         17           Elwindari, Nastiti         BJ03         48         Fdez-gubieda, M. L.         QM06         23           Emam-ismail, M         QL21         231         Fear, Natalia         PK09, QM01	44 13 101, 366 19 117 104 184, 343 96
Eliseev, Andrey         SL27         360         Fang, Dong         E03         93           Elizabeth, Suja         PH04         147         Fang, Yikun         R007, SN11         30           Elkaim, Erik         RA11         251         Fantechi, Elvira         CB04         49           Eliks, Tim         JE01         101         Farhan, M. Arshad         QH27         21           Elmers, H J         QM12, RN12, RP20         233, 298, 306         Farle, M         R008         30           Elmers, Hans-joachim         AG04, CH04         38, 54         Fasebneder, Juergen         R111, SH04         28           Eloir, Jean-charles         KB01         107         Faulhaber, Enrico         QC24         19           Eloirdi, Rachel         RE03         266         Fauque, Benoit         K104         11           Elsayed, H. M.         RN01         296         Faugue, Benoit         K104         11           Elsayed, H. M.         RN01         296         Faugue, Benoit         K104         11           Elsayed, H. M.         RN01         233         73         Feder, G. H.         QM06         233           Emam-ismail, M         QL21         231         Fead, Ref Redrand H.	13 101, 366 19 117 104 184, 343 96
Elizabeth, Suja         PH04         147         Fang, Yikun         R007, SN11         30           Elkaim, Erik         RA11         251         Fantechi, Elvira         CB04         49           Ellins, Tim         JE01         101         Farhan, M. Arshad         QH27         21           Elmers, H J         QM12, RN12, RP20         233, 298, 306         Farle, M         RP08         30           Elmers, Hans-joachim         AG04, CH04         38, 54         Fassbender, Juergen         R111, SH04         226           Eloid, Jean-charles         KB01         107         Faulhaber, Enrico         QC24         19           Eloid, Rachel         RE03         266         Fauque, Benoit         KJ04         11           Elayed, H. M.         RN01         296         Fayzullin, Rafael         PM22, SO23         17           Elwindari, Nastiti         BJ03         48         Fde2-gubieda, M. L.         QM06         23           Emmel, Mirko         CH03         54         Fecher, G. H.         QG24, SK24         21           Emoto, Shun         PK12         162         Feder, R         RP20         30           Enders, Axel         AD03         35         Fecher, G. H.         CH03, QN19,	01, 366 9 17 04 84, 343 96
Elkaim, Erik         RA11         251         Fantechi, Elvira         CB04         455           Ellis, Tim         JE01         101         Farhan, M. Arshad         QH27         21           Elmers, H J         QM12, RN12, RP20         233, 298, 306         Farle, M         RP08         30           Elmers, Hans-joachim         AG04, CH04         38, 54         Fassbender, Juergen         R111, SH04         28           Eloid, Jean-charles         KB01         107         Faulhaber, Enrico         QC24         19           Eloidi, Rachel         RE03         266         Fauque, Benoit         KJ04         11           Elsayed, H. M.         RN01         296         Fayzulin, Rafael         PM22, S023         17           Elwindari, Nastiti         BJ03         48         Fdez-gubieda, M. L.         QM06         23           Ermam-ismail, M         QL21         231         Fear, Natalia         PK09, QM01         16           Ermoto, Akira         RP17, SK21         306, 353         Fecher, G. H.         QG24, SK24         21           Emoto, Shun         PK12         162         Feder, R         RP20         30           Endo, Yasushi         QM07, RK04         232, 286         Feigurin, Adria	19 117 104 184, 343 96
Ellis, Tim         JE01         101         Farhan, M. Arshad         QH27         21           Elmers, H J         QM12, RN12, RP20         233, 298, 306         Farle, M         RP08         30           Elmers, Hans-joachim         AG04, CH04         38, 54         Fassbender, Juergen         R111, SH04         28           Eloi, Jean-charles         KB01         107         Faulhaber, Enrico         QC24         19           Eloirdi, Rachel         RE03         266         Fauque, Benoit         KJ04         11           Elsayed, H. M.         RN01         296         Fayzullin, Rafael         PM22, S023         17           Elwindari, Nastiti         BJ03         48         Fde2-gubieda, M. L.         QM06         23           Emam-ismail, M         QL21         231         Fear, Natalia         PK09, QM01         66           Emmel, Mirko         CH03         S4         Fecher, G. H.         QG24, SK24         21           Emoto, Akira         RP17, SK21         306, 353         Fecher, R         RP20         30           Endor, Keita         PG04         143         Fedorov, An         S004         33           Endo, Keita         QM07, RK04         232, 286         Feiguin, Adrian	17 104 184, 343 96
Elmers, H J         QM12, RN12, RP20         233, 298, 306         Farle, M         RP08         30           Elmers, Hans-joachim         AG04, CH04         38, 54         Fassbender, Juergen         R111, SH04         28           Eloi, Jean-charles         KB01         107         Faulhaber, Enrico         QC24         19           Eloirdi, Rachel         RE03         266         Fauque, Benoit         KJ04         11           Elsayed, H. M.         RN01         296         Fayzullin, Rafael         PM22, SO23         17           Elwindari, Nastiti         BJ03         48         Fde2-gubieda, M. L.         QM06         23           Ermam-ismail, M         QL21         231         Fear, Natalia         PK09, QM01         16           Ermoto, Akira         RP17, SK21         306, 353         Fecher, G. H.         QG24, SK24         21           Ernoto, Shun         PK12         162         Fedorov, A.         SD04         33           Endor, Keita         PG04         143         Fedorov, A.         SD04         33           Endo, Keita         QM07, RK04         232, 286         Feigur, C.         QG24, SK24         21           Enoki, Kentaro         PD13         134         QN19, RC10	904 184, 343 96
Elmers, Hans-joachim         AG04, CH04         38, 54         Fassbender, Juergen         R11, SH04         28           Eloi, Jean-charles         KB01         107         Faulhaber, Enrico         QC24         19           Eloirdi, Rachel         RE03         266         Fauque, Benoit         KJ04         11           Elsayed, H. M.         RN01         296         Fayzullin, Rafael         PM22, SO23         17           Elwindari, Nastiti         BJ03         48         Fdez-gubieda, M. L.         QM06         23           Emmer, Mirko         CH03         54         Fear, Natalia         PK09, QM01         16           Emoto, Akira         RP17, SK21         306, 353         Fecher, G. H.         QG24, SK24         21           Emoto, Shun         PK12         162         Feder, R         RP20         30           Endor, Kata         PG04         143         Fedorov, A.         SD04         33           Endo, Yasushi         QM07, RK04         232, 286         Feiguin, Adrian         JD01         10           Endok, Yasushi         QM07, RK04         232, 286         Felser, C.         QG24, SK24         21           Englek, Josefin         SM02         361         Felser, C.	84, 343 96
Eloi, Jean-charles         KB01         107         Faulhaber, Enrico         QC24         19           Eloirdi, Rachel         RE03         266         Fauque, Benoit         KJ04         11           Elsayed, H. M.         RN01         296         Fayzullin, Rafael         PM22, SO23         17           Elwindari, Nastiti         BJ03         48         Fdez-gubieda, M. L.         QM06         23           Ernam-ismail, M         QL21         231         Fear, Natalia         PK09, QM01         16           Ernmel, Mirko         CH03         54         Fecher, G. H.         QG24, SK24         21           Ernoto, Akira         RP17, SK21         306, 353         Fecher, Gerhard H.         CH03, QN19, SD15         54           Ernoto, Shun         PK12         162         Feder, R         RP20         30           Enders, Axel         AD03         35         Fedorov, A.         SD04         33           Endo, Yasushi         QM07, RK04         232, 286         Feiguin, Adrian         JD01         10           Endok, Keita         BI02         47         Felser, Claudia         CH03, CH04, HP72, I03, 54         54           Enoki, Kentaro         PD13         134         QN19, RC10	96
Eloirdi, Rachel         RE03         266         Fauque, Benoit         K.104         11           Elsayed, H. M.         RN01         296         Fayzullin, Rafael         PM22, S023         17           Elwindari, Nastiti         BJ03         48         Fdez-gubieda, M. L.         QM06         23           Emam-ismail, M         QL21         231         Fear, Natalia         PK09, QM01         16           Emmel, Mirko         CH03         54         Fecher, G. H.         QG24, SK24         21           Emoto, Akira         RP17, SK21         306, 353         Fecher, Gerhard H.         CH03, QN19, SD15         54           Emoto, Shun         PK12         162         Feder, R         RP20         30           Enders, Axel         AD03         35         Fedorov, A.         SD04         33           Endo, Keita         PG04         143         Fedorov, Dmitry         EE01         67           Endoh, Tetsuo         BI02         47         Felser, C.         QG24, SK24         21           Enoki, Kentaro         PD13         134         QN19, RC10         QN19, RC10         24           Enoki, Kentaro         PB29, PB30, PB32, QH10         124, 125, 214         Feng, Chunmu         AB04 </td <td></td>	
Elsayed, H. M.         RN01         296         Fayzullin, Rafael         PM22, SO23         17           Elwindari, Nastiti         BJ03         48         Fdez-gubieda, M. L.         QM06         23           Emam-ismail, M         QL21         231         Fear, Natalia         PK09, QM01         16           Emmel, Mirko         CH03         54         Fecher, G. H.         QG24, SK24         21           Emoto, Akira         RP17, SK21         306, 353         Fecher, Gerhard H.         CH03, QN19, SD15         54           Emoto, Shun         PK12         162         Feder, R         RP20         30           Enders, Axel         AD03         35         Fedorov, A.         SD04         33           Endo, Keita         PG04         143         Fedorov, Dmitry         EE01         67           Endoh, Yasushi         QM07, RK04         232, 286         Feiguin, Adrian         JD01         10           Endoh, Tetsuo         BI02         47         Felser, C.         QG24, SK24         21           Enoki, Kentaro         PD13         134	14
Elwindari, Nastiti         BJ03         48         Fdez-gubieda, M. L.         QM06         23           Emam-ismail, M         QL21         231         Fear, Natalia         PK09, QM01         16           Emmel, Mirko         CH03         54         Fecher, G. H.         QG24, SK24         21           Emoto, Akira         RP17, SK21         306, 353         Fecher, Gerhard H.         CH03, QN19, SD15         54           Emoto, Shun         PK12         162         Feder, R         RP20         30           Enders, Axel         AD03         35         Fedorov, A.         SD04         33           Endo, Keita         PG04         143         Fedorov, Dmitry         EE01         67           Endor, Yasushi         QM07, RK04         232, 286         Feiguin, Adrian         JD01         10           Endor, Tetsuo         BI02         47         Felser, C.         QG24, SK24         21           Engelke, Josefin         SM02         361         Felser, Claudia         CH03, CH04, HP72, II03, 54         54           Enoki, Kentaro         PD13         124, 125, 214         Feng, Chunmu         AB04         34           Entari, Shiro         BH05, BH06, QL15         46, 47, 230         Feng, Kai <td>14</td>	14
Emam-ismail, M         QL21         231         Fear, Natalia         PK09, QM01         16           Emmel, Mirko         CH03         54         Fecher, G. H.         QG24, SK24         21           Emoto, Akira         RP17, SK21         306, 353         Fecher, Gerhard H.         CH03, QN19, SD15         54           Emoto, Shun         PK12         162         Feder, R         RP20         30           Enders, Axel         AD03         35         Fedorov, A.         SD04         33           Endo, Keita         PG04         143         Fedorov, Dmitry         EE01         67           Endoh, Yasushi         QM07, RK04         232, 286         Feiguin, Adrian         JD01         10           Endoh, Tetsuo         BI02         47         Felser, C.         QG24, SK24         21           Engelke, Josefin         SM02         361         Felser, Claudia         CH03, CH04, HP72, II03, 54         54           Enoki, Kentaro         PD13         134         QN19, RC10         24         24         24           Enoki, Masanori         PB29, PB30, PB32, QH10         124, 125, 214         Feng, Chunmu         AB04         34           Entani, Shiro         BH05, BH06, QL15         46, 47, 230	71, 370
Emmel, Mirko         CH03         54         Fecher, G. H.         QG24, SK24         21           Emoto, Akira         RP17, SK21         306, 353         Fecher, Gerhard H.         CH03, QN19, SD15         54           Emoto, Shun         PK12         162         Feder, R         RP20         30           Enders, Axel         AD03         35         Fedorov, A.         SD04         33           Endo, Keita         PG04         143         Fedorov, Dmitry         EE01         67           Endo, Yasushi         QM07, RK04         232, 286         Feiguin, Adrian         JD01         10           Enders, Josefin         SM02         361         Felser, C.         QG24, SK24         21           Engelke, Josefin         SM02         361         Felser, C.         QG24, SK24         24           Enoki, Kentaro         PD13         134         QN19, RC10         24         24         24           Entani, Shiro         BH05, BH06, QL15         46, 47, 230         Feng, Chunmu         AB04         34           Entani, Shiro         BH05, BH06, QL15         46, 47, 230         Feng, Kai         RA02         24           Entel, Peter         QH05         213         Feng, Cai         RA	32
Emoto, Akira         RP17, SK21         306, 353         Fecher, Gerhard H.         CH03, QN19, SD15         54           Emoto, Shun         PK12         162         Feder, R         RP20         30           Enders, Axel         AD03         35         Fedorov, A.         SD04         33           Endo, Keita         PG04         143         Fedorov, Dmitry         EE01         67           Endo, Yasushi         QM07, RK04         232, 286         Feiguin, Adrian         JD01         10           Endoh, Tetsuo         BI02         47         Felser, C.         QG24, SK24         21           Engelke, Josefin         SM02         361         Felser, Claudia         CH03, CH04, HP72, II03, 54         54           Enoki, Kentaro         PD13         134         QN19, RC10         24         24           Enoki, Masanori         PB29, PB30, PB32, QH10         124, 125, 214         Feng, Chunmu         AB04         34           Enteni, Shiro         BH05, BH06, QL15         46, 47, 230         Feng, Haibo         RO04, SN11         30           Entel, Peter         QH05         213         Feng, Kai         RA02         24           Eom, Jonghwa         PJ30, QK06         159, 225         Feng, Shipi	61, 231
Emoto, Shun         PK12         162         Feder, R         RP20         30           Enders, Axel         AD03         35         Fedorov, A.         SD04         33           Enders, Axel         PG04         143         Fedorov, Dmitry         EE01         67           Endo, Keita         PG04         232, 286         Feiguin, Adrian         JD01         10           Endoh, Tetsuo         BI02         47         Felser, C.         QG24, SK24         21           Engelke, Josefin         SM02         361         Felser, Claudia         CH03, CH04, HP72, II03, 54         54           Enoki, Kentaro         PD13         134         QN19, RC10         244         244           Entani, Shiro         BH05, BH06, QL15         46, 47, 230         Feng, Haibo         RO04, SN11         30           Entel, Peter         QH05         213         Feng, Kai         RA02         24           Eom, Jonghwa         PJ30, QK06         159, 225         Feng, Qi         RN18         29           Eom, Man Jim         SD07         331         Feng, Shiping         PB28         12	12, 354
Enders, Axel         AD03         35         Fedorov, A.         SD04         33           Endo, Keita         PG04         143         Fedorov, Dmitry         EE01         67           Endo, Yasushi         QM07, RK04         232, 286         Feiguin, Adrian         JD01         10           Endoh, Tetsuo         BI02         47         Felser, C.         QG24, SK24         21           Engelke, Josefin         SM02         361         Felser, Claudia         CH03, CH04, HP72, II03, 54         54           Enoki, Kentaro         PD13         134         QN19, RC10         244         244         244           Entani, Shiro         BH05, BH06, QL15         46, 47, 230         Feng, Chunmu         AB04         344           Entel, Peter         QH05         213         Feng, Kai         RA02         244           Eom, Jonghwa         PJ30, QK06         159, 225         Feng, Qi         RN18         29           Eom, Man Jim         SD07         331         Feng, Shiping         PB28         12	4, 239, 332
Endo, Keita         PG04         143         Fedorov, Dmitry         EE01         67           Endo, Yasushi         QM07, RK04         232, 286         Feiguin, Adrian         JD01         10           Endoh, Tetsuo         BI02         47         Felser, C.         QG24, SK24         21           Engelke, Josefin         SM02         361         Felser, Claudia         CH03, CH04, HP72, II03, 54         54           Enoki, Kentaro         PD13         134         QN19, RC10         54           Enoki, Masanori         PB29, PB30, PB32, QH10         124, 125, 214         Feng, Chunmu         AB04         34           Entani, Shiro         BH05, BH06, QL15         46, 47, 230         Feng, Haibo         R004, SN11         30           Entel, Peter         QH05         213         Feng, Kai         RA02         24           Eom, Jonghwa         PJ30, QK06         159, 225         Feng, Qi         RN18         29           Eom, Man Jim         SD07         331         Feng, Shiping         PB28         12	06
Endo, Yasushi         QM07, RK04         232, 286         Feiguin, Adrian         JD01         10           Endoh, Tetsuo         BI02         47         Felser, C.         QG24, SK24         21           Engelke, Josefin         SM02         361         Felser, Claudia         CH03, CH04, HP72, II03, 54         54           Enoki, Kentaro         PD13         134         QN19, RC10         74           Enoki, Masanori         PB29, PB30, PB32, QH10         124, 125, 214         Feng, Chunmu         AB04         34           Entani, Shiro         BH05, BH06, QL15         46, 47, 230         Feng, Haibo         R04, SN11         30           Entel, Peter         QH05         213         Feng, Kai         RA02         24           Eom, Jonghwa         PJ30, QK06         159, 225         Feng, Qi         RN18         29           Eom, Man Jim         SD07         331         Feng, Shiping         PB28         12	30
Endo, Yasushi         QM07, RK04         232, 286         Feiguin, Adrian         JD01         10           Endoh, Tetsuo         BI02         47         Felser, C.         QG24, SK24         21           Engelke, Josefin         SM02         361         Felser, Claudia         CH03, CH04, HP72, II03, 54         54           Enoki, Kentaro         PD13         134         QN19, RC10         24         124, 125, 214         Feng, Chunmu         AB04         34           Entani, Shiro         BH05, BH06, QL15         46, 47, 230         Feng, Haibo         RO04, SN11         30           Entel, Peter         QH05         213         Feng, Kai         RA02         24           Eom, Jonghwa         PJ30, QK06         159, 225         Feng, Qi         RN18         29           Eom, Man Jim         SD07         331         Feng, Shiping         PB28         12	7
Endoh, Tetsuo         BIO2         47         Felser, C.         QG24, SK24         21           Engelke, Josefin         SM02         361         Felser, Claudia         CH03, CH04, HP72, II03, 54         54           Enoki, Kentaro         PD13         134         QN19, RC10         54           Enoki, Masanori         PB29, PB30, PB32, QH10         124, 125, 214         Feng, Chunmu         AB04         34           Entani, Shiro         BH05, BH06, QL15         46, 47, 230         Feng, Haibo         RO04, SN11         30           Entel, Peter         QH05         213         Feng, Kai         RA02         24           Eom, Jonghwa         PJ30, QK06         159, 225         Feng, Qi         RN18         29           Eom, Man Jim         SD07         331         Feng, Shiping         PB28         12	00
Engelke, Josefin         SM02         361         Felser, Claudia         CH03, CH04, HP72, II03, 54         54           Enoki, Kentaro         PD13         134         QN19, RC10         34           Enoki, Masanori         PB29, PB30, PB32, QH10         124, 125, 214         Feng, Chunmu         AB04         34           Entani, Shiro         BH05, BH06, QL15         46, 47, 230         Feng, Haibo         R004, SN11         30           Entel, Peter         QH05         213         Feng, Kai         RA02         24           Eom, Jonghwa         PJ30, QK06         159, 225         Feng, Qi         RN18         29           Eom, Man Jim         SD07         331         Feng, Shiping         PB28         12	12, 354
Enoki, Kentaro         PD13         134         QN19, RC10           Enoki, Masanori         PB29, PB30, PB32, QH10         124, 125, 214         Feng, Chunmu         AB04         34           Entani, Shiro         BH05, BH06, QL15         46, 47, 230         Feng, Haibo         RO04, SN11         30           Entel, Peter         QH05         213         Feng, Kai         RA02         24           Eom, Jonghwa         PJ30, QK06         159, 225         Feng, Qi         RN18         29           Eom, Man Jim         SD07         331         Feng, Shiping         PB28         12	4, 32, 97, 239, 260
Enoki, Masanori         PB29, PB30, PB32, QH10         124, 125, 214         Feng, Chunmu         AB04         34           Entani, Shiro         BH05, BH06, QL15         46, 47, 230         Feng, Haibo         RO04, SN11         30           Entel, Peter         QH05         213         Feng, Kai         RA02         24           Eom, Jonghwa         PJ30, QK06         159, 225         Feng, Qi         RN18         29           Eom, Man Jim         SD07         331         Feng, Shiping         PB28         12	, , , , , , ,
Entani, Shiro         BH05, BH06, QL15         46, 47, 230         Feng, Haibo         RO04, SN11         30           Entel, Peter         QH05         213         Feng, Kai         RA02         24           Eom, Jonghwa         PJ30, QK06         159, 225         Feng, Qi         RN18         29           Eom, Man Jim         SD07         331         Feng, Shiping         PB28         12	4
Entel, Peter         QH05         213         Feng, Kai         RA02         24           Eom, Jonghwa         PJ30, QK06         159, 225         Feng, Qi         RN18         29           Eom, Man Jim         SD07         331         Feng, Shiping         PB28         12	01, 366
Eom, Jonghwa         PJ30, QK06         159, 225         Feng, Qi         RN18         29           Eom, Man Jim         SD07         331         Feng, Shiping         PB28         12	49
Eom, Man Jim SD07 331 Feng, Shiping PB28 12	99
	24
Eom, Man Jin RB14 256 Feng, Wuwei QO22 24	44
	26
	50, 356
	32, 246
	152
-	.68
	31
	43
	29
	98, 232
	38
	.98
	67
	32
	33
Evans, Rfl FC01 73 Ferriani, Paolo IH02 96	
	51, 52, 93, 283
	13, 94, 28 75
	75
Ewings, Russell A. HA04 83 Feyerherm, Ralf EC04 66	
	150
	03
F Figueroa, Adriana I EF03 68	
Hjaikowski, Marcin PD2 I IB	35
Fabreges, Xavier     AA03, RQ16     33, 309     Fikacek, Jan     BB04     42       Fabreges, Xavier     AA03, RQ16     33, 309     Fikacek, Jan     B04     42	
Fabris, Frederik         PB10, PB33         121, 125         Filipek, Stanislaw M.         Ql06         21	18
	00, 138, 195, 275
• • •	00, 138, 195, 275 95
Falck, Augusto PB33 125 Filippetti, Alessio PF03 13	00, 138, 195, 275 95 95

Filing av Mardiania Daviantiala	DEOO	107
Filippov, Vladimir Borisovich	PE09	137
Finazzi, Marco	PL21	167
Finger, Thomas	AA05	33
Fink- Finowicki, Jan	JJ03	106
Fiorani, Dino	RN07	297
Fiorillo, Fausto	HH01	88
Fischbacher, Johann	DE01	59
Fischer, Gerda	QL05	228
Fischer, Peter	DE05, HG02, ID01, RJ16, RJ20, SH04	60, 88, 92, 285, 343
Fisher, B	PD03	132
Fisher, lan	HA05	84
Fisk, Zachary	GC03, JC01, PD23,	77, 100, 135, 28,
	PP03, QA11, QJ11,	184, 222, 263, 268,
	RD11, RE10, SB15, SF03	322, 337
Fitzsimmons, M.	BD01	43
Fitzsimmons, M. R.	QM18	234
Fitzsimmons, Mike	KD02	109
Flachbart, K.	QC22	195
Flachbart, Karol	QC20, QI02, RE13	195, 218, 268
Fleck, Catherine L	BF03	45
Fleet, Luke	CG02, PJ17, SK03	53, 157, 350
Fleury, Eric	AJ05	40
Flouquet, Jacques	JC02, PC17, RE10, SF03	100, 129, 268, 337
Folcke, Emeric	BB03, KB02	41, 107
Follath, Rolf	CD03	51
Fontes, Magda	SE22	336
Fontes, Magda B.	SM17	363
Forcen-vazquez, Elena	Q016	243
Forostiak, Serhiy	BH03	46
Fox, A. M.	RN18	299
Foyevtsova, Kateryna	SG22	343
Fraga, G L F	PG03	142
Fraile, A	SF09	338
Fraile, Arantxa	RN14	299
Franco, N.	QM13	233
Francoual, Sonia	CI04, QA01	55, 182
Franken, Jeroen	CH05, EI06, GE02, JE01	54, 71, 78, 101
Franken, Jeroen H.	IE02	93
		195
Franz, Christian	QC17	
Fratesi, Guido	PL21	167
Freeman, Arthur Freeman, Arthur J	EC02	65
	RB26	258
Freeman, Arthur J.	RC16, SO10	261, 368
Freeman, Paul G	PI03	151
Freeman, Paul G.	PC22	130
Freemnan, A. J.	RM04	295
Friedemann, Sven	FB04, PC09, QH12, RA19, RE15, SC16, SF05	73, 128, 214, 252, 268, 328, 337
Friedenberger, Nina	RP08	304
Friemel, Gerd	JC04	100
Frings, Paul	RQ16	309
Fritsch, Veronika	RE12	268
Frontera, Carlos	EC04, RA17	66, 252
Frost, Christopher D	IC05	92
Fruchart, Olivier	SH03, SH06	343, 344
Fu, Bin Hao	RD12	264
Fu, Hua-hua	QK18	227
Fu, Tsu-yi	IH06	96
Fuente, C De La	SF09	338

## AUTHOR INDEX

Fujieda, S Fujii, Yutaka Fujimori, Atsushi Fujimori, Hiroaki Fujimori, Shin-ichi Fujimoto, Kohdai Fujimoto, Satoshi Fujisawa, M. Fujisawa, Masashi Fujishiro, Hiroyuki Fujita, A Fujita, Asaya Fujita, Hirohito Fujita, Kazuhiro Fujita, Masaki Fujita, Takahito Fujiwara, Kosuke Fujiwara, Makoto Fujiwara, N. Fujiwara, Naoki Fujiwara, Tetsuya Fukada, Yukimasa Fukami, S. Fukami, Shunsuke Fukamichi, K Fukatani, Naoto Fukazawa, H. Fukazawa, Hideto Fukui, Takahiro Fukuma, Michinori Fukuma, Yasuhiro Fukunaga, Hirotoshi Fukushi, Syousuke Fukushi, Takanori Fukushima, A. Fukushima, Akio Fukushima, Tetsuya Fukuta, R. Fukuta, Ryuichiro Fukuyado, Junichi Fullerton, Eric E Fuminori, Honda Funaki, Nakaba Furdyna, J. K. Furomoto, Yoshitoki Furukawa, Kohsuke Furukawa, Nobuo Furukawa, S. Furukawa, Satoshi Furukawa, Yuji Furukawa, Yuta

Fujibayashi, David E.

PC26 131 RR08 312 QE16, QE21, QF09, 202, 205, 206, 281 QF17, RI15 56 DA02 93 ID04 RD25, SE23 266, 337 RP05 304 RC01, RD05, SD09 259, 262, 331 KI05 113 QE10, QE16 201, 202 342 SG18 312 RR08 SO24, SO26 370, 371 165, 353 PL12, SK18 CA02 48 BA04, PB29, PB30, PB32, 41, 124, 125, 276 RH12 QH02 213 351 SK08 SK13 352 QB23 190 QB16 189 QI18, SC24, SG09 220, 329, 341 RA18 252 RL07, SH07 291, 344 RL10 291 RR08 312 PJ12, PL12, RC18, SA07, 156, 165, 261, 317, SK18 353 191 QB28 187, 189, 198, 201 QB02, QB15, QD12, QE12 SD09 331 PB06, SG16 121, 342 HD04, HE05, QM10, 85, 86, 233, 287 RK11 104, 168, 301 JH01, PM08, PM09, RO03 JI01 105 366 SN14 CE03, CE06, ID02, RJ07 51, 52, 93, 283 BI01, HF01, IE04, KF01, 47, 87, 94, 110, 163 PK20 KF04 111 CF04 53 PH07, PI07 148, 151 SB24 324 PK09, QM01 161, 231 RD19 265 PN24 176 PN08, PN09, QL08, 173, 229, 232 QM04 QP09, QP10, QP11 246, 247 KC06 109 GB01, PA07 76, 117 RA22 253 DA04 56 QB15, QE12 189, 201 270 RF09

Eurucaka Michiro	RDO 1	202
Furusaka, Michiro Fuse, Takahiro	RP01 PF11	303 141
Fuseya, Yuki	KJ04	141
Fusil, Stephane	AE03	36
Futamoto, Masaaki	ED01, PM18, SM15	66, 170, 363
	2001, 11110, 31113	00, 170, 505
G		
Gabani, S.	QC22	195
Gabani, Slavomir	QC20, QD12, QI02, RE13	195, 198, 218, 268
Gaczynski, Piotr	RE03	266
Gahlot, Ajay Pratap Singh	SC02	325
Gaita-arino, A.	BC01	42
Galanakis, lossif	RP15	305
Galceran, Regina	IH01, SO18	95, 369
Gallagher, Andrew	KB04	108
Gallagher, Bryan L	SH09	344
Gallais, Yann	PD17	134
Galyas, Anatoly	QG05	209
Gambarelli, Serge	AI04	39
Gandhi, Ashish Chhaganlal	QP07	246
Gang, Qiu Xiang	SB01	320
Ganguly, Arnab	SH12	345
Gan'shina, Elena A	PN22	175
Ganz, Philipp R.	QL05	228
Gao, Bo	SO25	370
Gao, Guoying	PN10, RM03, SO01	173, 295, 366
Gao, Kaige	RA14	251
Gao, Liangqiu	RL05	290
Gao, Lu	AH04	38
Gao, Tie Ren	BD05	43
Garbarino, Gaston	JJ04, RB20	106, 257
Garcia - Garcia, A. Garcia - Munoz, Jose Luis	KC05, PJ23 PA13	109, 158 118
Garcia, Carlos	PJ25	158
Garcia, Fernando A.	PI19	158
Garcia, Flavio	RJ12, RJ13	284
Garcia, Jackeline Collave	SE22	336
Garcia, Javier	DG03	61
Garcia, Luis	EF03	68
Garcia, Luis Miguel	GG03, GG05, QI14	80, 220
Garcia, Noel A.	QB18	190
Garcia, Vincent	AE03	36
Garcia-flores, Ali F.	PI19	153
Garcia-garcia, Javier	PI18	153
Garcia-munoz, Jose Luis	EC04, RA17	66, 252
Garcia-saiz, A.	GH05	81
Garitaonandia, Jose Javier	QH16, RN13	215, 298
Garnweitner, Georg	JG06	104
Garst, Markus	DF02, RI12	60, 280
Gasche, T. P.	QM13	233
Gastelois, Pedro	EI01	70
Gataullin, Eduard	PE08	137
Gaudin, Gilles	SH05	344
Gaulin, Bruce D	JD01	100
Gaur, N.k.	QA14, RF04, RQ06	184, 269, 308
Gautam, Sanjeev	PG27, QO27	146, 245
Gavilano, Jorge	PO17, QF20	180, 207
Gavilano, Jorge L. Gavrichkov, Vladimir	CD04 PD01	51 132
Gavrilkin, S.	QC22	195
Gaviinnii, J.	QU22	

Gavrilkin, Sergey Yu.	QC19	195
Gawryluk, Dariusz J.	QB30	192
Gazo, Emil	QI02, RE13	218, 268
Ge, H.I.	RO14	303
Ge, Hongliang	RO13, RO15	302, 303
Geerts, Y. H.	EH02	70
Gegenwart, P	DC03	58
Gegenwart, Philipp	FB03	73
Gegenwart, Phlipp	SG13	341
Gehring, G. A.	RN18	299
Geibel, C.	QB14	189
Geibel, Christoph	AB02, CD03, FB04, JC01, KA03, PD12, QC09, QC18, QD07	34, 51, 73, 100, 107, 133, 193, 195, 197
Geim, Andre	PP07	29
Gemming, Sibylle	RJ11	284
Geng, Y. X.	PM03	167
Genossar, J	PD03	132
George, Jean - Marie	HE04	86
George, Jean-marie	Al04, CG03	39, 53
Georges, Antoine	AB04	34
Georgii, Robert	DF02	60
Gerasimov, Evgeny G.	Q108	219
Gerber, Alexander	DE02, RK01	59, 286
Gerber, Simon	CD04	51
Gercsi, Zsolt	RF15	271
Gerhardt, Theo	SH04	343
Geurts, Christian	JE01	101
Ghaffari, Muhammad	PN03	172
Ghiringhelli, Giacomo	BA03	41
Ghoshal, Sayak	PJ28	159
Ghoshray, A	QM24	235
Ghoshray, Amitabha	RB15	256
Ghoshray, K	QM24	235
Ghoshray, Kajal	RB15	256
Giamarchi, Thierry	CC03	50
Giannopoulos, George	DD06, QP12, QP13	59, 247
Giblin, S. R.	PI03	151
Giebels, F	RP20	306
Gillett, J.	GA01	75
Gillon, Beatrice	GF04	79
Gingras, M. J. P.	GF01	79
Ginting, Dianta	RA01	249
Giouroudi, Ioanna	SA02	316
Giovannini, Loris	GG02	80
Giovannini, Mauro	IA02, QD08, RD21	90, 197, 265
Giri, S	QL14, SO14	230, 369
Giri, S.	PL07	165
Giri, Saurav	RH02	275
Givord, D.	JH02	104
Givord, Dominique	EJ03	72
Gjoka, Margarit	BJ02, RO08, RQ10	47, 302, 308
Gkanas, Evagelos	SN15	366
Gladczuk, Leszek	QM17	234
Glavatskyy, I.n.	QG27	212
Gloos, Kurt	SB16, SK07	322, 351
Gloskovskii, A.	SK24	354
Gloter, Alexandre	AE03	36
Glushkov, V. V.	PE12, QC22, RH07	138, 195, 275
Glushkov, Vladimir	ED06, KI05, QC20	67, 113, 195, 340
	2000, 1100, 2020	5,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,

100, 107,

Glushkov, Vladimir V.	SG07, SG08	105
Giusrikov, vladimir v. Gnezdilov, Vladimir	QC19	195 203
Gnida, Daniel	QE27 AB03, SF12	203 34, 339
Go, Naoto	PK12	162
Godinho, M.	QM13	233
Godlewski, Marek	DD02	58
Goetzke Macedo, Daniela	PB10	121
Goff, Jon	DH03	62
Gofryk, K.	FB01	73
Gofryk, Krzysztof	SE19	336
Goq, Thomas	DC03	58
Goh, Jing Qiang	S109	350
Goh, Swee K	FB04	73
Goh, Swee K.	PC09, RQ09	128, 308
Goh, Swee Kuan	PC02	127
Goho, Takeshi	RG11	273
Goikhman, Alexander	DD04, SK26	59, 354
Goikolea, Eider	RN13	298
Gokhfeld, Denis	SB09	321
Goko, T.	GF01	79
Gollisch, H	RP20	306
Goltz, Till	AB02	34
Golubov, A. A.	GA03	75
Gomes, Angelo M	0009	242
Gomez Sal, J. C.	GH05, QD10	81, 198
Gomez, Maria Elena	SA12	318
Goncalves, Antonio	SG03	340
Gong, Junfeng	RM08	296
Gonzalez, Jorge L.	SM17	363
Gonzalez, Julian	CH01, CJ04	54, 56
Gonzalez, Julianmaria	DJ03	63
Gonzalez, Lorena	DJ03	63
Goo, Cheol Soo	RR04	311
Goodilin, Evgenii	ED06, KI05	67, 113
Gopalakrishnan, C	PO24, SA21, SA22,	181, 319, 320, 35
	SA23, SK32	
Goraus, Jerzy	PD21	135
Gorbunov, D. I.	Q103, Q109	218, 219
Gorbunov, Denis	SG03	340
Goremychkin, E A	SE13	335
Goremychkin, E. A.	RA22	253
Goremyhkin, E A	SF09	338
Gorria, Pedro	JJ04	106
Goryunov, Yuriy Vladimirovich	IG02, PE09, QG23	95, 137, 212
Goswami, Partha	PB01, RE02	120, 266
Goswami, Partho	SC02	325
Gotberg, Olof	HA03, SE08	83, 334
Goto, Junpei	PL06	164
Goto, Takayuki	RH03, RH11	275, 276
Goto, Terutaka	RB05	254
Gotou, Junpei	RC02	259
Gottlieb-schoenmeyer, Saskia	QH24	216
Gotze, Kathrin	QD07	197
Gouchi, Jun	QC03	192
Goumri-said, Souraya	QL17, SK12	230, 352
Gouvea, Cristol De Paiva	PB11	121
Goya, Tomoki	RB06	255
Goyal, Neeraj	QJ17	223

## AUTHOR INDEX

Gradhand, Martin Graf, Matthias Grams, Christoph Granado, Eduardo Granovsky, Alexander Granroth, Garrett E Granroth, Garrett E. Grbic, Michael Greene, L. H. Greene, Laura H. Gretarsson, Hlynur Greven, Martin Grigera, Santiago Grigoriev, Sergei Grigoriev, Sergey Grigorieva, Anastasiya Grigoryeva, Natalia Grigoryeva, Natalya Grimaldi, E. Grin, Yuri Grioni, Marco Grisewood, N. Grishin, Alexander Griveau, Jean-christophe Groessinger, Roland Grois, Andreas Grollier, J. Grollier, Julie Grosche, F Malte Grosche, F. Malte Grosche, Malte Gruber, Manuel Grundler, Dirk Grunin, Alexey Grytsyuk, Sergiy Gschneidner, Jr., Karl Gu, Bo Gu, G. Gu, Genda Gu, Han Guan, Wen Guari, Yannick Gubbiotti, Gianluca Gubkin, Andrey Gubkin, Andrey F. Gudim, I Gudim, Irina Gudim, Irina A Gudim, Irina A. Guduru, Veerendra Guduru, Veerendra Kumar Guenther, Marco Guerrero, Sebastian Guetari, Rim Guidi, T. Guillot, Maurice Guimaraes, Alberto Passos Guimaraes, Marcos H. D. Gukasov, Arsen Guleria, Anupam

EE01 KG01 QA22 PI19 CH01, SO04, SO27 PI07 RD14 IC02 GA01 RD09 DC03 CA03 EA04 QM08 DF04, SL27 ED06, KI05 SL27 QM08 CE06, ID02, RJ07 RD12 SE08 PJ26 CJ01 RE03, SE19 SA02 BE03 CE03, CE06, ID02, RJ07 HF01, ID03, IE04 FB04, QH12, RA19, SF05 73, 214, 252, 337 RE14, RE15 PC02, PC09 GI03 GG02 DD04 QP01, QP02 JJ01 JF03 GA01 HA03 RR29 SO25 RN15 GG02 QG04 QN10 QJ16 IJ01, RA23 CI04 PA02 DH02 SA10 PI05 QF20 PM06 RA22 QJ15 RJ13 GI05 SE16 QB14

Gumeniuk, Roman	RD12	264	Ham,
Gunduz Akdogan, Nilay	EJ04	72	Hama
Guo, Guan Yu	JI05	105	Hama
Guo, Hong	QK02	224	Hama
Guo, Huaiming	SD18	333	Hama
Guo, Lu Kai	RD12	264	Hama
Guo, Shuang	RK16	288	Hama
Guo, Zhaohui	PK18, RO04, RO07	163, 301	Hamr
Gurevich, Alexander	JA02	98	Han,
Guritanu, Violeta	QD07	197	
Guryeva, Tatyana	QM19	234	Han,
Gusakova, Daria	SI09	347	Han,
Gusev, Sergey	SG17	342	Han,
Gusin, Pawel	PG26	146	Han,
Gutersloh, Daniel	PE07	137	Han,
Gutfleisch, Oliver	DJ04, KH01, SO19	63, 112, 369	Han,
Guth, Konrad	KH01	112	Han,
Gutmann, Matthias	SD13	332	Han,
Gutowska, M. U.	PH06, PM05	147, 168	Han,
Guziewicz, Ela	DD02	58	Han,
Gvasaliya, Severian	CC03, IJ06	50, 98	- 1
Gwon, M. J.	PO23	181	Han,
Gwon, Minji	RK15	288	Han,
Gwon, Yoonjung	PN08, PN09, QM04	173, 232	Han,
Gyawali, Parshu R.	PJ01	154	Han,
Gyuwan, Farsha K.	1501	191	Han,
Н			Han,
Haam, So Young	JB02	99	Han,
Haase, Juergen	RQ09	308	Hana
Haazen, Pascal P.	IE02	93	Hand
Haberkorn, N.	QB31, SB12	192, 322	Hand
Habermeier, Hu.	ED05	67	Hand
Habermeier, Hanns-ulrich	RB29	258	Hand
Hachiuma, Suguru	QA05	183	Hane
Hackl, Rudi	CA03	49	Hanfl
Haddon, Robert C.	PN16	174	Hanh
Hadiyawarman, H	QL09	229	Hans
-			
Hadjipanayis, George	BJ02, CB05, DD06, EJ04, KB04, RO08	47, 49, 59, 72, 108, 302	Hany Hao,
Hadjipanayis, George C.	GJ01	82	Нарр
Haenni, Nora	QG14	210	Hara,
Haga, Yoshinori	GC03, HP11, IG04,	77, 30, 95, 131, 133,	Hara,
naga, toshinon	PC28, PD09, PD13,	134, 135, 192, 262,	Harad
	PD23, QC03, RD01,	263, 266, 268, 322,	Harad
	RD11, RD25, RE10,	337	
	SB15, SE23, SF03		Harad
	0.642	404	Harad
Hagemann, Hans	QC13	194	Harad
Hagihala, Masato	QE18	202	Harag
Hagiwara, M.	CF01	52	Haral
Hagiwara, Makoto	PB25	124	Haral
Hagiwara, Masayuki	PD13, PI06, QB09, QB24,	134, 151, 188, 191,	Hardi
	QF06, QH02, RH25	204, 213, 278	Hardy
Hagymasi, Imre	PD15	134	Hardy
Hahn, Byung-dong	DJ02	63	Hari,
Hajiri, Tetsuya	AD04, QD18	35, 199	Harin
Hakimi, A. M. H. R.	RN18	299	
Halbedel, Bernd	Q002	241	Harna
Hall, Anders	QH11	214	Harris
Hall, Michael	HH02, KH03	88, 112	Harris
Halue, K.	PE17	138	Harris

Ham, Chan	PO13	179
Hamada, Rika	QH28	217
Hamada, Tomoyuki	RF07	270
Hamamoto, Kenta	SA03	316
Hamano, Suguru	IG05	95
Hamasaki, Tatsuichi	RH19	277
Hamaya, Kohei	HE01, RK10	86, 287
Hamrle, Jaroslav	CH03, QH19	54, 216
Han, Dong-soo	DE05, RJ16, RJ17, RJ18, SH14	60, 285, 345
Han, Guchang	QP06	246
Han, Jeong-ho	RR14, SN04	313, 364
Han, Jinhee	RC15	261
Han, Jung Hoon	DF01	60
Han, Jung-ho	SN02	364
Han, Jungmin	RL14	292
Han, S. W.	PG11	144
Han, Sang Wook	PB12	122
Han, Seungkyu	RA26	253
Han, Suk Hee	QK11, QK12, QK13, RL28	225, 226, 294
Han, Sukhee	QK10	225
Han, Suk-hee	PJ30	159
Han, Tianheng	QE27	203
Han, Xiufeng	QK02, SA19, SA20	224, 319
Han, Xiu-feng	RM01	295
Han, Yibo	QF06	204
Han, Young Ho	BJ04, QN13	48, 238
Hanazawa, Atsufumi	RI17	281
Handayani, Ismudiati Puri	PH17	149
Handayani, Puri I.	PA24	119
Handoko, Djati	RP03, SH10	303, 344
Handoko, Erfan	RA15	251
Haney, Paul M.	EE02	67
Hanfland, Michael	RB20	257
Hanh, Duong Thi	HC02, PE15	84, 138
Hanson, Maj	SL15	358
Hanyu, Takahiro	BI02	47
Hao, Lin	PK03	160
Happo, Naohisa	SK13	352
Hara, Masahiro	KC06	109
Hara, Shigeo	QE17, QF09	202, 205
Harada, Eiichiro	RD06	263
Harada, Kazunori	KF05	111
Harada, Keitaro	SO515	369
Harada, Ken	AF03, PG15	37, 144
Harada, Masashi	PO17	180
Haraguchi, Shinya	AE02, PL06, RC02, RL02	36, 164, 259, 290
Haraldsen, Jason	KG01	111
Haraldsen, Jason T	EB04	65
Harding, Phimphaka	SL10	357
Hardy, Vincent	JB05	100
Hardy, Walter	BA01	40
Hari, Murali Krishnan	Q024	244
Harima, Hisatomo	HP11, RD16, RG11, SB05, SG14	30, 264, 273, 321, 341
Harnagea, L.	RB17	256
Harrison, N.	SC03	326
Harrison, Neil	BA01, EA03, PB02, QB13	40, 64, 120, 189
Harrison, Nicholas	RF15	271

Harumori, Kohei	SK23	354
Haruna, Daiki	QG10	209
Haryono, Suprijadi	RC02	259
Hasan, M. Z.	SD04	330
Hasan, M. Zahid	HI02	89
Hase, Masashi	PH02, QA05, QA07, QA09, RH11	147, 183, 184, 276
Hasegawa, Kota	RL15	292
Hasegawa, Takahiro	QI17	220
Hasegawa, Takashi	SL02	356
Hasegawa, Takumi	PC05, PG16, RD10, SG15	127, 145, 263, 342
Hasegawa, Takuya	QA05, QA07, QA09	183, 184
Hashemi, S. E.	PO29	182
Hashemi, Sayed Ebrahim	PO27	181
Hashi, Shuichiro	PO20, SM08, SN14	180, 362, 366
Hashim, Mansor	BJ06, RQ22	48, 310
Hashimoto, Susumu	CE05	52
Hashimoto, Takuya	BG04	46
Hashmi, Arqum	RM07	296
Hassinger, Elena	JC02	100
Hasuo, Tadahiko	SG09	341
Hatakeyama, Kenichi	QN02	236
Hatakeyama, Kodai	RP06	304
Hatayama, Nobukuni	PF08	140
Hatayama, Yuki	RO09	302
Hatch, R.	1102	97
Hattori, Taisuke	RD05, RD08	262, 263
Hauguel, Tony	SA24	320
Havela, Ladislav	QI19, RD26, RE09, SG03	220, 266, 267, 340
Haverkort, Maurits Wim	QC09	193
Haw, S. C.	QB21, SM11	190, 362
Hayakawa, J.	IF01	94
Hayashi, Hisashi	QD17	199
Hayashi, Junichi	PD10	133
Hayashi, Kouichi	SK13	352
Hayashi, Masamitsu	IF05	94
Hayashi, Misaki	RP07	304
Hayashida, Minami	QC18	195
Hayato, Miyagawa	PN24	176
Hayden, Stephen	RE14	268
Hayn, Roland	CD03	51
He, Chenchong	AJ03, JI06	40, 105
He, Chunyong	DH01	62
He, Junfeng	SB31	325
He, Long	SB13	322
He, Pan	JI05	105
He, Zhangzhen	PA18, QF24	118, 207
Heald, S. M.	RN18	299
Hedo, Masato	PD14, QC11, QI12, QI13	134, 194, 219
Heger, Gernot	RP22	306
Hehn, Michel	KF02, SH02	110, 343
Heiliger, Christian	DI03	63
Heinze, Stefan	AF02, IH02, IH03, PG12	36, 96, 144
Heiss, Dominik	IE03	93
Hemberger, Joachim	QA22	186
Hemmi, K.	CF04	53
Hemmi, Kazuhiro	PH07	148
Hendry, Euan	GD04, RK20	78, 289
Henkie, Zygmunt	PE19, SB26	139, 324

Henriques, Margarida Her, Eun Ju Hering, Eduardo Hering, Eduardo N Heringa, Jouke R Hernandez, Yenny Hernandez-velasco, Jorge Hernando, Blanca Herng, Tun Seng Herrero Albillos, Julia Herrero-albillos, Julia Herrero-martin, Javier Herringer, S N Herrmannsdoerfer, Thomas Herschbach, Christian Hervieu, Maryvonne Herynek, Vit Herzer, Giselher Herzog, Gabriela Hewson, Alex Heyderman, Laura Heyderman, Laura J. Hicks, Trevor Hidaka, Hiroyuki Hien, Nguyen Thi Minh Hierro Rodriguez, Aurelio Hiess, Arno Higasginaka, Ryuji Higashinaka, Ryuji Hikihara, Toshiya Hild, K Hilgenkamp, Hans Hill, John Hill, R. W. Hillebrands, Burkard Hiller, Adrian Hillier, A Hillier, A D Hillier, A. D. Hillier, Adrian Hillion, Arnaud Hilscher, Gerfried Hindmarch, Aidan T Hinks, D G Hinzke, Denise Hirai, Daigorou Hirai, Daishi Hiraka, Haruhiro Hirakawa, Sentaro Hiramatsu, Ryo Hirano, Masahiro Hirano, Masanori Hiraoka, N. Hiraoka, Nozomu Hirashima, Dai Hirata, Wataru Hirata, Yasuyuki Hiratsuka, Yusuke Hirayama, Shigeyuki Hirayama, T.

SG03 340 CB01 49 SE22 336 SE15 335 281 RI18 226 OK14 PI18 153 63 DJ03 58 DD03 GG03, RP04 80, 303 EF03, GG05, QI14, RP08 68, 80, 220, 304 EC04, PA15, RA17 66, 118, 252 42 BC05 131, 331 PC27, SD07 67 EE01 RA11 251 46 BH03 PM04 168 HF03, RP16 87, 305 SF15 339 HP62 31 RK17 288 IJ03 97 QC25, SE07 196, 334 SA16 318 GG03 80 AA05, RB07 33, 255 QE13 201 RD17 264 101 JD04 233 QM12 DH02, SA10 62, 317 58 DC03 84 HA05 85, 348 HD01, SI15 128 PC12 338 SF09 333, 334, 335 SE03, SE05, SE13 107 KA02 99, 325 JA05, SC01 237, 286 QN06, RK03 SC01 325 SH09 344 SC05, SE14 326, 335 FA02, FC03, RJ04 72, 74, 283 RB04 254 QD06 197 QG18, RP01 211, 303 QI13 219 RL07 291 QC08 193 QB15 189 QB21 190 QC09 193 PF19 142 QB24 191 Q013 242 DE04 60 SK31 355 RD15 264

Hirohata, A.

Hiroi, Zenji

Hirono, Toko

Hirose, Yusuke

Hirsch, Konstantin

Hirschberger, Max

Hirtenlechner, Eva

Hisada, Akihiko

Hisamatsu, Toru

Hisayoshi, Keiji

Ho, Huei-ying

Hjorvarsson, Bjorgvin

Hjorvarsson Bjorgvin

Hoa Hong, Nguyen

Hoang, Nam-nhat

Hodgson, Matthew

Hoeijmakers, Mark

Hoffmann, Germar

Hofmann, Michael

Hofmann, Philip

Hoglin, Viktor

Holbein, Simon

Holder, Matthias

Holler, Robert

Hodovanets, H.

Hirohata, Atsufumi

Hiroi, Masahiko

Hollmann, Nils
Holmes, Stuart N
Honda, Fuminori
Honda, Hiroyuki
Honda, S.
Honda, Syuta
Honda, Zentaro
Hong, Cheng Hai
Hong, Do-kwan
Hong, Hongpei
Hong, Jeongmin
Hong, Jin Pyo
Hong, Jinki
Hong, Jisang
Hong, Jongill
Hong, K P.
Hong, Kwang Pyo
Hong, S. C.
Hong, Seung Hwan
Hong, Soon Cheol
Hong, Sung-hak
Hong, Tae Min
Hong, Tao
Hong, Yang-ki
Honma, Yuki
Hono, K
Hono, Kazuhiro
Honolka, Jan

londa, Zentaro	QF06, RH25
long, Cheng Hai	QN03
long, Do-kwan	PM25, RR31
long, Hongpei	SO01
long, Jeongmin	PN16
long, Jin Pyo	RK19
long, Jinki	PJ31, PK21, QK10
long, Jisang	JH03, RM07
long, Jongill	JI02, RL30
long, K P.	PH03
long, Kwang Pyo	QN09
long, S. C.	PG11
long, Seung Hwan	SB32
long, Soon Cheol	KI01, PB12
long, Sung-hak	PB21, RB27
long, Tae Min	QN22
long, Tao	BC05
long, Yang-ki	JH04, KH04
lonma, Yuki	PH05
łono, K	BG01
lono, Kazuhiro	IF03, SK31
Ionolka, Jan	AD03
The state of the s	
The 19 <sup>th</sup> International Conference	ence on Magnetism

SK03	350	Hor, Y. S.
PJ17, SK11	157, 352	Hord, Roland
QG13, QG15, QH08,	210, 214, 239, 354	Hori, Akihiro
QN20, SK23		Hori, Shiori
JA04, PC17, PG16,	99, 129, 145, 199,	Hori, Yusuke
QE03, QE05, QE08,	200, 201, 210, 321,	Horigane, K.
QE13, QG16, SB05, SG04, SG14	340, 341	Horihata, A.
		Horng, Herng-er
RP07	304	Horng, Lance
PB27, PD08, PD13, RD19	124, 133, 134, 265	Horsch, Peter
KD01	109	Horvatic, Mladen
SF05	337	Hosaka, Tomohiro
QG14	210	Hosokawa, Shinya
SE01	333	Hosono, H.
QG13 PN11, RQ03	210 173, 307	Hosono, Hideo Hosono, Kazuhiro
AF05, RP15	37, 305	Hossain, Z.
QH11	214	Hossinzadeh, Nahid
PL03, PL13, QM09	164, 166, 233	Hotta, Chisa
SM07, SM10	362	Hotta, Koji
	251, 269, 341, 368	Hotta, Takashi
QK09	225	Hou, Yanglong
RD02	262	Howard, A.
GE02, JE01	78, 101	Hrabec, Ales
GH02	81	Hradil, K.
EF02, QG26	68, 212	Hsiang, Hsing-i
1102	97	Hsiao, S. N.
PM10	169	Hsieh, C.c.
AA05	33	Hsu, Chuang-han
CD03	51	Hsu, Jen-hwa
RG15	274	Hsu, Pin-jui
PI08, PI17	152, 153	Hsu, Wei-hung
PN23	175	Hsu, Y. Y.
HP11, PB27, PD08,	30, 124, 133, 134	Hsueh, Wang-jung
PD13, PD16		Hu, Chong Der
PG22	146	Hu, Fengxia
CG02	53	Hu, Guanghui Hu, Hai Ning
RJ08, SK20	283, 353	Hu, Jin
QF06, RH25	204, 278 236	Hu, R.
QN03 PM25, RR31	171, 316	Hu, Shaojie
SO01	366	Hu, Xinghao
PN16	174	Hu, Zhiwei
RK19	289	Huang, Chao-hsien
PJ31, PK21, QK10	159, 163, 225	Huang, Chia-sheng
JH03, RM07	104, 296	Huang, Chiuan-fa
JI02, RL30	105, 294	Huang, Chonghui
PH03	147	Huang, J. H.
QN09	237	Huang, Jin-hua
PG11	144	Huang, K.w.
SB32	325	Huang, Ssu-yen
KI01, PB12	112, 122	Huang, Wen-min
PB21, RB27	123, 258	Huang, Xiao-lan
QN22	240	Huang, Xuelian
BC05	42	Huber, Rupert
JH04, KH04	104, 112	Hudl, Matthias
PH05	147 45	Hueso, Luis E. Hugli, Remo Viktor
BG01	45 94 355	Hugii, kemo viktor Huh, Mooyoung
IF03, SK31 AD03	94, 355 35	Huhtinen, Hannu
		. ananany norma

SD04	330
PI05	151
HC03, QF25	85, 208
QM16	234
PD10	133
QE15	201
CG02	53
AH03, GH01, QM27	38, 80, 236
PC30, SB19, SH08, SL18	131, 323, 344, 359
RG18	274
IC02	92
	92 183, 184
QA05, QA07, QA09	
SK13	352
QB23	190
PC23, QB16, QC08	130, 189, 193
PJ18	157
QB14	189
PB15	122
BC05	42
RM05	296
PF11	141
SL12	358
QB12	189
SH05	344
PI03	151
BJ01, QN24, QN25	47, 240
PK06, QP17	161, 248
SN06	365
AD01	35
PK06, QP17	161, 248
AD01	35
QN25	240
SB04	320
AD01	35
QF26	208
SO25	370
QJ06	221
SL16	358
RB12	256
QC23	196
RL20	293
PO19, SN13	180, 366
PI08, PI17	152, 153
SH08, SL18	344, 359
RR26	315
PK06, QP17	161, 248
PG21	145
PM03	167
PC14	129
AH03	38
PC14	129
RE11, SB25	268, 324
IH06	96
SM14	363
HA02	83
EI03	71
HE02	86
GF02	79
PM23	171
PB20	123
1020	U21

Huijben, Mark	DH02, SA10	62, 317
Hung, C. H.	SB04	320
Hung, Dung S	RK07	287
Hunt, C. R.	GA01	75
Hur, Nam Jung	SE02	333
Hur, Namjung	PI10, QA24, QM23	152, 186, 235
Husek, I	RB08	255
Hutanu, Vladimir	PA14, RP22	118, 306
Huxley, A. D.	GC01	77
Huxley, Andrew	GC04	77
Huxley, Andrew D	SE12	335
Huynh, Chinh Dang	RN19	299
Hwang, Chan Yong	PJ14	157
Hwang, Chanyong	BD03, IH04, IH05, QM26, RK18, RK19,	43, 96, 235, 288, 289, 307, 345
	RP23, SH13	
Hwang, Do Guwn	PO09, PO10	178
Hwang, Do Gwen	PO06	178
Hwang, Harold	BA02	41
Hwang, J. S.	SL08	357
Hwang, Jihoon	PB12, PG11, PH14, QA04	122, 144, 149, 183
Hwang, Jungmin	PA21	119
Hwang, Seong-min	RQ19, RQ20	310
Hwang, Su-jin	RR10	312
Hyun, Seung Ill	QH25	217
Hyun, Sung Wook	PA22	119
Hyun, Young-hwan	SD03	330
1		
•	DNOO	200
Ibarra, M R	RN08	298
Ibarra, M. R.	KC05, PJ23	109, 158
Ibarra, Manuel Ricardo Ibragimova, Elvira	DG05, PG08, SM16 RR28	61, 143, 363 315
Ibrahim, I. A.	RN01	296
Ichihara, Masaki	SB14	322
Ichikawa, Akihiro	QB26	191
Ichimura, Masahiko	PJ07, RL01	155, 290
Ichiyanagi, Yuko	PO02, QL03, RN02	177, 228, 297
Ide, Tetsuto	PH15	149
Ideta, Yukiichi	QF24	207
Idigoras, Olatz	SM04	361
ldo, Masayuki	PB08	121
ldutsu, Yuichi	RH25	278
Idzikowski, Bogdan	Q017	243
Idzuchi, Hiroshi	HE05	86
leda, Jun'ichi	JE02, JF02, KF03	102, 111
lenaga, Koichiro	QL11, QL12, QL13	229, 230
lga, Fumitoshi	PC05, PE10, RH03	127, 137, 275
Igarashi, Hideki	QC25	196
Igarashi, Kazuhiro	QF15	206
Igarashi, Suguru	PB06	121
Igawa, Naoki	JJ05	106
Iglesias, Oscar	KB06, RN14	108, 299
Iglesias-silva, E.	CB02	49
Ignatchik, Oleg	QD07	197
Iguchi, Daisuke	PC17	129
Iguchi, Satoshi	GB04	76
Ihara, Yoshihiko	PC15, RD05	129, 262

## AUTHOR INDEX

lida, Hiroki likubo, Satoshi limura, S. lizuka, T lizuka, Takuya Ikeda, Hiroaki Ikeda, Masami Ikeda, Naoshi Ikeda, Shoji Ikeda, Shugo Ikeda, Yoichi lkedo, Yutaka lkehara, Yuki lkeuchi, K. Ikeuchi, Kazuhiko Ikeura, Tohru Ikuo, Nakai Ilyn, Maxim lm, Dae Yeong lm, H. J. lm, Jino lm, Mi-young Imada, Masatoshi

Imai, Yoji Imai, Yoshiki Imai, Yuya Imakyurei, Takumi Imamura, Hiorshi Imamura, Hiroshi Imamura, Masaaki Imanaga, Yukihiro Imperia, Paolo Imura, Keiichiro Inaba, Kensuke Inada, Yoshihiko Inagaki, Keima Inagaki, T. Inagaki, Yuji

Inami, Nobuhito Inami, Toshiya Inamura, Yasuhiro Indenbom, Mikhail Infante, Ingrid Ingvarsson, Snorri Innocenti, Claudia Ino, Hiromitsu Inomata, K Inomata, Koichiro Inosov, Dmytro Inoue, J. Inoue, Katsuya

Inoue, Mitsuteru Inoue, T. Inoue, Y. F. Inoue, Yukihiro F Insausti, Maite

QC10, QC11 193, 194 PB32, QH09, QH10 125, 214 QB23 190 QJ05 221 197 QD07 OC04, OC08 192, 193 OB24, OF06 191, 204 RA18 252 247, 348 QP14, SI17 IG04, RB03, RD11, RD13 95, 254, 263, 264 195, 334 QC18, SE10 180 PO17 GJ02 82 53 CF04 56, 151, 276 DA01, PI07, RH12 PD16 134 192 QC01 SO04, SO27 367, 371 QO20 244 SG11 341 RC16 261 DE05, HG02, RJ16, RJ20, 60, 88, 285, 343 SH04 FE01, QB10 74, 188 KF05 111 PF17, SD05 142, 330 PD20 135 PG09 143 RJ06 283 87, 103, 250, 348 HF02, JG03, RA09, SI16 176, 368 PN26, SO09 DE04 60 179 PO12 199 QD18 139 PF02 130, 222 PC20, QJ07 353 SK18 214 QH09 202, 206, 229, 230, QE18, QF14, QL11, 307, 341 QL12, QL13, RQ05, SG09 EI02, SI17 70, 348 IG03, QI10 95, 219 RH12 276 SA24 320 QA27 187 GH03 81 CB04 49 PG05 143 SK29 355 EH01, IF03, IF05 69, 94 JC04 100 CG02 53 PG09, PG10, PG22, 143, 144, 146, 211 QG22 QP15, QP16 247, 248 SG02 339 EF01 68 341 SG12 RN13 298

Inui, Ken	PN07	173	Isomura, Shinsaku	PK07, QP08	161, 246	Jaim, H. N		5B06	321
Inumaru, Kei	PE10	137	Isono, Takayuki	PC17	129	Jaime, Ma			100
Ioannidou, Alexandra	RQ10	308	Isozaki, Toshiyuki	QK19	227	Jain, Abhi			39, 53
lon, Lucian	QK20	227	Itai, Kazumasa	PD15	134	Jakiela, Ra		DD02	58
Ionescu, Adrian	PN23	175	Itkis, Mikhail E.	PN16	174	Jakob, Ge		CH03	54
Ipatov, Mihail	AJ02, CJ02, PJ25	40, 55, 158	Ito, Eisuke	CB01	49	Jal, Emma		AD02	35
Ipatov, Mikhail	SO27	371	lto, Hiroyoshi	RK10	287	Jalli, Jeeva		H04	104
Ipponjima, Tsukasa	SG16	342	Ito, Keigo	SD03	349	Jamali, Ma	ahdi E <sup>i</sup>	EG03, GE03, PJ20, RL19	69, 78, 158, 293
Iqbal, Muhammad Waqas	QK06	225	lto, Keita	KF05	111	Jamet, Ma			39, 53, 58, 71, 78,
Iqbal, Muhammad Zahir	QK06	225	lto, Masakazu	QG13, QG15, QN20,	210, 239, 354		G	GE01, QK16	226
Irie, Ryotaro	PH15	149		SK23		Jang, Dae	kyu R	RR17	313
Irifune, T.	SG02	339	Ito, T.	RM04	295	Jang, H.	P	PH10	148
Irifune, Tetsuo	PE10	137	Ito, Takahiro	AD04, QD18	35, 199	Jang, Hyu	k-jae A	AE04	36
Iritani, Kensuke	QC08	193	Ito, Tomonori	RM05, SO10, SO11	296, 368	Jang, Ik-sa	ing R	RR14, SN02	313, 364
Irusta, Silvia	RN09	298	Ito, Toshimitsu	QA05, QA07, QA09,	183, 184, 276	Jang, Jin F	lak R	R24	314
Irvine, Andrew C.	IE03	93		RH11		Jang, Jung	j-tak C	CB01, PO15, PO16	49, 179
Isasa, Miren	HE02	86	Ito, Wataru	PB06, QG08, RP07	121, 209, 304	Jang, Jung	jwon Q	QO10, SD11	242, 331
Isayama, Akira	Q018	243	Ito, Yuichiro	PH15	149	Jang, Moo	ongyu R	RL21	293
Ishaque, Zahid	SH03, SH06	343, 344	Ito, Yuzo	SE10	334	Jang, Pyur	ngwoo P	PM20	170
Ishibashi, Akira	SM09, SO07	362, 367	Itoh, Akiyoshi	BG03, RK08, RK17	45, 287, 288	Jang, Seol	<-myeong R	R04	311
Ishibashi, Hiroki	PI02	151	Itoh, Atsushi	PC28	131	Jang, Tae-	suk R	RO01, RO02	300, 301
Ishibashi, Shoji	GB03, RG07	76, 272	Itoh, Hiroyoshi	RJ08, SK20	283, 353	Jang, Y. S.	. R	RQ08	308
Ishibashi, Shota	KF01	110	Itoh, Masayuki	DA01, PH19, PH20,	56, 150, 187, 191,	Jang, You	ngjae P	PK21	163
Ishibashi, Takayuki	KD03, RP17, SK21	109, 306, 353		PH23, QB04, QB26,	206, 254	Jang, Z. H	. Q	QF13	206
Ishibashi, Yusuke	RR05	311		QF15, RB02		Jang, Zeeł	noon R	RN10, RP13	298, 305
lshida, K.	KJ01	113	Itoh, Mitsuru	QF24, QK19, RL25,	207, 227, 294, 364	Jangid, S.	R	RA03	249
lshida, Kenji	QC08, QH13, RD05,	193, 215, 262, 263		SM20		Janjan, S.	M. Q	QM05	232
	RD08		Itoh, Ryusuke	QA05, QA07, QA09	183, 184	Jansen, Ro	on A	AI05, DI02	39, 63
Ishida, Yuko	RI17	281	Itoh, Shinichi	PI07, RH13, RI10	151, 276, 280	Janson, O	. Q	QF22	207
Ishihara, Sumio	GB02, PI07, RH10	76, 151, 276	Itoi, Chigaku	RI11	280	Jantaratar	na, Pongsakorn Sl	5L10	357
Ishii, Akira	PN13	174	Itou, Masayoshi	BA04, PG05, PG13,	41, 143, 144, 162,	Janutka, A	Andrzej S	SH15	345
Ishii, Fumiyuki	RC09, RC19	260, 262		PK12, SK04	351	Jaramillo,	Juan Fernando SI	5M04	361
lshii, H.	QB21	190	Itou, Tetsuaki	BC04, QE22	42, 203	Jardim, Da	aniela Regina P	800	178
Ishii, Rikako	QF21	207	Ivanov, Alexandre	JC04	100	Jascur, Mie	chal R	R109	280
Ishii, Yasuyuki	QB02	187	Ivanov, Vsevolod	SG07	340	Jastrabik,	Lubomir A	AH02	38
Ishii, Yui	SG14	341	Ivanov, Vsevolod Yu.	QC19	195	Javanpara	st, B. G	GF01	79
Ishijima, Hiroyuki	PN07	173	Ivanshin, Vladimir	PE08	137	Javorsky, F	'avel C	QI01, SB08	217, 321
Ishikawa, Hajime	QE05, QE08	200	Iversen, B.	1102	97	Jay, Jean-p	ohilippe S	5A24	320
Ishikawa, Jun J	PG14	144	Iwakawa, Ken	PD13, QH03	134, 213	Jean, Mali	.ck D	DG02, PJ22	61, 158
Ishikawa, Takuya	RB06	255	Iwama, Hiroki	SN03	364	Jendelova	, Pavla B	3H03	46
Ishikawa, Yosuke	RI01	279	Iwamoto, Takashi	RN03	297	Jeng, Ht.	. S'	5D04	330
Ishikawa, Yusuke	PD14	134	Iwasa, A.	Q103	218	Jeon, Byeo	ong Jo Jf	B02	99
Ishikawa, Yuya	QF09, QF17	205, 206	Iwasa, K.	CF04, PL14	53, 166	Jeon, Byur	ng-gu JE	B02	99
Ishimatsu, Naoki	PN20	175	Iwasa, Kazuaki	PI07, RD14	151, 264	Jeon, Eun-	-kyoung S	SD12, SD14	332
Ishio, Shunji	SL02	356	Iwasa, Yoshihiro	AE05	36	Jeon, Gi-v	van P	N25, SK28	176, 355
Ishiwata, N.	RL07	291	Iwasaki, Hitoshi	CE05	52	Jeon, Gun	Sang R	RA22, RC08, RE07	253, 260, 267
Ishiwata, Nobuyuki	RL10	291	Iwasawa, Kotaro	PE11	137	Jeon, Kun	-rok R	RL03, SM01	290, 361
Ishiwata, S.	RA05	250	Iwase, Hiroaki	PD20	135	Jeon, Kyu	ng-won R	RR21, RR22	314
Ishiwata, Shintaro	PA12	117	Iwata, Satoshi	RL08	291	Jeon, Seur	ngmok E	EH05	70
Ishiyama, Kazushi	PO20, RP11, SM08,	180, 305, 362, 366	lyama, Ayato	QA12	184	Jeon, Sun	g Jae P	PM28	172
	SN14		lyo, A.	QB28	191	Jeong, Ho	gyun R	RC13, RC17	261
Ishizuka, Hiroaki	RG03	272	lyo, Akira	DA04, PB05, PB26,	56, 120, 124, 187,	Jeong, Ilgy	yo P	PO19	180
Isikawa, Yosikazu	RD14, SG05	264, 340		QB02, QB15, QB29	189, 191	Jeong, In-I	bum IE	B03	91
Islam, A. T. M. Nazmul	PA23, RH24	119, 278	Izawa, Koichi	PC17, PC28, PD16	129, 131, 134	Jeong, Jae	hong R	RA22	253
Ismail, Ismayadi	BJ06, RQ22	48, 310				Jeong, Jin	-	HJ05	90
Isnard, O.	QG07, QI03	209, 218	J			Jeong, Se		PH22, PN20, PN21, QL18	150, 175, 231
Isnard, Olivier	EJ03	72	Jablonski, A.	DD02	58	Jeong, Su		RR27	315
Isobe, Masaaki	PC21, RH14, SF02	130, 277, 337	Jacko, Anthony	RG14	274	Jeong, Tae	-		244, 343
Isobe, Masahiko	QD03, RH16, SB14, SF04	196, 277, 322, 337	Jaffres, Henri	AI04, CG03, HE04	39, 53, 86	Jeong, Tae			313, 364
Isoda, Makoto	QE11	201	Jaglicic, Zvonko	SL11	357	Jeong, Y. I			171
						5.			

## AUTHOR INDEX

Jeong, Yeon-ho Jeong, Yoon Hee Jeong, Yoonhee Jeong, Yoon-hee Jeong, Yun-ho Jesche, Anton Jeschke, Harald Jeschke, Harald O. Jhang, Hau-chun Ji, Gaofeng Ji, Hyo Seok Ji, Lina Ji, Qing Ji, Sungdae Jia, Q. Jia, Quanxi Jian, Ming-hung Jiang, Zhaozhen Jiao, L. Jiao, Lin Jiao, Wenhe Jin, Hosub Jin, Jung-il Jin, X. R. Jin, Xiaofeng Jingwei, Cui Jo, Euna Jo, Jeong Hong Jo, Nahyun Jo, W. Jo, William Jo, Woo Seong Jo, Younghun Jo, Young-hun Jo, Younjung Johansson, Borje Johnston, David C Johrendt, Dirk Jomura, Mitsuhiro Jonietz, Florian Jono, Youhei J onsson, Petra Joo, Dae-suk Joo, Sungjung Joo, Younghoon Joseyphus, Raphael Justin Jourdan, M Jourdan, Martin Joy, P A Joy, P. A. Ju, Hye Sun Ju, Sol Jue, Emilie Julian, Stephen Jun, Akedo Jung, Chang Uk Jung, Dahee Jung, Heejong Jung, Hyunsung

Jung, Jin-seung

PM25 171 PA11, PN28, QA16 117, 176, 185 SM13 363 SA18 319 RR22 314 51 CD03 339, 343 SF14, SG22 274 RG14 164 PL03 QB07 188 RB22, RB23 257 CA03 49 PJ05 155 201 QE15 QB31 192 65 EB04 RR26 315 315 RR29 189 QB12 RD12, SB29 264, 325 RB13 256 RC13, RC16, RC17 261 EH02 70 SL09 357 JF01, JF05 102, 103 QJ01 221 BF05, JG05, PA09, PH18 45, 104, 117, 149 PL16 166 RC07, SD10, SD11 260, 331 317 SA06 SA05 317 RI20 282 87, 290 HF05, RL06 52, 290 CE04, RL03 RB14 256 65 EC01 201 QE12 190 QB22 233 QM10 60 DF02 255 RB06 QH11 214 PM25 171 PJ31, PK21, QK10 159, 163, 225 PO13 179 AJ04 40 RP20 306 CH04 54 QN12 238 SL13 358 QN05 237 PC19 129 SH05 344 EA04 64 PN14 174 RA28 254 HJ05 90 HB02, PM16 84, 170 DE05, RJ16, RJ18, SH14 60, 285, 345 RN20 300

Author Index

Jung, Jonghoon	PL19	167
Jung, Jong-suck	QA12	184
Jung, K.y.	PJ31	159
Jung, Kuyeol	PK21	163
Jung, M. H.	RK19	289
Jung, Min Hwa	PN28	176
Jung, Min-cherl	RC14	261
Jung, Min-hwa	SA18	319
Jung, Mung-hwa	PJ31	159
Jung, Myung- Hwa	RC07	260
Jung, Myung-chul	SO06	367
Jung, Myung-haw	Q011	242
Jung, Myunghwa	SM13	363
Jung, Myung-hwa	PC12, PK17, PN19, QH07, QO10, RK06	128, 163, 175, 214, 242, 286, 331, 349
	SD10, SD11, SJ02	
Jung, S. H.	RK19	289
Jung, Sanghoon	RK06	286
Jung, Sang-yong	RR11, RR18, RR21, RR22	312, 313, 314
Jung, Sol	PM15	169
Jung, Soon-gil	GA04, RB27, SB28	76, 258, 324
Jung, Sung Won	RC14	261
Jungwirth, T.	IF01	94
Jungwirth, Tomas	BE05, BI03, IE03, SK22	44, 47, 93, 354
Juranyi, Fanni Juraszek, Jean	PO17	180
Juraszek, Jean Jurchescu, Oana D.	QA27 AE04	187 36
Jurchescu, Oana D.	AE04	30
К		
K, Ramesh Kumar	SE21, SO29	336, 371
K. Sato, Noriaki	QD18	199
K., Samatha	BJ05	48
Kabiraj, D.	PG27	146
Kachi-terajima, Chihiro	QF21	207
Kaczorowski, Dariusz	AB03, QD05, SE16, SE17, SF12	34, 197, 335, 336, 339
Kaczorowski, Dariuzs	RD23	265
Kadowaki, Hiroaki	QE13	201
Kadowaki, K.	KJ01	113
Kadowaki, Kazuo	RB06	255
Kagayama, Tomoko	PB05, PB27, PC23, SG20	120, 124, 130, 342
Kageyama, Ryu Kajiyu, Hidaa	RP14 SM09, SO07	305
Kaiju, Hideo Kaipuma, Byosuko	PG04, PG07, QG08,	362, 367
Kainuma, Ryosuke	RP07, SK23	143, 209, 304, 354
Kajimoto, Ryoichi	DA01, RH12	56, 276
Kajino, J	SE03, SE05, SE13	333, 334, 335
Kajino, J.	KA02	107
Kajino, Jumpei	PD24, RC07	135, 260
Kajino, Junpei Kakazai C. N	QH07	214
Kakazei, G. N. Kakabashi, Yoshiro		109
Kakehashi, Yoshiro Kalon, Goninadhan	EF04, PB24, PF05, QG03 SK05	68, 124, 140, 208 351
Kalon, Gopinadhan Kaloni, Thaneshwor Prashad	PN02	351 172
Kaioni, maneshwor Prashau Kamada, Yasuhiro	RR13	313
Kamada, Yukihiro	PB05	120
Kamal, Khaja Mohaideen	QN12	238
Kamal, Mydeen	IA02	90
Kamarad, J.	PJ23	158

1 K 1	21.4.4	4.60
Kamarad, Jiri	PM11	169
Kamazawa, Kazuya	PO17	180
Kambale, Rahul Chandrakant		235
Kambe, Shinsaku	PD09	133
Kambe, Takashi	RA18, SC04	252, 326
Kameli, P.	PI24	154
Kameli, Parviz	PI23	154
Kamenev, K. V.	GC01	77
Kami, Daisuke	PO03	177
Kamihara, Yoichi	QB16, QC07, SE09	189, 193, 334
Kamihara, Youichi	QC08	193
Kamimori, Tastuo	PG05	143
Kamimori, Tatsuo	QG10	209
Kamionka, Thomas	FD02	74
Kampert, Erik	KB06, PB11, SB18, SD07	108, 121, 323, 331
Kanai, Noriko	QD17	199
Kanai, Yasushi	BG04	46
Kanatzidis, M. G.	GA01	75
Kanazawa, Ikuzo	PB31	125
Kane, Shashank N.	AJ05	40
Kanehira, Toma	GE04, JE04	78, 102
Kaneko, Chikafumi	PH05	147
Kaneko, Koji	JJ05, QJ09, SC23	106, 222, 329
Kaneko, Masaki	QD18	199
Kaneko, Takejiro	Q105	218
Kaneko, Tatsuya	PF09, PF10	140
Kaneko, Ulisses F	PI19	153
Kaneko, Y.	CG02, RA05	53, 250
Kaneko, Yoshio	PH08	148
Kaneshita, Eiji	CA01, SC10	48, 327
Kanetsuki, Hiroaki	PG05	143
Kang, Boyoun	QI15	220
Kang, Byeongki	BF05, JG05, PA09, PH18	45, 104, 117, 149
Kang, Byeongwon	PB03	120
Kang, Byung-sub	PN31	177
Kang, Chan Seok	RQ19, RQ20	310
Kang, Chang-jong	SC12	327
Kang, Chong Yun	RA21	252
Kang, Han Sam	QN15	238
Kang, Hanhim	QH26	217
Kang, J. H.	RA07	250
Kang, J.k.	Q027	245
Kang, Js.	BB02, PG11, PH14,	41, 144, 149, 183
	QA04	
Kang, Jeong Soo	PB12	122
Kang, Jun-goo	QL15	230
Kang, Min Gyu	RA21	252
Kang, Sun Hee	PA11, QA16	117, 185
Kang, Sung	QG16, QH15	210, 215
Kang, Sunhee	SM13	363
Kang, W. N.	GA04, PB03	76, 120
Kang, Won Nam	RB27, SB28	258, 324
Kang, Woun	КЈ04	114
Kanjilal, D.	PG27	146
Kanke, Yasushi	PI16	153
Kanki, Teruo	SK16	353
Kanoda, Kazushi	BC03	42
Kanomata, Takeshi	PG04, PG07, QH21	143, 216
Kanou, Manabu	RC12	260
Kanoun, Mohammed Benali	QL17, SK12	230, 352
		,

Kanter, J.	PI15	153
Kanter, Jakob	RB21, SC16	257, 328
Kao, Cc.	PH10, RA04	148, 249
Kao, Chih-kuei	IH06	96
Kapaklis, Vassilios	AF05, EI03, RP15	37, 71, 305
Karaki, Y.	EF01	68
Karaki, Yoshitomo	CC05	50
Karamat, Shumaila	PN03, QL02	172, 228
Karipoth, Prakash	AJ04	40
Karla, J. Merazzo	GG05	80
Karpinski, J.	PC01	126
Karpinski, Janusz	RB19, RB20, RB21	257
Karthigeyan, A	SA21	319
Karthik, T	PN18	175
Karube, Kosuke	RD05, RD08	262, 263
Kasagi, Teruhiro	QN02	236
Kasai, S	SK29	355
Kasai, Shinya	HG02, ID04, IF05, SK31	88, 93, 94, 355
Kasatani, Yuichi	RK10	287 113
Kashiwagi, T. Kashiwagi, Takanari	KJ01 RB06	255
Kassir-bodon, Zoukaa	SH06	344
Kastil, Jiri	PM11	169
Kataev, V.	QE23, QF22	203, 207
Katano, Susumu	RH01	274
Katayama-yoshida, Hiroshi	KF04	111
Kato, Akihiko	RH17	277
Kato, Harukazu	QD06	197
Kato, Haruki	PL06, RL02	164, 290
Kato, Hiroaki	EJ01, RO05, RO06	71, 301
Kato, Kenichi	QF25	208
Kato, Masaru	PC26	131
Kato, Motoharu	SG18	342
Kato, Reizo	BC04, SF01	42, 337
Kato, Takashi	PO18	180
Kato, Takeshi	RL08	291
Kato, Tomoki Katob, Hicachi	PC23	130
Katoh, Hisashi Katoh, Kapishi	PN07	173
Katoh, Kenichi Katori, Hiroko Aruga	QD09 QD13	197 198
Katou, Haruki	RC02	259
Katrych, Sergiy	RB21	257
Katsufuji, T.	PA04	116
Kaul, S. N.	QD10, QM06, SL07	198, 232, 357
Kawada, Yuki	SI03, SI17	346, 348
Kawae, Tatsuya	QE18, QF14, QL11,	202, 206, 229, 230,
	QL12, QL13, RE04, SG09	267, 341
Kawagoe, T.	PL04	164
Kawaguchi, Masashi	RL10	291
Kawaguchi, Shogo	PI02	151
Kawahara, Shin-ichi	EI05	71
Kawai, Tetsuroh	PM18, SM15	170, 363
Kawakami, Norio	DB05, HI04, PF16, RC01, RC04, RD05, RF01, RF09	57, 89, 141, 259, 262, 269, 270
Kawamata, Takayuki	DA01, QF07, QF08, QF10	56, 205
Kawamoto, Atsushi	PC15	129
Kawamura, Naomi	IG03, QD17	95, 199
Kawamura, Yukihiko	JJ05	106
Kawamura, Yukihiro	PD10, PD22, QD06	133, 135, 197

## AUTHOR INDEX

Kawana, D.	CF04	53
Kawana, Daichi	PI07, RD14, RH13	151, 264, 276
Kawanaka, Hirofumi	PM13	169
Kawanishi, Yohei	HE03	86
Kawar, Tatsuya	RQ05	307
Kawasaki, Ikuto	RD25, SE23	266, 337
Kawasaki, M.	RG10	273
Kawasaki, S	QB11	188
Kawasaki, Y.	SG02	339
Kawasaki, Yu	QF24	207
Kawashima, Kenji	PB06, PB07, SG16	121, 342
Kawashima, Naoki	DC04	58
Kawata, K.	SG11	341
Kawata, Tomokazu	PD22	135
Kawata, Toru	RD19	265
Kazakov, Alexander	CH01	54
Kazuhei, Wakiya	QC16	194
Kazunori, Umeo	QC16	194
Kazuo, Watanabe	QJ01	221
Kazuyuki, Matsubayashi	QI17, QJ01	220, 221
Ke, W. P. Ke	IC03	92
Keiderling, Uwe	DF02	60
Keimer, B.	GA03	75
Keimer, Bernhard	BA03, CA03, IJ02, JC04,	41, 49, 97, 100, 257
	RB20	
Keisuke, Matsumoto T	QC16	194
Kelekci, Ozgur	PJ24	158
Keller, Lukas	PH02, PI05	147, 151
Keller, Thomas	GC04	77
Kennedy, Shane J	QG09	209
Kennedy, Shane Joseph	EF02	68
Kenzelmann, M.	QE19	202
Kenzelmann, Michel	CD04, QF20	51, 207
Кера, М.	GC01	77
Kezsmarki, Istvan	PA14	118
Khademi, F.	PI24	154
Khajehnezhad, A.	QP23	249
Khalil, H. M. Waseem	PJ24	158
Khaliullin, Giniyat	BA03	41
Khalyavin, D. D.	KA02	107
Khan, Mohammed N. Islam	EI02	70
Khan, Nawazish Ali	PB13	122
Khan, Tahir Rao	IA02	90
Khan, Yasmin	Q025	244
Kharchenko, M. F.	PH06	147
Kharchenko, Yu.	PH06	147
Khasanov, Rustem	CD04, IA02, QB22	51, 90, 190
Khavronin, V. P.	RQ13	309
Khim, Seunghyun	JA03	98
Khizroev, Sakhrat	PN16	174
Khmelevskyi, Sergii	EF05	68
Khodaei, M.	QM28, SA04	236, 316
Khokhlov, Alexey R	Q014	243
Khomskii, Daniel	PA17	118
Khovaylo, Vladimir	DJ04	63
Khovaylo, Vladimir V.	SO19	369
Khusainov, Mansur G.	QA26	186
Khusainov, Marat M.	QA26	186
Khvalkovskiy, A.v.	ID02	93
Khvalkovskiy, Alexey	ID03, IE04	93, 94

Khym, S.	QL08	229
Ki, Sanghoon	PL15, QM21	166, 235
Kida, Takanori	QB09, QF06	188, 204
Kiefer, K	RI05	279
Kihara, Takumi	Pa06, Qa09, QG08, Ra06, RH04	116, 184, 209, 250, 275
Kihou, K.	QB28	191
Kihou, Kunihiro	QB02, QB15	187, 189
Kijima, Hanae	JI01	105
Kikkawa, Akiko	Q109	222
Kikoin, Konstantin	FE04	75
Kikuchi, Hikomistu	QF17 RI15	206 281
Kikuchi, Hikomisu Kikuchi, Hikomitsu		202, 204, 205
Kikuchi, Hiroaki	QE16, QE21, QF04, QF09 RR13	313
Kikuchi, Hiroshi	SK21	353
Kikuchi, Yoshihiro	PF17	142
Kikuchi, Yuhei	QB27	191
Kim Anh, Do Thi	AJ06	40
Kim, Aaram Joo	RE07	267
Kim, Ahri	PM15	169
Kim, Beom Jun	RI19, RI20, RI21	282
Kim, Bongho	QP19, RL22, SL19	248, 293, 359
Kim, Bongjae	PG11	144
Kim, Bum-su	QL18	231
Kim, Bun-su	PN20	175
Kim, Byeong-geon	PL15, QM21	166, 235
Kim, Changsoo	BF05, JG05, PA09, PH18	45, 104, 117, 149
Kim, Chang-yeoul	RN24	300
Kim, Chanhee Kim, Chalai	RQ11	308 243
Kim, Chelgi Kim, Cheol Gi	QO19 SN13	366
Kim, Cheolqi	PO19, QN22, RN21, SL22	
Kim, Chin Mo	PA22	119
Kim, Choong H.	RC13, RC17	261
Kim, Chul Sung	PA22, PI13, PI14, RA26	119, 152, 153, 253
Kim, Chung Koo	CA02	48
Kim, Chung Man	QH07	214
Kim, Chungman	Q011	242
Kim, D. H.	PG11	144
Kim, Dw.	PO23	181
Kim, D.h.	PH14, QA04	149, 183
Kim, D.y.	EC01	65
Kim, Daehong	QP19, RL22, SL19	248, 293, 359
Kim, Daehyun Kim, Danbee	PB12	122
Kim, Dahbee Kim, Dohyang	RQ07 PM17	308 170
Kim, Dong H.	QL16	230
Kim, Dong Young	PM28	172
Kim, Dong-hyun	RP03, SH01, SH10	303, 343, 344
Kim, Dongseok	PK21	163
Kim, Dongsoo	RN23	300
Kim, Dongyoo	JH03, RM07	104, 296
Kim, Do-yeon	RN20	300
Kim, Eun-mee	RN20	300
Kim, Eunseong	PC25	130
Kim, Eun-young	RE01	266
Kim, G C	QB25	191
Kim, G. B.	RQ08	308
Kim, G. W.	PM14	169

Kim, Gwang-hee	SIO1	346
Kim, Gyutae	PJ30	159
Kim, H.	IA05, QC23, SB30	91, 196, 325
Kim, Hae Jin	SL11	357
Kim, Hanbit	JB02	99
Kim, Hee Seung	PI14	153
Kim, Heung Sik	RC13	261
Kim, Heungsik	DC03	58
Kim, Heung-sik	RC17	261
Kim, Hojung	PH21	150
Kim, Hong-seok	SD12, SD14, SO21	332, 370
Kim, Ho-sup	QM22	235
Kim, Hui Min	RR07	312
Kim, Hwijun Kim, Hwiseok	PM17, PM23, PM24 PO14	170, 171 179
Kim, Hyeong-do	RG04	272
Kim, Hyo-jin	PJ19, QO28, RH18	157, 245, 277
Kim, Hyung Joon	JB02	99
Kim, Hyung Jun	PK10	162
Kim, Hyungjun	PJ16	157
Kim, Hyung-jun	KC04, PL16, QK11,	109, 166, 225, 226,
,,	QK12, QK13, RL28, RL31	
Kim, Hyungsuk K. D.	QA15	185
Kim, Hyungyu	QP19	248
Kim, Hyun-jung	SD06	331
Kim, Ill Won	PA11, QA16	117, 185
Kim, Illwon	SM13	363
Kim, In Gee	QG20	211
Kim, Ingyu	JB02	99
Kim, In-seon	SC15	328
Kim, J.	RA07	250
Kim, Jy.	BD03, PG11, QA04	43, 144, 183
Kim, Jae Nyeong	QJ19	224
Kim, Jae Young	PA11	117
Kim, Jaehyun	PO06	178
Kim, Jae-sung	AD03	35
Kim, Jaewook	RQ11	308
Kim, Jaeyeong	RA21	252
Kim, Jae-yeong	SO03	367
Kim, Jaeyoung	SM13	363
Kim, Jae-young	AF04	37
Kim, Jeehoon Kim, Ji Wan	QB31, SB12 QM20	192, 322 234
Kim, Ji-hye	RJ18	234
Kim, Jimin	RQ11	308
Kim, Jin Mok	RQ17	309
Kim, Jin Woo	RO10	302
Kim, Jinhee	QO10, SD11, SD14	242, 331, 332
Kim, Jin-hee	CF03	52
Kim, Jinsung	SN07	365
Kim, Jiwan	RK13	288
Kim, Ji-wan	FC02, RK09	73, 287
Kim, Jiyeon	PI10	152
Kim, Jong Hee	QN22	240
Kim, Jong-hwa	RR27	315
Kim, Jongryoul	HB02, PM16	84, 170
Kim, Jong-woo	DJ02	63
Kim, Joon-il	QK23	227
Kim, Ju Young	QH29	217
Kim, Ju_young	SC08	326

Kim, Ju-jin	SD12	332
Kim, Jun Sung	RB14, SD07	256, 331
Kim, June Seo	FD03, QK14, SH16	74, 226, 345
Kim, Jung Dong	PM28	172
Kim, Jungeun	HC03	85
Kim, Jungho	DC03	58
Kim, Junghwa	IJO2	97
Kim, Jun-young	PN23	175
Kim, Juyoung	QI15	220
Kim, K. W. Kim, Kab-jin	SL08, SL09 KE03, SH07	357
Kim, Kee Hoon	CI03, JA03, JB02, PA04,	110, 344 55, 98, 99, 116, 117,
Kin, Kee hoon	PA09, PB23	123
Kim, Keun-ho	RR15	313
Kim, Ki Yeon	QM18, QM20	234
Kim, Ki-chan	RR09, RR10	312
Kim, Kiwoong	PO11, RQ17, RQ18, RQ19, RQ20	179, 309, 310
Kim, Kwang Hee	RR19	314
Kim, Kwang Heon	QO20, SG21	244, 343
Kim, Kwang Lok	PB16, PB18, PB35, PB36, PB38	
Kim, Kwang S	QH26	217
Kim, Kyeong-sup	PN31	177
Kim, Kyoo	SC12	327
Kim, Kyoungchul	SA09	317
Kim, Kyoung-whan	EE05, HP21	68, 30
Kim, Kyung Ho	KC04	109
Kim, Kyung-ho	PL16, QK12, RL31	166, 226, 295
Kim, Kyunghyun	RN10	298
Kim, Kyungwan	HA02	83
Kim, Kyung-whan	PJ33	160
Kim, Kyu-tae	RQ07	308
Kim, Mi Jung	SN04	364
Kim, Mi-jung	RR14, SN02	313, 364
Kim, Min-ho	RR23	314
Kim, Minseon	RA26	253
Kim, Miryeon Kim, Miyoung	RL17 RH20, SK27	292 278, 354
Kim, Mun-seog	RQ07	308
Kim, N.	QF13	206
Kim, Ran-hyang	BH02	46
Kim, Rokyeon	GF05	79
Kim, S. B.	PG11	144
Kim, S.b.	PH10	148
Kim, Sam Jin	RA26	253
Kim, Sang-beom	PM16, PM19, PM21	170
Kim, Sanghoon	RL30	294
Kim, Sang-il	RL06	290
Kim, Sang-koog	DE05, RJ14, RJ15, RJ16, RJ17, RJ18, SH14, SJ07	60, 284, 285, 345, 350
Kim, Seong-gon	KH04	112
Kim, Seung Hyun	RK06	286
Kim, Seunghyun	RN10	298
Kim, Shin Ae	RP01	303
Kim, Shin-ae	PI10, RA22	152, 253
Kim, Soo Hyun	QH07	214
Kim, Soohyun	PC12	128
Kim, Soo-whan Kim, Soveon	PC12 RH08	128 276
Kim, Soyeon	NIIVO	270

## AUTHOR INDEX

Kim, Suhyun Kim, Sumin Kim, Sun Yung Kim, Sung Baek Kim, Sung Hoon Kim, Sung Wng Kim, Sung Yong Kim, Sung-beak Kim, Sungho Kim, Sunghyun Kim, Sung-jin Kim, Sungyong Kim, Tae Hee Kim, Tae Wan Kim, Tae-hoon Kim, Tae-suk Kim, Taewan Kim, Taewoo Kim, Tae-young Kim, Taeyueb Kim, Wan-seop Kim, Won Dong Kim, Wondong Kim, Won-kyung Kim, Y C Kim, Y. H. Kim, Y.i. Kim, Yong Baek Kim, Yong-jae Kim, Yongjin Kim, Yongmin Kim, Yoonbai Kim, Young Do Kim, Young Keun Kim, Young Kwang Kim, Young-gyun Kim, Young-hak Kim, Young-june Kim, Youngkwang Kim, Yun Won Kim, Yun-won Kimel, Alexey V. Kimishima, Yoshihide Kimling, Judith Kimura, Akio Kimura, Kenta Kimura, Nobuyuki Kimura, Noriaki Kimura, S. Kimura, Shin-ichi

Kimura, Shin-ichi Kimura, Shoji Kimura, Shojiro Kimura, Takashi

Kimura, Tsuyoshi

Q011 242 PM15 169 SL01 356 PI13 152 PO20, SM08, SN14 180, 362, 366 130 PC23 162 PK10 PA04 116 KH04 112 PO06 178 RR11 312 157 PJ16 CB01, PJ17, QP18 49, 157, 248 PJ24 158 RO01, RO02 300, 301 159 PJ31 157, 162, 163 PJ16, PK10, PK17 313, 365 RR16, SN05 SM10 362 QK10 225 RQ07 308 RK19 289 BD03, IH04, QM26, 43, 96, 235, 288, RK18, RP23 307 150, 175 PH22, PN20, PN21 QB25 191 RQ08 308 123 PB19 RA22 253 312, 314 RR11, RR21, RR22 178 PO06 179, 181, 242, 310, PO14, PO26, QO08, RQ23, RR12 312 SD03 330 RO10 302 KC04, PJ14, RK06, SH13 109, 157, 286, 345 PK10 162 PB17, PB37, PB40 122, 126 PO14, PO26, RQ23, 179, 181, 310, 315 RR30 DC03 58 PJ16 157 RQ15 309 SC14 327 RK17 288 SB03 320 RP04 303 KF05, SD11 111, 331 CC05 50 EF04 68 131, 192, 193, 194, PC28, QC03, QC10, QC11, RD08 263 KA02, SG11 107, 341 AD04, QD07, QD18 35, 197, 199 QI18 220 151, 210, 278 PI06, QG15, RH25 AG03, HE01, KC06, 37, 86, 109, 228, QK24, QL11, RJ08, 229, 283, 289, 293, RK21, RL20, RL23, SI12 347 JB01, QA12 99, 184

Kindo, K.	GJ03, QI03	82, 218	Knight, I
Kindo, Koichi	HC03, PA18, PD13,	85, 118, 134, 137,	Knowles
	PE10, PG17, QE16, QF11, QF14, QF17, RI15,	145, 202, 205, 206, 281, 304, 307	Ko, Jung
	RP07, RQ05	201, 304, 307	Ko, Kt.
		07	Ko, Kow
King, Phil	1102	97	Ko, Kyou
Kino, Ryo	QA09	184	Ko, Rock
Kinoshita, Toyohiko	RP07	304	Ko, Seur
Kinouchi, Hiroaki	QB29	191	Ko, Sung
Kiparisov, Semion Kirchner, Stefan	SM19	364	Ko, Won
Kirillov, Oleg A.	JC01 AE04	100 36	Kobayas
Kirilyuk, Andrei	RK17	288	Kobayas Kobayas
Kirschner, J	RP20	306	Kobayas
Kirschner, Jurgen	AD02, AE01, El01	35, 70	Kobayas
Kiryukhin, V.	RA22	253	Kobayas
Kiseki, Kohei	AG03, SI12	37, 347	Kobayas
Kishimoto, Kengo	QM10	233	Kobayas
Kishimoto, Y.	SG02	339	Kobayas
Kishimoto, Yutaka	QF24	207	Kobayas
Kishine, Jun-ichiro	DF03, PG09, PG10,	60, 143, 144, 211,	Kobayas
	QG22, QG28	212	Kobayas
Kisielewski, Marek	AD05		Kobayas
Kiss, T.	QB28	191	Kobayas
Kiswandhi, Andhika	SB05	321	Kobayas
Kita, Eiji	PG05, PI16, QI04, QM16,	143, 153, 218, 234,	Kobayas
	SJ06	349	Kobayas
Kita, Tomoko	RF09	270	Kobayas
Kitagawa, Jun	EG05	69	Kobayas
Kitagawa, Shunsaku	QC08	193	
Kitagawa, Susumu	HC03, QF25	85, 208	Kobayas
Kitagawa, Yutaro	QA10	184	Kobiyam
Kitai, Tetsuo	Q105	218	Kodama
Kitamura, T.	KJ01	113	Kodama
Kitanishi, Fumiya	SO11	368	Kodama
Kitaoka, Yoshio	DA04, PB26, QB29	56, 124, 191	Kodderit
Kitaoka, Yukie	SO10	368	Kodera,
Kitazawa, Genki	SM08	362	Koehler,
Kitazawa, Hideaki	JJ05, PH02, PH09, QA09, QJ02	106, 147, 148, 184,	Koehler,
Kito, Hijiri	Q802 QB02	221 187	Koenraa
Kittaka, S.	EF01	68	Koerner,
Kittaka, Shunichiro	QJ08	222	Koga, Al
Kiyanagi, Yoshiaki	RP01	303	Koga, M Koh, Eui
Kiyokawa, Takayasu	QA02	182	Kohama
Klaeui, Mathias	FD03, QK14, SH16	74, 226, 345	Kohara,
Klaui, Mathias	CH04, JE03	54, 102	Kohda, M
Klauss, Hans-henning	AB02, PI05, QB22, QJ03	34, 151, 190, 221	Kohlbred
Klein, Olivier	ID03	93	Kohler, S
Klemm, R. A.	KJ01	113	Kohlhep
Klicpera, Milan	QI01	217	Kohmoto
Klimczuk, Thomasz	SE19	336	Kohn, Ka
Klimin, Sergey A.	RA24	253	Kohno, A
Klingeler, R.	QE23, QF18, QF22	203, 206, 207	Kohno, I
Klintberg, Lina	FB04, PC09	73, 128	Kohori, Y
Klintberg, Lina Esther	PC02	127	Kohori, Y
-		26	
Kloc, Christian	AE05	36	
Kloc, Christian Klyatskaya, Svetlana	AE05 GH02	81	Koichi, T
			Koichi, Ta Koike, Ka

, Kevin	SD13	332
es, Elisabeth S.	KB03	107
ngho	JI02, RL30	105, 294
t.	PH10	148
wan-young	QG01	208
oung Chul	QG17	211
ck Kil	QB25	191
ung Hyo	CB01	49
ng-won	QG01	208
onsuk	PO13	179
ashi, H.	CG02	53
ashi, Hiroki	RD14	264
ashi, Hisao	QD11, RB03, RD13	198, 254, 264
ashi, K.	SH07	344
ashi, Keisuke	RD20	265
ashi, Kensuke	RL10, SE07	291, 334
ashi, Kohei	SA07	317
ashi, M.	PC01, PC29, PI15	126, 131, 153
ashi, Masaki	HA03	83
ashi, Nobukiyo	JIO1	105
ashi, Norio	QF10, SG18	205, 342
ashi, Riki	SC23	329
ashi, Satoru	PH05, PI22, RR05, RR13	147, 154, 311, 313
ashi, Tatsuo C	QF25, SE10	208, 334
ashi, Tatsuo C.	HC03, QC18	85, 195
ashi, Tatsuya	QB24	191
ashi, Toshiaki	RH01	274
ashi, Tsukasa	SK21	353
ashi, Yoshiaki	DA01, QB04, QB26,	56, 187, 191, 254
	RB02, RB05	
ashi, Yuta	SL02	356
ma, Takuya	QD13	198
na, Kenji	RP07	304
ia, Motoi	RJ10	284
na, Yuta	QE20, QF21	202, 207
ritzsch, Diemo	1103	97
a, Mitsuru	EC03	65
r, Daniel	PC27	131
r, Juergen	QJ18	223
aad, Paul	BE05	44
er, Michael	RJ11	284
Akihisa	SC13	327
Mikito	SB17	323
ui-kwan	EH02	70
na, Yoshimitsu	SF13	339
a, Takao	QD12, QD15, RD24	198, 266
, Makoto	KC03	108
echer, Joachim	QF20	207
Sebastian	RI13	281
pp, Juergen	JE01	101
oto, Osamu	SL03	356
Kay	PI16	153
, Atsushi	SA03	316
, Masanori	RE05	267
, Y	QB28	191
, Yoh	QB02, QB15, QD12,	187, 189, 198, 201
	QE12	
Takase	QD14	198
Kazuyuki	HG01, RP14	87, 305
Kunihiro	EJ01, RO05, RO06	71, 301

	DI O A	200
Koike, Mariko Koike, Yoji	RL04 QF07, QF08, QF10	290 205
Koinuma, Hideomi	PH22, PN21, QL18	150, 175, 231
Kokorina, E. E.	RF16	271
Kokubu, Mitutane	QD12	198
Kolano - Burian, A.	PM05	168
Kolbe, M	RP20	306
Kolbe, Michaela	CH04	54
Koledov, Victor V.	5019	369
Kolesnikov, Alexander I.	RD14	264
Kollath, Corinna	CC03	50
Kolodiazhnyi, Taras	RH06	275
Kolomiets, A.	RD26	266
Kolomiets, Aleksandre	RE09	267
Kolomiets, Olexander	QI19	220
Komaki, Yasuhiro	QD12, QE12	198, 201
Komarek, Alexander	PI17	153
Komatsubara, Takemi	RD08	263
Komorida, Yuki	KB06	108
Kondo, A.	QI03	218
Kondo, Akinori	PE10	137
Kondo, Kenji	QK05, SM09, SO07	224, 362, 367
Kondo, Takaya	QL03, RN02	228, 297
Konishi, K.	SL14	358
Konishi, Takehisa	KG03, PF10, PF15, SF04	111, 140, 141, 337
Konno, Rikio	PF08	140
Konno, Toyohiko J.	KB05	108
Kono, Hiroshi	RD18	265
Konstantinovic, Zorica	IH01, SO18	95, 369
Koo, Bon Heun	PM14, PN15	169, 174
Koo, Dae-hyun	PM25, RR31	171, 316
Koo, Hyun Cheol	PJ30, QK11, QK12, QK13, RL28	159, 225, 226, 294
Koo, Hyuncheol	QK10	225
Koo, Hyun-joo	PH11	148
Koo, J.w	SK29	355
Koo, Tae Yeong	PA11, QA16	117, 185
Koo, Taeyoung	SM13	363
Koopmans, Bert	CH05, EI06, GE02, IE02, JE01, SI07	54, 71, 78, 93, 101, 347
Kopeliovich, Alexander I	PJ02	155
Korbecka, A.	BE02	44
Korekawa, Tomoya	PN07	173
Korelis, Panagiotis	RP15	305
Korolev, Alexey	PB14	122
Korolyov, Alexandre	QL07	229
Korotana, Romi Kaur	RF15	271
Kosaka, Masashi	QD13, QI16	198, 220
Koseki, Gen	PD11	133
Koseki, Minoru	PD11	133
Kosel, Juergen	SA02	316
Kosel, Jurgen	PJ06, PM27, RQ02	155, 171, 307
Koshiba, Shyun Kossick, Martin	PN24 KD01	176 109
Koster, Gertjan	DH02, SA10	62, 317
Kota, Yohei	PJ10, PL17	156, 166
Kotaka, Hiroki	RC09, RC19	260, 262
Kotani, Y.	QB28	191
Koteawa, Hisashi	QB17	189
Kotegawa, Hisahi	SG12	341
<u> </u>		

Kotetes, Panagiotis Kothapalli, Karunakar Kothari, D. C. Kotliar, G. Kotnala, R. Kototani, Shouhei Koubaa, Mohamed Kouh, Taejoon Kousaka, Yusuke Koutroulakis, Georgios Kovac, Jozef Kovac, P Kovylina, M. Kovylina, Miroslavna Kowalczyk, M. Kowalik, Iwona Koyama, Keiichi Koyama, Kuniyuki Koyama, T. Koyama, Takehide Koyama, Tsukasa Koyanagi, Tsuyoshi Koyota, Yuya Koza, Michael M. Kozono, Yuuki Kraemer, Karl Kraken, Mathias Kramar, Oleksandr Kramarenko, Elena Yu Kramer, K W Kramer, Steffen Kratochvilova, Marie Krause, Stefan Krawczyk, M. Krawczyk, Maciej Krellner, Cornelius Kremer, Reinhard Kremer, Reinhard K Kreuzpaintner, Wolfgang Krinitsina, Tatiana Krinitsina, Tatiana P. Krishnakumar, Varadharajan Krishnan, Murali Krive, Ilya V. Krockenberger, Y. Kroha, Johann Krompiewski, Stefan Kronast, Florian Krotscheck, Eckhard Krueger, Benjamin Krug V Nidda, Hans-albrecht Kruger, Natalia Kruglyak, Volodymyr Kruglyak, Volodymyr V.

Kotegawa, Hisashi

PD25, SF03 136, 337 1104 97 SF13 339 299 RN16 PE17 138 PG27 146 187, 254 OB04, RB02 64 DJ05 PA22 119 DF03, PG09, QG22, 60, 143, 211, 212 QG28 127, 337 PC03, SF01 IB01, QO17 91, 243 RB08 255 BD01 43 BD04 43 PH06, PM05 147, 168 DD02 58 GJ02, QG15, QH08, 82, 210, 214, 218, QI05, QN20, SK23 239, 354 PB25, QC05 124, 193 RL07 291 QD15, RD24 198, 266 AF03, PG15, QG22, 37, 144, 211, 212 QG28 QM10 233 PD11 133 268 RE12 52 CE05 210 QG14 104, 299 JG06, RN15 274 RG13 243 Q014 RI05 279 IC02 92 PC10, PD18, PD19, 128, 134, 135, 223, 321 QJ14, SB08 HF03, HF04, RP16 87, 305 DD02 58 Q005 241 AB02, CD03, FB04, 34, 51, 73, 100, 133 JC01, PD12 JB04 99 PH11 148 DG01 61 PB22, SK02 123, 350 SK01 350 QC13 194 163 PK16 RG17 274 RG10 273 PF07 140 QK04 224 FD03, GG05, RP04, RP08 74, 80, 303, 304 274 RG15 74 FD02 278 RH24 201 QE14 70, 85, 103 EH04, HD03, JG01 78, 241, 289 GD04, QO05, RK20

Kruize, Michelle	DH02, SA10	62, 317	Kumar, Prade
Krychowski, Damian	PE06	137	Kumar, Rajes
Krycka, Kathryn	KD02	109	Kumar, Ravi
Krzysteczko, Patryk	IF04	94	Kumar, Shale
Ku, H. C.	SB04	320	Kumar, Vino
Ku, Huan-chiu	RR01	311	Kumari, Poo
Ku, Jang Hae	PJ30	159	Kummer, Ku
Kuang, Lulin	PB28	124	Kuo, C. N.
Kubetzka, Andre	IH02, IH03	96	Kuo, K. M.
Kubinova, Sarka	BH03	46	Kuo, Nan-ho
Kubo, Hidenori	RH19, RP12	277, 305	Kuo, Yung-k
Kubo, Katsutaka	PG17, QF11	145, 205	Kupriyanova
Kubo, Kazuki	AI02	39	Kuramochi, I
Kubo, Kazuya	BC04	42	Kurde, Julia
Kubo, Yasunori	SE01	333	Kurihara, Ryo
Kubota, H.	CE03, CE06, ID02, RJ07	51, 52, 93, 283	Kurisu, Maki
Kubota, Hitoshi	BI01, HF01, IE04, KF01,	47, 87, 94, 110, 163	Kurita, Kohe
	PK20		Kurita, Nobu
Kubota, Kazuhiro	RD01	262	Kuroda, Keis
Kubota, M.	EF01	68	Kuroe, Harul
Kubota, Miho	RL13	292	
Kubota, T.	SK24	354	Kurokawa, A
Kubota, Takahide	CH02	54	Kurosawa, To
Kubota, Tatsuro	BC04	42	Kusano, Taka
Kubota, Yoshiki	HC03, PI02, QF25	85, 151, 208	Kusunose, H
Kubozono, Yoshihiro	SC04	326	Kutnyakhov,
Kucherenko, Yuri	PD12	133	Kuwada, Kei
Kuchko, Andriy	JG01	103	Kuwahara, H
Kuchler, Robert	AB02	34	
Kudo, Kazutaka	QF10	205	Kuwai, Tomo
Kudrevatykh, N. V.	Q103	218	Kuwana, Kai
Kudrevatykh, N.v.	GJ03	82	Kuwano, Ke
Kudryavtsev, Y. V.	QG27	212	Kuzian, R. O
Kuepferling, Michaela	HH01	88	Kuzmin, D. A
Kuga, Kentaro	CD02, IG03, QD01,	51, 95, 196, 264	Kuz'min, M.
	RD13		Kuznetsov, A
Kuga, Kiyoshi	SK21	353	Kuznetsova,
Kuhns, Phil	JA02	98	Kwilu, Augu
Kuiper, Koen	SI07	347	Kwon, Hae-v
Kulemanov, Ivan V	PN22	175	Kwon, Hyuc
Kulinich, Sergey I.	RG17	274	Kwon, Jae H
Kulka, A.	PH06	147	Kwon, Joonh
Kulkarni, Prabhanjan Dilip	PK16	163	Kwon, O. Y.
Kulkarni, Ruta	QJ17, RB24	223, 258	Kwon, Oh Ju
Kulkarni, Ruta N	SE21	336	Kwon, Oryor
Kumada, Nobuhiro	QE10	201	Kwon, Sangi
Kumagai, Keigou	SG05	340	Kwon, Seun
Kumagai, Ken-ichi	SB27	324	Kwon, Woo
Kumai, R.	CF04	53	Kwon, Yong
Kumai, Reiji	PH07, SE07	148, 334	Kwon, Youn
Kumano, Yuta	RI02	279	Kyakuno, Ha
Kumar, Anil P. S	QM03	232	Kyomi, Takuy
Kumar, Anil P. S.	CG04, SJ08, SL23	53, 350, 359	L
Kumar, Dheeraj	HD04, RK11	85, 287	
Kumar, G. Ravi	PI12	152	Labarta, A.
Kumar, Manish	SA14	318	Labarta, Ami
Kumar, Manoj	RA27	253	Lacour, Dani
Kumar, Neeraj	RB24	258	Laczkowski,
Kumar, P. S. Anil	PJ28, SL21	159, 359	
Kumar, Pawan	QL10, SM18	229, 363	Lai, Bin

Pradeep	SE14	335
Rajesh	PH16	149
Ravi	PH16	149
Shalendra	PM14, PN15	169, 174
Vinod	PH16	149
, Poonam	PI12	152
er, Kurt	PD12	133
. N.	PH01, RP02	147, 303
М.	QM25	235
an-hong	QF26	208
ung-kang	PE02	136
nova, Galina	SK26	354
ochi, Hiromi	PG06, RP12	143, 305
Julia	RP08	304
a, Ryosuke	RB05	254
Makio	PN07, QG11, SM10	173, 210, 362
Kohei	PD11	133
Nobuyuki	SB05	321
, Keisuke	PD25	136
Haruhiko	PA19, QA05, QA07,	119, 183, 184, 276
	QA09, RH11	
wa, Akinobu	PO02, QL03, RN02	177, 228, 297
wa, Tohru	PB08	121
, Takanao	RO05	301
ose, Hiroaki	SB17	323
khov, D	RN12, RP20	298, 306
a, Kenta	RQ03	307
ara, Hideki	PA06, PA19, QA05, RA06, SB24	116, 119, 183, 250, 324
Tomohiko	SG05	340
a, Kanako	QF25	208
o, Keisuke	RB04	254
R. O.	QF18	206
n, D. Aleksandrovich	OH17	215
n, M. D.		82, 218, 219
sov, Alexey V.	QC19	195
sova, Elena	PB22	123
5	KD04	110 302
Hae-woong	RO12	
Hyuckchan	RQ17	309
Jae Hyun	EG03, HD02, PJ08, SI02	69, 85, 156, 346
Joonhyun	PK04, QK03	161, 224
0. Y.	PM26	171
Oh Jun	PB16, PB18, PB35	122, 123, 125
Oryong	KIO1	112
Sangil Saura a Uhudu	BF05, JG05, PA09, PH18	
Seung Hyuk	QN09	237
Woo Jun	PI14	153
Yong Seung	CF03	52
Young-wan	EH02	70
io, Haruka Taluara	SO515	369
Takuya	PO18	180
, A.	BD01	43
, Amilcar	BD04, KB06, RN14	43, 108, 299
Daniel	KF02, SH02	110, 343
wski, Piotr	GE01, QK15, QK16,	78, 226, 360
	QK17, SL24	266
	SN12	366

Lake, Bella	IC05, PA23, RH24	92, 119, 278
Lam, Duong Duc	PK20	163
Lam, V. D.	RA01	249
Landa-canovas, Angel	PI18	153 42
Landee, C P Landee, C. P.	BC05 QF02	42 204
Landge, Kalpana Kamalkishor	SK19	353
Landsgesell, Sven	RB01, RB07	254, 255
Lang, Michael	HC04, QE14, QJ11, RI13,	
Lang, Michael	SG19	342
Langbehn, Bruno	KD01	109
Langenberg, Andreas	KD01	109
Langie, Douglas	PB10	121
Langridge, Sean	KE01, PN23	110, 175
Larbalestier, David	JA02	98
Larde, Rodrigue	BD02, DG02, EJ03, KB02, PJ22	43, 61, 72, 107, 158
Lari, Leonardo	SK11	352
Larionova, Joulia	RN15	299
Larrea J., J.	CD01, SF10	50, 338
Lartigue, Lenaic	RN15	299
Lascialfari, Alessandro	RN15	299
Lau, Tobias	KD01	109
Laubschat, Clemens	CD03, PD12	51, 133
Lausberg, Stefan	AB02	34
Laver, Mark	KD02, QF20	109, 207
Lawicki, Arkadiusz	KD01	109
Lazarov, Vlado K	SK11	352
Lazo, Cesar	AF02, PG12	36, 144
Lazuta, A. V.	RQ13	309
Le Breton, Jean-christophe	AI04	39
Le Breton, Jean-christophe Le Breton, Jean-marie		39 41, 43, 61, 72, 107, 158, 278
	AI04 BB03, BD02, DG02,	41, 43, 61, 72, 107, 158, 278 288
Le Breton, Jean-marie Le Guyader, Loic Le Tacon, Mathieu	Al04 BB03, BD02, DG02, EJ03, KB02, PJ22, RH22 RK17 CA03, RB20	41, 43, 61, 72, 107, 158, 278 288 49, 257
Le Breton, Jean-marie Le Guyader, Loic Le Tacon, Mathieu Le Tacon, Mathieu	Al04 BB03, BD02, DG02, EJ03, KB02, PJ22, RH22 RK17 CA03, RB20 BA03	41, 43, 61, 72, 107, 158, 278 288 49, 257 41
Le Breton, Jean-marie Le Guyader, Loic Le Tacon, Mathieu Le Tacon, Matthieu Le, Cuong V.	Al04 BB03, BD02, DG02, EJ03, KB02, PJ22, RH22 RK17 CA03, RB20 BA03 QL16	41, 43, 61, 72, 107, 158, 278 288 49, 257 41 230
Le Breton, Jean-marie Le Guyader, Loic Le Tacon, Mathieu Le Tacon, Matthieu Le, Cuong V. Le, Tuan Tu	Al04 BB03, BD02, DG02, EJ03, KB02, PJ22, RH22 RK17 CA03, RB20 BA03 QL16 SL22	41, 43, 61, 72, 107, 158, 278 288 49, 257 41 230 359
Le Breton, Jean-marie Le Guyader, Loic Le Tacon, Mathieu Le Tacon, Matthieu Le, Cuong V. Le, Tuan Tu Lebecki, Kristof M.	Al04 BB03, BD02, DG02, EJ03, KB02, PJ22, RH22 RK17 CA03, RB20 BA03 QL16 SL22 RJ04	41, 43, 61, 72, 107, 158, 278 288 49, 257 41 230 359 283
Le Breton, Jean-marie Le Guyader, Loic Le Tacon, Mathieu Le Tacon, Matthieu Le, Cuong V. Le, Tuan Tu Lebecki, Kristof M. Lebreton, Jean-christophe	Al04 BB03, BD02, DG02, EJ03, KB02, PJ22, RH22 RK17 CA03, RB20 BA03 QL16 SL22 RJ04 CG03	41, 43, 61, 72, 107, 158, 278 288 49, 257 41 230 359 283 53
Le Breton, Jean-marie Le Guyader, Loic Le Tacon, Mathieu Le Tacon, Matthieu Le, Cuong V. Le, Tuan Tu Lebecki, Kristof M. Lebreton, Jean-christophe Lechevallier, Luc	Al04 BB03, BD02, DG02, EJ03, KB02, PJ22, RH22 RK17 CA03, RB20 BA03 QL16 SL22 RJ04 CG03 BD02, RH22	41, 43, 61, 72, 107, 158, 278 288 49, 257 41 230 359 283 53 43, 278
Le Breton, Jean-marie Le Guyader, Loic Le Tacon, Mathieu Le Tacon, Matthieu Le, Cuong V. Le, Tuan Tu Lebecki, Kristof M. Lebreton, Jean-christophe Lechevallier, Luc Ledue, Denis	Al04 BB03, BD02, DG02, EJ03, KB02, PJ22, RH22 RK17 CA03, RB20 BA03 QL16 SL22 RJ04 CG03 BD02, RH22 EJ03	41, 43, 61, 72, 107, 158, 278 288 49, 257 41 230 359 283 53 43, 278 72
Le Breton, Jean-marie Le Guyader, Loic Le Tacon, Mathieu Le Tacon, Mathieu Le, Cuong V. Le, Tuan Tu Lebecki, Kristof M. Lebreton, Jean-christophe Lechevallier, Luc Ledue, Denis Lee, B. W.	Al04 BB03, BD02, DG02, EJ03, KB02, PJ22, RH22 RK17 CA03, RB20 BA03 QL16 SL22 RJ04 CG03 BD02, RH22 EJ03 RA07	41, 43, 61, 72, 107, 158, 278 288 49, 257 41 230 359 283 53 43, 278 72 250
Le Breton, Jean-marie Le Guyader, Loic Le Tacon, Mathieu Le Tacon, Mathieu Le, Cuong V. Le, Tuan Tu Lebecki, Kristof M. Lebreton, Jean-christophe Lechevallier, Luc Ledue, Denis Lee, B. W. Lee, B.c.	Al04 BB03, BD02, DG02, EJ03, KB02, PJ22, RH22 RK17 CA03, RB20 BA03 QL16 SL22 RJ04 CG03 BD02, RH22 EJ03 RA07 PJ31	41, 43, 61, 72, 107, 158, 278 288 49, 257 41 230 359 283 53 43, 278 72 250 159
Le Breton, Jean-marie Le Guyader, Loic Le Tacon, Mathieu Le Tacon, Matthieu Le, Cuong V. Le, Tuan Tu Lebecki, Kristof M. Lebreton, Jean-christophe Lechevallier, Luc Ledue, Denis Lee, B. W. Lee, B. C. Lee, Bo Wha	Al04 BB03, BD02, DG02, EJ03, KB02, PJ22, RH22 RK17 CA03, RB20 BA03 QL16 SL22 RJ04 CG03 BD02, RH22 EJ03 RA07 PJ31 RA21	41, 43, 61, 72, 107, 158, 278 288 49, 257 41 230 359 283 53 43, 278 72 250 159 252
Le Breton, Jean-marie Le Guyader, Loic Le Tacon, Mathieu Le Tacon, Matthieu Le, Cuong V. Le, Tuan Tu Lebecki, Kristof M. Lebreton, Jean-christophe Lechevallier, Luc Ledue, Denis Lee, B. W. Lee, B. C. Lee, Bo Wha Lee, Bowha	Al04 BB03, BD02, DG02, EJ03, KB02, PJ22, RH22 RK17 CA03, RB20 BA03 QL16 SL22 RJ04 CG03 BD02, RH22 EJ03 RA07 PJ31 RA21 QL09	41, 43, 61, 72, 107, 158, 278 288 49, 257 41 230 359 283 53 43, 278 72 250 159 252 229
Le Breton, Jean-marie Le Guyader, Loic Le Tacon, Mathieu Le Tacon, Mathieu Le, Cuong V. Le, Tuan Tu Lebecki, Kristof M. Lebreton, Jean-christophe Lechevallier, Luc Ledue, Denis Lee, B. W. Lee, B. C. Lee, Bo Wha Lee, Bowha Lee, Bo-wha	Al04 BB03, BD02, DG02, EJ03, KB02, PJ22, RH22 RK17 CA03, RB20 BA03 QL16 SL22 RJ04 CG03 BD02, RH22 EJ03 RA07 PJ31 RA21 QL09 SO03	41, 43, 61, 72, 107, 158, 278 288 49, 257 41 230 359 283 53 43, 278 72 250 159 252 229 367
Le Breton, Jean-marie Le Guyader, Loic Le Tacon, Mathieu Le Tacon, Matthieu Le, Cuong V. Le, Tuan Tu Lebecki, Kristof M. Lebreton, Jean-christophe Lechevallier, Luc Ledue, Denis Lee, B. W. Lee, B. C. Lee, Bo Wha Lee, Bowha Lee, Bo-wha Lee, Bumsung	Al04 BB03, BD02, DG02, EJ03, KB02, PJ22, RH22 RK17 CA03, RB20 BA03 QL16 SL22 RJ04 CG03 BD02, RH22 EJ03 RA07 PJ31 RA21 QL09 SO03 JA03, PA04	41, 43, 61, 72, 107, 158, 278 288 49, 257 41 230 359 283 53 43, 278 72 250 159 252 259 367 98, 116
Le Breton, Jean-marie Le Guyader, Loic Le Tacon, Mathieu Le Tacon, Mathieu Le, Cuong V. Le, Tuan Tu Lebecki, Kristof M. Lebreton, Jean-christophe Lechevallier, Luc Ledue, Denis Lee, B. W. Lee, B. C. Lee, Bo Wha Lee, Bowha Lee, Bowha Lee, Bowha Lee, Bumsung Lee, Byung Chan	Al04 BB03, BD02, DG02, EJ03, KB02, PJ22, RH22 RK17 CA03, RB20 BA03 QL16 SL22 RJ04 CG03 BD02, RH22 EJ03 RA07 PJ31 RA21 QL09 SO03	41, 43, 61, 72, 107, 158, 278 288 49, 257 41 230 359 283 53 43, 278 72 250 159 252 229 367
Le Breton, Jean-marie Le Guyader, Loic Le Tacon, Mathieu Le Tacon, Matthieu Le, Cuong V. Le, Tuan Tu Lebecki, Kristof M. Lebreton, Jean-christophe Lechevallier, Luc Ledue, Denis Lee, B. W. Lee, B. C. Lee, Bo Wha Lee, Bowha Lee, Bo-wha Lee, Bumsung	Al04 BB03, BD02, DG02, EJ03, KB02, PJ22, RH22 RK17 CA03, RB20 BA03 QL16 SL22 RJ04 CG03 BD02, RH22 EJ03 RA07 PJ31 RA21 QL09 SO03 JA03, PA04 SK14	41, 43, 61, 72, 107, 158, 278 288 49, 257 41 230 359 283 53 43, 278 72 250 159 252 229 367 98, 116 352
Le Breton, Jean-marie Le Guyader, Loic Le Tacon, Mathieu Le Tacon, Mathieu Le, Cuong V. Le, Tuan Tu Lebecki, Kristof M. Lebreton, Jean-christophe Lechevallier, Luc Ledue, Denis Lee, B. W. Lee, B. W. Lee, B. C. Lee, Bo Wha Lee, Bowha Lee, Bowha Lee, Bumsung Lee, Byung Chan Lee, Byungchan	Al04 BB03, BD02, DG02, EJ03, KB02, PJ22, RH22 RK17 CA03, RB20 BA03 QL16 SL22 RJ04 CG03 BD02, RH22 EJ03 RA07 PJ31 RA21 QL09 SO03 JA03, PA04 SK14 PK21	41, 43, 61, 72, 107, 158, 278 288 49, 257 41 230 359 283 53 43, 278 72 250 159 252 229 367 98, 116 352 163
Le Breton, Jean-marie Le Guyader, Loic Le Tacon, Mathieu Le Tacon, Mathieu Le, Cuong V. Le, Tuan Tu Lebecki, Kristof M. Lebreton, Jean-christophe Lechevallier, Luc Ledue, Denis Lee, B. W. Lee, B. W. Lee, B. W. Lee, Bo. Wha Lee, Bowha Lee, Bowha Lee, Bowha Lee, Bumsung Lee, Byung Chan Lee, Byungchan Lee, C. H.	Al04 BB03, BD02, DG02, EJ03, KB02, PJ22, RH22 RK17 CA03, RB20 BA03 QL16 SL22 RJ04 CG03 BD02, RH22 EJ03 RA07 PJ31 RA21 QL09 SO03 JA03, PA04 SK14 PK21 QB28	41, 43, 61, 72, 107, 158, 278 288 49, 257 41 230 359 283 53 43, 278 72 250 159 252 259 367 98, 116 352 163 191
Le Breton, Jean-marie Le Guyader, Loic Le Tacon, Mathieu Le Tacon, Mathieu Le, Cuong V. Le, Tuan Tu Lebecki, Kristof M. Lebreton, Jean-christophe Lechevallier, Luc Ledue, Denis Lee, B. W. Lee, B. W. Lee, B. W. Lee, B. Wha Lee, Bo Wha Lee, Bowha Lee, Bowha Lee, Bo-wha Lee, Byung Chan Lee, Byung chan Lee, C. H. Lee, Chang Hee	Al04 BB03, BD02, DG02, EJ03, KB02, PJ22, RH22 RK17 CA03, RB20 BA03 QL16 SL22 RJ04 CG03 BD02, RH22 EJ03 RA07 PJ31 RA21 QL09 SO03 JA03, PA04 SK14 PK21 QB28 RP01	41, 43, 61, 72, 107, 158, 278 288 49, 257 41 230 359 283 53 43, 278 72 250 159 252 229 367 98, 116 352 163 191 303
Le Breton, Jean-marie Le Guyader, Loic Le Tacon, Mathieu Le Tacon, Mathieu Le, Cuong V. Le, Tuan Tu Lebecki, Kristof M. Lebreton, Jean-christophe Lechevallier, Luc Ledue, Denis Lee, B. W. Lee, B. W. Lee, B. W. Lee, B. W. Lee, B. Wha Lee, Bowha Lee, Bowha Lee, Bowha Lee, Bowha Lee, Byung Chan Lee, Byung chan Lee, C. H. Lee, Chang Hee Lee, Chang Hoon	AI04 BB03, BD02, DG02, EJ03, KB02, PJ22, RH22 RK17 CA03, RB20 BA03 QL16 SL22 RJ04 CG03 BD02, RH22 EJ03 RA07 PJ31 RA21 QL09 SO03 JA03, PA04 SK14 PK21 QB28 RP01 EH02	41, 43, 61, 72, 107, 158, 278 288 49, 257 41 230 359 283 53 43, 278 72 250 159 252 229 367 98, 116 352 163 191 303 70
Le Breton, Jean-marie Le Guyader, Loic Le Tacon, Mathieu Le Tacon, Mathieu Le, Cuong V. Le, Tuan Tu Lebecki, Kristof M. Lebreton, Jean-christophe Lechevallier, Luc Ledue, Denis Lee, B. W. Lee, B. W. Lee, B. W. Lee, B. W. Lee, B. Wha Lee, Bowha Lee, Bowha Lee, Bowha Lee, Bowha Lee, Bowha Lee, Bowha Lee, Bumsung Lee, Byung Chan Lee, C. H. Lee, Chang Hee Lee, Chang Hoon Lee, Changhee	Al04 BB03, BD02, DG02, EJ03, KB02, PJ22, RH22 RK17 CA03, RB20 BA03 QL16 SL22 RJ04 CG03 BD02, RH22 EJ03 RA07 PJ31 RA21 QL09 SO03 JA03, PA04 SK14 PK21 QB28 RP01 EH02 PJ17	41, 43, 61, 72, 107, 158, 278 288 49, 257 41 230 359 283 53 43, 278 72 250 159 252 259 367 98, 116 352 163 191 303 70 157
Le Breton, Jean-marie Le Guyader, Loic Le Tacon, Mathieu Le Tacon, Mathieu Le, Cuong V. Le, Tuan Tu Lebecki, Kristof M. Lebreton, Jean-christophe Lechevallier, Luc Ledue, Denis Lee, Bowha Lee, B. W. Lee, B. C. Lee, Bo Wha Lee, Bowha Lee, Bowha Lee, Bowha Lee, Bowha Lee, Bowha Lee, Bowha Lee, Bumsung Lee, Bumsung Lee, Byung Chan Lee, Chang Hee Lee, Chang Heo Lee, Chang Hoon Lee, Changhee Lee, Changhoon	Al04 BB03, BD02, DG02, EJ03, KB02, PJ22, RH22 RK17 CA03, RB20 BA03 QL16 SL22 RJ04 CG03 BD02, RH22 EJ03 RA07 PJ31 RA21 QL09 SO03 JA03, PA04 SK14 PK21 QB28 RP01 EH02 PJ17 QJ18	41, 43, 61, 72, 107, 158, 278 288 49, 257 41 230 359 283 53 43, 278 72 250 159 252 229 367 98, 116 352 163 191 303 70 157 223 146, 176, 348, 355 129
Le Breton, Jean-marie Le Guyader, Loic Le Tacon, Mathieu Le Tacon, Mathieu Le, Cuong V. Le, Tuan Tu Lebecki, Kristof M. Lebreton, Jean-christophe Lechevallier, Luc Ledue, Denis Lee, Boutha Lee, B. W. Lee, B. W. Lee, B. C. Lee, Bo Wha Lee, Bowha Lee, Bowha Lee, Bowha Lee, Bowha Lee, Bowha Lee, Bowha Lee, Bowha Lee, Bowha Lee, Chang Hee Lee, Chang Hee Lee, Changhee Lee, Changhoon Lee, Cheol Eui Lee, Chi-hung Lee, Chul-ho	AI04 BB03, BD02, DG02, EJ03, KB02, PJ22, RH22 RK17 CA03, RB20 BA03 QL16 SL22 RJ04 CG03 BD02, RH22 EJ03 RA07 PJ31 RA21 QL09 SO03 JA03, PA04 SK14 PK21 QL09 SO03 JA03, PA04 SK14 PK21 QB28 RP01 EH02 PJ17 QJ18 PG25, PN25, SI14, SK28 PC18 QB02, QB15, QB29	41, 43, 61, 72, 107, 158, 278 288 49, 257 41 230 359 283 53 43, 278 72 250 159 252 229 367 98, 116 352 163 191 303 70 157 223 146, 176, 348, 355
Le Breton, Jean-marie Le Guyader, Loic Le Tacon, Mathieu Le Tacon, Mathieu Le, Cuong V. Le, Tuan Tu Lebecki, Kristof M. Lebreton, Jean-christophe Lechevallier, Luc Ledue, Denis Lee, Bo. W. Lee, B. W. Lee, B. W. Lee, B. W. Lee, Bo. Wha Lee, Bo Wha Lee, Bowha Lee, Bowha Lee, Bowha Lee, Bowha Lee, Bowha Lee, Bumsung Lee, Bumsung Lee, Bumsung Lee, Bumsung Lee, Bumsung Lee, Chang Hee Lee, Chang Heo Lee, Chang Heo Lee, Changhoon Lee, Changhoon Lee, Cheol Eui Lee, Chi-hung	AI04 BB03, BD02, DG02, EJ03, KB02, PJ22, RH22 RK17 CA03, RB20 BA03 QL16 SL22 RJ04 CG03 BD02, RH22 EJ03 RA07 PJ31 RA21 QL09 SO03 JA03, PA04 SK14 PK21 QL99 SO03 JA03, PA04 SK14 PK21 QB28 RP01 EH02 PJ17 QJ18 PG25, PN25, SI14, SK28 PC18	41, 43, 61, 72, 107, 158, 278 288 49, 257 41 230 359 283 53 43, 278 72 250 159 252 229 367 98, 116 352 163 191 303 70 157 223 146, 176, 348, 355 129

## AUTHOR INDEX

Lee, Donghae Lee, Dong-hae Lee, Dong-il Lee, Eunhee Lee, Eunsook Lee, Geunsik Lee, H. J. Lee, H. Y. Lee, H.k. Lee, Ha Na Lee, Haigun Lee, Hak Bong Lee, Hak-bong Lee, Hakjoon Lee, Hana Lee, Hanju Lee, Han-oh Lee, Hee-ju Lee, Heejung Lee, Hsin-han Lee, Hu - Jong Lee, Hunju Lee, Hunpyo Lee, Hun-sung Lee, Hyeong Jun Lee, Hyun Joon Lee, Hyun Jung Lee, Hyun Sook Lee, Hyungjun Lee, Hyungyu Lee, Hyun-woo Lee, In Kyu Lee, J H Lee, J. Lee, J. H. Lee, J. M. Lee, J. Y. Lee, J.-s. Lee, Jae II Lee, Jae-chul Lee, Jae-hyun Lee, Jaejin Lee, Jae-jun Lee, Jee Yong Lee, Jehyun Lee, Jeong Soo Lee, Jeonghyeon Lee, Jeong-o Lee, Jhinhwan Lee, Ji Hye Lee, Jin Bae Lee, Jin Seok Lee, Jin Yong Lee, Jinho

Lee, Jinju

Lee, Jinyong

Lee, Jinyoung

QH15 215 QG16 210 PM19, PM21 170 RN24 300 PG11, PH14, QA04 144, 149, 183 217, 257 QH26, QH27, RB22 RQ08 308 161, 248 PK06, QP17 123 PB19 PJ24 158 PB16, PB17, PB18, PB35, 122, 123, 125, 126 PB36, PB37, PB38, PB39, PB40 PA01 116 CI03, PA03 55, 116 PN09 173 PJ16, PK10 157, 162 SA09 317 FB01 73 231 QL18 RH08 276 QK22, SK30 227, 355 RB14 256 171 PM23 339, 343 SF14, SG22 290, 361 RL03, SM01 RC08 260 RQ18 310 JC03 100 PO06, PO09, PO10 178 QK21, RC15 227, 261 RL22, SL19 293, 359 EE02, EE05, HP21, KE03, 67, 68, 30, 110, 160 PJ33 PI14 153 QB25 191 QB31 192 SA06 317 QB21, SM11 190, 362 PB03 120 PH10, RA04 148, 249 RH20, SK19, SK27 278, 353, 354 KE03, SI13 110, 348 179 PO15, PO16 JH04, KH04 104, 112 RR14, SN04 313, 364 211 QG20 RJ15, SJ07 285, 350 QM18, QM20 234 QK03 224 SD12, SD14 332 CA02, RQ11 48, 308 SA05 317 357 SL11 SD14 332 QG17 211 CA02 48 RL17 292 PO09 178 PO10 178

Author Index	
Author	

Lee, Jongseok	PB17, PB36, PB37, PB39, RG10	122, 1
Lee, Joo In	RK19	289
Lee, Ju	RR14, SN02, SN04	313, 3
Lee, Juho	PM17	170
Lee, Jung Won	QN09	237
Lee, Junggoo	JH04, PM24	104, 1
Lee, Jung-goo	KH04, RN06, SL29	112, 2
Lee, Junghyon	РК04	161
Lee, Junhyun	IA03	91
Lee, K. B.	RQ08	308
Lee, Ki Doek	SN04	364
Lee, Ki-doek	RR14, SN02	313, 3
Lee, Kiejin	SA09	317
Lee, Ki-seung	RL26	294
Lee, Ki-suk	DE05, RJ14, RJ16, RJ18, SH14	60, 28
Lee, Kiyoung	PN23	175
Lee, Kwan-woo	PF06, SO05, SO06	140, 3
Lee, Kyeong-dong	RK09, RK22, SI04	287,2
Lee, Kyu Won	PG25, PN25, SI14, SK28	146, 1
Lee, Kyujoon	PC12, PK17, PN19, QO10, RC07, SD10, SD11	128, 1 260, 3
Lee, Kyung Jin	RL26	294
Lee, Kyungdong	QM23	235
Lee, Kyung-jin	CE04, EE02, EE05, HF05,	
	HP21, PJ33, RJ19, RL06, RL14, SI09	160, 2 347
Lee, M. K.	RQ08	308
Lee, N.	AA04	33
Lee, N. H.	PB03	120
Lee, Nam Hoon	GA04, RB27	76, 25
Lee, Nam-kyu	RR15	313
Lee, Nyun Jong	PJ17, QP18	157, 2
Lee, Paul	PN03, QL02	172, 2
Lee, S H.	PH03	147
Lee, S. J.	RQ08, SJ06, SL08	308, 3
Lee, S.s.	RK18	288
Lee, Sang Sun	RP23	307
Lee, Sangho	RL30	294
Lee, Sanghoon	BE01, PN08, PN09, QL08, QM04	43, 17
Lee, Sang-hyuk	RP03, SH10	303, 3
Lee, Sanghyun	PG24, PI10	146, 1
Lee, Sangji-suk	PL11	165
Lee, Sang-suk	BH02, RR15	46, 31
Lee, Sang-wha	PO04	177
Lee, Sangwon	RH08	276
Lee, Sangyeop	QM04	232
Lee, Sang-yun	PM16, PM19, PM21	170
Lee, Seok-hee	RL30	294
Lee, Seongjae	RL21	293
Lee, Seong-joo	RQ18, RQ19, RQ20	310
Lee, Seong-rae	RO01, RO02	300, 3
Lee, Seongsu	HJ05, PI10, QA24, RA20, RA22	90, 15 253
Lee, Seo-won	S109	347
Lee, Seung Goo	QN22	240
Lee, Seung-beck	QP19, RL22, SL19	248, 2

B39,	122, 126, 273	Lee, Seungho
		Lee, Seunghun
	289	Lee, Seung-hun
	313, 364	Lee, Seungjun
	170	Lee, Seungkyo
	237	Lee, Shang-fan
	104, 171	Lee, Soogil
	112, 297, 360	Lee, Soonchil
	161	Lee, Soon-gul
	91	Lee, Sujin
	308	Lee, Suk Ho
	364	Lee, Sung Hoon
	313, 364	Lee, Sungbin
	317	Lee, Sung-hyun
	294	Lee, Sungsu
118,	60, 284, 285, 345	Lee, Tae Ho
		Lee, Tae Young
	175	Lee, Taehoon
	140, 367	Lee, W. H.
	287, 289, 346	Lee, Wei-cheng
K28	146, 176, 348, 355	Lee, Woncheol
	128, 163, 175, 242,	Lee, Woo
	260, 331	Lee, Y. P.
		Lee, Yen-chen
	294	Lee, Ying-chen
	235	Lee, Yong Ho
IF05,	52, 67, 68, 87, 30,	Lee, Yong Woo
_06,	160, 285, 290, 292,	Lee, Yong-ho
	347	Lee, Young S.
	308	Lee, Yun-hi
	33	Lees, Martin R
	120	Lefevbre, Williams
	76, 258	Lefevre, C.
	313	Lefevre, Christophe
	157, 248	Legarra, Estibaliz
	172, 228	Legut, Dominik
	147	Leist, Jeannis Leitenstorfer, Alfred
	308, 349, 357	Leithe-jasper, Andreas
	288	
	307	Lejay, Pascal Lemaitre, Aristide
	294	Lemee-cailleau, Marie-hele
	43, 173, 229, 232	Lemmens, P.
	303, 344	Lemmens, Peter
	146, 152	Lemus, Jose
	146, 152	Lenertz, Marc
	46, 313	Lengyel, Edith
	177	Leon, Gladys
	276	Leon, L. M.
	232	Leonyuk, Nicolay I.
	170	Lepadatu, Serban
	294	Lesne, Edouard
	293	Lesnikov, Valeriy
	310	Letellier, Florent
	300, 301	Leung, Chi Wah
A20,	90, 152, 186, 252,	Levchenko, A. V.
0,	253	Levchenko, Anna Vasilievn
	347	Lewinska, S.
	240	Li, Anhua
	248, 293, 359	Li, Bin
	., .,	

	RR21	314
	PH22, PN20, PN21, QL18	150, 175, 231
	AC02	34
	RL17	292
	PK04, QK03	161, 224
	PC14, RK07, RL18	129, 287, 292
	RL30	294
	BF05, JG05, PA09, PH18	45, 104, 117, 149
	PB21, RB27, SC14	123, 258, 327
	SA09	317
	JB02	99
	PB21, RB27	123, 258
	JD03, PL08	101, 165
	SJ05	349
	QI15	220
	SO21	370
	RL28	294
	RN10	298
	RP10	304
	GA01	75
	JH04	104
	SO21	370
	QG27, RA07, SL08, SL09	212, 250, 357
	SB25	324
	PL03	164
	RQ17	309
	PN28	176
	RQ18, RQ19, RQ20	310
	QE27 RL31	203 295
	BF03	45
	KB02	107
	SA06	317
	SA05	317
	QH16	215
	CH03, QH19	54, 216
	AA05	33
I	HA02	83
as	RD12	264
	SG05	340
	CG03	53
rie-helene	SB08	321
	PI09	152
	QE27	203
	Q026	245
	PIO4	151
	AB02	34
	QB18, QB19	190
	CB02	49
	PA20	119
	Q024	244
	HE04	86
	SG17	342
	BD02	43
	DH04	62 105 275
asilievna	QC22, RH07 PE09	195, 275 137
	PH06, PM05	137
	RO04	301
	QB01	187
	200.	

Li, Bodong	PM27	171
Li, Carissa H.	KB03	107
Li, D. X.	QJ10	222
Li, Fashen	QN21, QN23	239, 240
Li, Feng	QK01	224
Li, Gang	JA02	98
Li, Hang	IE01	93
Li, J.	PL02, SM05	164, 361
Li, Jiajie	RO04	301
Li, Jian	RR19, RR24, RR25	314, 315
Li, Jichao C	KI02	113
Li, Li	PN05, QL04	172, 228
Li, Liang	PL06	164
Li, Lingwei	QJ06	221
Li, Long-jie	QP17	248
Li, Peisen	SA19, SA20	319
Li, Q.	GA01	75
Li, Qing	PO01	177
Li, Tian	BE03	44
Li, Tong	DD03	58
Li, W F	DD06	59
Li, Wanfeng	CB05, EJ04	49, 72
Li, Wei	PK18, QN21, RO04, RO07, SL12, SN11, SN12	163, 239, 301, 358, 366
Li Maimin		
Li, Weimin Li, Wen-hsien	QP22 PC18, PO25, QA20	249 129, 181, 185
Li, Xianhong	RL21	293
Li, Yanfeng	QN21	239
Li, Yingjie	PN07, QC01	173, 192
Li, Yong Hui	RA26	253
Li, Yuan	CA03, JC04	49, 100
Li, Yufan	JF05	103
Li, Yuke	AB04, IJ04	34, 98
Li, Z	QB11	188
Li, Zhenghua	RP06	304
Liang, Ruixing	BA01	40
Liang, Shiheng	QK02	224
Liao, Shu-hsien	QM27	236
Liao, Yen Fa	QC12	194
Liao, Yen-fa	QC09	193
Liborio, Leandro	RF15	271
Lichtenegger, Thomas	RG15	274
Lichtenwalner, Daniel J.	AH04	38
Lileev, Alexey	QJ13, SN09	223, 365
Lim, Byunghwa	PO19	180
Lim, Hyein	RL17	292
Lim, Jaein	Q019	243
Lim, Jung Tae	PA22, PI13	119, 152
Lim, Kwon Taek	RN05	297
Lim, Pang Boey	QP15, QP16	247, 248
Lim, Sang Ho	PL16, QK12	166, 226
Lim, Sang-ho	RQ23	310
Lim, Sanghyun	PO11	179
Lim, Seon-ho	PO14	179
Lim, Yong-sik	PN31	177
Lim, Yun-seog Lima Sharma, Ana	PM19, PM21	170 242
Lima Sharma, Ana Lin, C T	Q009 0820	242 190
Lin, C. M.	QB20 RR01	311
Lin, C. T.	GA03	75
C. II	2.100	

Lin, C. Y.	SB04	320
Lin, Chengtian	IJ02, RB20	97, 257
Lin, Chun-liang	SM06	361
Lin, H.	SD04	330
Lin, Hsin	HI02	89
Lin, Hsiu-hau	RE11, SB25	268, 324
Lin, Jianhua	PH02	147
Lin, Ko-wei	PL18	166
Lin, Kurt	PO13	179
Lin, Lu-kuei	PC14	129
Lin, Minn-tsong	AD01	35
Lin, Pang	RL18	292
Lin, Ruei-lin	QN24	240
Lin, T.	GF01	79
Lin, Wei-yan	QM09	233
Lin, Z. W.	GA01	75
Ling, D. C.	PC13	128
Liou, Sy-hwang	RO07	301
Liou, Yung	QM09	233
Lipinski, Stanislaw	PE06	137
Lisboa-filho, Paulo Noronha	QA06	183
Lisenkov, Sergey	QA27	187
Lisowski, W.	DD02	58
Lisyansky, Alexander A.	HD05	86
Litterst, Jochen	JG06, RN15	104, 299
Litvinova, Tatyana	PE08	137
Liu, C.	SD04	330
Liu, C. Y.	PC13	128
Liu, Chang	HI02	89
Liu, Chieh-wen	QM27	236
Liu, Chi-hsin	PL18	166
Liu, Chunli	QL09	229
Liu, Dongping	QK02	224
Liu, H. F.	RP02	303
Liu, H. L.	PC13	128
Liu, Hsin-tzu	RR01	311
Liu, J. Ping	EJ02	71
Liu, Jun-ming	PA05	116
Liu, Liying	SM17	363
Liu, Mei	QB01	187
Liu, Qiang	PI17	153
Liu, Qiangchun	SA01	316
Liu, Tao	PK18	163
Liu, Tijiang	RB12	256
Liu, Wei	PN04	172
Liu, Weiqiang Liu, X.	RO15, SN01 PN08, PN09, QL08,	303, 364 173, 229, 232
Liu, A.	QM04	175, 229, 252
Liu, Xi	SL02	356
Liu, Xiaoli	SL06	357
Liu, Xiaoxi	PK07, QP08, RL05	161, 246, 290
Liu, Xuerong	DC03	58
Liu, Yi	JB03	99
Liu, Ying Dan	QN03	236
Liu, Zuli	PO01	177
Liu, Zu-li	QK18	227
Liuqiang, Wen	Q015	243
Lobanova, Inna	SG07	340
Locatelli, Andrea	SH06	344
Locatelli, Nicolas	ID03	93

Lockwood David	0612
Lockwood, David	QG12
Lodya, Lonzeche J. Augustin	QH22
Loew, U.	QE23
Loewenhaupt, Michael	PH04, QC24
Logg, Peter	PC02
Lohneysen, Hilbert V.	EA01, EG02
Loidl, Alois	PA23, RH24
Lombardo, Pierre	CD03
Long, Jerome	RN15
Long, Phan The	AJ06
Lonzarich, Gil	BA01
Lopera, Wilson	SA12
Lopes, Rovan Fernandes	PB09, RB09
Lopez-diaz, Luis	SH16
Lopez-quintela, M. A.	CB02
Lora-serrano, Raimundo	PI19, SE15
Lord, James S.	PO17
Lorenz, Thomas	RI12
Lorenz, W. E. A.	QF18
Lorenzer, K A.	CD01, SF10
Lortz, Rolf	DH04
Loshkareva, Natalya	SL17
Lott, Dieter	DG01, QM02, QM08
Louahadj, Lamis	AI04
Louca, Despina	BF02
Lowe, Konrad	KH01
Lu, C. R.	PC13
Lu, Chun-i	AD01
Lu, Hantao	EB03
Lu, Jun	PA21, QK23, RA02,
	RQ01
Lu, K. T.	QB21
Lu, Kueih-tzu	SM11
Lu, Wei	AJ03, JI06
Lu, Xin	FB01, SE02
Lu, Yunhao	SK34
Lubhelwane, Siyanda	QH22
Lucas, Irene	SM16
Ludwig, Frank	AH01
Lue, C. S.	PE02, PH01
Lue, Chin Shan	RP02
Luetkens, Hubertus	AB02, PI05, PI17, QB22,
	QJ03
Luis, F.	BC01
Lukasiewicz, Malgorzata I	DD02
Luke, G. M.	GF01
Lummen, Tom T. A.	HG04, PA24
Lumpkin, Greg	PO12
Lunov, Oleg	AH02
Luo, Hui Quian	RB10
Luo, Huigian	RB11
Luo, Yongkang	AB04, IJ04, RB13
Luong, Dien Xuan	RN19
Lupi, Stefano	SA08
Luque, F.j.	DD02
Luque, r.j. Lushchyk, P	RP20
Lushchyk, P Lutsev, Leonid	QM08
Luzon, Javier	GF04
Luzon, Javier Lv, Bin	
	JA02, PN05, QL04 IJ04
Lv, Chen	UV4

210	Lym, Dae Young	SG21
216	Lynn, Jeffrey W.	QA20
203		
147, 196	Μ	
127	M, Aparnadevi	PI01
64, 69	M, Manivelraja	QD04
119, 278	M, Shalini	RN16
51	M. Tachiki, M.	KJ01
299	M., Ramesh	BJ05
40	Ma, Li	JI05
40	Ma, Long	QB07
318	Ma, Qinli	CH02
121, 255	Ma, Xiao-ping	SH10
345	Maan, Jan Kees	DH02, SA10
49	Maca, Frantisek	BE05
153, 335	Macedo, Claudio Andrade Macedo, Marcelo	PF13
180 280	Machida, Kenji	QL01 SK21
206	Machida, Yo	PC17, PC28, F
50,338	Mackenzie, Andrew	EA04
62	Maclaughlin, Douglas E.	CD04
358	Madami, Marco	GG02
61, 231, 232	Maeda, Youichi	QD02
39	Maegawa, Satoru	BC04, QE22
44	Maehira, Takahiro	QC14, QC15
112	Maekawa, Sadamichi	FA03, JE02, JF
128		KF03, PJ07, PJ
35		RF07, RL01
65	Maeno, Yoshiteru	PC20
119, 227, 249, 307	Maeta, Takashi	PB27
	Maeter, Hemke	PI05, QJ03
190	Magalhaes, S. G.	QD10
362	Magalhaes, Sergio	PE01
40, 105	Magen, Cesar	DG05, GG05,
73, 333		RN09, SM16
356	Magishi, Ko-ichi	QC05
216	Magnavita, E. Thizay	Q003
363	Magni, Alessandro	HH01
38	Mahadevan, Priya	HC01, RF12
136, 147	Mahajan, Mangesh B. Mahdavifar, Saeed	SL13 RG16
303	Mahendiran, Ramanathan	PI01, QL10, R/
34, 151, 153, 190, 221	Mahlubi, Zwelithini Melford	RD23
42	Mahmoodi, M.	PO27, PO29, F
58	Mahmoodi, Mostafa	PO28
79	Maignan, Antoine	BB03, JB05, R
88, 119	Maikov, Vladislav	QL07
179	Maisuradze, Alexander	CD04
38	Maiti, Kalobaran	PE14, SC06
255	Maitre, Adeline	EJ03
255	Majetich, Sara	GG01
34, 98, 256	Majewski, J. A.	BE02
299	Maji, Bibekananda	Q004
317	Majumdar, Pinaki	RG12
58	Majumdar, S	QL14, SO14, F
306	Majumder, M	QM24
232	Majumder, Mayukh	RB15
79	Makarov, Denys	SJ07
98, 172, 228	Maki, Makoto Makio, Kurisu	PB04
98	Makio, Kurisu	SM07

ae Young	SG21	343
effrey W.	QA20	185
arnadevi	PI01	150
nivelraja	QD04	197
lini	RN16	299
hiki, M.	KJ01	113
mesh	BJ05	48
	JI05	105
ng	QB07	188
nli	CH02	54
ao-ping	SH10	344
Jan Kees	DH02, SA10	62, 317
Frantisek	BE05	44
lo, Claudio Andrade	PF13	141
lo, Marcelo	QL01	228
da, Kenji	SK21	353
da, Kenji da, Yo	PC17, PC28, PD16	129, 131, 134
	EA04	
nzie, Andrew		64
ughlin, Douglas E.	CD04	51
ni, Marco	GG02	80
i, Youichi	QD02	196
iwa, Satoru	BC04, QE22	42, 203
ra, Takahiro	QC14, QC15	194
wa, Sadamichi	FA03, JE02, JF02, JF03,	72, 102, 111, 155,
	KF03, PJ07, PJ34, PP02,	160, 28, 270, 290
	RF07, RL01	
o, Yoshiteru	PC20	130
, Takashi	PB27	124
	PI05, QJ03	151, 221
haes, S. G.	QD10	198
haes, Sergio	PE01	136
n, Cesar	DG05, GG05, RN08, RN09, SM16	61, 80, 298, 363
ni, Ko-ichi	QC05	193
avita, E. Thizay	Q003	241
, Alessandro	HH01	88
levan, Priya	HC01, RF12	84, 270
an, Mangesh B.	SL13	358
avifar, Saeed	RG16	274
idiran, Ramanathan	PI01, QL10, RA08	150, 229, 250
bi, Zwelithini Melford	RD23	265
oodi, M.	PO27, PO29, PO31	181, 182
oodi, Mostafa	PO28	181
an, Antoine	BB03, JB05, RA11, RH22	41, 100, 251, 278
v, Vladislav	QL07	229
adze, Alexander	CD04	51
Kalobaran	PE14, SC06	138, 326
, Adeline	EJ03	72
ch, Sara	GG01	80
iski, J. A.	BE02	44
ibekananda	Q004	241
ndar, Pinaki ndar, S	RG12	273
idar, S	QL14, SO14, PL07, RH02	165, 230, 275, 369
ider, M	QM24	235
ider, Mayukh	RB15	256
ov, Denys	SJ07	350
Makoto	PB04	120
, Kurisu	SM07	362

Makiyama, Shun	SG09	341
Makridis, Sofoklis	SN15	366
Maksimova, Ksenia	SK26	354
Maleyev, Sergey V.	DF04	60
Malikov, Shavkat	RR28	315
Malinowski, Gregory	SH02	343
Maljuk, Andrey	U02	97
Mallia, Giuseppe	RF15	271
Malone, Liam	JC02	100
Mal'tsev, Viktor V.	PA20	119
Malyuk, Andrey	PA10	117
Mamiya, Hiroaki	JJ05, RI07	106, 280
Manabe, N.	SL14	358
Manaf, Azwar	BJ03, QN16, RA15	48, 239, 251
Manago, Takashi Manaka, Hirotaka	PG06, RK21, RP12	143, 289, 305
Manami, Onizawa	QE09, QH02, RI06 OD14	200, 213, 279 198
Manchon, Aurelien	CG05, EE03, IE01, PJ06,	53, 67, 93, 155, 156,
Maricion, Autelien	PJ13, PJ29, PJ34, QL17,	159, 160, 230, 285,
	RJ19, SK12	352
Mandal, Ruma	HD04	85
Manivannan, Nallayian	PB23	123
Maniwa, Yutaka	SO515	369
Mankey, Gary J.	QM02	231
Manna, P K	QP05	246
Manna, P. K.	RN04	297
Manna, Rudra Sekhar	QJ11, RI13	222, 281
Manni, Soham	SG13	341
Manske, Dirk	CA03, KJ03	49, 114
Mansson, M.	PC29, PI15	131, 153
Mansson, Martin	HA03, PI06, PO17, QF20,	83, 151, 180, 207,
	SE08	334
Mantovan, Roberto	SK26	354
Manzoor, Sadia	SA15	318
Mao, Zhiqiang	RB12	256
Maple, M. Brian	BB01	41
Marcatoma, Justiniano	PL09	165
Marchetti, Gionni	QK09	225
Marcin, Jozef	IB01, QO17	91, 243
Margaris, George Mari, Kenichiro	CB03 SK18	49
Marin, Lorena	SM16	353 363
Maritato, Luigi	SA08	317
Markiewicz, R. S.	HI02	89
Markiewicz, Robert S.	BA04	41
Marques-ferreira, Pablo	PI19	153
Marquina, Clara	RN08, RN09	298
Marrows, Chris	Q024	244
Marrows, Christopher	KE01	110
Marsilius, Mie	PM04	168
Martens, Michael	FD02, RJ05, RP04	74, 283, 303
Martens, Stephan	RP04	303
Marti, X.	IF01	94
Marti, Xavier	BE05	44
Martin, Christine	JB05, RA11	100, 251
Martin, I.	FB01	73
Martin, Jose Ignacio	GG03	80
Martinez, Benjamin	IH01, SO18	95, 369
Martinez, Eduardo	SH16	345
Martinez-perez, M. J.	BC01	42

## AUTHOR INDEX

Martin-rodriguez, Damian Martins, Maximiliano Delany Marty, Alain Masek, Jan

Maskova, Silvie Masoudpanah, S. M. Massboeuf, Aurelien Masthoff, Ingke Christine Masti, S A Masuda, Keisuke Masuda, Ryo Masuda, Takatsugu Masui, T. Masumoto, Hiroshi Mathai, Cijy Mathon, Olivier Mathur, Neil D Matiks, Y. Matoba, Masanori Matsubara, Takeshi Matsubayashi, K. Matsubayashi, Kazuyuki

Matsuda, Kazuyuki Matsuda, Ken-ichi Matsuda, M. Matsuda, Masaaki Matsuda, Ryotaro Matsuda, Saori Matsuda, Tatsuma Matsuda, Tatsuma D Matsuda, Tatsuma D. Matsuda, Tatsuma Daruma Matsuda, Yasuhiro H. Matsuda, Yuji Matsudaira, Ken-ichiro Matsueda, Hiroaki Matsuhira, Kazuyuki Matsui, Fumihiko Matsui, Kazuki Matsuishi, S. Matsuishi, Satoru Matsukawa, Michiaki Matsukura, F. Matsukura, Fumihiro Matsumoto, Akiyo Matsumoto, K. T. Matsumoto, Kazuyuki Matsumoto, Keisuke Matsumoto, Keisuke T Matsumoto, Keisuke T. Matsumoto, Masashige Matsumoto, Naoki Matsumoto, Rie Matsumoto, Sayaka Matsumoto, Shinji

Matsumoto, Shoutaro

QH16 215 284 RJ12 EI04, GE01, QK15, 71, 78, 226, 293, QK16, QK17, RL24, SL24 360 BE05, SK22 44, 354 QI19, RE09 220, 267 QM28 236 344 SH06 JG06 104 238 ON14 SB23 324 QM16 234 QA02 182 HA01 83 JI01 105 CG04 53 HG04, RH23 88, 278 AE03 36 GA03 75 QC07, QC08, SE09 193, 334 PC17 129 214, 264 QH09, RD15 PD14, QD17, QI18, 134, 199, 220, 239, 263, 265, 267, 329, QN20, RD06, RD22, RE04, SC24, SE01, SE11 333, 335 SO515 369 SK06 351 PH03 147 HJ03 90 HC03, QF25 85, 208 QJ08 222 JC02 100 IG04, PD09, PD23, RD25, 95, 133, 135, 266, SB15 322 RD01, RD11, RE10, SF03 262, 263, 268, 337 HP11 30 IG03 95 IA01, QC04, SB27 90, 192, 324 PH23 150 EB03 65 CC05, QE13 50, 201 SK13 352 PD10 133 QB23 190 PC23, QB16 130, 189 PI22 154 BE02 44 QP14 247 278 RH21 EF01 68 RI01 279 SG15 342 RD04 262 SG12 341 SB17 323 QM10 233 HF01, IE04 87, 94 PM02 167 SN03 364 PN24 176

Matsumoto, Tatsuya	QM29	236
Matsumoto, Yoshihiro	BH05, BH06	46, 47
Matsumoto, Yosuke	CD02, QD01	51, 196
Matsumoto, Yuji	PD09, PD23	133, 135
Matsumura, Masahiro	QD06, RD10	197, 263
Matsumura, Takeshi	QI10	219
Matsunaga, Takahiro	KC06	109
Matsunami, Masaharu	AD04, QD18	35, 199
Matsuo, A.	Q103	218
Matsuo, Akira	HC03, PG17, QE16, QF11, QF14, QF17, RI15, RQ05	85, 145, 202, 205, 206, 281, 307
Matsuo, Mamoru	KF03	111
Matsuoka, Eiichi	PE11, PE13, QF04, QI10, QI16, RD01	137, 138, 204, 219, 220, 262
Matsuoka, Yuki	QH20	216
Matsushima, Yasushi	SL03	356
Matsushita, Takuya	Q107	218
Matsushita, Tomohiro	SK13	352
Matsuura, Keisuke	PG17, QF14	145, 206
Matsuura, Masato	PB30, PB32	125
Matsuyama, Hideo	RP14	305
Matsuyama, Kimihide	QM29, QP09, QP10, QP11, SJ03	236, 246, 247, 349
Matsuzaka, S.	CG02	53
Mattauch, Stefan	DG01	61
Mattenberger, K.	PI15	153
Mattenberger, Kurt	RB21	257
Matusiak, Marcin	AB03, SF12	34, 339
Matuura, Masato	PB29	124
Mayoral, Alvaro	RN09	298
Mayoral, Isabel	Q016	243
Mayr, Franz	PA23, RH24	119, 278
Maziarz, W.	PM05	168
Mazidian, Bayan	JA05	99
Maziewski, Andrzej	AD05	
Mazrouei Sebdani, S.	PO27, PO29, PO30, PO31	181, 182
Mazrouei Sebdani, Saeed	PO28	181
Mazumdar, Chandan	RB15	256
Mazzoli, Claudio	BF03, PA15	45, 118
Mc Gill, S.	AA04	33
Mc.donald, Ross	RB21	257
Mccollam, Alix	DH02, SA10	62, 317
Mccord, Jeffrey	SH04	343
Mcdonald, Ross	GC03	77
Mcdonald, Ross D	QB13	189
Mcgee, Fintan	SH09	344
Mcintyre, Garry	IJ03, PO25	97, 181
Mcmichael, Robert D	EG01	69
Mcmorrow, D F	RI05	279
Mcmorrow, D. F.	QN01	236
Mcmorrow, Des	QG14	210
Mcmorrow, Desmond	AF01	36
Md Nor, Anis Faridah	QM29	236
Measson, Marie-aude	PD17	134
Medarde, Marisa	PO17	180
Medicherla, V R R	PE14	138
Medjanik, K Madanas Farraira, Laticia	RP20	306
Medonca Ferreira, Leticie	SE15	335

	2546	
Medvedev, M. V.	RF16	271
Mehrafarin, Mohammad	SB07	321
Mei, Li-then	QN25	240
Meier, Guido	FD02, HG02, RJ05, RP04, SH04	74, 88, 283, 303, 343
Meier, Johan G.	RN09	298
Meindl, Manfred	AH01	38
Meisel, Mark W.	KB03	107
Meissner, Thomas	RQ09	308
Melander, Emil	AF05	37
Melisek, T	RB08	255
Melkov, Gennady	JG02	103
Melnychenko, Nataliya	SC01	325
Meloche, Eric	QG12	210
Mena, Mattia	QG14	210
Mendonca, Ana P. Aguiar De	PB09, RB09	121, 255
Meng, Kangkang	QK23	227
Mengotti, Elena	RK17	288
Mentes, Onur	SH06	344
Meny, C.	SA06	317
Meny, Christian	AJ05, SA05	40, 317
Menzel, Dirk	QM08, SM02	232, 361
Menzel, Matthias	1H03	96
Merazzo, Karla J.	GG04	80
Merchant, P.	QE19	202
Merkt, Ulrich	RP04	303
Mertig, Ingrid	EE01	67
Merlig, Inghu Meshcheriakova, Olga	QN19	239
Mesler, Brooke	RJ20	285
		126
Mesot, J.	PC01	
Mesot, Joel	IC02, QF20	92, 207
Metaxas, Peter	IEO4	94
Methfessel, Torsten	AG04	38
Metoki, Naoto	IG04, JJ05, PH09, QJ09	95, 106, 148, 222
Mettus, Denis	CH01	54
Meven, Martin	PA14, RP22	118, 306
Meziere, Cecile	SB20	323
Mi, J.I.	1102	97
Mibu, Ko	QM16	234
Michel, Ann Kathrin	RN15	299
Michel, Anny	QP18	248
Michihiro, Yoshitaka	QF24	207
Michimura, Shinji	IG03, RH03	95, 275
Michimura, Sinji	QI10	219
Michiue, Yuichi	RH06	275
Michlmayr, Thomas	RK24	289
Michor, Herwig	IA02, SC01	90, 325
Micklitz, Hans	RB10, RB11	255
Miclea, C.	SA11, SC03	318, 326
Miclea, C. F.	SA11, SC03	318, 326
Miclea, C. T.	SA11	318
Middey, S.	RF12	270
Mielke, Chuck	BA01, GC03	40, 77
Mignot, Jean - Michel	RD14	264
Mihashenok, Natalia	PI11	152
Miike, Kazunari	QL03, RN02	228, 297
Mijnarends, Peter E.	BA04	41
Mikhaylovskiy, Rostislav V.	GD04, QO05, RK20	78, 241, 289
Miki, Hiroyuki	DJ04, QG18	63, 211
Millan, Angel	KB06	108

Milyaev, Michael A.	SK01	350
Milyaev, Mikhail	SK02	350
Mimura, Kojiro	RD20	265
Min, B. I.	PE17, PG11, RG02, SC12	138, 144, 272, 327
Min, Byoung Chul	QM18	234
Min, Byoung-chul	CE04, RK09, RL03, RL06, RL14, SL26	52, 287, 290, 292, 360
Min, Byung Chul	RL26	294
Min, C	BG01	45
Min, Gyeong Im	CF03	52
Min, Jin Hong	QN15	238
Min, K. J.	PJ24	158
Minh, Nguyen Van	PH14	149
Minnullin, Arthur	QA26	186
Mino, Michinobu	RK23	289
Minola, Matteo	BA03	41
Miotkowski, Ireneusz	HI02	89
Mirebeau, Isabelle	AA03	33
Miron, Mihai	SH05	344
Misawa, Takahiro	PE18, QB10	139, 188
Misek, Martin	BB04, QJ14	42, 223
Misiorny, Maciej	DB03	57
Misoka, Keita	RH11	276
Misu, Tekeshi	KC02	108
Misuraca, Jennifer	QK23	227
Mita, Yoshimi	SG20	342
Mitani, S	SK29	355
Mitani, Seiji	EH01, IF03, IF05, SK11, SK31	69, 94, 352, 355
Mitchell, J N	GC03	77
Mito, Masaki	KB06, PB25, PG09, PG10, QG19, RN03, SA03	108, 124, 143, 144, 211, 297, 316
Mito, Takeshi	QD12, QD15, RD24	198, 266
Mitrelias, Thanos	BH01	46
Mitsen, Kirill V.	QC19	195
Mitsuda, Akihiro	IG05, QD03, QD15, SB14	95, 196, 198, 322
Mitsuda, Setsuo	PH05	147
Mitsui, Takaya	QM16	234
Mitsui, Yoshifuru	GJ02, QH08	82, 214
Mitsumata, Chiharu	KH02, SJ06	112, 349
Mitsumoto, Keisuke	RB05	254
Mitsunaga, Daisuke	QH08, QN20	214, 239
Mitsuzuka, Yoshio	RK04	286
Miura, Daisuke	RL09, SI06	291, 346
Miura, Hideo	RP11	305
Miura, Katsuya	QP14	247
Miura, Tomohiro	PB29	124
Miura, Tomoya	PG04	143
Miura, Yasunao	PD13	134
Miura, Yoko	QE09, RI06	200, 279
Miura, Yoshio	IF02	94
Miwa, Shinji	KF01, PK20, RL11, SI05, SK10	110, 163, 291, 346, 352
Miyagawa, Kazuya	BC03	42
Miyahara, Shin	GB01, PA07	76, 117
Miyake, Atsushi	PB27, PC21, PC23	124, 130
Miyake, Kazumasa	IG01	95
Miyake, Kohsaku	CE05	52
Miyake, Takashi	RG07	272

## AUTHOR INDEX

Miyamachi, Toshio Miyamoto, Koji Miyamoto, Souta Miyao, Masanobu Miyasaka, Hitoshi Miyasaka, S. Miyasaka, Shigeki Miyasaka, Toshiki Miyata, Atsuhiko Miyata, Masahiko Miyawaki, Tetsuya Miyazaki, H. Miyazaki, Hidetoshi Miyazaki, Ryoichi Miyazaki, Shota Miyazaki, Takamichi Miyazaki, Terunobu Mizokawa, T. Mizuguchi, Masaki Mizuguchi, Y. Mizuguchi, Yoshikazu Mizukami, Shigemi Mizumaki, Masaichiro Mizuno, Fumio Mizuno, Hiroki Mizuno, Yoshiyuki Mizunuma, Kotaro Mizuochi, Norikaze Mizuochi, Norikazu Mizusaki, Soichiro Mizushima, Toshio Moeller, A. Moeller, Thomas Moessner, R. Mohammad Shafique, Anwar Mohan, Ashwin Mohn, Peter Mohseni, S. M. Mokrousov, Yuriy Molavi, Mostafa Mole, R. A. Mole, Richard A Molegraaf, Hajo Molenda, J. Molenkamp, Laurens Wigbolt AC01 Molkanov, P. L. Moll, Philip J. W. Moll, Philip J.w. Molodtsov, Serguei L. Molokeev, Maxim Mombetsu, Shota Momoi, Tsutomu Momono, Naoki

Monaco, Giulio Mondragon, Reginaldo Montaigne, Francois Montbrun, Richard Monteiro, Pedro M.d.s.

PN23

GI03 82 KF05 111 340 SG05 HE01, RK10 86, 287 207 QF21 CF04, HA01 53, 83 PH07, PI07, QB24 148, 151, 191 QL03, RN02 228, 297 AC03 34 37 AG03 PJ12, PL12, RC18, SA07, 156, 165, 261, 317, SA17, SK18, SO20 319, 353, 370 SG11 341 AD04, QH28 35, 217 RD17 264 SG16 342 EJ01, RO06 71, 301 CH02 54 RB28 258 105, 230, 292 JI04, QL15, RL15 258 RB28 188, 189 QB09, QB17 CH02, EJ01 54, 71 IG03, QD17 95, 199 RH12 276 RI01 279 301 RO06 SI17 348 SI05 346 291, 352 RL11, SK10 144, 351 PG13, SK04 SG05 340 QE23 203 KD01 109 CC04 50 PM14 169 SO22 370 EF05 68 CE02, PK13 51, 162 IH03 96 PC16 129 PI03 151 QG26 212 DH02 62 PH06 147 34 309 RQ13 SC16 328 RB19, RB21 257 CD03, PD12 51, 133 152 PI11 263 RD07 101 JD04 PB08, QA19, QJ04, 121, 185, 221, 263, RD06, RD22, SE06 265, 334 258 RB28 Q026 245 KF02, SH02 110, 343 EG02 69

Author Index

Monteiro, Rodrigo S.	Q003
Montoncello, Federico	GG02
Moon, Eun-gook	JD05
Moon, Han Seb	RQ18
Moon, Jung-hwan	CE04, EE05, HP21, RJ19
Moon, Kyoung-woong	KE03
Moon, Seung Ho	CB01, PO15, PO16
Moon, Wonha	RH08
Morais, Paulo Cesar	PO08
Morales, R.	BD01, ED03
Morales, Rafael	BD04
Morales, Ricardo	Q026
Morbec, J. M.	RM02
Morellon, Luis	SM16
Morenzoni, Elvezio	PB20, QL19
Moretti Sala, Marco	BA03, RB28
Morgan, Jason	KE01
Mori, Shigeo	AF03, PG15, QG22, QG28
Moriki, Takahiro	QH15
Morimoto, Kazuhiro	AE05
Morimoto, Y.	SL14
Morinaga, Rei	QE13
Morioka, T.	PG24
Morisako, Akimitsu	PK07, QP08, RL05
Morishita, Kohei	RP01
Morita, Kyohei	PD25
Moriyama, Takumi	KD03
Moro, Maria	PG08
Morodomi, Hiroki	QE18, QF14
Morosov, Alexander I	PL05
Mosendz, O.	BG02
Moshchalkov, Victor	QC19, QC20
Moskvin, Alexander	PB14
Moskvin, Evgeny V.	DF04
Mostovoy, M.	RA05
Mostovoy, Maxim	CI01, DF01
Mostovshchikova, Elena	SL17
Mota, A. C.	SC03
Motome, Yukitoshi	JD02, PE18, PG23, RG03
Motonami, Satoru	RD20
Motoya, Kiyoichiro	HJ03, PH15, QE13
Motoyama, Gaku	QC03
M. Ghahfarokhi, Ebrahim	RN22
Movshovich, R.	QB31, SC03
Movshovich, Roman	SB12
Moya, Carlos	RN14
Moya, Maria Dolores Salvador	RO08
Moya, Xavier	AE03
Moyoshi, Taketo	HJ03, PH15
Mruczkiewicz, Michal	Q005
Mryasov, Oleg	JIO3
Mryasov, Oleg N	KH04
Mryasov, Oleg N.	JF04
Msomi, Justice	QL20
Mudry, Christopher	QF20
Mudryk, Yaroslav	1101
	JJ01
Muechler, Lukas	RC10
Muehlbauer, S. Clemens	
	RC10

	241
	80
	101
	310
, HP21, RJ19	52, 68, 30, 285
	110
5, PO16	49, 179
	276
	178
3	43, 66
	43
	245
	295
	363
	123, 231
3	41, 258
	110
5, QG22,	37, 144, 211, 212
	215
	36
	358
	201
	146
3, RL05	161, 246, 290
	303
	136
	109
4	143
1	202, 206
	164
0	45
.0	195
	122 60
	250
	54, 60
	358
	326
, PG23, RG03	
, 1 025, 11005	265
, QE13	90, 149, 201
	192
	300
3	192, 326
	322
	299
	302
	36
	90, 149
	241
	105
	112
	103
	231
	207
	106
	260
	98
	226

Mueller, Jens	QA11, QJ11	184, 222
Mueller, M.	CD01	50
Mueller, Marc	EG02	69
Mufti, Nadang	AB02	34
Mufti, Nandang	PH17	149
Muhammad, Kamran	SB01	320
Muhlbauer, Sebastian	CC03, DF02	50, 60
Muir, Jack	PO12	179
Mukadam, M. D. Mukai, Kazuhiko	PN17 PO17	174 180
Mukhamedshina, Nuranya	RR28	315
Mukherjee, Sankha Subhra	EG03, HD02, PJ08, SI02	69, 85, 156, 346
Mukherjee, Sudip	RN17	299
Mukhin, A. A.	PA15	118
Mukhin, Alexander	PA13	118
Mukhopadhyay, Sutirtha	IC02	92
Mukovskii, Ya. M.	RQ13	309
Mukuda, Hidekazu	DA04, PB26, QB29	56, 124, 191
Mulders, A.m.	SO08	368
Muller, M.	SF10	338
Muller, Ralf	RI12	280
Mun, E. D.	IA05	91
Mun, Eundeok	QB13	189
Munchenberger, Jana	IF04	94
M. Cagigas, Julian Andres	RB10	255
Munevar, Julian	RB11	255
Munoz, Carmen	KI02, SD16	113, 332
Munzenberg, Markus	DI03	63
Murakami, Naoki	RP01	303
Murakami, Shuichi	RL16	292
Murakami, Y.	CF04, PL14	53, 166
Murakami, Yoichi	SE07	334
Murakami, Youichi Murakami, Yuta	PH07, PH08, PI07	148, 151
Murakami, Yuta Murakawa, H.	RI14 PA14	281 118
Muramoto, Shingo	IF02	94
Muranoto, Shingo Murasawa, Y.	PL04	164
Murata, Atsushi	QM16	234
Murayama, Shigeyuki	QA19, QC01, QJ04,	185, 192, 221, 263,
, , , , ,	RD06, RD22, SE06	265, 334
Murayama, Sigeyuki	PB08	121
Mure, Elena	IE02, JE01	93, 101
Muro, Y	SE03, SE05, SE13	333, 334, 335
Muro, Y.	KA02, QC12	107, 194
Muro, Yuji	PD24, QE13	135, 201
Muromachi, Eiji	PC21	130
Murphy, J.	QC23	196
Murphy, T.	IA05	91
M. Ghahfarokhi, S Ebrahim	PB15	122
Musfeldt, J. L.	AA04	33
Mutka, Hannu	BC02	42
Muto, Tomoki	RH03	275
Myakon'kikh, Veronika	PC07	127
Mysik, Alexey	QL07	229
Ν		
Na, Sewoong	RB14	256
Naaman, Ron	GI02	81
Naastepad, F.	RQ12	309
Nafradi, Balint	PA14	118

Nagai Kousuka	PG09	142
Nagai, Kousuke Nagalakshmi, Ramamoorthi	QC13	143 194
Nagano, Katsuaki	RD10	263
Nagano, Takuma	PG10	144
Nagano, Y.	SL14	358
Nagano, Yutaro	SB14	322
Naganuma, Hiroshi	CH02, El02, KD04, RL13,	
Naganuma, mitosm	SI03, SI17, SK08	346, 348, 351
Nagaosa, N.	RA05	250
Nagaosa, Naoto	AE05, DF01	36, 60
Nagasawa, N	PC04	127
Nagasawa, N.	EF01	68
Nagasawa, Naohiro	RD04	262
Nagashima, Ryo	QB15	189
Nagata, Shoichi	QH14	215
Nagata, Tomoko	RA18	252
Nagata, Yujiro	PG13, SK04	144, 351
Naghavi, S. S.	QG24	212
Nair, K. G. M.	PN01	172
Naito, Makoto	PG13	144
Naito, Tomoyuki	SG18	342
Naka, Makoto	GB02	76
Naka, Takashi	RN11	298
Nakada, Kengo	PN13	174
Nakagawa, Shigeki	HB03	84
Nakagawa, T	QM12	233
Nakagawa, Takeshi	AG04, El01	38, 70
Nakahara, Mako	SG05	340
Nakai, Ikuo	PN07	173
Nakai, Yusuke	QC08, RD05, SO515	193, 262, 369
Nakajima, K.	RA22	253
Nakajima, Kazuo	RP01	303
Nakajima, Kenji	RH12	276
Nakajima, Taro	PH05	147
Nakama, Akihiro	RD17	264
Nakama, Takao	PD14, QC11, QI12, QI13	134, 194, 219
Nakamoto, Atsuhiko	RP14	305
Nakamoto, Go	QG11	210
Nakamura, Ai	QI12, QI13	219
Nakamura, Hiroyuki	HJ02, QA12, QH13	90, 184, 215
Nakamura, Kazuma	QB10	188
Nakamura, Kohji	EC02, RM04, RM05, SO10, SO11	65, 295, 296, 368
Nakamura, Mitsutaka	RH12	276
Nakamura, Mitsuteru	PD11	133
Nakamura, Saori	QH14	215
Nakamura, Shintaro	SL04	356
Nakamura, Takuya	RP14	305
Nakamura, Tetsuya	RP07	304
Nakamura, Toshiyuki	IG03	95
Nakane, Takayuki	RN11	298
Nakanishi, Yoshiki	PD11	133
Nakano, Hiroki	QE07, QE11, RI04, RI16	200, 201, 279, 281
Nakano, Kyohei	PF12, RF03	141, 269
Nakano, Masaki	JH01, PM08, PM09, RO03	104, 168, 301
Nakano, Takehito	HC02, PE15, RI17	84, 138, 281
Nakano, Tomohito	RD24	266
Nakao, A.	CF04	53
Nakao, Akiko	PH07, SF04	148, 337

## AUTHOR INDEX

Nakao, H. Nakao, Hironori

Nakao, Yuya Nakashima, Miho Nakashima, Naoya Nakata, Hayato Nakatani, Ryoichi Nakatani, Y. Nakatani, Yoshinobu

Nakatsuji, S. Nakatsuji, Satoru

Nakayama, T. Nakayama, Takayoshi Nakmura, Tetsuro Nakotte, Heinrich Nakotte, Heinz Naletov, Vladimir V. Nallamuthu, Sengodan Nam, Chunghee Nam, Hyoung Uk Nam, Moon-sun Nam, Nguyen Hoang Nam, Yoon Jae Namatame, Hirofumi Namkung, Seok Nandy, Ashis Kumar Nanto, Dwi Naohiro, Nagasawa Napolskiy, Kirill Nara, Daichi Nara, So Naramoto, Hiroshi Narayanapillai, K. Nardone, Marc Narikiyo, Kurahito Nartowt, B Narumi, Y. Narumi, Yasuo Naruse, Koki Nascimento, Valberto Pedruzi Nasirpouri, Farzad Nasreen, Farzana Nasu, Joji Nateprov, Alexandr Nikolaevich Nath, Ramesh Naumov, Sergey Naumova, Larisa Naumova, Larisa I. Navarro, Oracio Navarro-quezada, Andrea Navas, D. Nazmunnahar, Mst. Nedelko, N. Nekrasov, I. A.

Nelson- Cheeseman, Brittany

CF04, PL14 53, 166 PH07, PH08, PI07, SE07, 148, 151, 334, 337 SF04 PG09 143 QH04, RD19 213, 265 230 QL12, QL13 QE21 202 DE04, EI05, PK08, SK33 60, 71, 161, 355 RL07 291 HG02, QP03, RJ01, RJ02, 88, 245, 282, 284, RJ09, SJ01 349 CF01 52 CC05, CD02, IG03, 50, 51, 95, 144, 196, PG14, QD01, RD13, 264, 333, 338 SE01, SF07 SK03 350 PF15 141 QC07 193 SF13 339 RE09 267 93 ID03 194 QC13 KC01 108 244, 343 QO20, SG21 323 SB20 84, 138 HC02, PE15 166 PL16 265 RD20 300, 301 RO01, RO02 270 RF12 OH01 213 QC16 194 SL27 360 RP14 305 QG18 211 BH05, BH06 46, 47 GE03, PJ20, PJ21, RL19 78, 158, 293 RQ16 309 SB24 324 SE14 335 GJ03 82 PG17, QF09, RP07 145, 205, 304 QF10 205 234 QM15 232 QM05 SF13 339 RH10 276 IG02, QG23 95, 212 QE12 201 358 SL17 350 SK02 350 SK01 211, 245, 352 QG21, QO26, SK15 BE03 44 66 ED03 DJ03 63 PH06, PM05 147, 168 RF16 271 DH01 62

Nelson, C.s.	PH10
Nelson, Christie S.	PA02
Nemnes, George Alexandru	QK20
Nemoto, Yuichi	RB05, RD07
Nenert, G.	QF22
Nenert, Gwilherm	PI17, QJ02
Nepijko, S A	QM12, RN12, RP20
Neu, Volker	DG01
Neubauer, Andreas	DF02, QH12, SF05
Neumann, I.	DI01
Neupane, Madhab	HI02, SD04
Neusser, Sebastian	GG02
Ng, Dickon H.I.	BE04
Ngubane, Jeanette	QH22
Nguyen, Anh T. N	CE02
Nguyen, Anh Tuan	QF16
Nguyen, Duc-tho	RA10
Nguyen, Hoa	SL05
Nguyen, Hong T. M.	QL16
Nguyen, Khanh Duy	QA10, QA17
Nguyen, Sinh Huy	CJ05
Nguyen, Thanh Van	5013
Nguyen, Thuy-trang	RF02
Nguyen, Tiep H.	QL16
Nguyen, Tuan Anh	SO13
Nguyen, Van Dai	GE01, QK15, QK16,
Nguyen, van Dar	QK17, RL24, SL24
Nguyen, Van Thanh	QF16
Nhlapo, Ta	QL20
Ni, Ni	QB13
Niarchos, Dimitrios	EJ04
Niarchos, Dimitrios G	RO08
Niarchos, Dimitris G	BJ02, DD06, QP12,
Marchos, Dimitris G	QP13, RQ10
Nicklas, M.	SC03
Nicklas, Michael	AB02, IA02, RD12, SE
Nie, Wenxing	PF18
Nie, Yuefeng	HA03
Niedermayer, Ch.	OE19
Niedermayer, Christof	HA03
Nielsch, Kornelius	RP04
Niemeyer, Markus	KD01
Niermann, Daniel	QA22
Nigam, A K	Q004
Nigam, A. K.	QH06, SO22
Nigam, A.k	PG19
Nii, Yoichi	QA10, QA13
Niimi, Yasuhiro	HE03, ID04
Niimi, Yohei	RC18
Niizeki, T	SK29
Niizeki, Tomohiko	IF03
Niki, Haruo	QH14
Nikolaev, Pavel	PJ32
Nikolenko, Victor V.	Q012
Nikseresht, Neda	QF23
Nilsen, Gøran J.	QE05, QE08
Nimori, S.	QJ10

	148	Ninios, Kostas
	116	Ninomiya, Keisuke
	227	Niratisairak, Sanyalak
	254, 263	Nishi, Yusuke
	207	Nishihara, Sadafumi
	153, 221	Nishihara, Yoshikazu
)	233, 298, 306	Nishikawa, Takuo
	61	Nishikawa, Yunori
	60, 214, 337	Nishimoto, Satoshi
	62	Nishimura, Katsuhiko
	89, 330	Nishimura, Naoto
	80	Nishimura, Taizo
	44	Nishino, Yoichi
	216	Nishioka, Takashi
	51	Nishisako, Yuya
	206	Nishishita, Muneyuki
	251	Nishitani, Koji
	356	Nishitani, Y.
	230	Nishiumi, Naoto
	184, 185	Nishiyama, Kouhei
	56	Nishiyama, Masahide
	368	Nishiyama, N.
	269	Nishiyama, Norimasa
	230	Nishizaki, Terukazu
	368	Nitta, Junsaku
,	78, 226, 293, 360	Niyogi, Sandip
		Nizhankovskii, Victor
	206	Nizzoli, Fabrizio
	231	No, Young-il
	189	Noad, H. M. L.
	72	Noda, Kazuto
	302	Nogami, A.
	47, 59, 247, 308	Noguchi, Tomoaki
	326	Noh, Han-jin Noh, Hwayong
SB29	34, 90, 264, 325	Noh, Hyunho
5025	142	Noh, Seung-hyun
	83	Noh, Su Jung
	202	Noh, T. W.
	83	Nohara, Hiroki
	303	Noji, Takashi
	109	Nojima, Tsutomu
	186	Nojiri, H.
	241	Nojiri, Hiroyuki
	213, 370	Nolas, George S.
	145	Noll, T
	184, 175	Nolting, Frithjof
	86, 93	Nomura, Hikaru
	261	Nomura, Kentaro
	355	Nomura, Ryuya
	94	Nomura, Tatsuya
	215	Nonoguchi, Seiji
	160	
	242	Nordblad, Per
	207	Noro, Yoshihiko
	200	Notin, Lucien
	222	
	154, 280	Novais, Erico
	79	Novak, Vit
	315	Noverola, Humberto

s	BC05	42
eisuke	PN26	176
Sanyalak	RP17	306
2	QD18	199
dafumi	QG22	211
shikazu	PM13	169
akuo	SK08	351
unori	SF15	339
atoshi	PF09, PF12	140, 141
atsuhiko	QI04, QJ06	218, 221
aoto	PD13	134
izo	PG17	145
hi	QH28	217
ashi	PD10, QD06, RD10	133, 197, 263
iya	SK23	354
uneyuki	PF19	142
i	QD15	198
	BE02	44
oto	QC18	195
ouhei	QD12	198
Iasahide	BC04, QE22	42, 203
lasaniae	SG02	339
orimasa	PE10	137
ukazu	SG18	342
ukazu U	KC03	108
	PN16	174
p i, Victor	JJ03	106
		80
zio	GG02	46
	BH02	
L.	GF01	79
)	PF16	141
L:	PA04	116
noaki	PE10	137
	HJ05, PB12	90, 122
ng	PJ16, PJ24	157, 158
)	SD14	332
nyun	PO15, PO16	179
I	PJ14, SH13	157, 345
	HP12, SA16	30, 318
ki	SF03	337
	QF07	205
omu	SL04	356
	PG24	146
ki	QF09, RP07	205, 304
e S.	DJ01	63
	RP08	304
jof	BC02, RK17	42, 288
aru	DE04	60
itaro	AE05	36
ya	QL11	229
iuya	KC06, QK24, RL20, RL23	109, 228, 293
eiji	KC06, QK24, QL11, RL20, RL23	109, 228, 229, 293
r	PM10, QH11, RN07	169, 214, 297
iko	PG13, SK04	144, 351
	EI04, GE01, QK15, OK17, RL24	71, 78, 226, 293
	QK17, RL24	204
	RJ13	284
	BE05	44
mberto	Q026	245

Novitskiy, Nicolay	QM08	232
Nowak, Uli	FC03	74
Nowak, Ulrich	FA02, RJ04	72, 283
Nozaki, Hiroshi	PIO6, PO17	151, 180
Nozaki, Taiichiro	RL25	294
Nozaki, Takayuki Nozaki, Tamahira	BI01, KF01, PK20, SK10	47, 110, 163, 352
Nozaki, Tomohiro Nozaki, Yukio	PL10, SM03 JG03, QP09, RJ10, RK10, SI08	165, 361 103, 246, 284, 287, 347
Nozawa, Akihiko	RB06	255
Nozue, Yasuo Nucara, Alessandro	HC02, PE15, RI17	84, 138, 281
Nugroho, Agung	SA08 CC05, PH17	317 50, 149
Nunes, Sabrina Esperanca	SB18	323
0		
-		
O' Farrell, Eoin	HA05	84
O' Farrell, Eoin T. C.	CD02	51
O'farrell, Eoin	SF07	338
Obi, Yoshihisa	SEO6 SI15	334 348
Obry, Bjoern Obry, Bjorn	HD01	85
Ochai, Akira	QD11	198
Ochiai, Akira	QJ08	222
Ochoa, Martha Teresita	Q026	245
Oda, Migaku	PB08	121
Oda, Tatsuki	AE02, PL06, RC02, RL02	
Odkhuu, Doji	Q023	244
Odkhuu, Dorj	KI01	112
Oezelt, Harald	DE01	59
O'farrell, Eoin Conor	QD01	196
Ogata, Masao	SE04	333
Ogawa, Daisuke	EJ01, RO05, RO06	71, 301
Ogawa, Satoru	SG18	342
Ogita, Norio	PC05, PG16, RD10, SG15	127, 145, 263, 342
Ognev, Alexey	AG05, QM05	38, 232
Ogrin, Feodor Y.	RK20	289
Oguchi, Tamio	EC03	65
Oguri, Akira	PE16	138
Ogusu, Hiroki	PC17	129
Oh, Hangdeok	PO04	177
Oh, Hyungju	RB16	256
Oh, Ok Kyun	QN17	239
Oh, S.	SB28	324
Oh, S. H.	SA06	317
Oh, Sang-soo	QM22	235
Oh, Se-jung	GF05	79
Oh, Seol Hee	SA05	317
Oh, Seung-lim	RN20	300
Oh, Suhk Kun	SH01	343
Oh, Suhk-kun	QH01	213
Oh, Sunjong	Q019	243
Oh, Yoon Seok	JB02	99
Oh, Young Kun	PB40	126
Ohara, Jun	IJ05	98
Ohara, S.	RD15	264
Ohara, Satoshi	RN11	298
Ohara, Shigeo	QD17, RD18, RD20,	199, 265, 335

Nimori, Shigeki

Ning, F. L.

Ning, Sun

PI22, RI07

GF01

RR25

## AUTHOR INDEX

Ohashi, Masashi Ohashi, Takuma Ohashi, Tatsuro Ohdaira, Yusuke Ohe, Junichiro Ohe, Jun-ichiro Ohgushi, K. Ohgushi, Kenya Ohishi, Kazuki Ohkawara, Manabu Ohkubo, Tadakatsu Ohmichi, Eiji Ohnish, Tomohiro Ohnishi, Nozomi Ohnishi, Tomohiro Ohno, H. Ohno, Hideo Ohno, Masumi Ohno, T. Ohno, Takahisa Ohno, Takashi Ohno, Y. Ohnuma, Shigehiro Ohoyama, Kenji Ohsawa, Takashi Ohsuka, T Ohsumi, Hiroyuki Ohta, Hitoshi

Ohta, Tsuyoshi Ohta, Yukinori

Ohtake, Mitsuru Ohtani, H. Ohtani, Taiki Ohtomo, Manabu Ohtori, Hiroyuki Ohuchi, Syuya Ohya, Masahiro Ohzuku, Tsutomu Oikawa, Hiroto Oka, Kengo Oka, Kunihiko

Oka, Takahi Oka, Takashi Oka, Toshihide Okabe, Hirotaka Okada, Hironari Okada, Kyoko Okada, Takuya Okamoto, Hiroaki Okamoto, J. Okamoto, Kazuaki Okamoto, Kiyomi Okamoto, Yoshihiko QJ12 222 RF09 270 KC03 108 EI02 70 294 RL27 FA03, FD04 72, 74 339 SG02 IA04, QO13, SG04 91, 242, 340 QB02, QB04 187 RP01 303 94 IF03 RP09, RP19 304, 306 351 SK04 201 QE10 PG13 144 44, 53, 344 BE02, CG02, SH07 47, 31, 247, 348 BI02, HP52, QP14, SI17 QF10 205 SG02 339 RF07 270 QF24 207 CG02 53 JI01, SL04 105, 356 QG18, RP01 211, 303 BI02 47 PC04 127 IA04, QO13 91, 242 KI05, PE11, QE03, QE10, 113, 137, 199, 201, QF04, RD01, RP05, RP09, 204, 262, 304, 306 RP19 QI04 218 KG03, PF09, PF10, PF12, 111, 140, 141, 267, PF14, PF15, RE08, RF03, 269, 337 SF04 66, 170, 363 ED01, PM18, SM15 214 QH09 214 QH10 363 SM15 46, 47 BH05, BH06 250 RA09 PI22 154 PD08 133 PO17 180 EI05 71 RA06 250 QA05, QA07, QA09, 183, 184, 276 RH11 RI14 281 EB02, HI05 65, 89 QB11 188 PC21, SF02 130, 337 PG07, QH21 143, 216 SK13 352 279 RI01 263 RD06 166 PL14 111 KF05 280 RI11 JA04, QE03, QE05, 99, 199, 200, 203, QE08, QE27, SG14 341

Olympia Tatawa		266 227
Okane, Tetsuo	RD25, SE23	266, 337
Okawauchi, Yasuharu	SC13	327
Okazaki, K. Okazaki, Buuii	QB28	191
Okazaki, Ryuji Oki, Soichiro	PH12, QC06 HE01	148, 193 86
Okita, Kazuhiko	RP11	305
Okubo, Akinari	SK23	354
Okubo, Kazuma	RP01	303
Okubo, S.	KI05	113
Okubo, Susumu	QE03, QE10, QF04, RP05	199, 201, 204, 304
Okuda, Mitsuhiro	KB01	107
Okuda, Taichi	KF05	111
Okuda, Tetsuji	RH04, SB24	275, 324
Okunishi, Kouichi	RI11	280
Okutomi, Yoshihito	CE05	52
Okuyama, Daisuke	PA12	117
Okuyama, Hiroaki	SG05	340
Oles, Andrzej	RG18	274
Olivetti, Elena Sonia	DG04	61
Omura, Kumiko	PI06	151
Oner, Yildirhan	QG06, QJ15	209, 223
Ong, Chin Shen	SL06	357
Ong, N. P.	AB04	34
Onimaru, T	PC04	127
Onimaru, Takahiro	EF01, PD24, RD04,	68, 135, 262, 341,
	SG12, SG15	342
Ono, Akiko	PC20, QJ07	130, 222
Ono, Hiroaki	QP20	248
Ono, Hiroshi	SK21	353
Ono, Shimpei	RL10	291
Ono, T.	RL07, SH07	291, 344
Ono, Teruo	GE05, HG02, RJ02, RL10	79, 88, 282, 291
Ono, Toshio	QE22	203
		101
Onoda, Shigeki	JD03	
Onodera, Hideya	QI10, QI16	219, 220
Onodera, Hideya Onosaka, Atsuhi	QI10, QI16 JA04	219, 220 99
Onodera, Hideya	QI10, QI16 JA04 HP11, PB27, PC28,	219, 220 99 30, 124, 131, 133,
Onodera, Hideya Onosaka, Atsuhi	QI10, QI16 JA04	219, 220 99
Onodera, Hideya Onosaka, Atsuhi	QI10, QI16 JA04 HP11, PB27, PC28, PD08, PD13, PD16, QC03, QH03, RD11, RD25, RE10, SB15, SE23,	219, 220 99 30, 124, 131, 133, 134, 192, 213, 263,
Onodera, Hideya Onosaka, Atsuhi	QI10, QI16 JA04 HP11, PB27, PC28, PD08, PD13, PD16, QC03, QH03, RD11,	219, 220 99 30, 124, 131, 133, 134, 192, 213, 263,
Onodera, Hideya Onosaka, Atsuhi	QI10, QI16 JA04 HP11, PB27, PC28, PD08, PD13, PD16, QC03, QH03, RD11, RD25, RE10, SB15, SE23,	219, 220 99 30, 124, 131, 133, 134, 192, 213, 263,
Onodera, Hideya Onosaka, Atsuhi Onuki, Yoshichika	QI10, QI16 JA04 HP11, PB27, PC28, PD08, PD13, PD16, QC03, QH03, RD11, RD25, RE10, SB15, SE23, SF03	219, 220 99 30, 124, 131, 133, 134, 192, 213, 263, 266, 268, 322, 337
Onodera, Hideya Onosaka, Atsuhi Onuki, Yoshichika Onuma, Kazuki	QI10, QI16 JA04 HP11, PB27, PC28, PD08, PD13, PD16, QC03, QH03, RD11, RD25, RE10, SB15, SE23, SF03 QL03, RN02	219, 220 99 30, 124, 131, 133, 134, 192, 213, 263, 266, 268, 322, 337 228, 297
Onodera, Hideya Onosaka, Atsuhi Onuki, Yoshichika Onuma, Kazuki	QI10, QI16 JA04 HP11, PB27, PC28, PD08, PD13, PD16, QC03, QH03, RD11, RD25, RE10, SB15, SE23, SF03 QL03, RN02 CH02, EI02, EJ01, RL13,	219, 220 99 30, 124, 131, 133, 134, 192, 213, 263, 266, 268, 322, 337 228, 297 54, 70, 71, 292, 346,
Onodera, Hideya Onosaka, Atsuhi Onuki, Yoshichika Onuma, Kazuki Oogane, Mikihiko	QI10, QI16 JA04 HP11, PB27, PC28, PD08, PD13, PD16, QC03, QH03, RD11, RD25, RE10, SB15, SE23, SF03 QL03, RN02 CH02, EI02, EJ01, RL13, SI03, SI17, SK08	219, 220 99 30, 124, 131, 133, 134, 192, 213, 263, 266, 268, 322, 337 228, 297 54, 70, 71, 292, 346, 348, 351
Onodera, Hideya Onosaka, Atsuhi Onuki, Yoshichika Onuma, Kazuki Oogane, Mikihiko Oogane, Mikiko	QI10, QI16 JA04 HP11, PB27, PC28, PD08, PD13, PD16, QC03, QH03, RD11, RD25, RE10, SB15, SE23, SF03 QL03, RN02 CH02, EI02, EJ01, RL13, SI03, SI17, SK08 KD04	219, 220 99 30, 124, 131, 133, 134, 192, 213, 263, 266, 268, 322, 337 228, 297 54, 70, 71, 292, 346, 348, 351 110
Onodera, Hideya Onosaka, Atsuhi Onuki, Yoshichika Onuma, Kazuki Oogane, Mikihiko Oogane, Mikiko Oogane, Yuta	QI10, QI16 JA04 HP11, PB27, PC28, PD08, PD13, PD16, QC03, QH03, RD11, RD25, RE10, SB15, SE23, SF03 QL03, RN02 CH02, EI02, EJ01, RL13, SI03, SI17, SK08 KD04 QD06	219, 220 99 30, 124, 131, 133, 134, 192, 213, 263, 266, 268, 322, 337 228, 297 54, 70, 71, 292, 346, 348, 351 110 197
Onodera, Hideya Onosaka, Atsuhi Onuki, Yoshichika Onuma, Kazuki Oogane, Mikihiko Oogane, Mikiko Oogane, Yuta Ooi, Takahiro	QI10, QI16 JA04 HP11, PB27, PC28, PD08, PD13, PD16, QC03, QH03, RD11, RD25, RE10, SB15, SE23, SF03 QL03, RN02 CH02, EI02, EJ01, RL13, SI03, SI17, SK08 KD04 QD06 RF17	219, 220 99 30, 124, 131, 133, 134, 192, 213, 263, 266, 268, 322, 337 228, 297 54, 70, 71, 292, 346, 348, 351 110 197 271
Onodera, Hideya Onosaka, Atsuhi Onuki, Yoshichika Onuma, Kazuki Oogane, Mikihiko Oogane, Mikiko Oogane, Yuta Ooi, Takahiro Oomaru, Kiyoaki	QI10, QI16 JA04 HP11, PB27, PC28, PD08, PD13, PD16, QC03, QH03, RD11, RD25, RE10, SB15, SE23, SF03 QL03, RN02 CH02, EI02, EJ01, RL13, SI03, SI17, SK08 KD04 QD06 RF17 SJ01	219, 220 99 30, 124, 131, 133, 134, 192, 213, 263, 266, 268, 322, 337 228, 297 54, 70, 71, 292, 346, 348, 351 110 197 271 349
Onodera, Hideya Onosaka, Atsuhi Onuki, Yoshichika Onuma, Kazuki Oogane, Mikihiko Oogane, Mikiko Oogane, Yuta Ooi, Takahiro Oomaru, Kiyoaki Oomi, Gendo	Ql10, Ql16 JA04 HP11, PB27, PC28, PD08, PD13, PD16, QC03, QH03, RD11, RD25, RE10, SB15, SE23, SF03 QL03, RN02 CH02, El02, EJ01, RL13, Sl03, Sl17, SK08 KD04 QD06 RF17 SJ01 PJ03	219, 220 99 30, 124, 131, 133, 134, 192, 213, 263, 266, 268, 322, 337 228, 297 54, 70, 71, 292, 346, 348, 351 110 197 271 349 155
Onodera, Hideya Onosaka, Atsuhi Onuki, Yoshichika Onuma, Kazuki Oogane, Mikihiko Oogane, Mikiko Oogane, Yuta Ooi, Takahiro Oomaru, Kiyoaki Oomi, Gendo Oosawa, Akira	Ql10, Ql16 JA04 HP11, PB27, PC28, PD08, PD13, PD16, QC03, QH03, RD11, RD25, RE10, SB15, SE23, SF03 QL03, RN02 CH02, El02, EJ01, RL13, SI03, SI17, SK08 KD04 QD06 RF17 SJ01 PJ03 RH03	219, 220 99 30, 124, 131, 133, 134, 192, 213, 263, 266, 268, 322, 337 228, 297 54, 70, 71, 292, 346, 348, 351 110 197 271 349 155 275
Onodera, Hideya Onosaka, Atsuhi Onuki, Yoshichika Onuma, Kazuki Oogane, Mikihiko Oogane, Mikiho Oogane, Yuta Ooi, Takahiro Oomaru, Kiyoaki Oomi, Gendo Oosawa, Akira Oozono, Satoshi Openeer, Peter Oppeneer, P M	QI10, QI16 JA04 HP11, PB27, PC28, PD08, PD13, PD16, QC03, QH03, RD11, RD25, RE10, SB15, SE23, SF03 QL03, RN02 CH02, EI02, EJ01, RL13, SI03, SI17, SK08 KD04 QD06 RF17 SJ01 PJ03 RH03 RH04 FC03 QM12	219, 220 99 30, 124, 131, 133, 134, 192, 213, 263, 266, 268, 322, 337 228, 297 54, 70, 71, 292, 346, 348, 351 110 197 271 349 155 275 275 74 233
Onodera, Hideya Onosaka, Atsuhi Onuki, Yoshichika Onuma, Kazuki Oogane, Mikihiko Oogane, Mikiho Oogane, Yuta Ooi, Takahiro Oomaru, Kiyoaki Oomi, Gendo Oosawa, Akira Oozono, Satoshi Openeer, Peter Oppeneer, P M Oppeneer, P. M.	QI10, QI16 JA04 HP11, PB27, PC28, PD08, PD13, PD16, QC03, QH03, RD11, RD25, RE10, SB15, SE23, SF03 QL03, RN02 CH02, EI02, EJ01, RL13, SI03, SI17, SK08 KD04 QD06 RF17 SJ01 PJ03 RH03 RH04 FC03 QM12 RG05	219, 220 99 30, 124, 131, 133, 134, 192, 213, 263, 266, 268, 322, 337 228, 297 54, 70, 71, 292, 346, 348, 351 110 197 271 349 155 275 275 74 233 272
Onodera, Hideya Onosaka, Atsuhi Onuki, Yoshichika Onuma, Kazuki Oogane, Mikihiko Oogane, Mikiho Oogane, Yuta Ooj, Takahiro Oomaru, Kiyoaki Oomi, Gendo Oosawa, Akira Oozono, Satoshi Openeer, Peter Oppeneer, P. M. Oppeneer, Peter	QI10, QI16 JA04 HP11, PB27, PC28, PD08, PD13, PD16, QC03, QH03, RD11, RD25, RE10, SB15, SE23, SF03 QL03, RN02 CH02, EI02, EJ01, RL13, SI03, SI17, SK08 KD04 QD06 RF17 SJ01 PJ03 RH03 RH04 FC03 QM12 RG05 BC02	219, 220 99 30, 124, 131, 133, 134, 192, 213, 263, 266, 268, 322, 337 228, 297 54, 70, 71, 292, 346, 348, 351 110 197 271 349 155 275 275 275 74 233 272 42
Onodera, Hideya Onosaka, Atsuhi Onuki, Yoshichika Onuma, Kazuki Oogane, Mikihiko Oogane, Mikiko Oogane, Yuta Ooj, Takahiro Oomaru, Kiyoaki Oomi, Gendo Oosawa, Akira Oozono, Satoshi Openeer, Peter Oppeneer, P M Oppeneer, P. M.	QI10, QI16 JA04 HP11, PB27, PC28, PD08, PD13, PD16, QC03, QH03, RD11, RD25, RE10, SB15, SE23, SF03 QL03, RN02 CH02, EI02, EJ01, RL13, SI03, SI17, SK08 KD04 QD06 RF17 SJ01 PJ03 RH03 RH03 RH04 FC03 QM12 RG05 BC02 RK12	219, 220 99 30, 124, 131, 133, 134, 192, 213, 263, 266, 268, 322, 337 228, 297 54, 70, 71, 292, 346, 348, 351 110 197 271 349 155 275 275 275 275 74 233 272 42 287
Onodera, Hideya Onosaka, Atsuhi Onuki, Yoshichika Onuma, Kazuki Oogane, Mikihiko Oogane, Mikihiko Oogane, Yuta Ooi, Takahiro Oomaru, Kiyoaki Oomi, Gendo Oosawa, Akira Oozono, Satoshi Openeer, Peter Oppeneer, P. M. Oppeneer, P. M. Oppeneer, Peter Oppeneer, Peter Oppeneer, Peter M Oppeneer, Peter M	QI10, QI16 JA04 HP11, PB27, PC28, PD08, PD13, PD16, QC03, QH03, RD11, RD25, RE10, SB15, SE23, SF03 QL03, RN02 CH02, EI02, EJ01, RL13, SI03, SI17, SK08 KD04 QD06 RF17 SJ01 PJ03 RH03 RH03 RH04 FC03 QM12 RG05 BC02 RK12 SA12	219, 220 99 30, 124, 131, 133, 134, 192, 213, 263, 266, 268, 322, 337 228, 297 54, 70, 71, 292, 346, 348, 351 110 197 271 349 155 275 275 275 275 74 233 272 42 287 318
Onodera, Hideya Onosaka, Atsuhi Onuki, Yoshichika Onuma, Kazuki Oogane, Mikihiko Oogane, Mikiko Oogane, Yuta Ooi, Takahiro Oomaru, Kiyoaki Oomi, Gendo Oosawa, Akira Oozono, Satoshi Openeer, Peter Oppeneer, P M Oppeneer, P. M. Oppeneer, Peter Oppeneer, Peter Oppeneer, Peter M Ordonez, John Edward Orel, Valerii	QI10, QI16 JA04 HP11, PB27, PC28, PD08, PD13, PD16, QC03, QH03, RD11, RD25, RE10, SB15, SE23, SF03 QL03, RN02 CH02, EI02, EJ01, RL13, SI03, SI17, SK08 KD04 QD06 RF17 SJ01 PJ03 RH03 RH03 RH04 FC03 QM12 RG05 BC02 RK12 SA12 BH01	219, 220 99 30, 124, 131, 133, 134, 192, 213, 263, 266, 268, 322, 337 228, 297 54, 70, 71, 292, 346, 348, 351 110 197 271 349 155 275 275 275 275 74 233 272 42 287 318 46
Onodera, Hideya Onosaka, Atsuhi Onuki, Yoshichika Onuma, Kazuki Oogane, Mikihiko Oogane, Mikihiko Oogane, Yuta Ooi, Takahiro Oomaru, Kiyoaki Oomi, Gendo Oosawa, Akira Oozono, Satoshi Openeer, Peter Oppeneer, P. M. Oppeneer, P. M. Oppeneer, Peter Oppeneer, Peter Oppeneer, Peter M Oppeneer, Peter M	QI10, QI16 JA04 HP11, PB27, PC28, PD08, PD13, PD16, QC03, QH03, RD11, RD25, RE10, SB15, SE23, SF03 QL03, RN02 CH02, EI02, EJ01, RL13, SI03, SI17, SK08 KD04 QD06 RF17 SJ01 PJ03 RH03 RH03 RH04 FC03 QM12 RG05 BC02 RK12 SA12	219, 220 99 30, 124, 131, 133, 134, 192, 213, 263, 266, 268, 322, 337 228, 297 54, 70, 71, 292, 346, 348, 351 110 197 271 349 155 275 275 275 275 74 233 272 42 287 318

Orihashi, Hiroki	QH08, QN20
Orii, Daisuke	PC21
Orlov, Andrey F	PN22
Ortega, Daniel	RN13
Ortiz, Christian	PJ29
Osakabe, Toyotaka	QI16
Osaki, Soichiro	SK33
Oshikawa, Masaki	PF18, QF11, RI02
Oshino, Yuuki	SK21
Oshiro, Morihito	QH14
Ostler, Thomas A.	RK17
Ota, Satoshi	PO03
Ota, Y.	QB28
Otani, Yoshichika	HD04, HE03, HE05,
	ID04, RK11
Otsubo, Toru	QI10
Otsuka, Hiromi	JD06
Otsuka, Shintaro	PO18
Otsuka, Yuichi	RG09
Otsuka, Yuto	QP09, QP10, QP11
Ott, Hans Rudolf	IC02
Ou, Z. Q.	SO16
Ou, Zhiqiang	S008
Ouardi, S.	SK24
Ouardi, Siham	CH03, SD15
Ouladdiaf, B.	RB17
Ouladdiaf, Bachir	SB08
Ouyang, Yuyu	AH04
Ovchinnikov, Alexander	DF03
Ovchinnikov, Alexander Sasha	
Ovchinnikov, S. G.	QK07
Oviedo Roa, Raul Owaki, Yasuhiro	SK15 PE15
Oyamada, Akira	BC04, QE22
Ozawa, Kiyoshi	PH02
Ozawa, Tadashi C	SK04
Ozdogan, K.	RP15
Ozerov, M.	QF02, QF03
020104, 141.	Q102, Q105
P	
P, D.	PC22
P, Suresh	RA16
P. D., Babu	QD04
P. Del Real, Rafael	GG04
Padilla, Jessica	EC04
Padilla-pantoja, Jessica	RA17
Paglione, J.	QB07
Pagliuso, Pascoal	SE22
Pagliuso, Pascoal G.	SE15
Pal, A	QM24
Pal, Semanti	HD04, RK11
Palacio, Fernando	KB06, PG22, QO16
Pal'chik, Marina	SM19
Palm, Eric	GJ02
Palstra, T.t.m	PH17
Pan, M.x.	RO14
Pan, Minxiang	RO13, RO15
Pan, Rui	SN01
Pan, Wei	ED02, SN11
Pandey, Srikrishna	PJ32

85, 86, 93, 287

Danday Cudhaliar	DE10	142
Pandey, Sudhakar Pandiyan, Manoj	PF19 DH03	142 62
Panholzer, Martin	RG15	274
Panigrahi, B. K.	PN01	172
Pankrats, A	QJ16	223
Pankrats, Anatolii	U01	97
Paolasini, Luigi	AF01	36
Papaioannou, Evangelos	AF05	37
Parakkat, Vineeth Mohanan	QM03, SJ08, SL23	232, 350, 359
Pardo, J. A.	KC05, PJ23	109, 158
Pareek, S. P.	SL25	360
Park, Bg.	PG11, QA04	144, 183
Park, Byong-guk	IF01	94
Park, Chan Woo	RL21	293
Park, Chang-yup	RL03, SM12	290, 362
Park, D.g	HH03	88
Park, D.g.	PM07, PO05, RQ04, SK17	168, 178, 307, 353
Park, Dong-soo	DJ02	63
Park, Eunsung	SE02	333
Park, Gwan Soo	QN04, QN05, QP04, RR07	237, 246, 312
Park, Gwang-seo	Q010	242
Park, Gwansoo	SN07	365
Park, H. L.	PM26	171
Park, Hae Jung	QN07, QN17	237, 239
Park, Hee Yeon	SB22	323
Park, Insuk	QB25	191
Park, Insung	PJ16, PK10, PK17	157, 162, 163
Park, J G.	PG24	146
Park, J. H.	IA05	91
Park, J. K.	RP10	304
Park, J. M. S.	PH03	147
Park, J. S.	RA07	250
Park, Jh.	PH10	148
Park, Jae - Hyun	RB14	256
Park, Jae Hyun	CE04	52
Park, Jae Y.	CJ03	55
Park, Jae-hoon	AA01, AF04, JB02	33, 37, 99
Park, Jae-hyuk	PN14	174
Park, Je Geun	HJ01	89
Park, Je-geun	GF05, PI10, RA22	79, 152, 253
Park, Jewook	RC14	261
Park, Ji Hun	PN20, PN21	175
Park, Jihoon	JH04, KH04	104, 112
Park, Jin-hong	DF01	60
Park, Jong C.	CJ03	55
Park, Joonbum	SD07	331
Park, Jun Kue	SI14	348 146
Park, Junghwan Park, Ju-young	PG24	99
Park, Kwang-jun	JBO2 PL11	165
Park, Lee-seul	SD14	332
Park, Po Gyu	RQ07	308
Park, San Youn	QA16	185
Park, Sang Kook	SB22	323
Park, Seonghun	AF04	37
Park, Seung-young	CE04, HF05, RL06	52, 87, 290
Park, Sungkyun	QM18, QM20	234
Park, Sungmin	Q010	242

## AUTHOR INDEX

Park, Sungyu	PC25	130
Park, Tae-eon	SL26	360
Park, Tuson	FB01, PC19, SE02	73, 129, 333
Park, W. K.	GA01	75
Park, Wan Kyu	RD09	263
Park, Wongoo	GF05	79
Park, Yeonjoo	PB18, PB37, PB39, PB40	123, 126
Park, Youn Ho	QK11	225
Park, Young Woo	QN15, QN17	238, 239
Park, Young-woo	QN07	237
Park, Younho	SL26	360
Park, Youn-ho	RL03	290
Park, Yu Seop	RR04	311
Parker, Dinah. R.	QG19	211
Parker, G.	BG02	45
Parkes, Duncan E	SH09	344
Parkin, Stuart	HP41	31
Parkin, Stuart S. P.	SM01	361
Parkin, Stuart S.p.	RK22	289
Parreiras, Sofia De Oliveira	RJ12	284
Parwanta, Kadek J.	QL09	229
Parwanta, Kadek Juliana	RA21	252
Pascarelli, Sakura	RH23	278
Paschen, S.	CD01, SF10	50, 338
Pascua, Gwendolyne	PI05, PI17, QB22, QJ03	151, 153, 190, 221
Pascual, Jose Ignacio	PG08	143
Pashkevich, Yurii G.	QJ03	221
Pashkin, Alexej	HA02	83
Pasquale, Massimo	DG04	61
Passamani, Edson Caetano	QM15	234
Paszkowicz, Wojciech	DD02	58
Patel, Sheena K. K.	PK09, QM01	161, 231
Pathak, Arjun Kumar	CH01	54
Patil, Sujata M.	QD16	199
Patil, Swapnil	PE14	138
Patlagan, L.	PD03	132
Patoary, M. Atiqur R.	PB24, PF05	124, 140
Patra, Ajit K	QK14	226
Patrin, Gennady	PJ27, SM19, SO28	159, 364, 371
Patrin, Konstantin	PJ27, SM19, SO28	159, 364, 371
Patthey, L.	PC01	126
Patthey, Luc	PD12	133
Paudyal, Durga	JJ01	106
Paulose, P. L.	PC11, QD16	128, 199
Paulose, P.I.	QB14	189
Pautrat, Alain	DG02, EH03	61, 70
Pavlova, Ekaterina	SG17	342
Pavlukhina, Oksana	PI20	154
Pecharsky, Vitalij	JJ01	106
Peddis, Davide	RN07	297
Pedersen, Bjoern	QH24	216
Pedersen, Kasper Steen	BC02	42
Pedrero, Louis	SF05	337
Pedrero, Luis	AB02	34
Peets, Darren	102	97
Pei, Wenli	SN10	365
Peiro, Julian	Al04, CG03	39, 53
Pelegrini, Fernando	PL09, QM14, QM15	165, 233, 234
Pelennen, Anatoliy P.	SO19	369
Pelevin, Ivan	SN09	365
	5,105	505

Pena, Luis	IH01, SO18	95, 369
Pena, Octavio	QA06	183
Penaranda-foix, F L	RQ10	308
Penaranda-foix, Felipe L	RO08	302
Pepe, Giulio	IJ03	97
Pepin, Catherine	PD04	132
Peprah, Marcus K.	KB03	107
Peratheepan, P	SE03, SE05	333, 334
Pereira, Andre M	SO31	371
Perez - Mato, Juan Manuel	PA13	118
Perez, Nicolas	RN14	299
Perez, Olivier	EH03	70
Perkert, Sandra	AG04	38
Perov, Nikolai	CH01, PN22, QO14	54, 175, 243
Perring, Toby G.	HA04	83
Perry, Robin	EA04	64
Persson, Jrg	QA23	186
Perucchi, Andrea	SA08	317
Peters, Robert	DB05	57
Peterson, Joe	RE09	267
Petit, S.	QF18	206
Petit, Sylvain	AA03, DD05, JB05, RA11	33, 59, 100, 251
Petrakovskii, G. A.	SG20	342
Petrenko, O. A.	CC04	50
Petrenko, Oleg A	BF03 GF03	45 79
Petrenko, Oleg A. Petrova, Elena A	PN22	175
Petrovic, C.	QC23	196
Petti, Daniela	AI03	39
Petzold, Vivien	QD07	197
	Q.5 07	157
Pfleiderer, C.	PI09	152
Pfleiderer, C. Pfleiderer, Christian	PI09 DF02, GC04, QH12,	152 60, 77, 214, 216,
Pfleiderer, C. Pfleiderer, Christian	Pl09 DF02, GC04, QH12, QH24, SF05	152 60, 77, 214, 216, 337
	DF02, GC04, QH12,	60, 77, 214, 216,
Pfleiderer, Christian	DF02, GC04, QH12, QH24, SF05	60, 77, 214, 216, 337
Pfleiderer, Christian Pham, Duc-thang	DF02, GC04, QH12, QH24, SF05 RA10	60, 77, 214, 216, 337 251
Pfleiderer, Christian Pham, Duc-thang Pham, Huyen-yen	DF02, GC04, QH12, QH24, SF05 RA10 SO12	60, 77, 214, 216, 337 251 368
Pfleiderer, Christian Pham, Duc-thang Pham, Huyen-yen Pham, Huyen-yen Duc	DF02, GC04, QH12, QH24, SF05 RA10 SO12 RA10	60, 77, 214, 216, 337 251 368 251
Pfleiderer, Christian Pham, Duc-thang Pham, Huyen-yen Pham, Huyen-yen Duc Pham, Thang D.	DF02, GC04, QH12, QH24, SF05 RA10 SO12 RA10 QL16	60, 77, 214, 216, 337 251 368 251 230
Pfleiderer, Christian Pham, Duc-thang Pham, Huyen-yen Pham, Huyen-yen Duc Pham, Thang D. Pham, The-tan	DF02, GC04, QH12, QH24, SF05 RA10 SO12 RA10 QL16 SO12	60, 77, 214, 216, 337 251 368 251 230 368
Pfleiderer, Christian Pham, Duc-thang Pham, Huyen-yen Pham, Huyen-yen Duc Pham, Thang D. Pham, The-tan Pham, Van-ben	DF02, GC04, QH12, QH24, SF05 RA10 SO12 RA10 QL16 SO12 SG10	60, 77, 214, 216, 337 251 368 251 230 368 341
Pfleiderer, Christian Pham, Duc-thang Pham, Huyen-yen Pham, Huyen-yen Duc Pham, Thang D. Pham, Thang D. Pham, The-tan Pham, Van-ben Phan, Huong Manh Phan, Manh-huong Phan, T. L.	DF02, GC04, QH12, QH24, SF05 RA10 SO12 RA10 QL16 SO12 SG10 RN19 DJ01 RA01	60, 77, 214, 216, 337 251 368 251 230 368 341 299
Pfleiderer, Christian Pham, Duc-thang Pham, Huyen-yen Pham, Huyen-yen Duc Pham, Thang D. Pham, The-tan Pham, Van-ben Phan, Huong Manh Phan, Manh-huong Phan, T. L. Phase, D.m	DF02, GC04, QH12, QH24, SF05 RA10 SO12 RA10 QL16 SO12 SG10 RN19 DJ01 RA01 SA14	60, 77, 214, 216, 337 251 368 251 230 368 341 299 63 249 318
Pfleiderer, Christian Pham, Duc-thang Pham, Huyen-yen Pham, Huyen-yen Duc Pham, Thang D. Pham, The-tan Pham, Van-ben Phan, Van-ben Phan, Huong Manh Phan, Manh-huong Phan, T. L. Phase, D.m Philip, Geo George	DF02, GC04, QH12, QH24, SF05 RA10 SO12 RA10 QL16 SO12 SG10 RN19 DJ01 RA01 SA14 SL28	60, 77, 214, 216, 337 251 368 251 230 368 341 299 63 249 318 360
Pfleiderer, Christian Pham, Duc-thang Pham, Huyen-yen Pham, Huyen-yen Duc Pham, Thang D. Pham, The-tan Pham, Van-ben Phan, Van-ben Phan, Huong Manh Phan, Manh-huong Phan, T. L. Phase, D.m Philip, Geo George Phillip, Geo George	DF02, GC04, QH12, QH24, SF05 RA10 SO12 RA10 QL16 SO12 SG10 RN19 DJ01 RA01 SA14 SL28 QA25	60, 77, 214, 216, 337 251 368 251 230 368 341 299 63 249 318 360 186
Pfleiderer, Christian Pham, Duc-thang Pham, Huyen-yen Pham, Huyen-yen Duc Pham, Thang D. Pham, The-tan Pham, Van-ben Phan, Van-ben Phan, Huong Manh Phan, Manh-huong Phan, T. L. Phase, D.m Philip, Geo George Phillip, Geo George Phillips, P.	DF02, GC04, QH12, QH24, SF05 RA10 SO12 RA10 QL16 SO12 SG10 RN19 DJ01 RA01 SA14 SL28 QA25 GA01	60, 77, 214, 216, 337 251 368 251 230 368 341 299 63 249 318 360 186 75
Pfleiderer, Christian Pham, Duc-thang Pham, Huyen-yen Pham, Huyen-yen Duc Pham, Thang D. Pham, The-tan Pham, Van-ben Phan, Van-ben Phan, Huong Manh Phan, Manh-huong Phan, T. L. Phase, D.m Philip, Geo George Phillip, Geo George Phillip, P. Phromsuwan, Udomchok	DF02, GC04, QH12, QH24, SF05 RA10 SO12 RA10 QL16 SO12 SG10 RN19 DJ01 RA01 SA14 SL28 QA25 GA01 RP18	60, 77, 214, 216, 337 251 368 251 230 368 341 299 63 249 318 360 186 75 306
Pfleiderer, Christian Pham, Duc-thang Pham, Huyen-yen Pham, Huyen-yen Duc Pham, Thang D. Pham, The-tan Pham, Van-ben Phan, Van-ben Phan, Huong Manh Phan, Manh-huong Phan, T. L. Phase, D.m Philip, Geo George Philip, Geo George Philips, P. Phromsuwan, Udomchok Piamonteze, Cinthia	DF02, GC04, QH12, QH24, SF05 RA10 SO12 RA10 QL16 SO12 SG10 RN19 DJ01 RA01 SA14 SL28 QA25 GA01 RP18 BC02	60, 77, 214, 216, 337 251 368 251 230 368 341 299 63 249 318 360 186 75 306 42
Pfleiderer, Christian Pham, Duc-thang Pham, Huyen-yen Pham, Huyen-yen Duc Pham, Thang D. Pham, The-tan Pham, Van-ben Pham, Van-ben Phan, Huong Manh Phan, Manh-huong Phan, T. L. Phase, D.m Philip, Geo George Philip, Geo George Philips, P. Phromsuwan, Udomchok Piamonteze, Cinthia Piantek, Marten	DF02, GC04, QH12, QH24, SF05 RA10 SO12 RA10 QL16 SO12 SG10 RN19 DJ01 RA01 SA14 SL28 QA25 GA01 RP18 BC02 PG08	60, 77, 214, 216, 337 251 368 251 230 368 341 299 63 249 318 360 186 75 306 42 143
Pfleiderer, Christian Pham, Duc-thang Pham, Huyen-yen Pham, Huyen-yen Duc Pham, Thang D. Pham, The-tan Pham, Van-ben Pham, Van-ben Phan, Huong Manh Phan, Manh-huong Phan, T. L. Phase, D.m Philip, Geo George Philip, Geo George Philips, P. Phromsuwan, Udomchok Piamonteze, Cinthia Piantek, Marten Piao, Hong-guang	DF02, GC04, QH12, QH24, SF05 RA10 SO12 RA10 QL16 SO12 SG10 RN19 DJ01 RA01 SA14 SL28 QA25 GA01 RP18 BC02 PG08 SH01, SH10	60, 77, 214, 216, 337 251 368 251 230 368 341 299 63 249 318 360 186 75 306 42 42 143 343, 344
Pfleiderer, Christian Pham, Duc-thang Pham, Huyen-yen Pham, Huyen-yen Duc Pham, Thang D. Pham, The-tan Pham, Van-ben Phan, Huong Manh Phan, Huong Manh Phan, Manh-huong Phan, T. L. Phase, D.m Philip, Geo George Philip, Geo George Philip, Geo George Philip, Geo George Philip, P. Phromsuwan, Udomchok Piamonteze, Cinthia Piantek, Marten Piao, Hong-guang Picone, Andrea	DF02, GC04, QH12, QH24, SF05 RA10 SO12 RA10 QL16 SO12 SG10 RN19 DJ01 RA01 SA14 SL28 QA25 GA01 RP18 BC02 PG08 SH01, SH10 PL21	60, 77, 214, 216, 337 251 368 251 230 368 341 299 63 249 318 360 186 75 306 42 143 343, 344 167
Pfleiderer, Christian Pham, Duc-thang Pham, Huyen-yen Pham, Huyen-yen Duc Pham, Thang D. Pham, The-tan Pham, Van-ben Pham, Van-ben Phan, Manh-huong Phan, T. L. Phase, D.m Philip, Geo George Philip, Geo George Philip, Geo George Philip, Geo George Philips, P. Phromsuwan, Udomchok Piamonteze, Cinthia Piantek, Marten Piao, Hong-guang Picone, Andrea Picozzi, Silvia	DF02, GC04, QH12, QH24, SF05 RA10 SO12 RA10 QL16 SO12 SG10 RN19 DJ01 RA01 SA14 SL28 QA25 GA01 RP18 BC02 PG08 SH01, SH10 PL21 AA02, PA12	60, 77, 214, 216, 337 251 368 251 230 368 341 299 63 249 318 360 186 75 306 42 143 343, 344 167 33, 117
Pfleiderer, Christian Pham, Duc-thang Pham, Huyen-yen Pham, Huyen-yen Duc Pham, Thang D. Pham, The-tan Pham, Van-ben Pham, Van-ben Phan, Manh-huong Phan, T. L. Phase, D.m Philip, Geo George Philip, Geo George Philip, Geo George Philip, Geo George Philips, P. Phromsuwan, Udomchok Pianonteze, Cinthia Piantek, Marten Piao, Hong-guang Picone, Andrea Picozzi, Silvia	DF02, GC04, QH12, QH24, SF05 RA10 SO12 RA10 QL16 SO12 SG10 RN19 DJ01 RA01 SA14 SL28 QA25 GA01 RP18 BC02 PG08 SH01, SH10 PL21 AA02, PA12 QB30	60, 77, 214, 216, 337 251 368 251 230 368 341 299 63 249 318 360 186 75 306 42 143 343, 344 167 33, 117 192
Pfleiderer, Christian Pham, Duc-thang Pham, Huyen-yen Pham, Huyen-yen Duc Pham, Thang D. Pham, The-tan Pham, Van-ben Pham, Van-ben Phan, Manh-huong Phan, T. L. Phase, D.m Philip, Geo George Philip, Geo George Philip, Geo George Philip, Geo George Philip, Geo George Philip, P. Phromsuwan, Udomchok Piamonteze, Cinthia Piantek, Marten Piao, Hong-guang Picone, Andrea Picozzi, Silvia Pietosa, Jaroslaw Pikul, Adam	DF02, GC04, QH12, QH24, SF05 RA10 SO12 RA10 QL16 SO12 SG10 RN19 DJ01 RA01 SA14 SL28 QA25 GA01 RP18 BC02 PG08 SH01, SH10 PL21 AA02, PA12 QB30 QD05	60, 77, 214, 216, 337 251 368 251 230 368 341 299 63 249 318 360 186 75 306 42 143 343, 344 167 33, 117 192 197
Pfleiderer, Christian Pham, Duc-thang Pham, Huyen-yen Pham, Huyen-yen Duc Pham, Thang D. Pham, Thang D. Pham, The-tan Pham, Van-ben Pham, Van-ben Phan, Manh-huong Phan, T. L. Phase, D.m Philip, Geo George Philip, Geo George Philip, Geo George Philip, Geo George Philips, P. Phromsuwan, Udomchok Piamonteze, Cinthia Piantek, Marten Piao, Hong-guang Picone, Andrea Picozzi, Silvia Pietosa, Jaroslaw Pikul, Adam Pikul, Valeriy	DF02, GC04, QH12, QH24, SF05 RA10 SO12 RA10 QL16 SO12 SG10 RN19 DJ01 RA01 SA14 SL28 QA25 GA01 RP18 BC02 PG08 SH01, SH10 PL21 AA02, PA12 QB30 QD05 RR28	60, 77, 214, 216, 337 251 368 251 230 368 341 299 63 249 318 360 186 75 306 42 143 343, 344 167 33, 117 192 197 315
Pfleiderer, Christian Pham, Duc-thang Pham, Huyen-yen Pham, Huyen-yen Duc Pham, Thang D. Pham, Thang D. Pham, The-tan Pham, Van-ben Pham, Van-ben Phan, Manh-huong Phan, T. L. Phase, D.m Philip, Geo George Philip, Geo George Philip, Geo George Philip, Geo George Philip, Geo George Philip, Geo George Philip, P. Phromsuwan, Udomchok Piamonteze, Cinthia Piantek, Marten Piao, Hong-guang Picone, Andrea Picozzi, Silvia Pietosa, Jaroslaw Pikul, Adam Pikul, Valeriy Piligkos, Stergios	DF02, GC04, QH12, QH24, SF05 RA10 SO12 RA10 QL16 SO12 SG10 RN19 DJ01 RA01 SA14 SL28 QA25 GA01 RP18 BC02 PG08 SH01, SH10 PL21 AA02, PA12 QB30 QD05 RR28 BC02	60, 77, 214, 216, 337 251 368 251 230 368 341 299 63 249 318 360 186 75 306 42 143 343, 344 167 33, 117 192 197 315 42
Pfleiderer, Christian Pham, Duc-thang Pham, Huyen-yen Pham, Huyen-yen Duc Pham, Thang D. Pham, Thang D. Pham, The-tan Pham, Van-ben Pham, Van-ben Phan, Manh-huong Phan, T. L. Phase, D.m Philip, Geo George Philip, Geo George Philip, Geo George Philip, Geo George Philip, Geo George Philip, Geo George Philip, P. Phromsuwan, Udomchok Piamonteze, Cinthia Piantek, Marten Piao, Hong-guang Picone, Andrea Picozzi, Silvia Pietosa, Jaroslaw Pikul, Valeriy Piligkos, Stergios Pineider, Francesco	DF02, GC04, QH12, QH24, SF05 RA10 SO12 RA10 QL16 SO12 SG10 RN19 DJ01 RA01 SA14 SL28 QA25 GA01 RP18 BC02 PG08 SH01, SH10 PL21 AA02, PA12 QB30 QD05 RR28 BC02 CB04, SI11	60, 77, 214, 216, 337 251 368 251 230 368 341 299 63 249 318 360 186 75 306 42 143 343, 344 167 33, 117 192 197 315
Pfleiderer, Christian Pham, Duc-thang Pham, Huyen-yen Pham, Huyen-yen Duc Pham, Thang D. Pham, Thang D. Pham, The-tan Pham, Van-ben Pham, Van-ben Phan, Manh-huong Phan, T. L. Phase, D.m Philip, Geo George Philip, Geo George Philip, Geo George Philip, Geo George Philip, Geo George Philip, Geo George Philip, P. Phromsuwan, Udomchok Piamonteze, Cinthia Piantek, Marten Piao, Hong-guang Picone, Andrea Picozzi, Silvia Pietosa, Jaroslaw Pikul, Adam Pikul, Valeriy Piligkos, Stergios	DF02, GC04, QH12, QH24, SF05 RA10 SO12 RA10 QL16 SO12 SG10 RN19 DJ01 RA01 SA14 SL28 QA25 GA01 RP18 BC02 PG08 SH01, SH10 PL21 AA02, PA12 QB30 QD05 RR28 BC02	60, 77, 214, 216, 337 251 368 251 230 368 341 299 63 249 318 360 186 75 306 42 143 343, 344 167 33, 117 192 197 315 42 49, 347
Pfleiderer, Christian Pham, Duc-thang Pham, Huyen-yen Pham, Huyen-yen Duc Pham, Thang D. Pham, Thang D. Pham, The-tan Pham, Van-ben Pham, Van-ben Phan, Huong Manh Phan, Manh-huong Phan, T. L. Phase, D.m Philip, Geo George Philip, R. Phromsuwan, Udomchok Piamonteze, Cinthia Piantek, Marten Piao, Hong-guang Picone, Andrea Picozzi, Silvia Pietosa, Jaroslaw Pikul, Valeriy Piligkos, Stergios Pineider, Francesco Piquini, Lucas K.	DF02, GC04, QH12, QH24, SF05 RA10 SO12 RA10 QL16 SO12 SG10 RN19 DJ01 RA01 SA14 SL28 QA25 GA01 RP18 BC02 PG08 SH01, SH10 PL21 AA02, PA12 QB30 QD05 RR28 BC02 CB04, SI11 QO03	60, 77, 214, 216, 337 251 368 251 230 368 341 299 63 249 318 360 186 75 306 42 143 343, 344 167 33, 117 192 197 315 42 49, 347 241

Disease C	PC02	45
Pisana, S. Piskupaua, Nataku kanauna	BG02	45
Piskunova, Nataly Ivanovna Pistora, Jaromir	PF01 CH03, QH19	139 54, 216
Pitcher, Michael. J.	QG19	211
Pizzini, Stefania	RL24, SH03, SH05, SH06	293, 343, 344
Plazaola, Fernando	QH16, RN13	215, 298
Pleiderer, Christian	QC17	195
Plumb, N. C.	PC01	126
Poddar, Asok	RB15	256
Podgornykh, S. Mikhailovich	PC07	127
Podolskii, Vitaliy	SG17	342
Pogorelov, Yu. G.	KC05	109
Pogossian, Souren P	SA24	320
Poienar, Maria	JB05, RA11	100, 251
Pokharel, Ramesh K.	PJ01	154
Polak, Christian	PM04	168
Politi, Paolo	KE01	110
Polyakov, Andrey	PD05	132
Polyakov, Victor	SO28	371
Polyakova, Klaudia	SO28	371
Pomar, Alberto	SO18	369
Pomjakushin, Vladimir	PI05, PO17	151, 180
Pomjakushin, Vladimir Yu.	PH02	147
Pomjakushina, E.	QF23	207
Pomjakushina, Ekaterina	IJ06, PI21, PO17	98, 154, 180
Pong, W. F.	PB12	122
Ponomaryov, A. N.	QF13	206
Ponomaryov, Alexey	RP13	305
Pons, Stephane	SE08	334
Pop, Viorel	EJ03	72
Popkov, Sergei	PI11	152
Popova, Marina N.	PA20, RA24	119, 253
Porer, Michael	HA02	83
Portemont, Celine	AI04	39
Porter, Daniel Graham	DH03	62
Pospisil, Jiri	HJ04, PC10, QC21, QJ02	90, 128, 195, 221
Postava, Kamil	CH03	54
Potapova, Nadezhda M.	DF04	60
Potthoff, Michael	PE05, PE07	136, 137
Poudyal, Narayan	EJO2	71 71
Pouget, Stephanie	EIO4	
Poulopoulos, Panagiotis Pourovskii, Leonid	RP15 AB04	305 34
Prabhakaran, D.	AB04 DH03, PI03, QN01	34 62, 151, 236
Prabhakaran, D. Prabhakaran, Dharmalingam	AF01	36
Pradhan, N. S.	PN01	172
Pramuji, Hinu	BJ03	48
Prasad, Arun S.	SL25	360
Praveena, K	GJ05, QN18, SN08	83, 239, 365
Prchal, J.	RD26	266
Prchal, Jiri	BB04, PD18, PD19, QJ14	
Prestat, Eric	Al04	39
Prestigiacomo, Joseph	CH01	54
Prevost, B.	PD05	132
Prida, Victor	DG03	61
Prieto, Pedro	SA12	318
Primo, Fernando Lucas	PO08	178
Prinsloo, Aletta Roletta	RH15	277
Pristas, Gabriel	QD12, QI02, RE13	198, 218, 268
Profeta, Gianni	QG25	212
	2020	£12

Droghado Madimir V	SK01	350
Proglyado, Vladimir V. Proglyado, Vyacheslav	SK02	350
Prokes, Karel	QG26, RB01, RB07,	212, 254, 255, 256,
Tronco, runci	RB17, RQ21	310
Prokhnenko, O.	RQ21	310
Prokleska, Jan	BB04, HJ04, PC10,	42, 90, 128, 195,
	QC21, QJ02, SB08	221, 321
Prokofiev, A.	CD01	50
Prokscha, Thomas	PB20, PN23, QL19	123, 175, 231
Proselkov, O.	BE02	44
Proselkov, Oleg	DD02	58
Proshin, Yurii N.	QA26, SC25, SC26	186, 329
Proshkin, Alexey	QG04	208
Proshkin, Alexey V.	Q108	219
Proskurin, Igor	DF03	60
Proskurina, Ekaterina	QG04	208
Protsenko, I E	RN12	298
Prozorov, R.	IA05, QC23, SB30	91, 196, 325
Prsa, Krunoslav	PI06, QF20, QF23, SE08	151, 207, 334
Prudnikov, Valerii	CH01	54
Pruschke, Thomas	FE03, PE03	75, 136
Przybylski, Marek	AD02, AE01, El01	35, 70
Przysłupski, Piotr	QM17	234
Puente Orench, Ines Puente-orench, Ines	KB06 JJ04	108 106
Pujari, P K	PM01	167
Pureur, Paulo	PB09, PB10, PB33, PG03,	
	RB09	121, 123, 142, 233
Putri, Witha B. K.	PB03	120
Puzniak, Roman	QB30, SB26	192, 324
Pyshkin, Pavel V	PJ02	155
•		
Q		
Qamar Ul Hassan, Syed	SA15	318
Qian, Bin	RB12	256
Qiu, Xue Peng	BD05	43
Qiu, Xuepeng	PJ20, PJ21, RL19	158, 293
Qiu, Z. Q.	IH05	96
Qiu, Zq.	BD03, QM26	43, 235
Quan, Zhiyong	PN04	172
Quang, Pham Hong	SL22	359
Quang, Tran Van	RH20	278
Quintanilla, Jorge	JA05	99
Quintero-castro, Diana	RH24	278
Quitmann, Christoph Qureshi, Navid	RJ11 PI08	284 152
Quiesiii, Naviu	FIUO	152
R		
R, Ramasamy	PO21, QN26	180, 240
Raabe, Joerg	RJ11	284
Radaelli, Greta	DE03	59
Radovic, M.	PC01	126
Radu, Florin	EC04	66
Rafique, Mohsin		318
	SA15	
Rafiqul Islam, Md.	RN05	297
Raghavender, A T	RN05 SM07	297 362
Raghavender, A T Raghavender, Anupati Telugu	RN05 SM07 SM10	297 362 362
Raghavender, A T Raghavender, Anupati Telugu Rahman, Gul	RN05 SM07 SM10 RM02	297 362 362 295
Raghavender, A T Raghavender, Anupati Telugu	RN05 SM07 SM10	297 362 362

Ramakrishnan, S.	OB1/	189
	QB14	
Ramanathan, Mahendiran	QA18	185
Ramanujan, Raju V	QL02	228
Ramasamy, Shivaraman	PO21	180
Ramaswamy, Shivaraman	PO24, QN26, QP24,	181, 240, 249, 319,
	SA21, SA22, SA23, SK32, SL28	320, 355, 360
Ramesh, Palanisamy	PN16	174
Ramos, Scheilla Maria	SE22	336
Ramulu, T.s.	SL22	359
Ran, S.	GA01	75
Rana, Bivas	RK11	287
Rangappa, Dinesh	RN11	298
Rani, Manju	SC02	325
Ranjan, Vishal	GI05	82
Ranjbar, Fereshte	RN22	300
Ranot, Mahipal	SB28	324
Ranzieri, P.	AG02	37
Rao, G.s.n.	BJ05	48
Rao, Siddharth	PJ08, SI02	156, 346
Rashad, Mohamed M.	RN01	296
Rasing, Theo	GD01, RK17	77, 288
Rasly, M.	RN01	296
Rawat, Rajdeep Singh	PN03, QL02	172, 228
Ray, R.	PL07	165
Raychaudhuri, Arup Kumar	HD04	85
Razzoli, E.	PC01	126
Rebizant, Jean	SE19	336
Reddy, Leelakrishna	RH15	277
Redinger, Josef	EF05	68
Regnat, Alexander	QC17, QH24	195, 216
Regnault, L. P.	QE19	202
Reichel, Franz	DE01	59
Reichlova, Helena	BE05	44
Reiffers, Marian	QC13, QD08, QD12,	194, 197, 198, 218,
nemers, manan	QI02, RD21	265
Reimann, Tommy	SM02	361
Reiner, J.	BG02	45
Reisner, G. M.	PD03	132
Reiss, Gunter	IF04	94 215 255
Reissner, Michael Ren, Chung-yuan	QH18, RB08	215, 255
	QE25	203
Ren, Wei	QA27	187
Ren, Xiaobing	JJ02, PG21	106, 145
Ren, Zhi	RC11	260
Repaka, D V Maheswar	QA18	185
Repi, V. Vekky R.	QN16	239
Ressouche, Eric	IJ03, IJ06, PA13	97, 98, 118
Reyes, Arneil P.	JA02	98
Reyren, Nicolas	HE04	86
Rezania, Hamed	QF12	205
Rhee, Chan Hyuk	PI13	152
Rhee, J. Y.	SL08, SL09	357
Rhee, Jang-rho	PL11	165
Rhee, Joo Yull	QH29, SC08	217, 326
Rhie, K.	PJ31	159
Rhie, Kungwon	PK21, QK10	163, 225
Rhie, Kung-won	SD12	332
Rhim, Sonny S. H.	RB26	258
Rho, Min-suk	BH02	46

Rhyee, Jong- Soo	RC05	259
Rhyee, Jong-soo	CF03	52
Ri, Hyeong-cheol	SB22	323
Ribeiro, Andre Neves	PF13	141
Ribeiro, B.	QM13	233
Ribeiro, Raquel A	SL20	359
Ribeiro, Raquel A.	Q003	241
Richter, Curt A.	AE04	36
Rijnders, Guus	DH02, SA10	62, 317
Rikken, Geert	RQ16	309
Rinaldi, Christian	AI03	39
Riseborough, Peter S	PD02, PE01	132, 136
Ritter, C	QJ16	223
Ritter, Clemens	U01	97
Ritter, Franz	QE14	201
Ritz, Robert	GC04	77
Ritzmann, Ulrike	FA02	72
Riva, Michele	PL21	167
Rivas, J.	CB02	49
Riyadi, Syarif	PH17	149
Rizal, Conrad	PJ01	154
Rizki, Youssef	BB03, RH22	41, 278
Rizwan, Syed	SA19, SA20	319
Robert, Julien	JB05, RA11	100, 251
Rodionov, Igor	CH01	54
Rodionova, Valeria	AJ02, DD04, SO04, SO27	40, 59, 367, 371
Rodmacq, Bernard	BD02	43
Rodriguez Fernandez, J.	GH05	81
Rodriguez Fernandez, Jesus	JJ04	106
Rodriguez, Luis Alfredo	DG05, GG05, SM16	61, 80, 363
Rodriguez-blanco, Clara	GF04	79
R. velamazan, Jose Alberto	GF04, PG22	79, 146
Roehlsberger, Ralf	QM19	234
Roemer, Florian	RP08	304
Roessler, Sahana	PH04	147
Rogalev, A.	QG07	209
Rogalev, Andrei	PA08, RP15, RP21	117, 305, 306
Rogalev, Andrey	PN22	175
Roger, Michel	DH03	62
Rogl, Peter	SC01	325
Rojas Sanchez, Juan-carlos	QK16, QK17, SL24	226, 360
Rojas, D. P.	QD10	198
Rojas, Juan-carlos	AI04	39
Rojas-sanchez, Juan Carlos	QK15, RL24	226, 293
Rokkaku, Tsugumi	QG13	210
Romanov, A	BH01	46
Romanova, Oksana	QG05	209
Ronning, F	EA03	64
Ronning, F.	FB01, QB12	73, 189
Ronning, Filip	RD09, SB12, SC16	263, 322, 328
Ronnow, H.	SF10	338
Ronnow, Henrik M	QF23	207
Ronnow, Henrik Moodysson	IC02	92
Rosales Rivera, Andres	SM04	361
Rosch, Achim	DF02	60 154
Roshanmehr, M.	PI24	154
Roshchin, I. V. Roshakov, Ilva	BD01	43
Roslyakov, Ilya Rospor H	SL27	360
Rosner, H. Rosner, Helge	QF22 AB02, QD07, RD12	207 34, 197, 264
Noshel, Helye	ADUZ, QUUI, INDIZ	54, 197, 204

Ross, Caroline A	HP51
Ross, Caroline A.	KC01
Roßner, Helge	FB04
Roth, Georg	PA14, RP22
Roth, Thomas	HG04, PA24, RH23
Rother, Niklas	PE05
Rothman, Johan	DD01
Rotter, Martin	PH04
Rougemaille, Nicolas Roulland, F.	SH03, SH05, SH06
Roulland, F. Roulland, Francois	SA06 SA05
	RA11
Rouquette, Jerome Rovezzi, Mauro	BE03
Rovira, Joan Josep Roa	PB11, SB18
Rowley, Stephen	AB04
Roy, Beas	QE12
Roy, M.	PN01, RA03
Roy, Pierre	IE03
Rubacheva, Anastasia D	PN22
Ruben, Mario	GH02
Ruck. Michael	PC27
Rudajevova, Alexandra	PD18
Rueffer, Rudolf	HG04
Ruegg, Ch	RIO5
Ruegg, Ch.	OE19
Ruegg, Christian	IC02, QF23, QG14
Rufinus, Jeffrey	SK25
Rugi, Francesco	CB04
Rui, Wenbing	PL01
Ruotolo, Antonio	DH04, RJ03
Rusek, Pawel	RI08
Rusek, Paweł	QE24
Rushforth, Andrew W	SH09
Rusponi, Stefano	BC02
Rusu, Calin	SE08
Ruzicka, Jiri	BH03
Ryabinkina, Ludmila	QG05
Ryll, H	RI05
Ryoichi, Ito	RP06
Ryu, Jisu	HP21, KE03, PJ33
Ryu, Jungho	DJ02
Ryu, Kwang-su	SM01
Ryzhov, V. A.	RQ13
Ryzhov, V. A. S	RQ13
	RQ13 SA21
S	·
<b>S</b> S, Anuraj	SA21
<b>S</b> S, Anuraj S, Radha	SA21 RN16, SO22
<b>S</b> S, Anuraj S, Radha S, Srinath	SA21 RN16, SO22 QD04, RA16
<b>S</b> S, Anuraj S, Radha S, Srinath S. Suzuki, Hiroyuki	SA21 RN16, SO22 QD04, RA16 QD18
<b>S</b> S, Anuraj S, Radha S, Srinath S. Suzuki, Hiroyuki Saadaoui, Hassan	SA21 RN16, SO22 QD04, RA16 QD18 PB20, QL19
<b>S</b> S, Anuraj S, Radha S, Srinath S. Suzuki, Hiroyuki Saadaoui, Hassan Sablina, Klara	SA21 RN16, SO22 QD04, RA16 QD18 PB20, QL19 PI11
<b>S</b> S, Anuraj S, Radha S, Srinath S. Suzuki, Hiroyuki Saadaoui, Hassan Sablina, Klara Sabyasachi, Sk.	SA21 RN16, SO22 QD04, RA16 QD18 PB20, QL19 PI11 RH02
<b>S</b> S, Anuraj S, Radha S, Srinath S. Suzuki, Hiroyuki Saadaoui, Hassan Sablina, Klara Sabyasachi, Sk. Sachdev, Subir	SA21 RN16, SO22 QD04, RA16 QD18 PB20, QL19 PI11 RH02 IA03
<b>S</b> S, Anuraj S, Radha S, Srinath S. Suzuki, Hiroyuki Saadaoui, Hassan Sablina, Klara Sabyasachi, Sk. Sachdev, Subir Sacramento, Pedro	SA21 RN16, SO22 QD04, RA16 QD18 PB20, QL19 PI11 RH02 IA03 RB25
<b>S</b> S, Anuraj S, Radha S, Srinath S. Suzuki, Hiroyuki Saadaoui, Hassan Sablina, Klara Sabyasachi, Sk. Sachdev, Subir Sacramento, Pedro Sacramento, Pedro D	SA21 RN16, SO22 QD04, RA16 QD18 PB20, QL19 PI11 RH02 IA03 RB25 SB21
<b>S</b> S, Anuraj S, Radha S, Srinath S. Suzuki, Hiroyuki Saadaoui, Hassan Sablina, Klara Sabjina, Klara Sabyasachi, Sk. Sachdev, Subir Sacramento, Pedro Sacramento, Pedro D Sacuto, Alain	SA21 RN16, SO22 QD04, RA16 QD18 PB20, QL19 P111 RH02 IA03 RB25 SB21 PD17
<b>S</b> S, Anuraj S, Radha S, Srinath S. Suzuki, Hiroyuki Saadaoui, Hassan Sablina, Klara Sabjasachi, Sk. Sachdev, Subir Sacramento, Pedro Sacramento, Pedro D Sacuto, Alain Sadhana, K	SA21 RN16, SO22 QD04, RA16 QD18 PB20, QL19 PI11 RH02 IA03 RB25 SB21 PD17 GJ05, QN18, SN08
<b>S</b> S, Anuraj S, Radha S, Srinath S. Suzuki, Hiroyuki Saadaoui, Hassan Sablina, Klara Sabjasachi, Sk. Sachdev, Subir Sacramento, Pedro Sacramento, Pedro D Sacuto, Alain Sadhana, K Sadowski, J.	SA21 RN16, SO22 QD04, RA16 QD18 PB20, QL19 P111 RH02 IA03 RB25 SB21 PD17 GJ05, QN18, SN08 BE02

31

108

73

136

58

147

317

317

251

44

34

201

93

175

81

131

134

88

279

202

354

49

164

280

203

344

42

334

46

209

279

304

63

361

309

319

199

152

275

258

323

134

44

240

146, 243

83, 239, 365

91

299, 370

197, 252

123, 231

30, 110, 160

62, 283

92, 207, 210

121, 323

172, 249

343, 344

118, 306

88, 119, 278

Sagar, J.         SK03         350           Sagar, James         SK11         352           Sagara, Naoto         QP16         248           Sagara, Naoto         QP16         289           Sagara, Naoto         QP16         289           Sagara, Harunobu         SC24         329           Sahade, Nishaina         CC05         361           Sahade, Nishaina         SC06         326           Sahade, Nishaina         CE05, P111, P110, SM03         52, 156, 165, 361           Sahade, Nishaina         CE05, P111, P110, SM03         52, 156, 165, 361           Sahon, Roshnee         QH06         213           Sahon, Roshnee         QH04         267           Saika, Nobutaka         PB08         121           Saika, Nuskate         QB04, RB02         187, 254           Saine, Paul         PC02         127           Sain, Naurang L.         RD28         258           Saito, Hidekazu         A05         39           Saito, Kotano         RD14         264           Saito, Kotano         RD14         264           Saito, Kotano         RD14         264           Saito, T         QE20, QE15         137, 112, 22	Safaeizadeh, Elham	PO27, PO29	181, 182
Sagara, Naoto         OP16         248           Sagawa, Harumobu         SC24         329           Sagawa, Harumobu         SC24         329           Saha, Sumita         HD04         85           Sahak, Sumita         HD04         85           Sahadevan, Ajeesh         GE03         78           Sahadevan, Ajeesh         GE03         78           Sahadevan, Ajeesh         GE03         213           Sahadovan, Ajeesh         GE03         2156, 165, 361           Sahac, Roshnee         QH06         213           Sahac, Noshnee         QH06         213           Sahat, Nobutaka         PB08         121           Saika, Naoka         QB07, R19         260, 262           Saito, Tikeka	Sagar, J.	SK03	350
Sagawa, Harunobu         SC24         329           Sagayama, Hajime         PG14, PG20, QA13         144, 145, 175           Saha, S. R.         QB07         188           Saha, S. R.         QB07         188           Sahade, Nishiana         SC06         326           Sahade, Nishiana         SC05         351           Sahade, Nishiana         CE05, P111, PL10, SM03         52, 156, 165, 361           Saho, Subrata Kumar         BH04         46           Saidaoui, H. B. Mohamed         P113         156           Saiki, Nubutaka         PB08         121           Saiki, Subuske         QB04, RB02         187, 254           Saiki, Naurang L.         RB28         258           Saiko, Seiya         RN03         297           Saito, Hitoshi         RP06         304           Saito, Kotano         RD14         264           Saito, Kotano         RD14         264           Saito, Mitsuru         QA17         185           Saito, Takashi         PE13, QC05         138, 193           Saito, Takashi         PD11         133           Saito, Takashi         PD13         134           Saito, Takashi         PD13	Sagar, James	SK11	352
Sagayama, Hajime         PG14, PG20, QA13         144, 145, 175           Saha, S. R.         QB07         188           Saha, Sumita         HD04         85           Sahadevan, Ajeesh         GE03         78           Sahadevan, Ajeesh         GE03         78           Sahadevan, Ajeesh         CE05, PJ11, PL10, SM03         52, 156, 165, 361           Sahoo, Roshnee         QH06         213           Saho, Subrata Kumar         BH04         46           Saidaou, H. B. Mohamed         P113         156           Saiga, Yuta         RE04         267           Saikai, Nobutaka         PB08         121           Saiki, Shunsuke         QB04, RB02         187, 254           Sains, Seiya         RN03         297           Saito, Hidekazu         A05         39           Saito, Kotaro         RD14         264           Saito, Kotaro         RD14         264           Saito, Mineo         RC09, RC19         260, 262           Saito, Takashi         PE13, QC05         138, 193           Saito, Takashi         PE13, QC05         138, 193           Saito, Takashi         PE13, QC11         202, 207           Saito, Takashi	Sagara, Naoto	QP16	248
Saha, S. R.         Q807         188           Saha, Susmita         HD04         85           Sahadevan, Ajeesh         GE03         78           Sahadevan, Ajeesh         SE05         351           Sahadevan, Ajeesh         SE05         351           Sahadiwan, Ajeesh         CE05, PJ11, PL10, SM03         52, 156, 165, 361           Sahao, Submashi         CE05, PJ11, PL10, SM03         52, 156, 165, 361           Sahao, Submashi         CE05, PJ11, PL10, SM03         52, 156, 165, 361           Sahao, Submashi         CE05, PJ11, PL10, SM03         52, 156, 165, 361           Sahao, Nasahi         CE05, PJ11, PL10, SM03         52, 156, 165, 361           Sahao, Nasahi         PE04         267           Saika, Nobutaka         PB08         121           Saiki, Shunsuke         Q804, RB02         187, 254           Saino, Fidekazu         A005         39           Saito, Hitoshi         RP06         304           Saito, Kesami         KC03         108           Saito, Kesami         KC03         188           Saito, Takashi         P11         133           Saito, Takashi         P11         133           Saito, Takashi         P11         133	Sagawa, Harunobu	SC24	329
Saha, Susmita         HD04         85           Sahadeva, Ajeesh         GE03         78           Sahadevan, Ajeesh         GE03         78           Sahadevan, Ajeesh         GE05         911           Sahadevan, Ajeesh         CE05, PJ11, PL10, SM03         52, 156, 165, 361           Sahashi, Masashi         CE05, PJ11, PL10, SM03         52, 156, 165, 361           Saha, Subrata Kumar         BH04         46           Saidaoui, H. B. Mohamed         PJ13         156           Saiki, Nobutaka         PB08         121           Saiki, Nobutaka         PB08         127           Saino, Seiya         RN03         297           Saito, Hitoshi         RP06         304           Saito, Hitoshi         RP06         304           Saito, Kotaro         RD14         264           Saito, Kotaro         RD14         264           Saito, T         QB28         191           Saito, Takashi         PE13, QC05         138, 193           Saito, Takashi         PE13, QC05         188, 193           Saito, Takashi         PE13, QC05         189           Saito, Takashi         PE13, QC05         189           Saito, Takashi <td< td=""><td>Sagayama, Hajime</td><td>PG14, PG20, QA13</td><td>144, 145, 175</td></td<>	Sagayama, Hajime	PG14, PG20, QA13	144, 145, 175
Sahadev, Nishaina         SCO6         326           Sahadevan, Ajeesh         GE03         78           Sahadevan, Ajeesh         SK05         351           Sahashi, Masashi         CEO5, P111, PL10, SM03         52, 156, 165, 361           Saho, Subrata Kumar         BH04         46           Saidaoui, H, B. Mohamed         P113         156           Saika, Nobutaka         PB08         121           Saiki, Shunsuke         QB04, RB02         187, 254           Sains, Shuarang L.         RB28         258           Saisho, Seiya         RN03         297           Saito, Hidekazu         A05         39           Saito, Hidekazu         A05         39           Saito, Kotano         RD14         264           Saito, Kotano         RD14         264           Saito, Kotano         RD14         264           Saito, Takashi         PD11         133           Saito, Takashi         QE02, QE15         187, 189           Saito, Takashi         PD11         133           Saito, Takashi         QE20, QF21         202, 207           Saito, Takashi         QE02, QE15         187, 189      Saito, Takashi         PD13         134	Saha, S. R.	QB07	188
Sahadevan, Ajeesh         GE03         78           Sahadevan, Ajeesh M         SK05         351           Sahashi, Masashi         CE05, PJ11, PL10, SM03         52, 156, 165, 361           Sahoo, Roshnee         QH06         213           Saho, Subrata Kumar         BH04         46           Saidaou, H. B. Mohamed         PJ13         156           Saiki, Nobutaka         PB08         121           Saiki, Shunsuke         QB04, RB02         187, 254           Saines, Paul         PC02         127           Saini, Naurang L.         RB28         258           Saito, Seiya         RN03         297           Saito, Hidekazu         A05         39           Saito, Kesami         KC03         108           Saito, Kesami         KC03         108           Saito, Kesami         KC03         108           Saito, Kesami         QA17         185           Saito, Kesami         QC05         193           Saito, Takashi         PE13, QC05         138, 193           Saito, Takashi         QE2, QF21         202, 207           Saito, Takashi         QE10, QF21         202, 207      Saito, Takashi         PE13, QC05         1	Saha, Susmita	HD04	85
Sahadevan, Ajeesh M         SK05         351           Sahashi, Masashi         CE05, PJ11, PL10, SM03         52, 156, 165, 361           Sahoo, Roshnee         OH06         213           Sahu, Subrata Kumar         BH04         46           Saidaoui, H. B. Mohamed         PJ13         156           Saikai, Nobutaka         PB08         121           Saiki, Shunsuke         QB04, RB02         187, 254           Saines, Paul         PC02         127           Saini, Naurang L.         RE28         258           Saiso, Seiya         RN03         297           Saito, Hidekazu         A05         39           Saito, Kesami         KC03         108           Saito, Kotaro         RD14         264           Saito, Mitsuru         QA17         185           Saito, Takashi         PE13, QC05         138, 193           Saito, Takashi         PD11         133           Saito, Takashi         QE20, QF21         202, 207           Saito, Takashi         PD13         134           Sakai, Yait         RD13         264           Sakai, Akito         CD02, SE01, SF07         51, 333, 338           Sakai, Akito         QU3	Sahadev, Nishaina	SC06	326
Sahashi, Masashi         CE05, P111, PL10, SM03         52, 156, 165, 361           Sahoo, Roshnee         QH06         213           Sahu, Subrata Kumar         BH04         46           Saidaoui, H. B. Mohamed         P113         156           Saika, Nobutaka         PB08         121           Saika, Nusuke         QB04, RB02         187, 254           Sains, Shunsuke         QB04, RB02         187, 254           Sains, Naurang L.         RB28         258           Saisho, Seiya         RN03         297           Saito, Hitoshi         RP06         304           Saito, Keasmi         KC03         108           Saito, Keasmi         KC03         108           Saito, Mineo         RC09, RC19         260, 262           Saito, T.         QB28         191           Saito, Takashi         PE13, QC05         193           Saito, Taku         QB02, QB15         187, 189           Saito, Taku         QB20, QE15         134, 189           Saito, Taku         QB25, SE23         266, 337           Sakajuchi, Jyunya         PD13         264           Sakai, Hito         CD02, SE01, SF07         51, 333, 338           Sakai, Hito	Sahadevan, Ajeesh	GE03	78
Sahoo, Roshnee         QHOE         213           Sahu, Subrata Kumar         BH04         46           Saidaoui, H. B. Mohamed         PJ13         156           Saiga, Yuta         RE04         267           Saiki, Shunsuke         QB04, RB02         187, 254           Saiki, Shunsuke         QB04, RB02         187, 254           Saini, Naurang L.         RE28         258           Saisho, Seiya         RN03         297           Saito, Hidekazu         Al05         39           Saito, Hitoshi         RP06         304           Saito, Kesami         KC03         108           Saito, Kesami         KC03         108           Saito, Kotaro         RD14         264           Saito, Kotaro         RD14         264           Saito, Kitsuru         QA17         185           Saito, T         QB28         191           Saito, Takashi         PE13, QC05         138, 193           Saito, Takashi         QE20, QF21         202, 207           Saito, Takashi         QE20, QF21         202, 207           Saito, Takashi         QE20, QF21         202, 207           Saito, Takashi         QE20, SE01, SF07         51, 333	Sahadevan, Ajeesh M	SK05	351
Sahu, Subrata Kumar         BH04         46           Saidaoui, H. B. Mohamed         P113         156           Saiga, Yuta         RE04         267           Saikai, Nobutaka         P008         121           Saiki, Shunsuke         QB04, RB02         187, 254           Sains, Saung L.         RE28         258           Saisho, Seiya         RN03         297           Saito, Hidekazu         Al05         39           Saito, Kesami         KC03         108           Saito, Kesami         KC03         108           Saito, Kesami         KC03         108           Saito, Kesami         QA17         185           Saito, Mitsuru         QA17         185           Saito, Takashi         PE13, QC05         138, 193           Saito, Takashi         PE13, QC05         138, 193           Saito, Taku         QB02, QB15         187, 189           Saito, Taku         QB02, QF1         202, 207           Sakaguchi, Jyunya         PD13         134           Sakaguchi, Yui         RD13         264           Sakai, Hideaki         PA12         117           Sakai, Hideaki         PA12         117	Sahashi, Masashi	CE05, PJ11, PL10, SM03	52, 156, 165, 361
Saidaoui, H. B. Mohamed         PJ13         156           Saiga, Yuta         RE04         267           Saiki, Nobutaka         PB08         121           Saiki, Shunsuke         QB04, RB02         187, 254           Saines, Paul         PC02         127           Saini, Naurang L.         RB28         258           Saisho, Seiya         RN03         297           Saito, Hidekazu         Al05         39           Saito, Hitoshi         RP06         304           Saito, Kesami         KC03         108           Saito, Mitoshi         RP06         260, 262           Saito, Mitsuru         QA17         185           Saito, T         QB28         191           Saito, Takashi         PE13, QC05         138, 193           Saito, Takashi         PD11         133           Saito, Taku         QB02, QB15         187, 189           Saito, Taku         QB02, QE15         184, 189           Saito, Taku         QB02, SE23         266, 337           Sakaguchi, Yui         RD13         264           Sakai, Akito         CD02, SE01, SF07         51, 333, 338           Sakai, Junya         QL11         229	Sahoo, Roshnee	QH06	213
Saiga, Yuta         RE04         267           Saiki, Nobutaka         PB08         121           Saiki, Shunsuke         QB04, RB02         187, 254           Saines, Paul         PC02         127           Saini, Naurang L.         RB28         258           Saisho, Seiya         RN03         297           Saito, Hidekazu         Al05         39           Saito, Hidekazu         Al05         304           Saito, Kesami         KC03         108           Saito, Kesami         RC09, RC19         260, 262           Saito, Mineo         RC09, RC19         260, 262           Saito, Mitsuru         QA17         185           Saito, T.         QB28         191           Saito, Takashi         PE13, QC05         138, 193           Saito, Takashi         PD11         133           Saito, Taku         QB02, QB15         187, 189           Saito, Taku         QB02, QB15         187, 189           Saito, Taku         QB02, QB15         137           Sakaguchi, Jyunya         PD13         134           Sakaguchi, Yui         RD13         264           Sakai, Hideaki         PA12         117	Sahu, Subrata Kumar	BH04	46
Saikai, Nobutaka         PB08         121           Saiki, Shunsuke         QB04, RB02         187, 254           Saini, Naurang L.         RB28         258           Saisho, Seiya         RN03         297           Saito, Hidekazu         Al05         39           Saito, Hitoshi         RP06         304           Saito, Kesami         KC03         108           Saito, Kotaro         RD14         264           Saito, Mineo         RC09, RC19         260, 262           Saito, Mineo         RC09, RC19         260, 262           Saito, Mineo         RC09, RC19         261, 262           Saito, T         Q828         191           Saito, Takahito         QC05         138, 193           Saito, Takashi         PE13, QC05         138, 193           Saito, Taku         Q802, QB15         187, 189           Saito, Taku         QE20, QF21         202, 207           Saito, Taku         QB2, SE23         266, 337           Sakaguchi, Yuji         RD13         264           Sakai, Akito         CD02, SE01, SF07         51, 333, 338           Sakai, Eijiro         QC14, QC15         194           Sakai, Junya         QL11	Saidaoui, H. B. Mohamed	PJ13	156
Saiki, Shunsuke         QB04, RB02         187, 254           Saines, Paul         PC02         127           Saini, Naurang L.         RB28         258           Saisho, Seiya         RN03         297           Saito, Hitdekazu         Al05         39           Saito, Hitdshi         RP06         304           Saito, Kotaro         RD14         264           Saito, Kotaro         RD14         264           Saito, Mitsuru         QA17         185           Saito, T         QB28         191           Saito, Takahito         QC05         193           Saito, Takashi         PE13, QC05         138, 193           Saito, Taku         QB02, QB15         187, 189           Saito, Taku         QB02, QB15         187, 189           Saito, Taku         QB02, QB15         134           Sakaguchi, Yuji         RD25, SE23         266, 337           Sakaguchi, Yuji         RD13         264           Sakai, Akito         CD02, SE01, SF07         51, 333, 338           Sakai, Akito         CD02, SE01, SF07         51, 333, 263           Sakai, Junya         QL11         229           Sakai, Junya         QL11         229 <td>Saiga, Yuta</td> <td>RE04</td> <td>267</td>	Saiga, Yuta	RE04	267
Saines, Paul         PCO         127           Saini, Naurang L.         RB28         258           Saisho, Seiya         RN03         297           Saito, Hidekazu         Al05         39           Saito, Hitoshi         RP06         304           Saito, Kesami         KC03         108           Saito, Kotaro         RD14         264           Saito, Mitsuru         QA17         185           Saito, T         QB28         191           Saito, Takashi         PE13, QC05         138, 193           Saito, Takashi         PE13, QC05         138, 193           Saito, Takashi         PE13, QC05         138, 193           Saito, Takashi         QE20, QF21         202, 207           Saito, Taku         QB02, QB15         187, 189           Saito, Taku         QB02, QB15         187, 189           Saito, Taku         QE20, QF21         202, 207           Saito, Taku         QB02, QB15         187, 189           Saito, Taku         QE02, QF21         202, 207           Saito, Taku         QE02, SE01, SF07         51, 333, 338           Sakai, Akito         CD02, SE01, SF07         51, 333, 263           Sakai, Hideaki <td< td=""><td>Saikai, Nobutaka</td><td>PB08</td><td>121</td></td<>	Saikai, Nobutaka	PB08	121
Saini, Naurang L.         RB28         258           Saisho, Seiya         RN03         297           Saito, Hidekazu         Al05         39           Saito, Hitoshi         RP06         304           Saito, Kesami         KC03         108           Saito, Kotaro         RD14         264           Saito, Kotaro         RC09, RC19         260, 262           Saito, Mitsuru         QA17         185           Saito, T         QB28         191           Saito, Takashi         PE13, QC05         138, 193           Saito, Takashi         PD11         133           Saito, Taku         QB02, QB15         187, 189           Saito, Toshiaki         QE20, QF21         202, 207           Saito, Toshiaki         QE20, QF21         202, 207           Sakou, Kito         CD02, SE01, SF07         51, 333, 338           Sakajuchi, Yui         RD13         264           Sakajuchi, Yui         RD13         264           Sakai, Hideaki         PA12         117           Sakai, Hideaki         PA12         117           Sakai, Junya         QL11         229           Sakai, Junya         QL11         229      S	Saiki, Shunsuke	QB04, RB02	187, 254
Saiko, Seiya         RN03         297           Saiko, Hidekazu         Al05         39           Saiko, Hitoshi         RP06         304           Saiko, Kesami         KC03         108           Saiko, Kesami         KC03         108           Saiko, Kotaro         RD14         264           Saiko, Kotaro         RD14         264           Saiko, Mitsuru         QA17         185           Saito, T         QB28         191           Saito, Takashi         PE13, QC05         138, 193           Saito, Takashi         PD11         133           Saito, Taku         QB02, QB15         187, 189           Saito, Toshiaki         QE20, QF21         202, 207           Saito, Toshiaki         QE20, QF21         202, 207           Saito, Toshiaki         QE20, QF21         202, 207           Sakaguchi, Yuji         RD13         264           Sakaguchi, Yui         RD13         264           Sakai, Hideaki         PA12         117           Sakai, Hideaki         PA12         117           Sakai, Junya         QL11         229           Sakai, Junya         QL11         229      Sakai, Seiji	Saines, Paul	PC02	127
Saito, Hidekazu         Al05         39           Saito, Hitoshi         RP06         304           Saito, Kesami         KC03         108           Saito, Kotaro         RD14         264           Saito, Mineo         RC09, RC19         260, 262           Saito, Mitsuru         QA17         185           Saito, T.         QB28         191           Saito, Takashi         PE13, QC05         138, 193           Saito, Takushi         PD11         133           Saito, Taku         QB02, QB15         187, 189           Saito, Toshiaki         QE20, QF21         202, 207           Saito, Toshiaki         QE20, QF21         202, 207           Sakor, Toshiaki         QE25, SE23         266, 337           Sakaguchi, Jyunya         PD13         134           Sakaguchi, Yui         RD13         264           Sakai, Hideaki         PA12         117           Sakai, Hinonori         IG04, PD09, RD11         95, 133, 263           Sakai, Junya         QL11         229           Sakai, Junya         QL11         201           Sakai, Toru         QE07, R111, R116         200, 280, 281           Sakai, Toru         QE07, R111, R116<	Saini, Naurang L.	RB28	258
Saito, Hitoshi         RP06         304           Saito, Kesami         KC03         108           Saito, Kotaro         RD14         264           Saito, Mineo         RC09, RC19         260, 262           Saito, Mitsuru         QA17         185           Saito, T.         QB28         191           Saito, Takahito         QC05         138, 193           Saito, Takashi         PE13, QC05         138, 193           Saito, Taku         QB02, QB15         187, 189           Saito, Toshiaki         QE20, QF21         202, 207           Saito, Toshiaki         QE20, QF21         202, 207           Sakaguchi, Jyunya         PD13         134           Sakaguchi, Jyunya         PD13         134           Sakaguchi, Yui         RD13         264           Sakai, Eijiro         QC14, QC15         194           Sakai, Hideaki         PA12         117           Sakai, Junya         QL11         229           Sakai, Joshi         QL03         228           Sakai, Joanu         PH02         147           Sakai, Yoshiki         RO06         301           Sakai, Yoshiki         RO06         301	Saisho, Seiya	RN03	297
Saito, Kesami         KC03         108           Saito, Kotaro         RD14         264           Saito, Mineo         RC09, RC19         260, 262           Saito, Mitsuru         QA17         185           Saito, T.         QB28         191           Saito, Takahito         QC05         193           Saito, Takashi         PE13, QC05         138, 193           Saito, Taku         QB02, QB15         187, 189           Saito, Toshiaki         QE20, QF21         202, 207           Saito, Taku         QB02, QB15         137, 189           Sakay, Toshiaki         QE20, QF21         202, 207           Saitoh, Fuji         Al02, FA01, KF03, RL15         39, 72, 111, 292           Sakay, Mito         CD02, SE01, SF07         51, 333, 338           Sakaguchi, Yui         RD13         264           Sakai, Hideaki         PA12         117           Sakai, Hideaki         PA12         117           Sakai, Junya         QL11         229           Sakai, N	Saito, Hidekazu	AI05	39
Saito, Kotaro         RD14         264           Saito, Mineo         RC09, RC19         260, 262           Saito, Mitsuru         QA17         185           Saito, T.         QB28         191           Saito, Takahito         QC05         193           Saito, Takashi         PE13, QC05         138, 193           Saito, Takashi         PD11         133           Saito, Taku         QB02, QB15         187, 189           Saito, Taku         QE20, QF21         202, 207           Saito, Toshiaki         QE20, QF21         202, 207           Saito, Toshiaki         QE20, QF21         202, 207           Saito, Toshiaki         QE20, QF21         202, 207           Sakot, Tyuji         RD25, SE23         266, 337           Sakaguchi, Yuji         RD13         264           Saka, Akito         CD02, SE01, SF07         51, 333, 338           Sakai, Akito         CD02, SE01, SF07         51, 333, 263           Sakai, Hironori         IG04, PD09, RD11         95, 133, 263           Sakai, Junya         QL11         229           Sakai, Josamu         PH02         147           Sakai, Seiji         BH05, BH06, QL15         46, 47, 230	Saito, Hitoshi	RP06	304
Saito, Mineo         RC09, RC19         260, 262           Saito, Mitsuru         QA17         185           Saito, T.         QB28         191           Saito, Takahito         QC05         193           Saito, Takashi         PE13, QC05         138, 193           Saito, Takashi         PD11         133           Saito, Taku         QB02, QB15         187, 189           Saito, Toshiaki         QE20, QF21         202, 207           Saito, Toshiaki         QE20, QF21         202, 207           Saito, Toshiaki         QE20, QF21         202, 207           Saito, Tuji         RD25, SE23         266, 337           Sakaguchi, Junya         PD13         134           Sakaguchi, Yui         RD13         264           Sakai, Fijiro         QC14, QC15         194           Sakai, Fijiro         QC14, QC15         194           Sakai, Hideaki         PA12         117           Sakai, Junya         QL11         229           Sakai, Naoki         QL03         228           Sakai, Samu         PH02         147           Sakai, Sakii, Samu         PH02         147           Sakai, Toru         QE07, R111, R116         200	Saito, Kesami	KC03	108
Saito, Mitsuru         QA17         185           Saito, T.         QB28         191           Saito, Takashi         PE13, QC05         193           Saito, Takashi         PE13, QC05         138, 193           Saito, Takashi         PD11         133           Saito, Takeshi         PD11         133           Saito, Taku         Q802, QB15         187, 189           Saito, Toshiaki         QE20, QF21         202, 207           Saitoh, Eiji         Al02, FA01, KF03, RL15         39, 72, 111, 292           Saitoh, Yuji         RD25, SE23         266, 337           Sakaguchi, Jyunya         PD13         134           Sakaguchi, Yui         RD13         264           Sakai, Akito         CD02, SE01, SF07         51, 333, 338           Sakai, Eijiro         QC14, QC15         194           Sakai, Hideaki         PA12         117           Sakai, Junya         QL11         229           Sakai, Junya         QL11         229           Sakai, Osamu         PH02         147           Sakai, Seiji         BH05, BH06, QL15         46, 47, 230           Sakai, Toru         QE07, R111, R116         200, 280, 281           Sakai, Toru <td>Saito, Kotaro</td> <td>RD14</td> <td>264</td>	Saito, Kotaro	RD14	264
Saito, T.         QB28         191           Saito, Takashi         QC05         193           Saito, Takashi         PE13, QC05         138, 193           Saito, Takashi         PD11         133           Saito, Taku         QB02, QB15         187, 189           Saito, Toshiaki         QE20, QF21         202, 207           Saito, Tyji         RD25, SE23         266, 337           Sakaguchi, Jyunya         PD13         134           Sakaguchi, Yui         RD13         264           Sakai, Akito         CD02, SE01, SF07         51, 333, 338           Sakai, Akito         CD02, SE01, SF07         51, 333, 263           Sakai, Hironori         IG04, PD09, RD11         95, 133, 263           Sakai, Junya         QL11         229           Sakai, Naoki         QL03         228           Sakai, Seiji         BH05, BH06, QL15         46, 47, 230           Sakai, Seiji         BH05, BH06, QL15         46, 47, 230           Sakai, Yoshiki         R006         301	Saito, Mineo	RC09, RC19	260, 262
Saito, Takahito         QC05         193           Saito, Takashi         PE13, QC05         138, 193           Saito, Takashi         PD11         133           Saito, Takeshi         PD11         133           Saito, Taku         Q802, QB15         187, 189           Saito, Toshiaki         QE20, QF21         202, 207           Saitoh, Eiji         Al02, FA01, KF03, RL15         39, 72, 111, 292           Saitoh, Yuji         RD25, SE23         266, 337           Sakaguchi, Jyunya         PD13         134           Sakaguchi, Yui         RD13         264           Sakai, Akito         CD02, SE01, SF07         51, 333, 338           Sakai, Hildeaki         PA12         117           Sakai, Hionori         IG04, PD09, RD11         95, 133, 263           Sakai, Junya         QL11         229           Sakai, Naoki         QL03         228           Sakai, Sakii, Sakii, Naoki         QL03         228           Sakai, Osamu         PH02         147           Sakai, Seiji         BH05, BH06, QL15         46, 47, 230           Sakai, Yoshiki         RO06         301           Sakai, Yoshiki         RO06         301      Sakaida, Masar	Saito, Mitsuru	QA17	185
Saito, Takashi         PE13, QC05         138, 193           Saito, Takeshi         PD11         133           Saito, Taku         QB02, QB15         187, 189           Saito, Toshiaki         QE20, QF21         202, 207           Saito, Toshiaki         QE20, QF21         202, 207           Saito, Toshiaki         QE20, QF21         202, 207           Saitoh, Eiji         Al02, FA01, KF03, RL15         39, 72, 111, 292           Saitoh, Yuji         RD25, SE23         266, 337           Sakaguchi, Jyunya         PD13         134           Sakaguchi, Yui         RD13         264           Sakai, Akito         CD02, SE01, SF07         51, 333, 338           Sakai, Akito         QC14, QC15         194           Sakai, Hironori         IG04, PD09, RD11         95, 133, 263           Sakai, Junya         QL11         229           Sakai, Naoki         QL03         228           Sakai, Osamu         PH02         147           Sakai, Seiji         BH05, BH06, QL15         46, 47, 230           Sakai, Toru         QE07, R111, R116         200, 280, 281           Sakai, Yoshiki         R006         301           Sakaida, Masaru         PF16         141	Saito, T.	QB28	191
Saito, Takeshi         PD11         133           Saito, Taku         Q802, QB15         187, 189           Saito, Toshiaki         QE20, QF21         202, 207           Saitoh, Eiji         Al02, FA01, KF03, RL15         39, 72, 111, 292           Saitoh, Fiji         RD25, SE23         266, 337           Sakaguchi, Jyunya         PD13         134           Sakaguchi, Yui         RD13         264           Sakai, Akito         CD02, SE01, SF07         51, 333, 338           Sakai, Akito         CD02, SE01, SF07         51, 333, 263           Sakai, Hideaki         PA12         117           Sakai, Hironori         IG04, PD09, RD11         95, 133, 263           Sakai, Junya         QL11         229           Sakai, Naoki         QL03         228           Sakai, Osamu         PH02         147           Sakai, Seiji         BH05, BH06, QL15         46, 47, 230           Sakai, Toru         QE07, R111, R116         200, 280, 281           Sakai, Toru         QE07, R111, R116         200, 280, 281           Sakaibara, To         CD01, EF01         50, 68           Sakakibara, T.         CD01, EF01         50, 201, 213, 222           Q108         Sakamoto, Naonor	Saito, Takahito	QC05	193
Saito, Taku         QB02, QB15         187, 189           Saito, Toshiaki         QE20, QF21         202, 207           Saitoh, Eiji         Al02, FA01, KF03, RL15         39, 72, 111, 292           Saitoh, Yuji         RD25, SE23         266, 337           Sakaguchi, Jyunya         PD13         134           Sakaguchi, Yui         RD13         264           Sakai, Akito         CD02, SE01, SF07         51, 333, 338           Sakai, Akito         QC14, QC15         194           Sakai, Hideaki         PA12         117           Sakai, Hironori         IG04, PD09, RD11         95, 133, 263           Sakai, Junya         QL11         229           Sakai, Naoki         QL03         228           Sakai, Osamu         PH02         147           Sakai, Seiji         BH05, BH06, QL15         46, 47, 230           Sakai, Toru         QE11         201           Sakai, Toru         QE07, RI11, RI16         200, 280, 281           Sakai, Masaru         PF16         141           Sakaibara, To         CD01, EF01         50, 68           Sakakibara, Toshiro         CC05, QE13, QH03, Q108         50, 201, 213, 222           Sakamoto, Naonori         KC02         10	Saito, Takashi	PE13, QC05	138, 193
Saito, Toshiaki         QE20, QF21         202, 207           Saitoh, Eiji         Al02, FA01, KF03, RL15         39, 72, 111, 292           Saitoh, Yuji         RD25, SE23         266, 337           Sakaguchi, Jyunya         PD13         134           Sakaguchi, Yui         RD13         264           Sakai, Akito         CD02, SE01, SF07         51, 333, 338           Sakai, Akito         QC14, QC15         194           Sakai, Hideaki         PA12         117           Sakai, Hironori         IG04, PD09, RD11         95, 133, 263           Sakai, Junya         QL11         229           Sakai, Naoki         QL03         228           Sakai, Osamu         PH02         147           Sakai, Seiji         BH05, BH06, QL15         46, 47, 230           Sakai, Toru         QE11         201           Sakai, Toru         QE07, RI11, RI16         200, 280, 281           Sakai, Yoshiki         RO06         301           Sakai, Yoshiki         RO06         301           Sakaida, Masaru         PF16         141           Sakaibara, T.         CD01, EF01         50, 68           Sakamoto, Naonori         KC02         108      Sakamoto, Naon	Saito, Takeshi	PD11	133
Saitoh, Eiji         Al02, FA01, KF03, RL15         39, 72, 111, 292           Saitoh, Yuji         RD25, SE23         266, 337           Sakaguchi, Jyunya         PD13         134           Sakaguchi, Yui         RD13         264           Sakai, Akito         CD02, SE01, SF07         51, 333, 338           Sakai, Akito         QC14, QC15         194           Sakai, Hideaki         PA12         117           Sakai, Hironori         IG04, PD09, RD11         95, 133, 263           Sakai, Junya         QL11         229           Sakai, Naoki         QL03         228           Sakai, Seiji         BH05, BH06, QL15         46, 47, 230           Sakai, Toru         QE11         201           Sakai, Toru         QE11         201           Sakai, Toru         QE07, RI11, RI16         200, 280, 281           Sakai, Yoshiki         RO06         301      Sakaida, Masaru         PF16         141      Sakaibara, T.         CD01, EF01         50, 68      Sakamoto, Hirotoshi         HC03         85      Sakamoto, Naonori         KC02         108      Sakamoto, Yasutaka         SG09         341      Sakanoto, Yasutaka         SG09         341 <td>Saito, Taku</td> <td>QB02, QB15</td> <td>187, 189</td>	Saito, Taku	QB02, QB15	187, 189
Saitoh, Yuji         RD25, SE23         266, 337           Sakaguchi, Jyunya         PD13         134           Sakaguchi, Yui         RD13         264           Sakai, Akito         CD02, SE01, SF07         51, 333, 338           Sakai, Akito         QC14, QC15         194           Sakai, Hideaki         PA12         117           Sakai, Hironori         IG04, PD09, RD11         95, 133, 263           Sakai, Junya         QL11         229           Sakai, Naoki         QL03         228           Sakai, Seiji         BH05, BH06, QL15         46, 47, 230           Sakai, Toru         QE11         201           Sakai, Toru         QE07, RI11, RI16         200, 280, 281           Sakai, Yoshiki         RO06         301           Sakaida, Masaru         PF16         141           Sakaibara, T.         CD01, EF01         50, 68           Sakamoto, Hirotoshi         HC03         85           Sakamoto, Yasutaka         SG09         341           Sakano, Rui         PE16         138           Sakano, Rui         PE16         138           Sakano, Rui         PE16         138           Sakano, Rui         PE16         <	Saito, Toshiaki	QE20, QF21	202, 207
Sakaguchi, Jyunya         PD13         134           Sakaguchi, Yui         RD13         264           Sakai, Akito         CD02, SE01, SF07         51, 333, 338           Sakai, Akito         QC14, QC15         194           Sakai, Fijiro         QC14, QC15         194           Sakai, Hideaki         PA12         117           Sakai, Hinonori         IG04, PD09, RD11         95, 133, 263           Sakai, Junya         QL11         229           Sakai, Naoki         QL03         228           Sakai, Seiji         BH05, BH06, QL15         46, 47, 230           Sakai, Seiji         BH05, BH06, QL15         46, 47, 230           Sakai, T^oru         QE11         201           Sakai, Toru         QE07, R111, R116         200, 280, 281           Sakaida, Masaru         PF16         141           Sakaida, Masaru         PF16         141           Sakaibara, T.         CD01, EF01         50, 68           Sakamoto, Hirotoshi         HC03         85           Sakamoto, Naonori         KC02         108           Sakamoto, Yasutaka         SG09         341           Sakano, Rui         PE16         138           Sakata, Masafumi	Saitoh, Eiji	AI02, FA01, KF03, RL15	39, 72, 111, 292
Sakaguchi, Yui         RD13         264           Sakai, Akito         CD02, SE01, SF07         51, 333, 338           Sakai, Akito         QC14, QC15         194           Sakai, Eijiro         QC14, QC15         194           Sakai, Hideaki         PA12         117           Sakai, Hideaki         PA12         117           Sakai, Hironori         IG04, PD09, RD11         95, 133, 263           Sakai, Junya         QL11         229           Sakai, Naoki         QL03         228           Sakai, Seiji         BH05, BH06, QL15         46, 47, 230           Sakai, Seiji         BH05, BH06, QL15         46, 47, 230           Sakai, T^oru         QE07, RI11, RI16         200, 280, 281           Sakai, Toru         QE07, RI11, RI16         200, 280, 281           Sakaida, Masaru         PF16         141           Sakaida, Masaru         PF16         141           Sakaibara, T.         CD01, EF01         50, 68           Sakamoto, Hirotoshi         HC03         85           Sakamoto, Naonori         KC02         108           Sakanoto, Naonori         KC02         108           Sakanoto, Rui         PE16         138           Sak	Saitoh, Yuji	RD25, SE23	266, 337
Sakai, Akito         CD02, SE01, SF07         51, 333, 338           Sakai, Eijiro         QC14, QC15         194           Sakai, Hideaki         PA12         117           Sakai, Hideaki         PA12         117           Sakai, Hironori         IG04, PD09, RD11         95, 133, 263           Sakai, Junya         QL11         229           Sakai, Naoki         QL03         228           Sakai, Osamu         PH02         147           Sakai, Seiji         BH05, BH06, QL15         46, 47, 230           Sakai, Toru         QE11         201           Sakai, Toru         QE07, RI11, RI16         200, 280, 281           Sakai, Toru         QE07, RI11, RI16         200, 280, 281           Sakaida, Masaru         PF16         141           Sakaibara, T.         CD01, EF01         50, 68           Sakaibara, Toshiro         CC05, QE13, QH03, QU08         50, 201, 213, 222           Sakamoto, Hirotoshi         HC03         85           Sakamoto, Naonori         KC02         108           Sakamoto, Yasutaka         SG09         341           Sakano, Rui         PE16         138           Sakata, Masafumi         PC21         130	Sakaguchi, Jyunya	PD13	134
Sakai, Eijiro         QC14, QC15         194           Sakai, Hideaki         PA12         117           Sakai, Hironori         IG04, PD09, RD11         95, 133, 263           Sakai, Hironori         IG04, PD09, RD11         229           Sakai, Junya         QL11         229           Sakai, Naoki         QL03         228           Sakai, Seiji         BH05, BH06, QL15         46, 47, 230           Sakai, Coamu         PH02         147           Sakai, T^oru         QE11         201           Sakai, Toru         QE07, RI11, RI16         200, 280, 281           Sakai, Yoshiki         RO06         301           Sakaida, Masaru         PF16         141           Sakaida, Masaru         PF16         141           Sakaibara, T.         CD01, EF01         50, 68           Sakamoto, Hirotoshi         HC03         85           Sakamoto, Naonori         KC02         108           Sakamoto, Naonori         KC02         108           Sakano, Rui         PE16         138           Sakata, Masafumi         PC21         130           Sakata, T.         PL04         164           Sakoda, Masahito         PE13, RD01         <	Sakaguchi, Yui	RD13	264
Sakai, HideakiPA12117Sakai, HironoriIG04, PD09, RD1195, 133, 263Sakai, JunyaQL11229Sakai, NaokiQL03228Sakai, OsamuPH02147Sakai, SeijiBH05, BH06, QL1546, 47, 230Sakai, T^oruQE11201Sakai, ToruQE07, RI11, RI16200, 280, 281Sakai, YoshikiRO06301Sakaida, MasaruPF16141Sakaibara, T.CD01, EF0150, 68Sakaibara, ToshiroCC05, QE13, QH03, QJ0850, 201, 213, 222Sakamoto, HirotoshiHC0385Sakamoto, NaonoriKC02108Sakano, RuiPE16138Sakata, MasafumiPC21130Sakata, T.PL04164Sakoda, MasahitoPE13, RD01138, 262	Sakai, Akito	CD02, SE01, SF07	51, 333, 338
Sakai, Hironori         IG04, PD09, RD11         95, 133, 263           Sakai, Junya         QL11         229           Sakai, Naoki         QL03         228           Sakai, Osamu         PH02         147           Sakai, Seiji         BH05, BH06, QL15         46, 47, 230           Sakai, T^oru         QE11         201           Sakai, Toru         QE07, R111, R116         200, 280, 281           Sakai, Yoshiki         R006         301           Sakaida, Masaru         PF16         141           Sakakibara, T.         CD01, EF01         50, 68           Sakamoto, Hirotoshi         HC03         85           Sakamoto, Naonori         KC02         108           Sakano, Rui         PE16         138           Sakano, Rui         PC21         130           Sakata, Masafumi         PC21         130           Sakata, T.         PL04         164           Sakoda, Masahito         PE13, RD01         138, 262	Sakai, Eijiro	QC14, QC15	194
Sakai, JunyaQL 11229Sakai, NaokiQL03228Sakai, OsamuPH02147Sakai, OsamuPH02147Sakai, SeijiBH05, BH06, QL1546, 47, 230Sakai, T^oruQE11201Sakai, ToruQE07, RI11, RI16200, 280, 281Sakai, ToruQE07, RI11, RI16200, 280, 281Sakai, YoshikiRO06301Sakaida, MasaruPF16141Sakaibara, T.CD01, EF0150, 68Sakakibara, ToshiroCC05, QE13, QH03, QJ0850, 201, 213, 222Sakamoto, HirotoshiHC0385Sakamoto, NaonoriKC02108Sakamoto, YasutakaSG09341Sakano, RuiPE16138Sakata, MasafumiPC21130Sakata, T.PL04164Sakoda, MasahitoPE13, RD01138, 262	Sakai, Hideaki	PA12	117
Sakai, Naoki         QL03         228           Sakai, Osamu         PH02         147           Sakai, Seiji         BH05, BH06, QL15         46, 47, 230           Sakai, T^oru         QE11         201           Sakai, Toru         QE07, R111, R116         200, 280, 281           Sakai, Yoshiki         RO06         301           Sakaida, Masaru         PF16         141           Sakakibara, T.         CD01, EF01         50, 68           Sakaibara, Toshiro         CC05, QE13, QH03, QJ08         50, 201, 213, 222           Sakamoto, Hirotoshi         HC03         85           Sakamoto, Yasutaka         SG09         341           Sakano, Rui         PE16         138           Sakata, Masafumi         PC21         130           Sakata, T.         PL04         164           Sakoda, Masahito         PE13, RD01         138, 262	Sakai, Hironori	IG04, PD09, RD11	95, 133, 263
Sakai, Osamu         PH02         147           Sakai, Seiji         BH05, BH06, QL15         46, 47, 230           Sakai, T^oru         QE11         201           Sakai, Toru         QE07, RI11, RI16         200, 280, 281           Sakai, Yoshiki         RO06         301           Sakaida, Masaru         PF16         141           Sakaibara, T.         CD01, EF01         50, 68           Sakaribara, Toshiro         CC05, QE13, QH03, QJ08         50, 201, 213, 222           Sakamoto, Hirotoshi         HC03         85           Sakamoto, Yasutaka         SG09         341           Sakano, Rui         PE16         138           Sakata, Masafumi         PC21         130           Sakata, T.         PL04         164           Sakoda, Masahito         PE13, RD01         138, 262	Sakai, Junya	QL11	229
Sakai, Seiji         BH05, BH06, QL15         46, 47, 230           Sakai, T^oru         QE11         201           Sakai, Toru         QE07, R111, R116         200, 280, 281           Sakai, Yoshiki         RO06         301           Sakaida, Masaru         PF16         141           Sakaibara, T.         CD01, EF01         50, 68           Sakaibara, Toshiro         CC05, QE13, QH03, QJ08         50, 201, 213, 222           Sakamoto, Hirotoshi         HC03         85           Sakamoto, Yasutaka         SG09         341           Sakaro, Rui         PE16         138           Sakata, Masafumi         PC21         130           Sakata, T.         PL04         164           Sakoda, Masahito         PE13, RD01         138, 262	Sakai, Naoki	QL03	228
Sakai, T^oru         QE11         201           Sakai, Toru         QE07, RI11, RI16         200, 280, 281           Sakai, Yoshiki         RO06         301           Sakaida, Masaru         PF16         141           Sakaibara, T.         CD01, EF01         50, 68           Sakaibara, Toshiro         CC05, QE13, QH03, QJ08         50, 201, 213, 222           Sakamoto, Hirotoshi         HC03         85           Sakamoto, Naonori         KC02         108           Sakano, Rui         PE16         138           Sakata, Masafumi         PC21         130           Sakata, T.         PL04         164           Sakoda, Masahito         PE13, RD01         138, 262	Sakai, Osamu	PH02	147
Sakai, Toru         QE07, RI11, RI16         200, 280, 281           Sakai, Yoshiki         RO06         301           Sakaida, Masaru         PF16         141           Sakaibara, T.         CD01, EF01         50, 68           Sakaibara, Toshiro         CC05, QE13, QH03, QJ08         50, 201, 213, 222           Sakamoto, Hirotoshi         HC03         85           Sakamoto, Naonori         KC02         108           Sakano, Rui         PE16         138           Sakata, Masafumi         PC21         130           Sakata, T.         PL04         164           Sakoda, Masahito         PE13, RD01         138, 262	Sakai, Seiji	BH05, BH06, QL15	46, 47, 230
Sakai, YoshikiRO06301Sakaida, MasaruPF16141Sakakibara, T.CD01, EF0150, 68Sakakibara, ToshiroCC05, QE13, QH03, QJ0850, 201, 213, 222Sakamoto, HirotoshiHC0385Sakamoto, NaonoriKC02108Sakamoto, YasutakaSG09341Sakata, MasafumiPC21130Sakata, T.PL04164Sakoda, MasahitoPE13, RD01138, 262	Sakai, T^oru	QE11	201
Sakaida, MasaruPF16141Sakakibara, T.CD01, EF0150, 68Sakakibara, ToshiroCC05, QE13, QH03, QJ0850, 201, 213, 222 QJ08Sakamoto, HirotoshiHC0385Sakamoto, NaonoriKC02108Sakamoto, NaonoriKC02108Sakamoto, YasutakaSG09341Sakano, RuiPE16138Sakata, MasafumiPC21130Sakata, T.PL04164Sakoda, MasahitoPE13, RD01138, 262	Sakai, Toru	QE07, RI11, RI16	200, 280, 281
Sakakibara, T.CD01, EF0150, 68Sakakibara, ToshiroCC05, QE13, QH03, QJ0850, 201, 213, 222Sakamoto, HirotoshiHC0385Sakamoto, NaonoriKC02108Sakamoto, YasutakaSG09341Sakano, RuiPE16138Sakata, MasafumiPC21130Sakata, T.PL04164Sakoda, MasahitoPE13, RD01138, 262	Sakai, Yoshiki	RO06	301
Sakakibara, ToshiroCC05, QE13, QH03, QJ0850, 201, 213, 222Sakamoto, HirotoshiHC0385Sakamoto, NaonoriKC02108Sakamoto, YasutakaSG09341Sakano, RuiPE16138Sakata, MasafumiPC21130Sakata, T.PL04164Sakoda, MasahitoPE13, RD01138, 262	Sakaida, Masaru	PF16	141
QJ08           Sakamoto, Hirotoshi         HC03         85           Sakamoto, Naonori         KC02         108           Sakamoto, Yasutaka         SG09         341           Sakano, Rui         PE16         138           Sakata, Masafumi         PC21         130           Sakata, T.         PL04         164           Sakoda, Masahito         PE13, RD01         138, 262	Sakakibara, T.	CD01, EF01	50, 68
Sakamoto, HirotoshiHC0385Sakamoto, NaonoriKC02108Sakamoto, YasutakaSG09341Sakano, RuiPE16138Sakata, MasafumiPC21130Sakata, T.PL04164Sakoda, MasahitoPE13, RD01138, 262	Sakakibara, Toshiro		50, 201, 213, 222
Sakamoto, NaonoriKC02108Sakamoto, YasutakaSG09341Sakano, RuiPE16138Sakata, MasafumiPC21130Sakata, T.PL04164Sakoda, MasahitoPE13, RD01138, 262	Sakamoto, Hirotoshi		85
Sakamoto, YasutakaSG09341Sakano, RuiPE16138Sakata, MasafumiPC21130Sakata, T.PL04164Sakoda, MasahitoPE13, RD01138, 262			
Sakano, Rui         PE16         138           Sakata, Masafumi         PC21         130           Sakata, T.         PL04         164           Sakoda, Masahito         PE13, RD01         138, 262			
Sakata, MasafumiPC21130Sakata, T.PL04164Sakoda, MasahitoPE13, RD01138, 262			
Sakata, T.         PL04         164           Sakoda, Masahito         PE13, RD01         138, 262			
Sakoda, Masahito PE13, RD01 138, 262			
	- and in a subalant diam		

Sakuma, Akimasa

Sakuma, Keita Sakurada, T. Sakurada, Takashi Sakurai, Hiroshi Sakurai, Hiroya Sakurai, Hiroyuki Sakurai, Shohei Sakurai, T. Sakurai, Takahiro Sakurai, Yoshiharu Sakurai, Yoshiki Salamakha, Leonid Salamati, H. Salamom, M. B. Salazar Mejia, Catalina Salazar, Nicolas Antonio Salce, Bernard Salles, B. Salman, Zaher Salvador, M D Samant, Mahesh Samanta, Tapas Samardak, Alexander Samardak, A. Sergeevich Samarin, N. A. Samarin, Nickolay Samariya, Arvind Samata, Hiroaki Sampaio, Joao Sampaio, Maria Sampathkumaran, E V Sampathkumaran, E. V. Sanada, Naoyuki Sanai, Tatsunori Sanchez Llamazares, Jose Luis JJ04 Sanchez Marcos, Jorge Sanchez-costa, Jose Sanda, Minoru Sandalov, Vladimir Sandig, Oliver Sandiumenge, Felip Sando, Daniel Sangregorio, Claudio Santamaria, Jesus Santava, Eva Santos, Tiffany Sarachik, Myriam P. Sarkar, Sujit Sarma, D. D. Sarrao, John L. Saruya, Takeshi Sasagawa, Takao Sasaki, Hiroko Sasaki, Hiromi Sasaki, N.

Sasaki, Ryo

CH02, KH02, PJ10, PL17, 54, 112, 156, 166, RL09, SI06 291, 346 PB34, PJ12, SA07 125, 156, 317 SG11 341 JI04 105 PK12, PM02 162, 167 PH19, RH06 150, 275 QP15, QP16 247, 248 RH09 276 KI05, SL14 113, 358 137, 199, 201, 204, PE11, QE03, QE10, QF04, RD01, RP05 262, 304 BA04, PG05, PG13, 41, 143, 144, 162, PK12, SK04 351 320 SB03 SC01 325 PI23, PI24 154 QB12 189 Q009 242 SM04 361 PC17 129 CE06 52 BC02, PB20, PN23, QL19 42, 123, 175, 231 RQ10 308 RN16 299 CH01 54 AG05 38 QM05 232 QC22, RH07 195, 275 67, 113, 340 ED06, KI05, SG08 360 SL25 351 SK04 87, 94 HF01, IE04 271 RF13 138 PE14 199 OD16 SC23 329 KF05 111 106 JJ04 106 79 GF04 PG17, QF11, QF14 145, 205, 206 RR28 315 RP08 304 IH01, SO18 95, 369 QA27 187 CB04, SI11 49, 347 RN09 298 QI01 217 AD03, BG02 35, 45 IC01 92 QF26 208 RF12 270 RD09 263 KF01 110 Q018 243 SG18 342 SL02 356 53 CF04 SE11 335

GB04

Sasaki, Satoru

Sasaki, Satoru	GB04	/6
Sasaki, Satoshi	RC11	260
Sasaki, Takahiko	GB04, RH03, RH11	76,
Sasano, Masashi	PN07	173
Saso, Tetsuro	PF17	142
Sassa, Y.	PC29, PI15	131
Sassa, Yasmine	HA03, SE08	83,
Sasso, Carlo Paolo	DG04, HH01	61,
Sato, Hideo	QP14, SI17	247
Sato, Hideyuki	PE13, QC05, RD17	138
Sato, Hirohiko	QE17, QF09	202
Sato, Hitoshi	RD20	265
Sato, Kazuhisa	KB05	108
Sato, Kazunori	KF04	111
Sato, Kei	RL13	292
Sato, Kentaro	PB29, PB30, PB32	124
Sato, Kiyoo	Q104	218
Sato, Koichi	RN11	298
Sato, Kota	SB05	321
Sato, Kotaro	SN03	364
Sato, Masahiro	JD04, RI11	101
Sato, Masatoshi	DA01, QB04, QB26,	56,
	RB02, RB05	
Sato, Mitsuhide	QF08, QF10	205
Sato, Mototsugu	QJ04	221
Sato, Nana	SI08	347
Sato, Noriaki	PD20	135
Sato, Noriaki K	RD05	262
Sato, Noriaki K.	PD14	134
Sato, Suguru	SJ06	349
Sato, Taku J	QE13	201
Sato, Tetsuya	RK08, SM20	287
Sato, Tomonori	HG02, RJ01, RJ02, RJ09	88,
Sato, Yoshiaki	SG09	341
Sato, Yuji	PL10	165
Satoh, Isamu	RD05	262
Satoh, Ryohei	RD24	266
Satter, A. A.	RN01	296
Savero-torres, Williams	GE01, QK15, QK16, QK17, SL24	78,
Savici, Andrei T	PI07	151
Savici, Andrei T.	RD14	264
Sawa, A.	PL14	166
Sawa, H.	CF01	52
Sawada, Kazuya	SM03	361
Sawada, Keisuke	RC19	262
Sawada, M.	PH14	149
Sawada, Masahiro	PN20	175
Sawami, Kazuma	QJ12	222
Sawamura, M.	KJ01	113
Sawatzki, Simon	KH01	112
Sawicki, M.	BE02	44
Sawicki, Maciej	BE03, DD02	44,
Sayanagi, Richard	PK09, QM01	161
Sazonov, Andrew	PA14	118
Schaadt, Daniel M.	QL05	228
Schaefer, R.	PM26	171
Schaefer, Rudolf	HB01	84
Schaf, Jacob	PBO9, PB10, PB11, PB33, RBO9, SB18	121

	76	Schafer, Hanjo	HA02	83
14.4	260	Schafers, Markus	IF04	94
111	76, 275, 276	Schalenbach, Maximilian	QA22	186
	173	Schau-magnussen, Magnus	BC02	42
	142	Schedler, Roland	QG02	208
	131, 153	Scheffer, Jurg	PA04	116
	83, 334	Schelhas, Laura	QA15	185
	61, 88	Schellekens, Adrianus	EIOG	71
	247, 348	Schellekens, Sjors	SIO7	347
)17	138, 193, 264	Schepotin, I	BH01	46
	202, 205	Schierle, E.	PC22	130
	265	Schlage, Kai	QM19	234
	108	Schlenhoff, Anika	HF03, HF04, RP16	87, 305
	111	Schlom, Darrell G	SA08	317
22	292	Schlottmann, Pedro	KA04, PA21, PE04	107, 119, 136
32	124, 125	Schmalzl, Karin	AA05, QM02	33, 231
	218	Schmidiger, David Jan	CC03	50
	298	Schmidt, Burkhard	DB02, QE04	57, 200
	321	Schmidt, Michael	PA23, RH24	119, 278
	364	Schmidt, Oliver G	SJO7	350
<b>D</b> 2 <i>C</i>	101, 280	Schmidt, Rene	AF02, HG03, PG12	36, 88, 144
B26,	56, 187, 191, 254	Schmidt, Wolfgang	QM02	231
	205	Schmiedeshoff, G. M.	IA05	91
	205	Schmitt, T.	PC01	126
	221	Schmitt, Thorsten	BA03	41
	347	Schmitz, Defleft	IH01	95
	135	Schneidewind, Astrid	QC24	196
	262	Schnelle, Frank	QJ11	222
	134	Schnelle, Walter	RD12	264
	349	Schoenemann, Rico Schoenhense, G	PC27	131
	201	Schon, Gerd	QM12, RN12, RP20 1104	233, 298, 306 97
	287, 364	Schonhense, Gerd	CH04	54
02, RJ09	88, 282, 284	Schotter, Joerg	AH01	38
	341 165	Schrefl, Thomas	DE01	59
		Schreyer, Andreas	DG01, QM02	61, 231
	262 266	Schrittwieser, Stefan	AH01	38
	296	Schroeder, Henning	KD01	109
K16,	78, 226, 360	Schroeder, Silke	IH02	96
K10,	78, 220, 300	Schuller, Ivan K	BD01	43
	151	Schuller, Ivan K.	BD04, ED03	43, 66
	264	Schvets, I V	PH16	149
	166	Schwabe, Andrej	PE05, PE07	136, 137
	52	Schwarz, Alexander	AF02, HG03, PG12	36, 88, 144
	361	Schwarz, Karlheinz	RB21	257
	262	Schwarzacher, Walther	KB01	107
	149	Schweitzer, Dieter	SG19	342
	175	Schweitzer, Sebastian	QK14	226
	222	Schwingenschlogl, U.	PN02	172
	113	Schwingenschlogl, Udo	QL17, QO01, QP01,	230, 240, 245, 352
	112		QP02, SK12	
	44	Schwobel, Jorg	GH02	81
	44, 58	Sebastian, S. E.	GA01	75
	161, 231	Sebastian, Suchitra	BA01	40
	118	Sebt, S. A.	QP23	249
	228	Sechovsky, Vladimir	BB04, HJ04, PC10, PD18,	42, 90, 128, 134,
	171	2* · · · ·	PD19, QC21, QJ02, QJ14	135, 195, 221, 223
	84	Sedlak, Kamil	AB02	34
11, PB33,	121, 125, 255, 323	Segal, Amir	DE02	59
		Segawa, Kouji	RC11	260
		segura, nouji		

Seike, Masayoshi	KF04	111
Seiro, Silvia Seki, Harumi	QC09, QC18, QD07 PC15	193, 195, 197 129
Seki, Kazuhiro	PC15 PF09, PF14, RE08	129 140, 141, 267
Seki, S.	RA05	250
Seki, Takeshi	JG03, KC03	103, 108
Sekiguchi, Koji	RJ10, SI08	284, 347
Sekiguchi, Toru	RD22	265
Sekimoto, S.	KJ01	113
Sekine, Akihiko	GB02	76
Sekine, Chihiro	PD10, PD11, PD22, QD06	133, 135, 197
Sekine, Tomoyuki	PA19, QA05, QA07, QA09, RH11	119, 183, 184, 276
Sekiyama, Akira	RD25	266
Sela, Eran	RI12	280
Selezneva, Nadezshda	QG04	208
Sellmyer, David	CB05	49
Seman, Hazrina Abu	RP12	305
Semeno, A. V.	PE12	138
Semeno, Alexey	ED06, KI05, SG07, SG08	67, 113, 340
Semeno, Alexey V.	QC19	195
Semisalova, Anna	Q014	243
Semisalova, Anna S	PN22	175
Senba, Shinya	QM10, SK13	233, 352
Sendilkumar, A	QD04	197
Sengupta, Pinaki	DC04	58
Senna, Mamoru	QL07	229
Seo, Ambrose	PA16	118
Seo, Dong-jea	SL26	360
Seo, Hitoshi	GB03, RG09	76, 273
Seo, Jeongdae	RC05	259
Seo, Juyeon	PO06	178
Seo, M. S. Seo, Soo-man	RA07	250
Seo, Soonbeom	HF05, HP21, PJ33 PC19	87, 30, 160 129
Seong, W. K.	PB03	129
Sera, Masafumi	QI10, RD10	219, 263
Seredkin, Vitalv	SO28	371
Serga, Aleksandr	HD01	85
Serga, Alexander A.	SI15	348
Sergey, Faleev	JF04	103
Serpico, C.	RJ07	283
Serrano-ramon, Luis	DG05	61
Serrate, David	IH02, PG08	96, 143
Sese, Javier	GG03, GG05	80
Sessi, Violetta	AD03	35
Seto, Makoto	QM16	234
Settai, Rikio	HP11, PB27, PD08, PD13, PD16, RD19, RD25	30, 124, 133, 134, 265, 266
Seung Mok, Jeon	PJ15	157
Severing, A.	KA02	107
Severing, Andrea	QC09, QC12	193, 194
Seyfarth, G.	PD05	132
Seyfarth, Gabriel	CD04	51
Seyyed Ebrahimi, S. A.	QM28, SA04	236, 316
Shamba, Precious	QG09	209
Shang, T.	QB12	189
Shang, Tian	SB29	325
Shao, Xiaoping	RQ01	307

Shaofeng, Lin	Q021	244
Sharbaf Zadeh, Asghar	PC16	129
Sharma, A. K	QN14	238
Sharma, Bharat Kumar	KI01	112
Sharma, Neetika	PI12	152
Sharma, P. K.	SL25	360
Sharma, Sandeep	AI05	39
Shavrov, Vladimir G.	SO19	369
Shavrov, Vladimir Grigor'evich	QH17	215
Shaw, C.c.	SN06	365
Shaykhutdinov, Kirill	QA08	183
She, Jian-huang	SE18	336
Sheikhzadeh, G. A.	PO27, PO29, PO30, PO31	181, 182
Sheikin, I.	GC01	77
Sheikin, Ilya	HP11	30
Shekhter, Robert I.	RG17	274
Shen, Baogen	RQ01, SO25	307, 370
Shen, Bao-gen	ED05	67
Shen, Chenyi	IJ04	98
Shen, Jian	AG01	37
Shen, Lei	GI04, QP06, SJ09	82, 246, 350
Sheppard, Charles Johannes	RH15	277
Shermadini, Zurab	QB22	190
Sherokalova, Elizaveta	QG04	208
Sherstobitova, Elena	QG02	208
Sherstobitova, Elena A.	QN10	238
Shevchenko, A	BH01	46
Shi, Chunlong	RA13, RA14	251
Shi, Jianzhong	QP22	249
Shi, M.	PC01	126
Shi, Ming	PD12	133
Shi, Yi	PN27	176
Shi, Z.	PL20	167
Shi, Zhong	BD05, JI05, SL16	43, 105, 358
Shibasaki, Seizi	SCO4	326
Shibauchi, Takasada	FB02, QC04, SB27	73, 192, 324
Shibuya, K.	RG10	273
Shick, Alexander	RE03	266
Shield, Jeffrey E	KB02	107
Shigeoka, Toru	QI18, SC24	220, 329
Shigeta, Iduru	QG13, QG15, SK23	210, 354
Shih, C.w.	SN06	365
Shih, Ying-ta	ED02	66 222
Shikama, T.	QJ10	
Shikoh, Eiji	AI02, ED04, PN12, RL04, SK16	39, 66, 174, 290, 353
Shim, In-bo	PI14	153
Shim, J. H.	PE17	138
Shim, J.h.	QH25	217
Shim, Je-ho	SH10	344
Shim, Jeong Hyun	BF05	45
Shim, Ji Hoon	JA03, QH27, RB22, RB23	98, 217, 257
Shim, Ji-hoon	QH26, QJ18, QJ19	217, 223, 224
Shim, Seong Hoon	RL31	295
Shim, Sung-ah	BH02	46
Shima, Toshiyuki	RO09, SN03	302, 364
Shimada, Daisuke	QN20	239
Shimada, Kenya Shimada, Yutaka	RD20	265
Shimada, Yutaka	QM07, RK04	232, 286

Author Index

Shimakawa, Yuichi	Q
Shimamoto, Kohei	E
Shimamura, Kazutoshi	RI
Shimano, Ryo	P/
Shimba, Kazuaki	Q
Shimidzu, Naoki	Sł
Shimizu, Katsuya	PE
Shimizu, Kenji	S
Shimizu, Ryutarou	RI
Shimizu, Sunao	PE
Shimizu, Tomohiro	PO
Shimizu, Yasuhiro	Pł
Shimizu, Yusei	Q
Shimojima, T.	Q
Shimokawa, Takaya	Pľ
Shimokawa, Tokuro	RI
Shimomura, Koji	Q
Shimomura, Naoki	PL
Shimotani, Hidekazu	A
Shimura, Y.	С
Shimura, Yasuyuki	С
Shin, Hy.	P
Shin, Hosun	SI
Shin, Hyungsoon	RI
Shin, Hyun-jin	P
Shin, Il Jae	Q
Shin, Il-jae	RI
Shin, Jaewon	P
Shin, Jihye	RI
Shin, Jinsik	Q
Shin, Jong Moon	Pľ
Shin, Jun-yeong	BI
Shin, K.h.	PJ
Shin, Kwang-ho	P
	R
Shin, Kyung Ho	Q
Shin, Kyungho	Q
Shin, Kyung-ho	С
	RI
Shin, R. H.	SA
Shin, Ran Hee	SA
Shin, S.	Q
Shin, Sang-hoon	Q
Shin, Sung Chul	Q
Shin, Sung-chul	RI
	SJ
Shin, Yooleemi	Eł
Shinaoka, Hiroshi	JC
Shinde, Kiran	Q
Shingubara, Shoso	P
Shinjo, Teruya	A
	RI
Shinkevich, Sergey	Q
Shinohara, Takafumi	Q
Shinoki, Takaho	Sł
Shinozaki, Kazuo	K
Shioda, Tatsutoshi	RI
Shiogai, Junichi	K
CI	~ .

QE10	201	Shiomi,
ED01	66	Shiota, N
RL10	291	Shiozaki
PA12	117	Shipton,
QN11	238	Shirage,
SK21	353	Shirage,
PB05, PB27, PC21, PC23	120, 124, 130	Shirahat
SG01	339	Shirai, N
RK08	287	Shiraishi
PB26	124	
PO18	180	Shiraishi
PH19, PH20, PH23, QF15	150, 206	Shirakas
QC25	196	Shirakav
QB28	191	Shirakav Shirane,
PM08, PM09	168	Shiratsu
RIO4	279	Shiro, K
QI06, QI07, QI08	218, 219 165, 361	Shiroka,
PL10, SM03 AE05	36	Shishido
CD01	50	Shitseva
CC05, QH03, QJ08	50, 213, 222	Shitseva
PO23	181	Shitseva
SD12	332	Shitseva
RL17	292	Shitseva
PB16, PB18, PB35, PB40		Shiue, C
QM18	234	Shivaran
RK09	287	Shlykov,
PO20, SM08	180, 362	Shnirma
RL17	292	Shoji, Ki
QM04	232	Shoshi, /
PN21, QL18	175, 231	Shue, Jy
BH02	46	Shueh, (
PJ31	159	Shukla,
PO14, PO26, QO08,	179, 181, 242, 310,	Shukla,
RQ23, RR12, RR30	312, 315	Shunich
Q011	242	Si, Q.
QK10	225	Si, Qimia
CE04, KE03, RL06, RL14,	52, 110, 290, 292,	Siahatga
RL26, SI13	294, 348	Sichelsch
SA06	317	Sidorenk
SA05	317	Sidorov,
QB28	191	Siemens
QK11	225	Sigov, A
	234 287, 289, 290, 346,	Sigrist, N Sikolenk
RK09, RK22, RL03, SI04, SJ05, SM01, SM12	287, 289, 290, 346, 349, 361, 362	Sikora, N
EH05, PN29, PN30	70, 176, 177	Silotri, S
JD02, RG07	101, 272	Silva, Nu
QM22	235	Silva, Pe
PO18	180	Simmet,
AI02, ED04, PN12, RL04,	39, 66, 174, 290,	Simonell
RL11, SI05, SK10, SK16	291, 346, 352, 353	Simonet
QF19	207	Singh, B
QC18	195	Singh, D
SK06	351	Singh, H
KC02	108	Singh, J.
RP17, SK21	306, 353	Singh, Ji
KC03	108	Singh, N
SE06	334	Singh, P
EG04	69	Singh, R

Shiomi, Mitsuhiro	QP03	245
Shiota, Yoichi	PK20, SI05	163, 346
Shiozaki, Ken	SD09	331
Shipton, E	PK09, QM01	161, 231
Shirage, Parasharam M	DA04	56
Shirage, Parasharam M.	QB29	191
Shirahata, Yasuhiro	QK19, RL25	227, 294
Shirai, Masafumi	IF02, KI03, RL02	94, 113, 290
Shiraishi, Masashi	AI02, ED04, PN12, RL04, SK16	39, 66, 174, 290, 353
Shiraishi, Takuya	SE11	335
Shirakashi, Jun-ichi	EG05	69
Shirakawa, Naoki	QD13	198
Shirakawa, Tomonori	PC08, PG18	128, 145
Shirane, Takashi	RH09	276
Shiratsuchi, Yu	EI05, PK08	71, 161
Shiro, K	SG02	339
Shiroka, Toni	IC02	92
Shishido, Hiroyuki	SB27	324
Shitsevalova, N. Yu.	JC04, PE12, QC22, RH07	100, 138, 195, 275
Shitsevalova, Natalia	QC20, QD12	195, 198
Shitsevalova, Natalja Yu.	QC19	195
Shitsevalova, Natalya	RE13	268
Shitsevalova, Natasha	Q102	218
Shiue, C.h.	PK11, PK15	162
Shivaram, B S	SC05, SE14	326, 335
Shlykov, Maxim Pavlovich	RI22	282
Shnirman, Alexander	1104	97
Shoji, Kimura	QJ01	221
Shoshi, Astrit	AH01	38
Shue, Jyh-ron	IH06	96
Shueh, Chin	PL18	166
Shukla, Dinesh K.	QA01	182
Shukla, Dinesh Kumar	CI04, PH16	55, 149
Shunichiro, Kittaka	QC16	194
Si, Q.	CD01	50
Si, Qimiao	AB01	33
Siahatgar, Mohammad	DB02, QE04	57, 200
Sichelschmidt, Jorg	QD07	197
Sidorenko, A.	CD01	50
Sidorov, V. A.	FB01	73
Siemensmeyer, Konrad	Q102	218
Sigov, Alexander S	PL05	164
Sigrist, Manfred	RF01, SD05	269, 330
Sikolenko, Vadim	QG02	208
Sikora, Marcin	HG04	88
Silotri, Salman	RE06	267
Silva, Nuno	KB06	108
Silva, Petrucio	QL01	228
Simmet, Thomas	AH02	38
Simonelli, Laura	RB28	258
Simonet, Virginie	AC04, IJ03	35, 97
Singh, Braj Bhusan	PG27	146
Singh, D. J.	GF01	79
Singh, H. K	SM18	363
Singh, J. P.	PK11, PK15	162
Singh, Jitendra Pal	PG27	146
Singh, Nirpendra Singh, Pankai	QO01 SC02	240 325
Singh, Pankaj Singh, R S	PE14	325 138
Singh, N.S		.50

Singh, Ravi	SD13	332
Singh, Surendra	QM18	234
Singh, Vikash	RA27	253
Singh, Yogesh	DC03, SG13	58, 341
Sinha, Bhavesh Bharat	QM22	235
Sinha, Brajalal Sinha, S. K	QO19, SN13	243, 366
Sinha, S. K.	BD01	43 284
Sinnecker, Joao Paulo Siratori, Kiiti	RJ13 PI16	153
Sirena, Martin	QL06	229
Sirisathitkul, Chitnarong	RP18, SL10	306, 357
Sirisathitkul, Yaowarat	RP18	306
Sirotti, Fausto	SH03, SH05	343, 344
Sitnikov, Maksim Nicolaevich	SO02	367
Skashita, Naomi	QM07	232
Skaugen, Arvid	CI04, QA01	55, 182
Skidanov, Vladimir	PG02	142
Skokov, Konstantin	DJ04	63
Skokov, Konstantin P.	SO19	369
Skorenkyy, Yuriy	RG13	274
Skorobagatko, Glib A	RG17	274
Skorvanek, Ivan	IB01, QO17	91, 243
Skourki, Y.	QF18	206
Skourski, Y.	GJ03, QI03	82, 218
Skourski, Yurii	PC27, RE09	131, 267
Skrotzki, Richard	PC27	131
Skumryev, Vassil	PA13	118
Slavin, Andrei	JG02	103
Slawska - Waniewska, A.	PH06, PM05	147, 168
Slebarski, Andrzej	PD21	135
Sluchanko, N. E.	JC04, PE12, QC22, RH07	100, 138, 195, 275
Sluchanko, Nickolay	ED06, KI05, QC20, QI02, RE13, SG07, SG08	67, 113, 195, 218, 268, 340
Sluchanko, Nickolay E.	QC19	195
Smadici, S.	PH10	148
Smeibidl, P.	RQ21	310
Smekhova, A.	QG07	209
Smekhova, Alevtina G	PN22	175
Sniadecki, Zbigniew	Q017	243
Snoeck, Etienne	DG05	61
So, Yoshitaka	PC28	131
Sobczak, W.	DD02	58
Soegijono, Bambang	SH01	343
Soh, Joon-young	PM16	170
Soh, Kwang-sup	BH02	46
Soh, Yeong-ah	PJ17	157
Sohn, Jeong Woo	RK09	287
Sohn, Jeong-woo	RK22	289
Sohrabi, Mahdi	PC16	129
Sokolov, D.	GC01	77
Sokolov, Dmitry	GC04	77
Sokolov, V. V.	SG20	342
Sokolovskiy, Vladimir	PI20, QH05	154, 213
Solin, Nikolay	SL17	358
Solorzano, G.	SM17	363
Soltan, Soltan	QL21, RB29	231, 258
Solyom, Jeno	PD15	134
Son, Derac	RQ14	309
Son, Il-ho	RR15	313

Shiojiri, Yasuhiro

Shiokawa, Yohei

Sondezi-mhlungu, B. M. Sone, Keita Song, Fengqi Song, Hoon Song, Hyonseok Song, Hyon-seok Song, Jaeyong Song, Ji Yeon Song, Jin Dong Song, Jindong Song, Jong Han Song, Jong Hyun Song, Jonghyun Song, Jong-hyun Song, Jung-bin Song, Kie-moon Song, Ki-myung Song, Sun Yong Song, Xiaoping Song, Young Sang Song, Young-joon Song, Young-sang Song, Young-yeal Song, Yun Sonntag, Andreas Soon Ok, Park Soriano, N. Sota, Shigetoshi Soulantika, Katerina Souliou, Sofia Michaela Sousa, Marcos Sousa, Marcos Antonio De Souza, Flavio L Sparks, Patricia D Spehling, Johannes Speliotis, Thanassis Spenato, David Spiesser, Aurelie Spinu, Leonard Springell1, R. Sri Ramulu, Torati Srikanth, Hariharan Srinivasan, Elaiyaraju Srinivasan, Radha Srivastava, Amit P Srivastava, Archana Srivastava, D Srivastava, M. K. Stadler, Shane Staerk, Martin Stamps, Robert Stanislavchuk, Taras N. Starikova, Anna Steadman, Paul Stebliy, Maxim Stefanoski, Stevce Stefanowicz, Wiktor Steffens, Paul

SF06 338 QD01, RD13 196, 264 PN27 176 PM07, PO05, RQ04, 168, 178, 307, 353 SK17 289 RK22 287, 346 RK09, SI04 SD12 332 BJ04, QN13 48, 238 QK11 225 QK10 225 Q027 245 117 PA11 SM13 363 SM12 362 PB17, PB36, PB37, PB38, 122, 126 PB39 PN31 177 QA24 186 RO10 302 JJ02, PG21, RM08, SO25 106, 145, 296, 370 PA01 116 PF06, SO06 140, 367 CI03, PA03, PA04 55, 116 PM12, PN31, SO03 169, 177, 367 SB13 322 HF03, HF04, RP16 87, 305 QP21 248 ED03 66 65 EB03 38 AH01 257 RB20 165 PL09 233 QM14 359 SL20 161, 231 PK09, QM01 34 AB02 59, 247 DD06, QP12, QP13 320 SA24 AI05 39 256 RB12 QN01 236 SN13 366 DJ01 63 RP14 305 PN06, QO25 173, 244 PM01 167 RQ06 308 PM01 167 SM18 363 CH01 54 FD03, SH16 74, 345 KE01 110 RA24 253 QJ13, SN09 223, 365 PA15 118 AG05 38 DJ01 63 44, 58 BE03, DD02 33 AA05

Steglich, F.	SC03	326	Sugano, Y	PC04	127
Steglich, Frank	AB02, FB04, JC01, KA03, QC09, QC18, RD12,	34, 73, 100, 107, 193, 195, 264, 325	Sugawara, Hitoshi	PD25, PE11, PE13, QC05, QF04, RD01	136, 137, 138, 19 204, 262
	SB29		Sugawara, Katsuya	SN03	364
Stein, Aaron	KE01	110	Sugawara, Tetsuya	QC10	193
Stein, Jonas	AA05	33	Sugihara, Atsushi	SK33	355
Steinigeweg, Robin	EC05	66	Sugimoto, Koudai	SC10	327
Steinke, Lucia	AB02	34	Sugimoto, Kunihisa	PA12, PG20	117, 145
Stelmakhovych, Olya	BE05	44	Sugimoto, Nozomi	RC18	261
Stepanov, Gennady V	Q014	243	Sugimoto, Satoshi	ID04, QN11	93, 238
Steppke, Alexander	AB02	34	Sugishima, Masaki	QD03, QD15	196, 198
Steren, Laura	QL06	229	Sugiura, Kyohei	RD24	266
Stewart, Greg	DA03	56	Sugiyama, Jun	PI06, PO17, SE08	151, 180, 334
Stewart, J R	SE03, SE05	333, 334	Sugiyama, Kiyohiro	HP11, PD13, QH03	30, 134, 213
Stiles, Mark D.	EE02	67	Sugiyama, Tomoyoshi	QJ08	222
Stobiecki, Feliks	AD05		Suh, B. J.	QF13	206
Stockert, Oliver	GC02, QC24, RE12	77, 196, 268	Suh, Byoung Jin	RN10	298
Stoeffler, Daniel	PI04	151	Suh, Byoungjin	RP13	305
Stoermer, Michael	DG01	61	Suh, Hwansoo	RQ11	308
Stognij, Alexandr	QM08	232	Sukegawa, H	SK29	355
Stoll, Hermann	FD01	74	Sukegawa, Hiroaki	EH01, IF03, IF05, SK11	69, 94, 352
Stollenwerk, Tobias	PF07	140	Sukhorukov, Yurii	SO17	369
Strache, Thomas	SH04	343	Sukovatitsina, Ekaterina	QM05	232
Strack, Christian	SG19	342	Sumimoto, Takaki	RR13	313
	IA03	91	Sumiyama, Akihiko	QC03	192
Strack, Philipp	HC04	85	-		
Strassel, Dominik			Sun, An-cheng	PK06, QP17	161, 248
Strecka, Jozef	QF06	204	Sun, D. L.	GA03	75
Streltsov, Sergey	PA17	118	Sun, Dan	EA04	64
Strempfer, Joerg	CI04, QA01	55, 182	Sun, H.d	QL02	228
Streubel, Robert	SJ07	350	Sun, Jian	RQ02	307
Strichovanec, P.	KC05	109	Sun, Jirong	SO25	370
Strichovanec, Pavel	GG05	80	Sun, Ji-rong	ED05	67
Strigari, Fabio	QC09, QC12	193, 194	Sun, Peijie	KA03	107
Strocov, V.	PC29, PI15	131, 153	Sun, Wei	PK18, QN21	163, 239
Strocov, V. N.	PC01	126	Sun, X. F.	IC03	92
Strohm, Cornelius	HG04, PA24, RH23	88, 119, 278	Sun, Xuefeng	SA19	319
Strokov, Vladimir	HA03	83	Sun, Young	JB03	99
Strydom, A M	SE03, SE05	333, 334	Sun, Yunlei	RB13	256
Strydom, A. M.	CD01, RG05	50, 272	Sundar, Anuraj	QA25	186
Strydom, Andre	SF08	338	Sung, Gun Yong	RL21	293
Strydom, Andre M.	SF06	338	Sung, Ki Young	RR09	312
Strydom, Andre Michael	RD23	265	Sung, Kil-dong	PL19	167
Strydom, Andre Micheal	RH15	277	Sung, N. H.	SB30	325
Stubna, Viliam	RI09	280	Sung, Nakheon	SB10	321
Stupakiewicz, Andrzej	AD05		Sunglae, Cho	PJ15	157
Su, C.w.	PK11, PK15	162	Sunyol, Joanjosep	DJ03	63
Su, Chien-yu	ED02	66	Suresh, K G	PM01, Q004	167, 241
Su, Yixi	QA23	186	Suresh, K. G.	QH06	213
Suard, Emmanuelle	IJ03	97	Suryanarayanan, Ramanathan	PI22	154
Suarez, Jaime Raul	Q026	245	Suter, Andreas	PI05, PN23	151, 175
Suber, Lorenza	RN07	297	Sutherland, Michael	HA05, RE14, SC16	84, 268, 328
Subias, Gloria	EC04	66	Suzuki, Hiroyuki S	JJ05	106
Suda, Ryutaro	EG05	69	Suzuki, Hiroyuki S.	Q109	222
Sudareva, Svetlana	PB22	123	Suzuki, Hisao	KC02	108
Suellow, Stefan	SM02	361	Suzuki, Insao Suzuki, Ippei	QK19, SM20	227, 364
		111		DA01	227, 364 56
Suemasu, Takashi Suergari, Christoph	KF05		Suzuki, Kazunori		
Suergers, Christoph	EG02, QL05	69, 228	Suzuki, Kazuya	QM16, SC23	234, 329
Suess, Dieter	SJ07	350	Suzuki, Kiyonori	RN13	298
Suga, Seiichiro	PF02, RF17	139, 271	Suzuki, Kosuke	PK12	162
Sugano, Ryoko	PJ07, RL01	155, 290	Suzuki, Kousuke	PM02	167

Suzuki, Michi-to	QC04	192
Suzuki, Ryosuke	SA17	319
Suzuki, S	RR08	312
Suzuki, Shota	RO09	302
Suzuki, Shugo	RM06	296
Suzuki, Takafumi	DC04	58
Suzuki, Takao	JI03	105
Suzuki, Yoshishige	KF01, PK20, RL11, SI05, SK10	110, 163, 291, 346 352
Suzuki, Yuri	DH01	62
Suzuki, Yuta	SC04	326
Svec, Peter	IB01, QO17	91, 243
Svoboda, Pavel	QI19	220
Swagten, H.j.m.	PK17	163
Swagten, Henk	CH05, El06, GE02, JE01	54, 71, 78, 101
Swagten, Henk J.	IE02	93
Swatek, Przemyslaw	SE16	335
Syassen, Karl	RB20	257
Sykova, Eva	BH03	46
Syljuasen, Olav F.	QF19	207
Syoji, Daiki	PH15	149
Syromyatnikov, Arseniy	SL27	360
Syrovets, Tatiana	AH02	38
Syudou, Mitsuhiro	PB25	124
Szalowski, Karol	RI03, RI09	279, 280
Szewczyk, A.	PH06, PM05	147, 168
Szlawska, Maria	SE17	336
Sztenkiel, Dariusz	DD02	58
Szymczak, Henryk	JJ03	106
Szymczak, Ritta	JJ03	106
т		
Tabata, Chihiro	QC25, SE07	196, 334
Tabata, Yoshikazu	HJ02, QH13	90, 215
Tacchi, Silvia	GG02	80
Tada, Yasuhiro	RD05	262
Tadataka, Watanabe	QD14	198
Tadera, Shin	SO515	369
Tadokoro, Tomoya	PA06, PA19	116, 119
Taetz, T.	QE23	203
Taga, Yuki	PD08	133
Taguchi, Katsuhisa	FD04	74
Taguchi, Minori	RN11	298
Taguchi, Yasujiro	PA12	117
Taguchi, Yuusaku	AE02, RL02	36, 290
Tahir, Mukarram A.	AH04	38
Tai, M. F.	SB04	320
Tai, Ming-fong	RR01	311
Taira, Tomoyuki	SK06	351
Tajima, S.	CF04	53
Tajima, Setsuko	HA01, PH07, PI07, QB24	83, 148, 151, 191
Tajiri, Takayuki	PG09, SA03	143, 316
Takabatake, T	PC04, QJ05, SE03, SE05, SE13	127, 221, 333, 334 335
		co
Takabatake, T.	EF01, KA02	68, 107
	EF01, KA02 PD24, PE10, QC12,	
Takabatake, T.		68, 107 135, 137, 194, 19 214, 260, 262, 27 341, 342

## AUTHOR INDEX

Takaesu, Yoshinao Takagi, H. Takagi, Hide Takagi, Hidenori Takagi, Seishi Takagi, Yasushi Takahara, Atsushi Takahashi, H. Takahashi, Hidefumi Takahashi, Hirokazu Takahashi, K. S. Takahashi, Ken-ichiro Takahashi, Kohki Takahashi, Masahiro Takahashi, R. Takahashi, Rui Takahashi, Saburo Takahashi, Shin Takahashi, Y K Takahashi, Yoshinori Takahashi, Yoshiyuki Takahashi, Youtarou Takahashi, Yukiko Takahiro, Onimaru Takamatsu, Shuhei Takamichi, Miyazaki Takanashi, Koki Takano, Atsuki Takano, Hideaki Takano, Y. Takano, Yasu Takano, Yoshihiko Takao, Sasagawa Takasaki, Aya Takase, S. Takashi, Sugawara Takashima, Teppei Takata, Emika Takata, Masaki Takatsu, Hiroshi Takayama- Muromachi, Eiji Takayama, Tomohiro Takayama-muromachi, Eiji Takeda, Hikaru Takeda, Hiroki Takeda, Keiki Takeda, Mahoto Takeda, Masataka Takeda, Naoya Takeda, Shogo Takeda, Yukiharu Takegahara, K. Takeguchi, Masaki Takehana, Kanji Takehara, Yuki Takemura, Yasushi Takenaka, K. Takenaka, Yuto

PD14, QC11, QI12, QI13 134, 194, 219 QH09 214 RH21 278 IA04, RB04, RH17 91, 254, 277 PB25, PG10 124, 144 70 EI01 RN03 297 IF01 94 QC06 193 PL17 166 273 RG10 PM08, PM09 168 QG08, QH08, QN20 209, 214, 239 193 QC05 339 SG02 149 PH15 FA03, PJ07, RL01 72, 155, 290 Q007 241 BG01 45 PF08 140 PO03 177 PA12 117 SK31 355 QC16 194 PC06 127 RO05 301 JG03, JI04, KC03, QL15, 103, 105, 108, 230, RL15, SE06 292, 334 QG16 210 PB08, QA19, QJ04, 121, 185, 221, 263, RD06, RD22, SE06 265, 334 258 RB28 42 BC05 QB09, QB17 188, 189 RC12 260 QH20 216 SG02 339 SM10 362 QM29, SJ03 236, 349 QF11 205 HC03, PA12, QF25 85, 117, 208 QE13 201 PH19 150 RB04, RH17, RH21 254, 277, 278 RH06 275 PH19 150 PB07 121 PD10 133 QG16, QH15 210, 215 QI12, QI13 219 RD24 266 PO03 177 RD25, SE23 266, 337 SG11 341 210 QG16 QA09 184 SK13 352 177 PO03 214 QH09 RC04 259

QD12

Takeshita, Nao

Takeshita, Nao	QD12
Takeshita, Soshi	IA04, QO13
Takeuchi, Akihito	EE04, RC06
Takeuchi, Hiromasa	PO02, QL03, RN02
Takeuchi, Tetsuya	HP11, PB27, PD08,
Takeuchi, Tsunehiro	PO17
Takeya, H.	RB28
Takeya, Hiroyuki	QB17
Takeyama, Shojiro	AC03
Takigawa, Masashi	QE08
Takimoto, Tetsuya	JC03, QC04
Takuya, Manabe	BC05
Talaat, Ahmed	AJ02, CJ02
Talham, Daniel R.	KB03
Tamaki, Yuki	PD14
Tamatsukuri, Hiromu	PH05
Tamion, Alexandre	QN06, RK03
Tamura, Dai	PD11
Tamura, Eiichi	KF01, SK10
Tamura, Eiiti	ED04
Tanaka, Arata	PI08, QC12
Tanaka, Daiki	RD10
Tanaka, Hidekazu	QC11, QE22, SK16
Tanaka, K.	HA01
Tanaka, Kenya	PE13, QC05
Tanaka, Koichi	SK13
Tanaka, Masaaki	QM16
Tanaka, Nobuo	RC18
Tanaka, Shigeki	PC23
Tanaka, Shinya	SE06
Tanaka, Takashi	RI15
Tanaka, Terumitsu	QM29, QP09, QP10 QP11, SJ03
Tanaka, Toshiki	SE01
Tanaka, Toshiyuki	QA19
Tanaka, Yasuhiro	RG08
Tanatar, M. A.	IA05, QC23, SB30
Tang, Shengjun	KD03
Tang, Xianwu	SA01
Tange, Hatsuo	PG05
Tanida, Hiroshi	QI10, RD10
Tanigaki, Katsumi	QI10
Taniguchi, Hiromi	GB04
Taniguchi, Kouji	QA13
Taniguchi, Masaki	RD20
Taniguchi, Tomohiro	HF02
Taniguchi, Toshifumi	QE21
Taniguchi, Yusuke	QH04
Taniyama, Tomoyasu	QK19, RL25, SM20
Tao, Lingling	QK02
Tao, Ling-ling	RM01
Tao, Qian	IJ04
Tarafder, K	QM12
Tarantini, C.	JA02
Tarasov, A. S.	QK07
Tarigan, Kontan	PM12
Tartakovskaya, Elena V.	QM02
Tashima, Kasumi	RB06
Tashiro, Hiroki	SA17
Tashiro, Hiroshi	QE22

	198	Tashiro, Kunihisa	RR06
	91, 242	Taskaev, Sergey	PM22, QH05, SO23
	67, 259	Taskaev, Sergey V.	QO12, SO19
02	177, 228, 297	Taskin, Alexey	RC11
8, PD13	30, 124, 133, 134	Tatara, Gen	EE04, FD04, RC06
	180	Tateiwa, Naoyuki	HP11, PD23, RD11,
	258	<b>T</b>	RE10, SB15
	189	Tatetsu, Yasutomi	QC14, QC15
	34	Tati Bismaths, Logane	RP04
	200	Tatsuoka, Sho	PE13, QC05
	100, 192	Tauber, Katarina	EE01
	42	Taufour, Valentin	JC02, SF03
	40, 55	Tayama, Takashi	QE13, SG05
	107	Taylor, A. J.	QB31
	134	Taylor, Alice Elizabeth	HA04
	147	Taylor, J	SE03, SE05, SF09
	237, 286	Taylor, Jon	SE20
	133	Tchoula Tchokonte, M. Bertin	
	110, 352	Tedesco, Antonio Claudio	PO08
	66	Tegus, O.	PM03
	152, 194	Teh, Geok Bee	SL01
	263	Teichert, Anke	QM20
16	194, 203, 353	Tekielak, Maria	AD05
	83	Telegin, Andrey	SL17, SO17
	138, 193	Telegin, Sergey	SL17
	352	Temerov, V	QJ16
	234	Temerov, Vladislav	RA23
	261	Temerov, Vladislav L	CI04
	130	Tennant, A.	RQ21
	334	Tennant, D. Alan	IC05
	281	Tenya, Kenichi	PD07
210,	236, 246, 247, 349	Terada, Noriki	JJ05, PH09, QJ09
		Terasaki, Akira	KD01
	333	Terasaki, Ichiro	PH12, QC06
	185	Terashima, Taichi	SB05
	273	Terazawa, Shinsuke	QH13
0	91, 196, 325	Terentev, Pavel B.	Q108
	109	Terent'ev, Pavel E.	QN10
	316	Tereshina, E. A.	QG07
	143	Terrade, Damien	CA03
	219, 263	Teruya, Atsushi	QI13
	219	Tetsuya, Fujiwara	QI17, QJ01
	76	Tewari, Bhagya Sindhu	PC31
	175	Teysser, J.	SF10
	265	Tezuka, Masaki	HI04
	87	Tezuka, Nobuki	QN11
	202	Thabet Mliki, Najeh	PM06
	213	Thach, Pham Van	SK10
20	227, 294, 364	Thakur, Indu Bala	RA03
	224	Thakur, Rajesh K.	QA14, RF04, RQ06
	295	Thakur, Rasna	QA14, RF04, RQ06
	98	Thakur, Sangeeta	SC06
	233	Thaler, A.	GA01
	98	Thalmeier, Peter	DB02, PD06, QE04
	225	Thamizhavel, A	SE21, SO29
	1.00	Thamizhavel, Arumugam	QJ17, RB24
	169		
	231	Thanh, Pham Thi	AJ06
		Thanh, Tran Dang	AJ06
	231		

311

260

67 100, 337 201, 340

192 83

265 178

167 356

234

223

253

55

310 92 133 106, 148, 222

168 352

249

184, 269, 308

184, 269, 308 326 75

57, 132, 200 336, 371 223, 258 40 40

319, 320 252

358, 369 358

333, 334, 338 336

171, 213, 370

242, 369

67, 74, 259

30, 135, 263, 268,

Thiaville, Andre
Thiem, Luu Van
Thiess, S.
Thiet, Duong Van
Thirumurugan, Arun
Thiruvadigal, John
Thomas, Andy
Thomas, Christopher
Thomasson, A.
Thomasson, Alexandre
Thompson, J D
Thompson, J.
Thompson, J. D.
Thompson, J.d.
Thompson, Joe D
Thompson, Joe D.
Thongmee, Sirikanjana
Thuesen, Christian Aa.
Thunstrom, Patrik
Tian, Jinzeng
Tiberkevich, Vasil
Tiberto, Paola
Timirgazin, Marat
Ting, Yi-wen
Tissier, Matthieu
Titvinidze, Irakli
Tjeng, Liu Hao
Tjernberg, O.
Tjernberg, Oscar
Tjia, May On
Tkachenko, Vira
Tobash, P H
Tobash, P.
Tobash, P. H.
Tobash, Paul H.
Tobola, Janusz
Toews, W. H.
Toga, Yuta
Togawa, Yoshihiko
Tohyama, Takami
Tokiwa, Yoshi
Toko, Kaoru
Tokuda, Yuki
Tokunaga, Masashi
Tokunaga, Y.
Tokunaga, Yo
Tokura, Y.
Tokura, Yoshinori
Tolbert, Sarah
Tomas, Amparo Borrell
Tomas, Milagros
Tomita, Hiroyuki
Tomita, Keitaro
Tomita, Norikazu
Tomita, Yusuke
Tomita, Yuusuke

AJ01, SH02	39, 343
SL22	359
SK24	354
PN30, QO22	177, 244
AJ04	40
SK32	355
DI03, IF04	63, 94
PD04	132
SA06	317
SA05	317
EA03	64
SC03, SC16	326, 328
	73, 189, 192, 322
PC19	129
JC01	100
GC03, PC03, RD09,	77, 127, 263, 333,
SE02, SF01	337
GJ04	83
BC02	42
RG06	272
PK03	160
JG02	103
AG02	37
BF04	45
PL18	166
SL24	360
PE05, PE07	136, 137
PI08, PI17, QC09, QC12	152, 153, 193, 194
PC29	131
HA03, SE08	83, 334
PH17	
	149
JG01	103
EA03	64
SC16	328
GC03	77
RD09	263
QH18	215
HA05	84
RF10, RF11	270
AF03, DF03, PG15,	37, 60, 144, 211,
QG22, QG28	212
CA01, EB03, SC10	48, 65, 327
FB03	73
KF05	111
RP09	304
PA06, PA18, QA09,	116, 118, 184, 185,
QA17, QG08, RA06,	209, 250, 275
RH04	
RA05	250
PD09	133
PA14, RA05, RG10	118, 250, 273
HP31, PA12, PH08, PI07	30, 117, 148, 151
QA15	185
RO08	302
Q016	243
KF01	110
SC04	326
PF04	139
JD02, RI02	101, 279
QB17	189
A.T. 1.1	

AJ01, SH02

39, 343

## AUTHOR INDEX

Tomitaka, Asahi	PO03	177
Tomiyama, Yo	PE11	137
Tomiyasu, Keisuke	RD14	264
Tompsett, David	PC02	127
Tompsett, David A.	PC09	128
Tong, Peng	BF02	44
Tonnerre, Jean-marc	AD02	35
Toperverg, Boris P.	DG01	61
Torikachvili, Milton	SF13	339
Toriumi, Shingo	RK08	287
Toriyama, Tatsuya	KG03, PF10, PF15	111, 140, 141
Torrejon, Jacob	AJ01, SH02	39, 343
Tortarolo, Marina	QL06	229
Toru, Shigeoka	QI17, QJ01	220, 221
Toshiro, Sakakibara	QC16	194
Toshiro, Takabatake	QC16	194
Toth, Sandor	PA23	119
Tou, Hideki	PD25, QB17, SF03, SG12	136, 189, 337, 341
Toulemonde, Olivier	PI17	153
Tournus, Florent	KB05, QN06, RK03	108, 237, 286
Toussaint, Jean-christophe	SH03, SH06	343, 344
Toyoda, Shingo	PG20	145
Toyokawa, Shuhei	ED04	66
Toyoki, Kentaro	PK08	161
Tozawa, Katsunori	PJ12	156
Tozer, Stan	PC09	128
Tran, Duc H.	PB03	120
Tran, Hoang Duc	RN19	299
Tran, L M	QB06	188
Tran, Vinh Hung	QB06	188
Tretiakov, Oleg	SD17	332
Trioni, Mario Italo	PL21	167
Trohidou, Kalliopi	CB03	49
Tronov, Anton	SG17	342
Trupina, L.	SA11	318
Tsai, C.c.	PK11, PK15	162
Tsai, Du-cheng	PK02	160
Tsai, M.c.	PK11	162
Tsarevskii, Sergey L.	SC25	329
Tsay, Jyh-shen	IH06, PK02, PL03, PL13,	96, 160, 164, 166,
	SM06	361
Tselepi, Marina	BH01	46
Tseng, Yu-chieh	SM06	361
Tsevelmaa, Tumurbaatar	KI01	112
Tsirlin, A.	QF22	207
Tsuchiura, Hiroaki	RF11	270
Tsuchiura, Hiroki	KH02, QB27, RF10	112, 191, 270
Tsuchiya, Katsuhiko Tsuchiya, Yashinari	RD24	266
Tsuchiya, Yoshinori	PH09	148
Tsuei, K. D. Tsuei, Ku Ding	QB21	190 194
. 5	QC12	194
Tsuei, Ku-ding Tsuji, Naoto	QC09 EB02, RF05	65, 269
-		
Tsujii, Hiroyuki Tsujikawa, Masahito	QL11, QL12, QL13 AE02, KI03, PL06, RL02	229, 230 36, 113, 164, 290
Tsujimoto, M.	AEU2, KIU3, PLU6, KLU2 KJ01	36, 113, 164, 290 113
Tsukahara, Hiroshi	RJ06, SI16	283, 348
Tsukamoto, Arata	BG03, RK08, RK17	265, 346 45, 287, 288
Tsunashima, Shigeru	RL08	43, 287, 288 291
Tsunekawa, M.	SG11	341
Sancharra, m.	2011	

Tsunoda, Masakiyo	PL17	166
Tsunoda, Yorihiko	PI16	153
Tsurkan, Vladimir	PA23	119
Tsurunaga, Aiko Tsuruta, Kausuka	QN02	236 202
Tsuruta, Kousuke		
Tsutaoka, Takanori Tsutsui, Kazuhiro	QI06, QI07, QI08, QN02	218, 219, 236
Tsutsui, Kazunno Tsutsui, Kenji	PJ18, RL16 CA01, SC10	157, 292 48, 327
Tsutsumi, Kenji	PB29, PB30, PB32	124, 125
Tsutsumi, S.	QB23	190
Tsutsumi, Tomoaki	R003	301
Tuan, Duong Anh	PN29, QO22	176, 244
Tugarinov, V	QJ16	223
Tumanov, Vadim	SC26	329
Tung, Min-jue	QN08	237
Turnbull, M M	BC05	42
Tusche, C	RP20	306
Tutsch, Ulrich	HC04	85
Tuuli, Elina	SB16, SK07	322, 351
	5510, 5107	522,551
U		
Ubana, Tatsuya	BG03	45
Uchida, Atsuko	JD01	100
Uchida, Kaori	QG15	210
Uchida, Ken-ichi	RL15	292
Uchida, Masakazu	SL03	356
Uchida, Shinichi	CA02	48
Uchida, Shin-ichi	PB05	120
Uchida, Takashi	EF04	68
Uchima, Kiyoharu	PD14, QI12, QI13	134, 219
Uchiyama, Toshiyuki	RD19	265
Uchiyama, Yu	RR06	311
Uda, Yuuki	PN24	176
Udagawa, Masafumi	PG23	146
Udagawa, Masayuki	PC05, PG16, RD10, SG15	127, 145, 263, 342
Udod. Lubov Victorovna	SO02	367
Uebayashi, Kazuhiko	QH23	216
Uechi, Ayaka	QH14	215
Ueda, H.	PH03	147
Ueda, Hiroaki	AC03	34
Ueda, Hiroaki T	RC06	259
Ueda, Kenji	PB34, PJ12, PL12, RC18,	125, 156, 165, 261,
	SA07, SA17, SK18, SO20	317, 319, 353, 370
Ueda, Koichi	QD15, RD24	198, 266
Ueda, Shigenori	KF05, RD20	111, 265
Ueda, Suguru	RF01	269
Ueda, Takuya	EC03	65
Ueda, Tomonari	QE10	201
Ueda, Y.	PH03	147
Ueda, Yutaka	IA04, PH20, PH23, QD03, QF24, RH16, SB14, SF04	91, 150, 196, 207, 277, 322, 337
Uehara, Masatomo	SB03	320
Uematsu, Daisuke	PG14	144
Uemura, Tetsuya	SK06	351
Uemura, Y. J.	GF01	79
Ueno, T.	PH14	149
Uesaka, Masanori	QF08	205
Uhlir, Vojtech	SH03, SH06	343, 344

Uhlirova, Klara	BE05, PD18, PD19	44, 134, 135
Uimin, Michael	QL07	229
Uji, Shinya	SB05	321
Ukleev, Victor	QM08	232
Ullrich, Carsten A.	RL12	291
Ulrich, V W	SC05	326
Ulysse, Christian	FD03	74
Um, Jaegon	RI19	282
Um, Youngje	QB03	187
Umehara, Izuru	QJ06	221
Umeo, K	PC04, QJ05	127, 221
Umeo, K.	EF01, KA02	68, 107
Umeo, Kazunori	QD09, RD04	197, 262
Umetsu, Hiroshi	QA13	175
Umetsu, Nobuyuki	SI06	346
Umetsu, Rie	RP07	304
Umetsu, Rie Y	QG08	209
Umetsu, Rie Y.	SK23	354
Umezawa, Jin	R006	301
Uneda, Kentaro	BC03	42
Unno, Tetsuya	PM02	167
Upton, Mary	DC03	58
Urbanek, Michal	SH06	344
Urbano, Ricardo R.	JA02	98
Urcelay - Olabarria, Irene	PA13	118
Urtizberea, Ainhoa	KB06	108
Useinov, Artur	PJ06	155
Ushiki, Yuji	Q007	241
Ustinov, Vladimir	SK02	350
Ustinov, Vladimir V.	SK01	350
Uthayakumar, S.	DH03	62
Uto, Haruka	PM09	168
Utsumi, Yuki	RD20	265
Utsumiya, Kazuhisa	JG03	103
Uvarov, N.v.	QG27	212
Uwatoko, Y.	QH09, RD15	214, 264
Uwatoko, Yoshiya	PD14, QD17, QI12, QI18,	134, 199, 219, 220,
	QN20, RD06, RD22,	239, 263, 265, 267,
	RE04, RP05, SC24, SE01, SE11	304, 329, 333, 335
	JETT	
Uyeda, Chiaki	PN11, RQ03	173, 307
Uykur, E.	HA01	83
Uykur, Ece	PH07	148
Uyyanonvara, Bunyarit	RP18	306
V		
V. Leitao, J.	RQ12	309
V. Loehneysen, Hilbert	RE12	268
V., R.reddy	QD04	197
Vaezi, H.	PI23	154
Vaghar, R.	SA04	316
Vaknin, David	RB18	257
Valencia, Sergio	AE03, EC04, IH01	36, 66, 95
Valenta, Jaroslav	QJ14	223
Valenti, Roser	RG14, SF14, SG22	274, 339, 343
Valenzuela, Belen	QB18, QB19	190
Valenzuela, S. O.	DI01	62
Valiska, Michal	HJ04, QC21	90, 195
Val'kov, Valery V.	QC02	192
Valldor, Martin	PI08	152

Vallejo, Emmanuel	QG21	211
Van Den Brink, Arno	EI06	71
Van Den Brink, Jeroen	DC02	58
Van Der Heijden, E.h.m.	PK17	163
Van Der Heijden, Mark	JE01	101
Van Der Linden, Peter J. E. M.	HG04, PA24, RH23	88, 119, 278
Van Dommelen, P.	RQ12	309
Van Eijck, L.	SO08	368
Van Lierop, Johan	PL18	166
Van Loosdrecht, P.h.m	PH17	149
Van Loosdrecht, Paul H. M.	HG04, PA24	88, 119
Van Wees, Bart J.	GI05	82
Vanacken, Johan	QC19, QC20	195
Vanecek, Vaclav	BH03	46
Varaprasad, B	BG01	45
Varga, Marek	IB01	91
Varga, Rastislav	PJ25	158
Varma, C. M.	SB31	325
Varnakov, S. N.	QK07	225
Vasambekar, P. N.	QN14	238
Vasilakaki, Marianna	CB03	49
Vasiliev, A.	QF22	207
Vasyuchka, Vitaliy	HD01	85
Vavassori, Paolo	GG05	80
Vavilova, E.	QE23, QF22	203, 207
Vazquez, Manuel	AJ01, DG03, GG04,	39, 61, 80
vazquez, manuer	GG05	55, 61, 66
Veber, Philippe	EH03	70
Vecchini, Carlo	DH03	62
Vedmedenko, Elena	IH03	96
Vehsthedt, Erin	IE03	93
Veis, Martin	KD03	109
Velazquez, Matias	EH03	70
Velicanov, Dmitrii Anotol'evich		367
Velikanov, D	QJ16	223
Velikanov, Dmitriy	PI11	152
Venkataiah, Gorige	RL25	294
Ventura, Cecilia I.	SI10	347
Venu, Reddy	SN13	366
Vera-marun, Ivan J.	GI05	82
Verba, Roman	JG02	103
Vergnaud, Celine	AI04, CG03, EI04	39, 53, 71
Verma, G. D.	SM18	363
Veron, Emmanuel	EH03	70
Vervacke, Celine	SJ07	350
Viart, N.	SA06	317
Viart, Nathalie	SA05	317
Vieira, V R	SB21	323
Vieira, Valdemar	PB10, PB33	121, 125
Vieira, Valdemar Das Neves	PB09, PB11, RB09, SB18	121, 125
Vijaysankar, Kalappattil	SL21	359
Vila, Laurent	Al04, GE01, QK15,	39, 78, 226, 293,
	QK16, QK17, RL24, SL24	360
Vilanova, Enrique	CH03	54
Vilas-vilela, J. M.	CB02	49
Villamor, Estitxu	HE02	86
Villegas, J. E.	ED03	66
Villegas, Javier E.	BD04	43
Vinai, Franco	AG02	37
Vinogradov, Alexey P.	HD05	86

## AUTHOR INDEX

Viraphong, Oudomsack Virgin, Lawrence. N. Vlachos, Athanasios Vlaic, Petru Vogel, Andreas Vogel, Jan Vojta, Matthias Volegov, Alexey Volkov, N. V. Volkov, Nikita Volkova, O. Volkova, Olga Vomir, M. Vomir, Mircea Von Bergmann, Kirsten Von Issendorff, Bernd Von Molnar, Stephan Vorderwisch, Peter Vorobiev, Alexei Vorotynov, Alexander Vovk, A. Vu, Khai Van Vundisa, Khuselwa Vuong, Nguyen Van Vyalikh, Denis V.

#### w

Wachnicki, Lukasz

Wada, Eiji Wada, Hirofumi Wada, Kakeru Wadley, Peter Wagh, Aditya A. Wagner, M. Wahlstrom, Erik Waintal, Xavier Wakabayashi, E. Wakabayashi, Katsunori Wakabayashi, Yusuke Waki, Takeshi Wakimoto, Shuichi Wakiwaka, Hiroyuki Wakiya, Kazuhei Wakiya, Naoki Waldmann, Oliver Walker, H. C. Walker, Helen Wallacher, Dirk Walsh, M. J. Walters, Andrew C. Wang, Chih-jen Wang, Chin-wei Wang, Daohan Wang, Fen Wang, Hailong Wang, Huijie Wang, James Wang, Jianli Wang, Jiawei

EH03 AH04 RP15 SK09 HG02, RJ05, SH04 SH03, SH05, SH06 EA02 QG04 QK07 PI11 QF22 QE27 PO23 FC02, RK05, RK13 IH02, IH03 KD01 QA11, QJ11, QK23 PO25 QM08 IJ01 KC05 CJ05 QH22 EJ02 CD03, PD12 DD02 QK19

227 AG03, IG05, QD03, QD15, SB14 322 353 SK21 44, 344 BE05, SH09 147 PH04 152 PI09 358 SL15 PJ13 156 PL04 164 SD05 330 QA12 184 HJ02, QH13 90, 215 BA04, SC23 41, 329 RR06 311 RD04 262 KC02 108 BC02 42 QN01 236 AF01 36 QG26 212 350 SK03 RB20 257 QA20 185 PC18, QA20 313 RR18 JB03 99 61 DG06 SN12 366 SL01 356 EF02, QG09, QG26 SA19, SA20 319

58 37, 95, 196, 198, 129, 185 68, 209, 212

Wang, Jingdai	PK18	163	Weil, M.	QF22	207
Wang, Junfeng	PA18, QF14	118, 206	Weinert, M.	RM04	295
Wang, Ke	Q024	244	Weinert, Michael	RB26, SO10	258, 368
Wang, L.m.	QM27	236	Weller, D.	BG02	45
Wang, Li Na	RD12	264	Wells, Barrett O.	HA03	83
Wang, N L	QB11	188	Wen, Hai-hu	JA01	98
Wang, Ning	RJ03	283	Wen, J. S.	GA01	75
Wang, S.	QE19	202	Wen, Z.c	SK29	355
Wang, S. J.	SB04	320	Wen, Zhechao	EH01	69
Wang, Shen	PN05, QL04	172, 228	Wen, Zhenchao	IF05	94
Wang, Sheng Hao	SB19	323	Wenderich, Sander	DH02, SA10	62, 317
Wang, Shuang	IC02, QF23	92, 207	Wenderoth, Martin	KA01	106
Wang, Tao	QN23	240	Wermeille, D.	QN01	236
Wang, W.	GC01	77	Werner, Michael	Q002	241
Wang, Weizhong	RL32	295	Werner, Philipp	EB02, RF05	65, 269
Wang, X. M.	IC03	92	Weymann, Ireneusz	DB03	57
Wang, Xiao Lei	DH04	62	Whangbo, Myung-hwan	PH11, QJ18	148, 223
Wang, Xiaolei	RJ03	283	Wheeler, Elisa	RB07	255
Wang, Xiaolin	SD01	330	White, Rudelle	QH22	216
Wang, Xiaoping	РК03	160	White, Steven	RF13	271
Wang, Xiuhe	RR18	313	Wieckowski, J.	PH06, PM05	147, 168
Wang, Xuefeng	PN27	176	Wienholdt, Sonke	FC03	74
Wang, Xuhui	EE03, PJ29, PJ34	67, 159, 160	Wierschem, Keola	DC04	58
Wang, Yin-gang	RA12	251	Wiesendanger, R	RN12	298
Wang, Ying-chieh	SM06	361	Wiesendanger, Roland	AF02, GH02, HF03,	36, 81, 87, 88, 96,
Wang, Yu	PG21, SO25	145, 370		HF04, HG03, IH02, IH03,	144, 28, 305
Wang, Yun-peng	RM01	295		PG12, PP05, RP16	
Wang, Z.s.	RO14	303	Wiesenmayer, Erwin	QB22	190
Wang, Zhe	PA23, RH24	119, 278	Wieser, Robert	IH03	96
Wang, Zhi-hong	ED05	67	Wijnheijmer, Ineke	BE05	44
Wanikawa, Yasushi	QF25	208	Wildes, Andrew	1)03	97
Warczewski, Jerzy	PG26	146	Wildes, Andrew R.	GF03, PG01	79, 142
Ward Jones, Sarah	KB01	107	Wilhelm, Fabrice	PA08, RP15, RP21	117, 305, 306
Ward, S	RI05	279	Wille, Hans-christian	QM19	234
Ward, Simon	QG14	210	Willers, Thomas	PI08, QC09, QC12	152, 193, 194
Warin, Patrick	EI04	71	Williams, T. J.	GF01	79
Watahiki, Masanori	QI10	219	Windsor, Yoav W	RK01	286
Watanabe, C.	KJ01	113	Winkler, H.	CD01	50
Watanabe, Eri	BC04	42	Winterlik, Jurgen	QN19	239
Watanabe, Hiroshi	PC08, PG18	128, 145	Wintz, Sebastian	RJ11, SH04	284, 343
Watanabe, Isao	PO17, QB02	180, 187	Wirth, Steffen	JC01, PH04	100, 147
Watanabe, Kazuo	GJ02, QG15, QH08,	82, 210, 214, 220	Wisniewski, Andrzej	QB30	192
Marco 1 Marco 11	QI18	477	Wisniewski, Piotr	SE16	335
Watanabe, Masatoshi	PO03	177	Witkowski, Bartlomiej S	DD02	58
Watanabe, Masayuki	KG03	111	Wittlin, Aleksander	QB30	192
Watanabe, S.	QB28	191	Woelfle, Peter	EB04	65
Watanabe, Shintaro	QI13	219	Woike, Theo	AB02	34
Watanabe, T.	SG11	341	Woitschach, Sarah	RE12	268
Watanuki, Ryuta	SC23	329	Wojcieszyk, Daniel	PG26	146
Wattiaux, Alain	PI17	153	Wojek, B. M.	PC29, PI15	131, 153
Wawryk, Ryszard	SB26	324	Wojek, Bastian M.	HA03, SE08	83, 334
Weber, Katharina	AB02	34	Wolf, Bernd	HC04, QE14, RI13	85, 201, 281
Wechke, E.	PC22	130	Wolf, Thomas	HA02	83
Wees, Bart J. Van	AI05	39	Wolff-fabris, Frederik	PB09, PB11, RB09, SB18,	121, 255, 323, 331
Wei, D. H.	PK01	160	Malfle Deter	SD07	111
Wei, Dahai	HE03	86	Wolfle, Peter	KG01	111
Wei, Jianqiang	QN23	240	Won, Choongjae	PI10	152
Weickert, Franziska	JD01	100	Won, Chung Jae	SE02	333
Weigand, Markus	RJ05	283	Won, Hyuk	QN04, QP04	237, 246
Weihe, Hoegni	BC02	42	Won, Jaehyuk	PN08, PN09, QM04	173, 232

Won, Jonghan	KD02	109
Won, Mi Hee	PA22	119
Wong, Franklin	DH01	62
Wong, Yat Choy	SLO1	356
Woo, Byung-chul	PM25, RR31 PO23	171, 316
Woo, J. H. Woodcock, Tom G	F023 KH01	181 112
Woodcock, Iom G Wosnitza, J.	CC04, GJ03, QF02,	50, 82, 204, 218
WUSHIIZA, J.	QF03, QI03	50, 82, 204, 218
Wosnitza, Joachim	PC27, PD05, RD07, RE09	131, 132, 263, 267
Wosnitza, Jochen	QD07, SD07	197, 331
Wray, L. A.	SD04	330
Wray, L. Andrew	HI02	89
Wrobel, Piotr	RF03	269
Wu, An-wei	SM06	361
Wu, Chau-chung	AH03	38
Wu, Chun-ming	PO25	181
Wu, Guang Heng	QG09	209
Wu, Han Chun	PJ14	157
Wu, Han-chun	SH13	345
Wu, Hsian-yuan	PK06	161
Wu, Hsing-hsuane	PL13	166
Wu, Hua	PIO8, SB02	152, 320
Wu, Hung-cheng	RN17	299
Wu, J. C.	SB19	323
Wu, J. W.	PO23	181
Wu, Jong-ching	PC30, SH08, SL18	131, 344, 359
Wu, Lin	JF05	103
Wu, Qiong	RO13, RO14	302, 303
Wu, Sc.	QG24	212 246
Wu, Sheng Yun Wu, T. C.	QP07 SB19	323
Wu, Te-ho	PL18	166
Wu, Tian-chiuan	PC30, SH08, SL18	131, 344, 359
Wu, Wenlong	EA04	64
Wu, Xiaoshan	PJ05, PK05, PN05, QL04,	
	RA13, RA14, RR29	251, 315
Wu, Yicheng	RA02	249
Wu, Yizheng	PL02, SM05	164, 361
Wulferding, Dirk	QE27	203
Wulfhekel, Wulf	GI03	82
Wunderlich, Joerg	IE03, IF01	93, 94
Wurmehl, S.	RB17	256
Wurmehl, Sabine	QH24	216
х		
Xaba, Bongani	QH22	216
Xavier, Allan M	SL20	359
Xia, Zhengcai	QI11	219
Xiang, Tao	FE02	75
Xianmin, Zhang	CH02	54
Xiao, F.	QF02	204
Xiao, Gang	GH03	81
Xiao, Jiang	PJ34	160
Xiao, Wen	DD03	58
Xiao, Yu	QK01	224
Xiaodi, Wen	Q015	243
Xiaopeng, Li	Q021	244
Xin, Yan	PA21	119
Xing, Y. T.	SM17	363

## AUTHOR INDEX

101

176

103

164

330

89

209

172

75

34, 98

256

96

164

243, 244

184, 222, 227

Xiong, Peng Xu, Cenke Xu, Jianbin Xu, Jianli Xu, Qingyu Xu, S.-y. Xu, Suyang Xu, Xiao Xu, Xiaohong Xu, Z. J. Xu, Zhu'an Xue, Qi-kun Xue, Xiaobo Xuliang, Lv

#### Υ

Y, Nguyen Thi Tinh Yadav, A. D. Yadav, C. S. Yadavalli, Tejabhiram Yagasaki, Katsuma Yagi, Takehiko Yakata, Satoshi Yako, Hitomi Yakovchuk, Viktor Yakovleva, M. Yakushiji, K. Yakushiji, Kay

Yamada, H. Yamada, I. Yamada, K. Yamada, Kazuyoshi

Yamada, Keisuke Yamada, Shigeki Yamada, Shinya Yamada, Tsutomu Yamada, Yuji Yamagami, Hiroshi Yamaguchi, Akira Yamaguchi, Masahiro Yamaguchi, Masahito Yamaguchi, Masanori Yamaguchi, Nobuhiro Yamaguchi, Shuhei Yamaguchi, Yasuo Yamaki, Yuki Yamamoto, Daisuke Yamamoto, Eetsuji Yamamoto, Etsuji

Yamamoto, Keiichi Yamamoto, Koji Yamamoto, Masafumi Yamamoto, Masaki Yamamoto, Shoji Yamamoto, Takashi QA11, QJ11, QK23 JD05 PN27 JF05 PL01 SD04 HI02 QG08 PN04 GA01 AB04, JJ04 RB13 II01 PL01 QO15, QO21

QE16 202 PN01 172 PC11 128 PO21, QN26 180, 240 QI12, QI13 219 PD11, PD22 133, 135 AG03, RJ08, RK21, SI12 37, 283, 289, 347 SO24, SO26 370, 371 PJ27 159 QE23 203 CE03, CE06, ID02, RJ07 51, 52, 93, 283 BI01, HF01, IE04, KF01, 47, 87, 94, 110, 163 PK20 PL14 166 SG02 339 201 QE15 BA04, PB29, PB30, PB32, 41, 124, 125, 211, QG18, RP01 303 HG02 88 QA02, QD02 182, 196 HE01, RK10 86, 287 PO03 177 QB15 189 HP11, RD25, SE23 30, 266, 337 QC03 192 QM07, RK04 232, 286 QK19 227 SG01 339 QG19 211 QG19 211 QG18, RP01 211, 303 148 PH08 SB23 324 322 SB15 HP11, IG04, PC28, PD13, 30, 95, 131, 134, PD23, QC03, RD11, 135, 192, 263, 268, RE10, SE23, SF03 337 133 PD10 135 PD20 351 SK06 215 QH13 98, 103 IJ05, JG04 QC07 193

Yamamoto, Yousuke	RK23	289	Yang, Shuo-wang
Yamamura, T.	QJ10	222	Yang, Tao
Yamanaka, Rina	RE04	267	Yang, Tim
Yamanaka, Shoji	PE10	137	Yang, Yang
Yamane, Yuta	JE02, JF02	102	Yang, Yuanjun
Yamanoi, Kazuto	RK21	289	Yang, Zhaorong
Yamanouchi, M.	SH07	344	Yano, Kazuo
Yamanouchi, Michihiko	QP14, SI17	247, 348	Yano, Midori
Yamasaki, Y.	CF04, PL14	53, 166	Yano, Shin-ichiro
Yamasaki, Yuichi	PH08	148	Yano, Shinya
Yamashita, Fumitoshi	JH01, RO03	104, 301	Yanoh, Takuya
Yamashita, Makoto	RF10, RF11	270	Yansen, Widi
Yamashita, T.	RD15	264	Yanushkevich, Kazimir
Yamashita, Tetsuro	QD17, RD18, RD20, SE11	199, 265, 335	Yao, Kailun
Yamauchi, Hiroki	QI16	220	Yao, Kai-lun
Yamauchi, Kunihiko			
	AA02, EC03, PA12	33, 65, 117	Yao, Y. D.
Yamauchi, Tooru Yamaura, Junish	SF04	337	Yao, Yeong-der
Yamaura, Junich	Q013	242	Yaresko, Alexander
Yamaura, Junichi	JA04, QE05, SG04, SG14	99, 200, 340, 341	Yarikov, Stanislav
Yamaura, Jun-ichi	IA04, PC17, PG16, SB05	91, 129, 145, 321	Yashima, Mitsuharu
Yamazaki, Yohei	PG07, QH21	143, 216	Yashin, Shadi
Yan, Biao	AJ03, JI06	40, 105	Yashina, Lada V
Yan, Binghai	RC10	260	Yasin, Shadi
Yan, Guo Jhih	RR26	315	Yasuda, Sho
Yan, Li Qin	PA04	116	Yasuda, Takashi
Yan, Li-qin	JB03	99	Yasufuku, Yoshimasa
Yan, Zhongjie	SL02	356	Yasuhiro, Morosawa
Yanagi, Kazuhiro	SO515	369	Yasui, Akira
Yanagihara, Hideto	PI16, QM16, SJ06	153, 234, 349	Yasui, Yukio
Yanagisawa, Tatsuya	QC25, RD07, SE07	196, 263, 334	Yasuoka, Hiroshi
Yanai, Takeshi	JH01, PM08, PM09,	104, 168, 301	Yeche, Nicolas
	RO03		Yehia, M.
Yanase, Youichi	PC06	127	Yelland, Ed
Yang, Cang-seob	RQ23	310	Yelland, Edward A
Yang, Chan-ho	GB05	76	Yellen, Benjamin B.
Yang, Charles Shieh-yueh	AH03	38	Yen, Nguyen Hai
Yang, Dong Gyu	PB17, PB36, PB37, PB38,	122, 126	Yeom, Han Woong
iang, bong oya	PB39	.22, .20	Yeom, Jaehoon
Yang, Dong-seok	QH01	213	Yeon, Duhyung
Yang, Hangfu	RO13	302	Yermakov, Anatoly
Yang, Hong-chang			Yi, Chang Hyun
5. 5 5	AH03, QM27	38, 236	
Yang, Hung-duen	QA20, RN17	185, 299	Yi, Hee-taek
Yang, Hyunsoo	EG03, GE03, HD02,	69, 78, 85, 156, 158,	Yi, Il Gu
	PJ08, PJ20, PJ21, RL19, SI02, SK05	293, 346, 351	Yi, Jiabao
			Yi, Su Do
Yang, In-sang	PH14, SA16	149, 318	Yildiz, Fikret
Yang, J. M.	RP10	304	Yim, H. I.
Yang, Jinhu	PC02, PC09	127, 128	Yimnirun, Rattikorn
Yang, Jong-heon	RL21	293	Yin, Xiaofei
Yang, L.	QB12	189	Yin, Xiaolu
Yang, Lei	РК03	160	Yoda, Yoshitaka
Yang, Lifeng	SA20	319	Yogi, Mamoru
Yang, Lin	SB29	325	Yokoi, Atsushi
Yang, Luke	RO15	303	Yokoo, Tetsuya
Yang, S.y.	QM27	236	Yokota, Saeko
Yang, Sangsun	SL29	360	Yokoyama, Chiori
Yang, See-hun	RK22	289	Yokoyama, Hisatoshi
Yang, Sen	JJ02, RM08, SO25	106, 296, 370	Yokoyama, Makoto
Yang, Seolun	AD03	35	Yokoyama, T
<i></i>			

SK34	356
PH02	147
KC03	108
SL06, SM14	357, 363
SA19	319
QA21	186
PG05, QI04	143, 218
QE22	203
QG28	212
PO02, QL03, RN02	177, 228, 297
PO02, QL03, RN02	177, 228, 297
RA21	252
QG05	209
PN10, PO01, RM03,	173, 177, 295, 366
SO01	
QK18	227
PK01	160
IH06	96
SC07	326
PJ27	159
DA04, QB29	56, 191
RD07	263
PN22	175
RE09	267
RM06	296
QH21 RP19	216 306
QD14	198
RD25, SE23	266, 337
DA01, PH12, QC06	56, 148, 193
PC03	127
AB02	34
QE23	203
GC01	77
SE12	335
AH04	38
AJ06	40
RC14	261
RH08	276
RQ14	309
QL07	229
QH29, SC08	217, 326
QF20	207
RI21	282 231
QL19 RI19, RI20	282
EI01	70
PM15	169
SL10	357
SN01	364
R007	301
QD11	198
QH14	215
JI01	105
PI07, RH13	151, 276
SK08	351
QF04	204
RF10	270
PD07, SE07	133, 334
QM12	233

Yokoyama, Takehito	PJ18, RC06	157, 259
Yokoyama, Toshihiko	AG04, El01	38, 70
Yonemitsu, Kenji	RG08	273
Yoneyama, Naoki	GB04	76
Yoo, Dongwon	PO16	179
Yoo, J. H.	RP10	304
Yoo, Kyung Jun	JA03	98
Yoo, Myoung-woo	DE05, RJ15	60, 285
Yoo, Sang-im	PJ19, QO28, RH18	157, 245, 277
Yoo, Taehee	PN08, QL08	173, 229
Yoo, Woosuk	PK17	163
Yoo, Y. J.	QG27, RA07	212, 250
Yoo, Yeong-min	BH02	46
Yooleemi, Shin	PJ15	157
Yoon, Dong Young	PA09	117
Yoon, Jungbum	FD03, RK02, RK06,	74, 286, 345, 349
	SH16, SJ02	
Yoon, Kwan-seob	PO14	179
Yoon, S.	PO23, QF13	181, 206
Yoon, Seok Soo	PM28	172
Yoon, Seungha	PK04, QK03	161, 224
Yoon, Sook Soo	SN13	366
Yoon, Sung Wook	PI13	152
Yoon, Sungwon	RN10	298
Yoon, W. S.	RQ08	308
Yoon, Woon-ha	DJ02	63
Yoshichika, Onuki	RD19	265
Yoshida, Hajime	Q105	218
Yoshida, Hiroyuki	QE03, QE05, QE08,	199, 200, 203, 277
	QE27, RH14	
Yoshida, K.	CG02	53
Yoshida, Kazuetsu	BG04	46
Yoshida, Makoto	QE08	200
Yoshida, Tohru	RH01	274
Yoshida, Tomohiko	KD03	109
Yoshida, Tsuneya	RC01	259
Yoshida, Yasuo	IH02	96
Yoshida, Yusuke	PG10	144
Yoshihara, Akira	SL04	356
Yoshihara, Tatsuhiko	RC18	261
Yoshihiro, Kuroiwa	QD14	198
Yoshii, Shunsuke	RH03	275
Yoshikawa, Takuro	PI16	153
Yoshiki, Takano	QD14	198
Yoshimi, Kazuyoshi	GB03	76
Yoshimoto, Tatsuya	RI01	279
Yoshimura, Kazuyoshi	PC02, PC09	127, 128
Yoshimura, Satoru	RP06	304
Yoshio, Miura	KI03	113
Yoshioka, Hideo	RG09	273
Yoshitake, Junki	PE18	139
Yoshiuchi, Shingo	PD08, PD13, PD16	133, 134
Yoshiya, Uwatoko	QI17, QJ01	220, 221
Yoshiyuki, Mizuno	RO05	301
Yoshizawa, Masahito	PD11	133
	PDTT PL01	133
You, Biao		
You, Chun Yeol	FD03, QM18, QM20, SH16	74, 234, 345
	5/110	

You, Chun-yeol

You, Y. B. Youn, S. J. Young, A. T. Young, Olga Yu, Han Young Yu, Ing-song Yu, J H Yu, Jaejun Yu, Jun Sang Yu, Kwon Kyu Yu, Kwon-kyo Yu, R. Yu, S. C. Yu, Seong Cho Yu, Seong-cho Yu, Unjong Yu, Weiqiang Yu, Xiuzhen Yu, Xuezhe Yu, Young-sang Yuan, Fu-te Yuan, H. Q. Yuan, Hongtao Yuan, Hui Qiu Yuan, Shujuan Yuan, Xueyong Yuan, Zhi-min Yuasa, S. Yuasa, Shinji Yue, Ming Yue, Zengji Yufu, Shigeyuki Yukawa, Koichi Yuki, Shozo Yukitake, Chikano

#### Ζ

Zablotskii, Vitalii Zagdoud, Amani Zagrebin, Mikhail Zaharko, O. Zaharko, Oksana

QE19

QF20

Yun, Sang-jun

Yun, Won Seok

Yunhyun, Cho

Yunoki, Seiji

Yuryev, Y. N.

Yusa, Hitoshi

Yusuf, S M

Yusuf, S. M.

Yutani, K.

Yusuke, Amakai

Yutani, Keisuke

Yuu, Okawara

RK02, RK06, RK09, 286, 287, 289, 346, RK22, SI04, SJ02, SJ05, 349, 361, 362 SM01, SM12 SB04 320 RB26 258 220 QI14 79 GF03 293 RL21 58 DD01 RO12 302 DC03, GF05, KG02, 58, 79, 111, 261 RC13, RC17 PO09 178 RQ17 309 RQ18 310 50 CD01 RA01 249 AJ06 40 169, 177, 213, 343, PM12, PN31, QH01, SH01, SO03 367 RF06 269 QB07 188 RA05 250 DG06 61 60, 285, 345 DE05, RJ16, SH14 PK06, QP17 161, 248 QB12 189 AE05 36 RD12, SB29 264, 325 SG01 339 PL01 164 SJ09 350 51, 52, 93, 283 CE03, CE06, ID02, RJ07 AI05, BI01, HF01, IE04, 39, 47, 87, 94, 110, KF01, PK20 163 RO15, SN01 303, 364 SD01 330 PB04 120 PC05 127 QN11 238 PB04 120 KE03, SI13 110, 348 KI01, QG20 112, 211 RR20 314 PC08, PG18 128, 145 RQ08 308 PE10 137 QO06, QP05 241, 246 DH05, PN17, RN04 62, 174, 297 QC01 192 KA02, QC12 107, 194 PD24 135 QJ01 221 AD05, AH02, BH03 38, 46 RK05 286 213 QH05

## Author Index

202

Zahradnik, Martin	KD03	109	Zhao, Yonggang	
Zakhidov, A.	QB12	189	Zhao, Yu-jun	
Zaleski, A J	QB06	188	Zhao, Z. Y.	
Zaleski, T A	QB06	188	Zhaoyang, Zeng	
Zamani, Atieh	QH11	214	Zheludev, Andrey	
Zamudio-bayer, Vicente	KD01	109	Zhen, X.g.	
Zang, Jiadong	DF01	60	Zheng, G -q	
Zanini, L. F.	JH02	104	Zheng, Guo-qing	
Zapf, Vivien	SF13	339	Zheng, H. Y.	
Zapf, Vivien S	QB13	189	Zherlitsyn, S.	
Zapotokova, Mariana	QD08, RD21	197, 265	Zherlitsyn, Sergei	
Zarefy, Amjaad	BD02	43	Zhigadlo, N. D.	
Zargar Shoushtari, Morteza	PB15, RN22	122, 300	Zhigadlo, Nikolai	
Zeiringer, Isolde	SC01	325	Zhigadlo, Nikolai D.	
Zeitler, Uli	KB06, SA10	108, 317	Zhigaldo, Nicolai	
Zeitler, Ulrich	DH02	62	Zhou, Dong	
Zemen, Jan	SK22	354	Zhou, Haidong	
Zeng, Minggang	GI04	82	Zhou, Mingge	
Zeng, Rong	QG09	209	Zhou, S. M.	
Zenkevich, Andrei	SK26	354	Zhou, Shaoxiong	
,	RH19		Zhou, Shiming	
Zenmyo, Kazuko Zenzai, Rvesuko		277 370	. 5	
Zenzai, Ryosuke	SO20		Zhou, Tiejun	
Zhang, Aimei	PK05, RR29	161, 315	Zhou, Tong	
Zhang, Caihong	BE04	44	Zhou, X. H.	
Zhang, Chenglin	RB10, RB11	255	Zhou, X. J.	
Zhang, Dongtao	SN01	364	Zhu, Hong	
Zhang, Fapei	QK01	224	Zhu, J.	
Zhang, Fengming	PN05, QL04	172, 228	Zhu, Jian Xin	
Zhang, Haijun	RC10	260	Zhu, Liming	
Zhang, Hyodo	RN13	298	Zhu, Minggang	
Zhang, J. L.	QB12	189		
Zhang, Jing Lei	RD12, SB29	264, 325	Zhu, Tao	
Zhang, Jiuxing	SN01	364	Zhu, Weihua	
Zhang, L	BG01	45	Zhu, Yao-hui	
Zhang, L.	SO08, SO16	368, 369	Zhu, Zengwei	
Zhang, P.z.	RO14	303	Zhut, Meng	
Zhang, Peng	SO03	367	Zhukov, Arcady	
Zhang, Pengyue	RO13, RO15	302, 303		
Zhang, Rong	PN27	176	Zhukova, Valentina	
Zhang, Sen	SA19, SA20	319	בווטונטים, ימוכוונווום	
Zhang, Shoucheng	HP71, RC10	32, 260	Zilch, Christian	
Zhang, Shuai	RD14	264	Ziman, Timothy	
Zhang, Weimin	QF04	204	Ziman, moury Zimmermann, Martin Vo	
Zhang, Wei-min Zhang, Wei-min		199, 201	Zitko, Rok	
-	QE03, QE10			
Zhang, Wenchen	SN11	366	Zitouni, Abdelaziz	
Zhang, Xiaozhong	HP61	31	Zivieri, Roberto	
Zhang, Xin	SA19	319	Zivkovic, Ivica	
Zhang, Y.	JH02	104	Zlatic, Veljko	
Zhang, Yi-jie	DB04	57	Zlotnikov, Anton O.	
Zhang, Ying	EJ02	71	Zomer, Paul J.	
Zhang, Ying-de	SO03	367	Zorko, Andrej	
Zhang, Yiwen	JIO1	105	Zou, Tao	
Zhang, Yu-zhong	SF14	339	Zou, Wenqin	
Zhao, Huaisong	PB28	124	Zou, Yang	
Zhao, Jianhua	DG06, QK23	61, 227	Zuberek, Ryszard	
Zhao, Pei	PK18	163	Zukeran, Chojun	
Zhao, X. R.	PL20	167	Zukovic, Milan	
		49		
Zhao, Xudona	CA03	49	Zuo, vveniidnu	
Zhao, Xudong Zhao, Yang	CA03 QA20	185	Zuo, Wenliang Zvezdin, K.a	

SA20

QA03

IC03

Q015

QE18

QB11

PC20

CC04

PC29

RB19 RB21

RB20

PA21 SL12

PL20

IB02

SJ09

SF01

RM09

SB31

PK03

KG01

QL03

PK14

PK05 QK08

KJ04

EG01

SO27

AH01 JF03

CI04

PE03

RQ16

GG02

QF20

IA02

QC02

GI05

AC04

JB03

JJ03

QI13

QN23

ID02

ID03, IE04

PN05, QL04

QE01, QE02

RE14, RE15, SF05

PL02, SM05

QN21, RO04, RO07,

SL12, SN11, SN12

PJ25, SO04, SO27

BD05, JI05

QN21, SL12

SL08, SL09

RD07, RE09, SD07

CC03, IJ06, PI06

319

183

92

243

202

188

130

357

50

131 257

257

257

358

167

91

350

337

296

325

160

111

228

162 161

225

114

69 AJ02, CH01, CJ02, CJ04, 40, 54, 55, 56, 158,

> 371 38

102

55

136

309

80

207

90

192

82

35

99

106

219

199

240

93

93, 94

172, 228

268, 337

AJ02, CJ02, PJ25, SO04, 40, 55, 158, 367,

367, 371

164, 361

239, 301, 358, 366

43, 105

239, 358 119

263, 267, 331

50, 98, 151

Zvyagin, A. A.	CC04, QF02	50, 204
Zvyagin, S.	QF03	204
Zvyagin, S. A.	QF02	204
Zwicknagl, Gertrud	DB02, PC24	57, 130

## AUTHOR INDEX

MEMO	MEMO

MEMO	_	MEMO
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